

Spectroscopy of Magnetized Black Holes and Topological Stars

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In collaboration with:
Marco Melis & Paolo Pani

Based on:
[arXiv:2406.19327](https://arxiv.org/abs/2406.19327) [gr-qc]



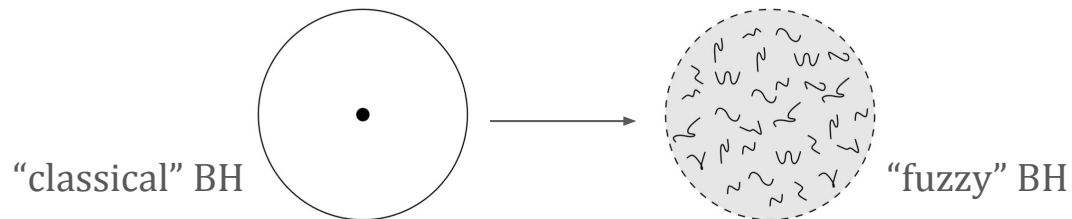
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Outline

- The Big Picture: BH microstructure
- Topological Stars as toy models
- Linear stability analysis:
 - Regge-Wheeler-Zerilli perturbation scheme
 - Linear response in t-domain: stability, echoes & QNM spectrum
- Summary & Follow-up

The big picture: BH microstructure

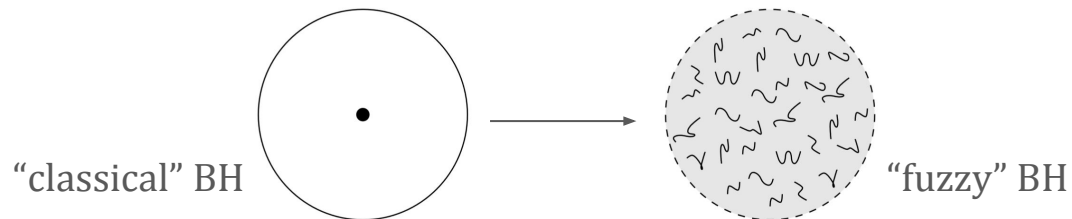


- **The Fuzzball program:** describing the BH microstructure
- BH as ensemble of **many, regular** and **horizonless** microstates
- Catch **B-H Entropy**, **Singularity** and **Information Paradox** with one stone!

Mathur (2025)

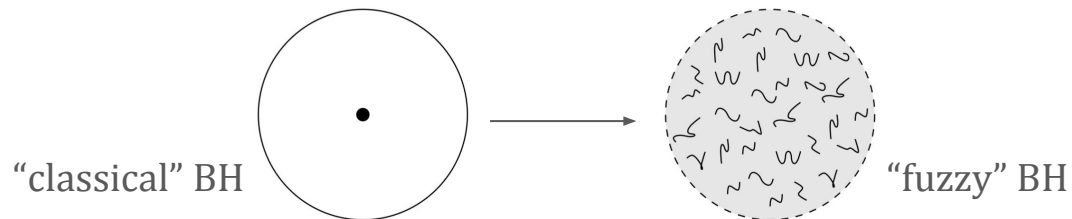
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The big picture: BH microstructure



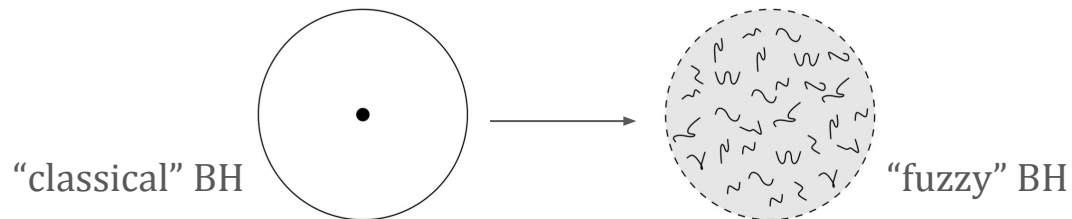
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- **BUT:** higher-D, complex topologies, supersymmetry, ...
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Bah & Heidmann (2021)

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**CHALLENGE
ACCEPTED!**

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Topological Stars in 5D

5D Einstein+Maxwell:
$$S_5 = \int d^5x \sqrt{-\mathbf{g}} \left(\frac{1}{2\kappa_5^2} \mathbf{R} - \frac{1}{4} \mathbf{F}_{AB} \mathbf{F}^{AB} \right)$$

$$ds^2 = -f_S(r) dt^2 + f_B(r) dy^2 + \frac{1}{h(r)} dr^2 + r^2 d\Omega_2^2 \quad F = P \sin \theta d\theta \wedge d\phi$$

$$f_B(r) = 1 - \frac{r_B}{r}, \quad f_S(r) = 1 - \frac{r_S}{r}, \quad h(r) = f_B(r) f_S(r), \quad P = \pm \kappa_5^{-1} \sqrt{\frac{3}{2} r_B r_S}$$

Bah & Heidmann (2021)

- **horizon** r_S
- **end-of-spacetime** r_B

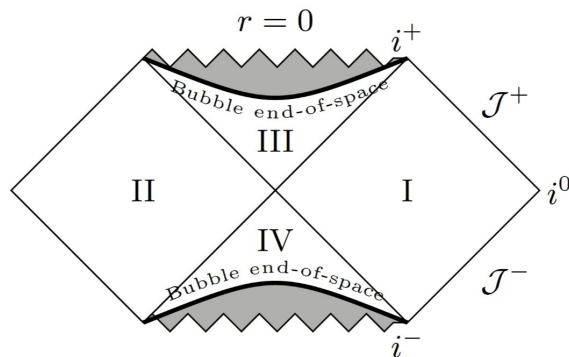
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- **Magnetized Black String:** $r_B < r_S$



Bah & Heidmann (2021)

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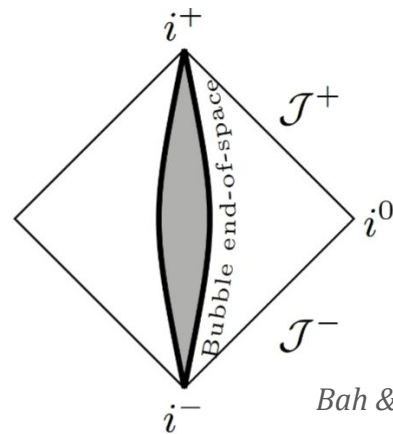
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- **Topological Star:** $r_B > r_S$



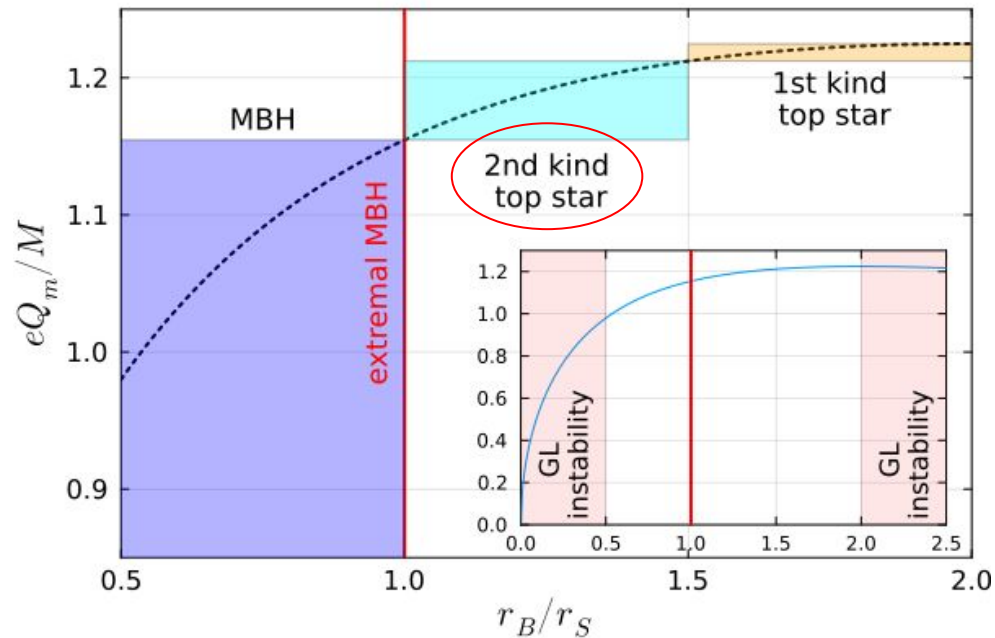
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Topological Stars in 5D

- 1st kind TS:** $\frac{3}{2}r_S \leq r_B \leq 2r_S$
 $r_{(ph)} = r_B$
 (unstable)
- 2nd kind TS:** $r_S < r_B < \frac{3}{2}r_S$
 $r_{(ph)} = \frac{3}{2}r_S$, $r_{(ph)} = r_B$
 (stable) (unstable)

Gregory-Laflamme instability:

$$r_B < \frac{1}{2}r_S, \quad r_B > 2r_S$$



Kaluza-Klein reduction to 4D

KK reduction (assume no y -dependence):

$$ds_5^2 = e^{-\frac{\sqrt{3}}{3}\Phi} ds_4^2 + e^{2\frac{\sqrt{3}}{3}\Phi} (dy + \cancel{A}_\mu dx^\mu)^2$$

decouple!

$$\mathbf{F}_{AB} dx^A dx^B = F_{\mu\nu} dx^\mu dx^\nu + (\partial_\mu \cancel{A}_\nu dx^\mu) \wedge (dy + A_\mu dx^\mu)$$

Einstein-Maxwell-scalar:

$$\mathcal{S} = \int dx^4 \sqrt{-g} \left[\frac{1}{2\kappa_4^2} \left(R - \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi \right) - \frac{1}{4e^2} e^{\frac{\sqrt{3}}{3}\Phi} F_{\mu\nu} F^{\mu\nu} \right]$$

Field equations:

$$G_{\mu\nu} - \frac{1}{2} \left(\partial_\mu \Phi \partial_\nu \Phi - \frac{1}{2} g_{\mu\nu} \partial_\rho \Phi \partial^\rho \Phi \right) + \frac{\kappa_4^2}{e^2} e^{\frac{\Phi}{\sqrt{3}}} \left(F_{\mu\rho} F^\rho{}_\nu + \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} \right) = 0$$

$$\square \Phi - \frac{\sqrt{3}\kappa_4^2}{6e^2} e^{\frac{\Phi}{\sqrt{3}}} F_{\mu\nu} F^{\mu\nu} = 0 \qquad \nabla^\rho \left(e^{\frac{\sqrt{3}}{3}\Phi} F_{\mu\rho} \right) = 0$$

Topological Star in 4D

Compactified TS(/MBH):

$$ds_4^2 = -f_S f_B^{1/2} dt^2 + \frac{1}{f_S f_B^{1/2}} dr^2 + r^2 f_B^{1/2} d\Omega_2^2 \quad F = \pm e Q_m \sin \theta d\theta \wedge d\phi \quad \Phi = \frac{\sqrt{3}}{2} \log f_B$$

Mass and charge:
$$M = \frac{2\pi}{\kappa_4^2} (2r_S + r_B), \quad Q_m = \frac{1}{\kappa_4} \sqrt{\frac{3}{2} r_S r_B}$$

- Background for perturbation scheme
- Warning! **Compactification introduces singularities!**
- 5D solution implies regularity boundary conditions

Regge-Wheeler-Zerilli perturbation scheme

magnetic charge background \rightarrow parity-mixing of perturbations

$$F = \pm e Q_m \sin \theta d\theta \wedge d\phi$$

See **Pereñiguez (2021)** for alternative approach

Type-I

axial gravitational + polar EM ($l \geq 1$)

$$h_{\mu\nu}^{\text{odd}} = \sum_{l,m} \begin{pmatrix} 0 & 0 & -h_0(t,r)/\sin\theta\partial_\phi & h_0(t,r)\sin\theta\partial_\theta \\ 0 & 0 & -h_1(t,r)/\sin\theta\partial_\phi & h_1(t,r)\sin\theta\partial_\theta \\ -h_0(t,r)/\sin\theta\partial_\phi & -h_1(t,r)/\sin\theta\partial_\phi & 0 & 0 \\ h_0(t,r)\sin\theta\partial_\theta & h_1(t,r)\sin\theta\partial_\theta & 0 & 0 \end{pmatrix} Y_{lm}(\theta, \phi)$$

$$f_{\mu\nu}^{\text{even}} = \sum_{l,m} \begin{pmatrix} 0 & f_{01}^+(t,r) & f_{02}^+(t,r)\partial_\theta & f_{02}^+(t,r)\partial_\phi \\ -f_{01}^+(t,r) & 0 & f_{12}^+(t,r)\partial_\theta & f_{12}^+(t,r)\partial_\phi \\ -f_{02}^+(t,r)\partial_\theta & -f_{12}^+(t,r)\partial_\theta & 0 & 0 \\ -f_{02}^+(t,r)\partial_\phi & -f_{12}^+(t,r)\partial_\phi & 0 & 0 \end{pmatrix} Y_{lm}(\theta, \phi)$$

Type-II

polar gravitational + polar scalar + axial EM ($l \geq 0$)

$$h_{\mu\nu}^{\text{even}} = \sum_{l,m} \begin{pmatrix} f_S f_B^{1/2} H_0(t,r) & H_1(t,r) & 0 & 0 \\ H_1(t,r) & f_S^{-1} f_B^{-1/2} H_2(t,r) & 0 & 0 \\ 0 & 0 & r^2 f_B^{1/2} K(t,r) & 0 \\ 0 & 0 & 0 & r^2 f_B^{1/2} \sin^2\theta K(t,r) \end{pmatrix} Y_{lm}(\theta, \phi)$$

$$f_{\mu\nu}^{\text{odd}} = \sum_{l,m} \begin{pmatrix} 0 & 0 & f_{02}^-(t,r)/\sin\theta\partial_\phi & -f_{02}^-(t,r)\sin\theta\partial_\theta \\ 0 & 0 & f_{12}^-(t,r)/\sin\theta\partial_\phi & -f_{12}^-(t,r)\sin\theta\partial_\theta \\ -f_{02}^-(t,r)/\sin\theta\partial_\phi & -f_{12}^-(t,r)/\sin\theta\partial_\phi & 0 & f_{23}^-(t,r)\sin\theta \\ f_{02}^-(t,r)\sin\theta\partial_\theta & f_{12}^-(t,r)\sin\theta\partial_\theta & -f_{23}^-(t,r)\sin\theta & 0 \end{pmatrix} Y_{lm}(\theta, \phi)$$

$$\delta\Phi = \sum_{l,m} \frac{\varphi(t,r)}{r} Y_{lm}(\theta, \phi)$$

Linear response in t-domain

Master equation for **Type-I ($l \geq 1$)** and **Type-II ($l = 0$)**: $\Psi = \{\varphi_{l=0}, \mathcal{E}_{l=1}, Z_{1;l>1}, Z_{2;l>1}\}$

$$\left[\frac{d^2}{dt^2} - \frac{d^2}{d\rho^2} + V_{\text{eff}} \right] \Psi(t, \rho) = 0$$

Generalized tortoise coordinate:

$$d\rho = \sqrt{\frac{g_{rr}}{g_{tt}}} dr = \frac{dr}{f_B^{1/2} f_S}$$

Effective potentials:

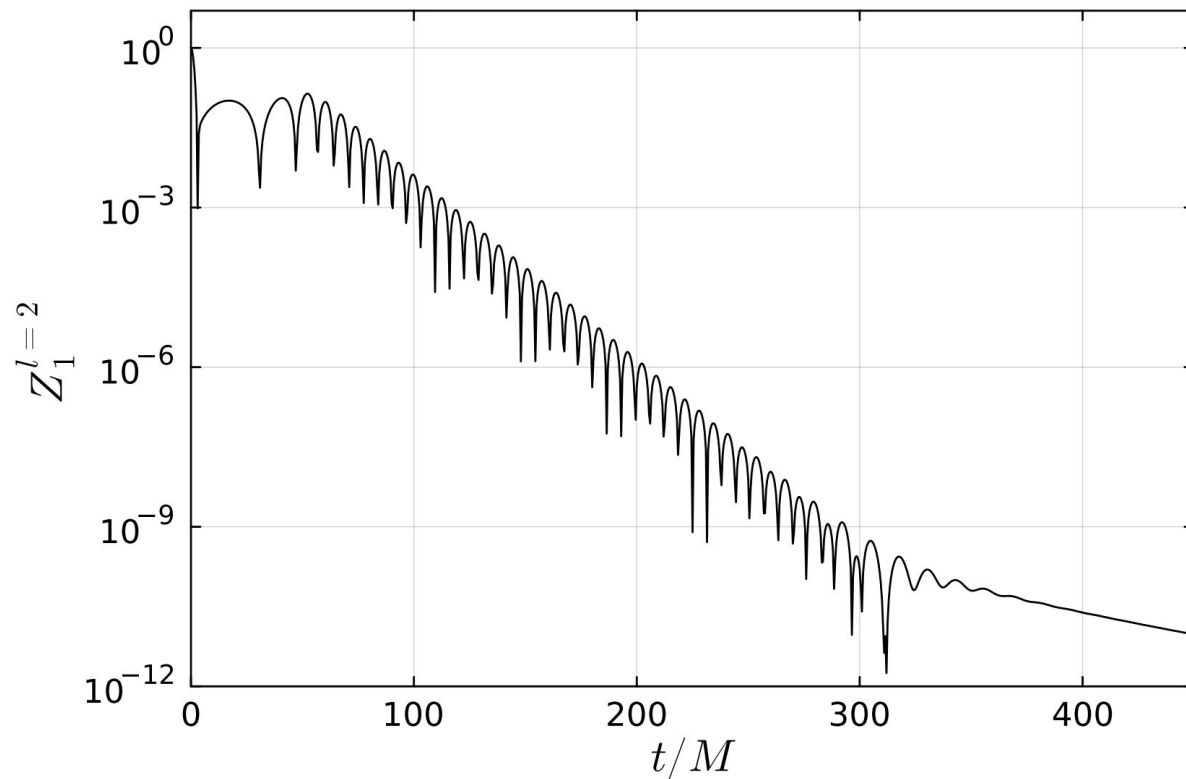
$l = 0$	$V_{\text{eff}}^{l=0} = \frac{(r-r_S)(32r^3 - 24r^2 r_B - 39r_B^2 r_S + r r_B (36r_S - 5r_B))}{16r^5(r-r_B)}$
$l = 1, \text{ Type-I}$	$V_{\text{eff}}^{l=1} = \frac{(r-r_S)(189r_B^4 r_S + 128r^4(r_B + 2r_S) + 64r^2 r_B^2(6r_B + 19r_S) - 16r^3 r_B(27r_B + 52r_S) - 9r r_B^3(9r_B + 92r_S))}{16r^5(4r - 3r_B)^2(r-r_B)}$
$l \geq 2, \text{ Type-I}$	$V_{\text{eff}}^{(1,2)} = \frac{(r-r_S)(16r^3 \Lambda - r^2(8r_B + 24r_S + 16\Lambda r_B) + r(11r_B^2 + 60r_B r_S) - 39r_B^2 r_S \mp 8r(r-r_B)\sqrt{(2r_B - 3r_S)^2 + 12r_B r_S \Lambda})}{16r^5(r-r_B)}$

Type-II ($l > 0$)? Coming soon!

Linear response in t-domain

Near-extremal MBH:

- (type-I) gravitational-induced perturbation
- fixed mass ($M=1$)
- ringdown + tail



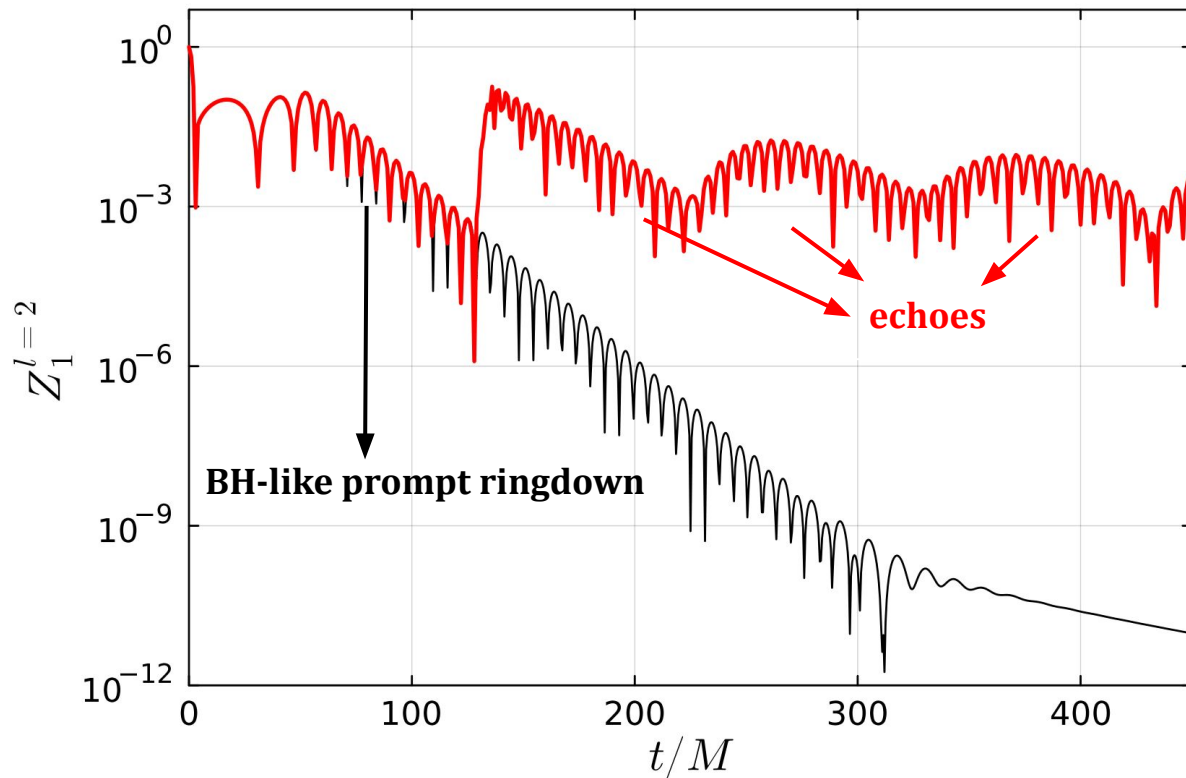
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2nd kind Top Star:

- fixed mass ($M=1$)
- very compact
- $e\Delta Q/M \approx 0.3\%$
- initial **BH-like ringdown**
- long-lived modes: **echoes!**



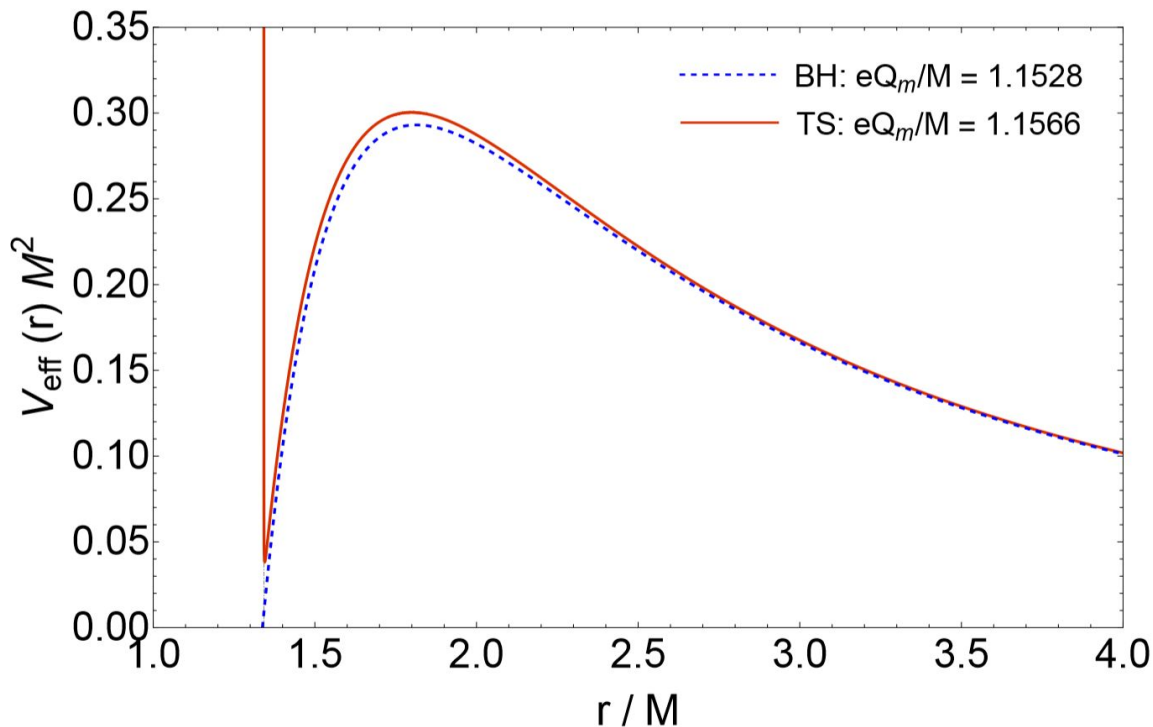
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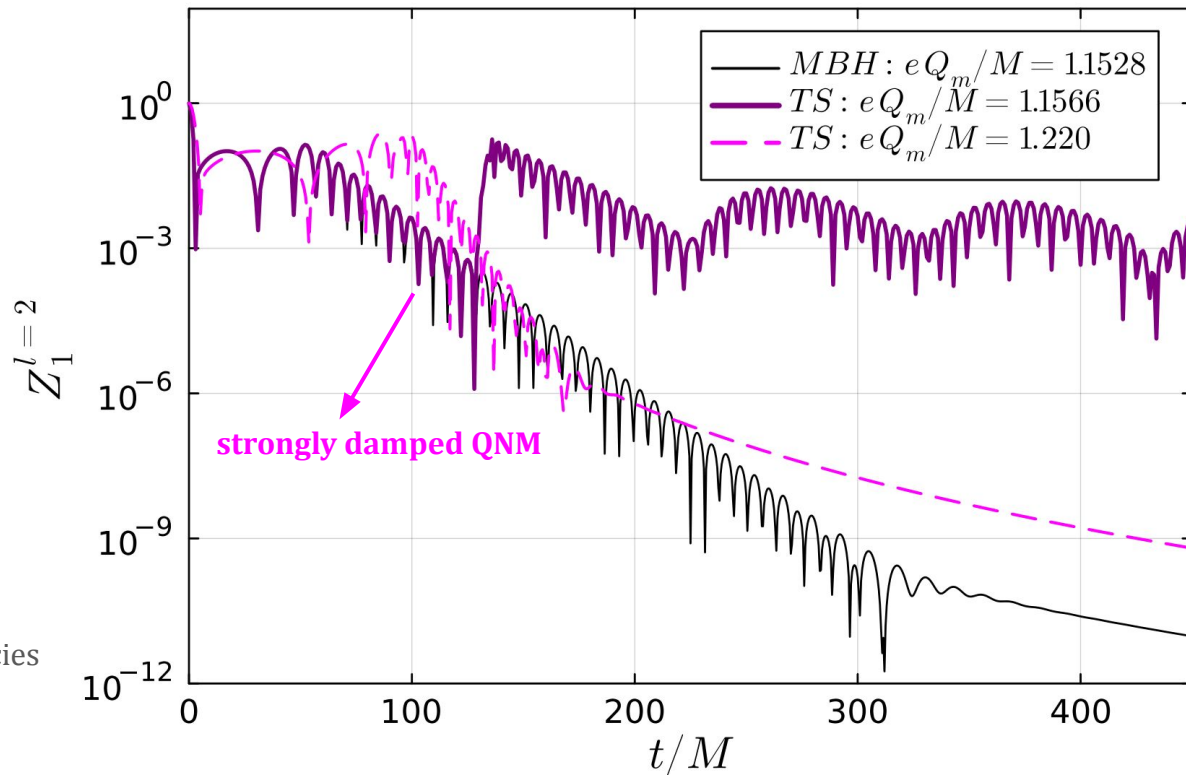
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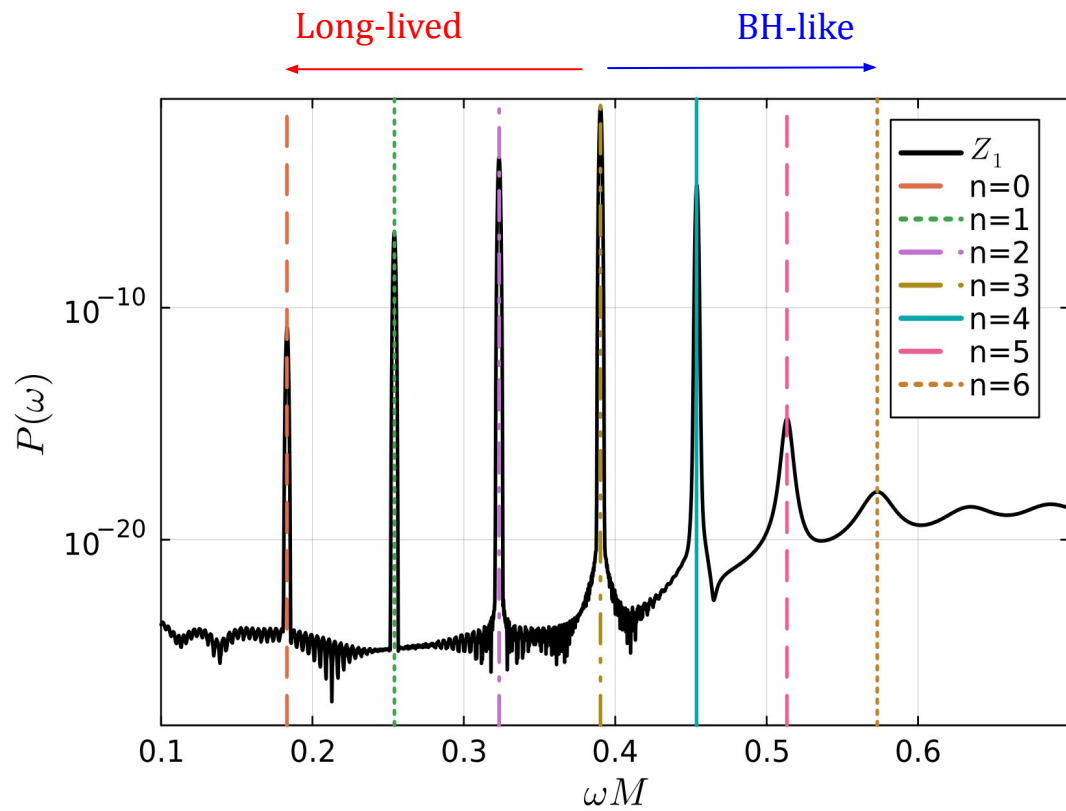
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- initial **BH-like ringdown**
- long-lived modes: **echoes!**

1st kind Top Star:

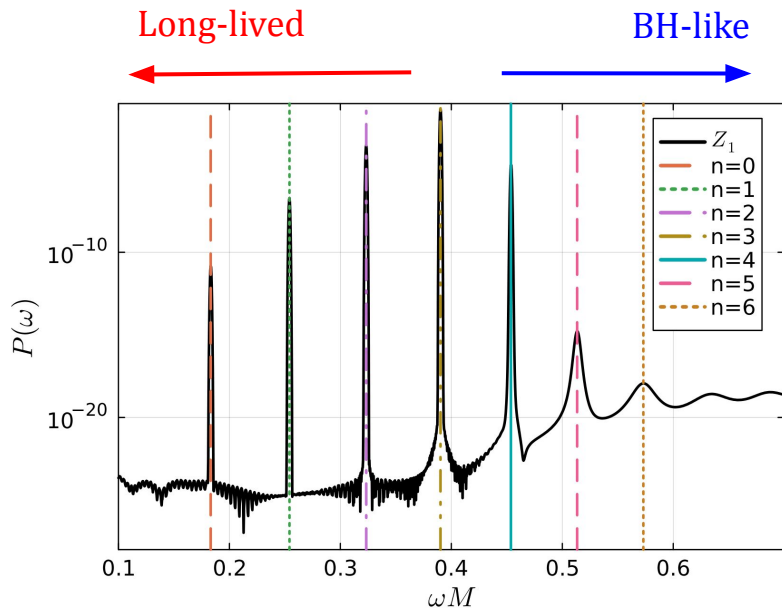
- fixed mass ($M=1$)
- star-like compactness
- highly-damped, higher frequencies (**w-modes**)



QNM Spectrum



QNM Spectrum



		Magnetized BH	TS, second kind	TS, first kind
$n = 0$	f-domain	$0.489568 - i 7.972 \times 10^{-2}$	$0.183217 - i 4.674 \times 10^{-10}$	$0.644348 - i 0.1551$
	t-domain	$0.489600 - i 7.978 \times 10^{-2}$	$0.183219 - i 3.349 \times 10^{-10}$	$0.643938 - i 0.1665$
$n = 1$	f-domain	-	$0.254071 - i 6.001 \times 10^{-8}$	-
	t-domain	-	$0.254084 - i 6.008 \times 10^{-8}$	-
$n = 2$	f-domain	-	$0.323219 - i 2.615 \times 10^{-6}$	-
	t-domain	-	$0.323263 - i 2.622 \times 10^{-6}$	-
$n = 3$	f-domain	-	$0.390169 - i 6.116 \times 10^{-5}$	-
	t-domain	-	$0.390256 - i 6.142 \times 10^{-5}$	-
$n = 4$	f-domain	-	$0.453786 - i 8.348 \times 10^{-4}$	-
	t-domain	-	$0.453832 - i 8.340 \times 10^{-4}$	-
$n = 5$	f-domain	-	$0.513765 - i 5.463 \times 10^{-3}$	-
	t-domain	-	$0.513375 - i 2.754 \times 10^{-3}$	-
$n = 6$	f-domain	-	$0.574947 - i 1.658 \times 10^{-2}$	-
	t-domain	-	$0.572869 - i 1.140 \times 10^{-2}$	-

Summary & Follow-up

- **Topological Stars** as **toy models** of BH microstate geometries
- Two-parameter solution that **interpolates** between **regular UCOs** and **BHs**
- **Linear response in time** (this talk) and **frequency domain** (attend Marco's talk!)
- Verified **linear stability** under **Type-I** perturbations and **(radial) Type-II**
- **QNM spectrum** of MBH and Top Stars
 - **BH-like QNMs**
 - **Long-lived modes (echoes)**
 - **Strongly-damped (w-modes)**
- Next steps:
 - **Type-II, $l > 0$**
 - **Nonlinear stability and 3+1 dynamics**
 - **More complex microstate geometries**

AD, M. Melis, P. Pani, 2406.19327
and

I. Bena, G. Di Russo, J. F. Morales, and
A. Ruiperez Vicente, 2406.19330