

# Perturbation of the Vaidya metric in the frequency domain: QNMs and Tidal Response

Lodovico Capuano

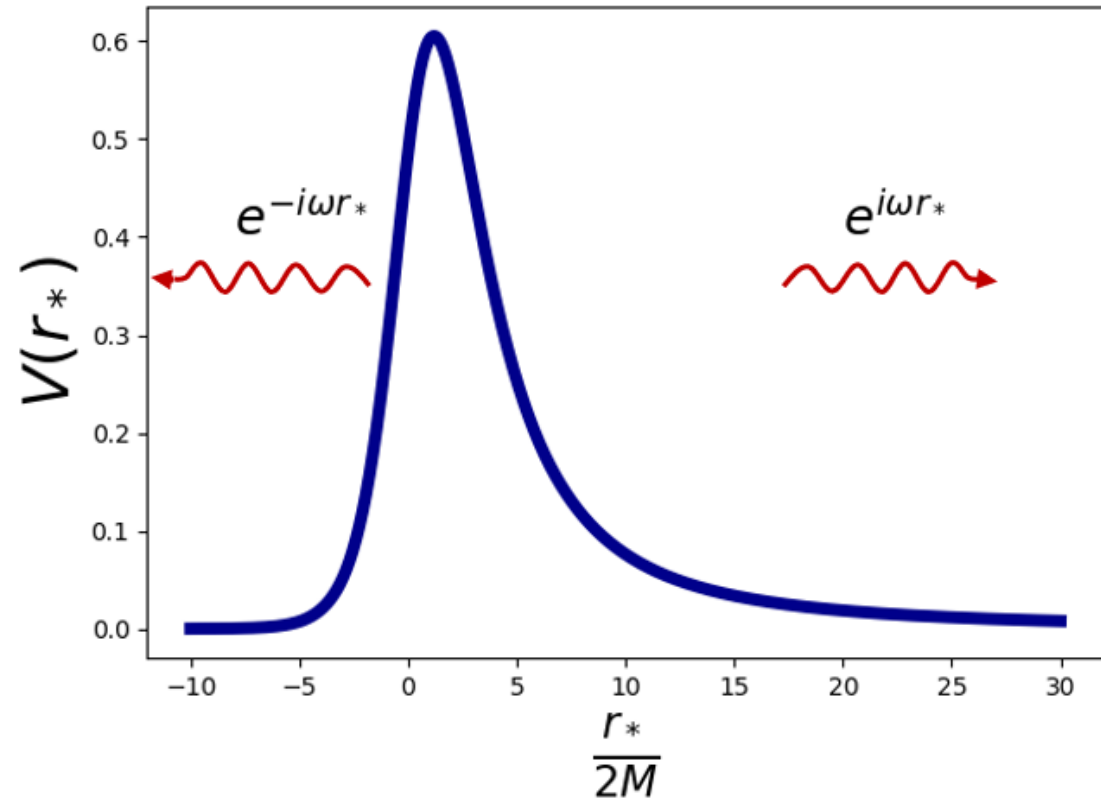
From arXiv:2407.06009  
LC, Luca Santoni, Enrico Barausse



# Quasi-Normal Modes

Schrödinger-like equation:

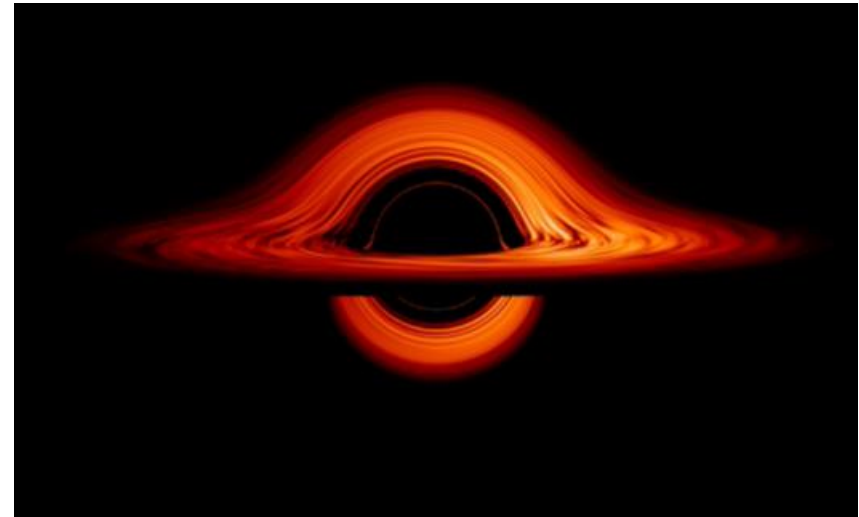
$$\left( \frac{d^2}{dr_*^2} - V(r_*) \right) \psi(r_*) = -\omega^2 \psi(r_*)$$



# Dynamical BHs

## Motivation

- **Accretion**
- **Hawking evaporation**
- **Superradiance**
- **GW absorption**



Credits: NASA

## The Vaidya solution

- **Metric in Eddington-Finkelstein coordinates:**

$$ds_V^2 = - \left( 1 - \frac{2M(u)}{r} \right) du^2 \pm 2du dr + r^2 d\Omega_{S^2}^2$$

- **Stress-Energy Tensor:**

$$T_{\mu\nu}^V = \pm \frac{M'(u)}{4\pi r^2} k_\mu k_\nu$$

## Time domain approach

- **Solve perturbation equation numerically as IVP**

Abdalla, Chirenti and Saa, 2006

Lin, Sun and Zhang, 2021

Redondo-Yuste, Pereñíguez and Cardoso, 2023

## Time domain approach

- **Solve perturbation equation numerically as IVP**  
Abdalla, Chirenti and Saa, 2006  
Lin, Sun and Zhang, 2021  
Redondo-Yuste, Pereñíguez and Cardoso, 2023
- **Any time-dependent mass function can be chosen**

## Time domain approach

- **Solve perturbation equation numerically as IVP**  
Abdalla, Chirenti and Saa, 2006  
Lin, Sun and Zhang, 2021  
Redondo-Yuste, Pereñiguez and Cardoso, 2023
- **Any time-dependent mass function can be chosen**
- **Doesn't allow high accuracy frequency extraction**

## Constant rate approximation

- **Rescaled coordinates:**

$$\begin{cases} T = \int \frac{du}{2M(u)} + \int \frac{dx}{f(x)} \\ x = \frac{r}{2M(u)} \end{cases}$$



## Constant rate approximation

- **Rescaled coordinates:**

$$\begin{cases} T = \int \frac{du}{2M(u)} + \int \frac{dx}{f(x)} \\ x = \frac{r}{2M(u)} \end{cases}$$

- **Conformal transformation of the metric:**

$$ds_V^2 = M(T, x)^2 \left[ -f(x)dT^2 + f(x)^{-1}dx^2 + x^2 d\Omega_{S^2}^2 \right]$$

## Constant rate approximation

- **Rescaled coordinates:**

$$\begin{cases} T = \int \frac{du}{2M(u)} + \int \frac{dx}{f(x)} \\ x = \frac{r}{2M(u)} \end{cases}$$

$$f(x) = 1 - \frac{1}{x} - 4|M'|x$$

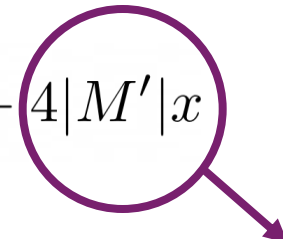
- **Conformal transformation of the metric:**

$$ds_V^2 = M(T, x)^2 \left[ -f(x)dT^2 + f(x)^{-1}dx^2 + x^2d\Omega_{S^2}^2 \right]$$

## Constant rate approximation

- **Rescaled coordinates:**

$$\begin{cases} T = \int \frac{du}{2M(u)} + \int \frac{dx}{f(x)} \\ x = \frac{r}{2M(u)} \end{cases}$$

$$f(x) = 1 - \frac{1}{x} - 4|M'|x$$


Similar to  
Schwarzschild-de Sitter

- **Conformal transformation of the metric:**

$$ds_V^2 = M(T, x)^2 \left[ -f(x)dT^2 + f(x)^{-1}dx^2 + x^2d\Omega_{S^2}^2 \right]$$

## Constant rate approximation

- **Rescaled coordinates:**

$$\begin{cases} T = \int \frac{du}{2M(u)} + \int \frac{dx}{f(x)} \\ x = \frac{r}{2M(u)} \end{cases}$$

$$f(x) = 1 - \frac{1}{x} - 4|M'|x$$

Similar to  
Schwarzschild-de Sitter

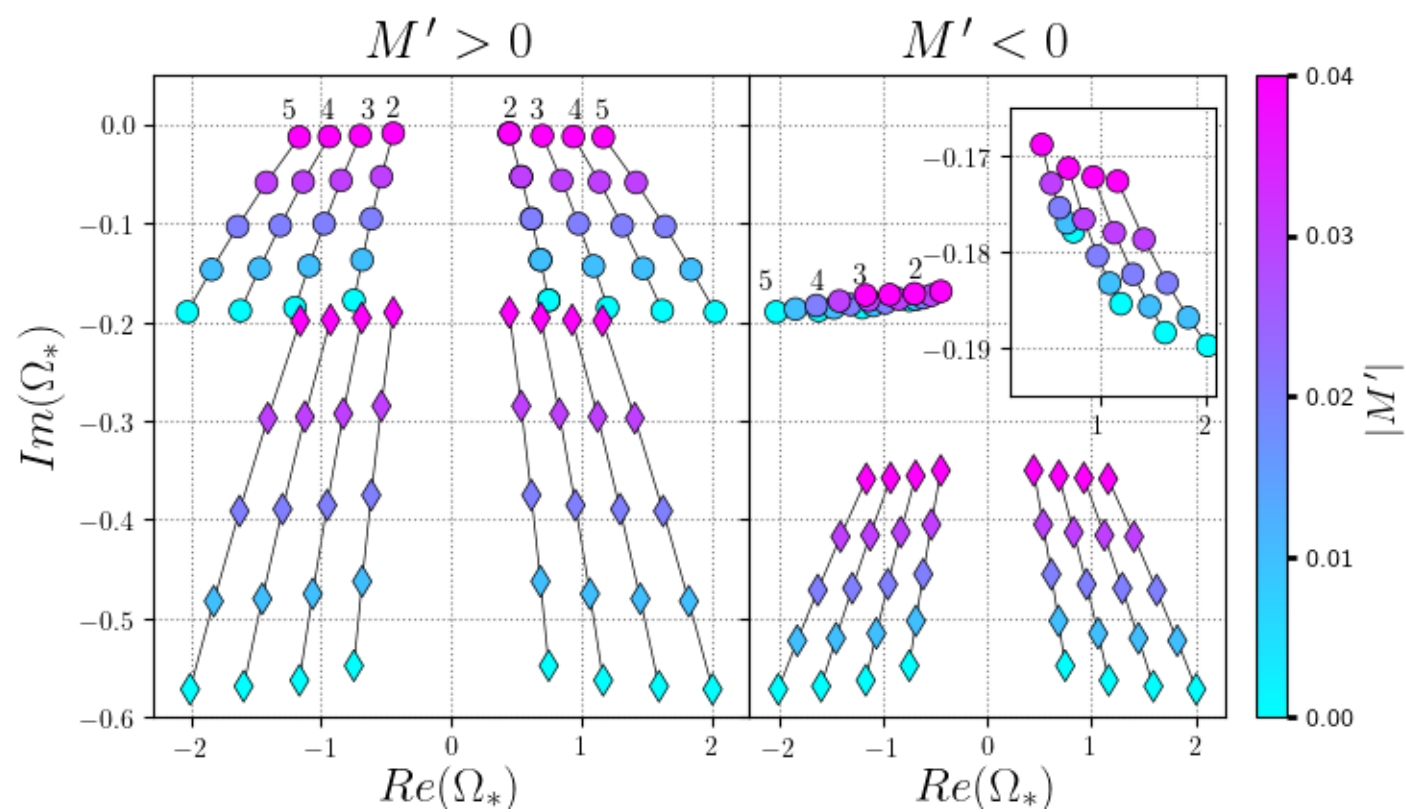
- **Conformal transformation of the metric:**

$$ds_V^2 = M(T, x)^2 \left[ -f(x)dT^2 + f(x)^{-1}dx^2 + x^2 d\Omega_{S^2}^2 \right]$$

- **Perturbation master equation:**

$$\left[ \frac{d^2}{dx_*^2} + \Omega^2 - f(x)V(x) \right] R(x) = 0$$

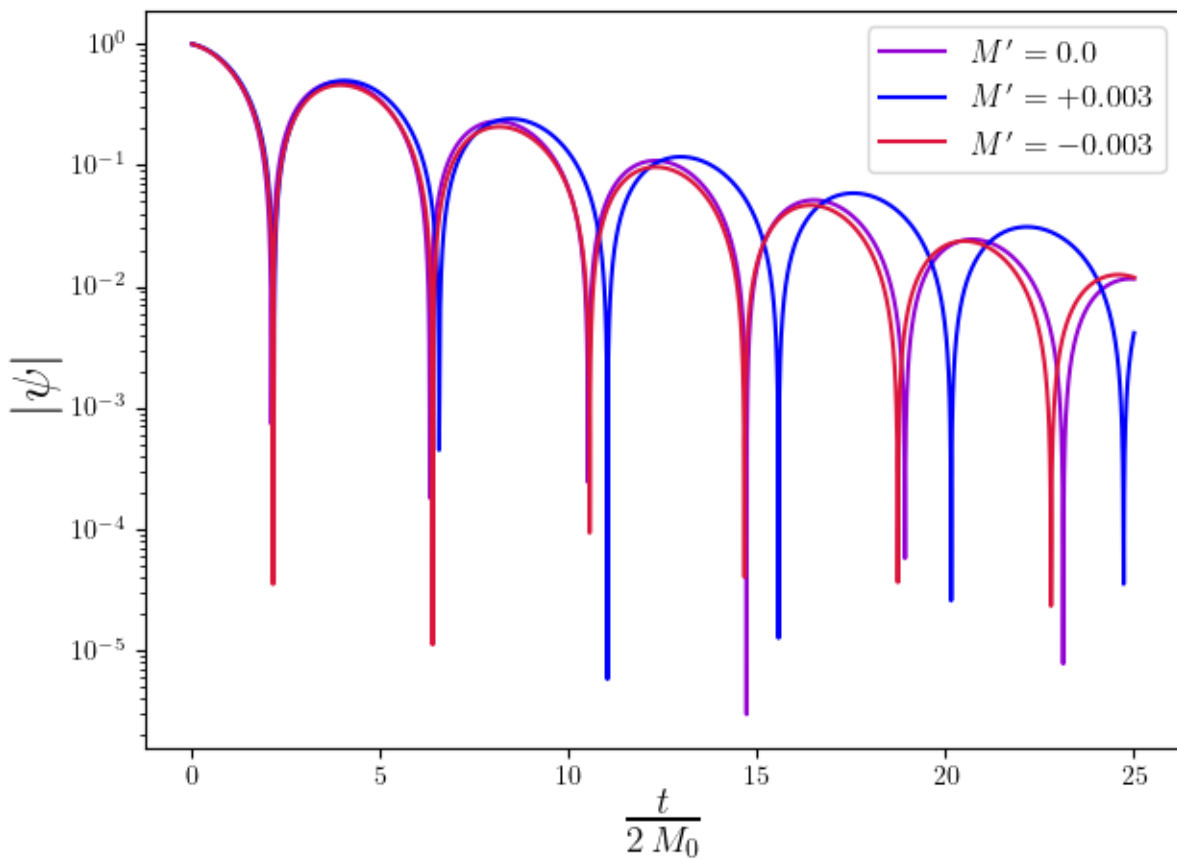
# RESULTS: QNMs



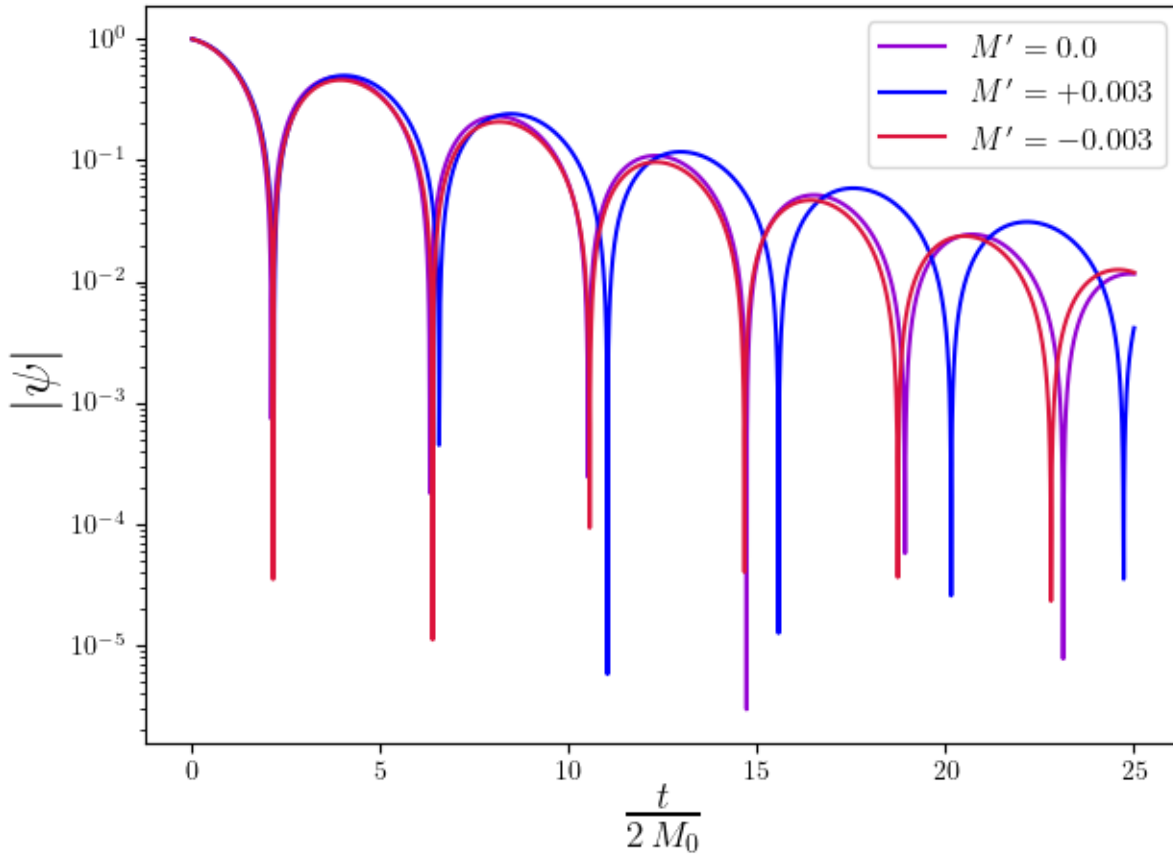
# INSTABILITIES

- **Scalar: possible for decreasing mass**
- **Electromagnetic: no instability**
- **Gravitational: possible for increasing mass**

## RESULTS: Time Domain



## RESULTS: Time Domain



$$\omega(u) \propto M(u)^{-1}$$



# TIDAL RESPONSE

## General treatment

**Static solution:**

$$\psi_{\text{static}} \sim x^{\ell+1} + k x^{-\ell}$$

# TIDAL RESPONSE

## General treatment

**Static solution:**

$$\psi_{\text{static}} \sim x^{\ell+1} + k x^{-\ell}$$

Tidal field



# TIDAL RESPONSE

## General treatment

**Static solution:**

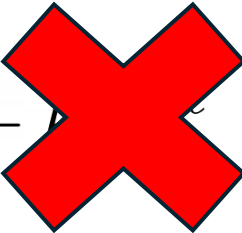
$$\psi_{\text{static}} \sim x^{\ell+1} + k x^{-\ell}$$

Tidal field                      Response

# TIDAL RESPONSE

## General treatment

**Static solution:**

$$\psi_{\text{static}} \sim x^{\ell+1} + \dots$$


In GR the LNs vanish!

# TIDAL RESPONSE

## The Vaidya case

$$\psi(u, r) \sim \left(\frac{r}{2M_0}\right)^{\ell+1} \left[ 1 + w(u, r) + \mathcal{O}\left(\frac{2M_0}{r}\right) + \right. \\ \left. + k \left(\frac{r}{2M_0}\right)^{-(2\ell+1)} \left( 1 + \mathcal{O}\left(\frac{2M_0}{r}\right) \right) \right].$$

# TIDAL RESPONSE

## The Vaidya case

Correction to tidal field

$$\psi(u, r) \sim \left(\frac{r}{2M_0}\right)^{\ell+1} \left[ 1 + w(u, r) + \mathcal{O}\left(\frac{2M_0}{r}\right) + \right. \\ \left. + k \left(\frac{r}{2M_0}\right)^{-(2\ell+1)} \left( 1 + \mathcal{O}\left(\frac{2M_0}{r}\right) \right) \right].$$

# TIDAL RESPONSE

## The Vaidya case

$$\psi(u, r) \sim \left(\frac{r}{2M_0}\right)^{\ell+1} \left[ 1 + w(u, r) + \mathcal{O}\left(\frac{2M_0}{r}\right) + \right. \\ \left. + k \left(\frac{r}{2M_0}\right)^{-(2\ell+1)} \left(1 + \mathcal{O}\left(\frac{2M_0}{r}\right)\right) \right].$$

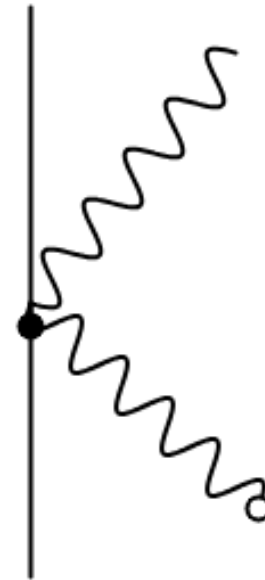
Correction to tidal field

Constant correction to Schwarzschild  
Love numbers

# Point-particle EFT

- **Full action:**

$$S_{\text{total}} = S_{\text{scalar}} + S_{\text{gravity}} + S_{\text{SET}} - \int dT e(x) \left[ M_0 + \sum_{\ell=1}^{\infty} \frac{\mu_{\ell}}{\ell!} \left( \partial_{(a_1} \dots \partial_{a_{\ell})} \varphi \right)^2 \right]$$





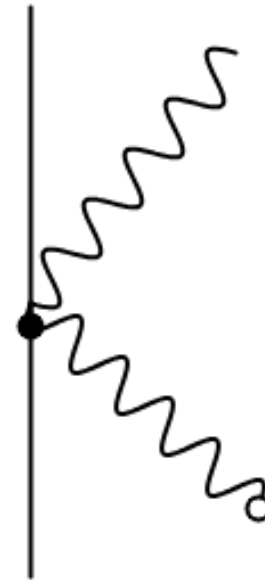
# Point-particle EFT

- **Full action:**

$$S_{\text{total}} = S_{\text{scalar}} + S_{\text{gravity}} + S_{\text{SET}} - \int dT e(x) \left[ M_0 + \sum_{\ell=1}^{\infty} \frac{\mu_{\ell}}{\ell!} \left( \partial_{(a_1} \dots \partial_{a_{\ell})} \varphi \right)^2 \right]$$

- **Solve KG equation perturbatively in the rate:**

$$\frac{1}{\sqrt{-g}} \partial_i \left( \sqrt{-g} g^{ij} \partial_j \varphi \right) = -\mu_{\ell} (-1)^{\ell} \partial^{(\ell)} \delta_D^{(3)}(x - \tilde{x})$$



# Point-particle EFT

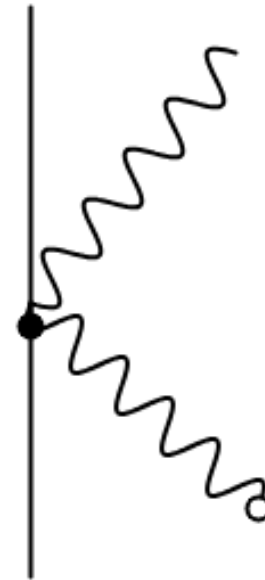
- **Full action:**

$$S_{\text{total}} = S_{\text{scalar}} + S_{\text{gravity}} + S_{\text{SET}} - \int dT e(x) \left[ M_0 + \sum_{\ell=1}^{\infty} \frac{\mu_{\ell}}{\ell!} \left( \partial_{(a_1} \dots \partial_{a_{\ell})} \varphi \right)^2 \right]$$

- **Solve KG equation perturbatively in the rate:**

$$\frac{1}{\sqrt{-g}} \partial_i \left( \sqrt{-g} g^{ij} \partial_j \varphi \right) = -\mu_{\ell} (-1)^{\ell} \partial^{(\ell)} \delta_D^{(3)}(x - \tilde{x})$$

$$g = \text{Minkowsky} + M' \delta g + \dots$$



# Point-particle EFT

- **Full action:**

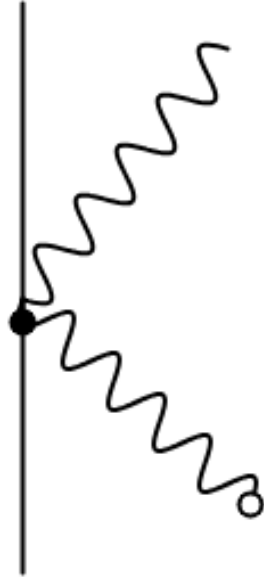
$$S_{\text{total}} = S_{\text{scalar}} + S_{\text{gravity}} + S_{\text{SET}} - \int dT e(x) \left[ M_0 + \sum_{\ell=1}^{\infty} \frac{\mu_{\ell}}{\ell!} \left( \partial_{(a_1} \dots \partial_{a_{\ell})} \varphi \right)^2 \right]$$

- **Solve KG equation perturbatively in the rate:**

$$\frac{1}{\sqrt{-g}} \partial_i \left( \sqrt{-g} g^{ij} \partial_j \varphi \right) = -\mu_{\ell} (-1)^{\ell} \partial^{(\ell)} \delta_D^{(3)}(x - \tilde{x})$$

$$g = \text{Minkowsky} + M' \delta g + \dots$$

$$\mu_{\ell} = \mu_{\ell}^{(0)} + M' \mu_{\ell}^{(1)} + \dots$$



# Point-particle EFT

- **Full action:**

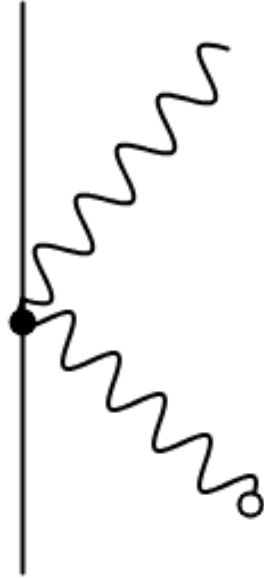
$$S_{\text{total}} = S_{\text{scalar}} + S_{\text{gravity}} + S_{\text{SET}} - \int dT e(x) \left[ M_0 + \sum_{\ell=1}^{\infty} \frac{\mu_{\ell}}{\ell!} \left( \partial_{(a_1} \dots \partial_{a_{\ell})} \varphi \right)^2 \right]$$

- **Solve KG equation perturbatively in the rate:**

$$\frac{1}{\sqrt{-g}} \partial_i \left( \sqrt{-g} g^{ij} \partial_j \varphi \right) = -\mu_{\ell} (-1)^{\ell} \partial^{(\ell)} \delta_D^{(3)}(x - \tilde{x})$$

$$g = \text{Minkowsky} + M' \delta g + \dots$$

$$\mu_{\ell} = \mu_{\ell}^{(0)} + M' \mu_{\ell}^{(1)} + \dots$$



- **Match couplings with other computation**

$$\varphi_{\text{resp}}^{(1)} \sim \mu_{\ell}^{(1)} x^{-\ell}$$

## CONCLUSIONS

- **Revisited perturbation of Vaidya metric with constant rate approximation**
- **Computed QNMs in the frequency domain**
- **Computed the linear Tidal Response**



*That's all Folks!*