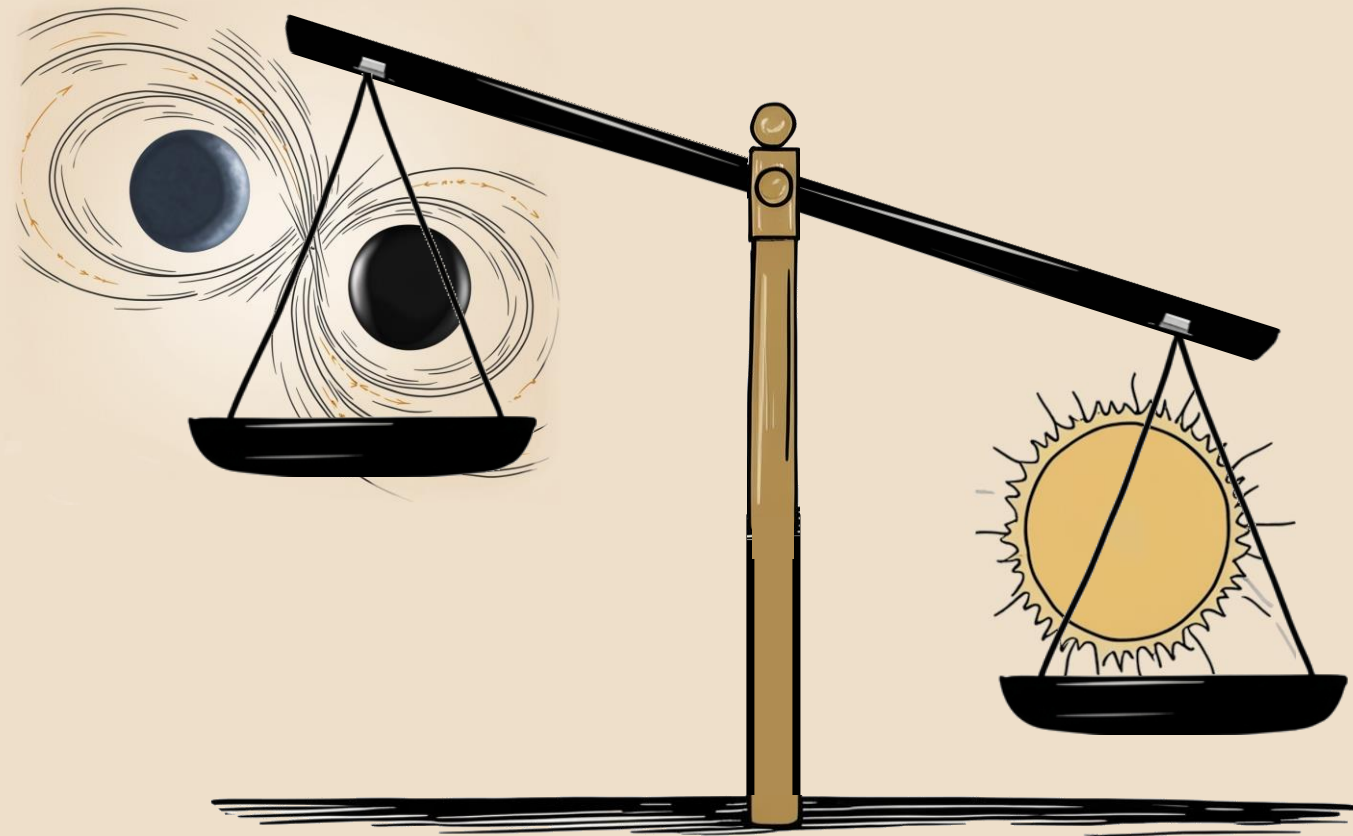


Gravitational waves from subsolar compact objects: implications for cosmology and high-density nuclear physics



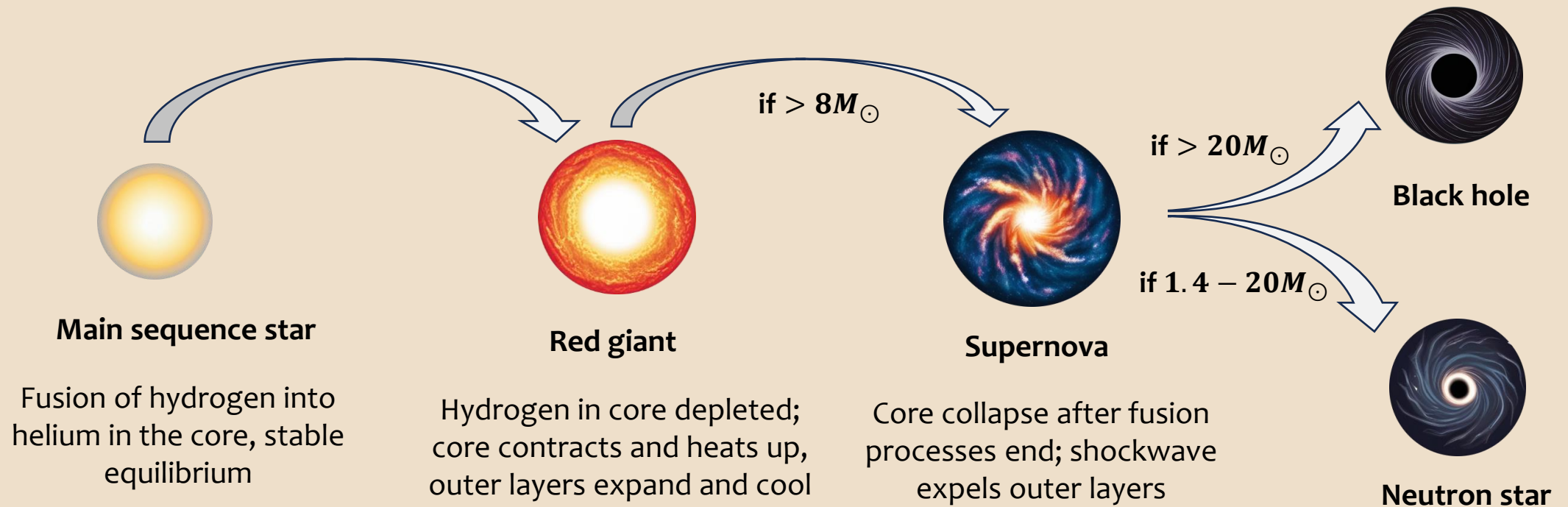
Massimo Vaglio 18/09/2024

1st TEONGRAV workshop on Gravitational Waves

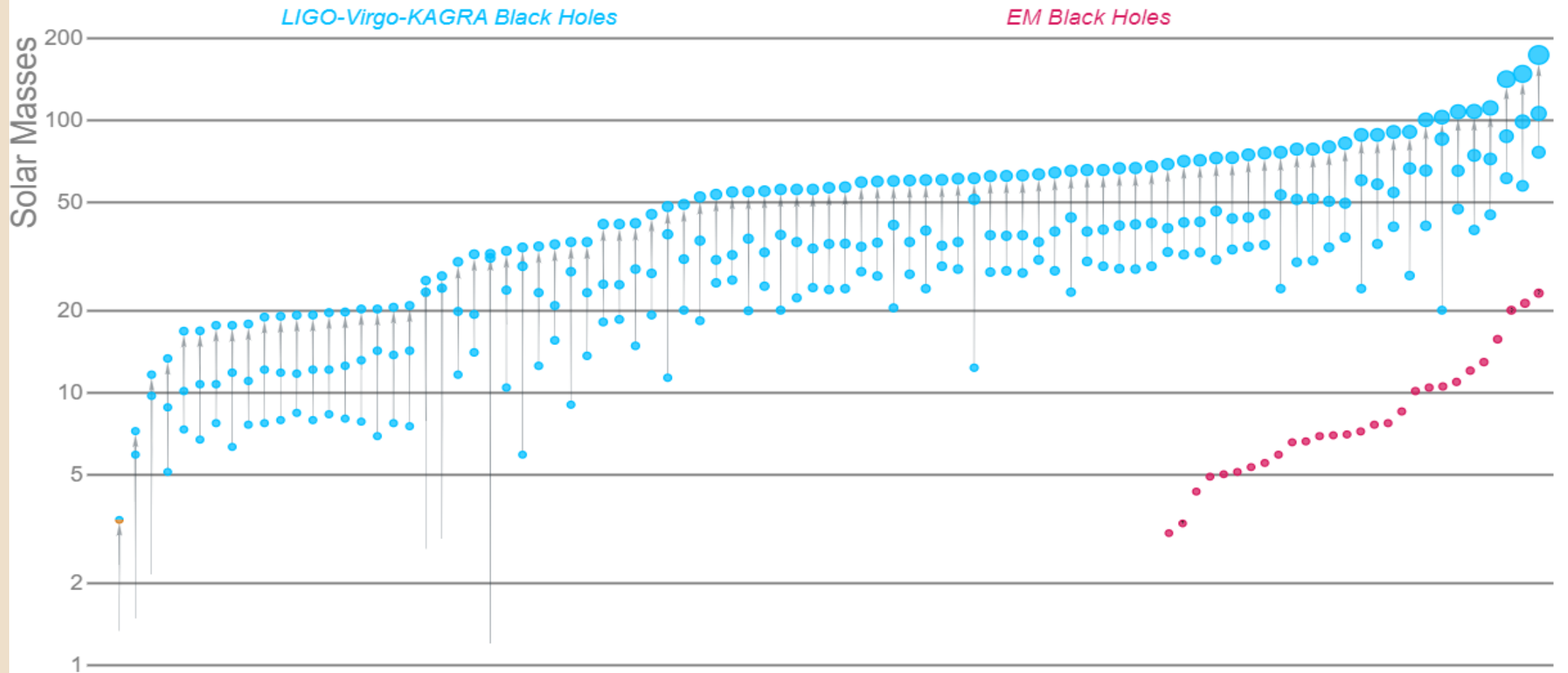
Based on: arXiv:2508.14287 – F.Crescimbeni, G.Franciolini, P.Pani, M.Vaglio

Why subsolar? Standard formation scenario of BH-NS

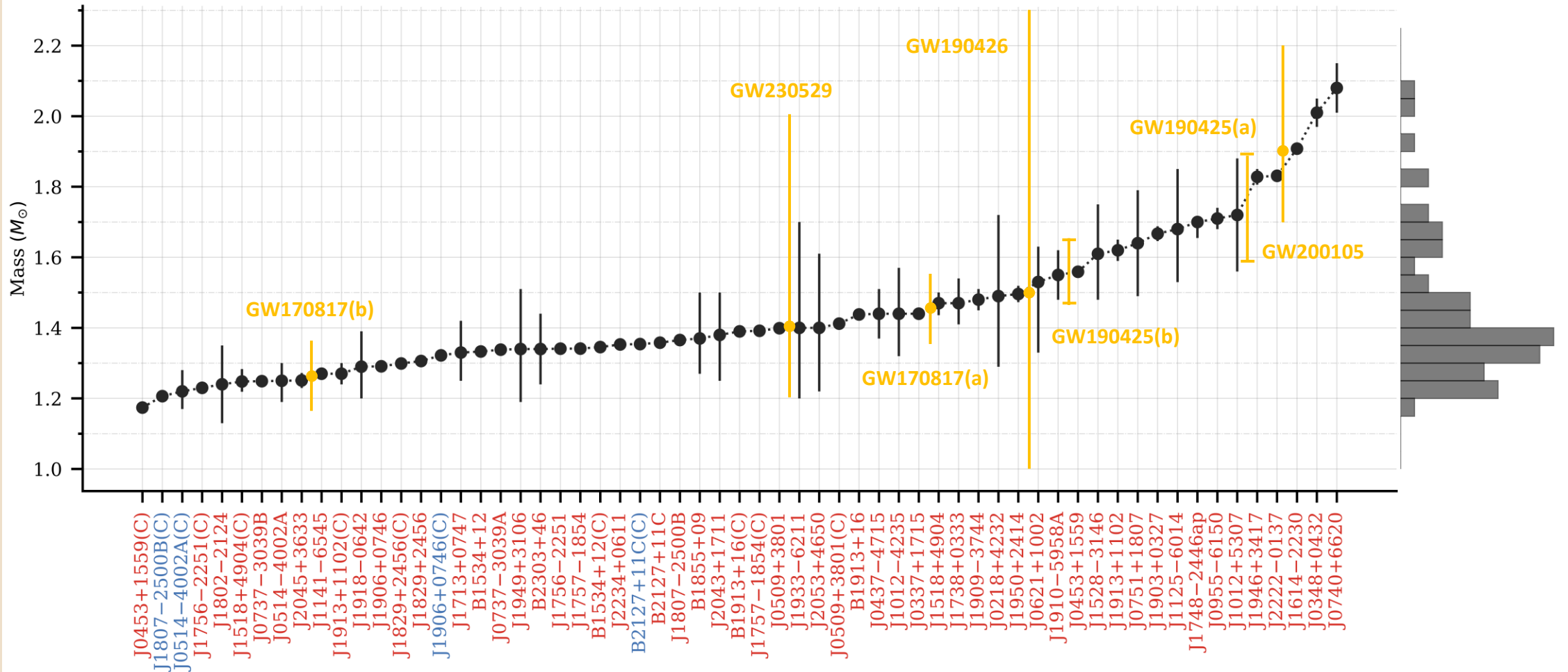
Conventional formation models suggest that black holes (BHs) and neutron stars (NSs), should have masses exceeding the solar mass



Mass distribution of black holes (EM+GW)



Mass distribution of neutron stars (binary pulsars + GW)



What would be the implications of a subsolar GW event?

Methodology & models

We performed parameter estimation on real and mock data and compute the Bayes factors to compare different models

$$p(\boldsymbol{\theta}|s) = \frac{\mathcal{L}(s|\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{\mathcal{Z}(s)}$$

Bayes formula

$$\mathcal{B}_{A/B} = \frac{\mathcal{Z}_A}{\mathcal{Z}_B}$$

Bayes factor

We used the public software BILBY [G. Ashton et al. (2019)] and a waveform template:

Frequency-domain TaylorF2

$$\tilde{h}(f, \boldsymbol{\theta}) = \mathcal{A}(f, \boldsymbol{\theta})e^{-i\phi(f, \boldsymbol{\theta})}$$

Quasicircular
Aligned spins
3.5 pN point particle
2pN spin corrections

Parameters: $\boldsymbol{\theta} = \{m_1, m_2, d_L, \theta, \phi, \theta_{JN}, \psi, t_c, \phi_c, \chi_1, \chi_2\}$

Analyzing the trigger event SSM200308

In the last concluded LVK observing run (O3b) three candidates of SSM binary BH events were identified, SSM200308 was the loudest.

IMRPhenomPv2						
Model	BH1	BH2	Agnostic	NS1	NS2	BS
$m_1 [M_\odot]$	$0.65^{+0.17}_{-0.15}$	$0.72^{+0.20}_{-0.17}$	$0.57^{+0.13}_{-0.10}$	$0.59^{+0.29}_{-0.08}$	$0.82^{+0.20}_{-0.14}$	$0.50^{+0.10}_{-0.07}$
$m_2 [M_\odot]$	$0.26^{+0.07}_{-0.04}$	$0.23^{+0.06}_{-0.04}$	$0.29^{+0.05}_{-0.05}$	$0.27^{+0.03}_{-0.08}$	$0.21^{+0.03}_{-0.03}$	$0.32^{+0.05}_{-0.05}$
χ_{eff}	$0.41^{+0.05}_{-0.04}$	$0.41^{+0.22}_{-0.05}$	$-0.13^{+0.08}_{-0.09}$	$0.15^{+0.16}_{-0.43}$	$0.72^{+0.07}_{-0.26}$	$0.36^{+0.25}_{-0.21}$
χ_p	$0.45^{+0.26}_{-0.26}$	-	-	-	-	-
$d_L [\text{Mpc}]$	80^{+37}_{-29}	83^{+41}_{-33}	97^{+45}_{-41}	110^{+139}_{-50}	76^{+37}_{-28}	106^{+84}_{-45}
$\Lambda_1 / 10^5$	-	-	-4^{+15}_{-10}	5^{+28}_{-3}	6^{+8}_{-5}	-
$\Lambda_2 / 10^7$	-	-	3^{+23}_{-12}	$1.3^{+0.6}_{-0.8}$	$0.3^{+0.3}_{-0.3}$	-
$\kappa_1 / 10^3$	-	-	15^{+347}_{-351}	-	-	-
$\kappa_2 / 10^3$	-	-	-287^{+114}_{-120}	-	-	-
$\log_{10} \tilde{\lambda}_f$	-	-	$-1.01^{+0.65}_{-0.42}$	-	-	-
$M_B [M_\odot]$	-	-	-	-	-	10^{+2}_{-2}
$\log_{10} \mathcal{B}$	-	0.31	-1.64	-2.68	0.22	-2.26

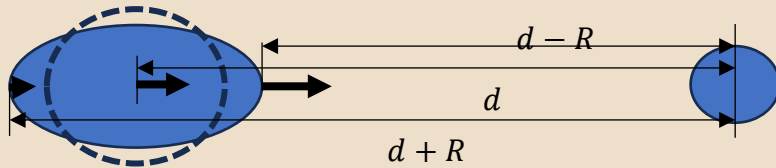
M. Prunier et al. (2023)

NS2 vs BH2

NS2 = BH2 + Tidal + Tapering

$$\theta_{NS2} = \theta_{BH2} \cup \{\Lambda_1, \Lambda_2\}$$

Tidal



$$g_{00} = -1 + \frac{2M}{r} + \frac{3Q_{ij}}{r^3} \left(n_i n_j - \frac{\delta_{ij}}{3} \right) + O\left(\frac{1}{r^4}\right) - \varepsilon_{ij} x_i x_j + O(r^3)$$

$$Q_{ij} = -\Lambda \varepsilon_{ij} \quad \text{Tidal deformability, 5pN \& 6pN}$$

Tapering

$$\tilde{h}(f) \rightarrow \tilde{h}(f) T(f|f_{\text{cut}}, f_{\text{slope}})$$

$$T(f|f_{\text{cut}}, f_{\text{slope}}) = \left[\frac{1 + e^{f_{\text{cut}}/f_{\text{slope}}}}{1 + e^{(f - f_{\text{cut}})/f_{\text{slope}}}} \right]$$

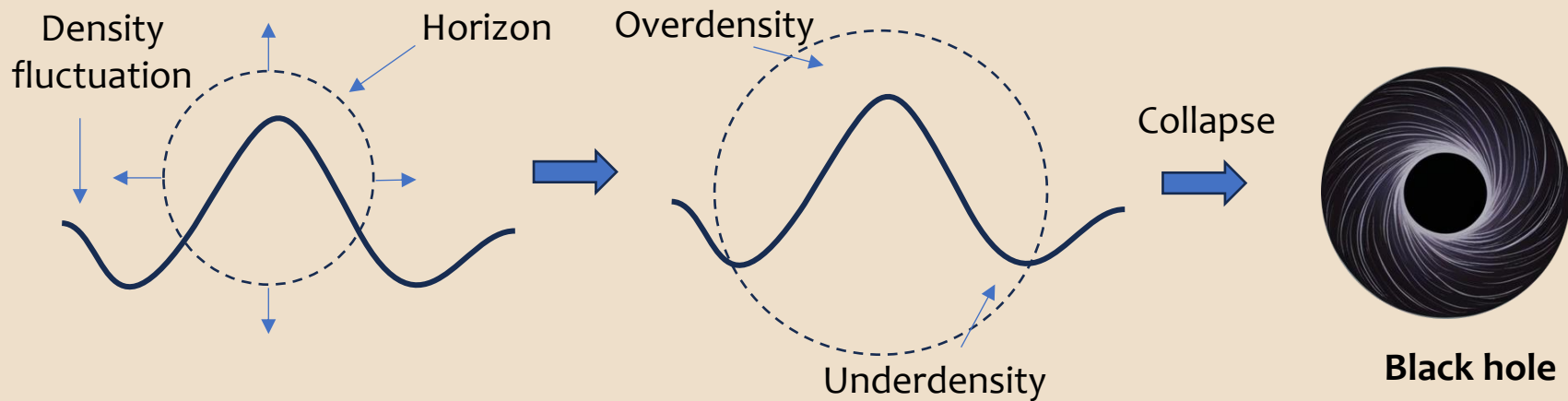
Tidal disruption

Detector	BH2 → NS2	NS2 → BH2
O4, SNR=15	-3.1	-2.9
O4, SNR=25	-5.1	-11.0

Implications for PBH: PBHs abundance

Primordial black holes are formed by collapse of overdense regions in in radiation dominated universe

$$M_{BH} \sim \frac{4\pi}{3} \rho_r H^{-3} \quad \text{PBH Mass} \sim \text{horizon mass at the formation}$$



Based on a model saturating the bounds from GWTC-3, to explain an event like SSM200308 one would need

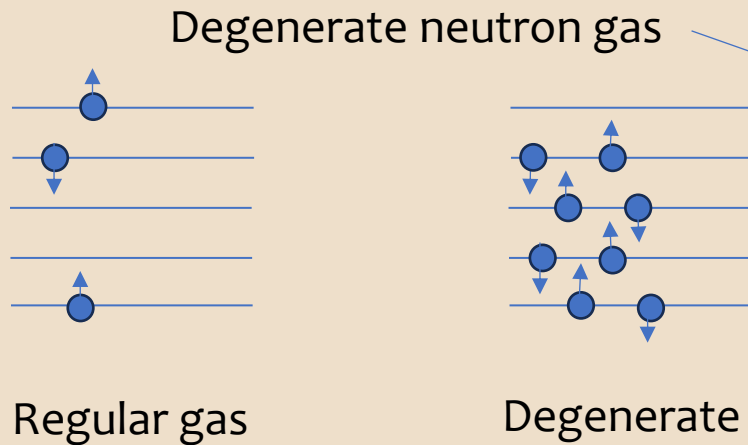
$$\frac{\rho_{PBH}}{\rho_{DM}} \equiv f_{PBH} \gtrsim \mathcal{O}(10^{-2})$$

NSs equation of state (EoS)

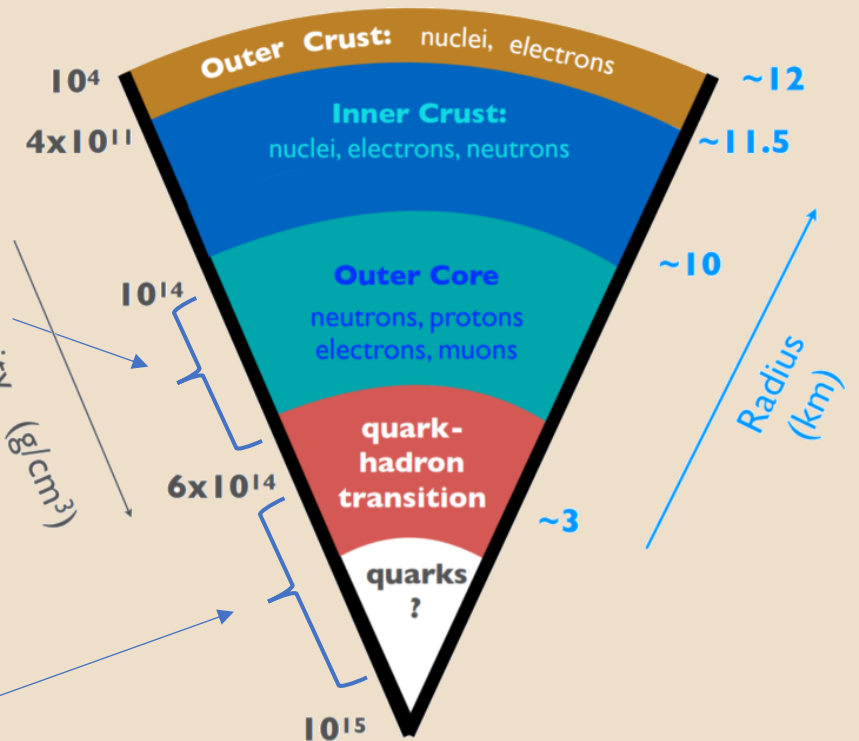
Neutron stars are modelled as relativistic perfect fluids

$$T_{\mu\nu} = (\epsilon + P)u_{\mu}u_{\nu} + P g_{\mu\nu}$$

$P(\epsilon)$ is the equation of state (EoS)



State of matter largely unknown at high densities



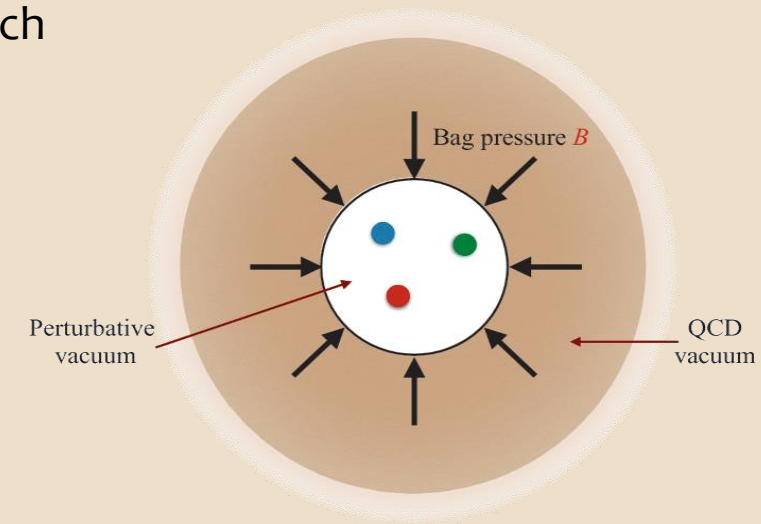
Strange quark-matter stars

Strange quark-matter models assumes a balance of up, down, and strange quarks.

EoS is obtained using perturbative QCD and an MIT-type bag model in which quarks are confined within a "bag" and interact via the strong force

J. M. Lattimer and M. Prakash (2001)

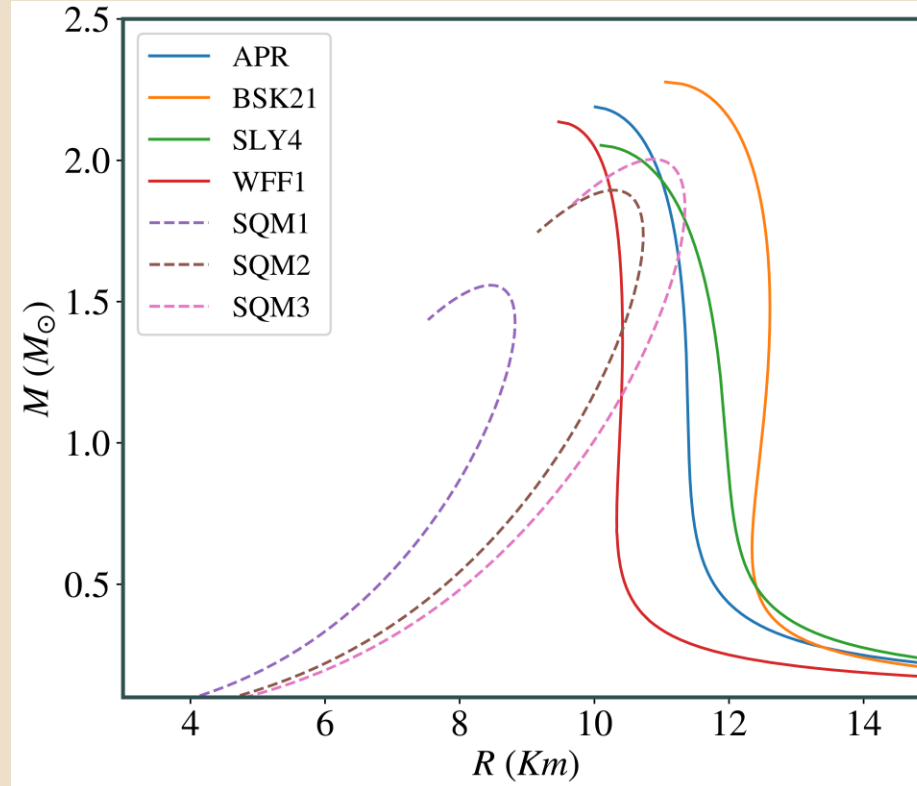
$$P(\epsilon) = \frac{\epsilon - B_{eff}}{3}$$



Features of Hadrons.
In: Journey to the Bound States.

This model often leads to a softer EoS compared to traditional neutron star matter because the pressure increase with density may be less steep

Implications for NSs: equation of state



EoS

APR, BSK21, SLY4, WFF1

SQM1-3

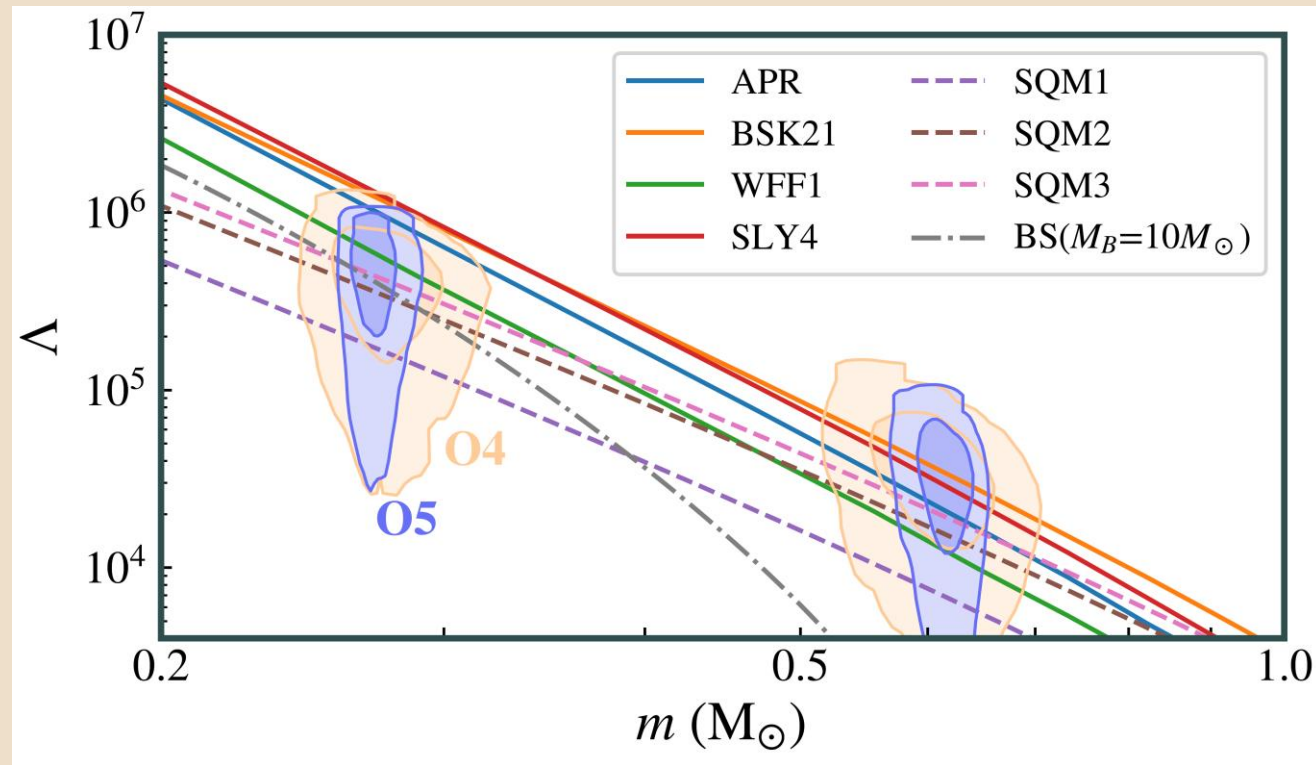
Strongly interacting components

Protons and neutrons

Quarks u,d,s

Implications for NSs: equation of state

SQM3 injected in O4 (O5) with SNR = 25 (SNR = 44)



The analysis is conservative: the inference considers Λ_1 and Λ_2 as independent

Implications for NSs: equation of state

Detectors	$m_1 [M_\odot]$	APR \rightarrow SQM3	SQM3 \rightarrow APR	WFF1 \rightarrow SQM3	SQM3 \rightarrow WFF1
O4, SNR = 25	0.63	-1.9	-3.8	0.1	-0.4
	0.27	-10.2	-19.9	-2.7	-5.0
O5, SNR = 44	0.63	-7.0	-12.3	-0.2	-1.0
	0.27	-37.5	-88.8	-11.3	-25.1

! With O4-O5 we will be able to distinguish between strange stars and neutron stars in subsolar range

!! Excluding WFF1 requires higher SNR or light comparable-mass SSM binaries

Conclusions



Already with current sensitivity one could decisively exclude some models of NSs and more exotic objects compared to the PBH hypothesis



In O4 this distinction would be crystal-clear, confirming or ruling out a putative PBH event



Using O4-O5 data it would be possible to use tidal-deformability measurements to decisively confirm/exclude light NSs with strange quark matter



This opportunity is provided by the fact that quark stars have a much smaller radius and tidal deformability than ordinary-EoS NSs in the SSM range, while overall, the tidal is very large