



### Scattering and dynamical capture of two black holes: synergies between numerical and analytical methods

S. Albanesi, A. Rashti, F. Zappa, R. Gamba, W. Cook, B. Daszuta, S. Bernuzzi, A. Nagar, and D. Radice

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#### Based on **2405.20398**:



## Introduction

- models for non-circularized orbits
- Focus on hyperbolic systems and their transition to bound orbits
  - scatterings  $\rightarrow$  dynamical captures
- Astrophysical motivations:
  - dense environments such as globular clusters and AGN
  - **GW190521** is compatible with a dynamical capture scenario [1]
- Theoretical motivations:
  - (elliptic) models

#### Detectable signals from binaries with residual eccentricity: need of GW

• Detection rates are highly uncertain, but populations of these systems are expected in

• Captures are the most delicate systems to model  $\rightarrow$  useful insights for eccentric

Post-Minkowskian descriptions for scatterings can be tested with NR simulations

[1] Gamba+:<u>2106.05575</u>









- Hyperbolic systems: initial positive binding energy ( $E_{ADM} > M$ )
- Some numerical relativity simulations, GR-Athena++ [2]
- Series of runs with  $E_0 = 1.011 M$  and decreasing angular mom  $p_{\omega}^0 \equiv J_0 / (\mu M)$









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![](_page_5_Picture_9.jpeg)

![](_page_5_Picture_10.jpeg)

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![](_page_6_Figure_4.jpeg)

![](_page_6_Picture_8.jpeg)

![](_page_6_Picture_9.jpeg)

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![](_page_7_Figure_4.jpeg)

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![](_page_8_Figure_4.jpeg)

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![](_page_9_Figure_4.jpeg)

![](_page_9_Picture_7.jpeg)

![](_page_9_Picture_8.jpeg)

# **Distinguishing scatterings and captures**

![](_page_10_Figure_4.jpeg)

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#### Same simulations as before with $E_0 = 1.011 M$ : scatterings, double encounters, and direct captures

![](_page_10_Picture_8.jpeg)

# Transition from unbound to bound

![](_page_11_Figure_1.jpeg)

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- Transition from scattering to merger
- Points/errors from scattering with lowest  $p_{\varphi}^{0}$  and merger with highest  $p_{\varphi}^{0}$
- How to distinguish scatterings from double encounters with huge apastra?
  - $\rightarrow$  check radiated energy\*
- Similar fits from GR-Athena++ and Kankani-McWilliams fit [3] but some differences
  - \* if computed from fluxes, require integration of  $\psi_4$

[3] Kankani,McWilliams 2404.03607

![](_page_11_Picture_11.jpeg)

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# **Orbital parameter space**

- How to sample/explore the orbital parameter space? • Eccentricity (typically) defined only for elliptic-like orbits dynamics-eccentricity cannot be defined in a gauge-invariant way • waveform-eccentricity cannot be generalized to single/short bursts

- Alternative: initial energy (ADM) and  $\mu$ -rescaled angular momentum ("ADM") of the system,  $(E_0, p_{\varphi}^0)$ 
  - used to sample the hyperbolic PS for the analysis of GW190521 as a dynamical capture [1]

#### Preliminary exploration of the PS using **TEOBResumS-Dalí** [4]

[1] Gamba+:2106.05575 [4] Nagar+:2009.12857

![](_page_12_Picture_12.jpeg)

![](_page_12_Picture_13.jpeg)

## Effective-one-body models

- •
- The effective metric is a  $\nu$ -deformation of Schwarzschild/Kerr, being  $\nu = m_1 m_2 / (m_1 + m_2)^2$  the symmetric mass ratio

Map from 2-body PN equations of motion to motion of a particle in an effective metric [5]

[5] Buonanno, Damour: gr-qc/9811091

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![](_page_13_Picture_8.jpeg)

![](_page_13_Figure_9.jpeg)

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## **Effective-one-body models**

- The effective metric is a  $\nu$ -deformation of Schwarzschild/Kerr, being  $\nu = m_1 m_2 / (m_1 + m_2)^2$  the symmetric mass ratio
- Three building blocks:
  - Hamiltonian:  $H_{\rm EOB} = M_{\rm V} / 1 + 2\nu (\hat{H}_{\nu}^{\rm eff} 1)$
  - Radiation reaction  $\mathcal{F}_{\varphi,r}$

 $\rightarrow$  solve Hamilton's equations

• Waveform  $h_{\ell m}^{\text{inspl}} = h_{\ell m}^{(\epsilon,N)_{c}} \hat{h}_{\ell m}^{(\epsilon,N)_{nc}} \hat{h}_{\ell m}^{c} \hat{h}_{\ell m}^{nc}$ 

 $h_{\ell m} = \theta(t - t_{\ell m}^{\text{match}}) h_{\ell m}^{\text{inspl}} \hat{h}_{\ell m}^{\text{NQC}} + \theta(t_{\ell m}^{\text{match}} - t) h_{\ell m}^{\text{ringdown}}$ 

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Map from 2-body PN equations of motion to motion of a particle in an effective metric [5]

non-spinning case, Schwarzschild  $\nu$ -deformation

1) 
$$\hat{H}_{\nu}^{\text{eff}} = \sqrt{A_{\nu} \left(1 + p_{\varphi}^{2} u^{2} + Q_{\nu}(r, p_{r_{*}})\right) + p_{r_{*}}^{2}}$$
$$A_{\nu} = 1 - 2u + 2\nu u^{3} + \dots$$

[5] Buonanno, Damour: gr-qc/9811091

![](_page_14_Picture_16.jpeg)

![](_page_14_Figure_17.jpeg)

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$$A_{\nu} = 1 - 2u + 2\nu u^3 + \dots$$

- What has to be changed for **non-circularized** inspirals?
  - Hamiltonian ok-ish modulo non-local terms and calibration
  - Radiation has to be generalized: generic **Newtonian prefactor** [6] and PN noncircular corrections

[5] Buonanno, Damour: gr-qc/9811091

[6] Chiaramello, Nagar: 2001.11736

![](_page_15_Picture_20.jpeg)

![](_page_15_Picture_21.jpeg)

![](_page_15_Picture_22.jpeg)

- Define effective potential as  $V = H_{EOB}(r)$
- Some equal mass non-spinning cases:

$$p_{\varphi}, p_{r_*} = 0; \nu$$

![](_page_16_Picture_6.jpeg)

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- Some equal mass non-spinning cases:

![](_page_17_Figure_3.jpeg)

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![](_page_17_Figure_6.jpeg)

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![](_page_18_Figure_3.jpeg)

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![](_page_18_Figure_6.jpeg)

![](_page_18_Picture_7.jpeg)

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![](_page_19_Figure_3.jpeg)

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![](_page_19_Figure_6.jpeg)

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![](_page_20_Figure_3.jpeg)

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![](_page_20_Figure_6.jpeg)

![](_page_20_Picture_7.jpeg)

# Parameter space: TEOBResumS-Dalí

![](_page_21_Figure_1.jpeg)

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peaks)

- Number of orbital frequency peaks Nas a proxy for the number of encounters
- Parameter space dominated by scatterings and direct captures (N = 1)
- Many encounters if the initial energy is close to the parabolic limit ( $E_0 = M$ ). **Double encounters** also permitted for higher energies: transition from scattering to capture
- Rich phenomenology produced by Dalí, but how accurate is this picture?

![](_page_21_Picture_7.jpeg)

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# Parameter space: TEOBResumS-Dalí and NR

(number of

 $\geq$ 

![](_page_22_Figure_1.jpeg)

- Numerical relativity simulations performed with GR-Athena++
- EOB/NR phenomenologies agrees for  $E_0 \lesssim 1.02 \ M (v_{\rm cm} \lesssim 0.2)$
- Dalí can **continuously span** the parameter space from quasi-circular binaries to hyperbolic systems
- EOB/NR deviations increase with the for higher energy/velocity (cfr. NR fits of the scattering-capture transition vs EOB one)
- Dalí built on PN results and informed only on quasi-circular NR simulations: large room for improvement

![](_page_22_Picture_9.jpeg)

![](_page_22_Picture_11.jpeg)

# Scatterings

- Scattering angles computed extrapolating relative tracks
- Each color corresponds to an energy series: markers for NR, lines for EOB
- Scattering-capture transition marked by vertical bands (NR) and dashed lines (EOB)
- As before, the EOB/NR disagreement increases for higher energies
- Agreement restored for high angular momentum (weak field)
- We also considered scatterings of unequal mass binaries (not shown in this plot)

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![](_page_23_Figure_8.jpeg)

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![](_page_23_Picture_10.jpeg)

### **EOB/NR mismatches**

![](_page_24_Figure_1.jpeg)

- (2,2) mismatches useful to quantify the accuracy of the model
- Configurations with  $E_0 \lesssim 1.02 \ M$
- Initial data optimization on a small region
- Nonspinning and spin-aligned  $(|\chi_i| = 0.5)$  equal mass + higher mass ratios up to q = 3 (nonspinning)
- Below/around 1% for most cases, some around 3%
- NR waveforms are far from being perfect! Issues when integrating the Weyl scalars. CCE / metric-perturbation-extraction?
- In the paper we also studied energetic curves

![](_page_24_Picture_10.jpeg)

![](_page_24_Picture_11.jpeg)

# Test-mass: useful insights

- We studied a few test-mass cases: dynamics driven by a PN-radiation reaction • Numerical waveform obtained by solving RWZ eqs. with RWZHyp [7,8]

![](_page_25_Figure_3.jpeg)

[7] Bernuzzi, Nagar: <u>1003.0597</u> [8] Bernuzzi, Nagar, Zenginoglu: <u>1107.5402</u>

![](_page_25_Picture_6.jpeg)

![](_page_25_Picture_7.jpeg)

# Test-mass: useful insights

![](_page_26_Figure_1.jpeg)

waveform associated to the last four close passages of the red trajectory

- We studied a few test-mass cases: dynamics driven by **PN-radiation reaction**
- Numerical waveforms obtained by solving RWZ eqs. in the time-domain

![](_page_26_Figure_6.jpeg)

- Different EOB waveforms computed on the same dynamics
- Noncircular corrections with explicit derivatives (not PN-expanded in the Newtonian contribution) seem more reliable. In particular, they catch the low-frequency signal generated at large separations

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# Conclusions

- EOB models can be used to describe these systems, with some caveats: avoid explicit dependance on eccentricity

  - low-energy regime ( $E_0 \leq 1.02 M$ ) or large angular momentum (weak-field) not too-high spins (to quantify), inherited from QC case phenomenological ringdown model still calibrated on QC data
- Next steps to increase/extend EOB accuracy:
  - inclusion of post-Minkowskian results in the dynamical sector (Lagrangian EOB?)
  - NR-information from non-circularized binaries (e.g. ringdown)
- Test-mass limit always useful to gain insights on EOB models, but also on full-NR simulations

#### Exploration of hyperbolic systems with NR and TEOBResumS-Dalí

![](_page_27_Picture_14.jpeg)

![](_page_27_Picture_15.jpeg)

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#### Exploration of hyperbolic systems with NR and TEOBResumS-Dalí

#### **Thank you for your attention!**

![](_page_28_Picture_16.jpeg)

![](_page_28_Picture_17.jpeg)

### **Backup slides**

![](_page_29_Picture_2.jpeg)

# **Analytical corrections to radiation**

- How can we describe noncircularized binaries?
  - 1. Hamiltonian: already generic, nothing to do (modulo calibration/non local terms) 📀
  - 2. Radiation reaction  $(\mathcal{F}_{\varphi}, \mathcal{F}_{r})$ : noncircular corrections in  $\mathcal{F}_{\varphi}$
  - 3. Waveform: noncircular corrections in each multipole
- Noncircular terms:
  - $\mathcal{F}_r$  from generic energy/momentum balance eqs. [6]

  - Extended noncircular corrections up to 2PN [9-12]: improvement of the waveform
  - [6] Bini, Damour: 1210.2834 [7] Chiaramello, Nagar : 2001.11736
  - [8] Albanesi+ : 2104.10559
  - [9] Khalil+:2104.11705

- [10] Placidi+ : 2112.05448
- [11] Albanesi+ : 2202.10063
- [12] Albanesi+ : 2203.16286

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• Generic "Newtonian" prefactor [7,8], e.g.  $\hat{h}_{22}^{(N,0)_{nc}} = 1 - \frac{\dot{r}^2}{2r\Omega^2} - \frac{\dot{r}^2}{2r\Omega^2} + \frac{2i\dot{r}}{r\Omega} + \frac{i\Omega}{2\Omega^2}$  (crucial thing:  $\dot{r}$ ,  $\dot{\Omega}$ ,...

$$h_{\ell m}^{\text{inspl}} = h_{\ell m}^{(\epsilon, N)_{\text{c}}} \hat{h}_{\ell m}^{\text{c}} \hat{h}_{\ell m}^{(\epsilon, N)_{\text{nc}}} \hat{h}_{\ell m}^{\text{nc}},$$
$$\mathcal{F}_{\varphi} = -\frac{32}{5} \nu r_{\Omega}^{4} \Omega^{5} \hat{f}_{\text{nc}_{22}},$$

$$\hat{h}_{\ell m}^{\text{nc}} = \hat{h}_{\ell m}^{\text{inst}} \hat{h}_{\ell m}^{\text{hered}}$$
$$\hat{f}_{\text{nc}_{22}} = \hat{F}_{22} \hat{f}_{\varphi,22}^{\text{N}_{\text{nc}}} + \hat{F}_{21} + \sum_{\ell=3}^{8} \ell = 3$$

![](_page_30_Picture_21.jpeg)

![](_page_30_Picture_22.jpeg)

![](_page_30_Picture_23.jpeg)

### Test: integration of test-mass waveforms

![](_page_31_Figure_1.jpeg)

![](_page_31_Picture_3.jpeg)

# **NR: integration from** $\psi_4$

![](_page_32_Figure_1.jpeg)

![](_page_32_Picture_3.jpeg)

# **EOB/NR time-domain comparison**

![](_page_33_Figure_1.jpeg)

- Initial conditions in an ideal world: same  $(E_0/M, J_0/M^2)$  for both EOB/NR,  $r_0^{\text{EOB}}$ from EOB/ADM 2PN coords transformation
- Practical issues, e.g. junk-radiation, introduce small variations and may lead to completely different phenomenologies
- Optimize EOB ICs to minimize unfaithfulness (mismatch), conceptually similar to what is done for elliptic binaries
- NR calibration on QC simulations (also for the ringdown)
- Artifacts in the NR waveform due to integration

![](_page_33_Figure_9.jpeg)

![](_page_33_Figure_10.jpeg)

![](_page_33_Figure_11.jpeg)

![](_page_33_Picture_12.jpeg)

![](_page_33_Picture_13.jpeg)

## Potentials for QC and elliptic cases

Evolution of Schwarzschild potentials under the effect of EOB radiation reaction

![](_page_34_Figure_2.jpeg)

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![](_page_34_Figure_5.jpeg)

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