

Compact objects in and beyond the Standard Model



Loris Del Grosso

Gravity theory group

<https://web.uniroma1.it/gmunu>



EuCAPT



in collaboration with **G. Franciolini, P. Pani and A. Urbano**

2015: a discovery that shook the world

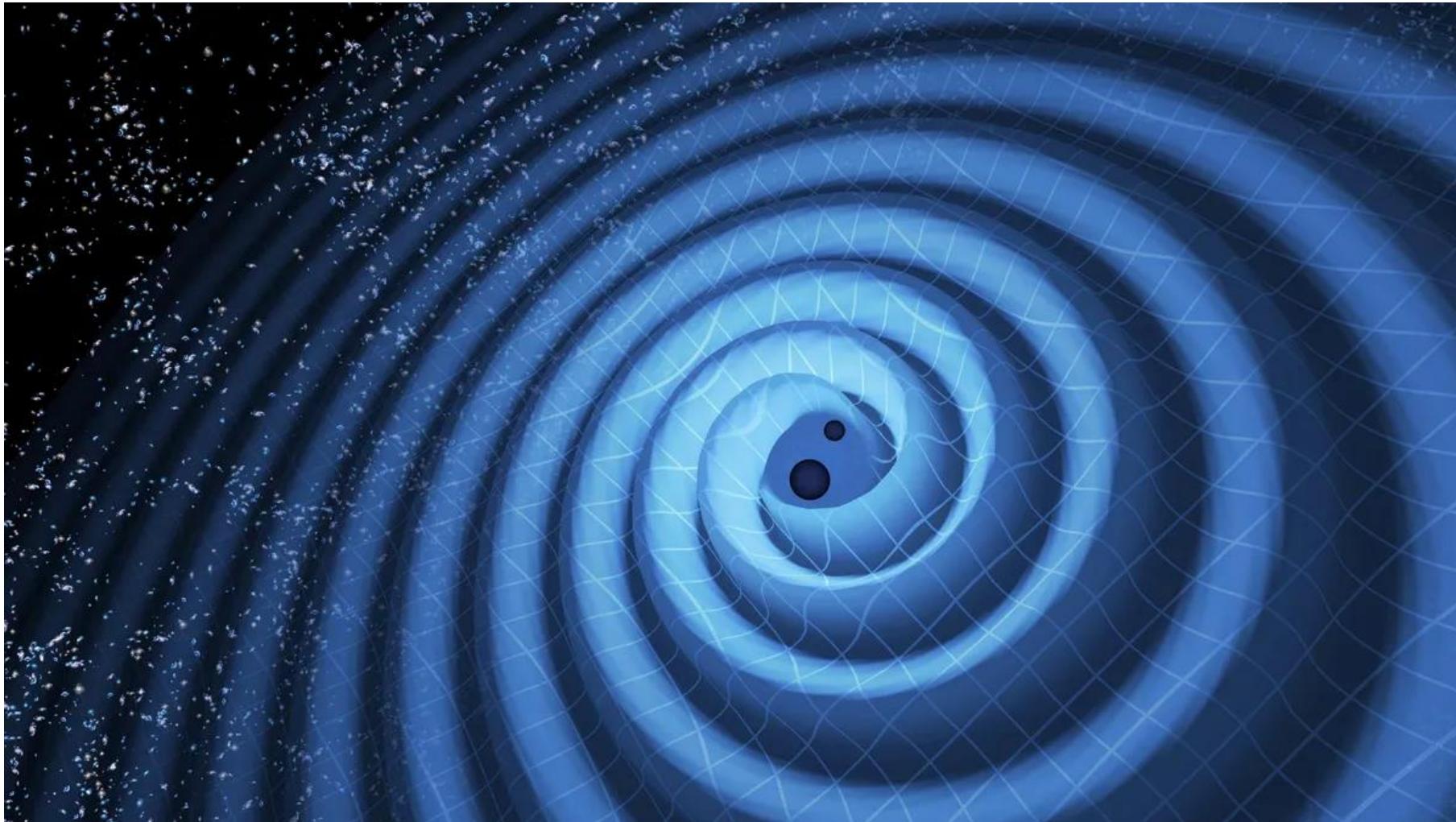


Image credit: LIGO/T. Pyle

Introduction

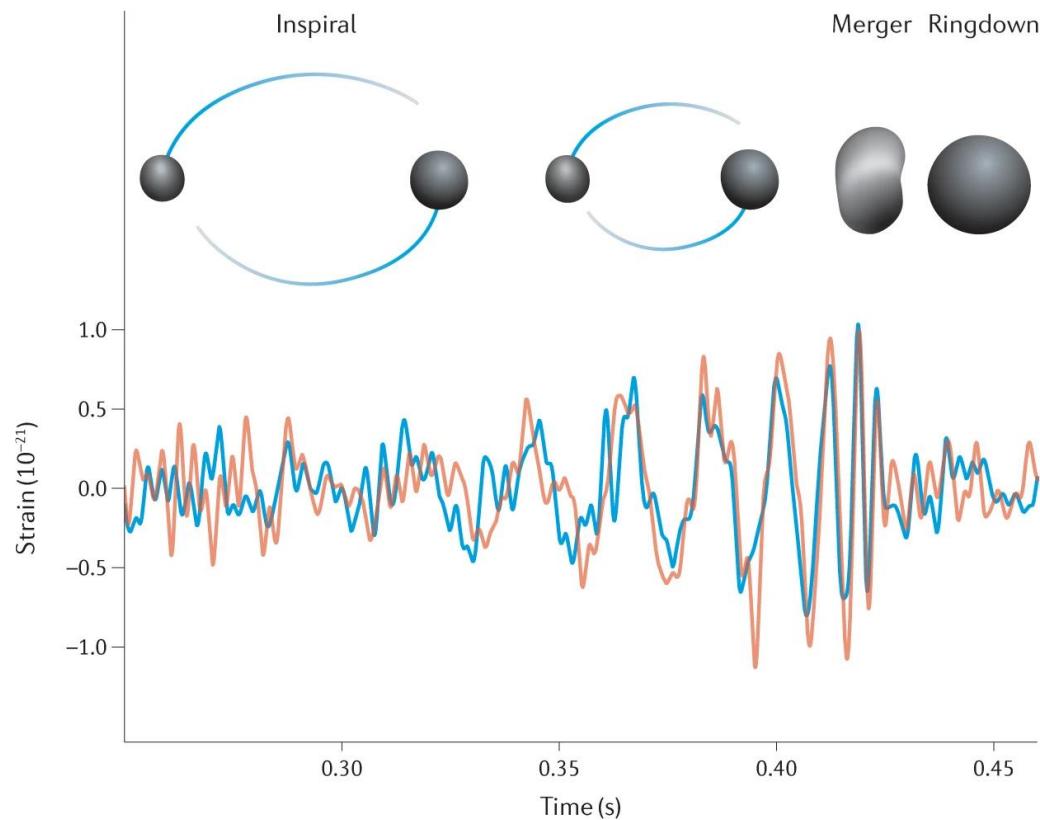


Image credit: M. Bailes et al., "Gravitational-wave physics and astronomy in the 2020s and 2030s," Nature Rev. Phys., 2021.

Introduction

The simplest source are binaries

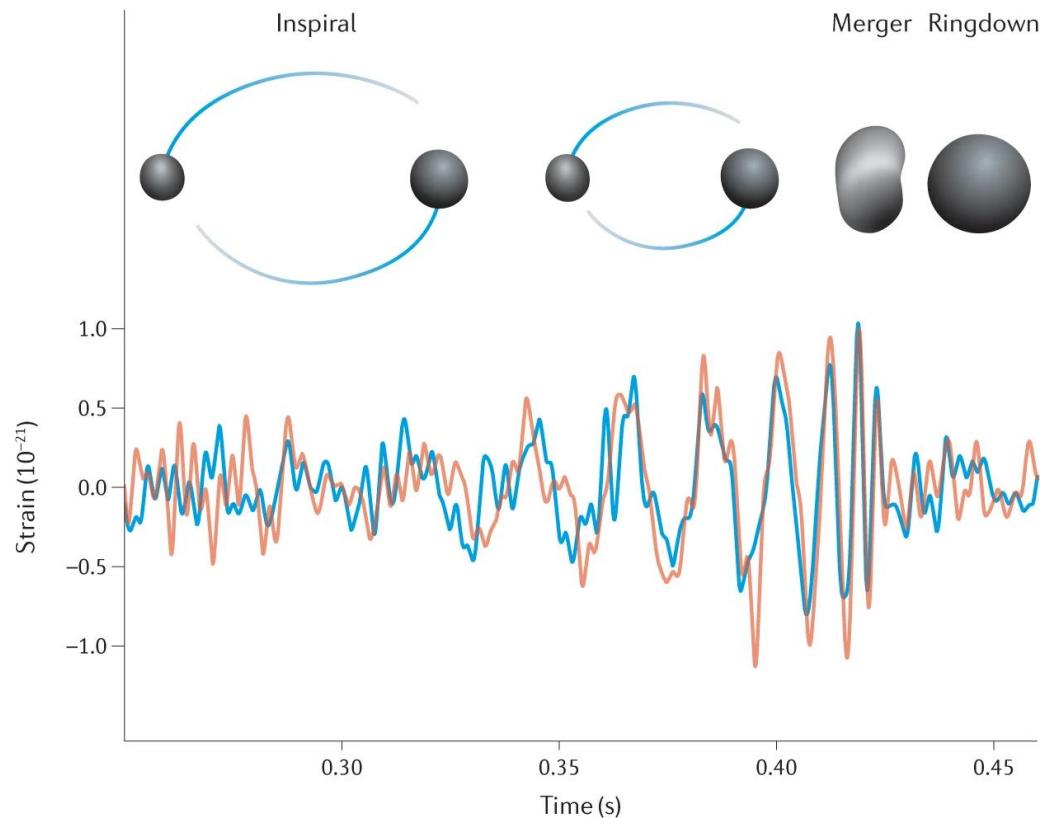


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Introduction

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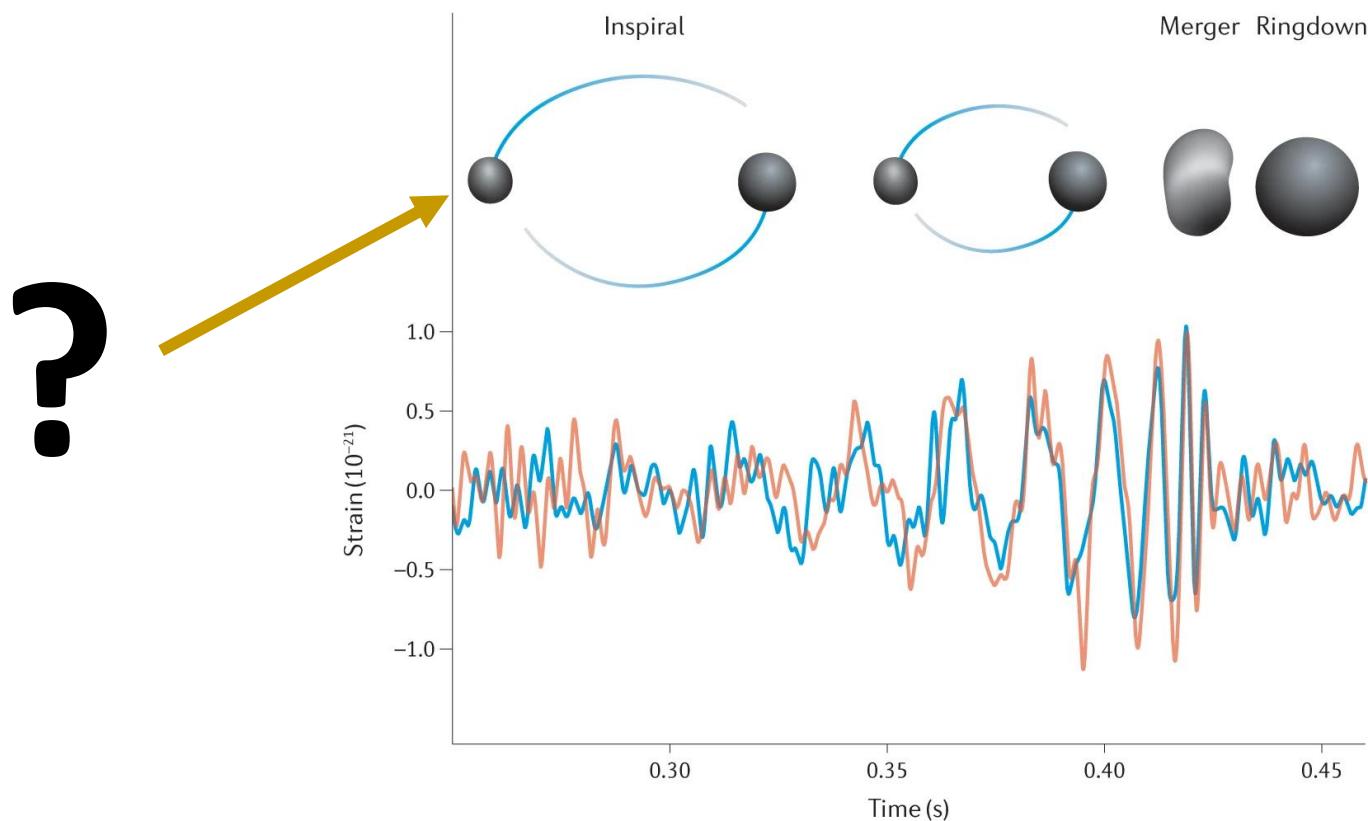
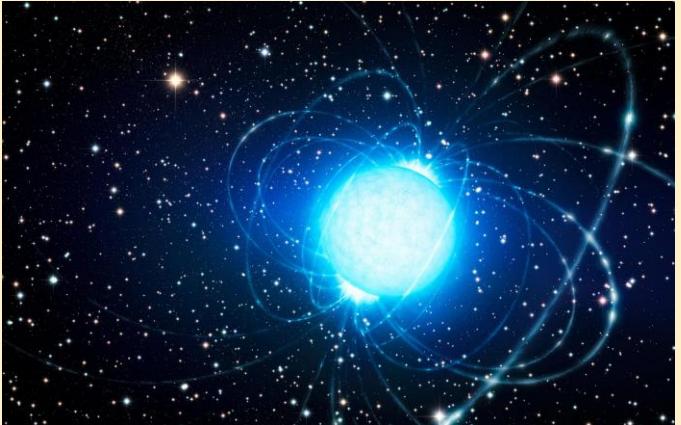


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Introduction

Neutron stars



Introduction

Neutron stars



$\sim (0.1 - 3) M_{\odot}$

Colpi, Shapiro, Teukolsky, "A Hydrodynamical Model for the Explosion of a Neutron Star Just below the Minimum Mass", *Astrophysical Journal* v.414, p.717, 1993.

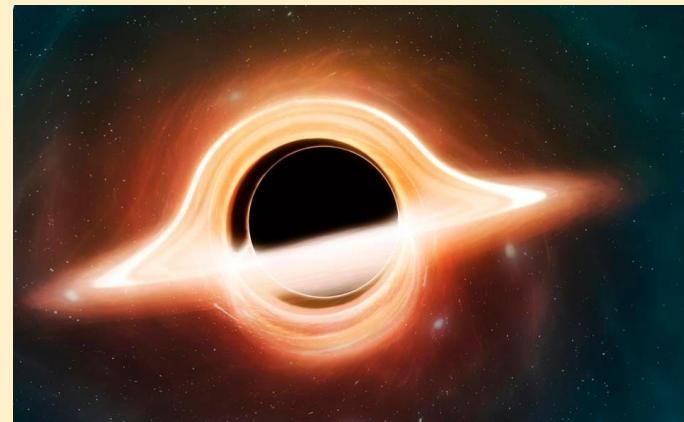
Introduction

Neutron stars



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Black holes



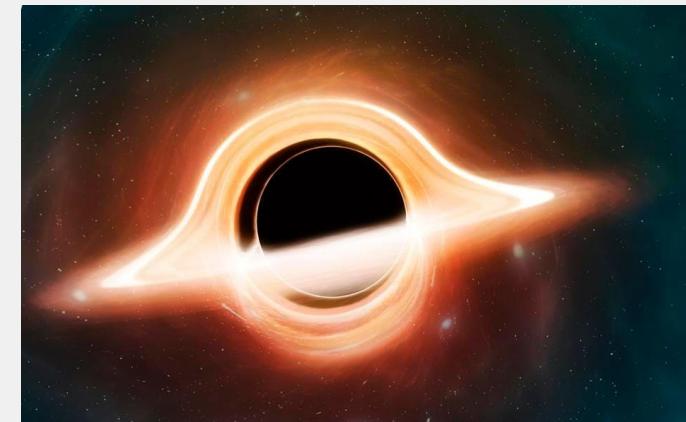
Introduction

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$\sim (0.1 - 3) M_{\odot}$

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$\sim (5 - 10^{10}) M_{\odot}$



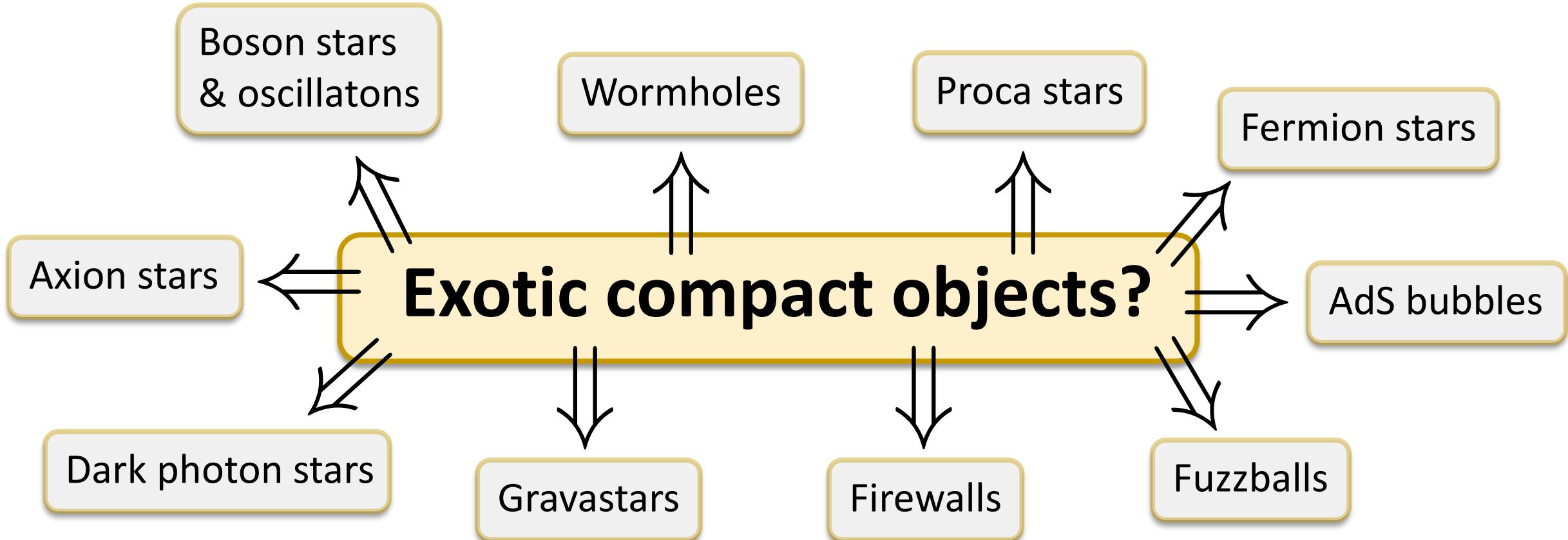
Exotic compact objects?

- Dark matter \Rightarrow new stable particles?

Exotic compact objects?

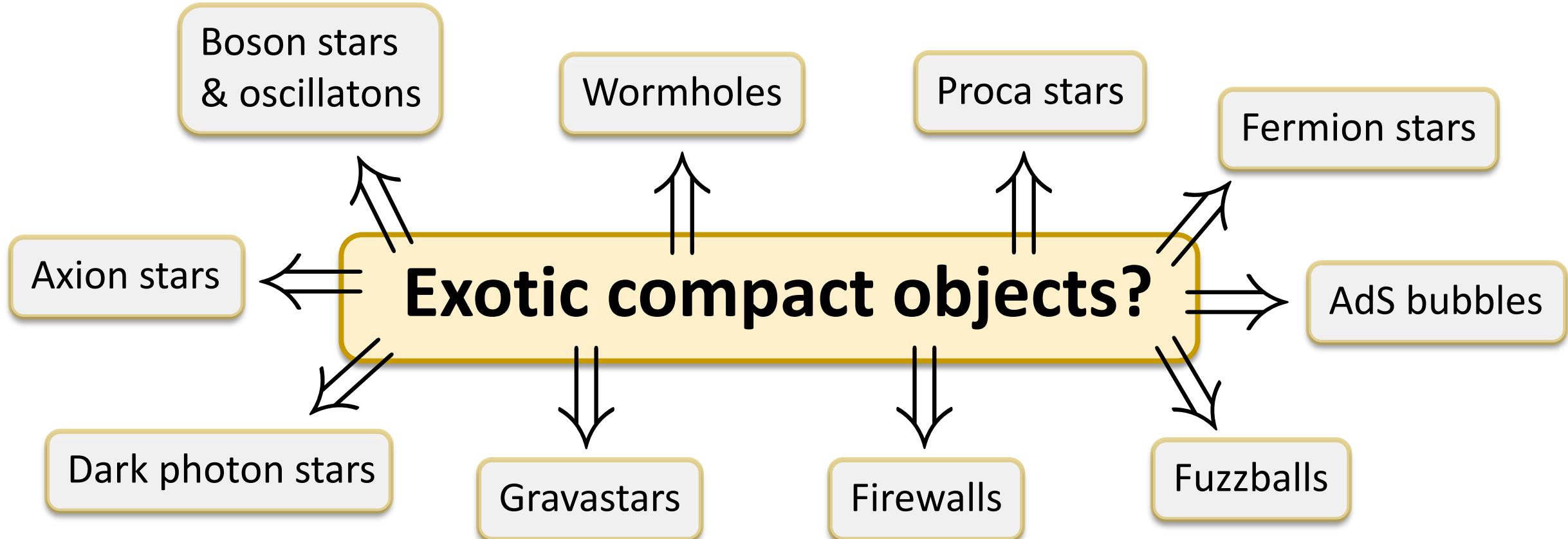
- Dark matter \Rightarrow new stable particles?
- New fields \Rightarrow ECOs arise

The ECO atlas

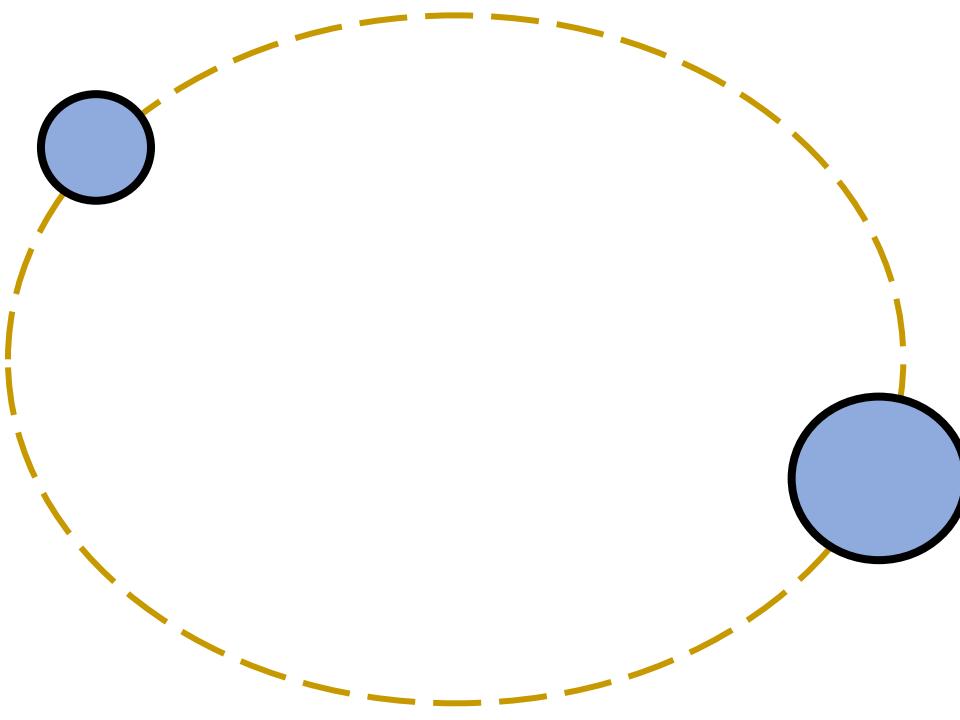


The ECO atlas

[Kaup, 1968; Ruffini, Bonazzola, 1969; Colpi, et. al., 1986]

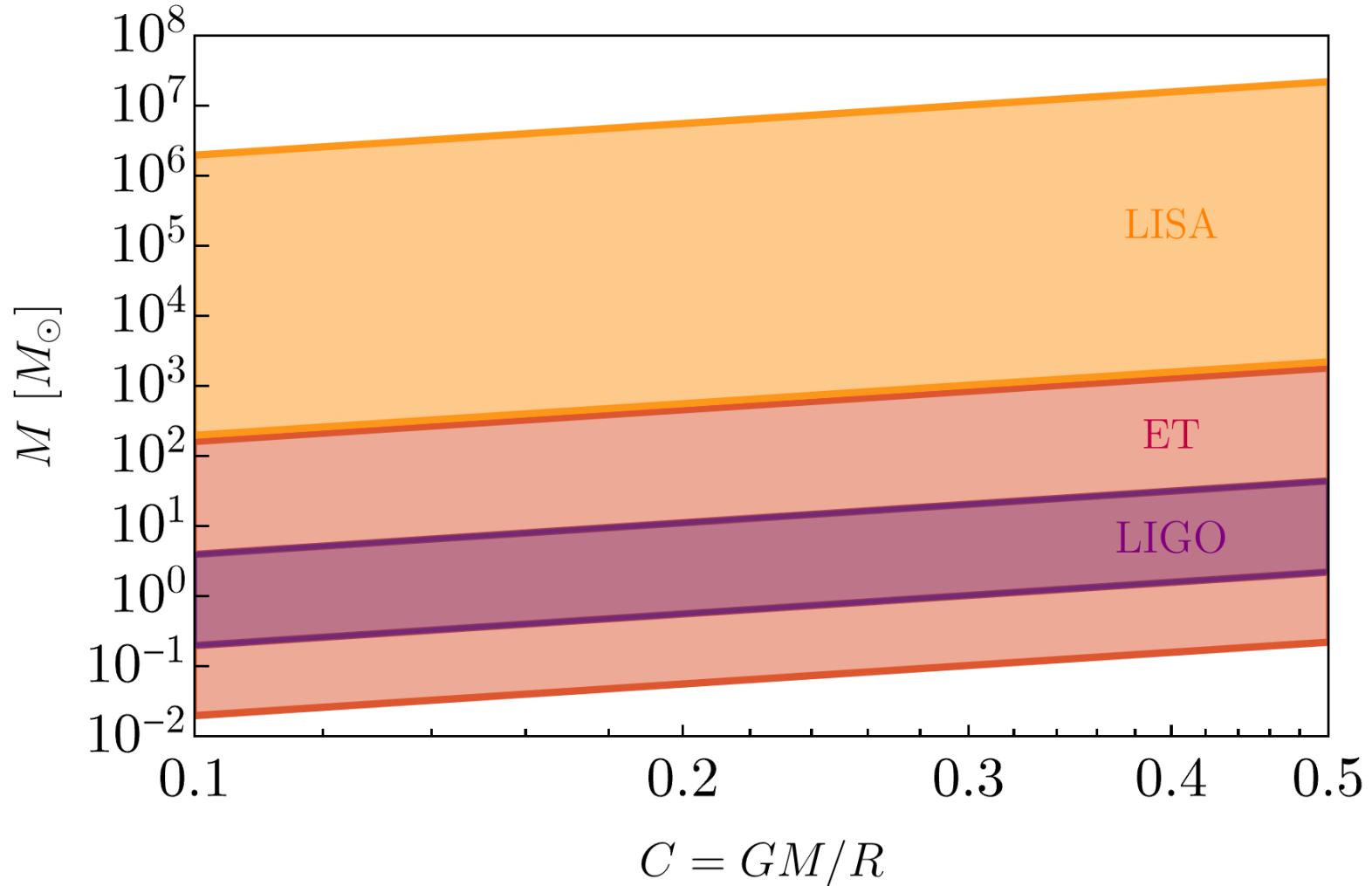


Testing ECOs with gravitational waves



$$f_{\text{ISCO}}^{\text{ECO}} = \frac{C^{3/2}}{3^{3/2} \pi G M_{\text{tot}}}$$

Testing ECOs with gravitational waves



Setup

$$\mathcal{L} = \frac{R}{16\pi G} - \frac{1}{2}\partial^\mu h \partial_\mu h - \frac{\lambda}{16}(h^2 - v^2)^2 - \bar{\psi}\gamma^\mu D_\mu \psi - \frac{f}{\sqrt{2}}h \bar{\psi}\psi$$

[Lee, Pang, 1987]

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[Lee, Pang, 1987]

- It is renormalizable (cf. solitonic boson stars)

[R. Friedberg, T. D. Lee, and Y. Pang, 1987]

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- It is local \Rightarrow causality respected

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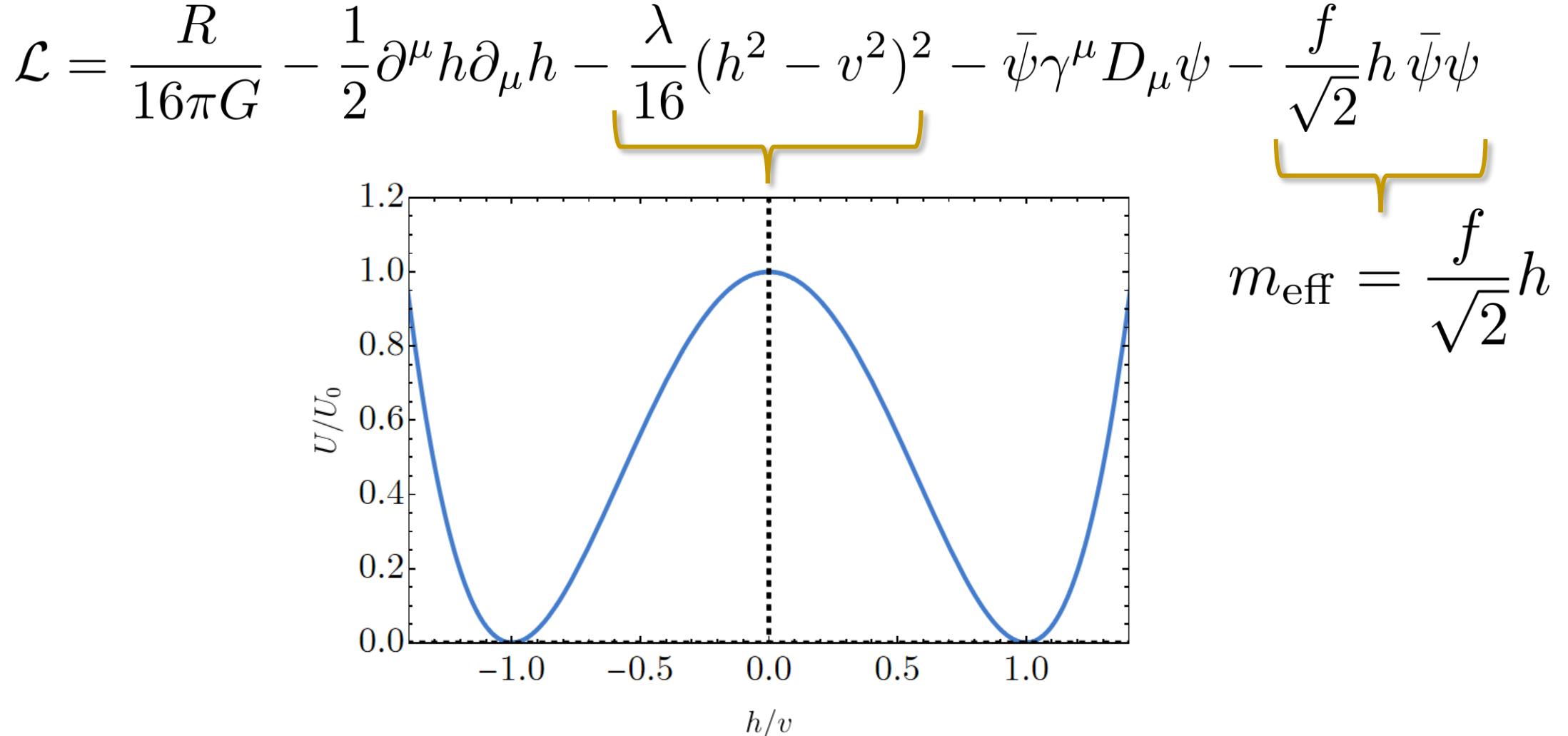
- It is local \Rightarrow causality respected
- Arises in several BSM extensions

Setup

$$\mathcal{L} = \frac{R}{16\pi G} - \frac{1}{2}\partial^\mu h \partial_\mu h - \frac{\lambda}{16}(h^2 - v^2)^2 - \bar{\psi} \gamma^\mu D_\mu \psi - \frac{f}{\sqrt{2}}h \bar{\psi} \psi$$

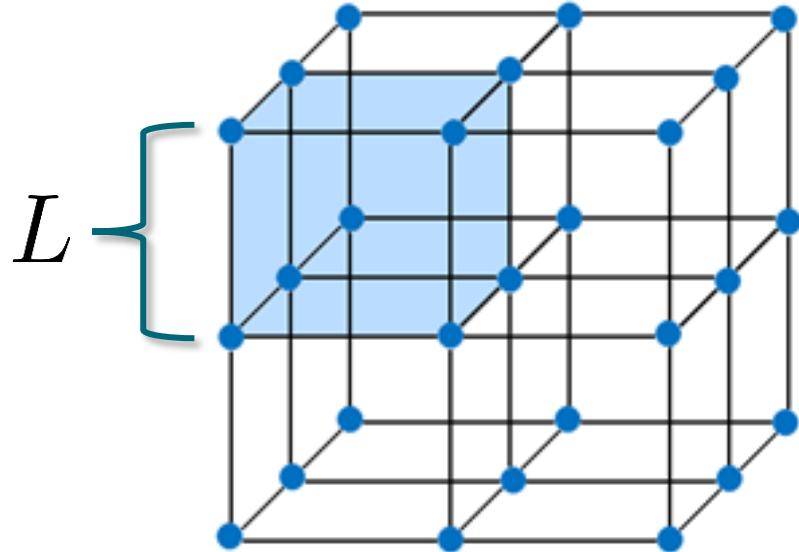
$$m_{\text{eff}} = \frac{f}{\sqrt{2}}h$$

Setup



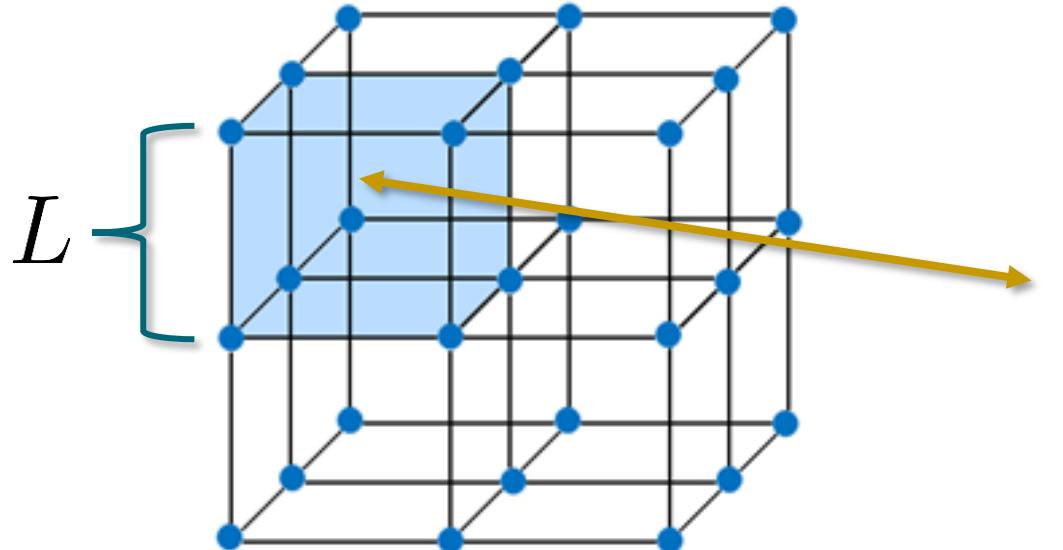
Thomas-Fermi approximation

- Devide the three-space into small cubes $L_{g_{\mu\nu}, h} \gg L \gg \lambda_B$



Thomas-Fermi approximation

- Devide the three-space into small cubes $L_{g_{\mu\nu},h} \gg L \gg \lambda_B$
- Fill each cube with a degenerate Fermi gas of Fermi momentum



- $W[k_F]$ (Energy density)
- $P[k_F]$ (Pressure)
- $S[k_F] = \langle \bar{\psi} \psi \rangle$ (Scalar density)

Equations of motion

$$1. \quad G_{\mu\nu} = 8\pi G \left(\partial_\mu h \partial_\nu h - \frac{1}{2} g_{\mu\nu} (\partial^\alpha h \partial_\alpha h + 2U) + (W + P) u_\mu u_\nu + Pg_{\mu\nu} \right)$$

$$= T_{\mu\nu}^{[h]}$$

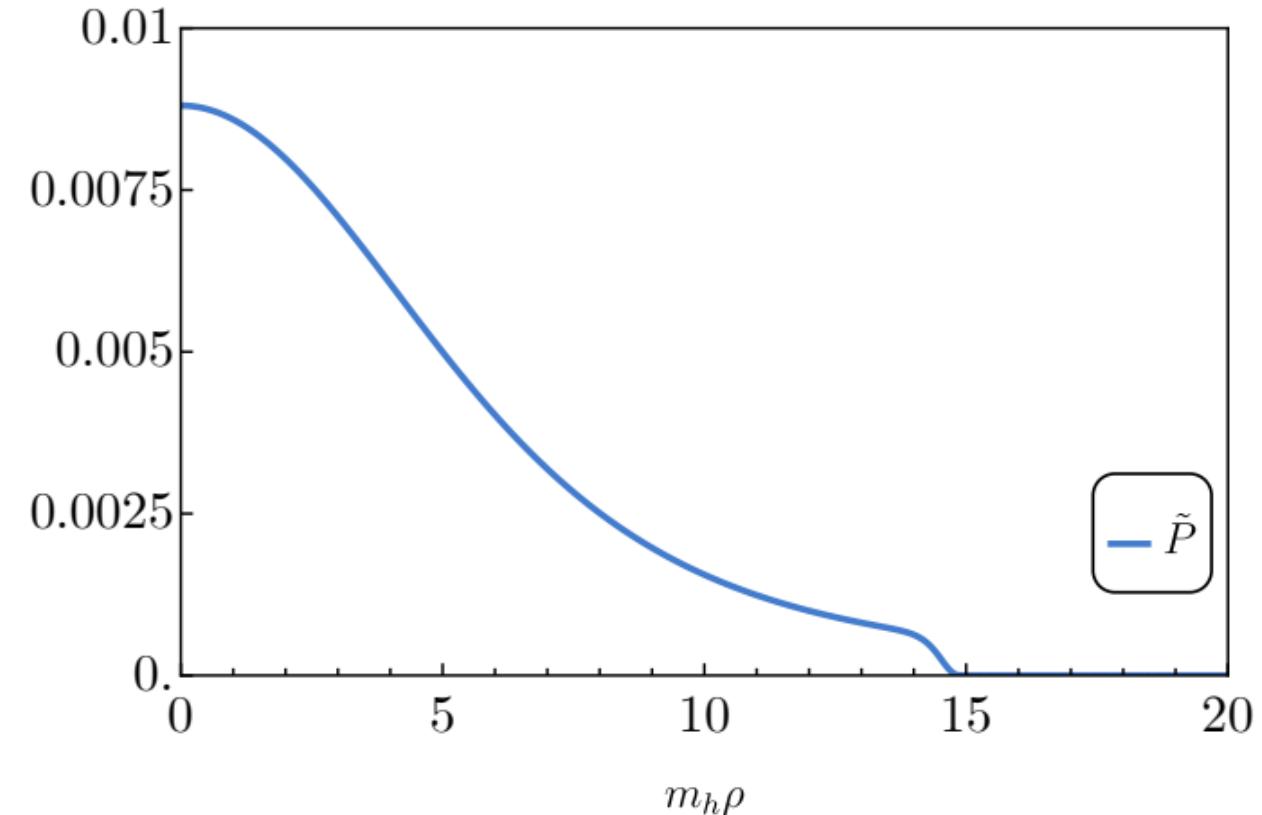
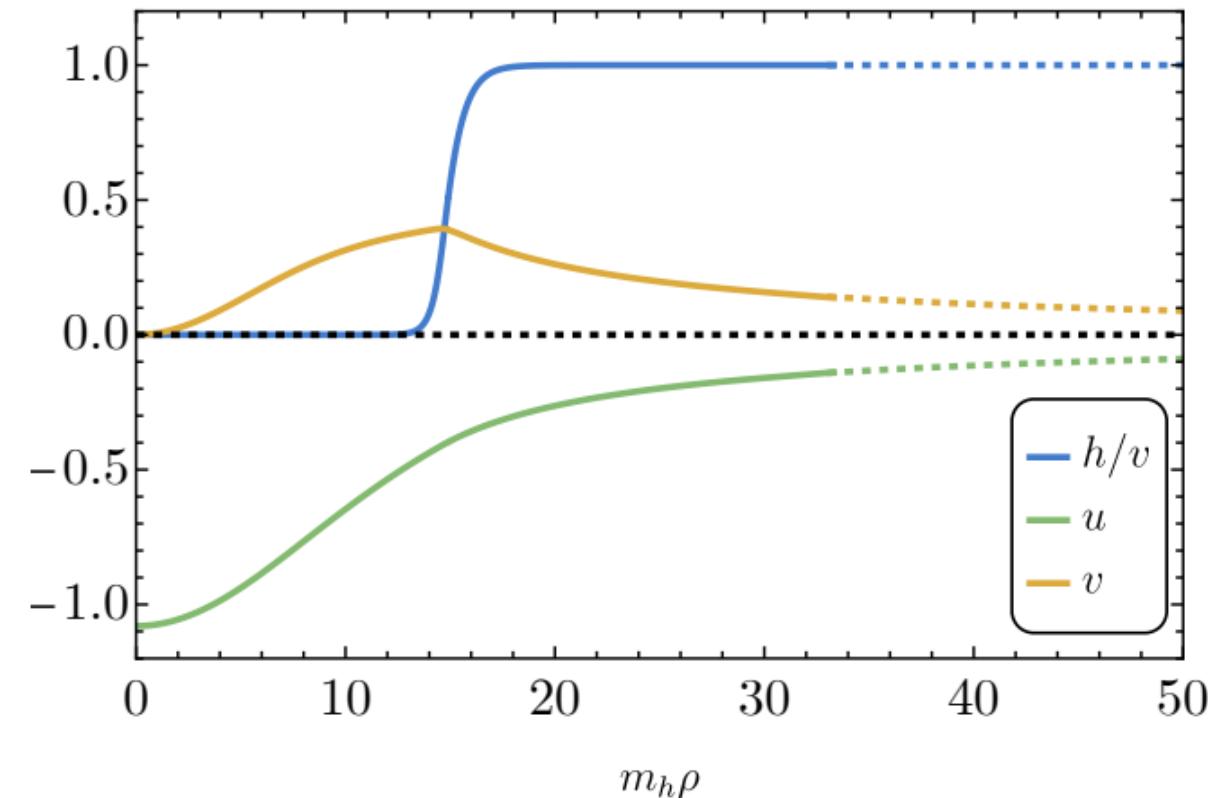
$$= T_{\mu\nu}^{[f]}$$

$$2. \quad \square h - \frac{\partial U}{\partial h} - \frac{f}{\sqrt{2}} S = 0$$

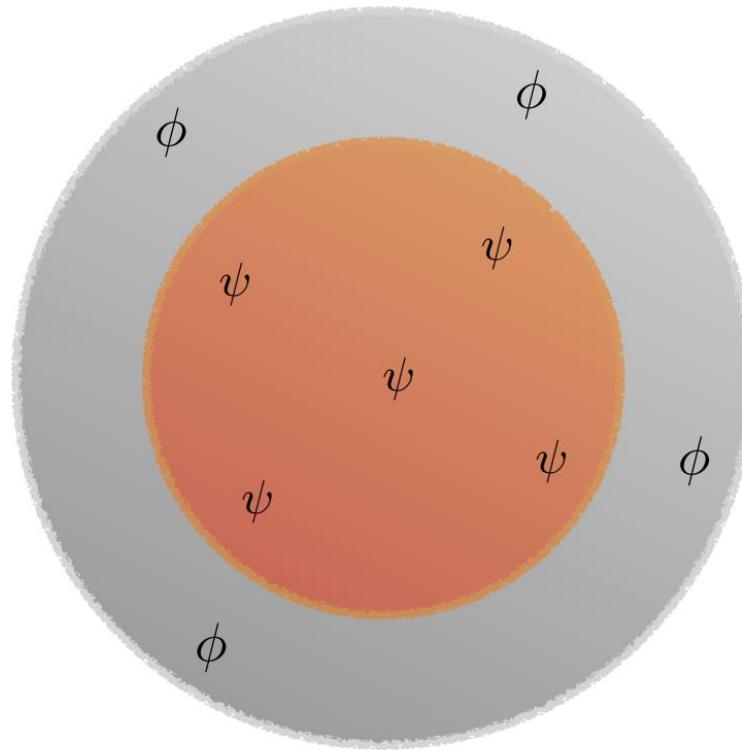
$$\propto (T^{[f]})_\mu^\mu = -W + 3P$$

$$3. \quad F(\underline{k}_F, g_{\mu\nu}, h) = 0$$

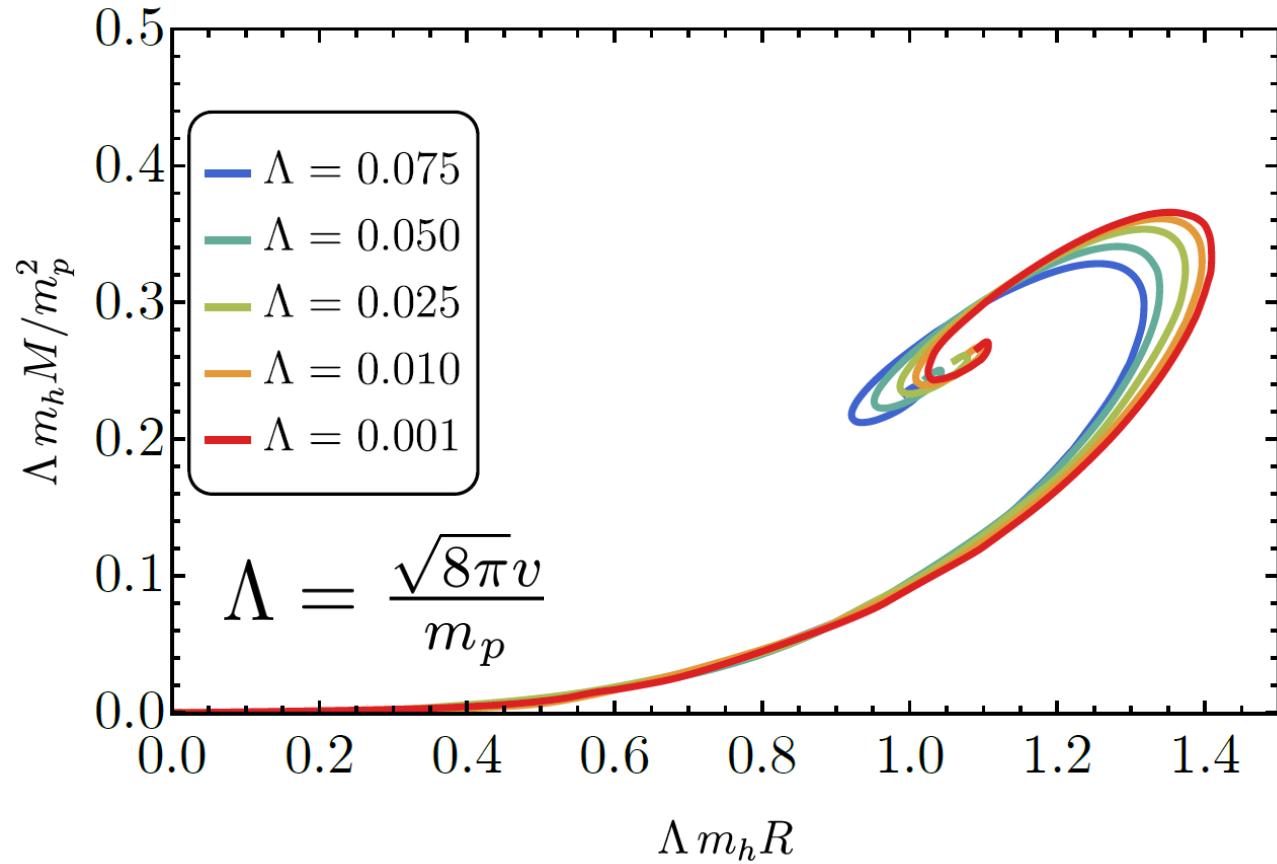
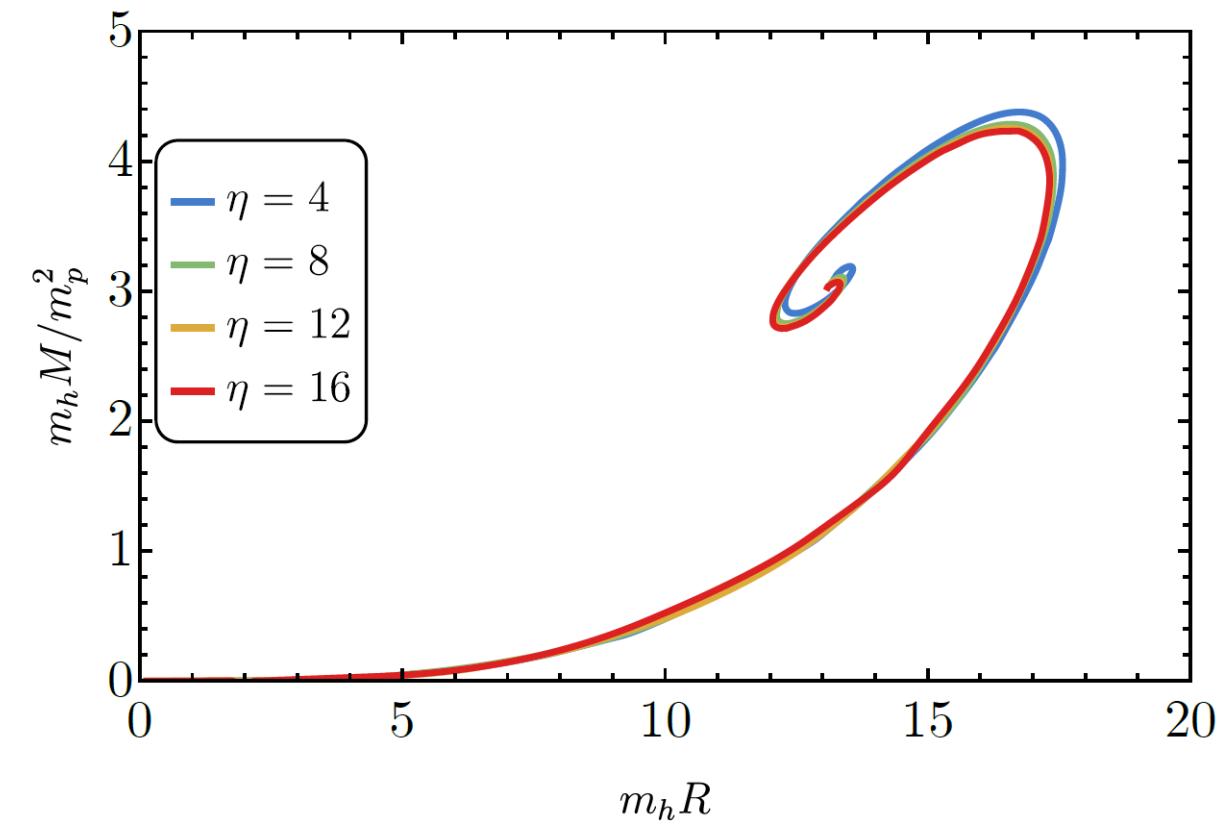
Example of solution



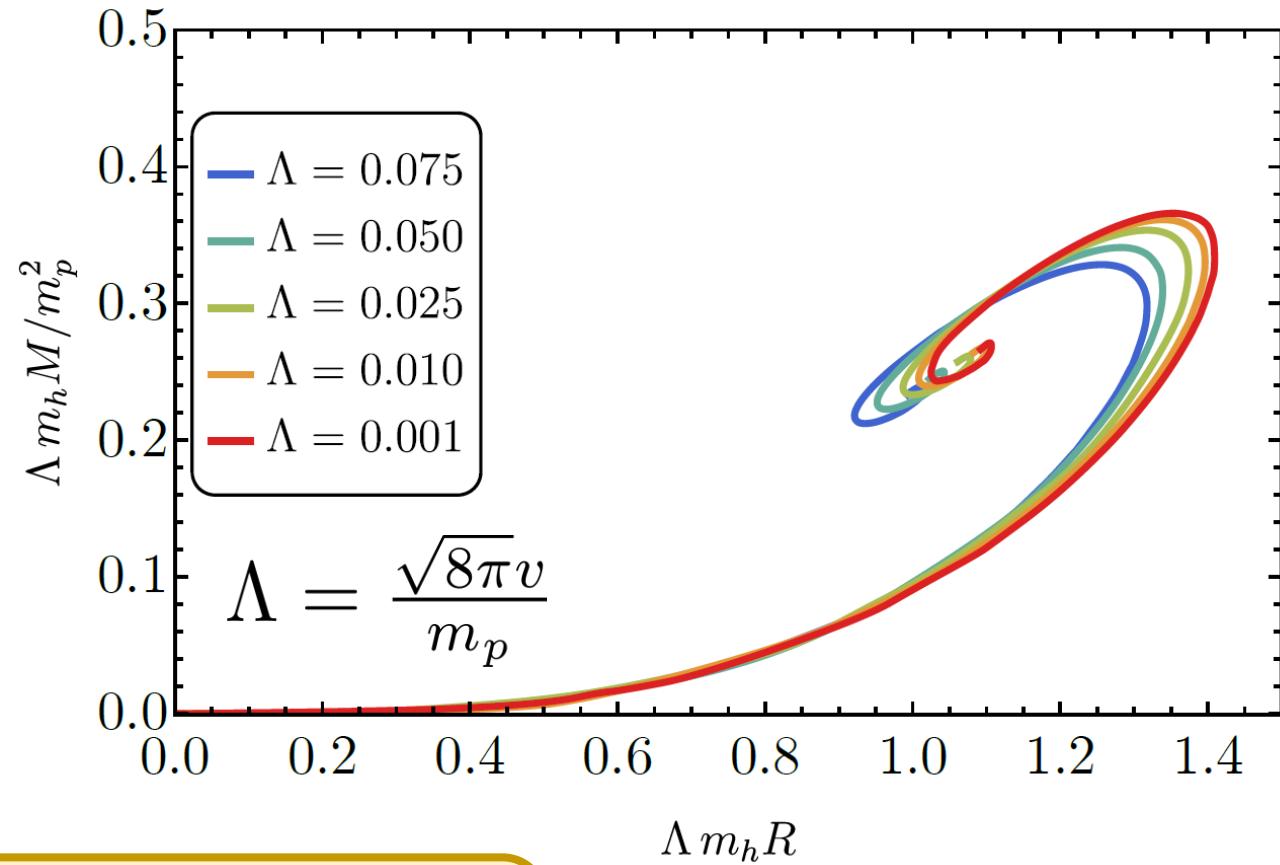
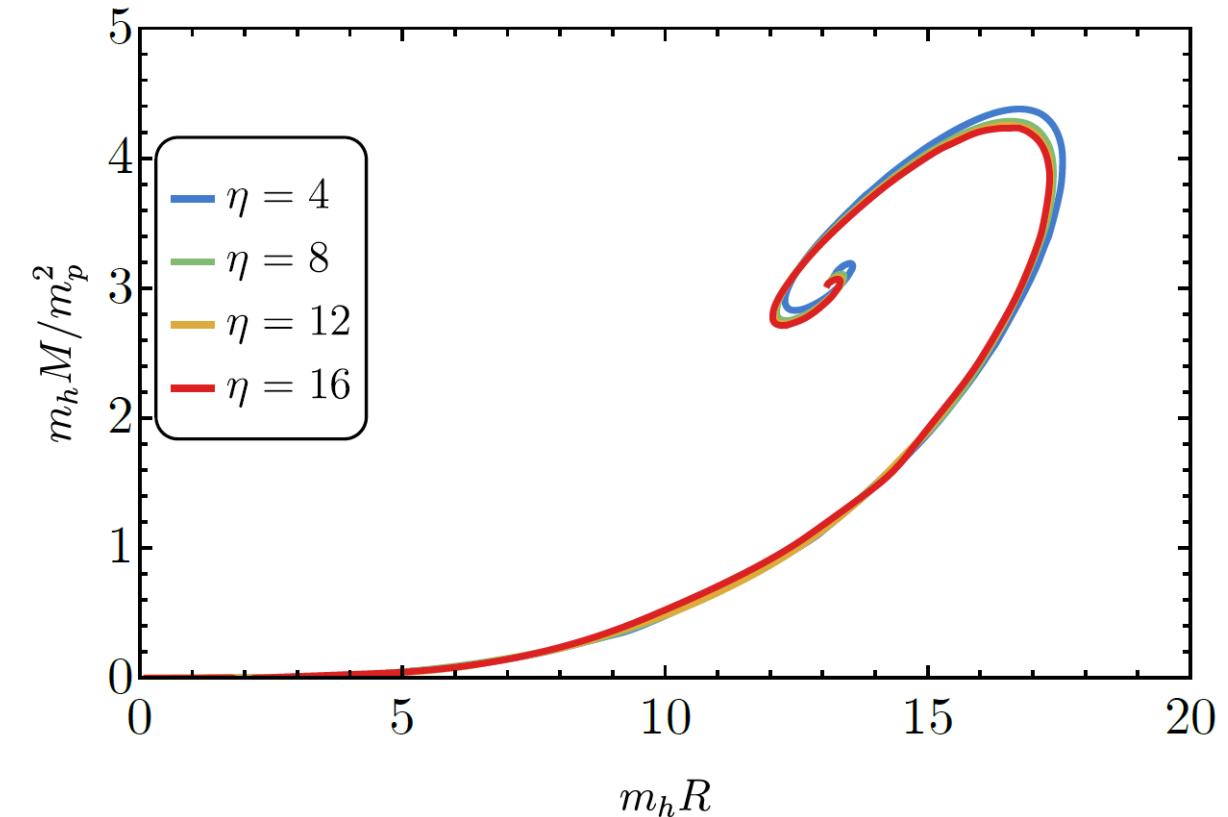
Pictorial illustration of a fermion soliton star



Mass-radius diagrams



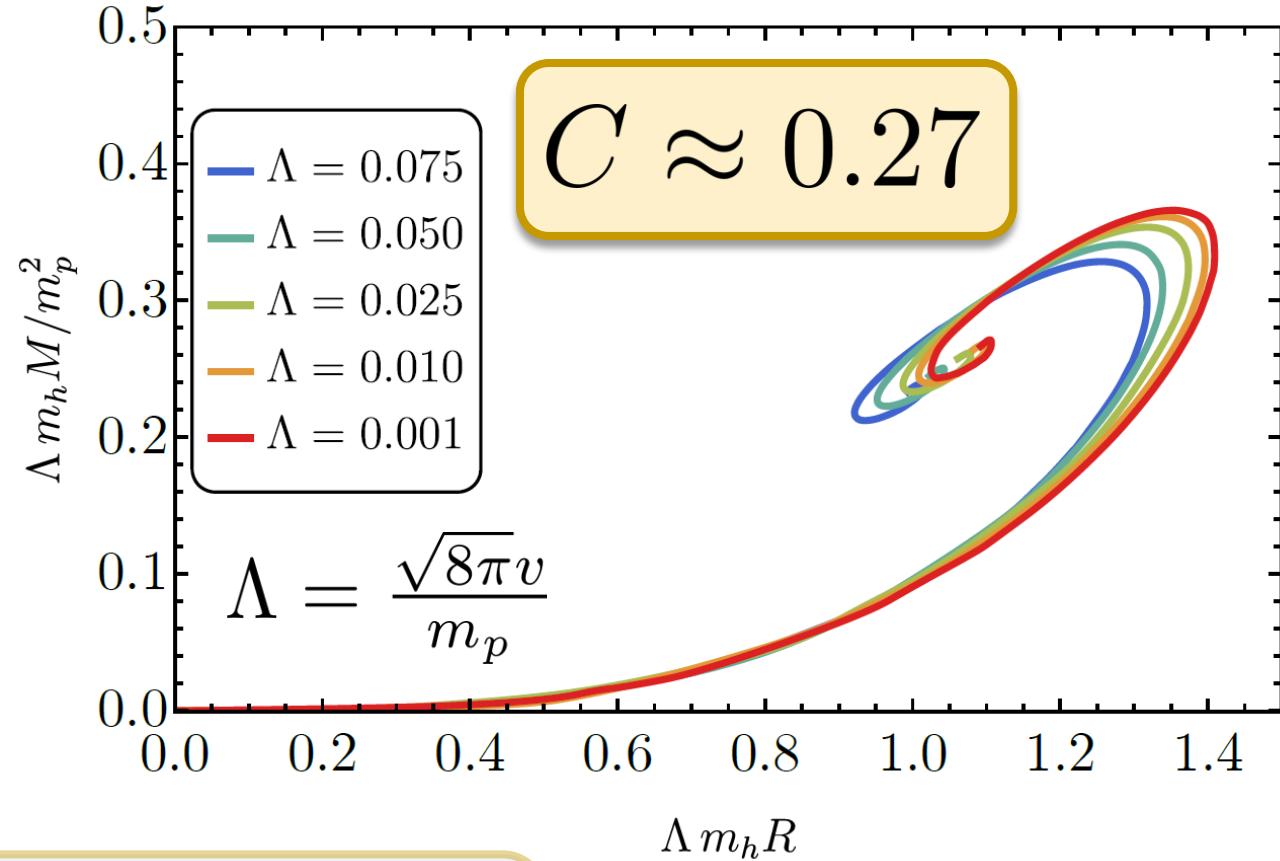
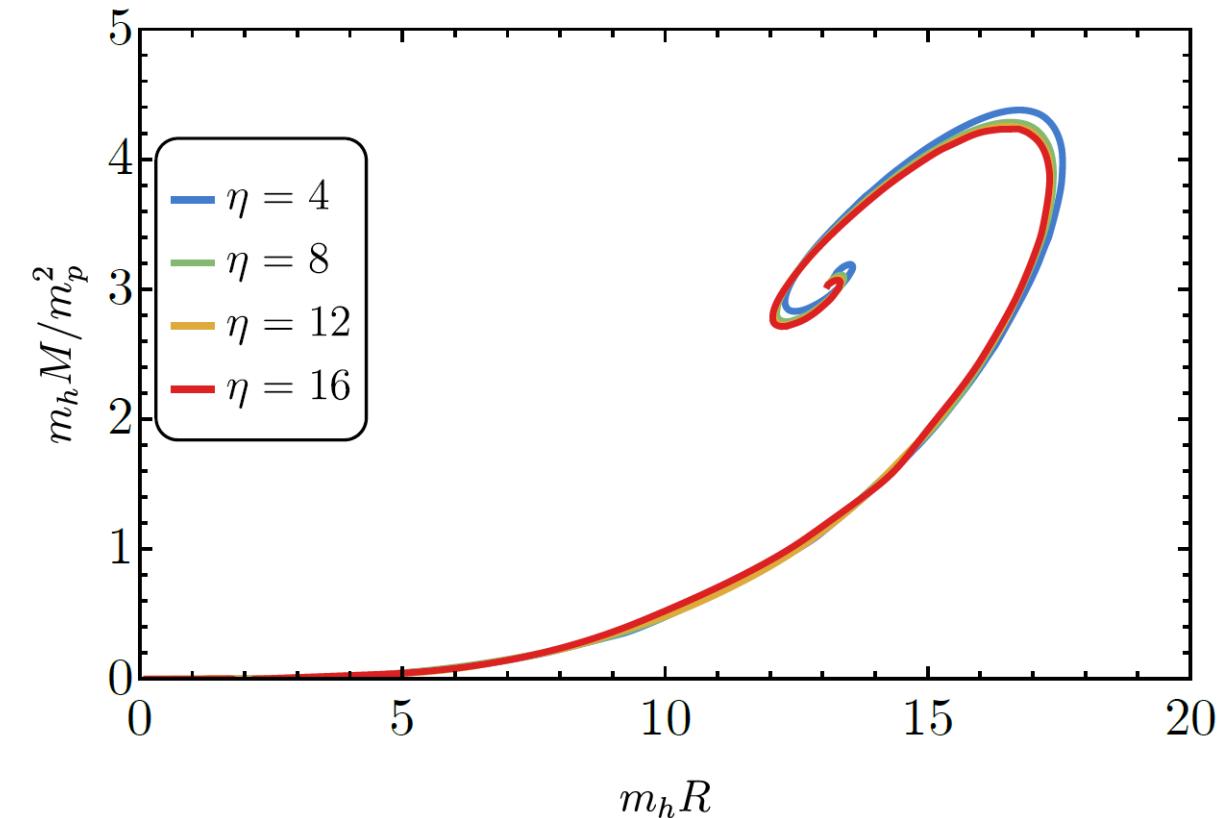
Mass-radius diagrams



$$M_c \sim M_\odot \left(\frac{0.34 \text{ GeV}}{q} \right)^2 \quad R_c \sim 5.5 \text{ km} \left(\frac{0.34 \text{ GeV}}{q} \right)^2$$

$$q \equiv (m_h v)^{1/2}$$

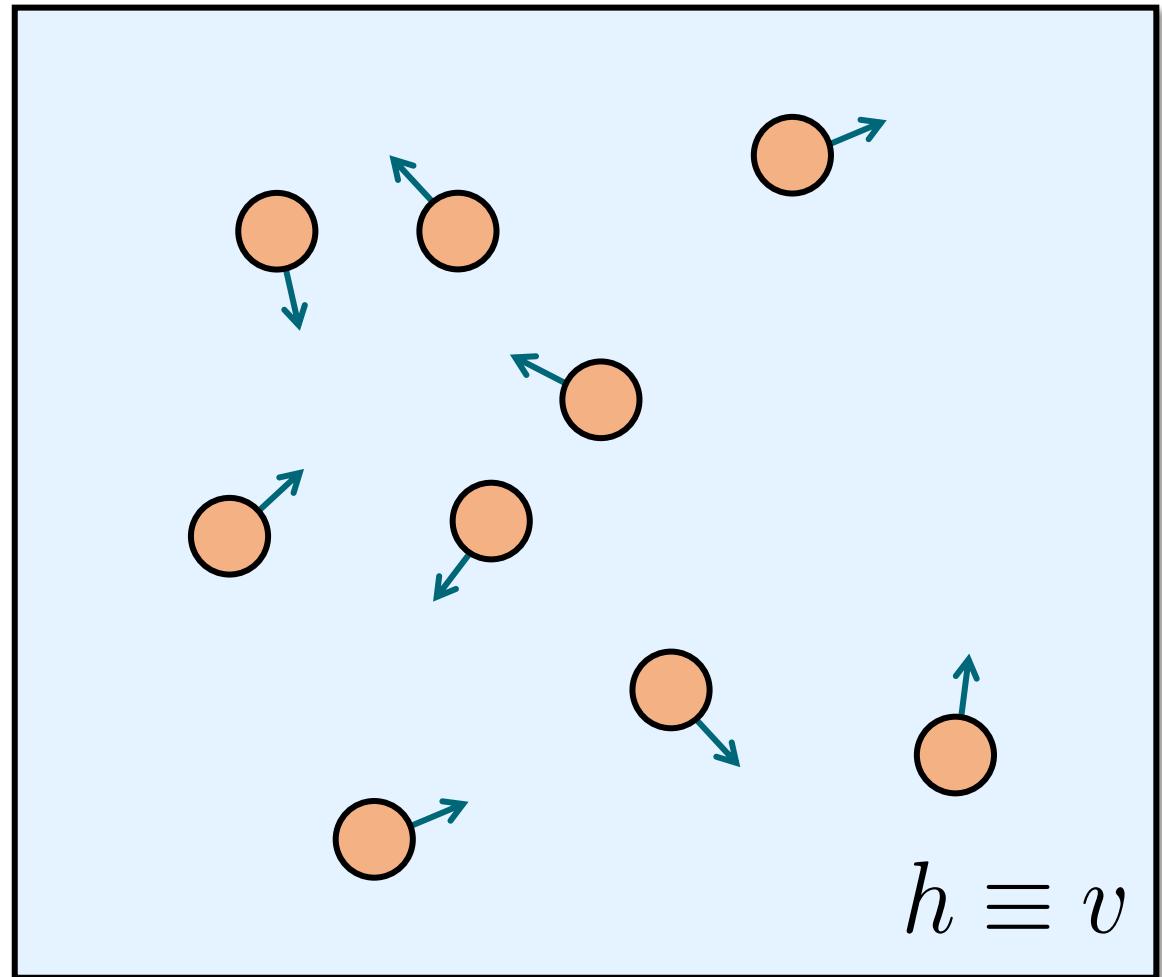
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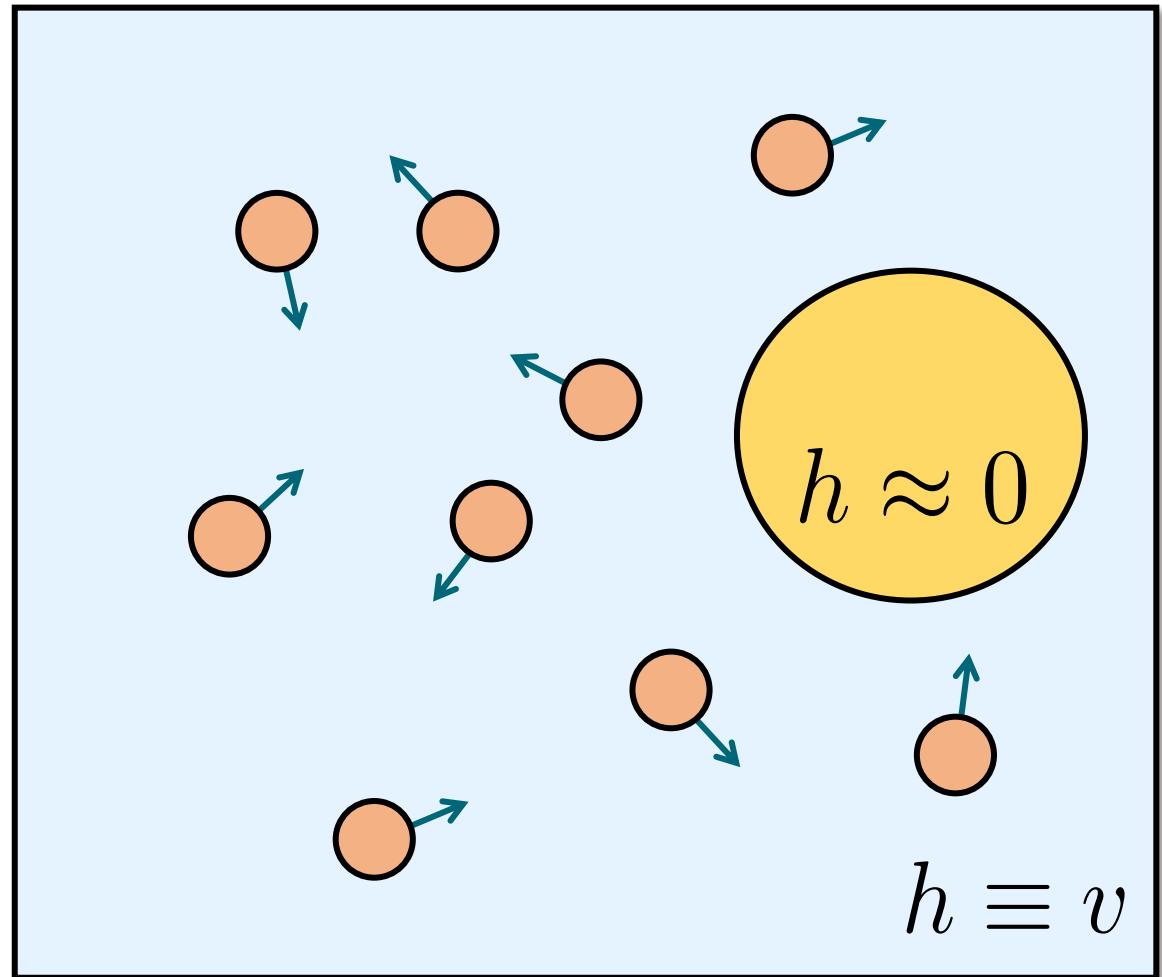
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What is the actual ground state?

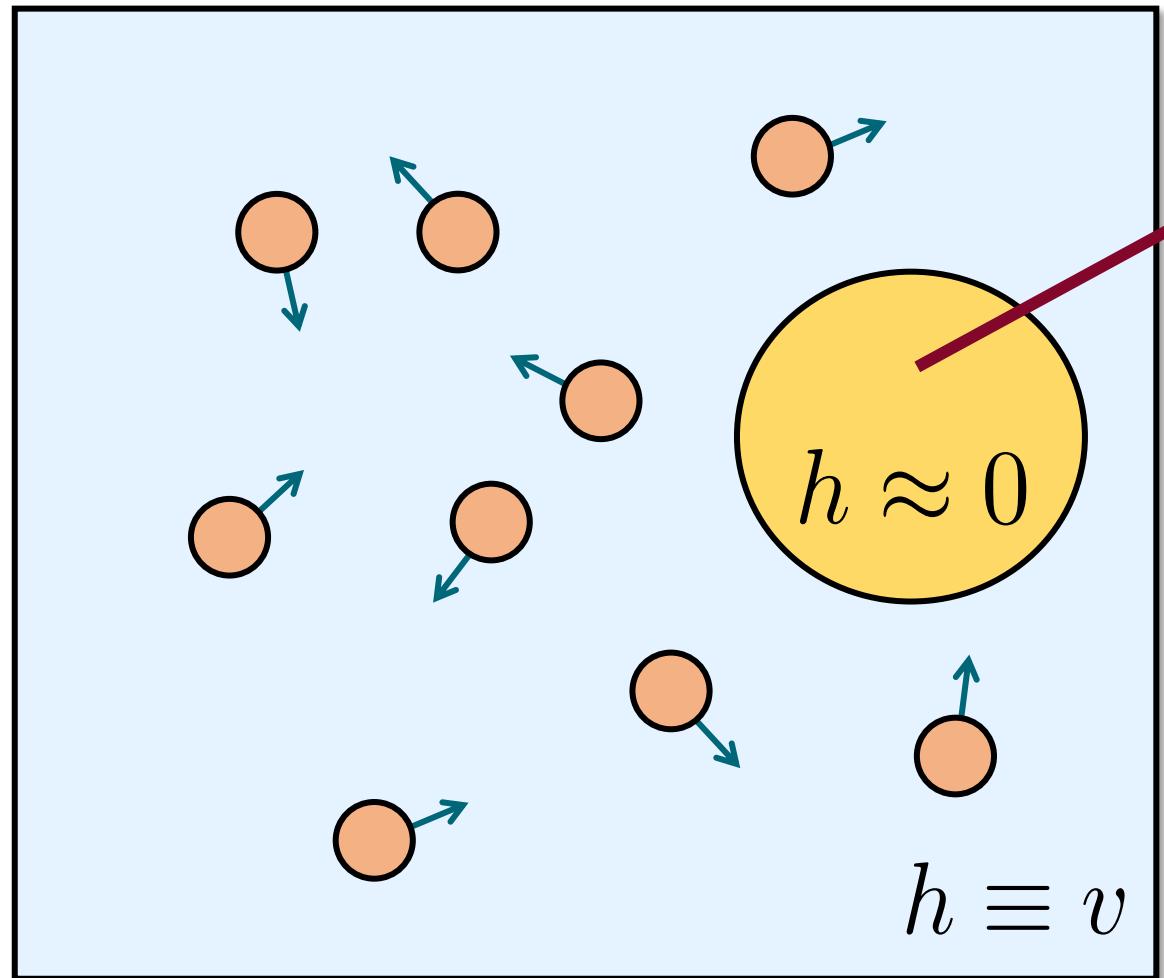


$$h \equiv v$$

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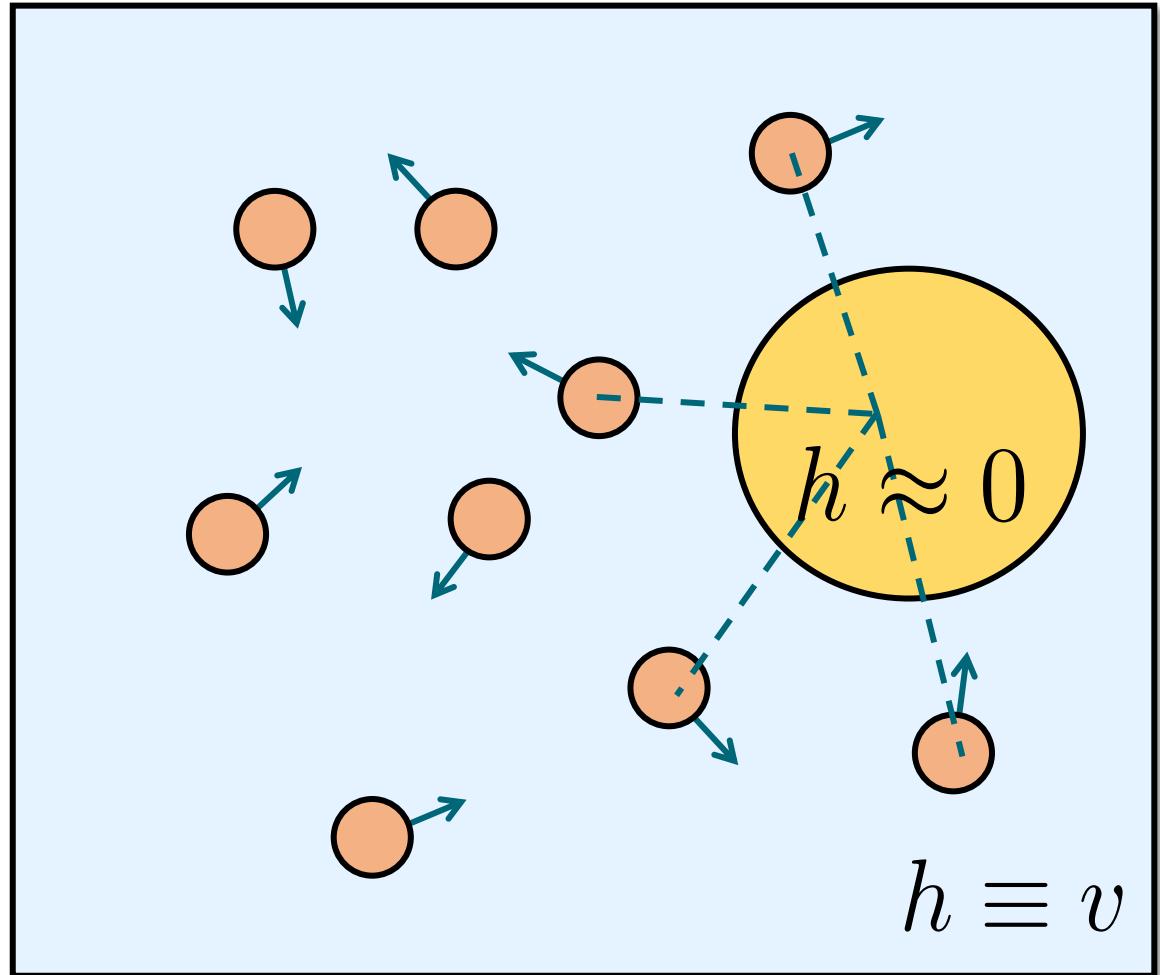


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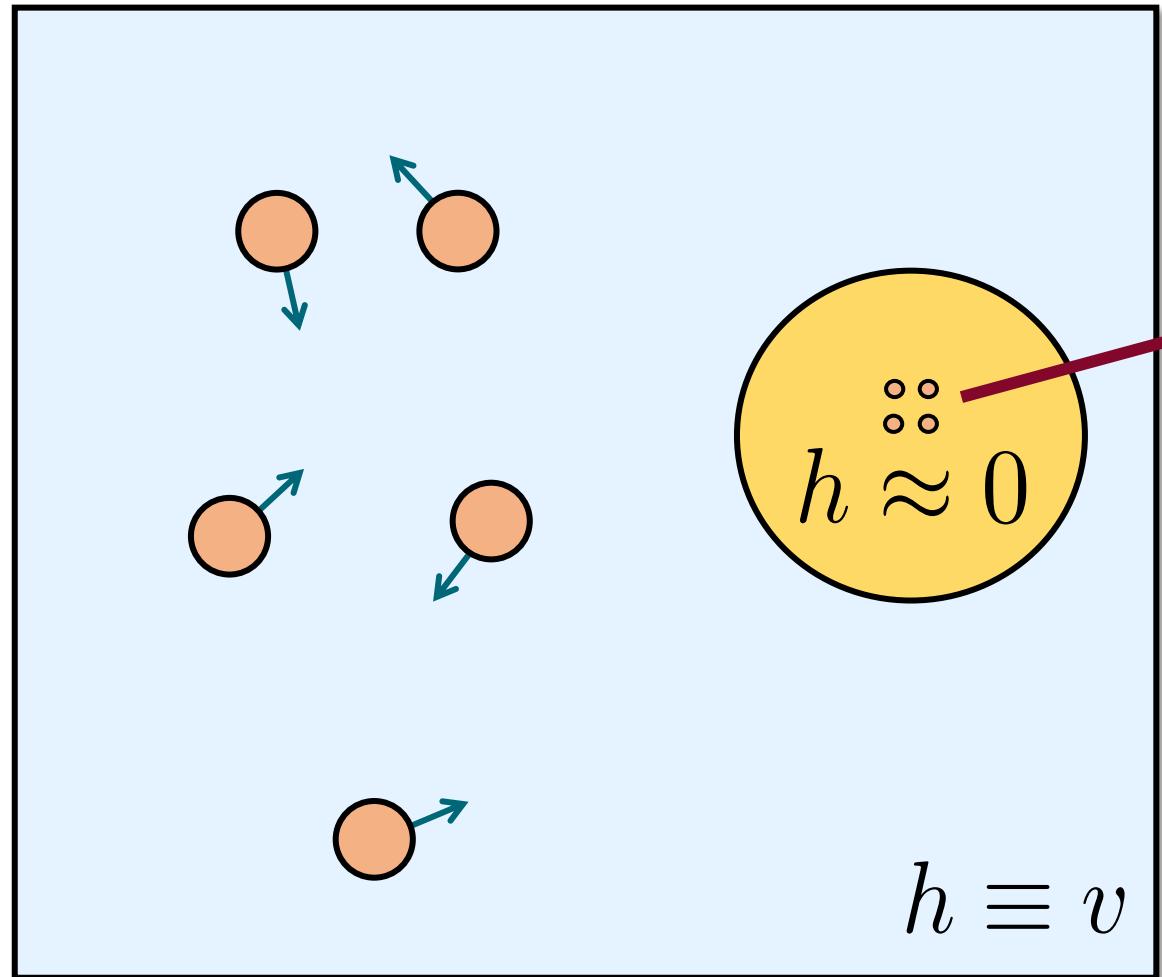
$$E_{\text{cost}} \sim m_h^2 v^2 R^3$$

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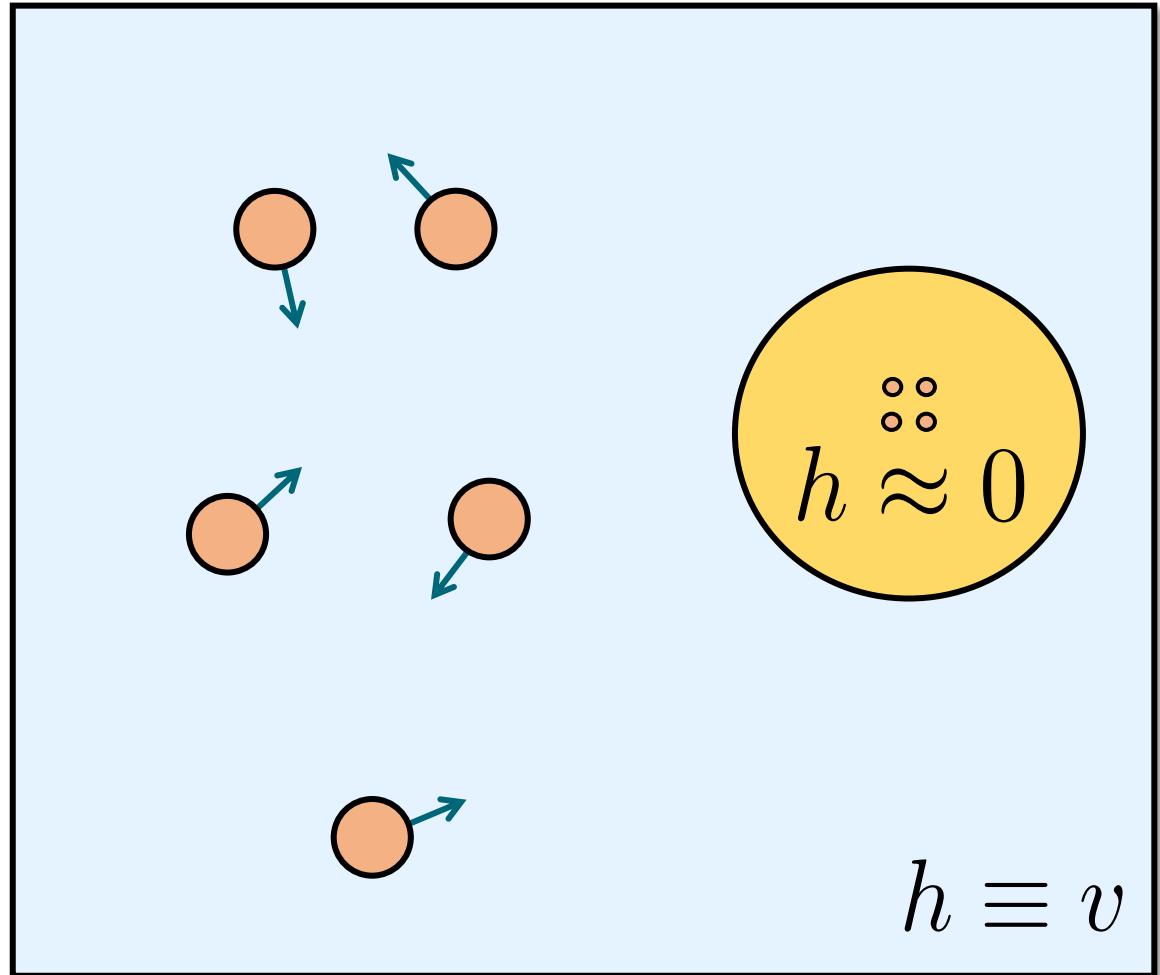
What is the actual ground state?



$$E_{\text{cost}} \sim m_h^2 v^2 R^3$$

$$E_{\text{gain}} \sim N m_f$$

What is the actual ground state?



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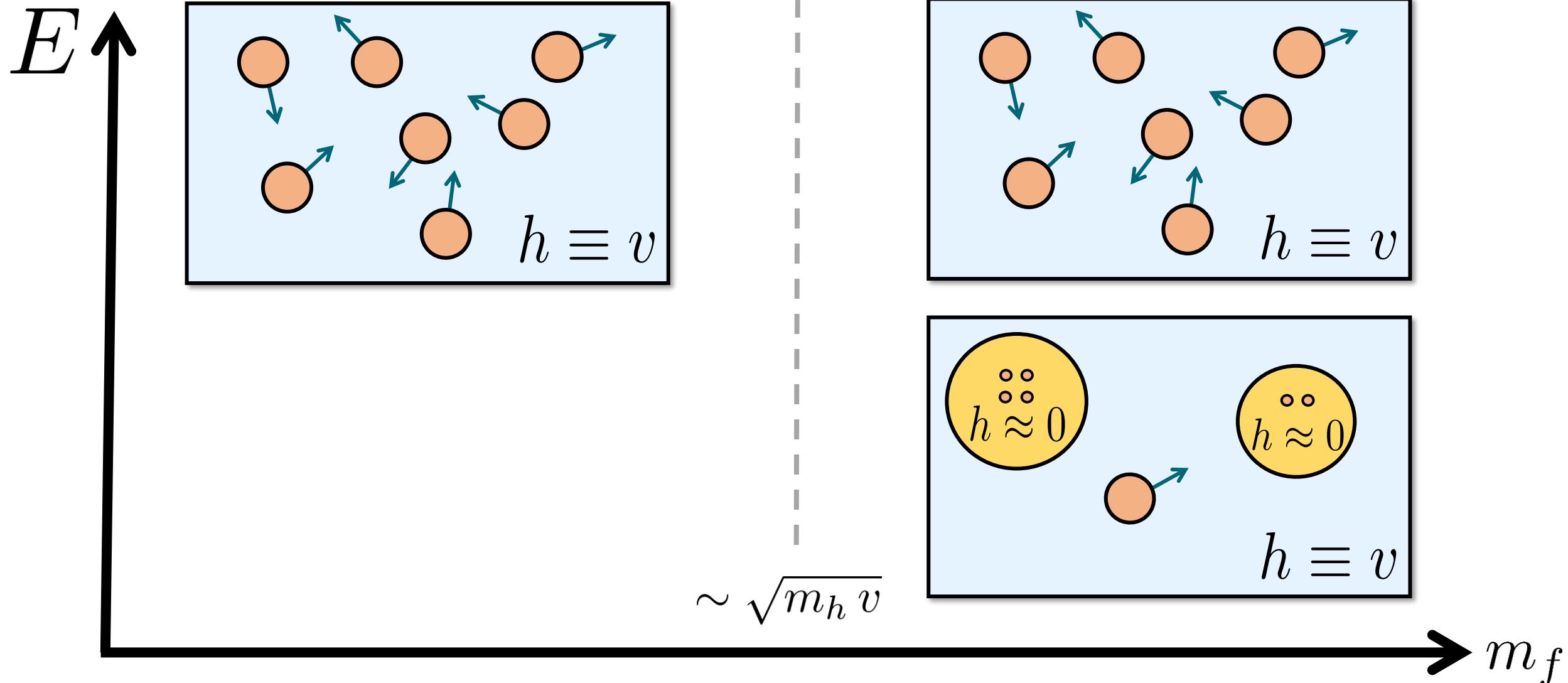
$$E_{\text{gain}} \sim N m_f$$

$$E_{\text{gain}} > E_{\text{cost}}$$



$$m_f \gtrsim \sqrt{m_h v}$$

Non-perturbative vacuum scalarization



Neutron soliton stars

Linear sigma model



Chiral symmetry breaking

Neutron soliton stars

Linear sigma model



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Chiral symmetry breaking

$$m_h \sim 500 \text{ MeV}$$

$$v \sim f_\pi \sim 130 \text{ MeV}$$

$$m_f \sim 1 \text{ GeV}$$

Neutron soliton stars

Linear sigma model



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$$M_c \approx 2 M_\odot$$

$$R_c \approx 10 \text{ km}$$

[R. Balkin, J. Serra, K. Springmann, S. Stelzl, and A. Weiler, 2023]

Other applications

Standard Model Higgs

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Standard Model Higgs



$$M_c \approx (4 \times 10^{-6}) M_{\odot}$$
$$R_c \approx 2 \text{ cm}$$

Other applications

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Dark soliton stars

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$$m_h \gtrsim m_{h_{\text{SM}}} \quad \lambda \sim \mathcal{O}(1)$$

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Dark soliton stars



$$m_h \gtrsim m_{h_{\text{SM}}} \quad \lambda \sim \mathcal{O}(1)$$



sub-solar regime

Thank you

Fermionic quantities

$$W = \frac{2}{(2\pi)^3} \int_0^{k_F(\rho)} d^3 k \sqrt{k^2 + (m_f - f\phi(\rho))^2} \quad (1)$$

$$P = \frac{2}{(2\pi)^3} \int_0^{k_F(\rho)} \frac{d^3 k \, k^2}{3\sqrt{k^2 + (m_f - f\phi(\rho))^2}} \quad (2)$$

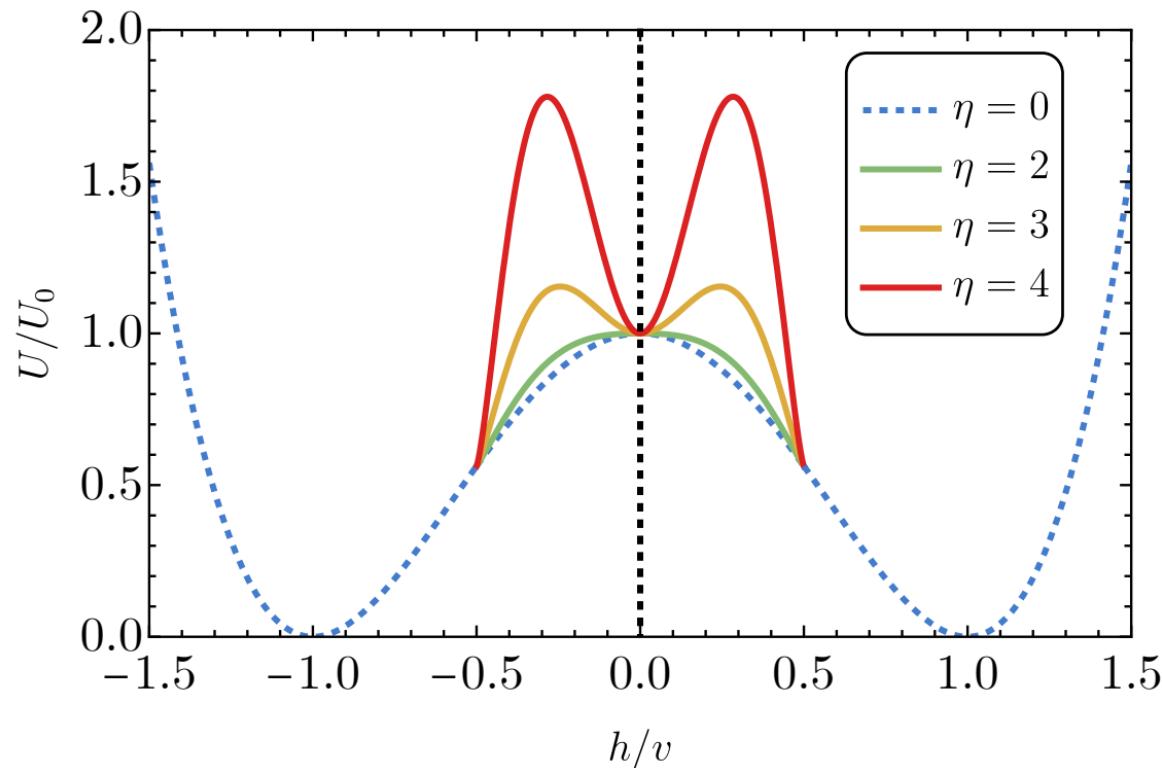
$$S = \frac{2}{(2\pi)^3} \int_0^{k_F(\rho)} d^3 k \frac{m_f - f\phi(\rho)}{\sqrt{k^2 + (m_f - f\phi(\rho))^2}} \quad (3)$$

Effective potential

$$U_{\text{eff}} = \frac{\lambda}{16}(h^2 - v^2)^2 + \frac{f}{\sqrt{2}}h <\bar{\psi}\psi>$$



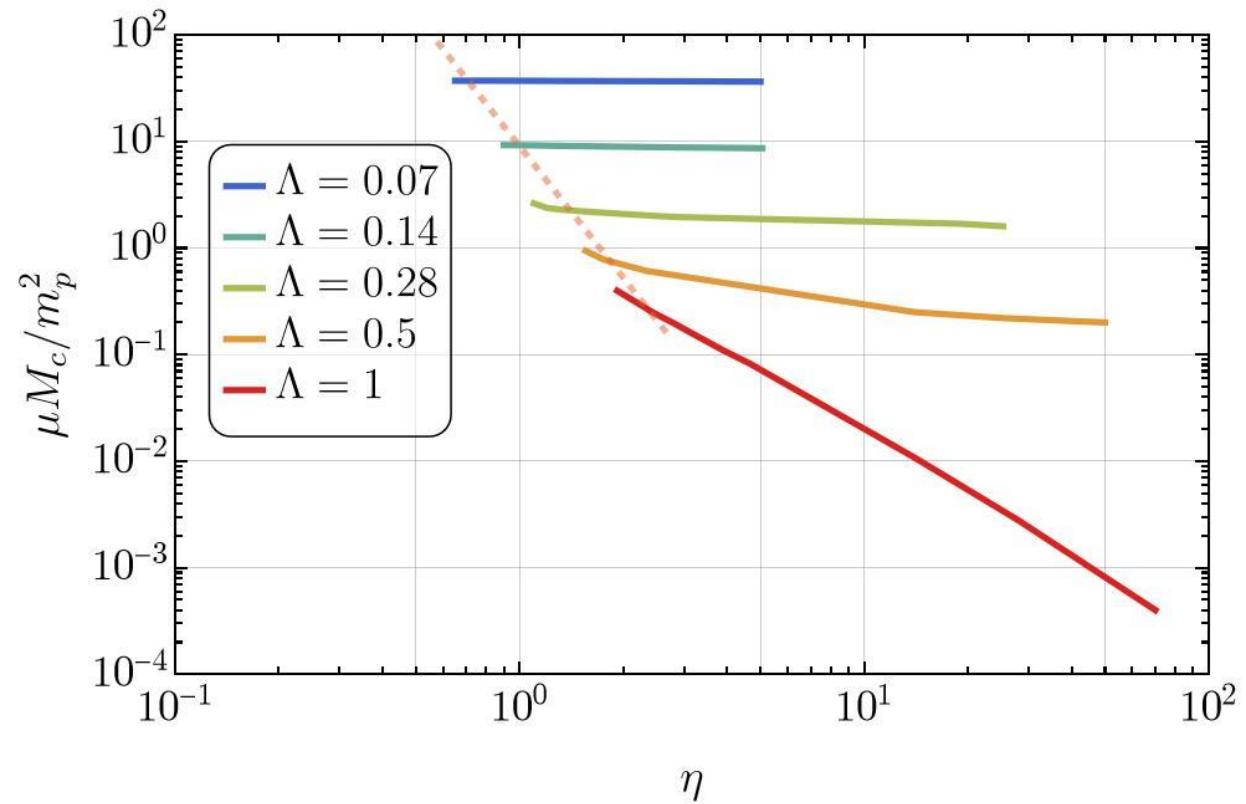
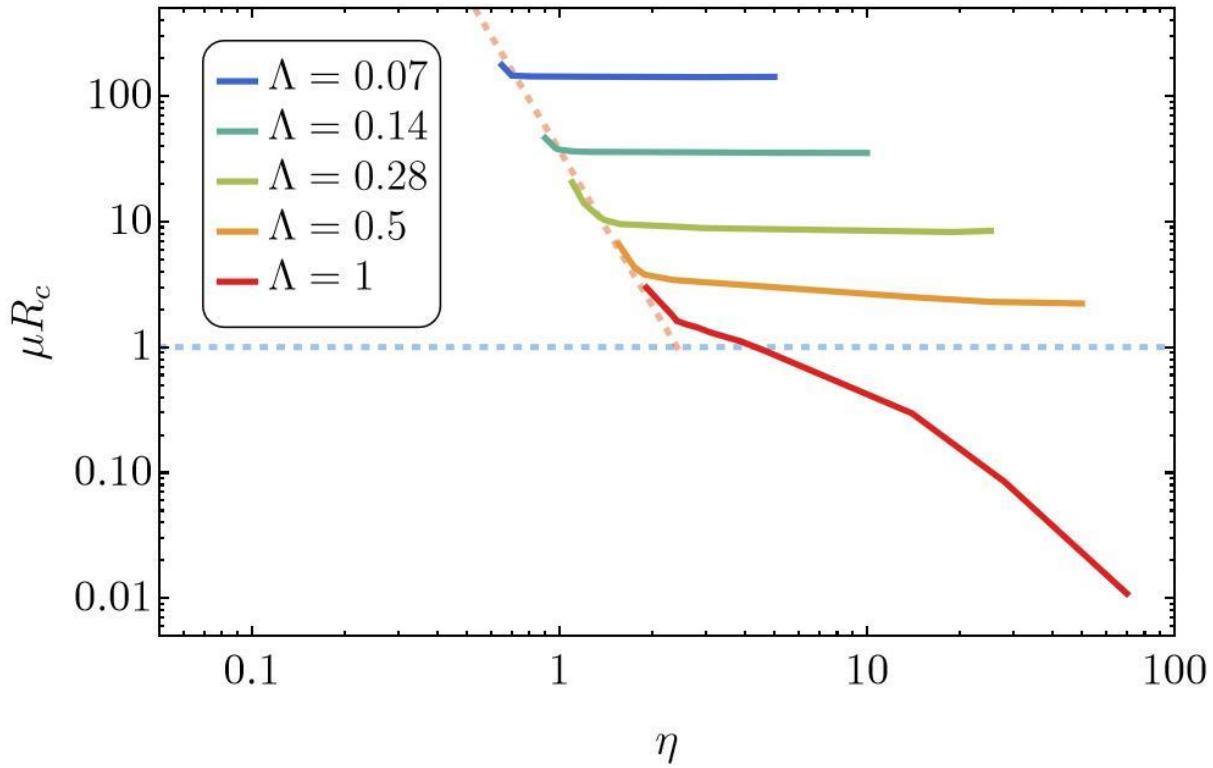
$$\eta = \frac{m_f}{m_h^{1/2}v^{1/2}}$$



False vacuum pockets if

$$m_f \gtrsim \sqrt{m_h v}$$

Confining regime



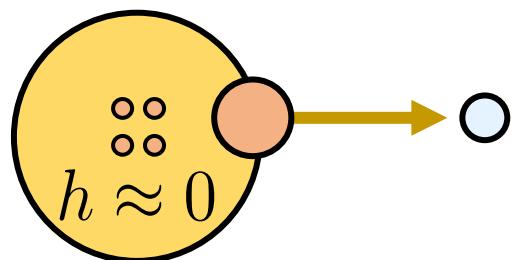
Issues

Numerical computations gives the actual condition

$$m_f > 2\sqrt{m_h v} \sim 350 \text{ GeV}$$

If the fermion is a weakly interacting particle

coupled to the SM $t_{decay} \sim \left(\frac{m_h v}{m_f^2} \right) \times 10^{-11} \text{ sec}$



Exotic phase?

Number density inside the pocket

$$n \sim (m_h v)^{3/2} \sim 7 \times 10^{47} \text{ cm}^{-3}$$

Nuclear matter density

$$n \sim 10^{38} \text{ cm}^{-3}$$

Colored superconductor?

PBH?

Higgs balls are produced in the radiation domination era with an initial mass M_i

$$GM = \frac{t}{1 + \frac{t}{t_i} \left(\frac{t_i}{GM_i} - 1 \right)}$$

$$M_i \gtrsim 3 \times 10^{-6} M_\odot$$

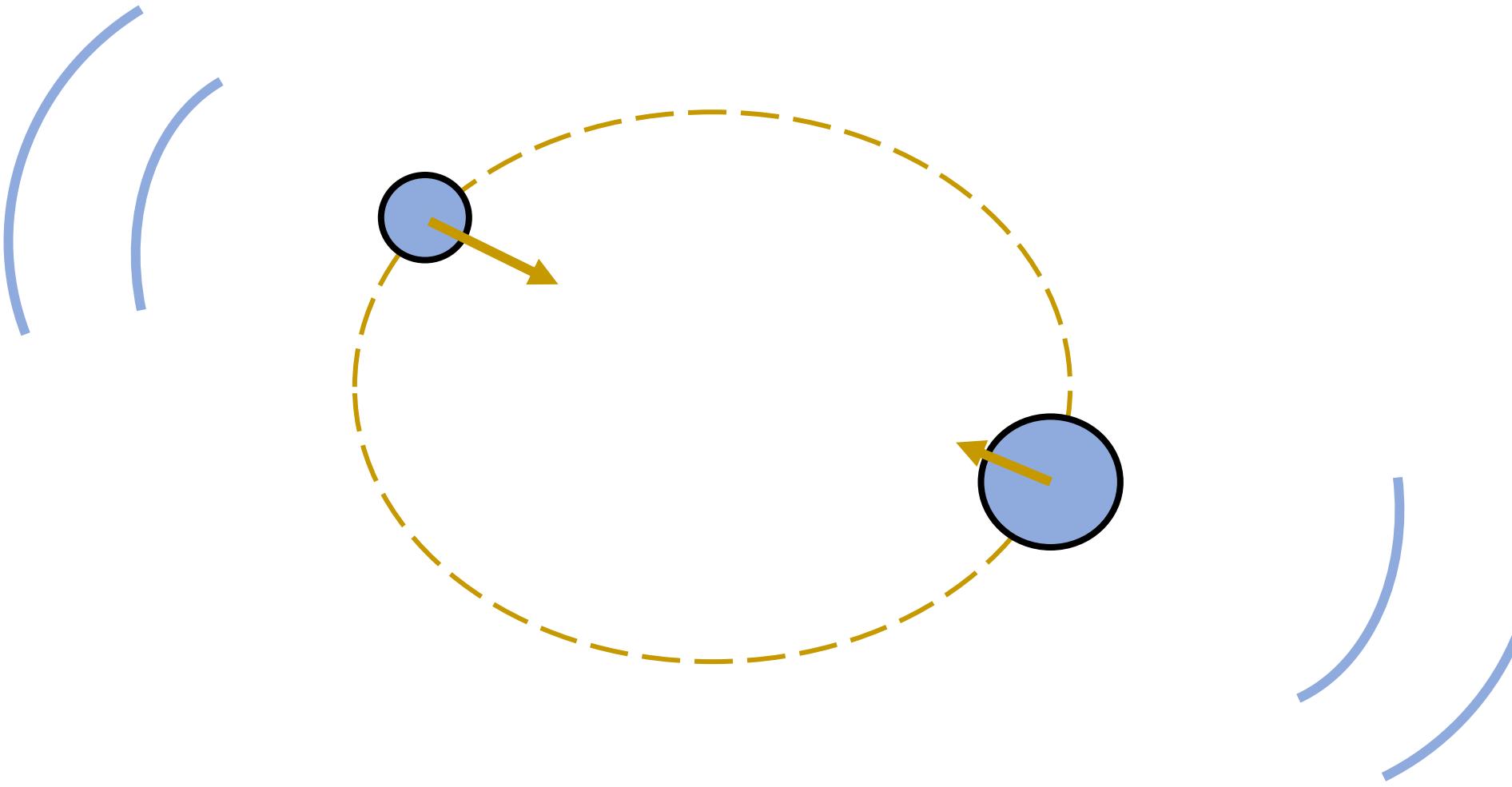
How to discriminate among compact objects?

Tidal deformability is encoded in the Love numbers

$$Q_{ij} = \lambda_2 E_{ij}$$

The spherical symmetry allows us to decompose the perturbation into polar and magnetic sectors

How to discriminate among compact objects?



How to discriminate among compact objects?

