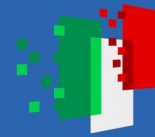




Finanziato
dall'Unione europea
NextGenerationEU



Ministero
dell'Università
e della Ricerca



Italiadomani
PIANO NAZIONALE
DI RIPRESA E RESILIENZA



Missione 4 Istruzione e Ricerca

Vasja Susič

(IPNP, Charles University)

25/01/2024

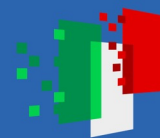
Grand Unified Theories:
A home to the axion?



Finanziato
dall'Unione europea
NextGenerationEU



Ministero
dell'Università
e della Ricerca



Italiadomani
PIANO NAZIONALE
DI RIPRESA E RESILIENZA



Grand Unified Theories: a home to the axion?

Vasja Susič

IPNP, Charles University



AxionOrigins: towards a complete theory for the origin of the axion

INFN, Laboratori Nazionali di Frascati

2024-01-25

- Motivation: what are GUTs and why consider them?
- Model building: possibilities and challenges
 - Choice of gauge group
 - Unification of gauge couplings
 - Symmetry breaking
 - Yukawa sector
- Some phenomenology [neutrinos, proton decay]
- Axion in GUTs

Motivation 1 — Unification of gauge couplings?

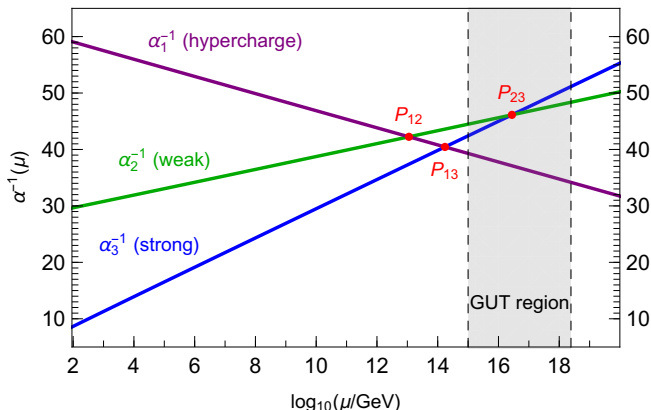
■ RG running of gauge couplings in the Standard Model (SM):

- Proton decay
(\approx exp. limit):

$$\mu_{\text{p-decay}} = 10^{15} \text{ GeV}$$

- Gravity
(reduced Planck scale):

$$\mu_{\text{Pl}} = 2.4 \cdot 10^{18} \text{ GeV}$$



■ Do the SM forces unify in one point at high energy M_{GUT} ?

→ GUT region: $\mu_{\text{p-decay}} < M_{\text{GUT}} < \mu_{\text{Pl}}$

→ Consistent with picture (**non-trivial indication!**)

Motivation 2 — What is a GUT?

- **Grand Unified Theory** (GUT): a Yang-Mills theory with a *simple* gauge factor G , such that

$$SU(3)_C \times SU(2)_L \times U(1)_Y \subset G \quad (1)$$

Motivation 2 — What is a GUT?

- **Grand Unified Theory (GUT)**: a Yang-Mills theory with a *simple* gauge factor G , such that

$$SU(3)_C \times SU(2)_L \times U(1)_Y \subset G \quad (1)$$

- Features of the GUT framework:

- (a) All 3 SM gauge couplings unify at high energy scale in group G
- (b) *Spontaneous symmetry breaking*: in one or multiple stages

$$G \rightarrow \dots G_i \dots \rightarrow G_{SM} \equiv SU(3)_C \times SU(2)_L \times U(1)_Y \quad (2)$$

- (c) SM (in detail) should be recovered at low (EW) scale

Motivation 3 —Why GUT?

- Historical proposal [1974]: Georgi-Glashow SU(5) model

$$\text{fermions : } 3 \times (10 \oplus \bar{5}) \quad (3)$$

$$\text{scalars : } 24 \oplus 5 \quad (4)$$

→ Predictions at M_{GUT} : $\sin^2 \theta_W = 3/8$, $M_D = M_E^T$

→ ruled out long ago (proton decay, Yukawa fit)

→ but other “realistic” GUT models proposed

Motivation 3 —Why GUT?

- Historical proposal [1974]: Georgi-Glashow SU(5) model

$$\text{fermions : } 3 \times (10 \oplus \bar{5}) \quad (3)$$

$$\text{scalars : } 24 \oplus 5 \quad (4)$$

→ Predictions at M_{GUT} : $\sin^2 \theta_W = 3/8$, $M_D = M_E^T$

→ ruled out long ago (proton decay, Yukawa fit)
→ but other “realistic” GUT models proposed

- Generic **benefits** of GUTs:

- 1 Unification of gauge interactions
- 2 Matter unification (at least partial)
- 3 Explain (hyper)charge quantization
- 4 Conceptually nicer/simpler than SM
- 5 Experimental sensitivity to super-high energies (e.g. proton decay)

Model building 1 — possibilities for unified group

- Requirements for a unified gauge group G :
 - (1) G is simple
 - (2) $G \supset G_{\text{SM}}$
 - (3) G has complex representations (since SM is chiral)

Model building 1 — possibilities for unified group

■ Requirements for a unified gauge group G :

- (1) G is simple
- (2) $G \supset G_{\text{SM}}$
- (3) G has complex representations (since SM is chiral)

■ Classification of simple finite-dimensional Lie algebras:

root system	name	comment	\mathbb{C} -irreps?
A_n	$SU(n+1)$	rotations in \mathbb{C}^{n+1}	all n
B_n	$SO(2n+1)$	rotations in \mathbb{R}^{2n+1}	/
C_n	$Sp(2n)$	rotations in \mathbb{H}^n	/
D_n	$SO(2n)$	rotations in \mathbb{R}^{2n}	odd n
E_6, E_7, E_8, F_4, G_2		exceptional	E_6

Model building 1 — possibilities for unified group

- Requirements for a unified gauge group G :

- (1) G is simple
- (2) $G \supset G_{\text{SM}}$
- (3) G has complex representations (since SM is chiral)

- Classification of simple finite-dimensional Lie algebras:

root system	name	comment	\mathbb{C} -irreps?
A_n	$SU(n+1)$	rotations in \mathbb{C}^{n+1}	all n
B_n	$SO(2n+1)$	rotations in \mathbb{R}^{2n+1}	/
C_n	$Sp(2n)$	rotations in \mathbb{H}^n	/
D_n	$SO(2n)$	rotations in \mathbb{R}^{2n}	odd n
E_6, E_7, E_8, F_4, G_2		exceptional	E_6

- Satisfying requirements:
 - (a) $SU(n)$, $n \geq 5$
 - (b) $SO(4n+2)$, $n \geq 2$
 - (c) E_6

- Minimal choices: $G_{\text{SM}} \subset SU(5) \subset SO(10) \subset E_6$

Model building 2 — embedding the SM fermions

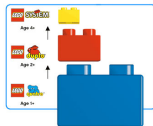
- Each generation in Standard Model $[SU(3)_C \times SU(2)_L \times U(1)_Y]$:

$$\begin{aligned} Q &\sim (3, 2, +\frac{1}{6}) & u^c &\sim (\bar{3}, 1, -\frac{2}{3}) & d^c &\sim (\bar{3}, 1, +\frac{1}{3}) \\ L &\sim (1, 2, -\frac{1}{2}) & e^c &\sim (1, 1, +1) \end{aligned} \quad (5)$$

Model building 2 — embedding the SM fermions

- Each generation in Standard Model $[SU(3)_C \times SU(2)_L \times U(1)_Y]$:

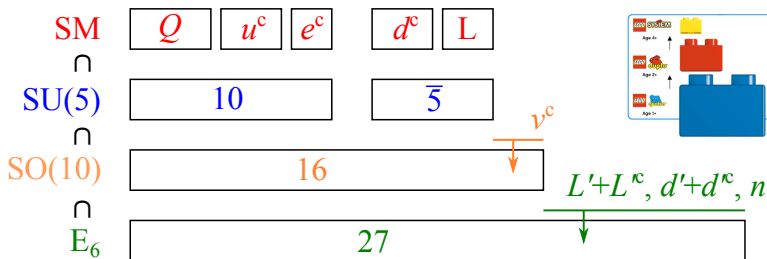
$$\begin{aligned} Q &\sim (3, 2, +\frac{1}{6}) & u^c &\sim (\bar{3}, 1, -\frac{2}{3}) & d^c &\sim (\bar{3}, 1, +\frac{1}{3}) & (5) \\ L &\sim (1, 2, -\frac{1}{2}) & e^c &\sim (1, 1, +1) \end{aligned}$$



Model building 2 — embedding the SM fermions

- Each generation in Standard Model $[SU(3)_C \times SU(2)_L \times U(1)_Y]$:

$$\begin{aligned}
 Q &\sim (3, 2, +\frac{1}{6}) & u^c &\sim (\bar{3}, 1, -\frac{2}{3}) & d^c &\sim (\bar{3}, 1, +\frac{1}{3}) \\
 L &\sim (1, 2, -\frac{1}{2}) & e^c &\sim (1, 1, +1)
 \end{aligned} \quad (5)$$



- $SO(10)$: matter unification, R-neutrinos
- E_6 : vector-like exotic fermions

Model building 3 — details of gauge coupling unification

- Coupling unification requires new states between M_{EW} and M_{GUT}

$$\frac{d}{dt} \alpha_i^{-1} = -\frac{1}{2\pi} \left(a_i + \frac{1}{4\pi} b_{ij} \alpha_j + \dots \right) \quad (6)$$

$$a_i = -\frac{11}{3} D_i(\mathcal{G}) + \frac{2}{3} D_i(\mathcal{F}) + \frac{1}{3} D_i(\mathcal{S}) \quad (7)$$

$\alpha_i := \frac{1}{4\pi} g_i^2$, D_i – Dynkin index;

\mathcal{G} – gauge bosons, \mathcal{F} – Weyl fermions, \mathcal{S} – \mathbb{C} scalars

Model building 3 — details of gauge coupling unification

- Coupling unification requires new states between M_{EW} and M_{GUT}

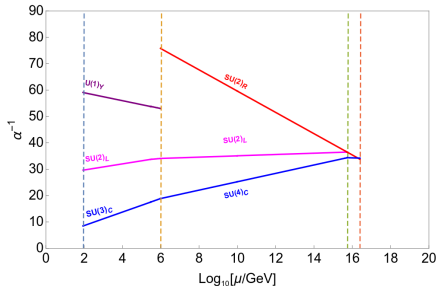
$$\frac{d}{dt} \alpha_i^{-1} = -\frac{1}{2\pi} \left(a_i + \frac{1}{4\pi} b_{ij} \alpha_j + \dots \right) \quad (6)$$

$$a_i = -\frac{11}{3} D_i(\mathcal{G}) + \frac{2}{3} D_i(\mathcal{F}) + \frac{1}{3} D_i(\mathcal{S}) \quad (7)$$

$\alpha_i := \frac{1}{4\pi} g_i^2$, D_i – Dynkin index;

\mathcal{G} – gauge bosons, \mathcal{F} – Weyl fermions, \mathcal{S} – \mathbb{C} scalars

- Example: $SO(10) \rightarrow SU(4)_C \times SU(2)_L \times SU(2)_R \rightarrow G_{SM}$



- Typical addition of fields: multi-stage breaking or \sim TeV SUSY

Model building 3 — details of gauge coupling unification

- Coupling unification requires new states between M_{EW} and M_{GUT}

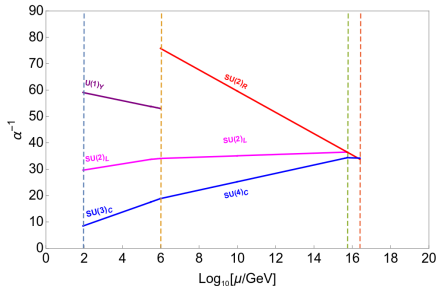
$$\frac{d}{dt} \alpha_i^{-1} = -\frac{1}{2\pi} \left(a_i + \frac{1}{4\pi} b_{ij} \alpha_j + \dots \right) \quad (6)$$

$$a_i = -\frac{11}{3} D_i(\mathcal{G}) + \frac{2}{3} D_i(\mathcal{F}) + \frac{1}{3} D_i(\mathcal{S}) \quad (7)$$

$\alpha_i := \frac{1}{4\pi} g_i^2$, D_i – Dynkin index;

\mathcal{G} – gauge bosons, \mathcal{F} – Weyl fermions, \mathcal{S} – \mathbb{C} scalars

- Example: $SO(10) \rightarrow SU(4)_C \times SU(2)_L \times SU(2)_R \rightarrow G_{SM}$



Properties:

- \mathcal{G} push upward \uparrow
 - \mathcal{F}, \mathcal{S} push downward \downarrow
 - $\alpha^{-1} = 0$: Landau pole
- limits matter content**
- 1-loop: straight lines
 - 2-loop: top-down viewpoint

- Typical addition of fields: multi-stage breaking or $\sim \text{TeV}$ SUSY

Model building 4 — GUT symmetry breaking

- Usual mechanism for SSB $G \rightarrow \dots \rightarrow G_{\text{SM}}$: analogous to EW-symmetry breaking in SM (**Higgs mechanism**)
 - Symmetry broken by a **VEV** of the scalar field
 - **Massive gauge bosons** from the coset G/G_{SM}

Model building 4 — GUT symmetry breaking

- Usual mechanism for SSB $G \rightarrow \dots \rightarrow G_{\text{SM}}$: analogous to EW-symmetry breaking in SM (**Higgs mechanism**)
 - Symmetry broken by a **VEV** of the scalar field
 - **Massive gauge bosons** from the coset G/G_{SM}
- Additional considerations:
 - Multi-stage breaking:
possible in **SO(10)** and **E₆**, no intermediate groups in **SU(5)**
example : $SO(10) \rightarrow 4_C 2_L 2_R \rightarrow 3_C 2_L 2_R 1_{B-L} \rightarrow G_{\text{SM}}$ (8)
 - Independent of SUSY breaking (if SUSY is present)

Model building 4 — GUT symmetry breaking

- Usual mechanism for SSB $G \rightarrow \dots \rightarrow G_{\text{SM}}$: analogous to EW-symmetry breaking in SM (**Higgs mechanism**)
 - Symmetry broken by a **VEV** of the scalar field
 - **Massive gauge bosons** from the coset G/G_{SM}
- Additional considerations:
 - Multi-stage breaking:
possible in **SO(10)** and **E₆**, no intermediate groups in **SU(5)**
example : $SO(10) \rightarrow 4_C 2_L 2_R \rightarrow 3_C 2_L 2_R 1_{B-L} \rightarrow G_{\text{SM}}$ (8)
 - Independent of SUSY breaking (if SUSY is present)
- Which scalar irreps can acquire a SM-preserving VEV?
→ They must contain **SM-singlets**:

$$SU(5) : 1, 5, 10, 15, 24, 35, 40, 45, 50, 70, 70', 75, \dots \quad (9)$$

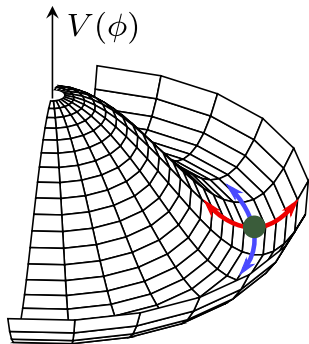
$$SO(10) : 1, 10, 16, 45, 54, 120, 126, 144, 210, 210', \dots \quad (10)$$

$$E_6 : 1, 27, 78, 351, 351', 650, \dots \quad (11)$$

Model building 5 — GUT symmetry breaking continued

■ Determining SSB for $G \rightarrow H$:

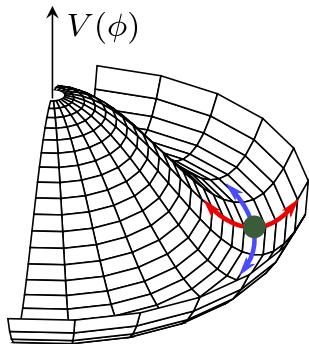
- Choose scalar representations ϕ of G
- Construct the most general scalar potential $V(\phi)$ using invariants of G
- Stationarity conditions (*singlet ansatz*):
solve $\partial_\phi V = 0$ to get $\phi_0 = \langle \phi \rangle$
- Mass matrix $M = \partial_\phi^2 V|_{\phi_0}$
 - **positive** mass: $\dim(\phi) - \dim(G/H)$
 - **zero** mass [WGBs]: $\dim(G/H)$



Model building 5 — GUT symmetry breaking continued

■ Determining SSB for $G \rightarrow H$:

- Choose scalar representations ϕ of G
- Construct the most general scalar potential $V(\phi)$ using invariants of G
- Stationarity conditions (*singlet ansatz*):
solve $\partial_\phi V = 0$ to get $\phi_0 = \langle \phi \rangle$
- Mass matrix $M = \partial_\phi^2 V|_{\phi_0}$
→ **positive** mass: $\dim(\phi) - \dim(G/H)$
→ **zero** mass [WGBs]: $\dim(G/H)$



■ Example: $SU(5) \rightarrow G_{SM}$ via $24 \sim \Sigma$

$$V(\Sigma) = -m^2 \text{Tr} \Sigma^2 + m' \text{Tr} \Sigma^3 + \lambda (\text{Tr} \Sigma^2)^2 + \lambda' \text{Tr} \Sigma^4 \quad (12)$$

$$\langle \Sigma \rangle = v \text{diag}(2, 2, 2, -3, -3) \quad (13)$$

$$\text{WGB} \sim (3, 2, -\frac{5}{6}) \quad (14)$$

■ Global minimum of V ? In general a hard problem!

Model building 6 — Yukawa sector

- Fermions — 3 generations of each (usually flavor \perp GUT):

$$\text{SU}(5) : 10, \bar{5} \quad \text{SO}(10) : 16 \quad \text{E}_6 : 27 \quad (15)$$

- Renormalizable Yukawa terms: FFS

Model building 6 — Yukawa sector

- Fermions — 3 generations of each (usually flavor \perp GUT):

$$\text{SU}(5) : 10, \bar{5} \quad \text{SO}(10) : 16 \quad \text{E}_6 : 27 \quad (15)$$

- Renormalizable Yukawa terms: FFS
- **Suitable** representations S of scalars to accommodate the Higgs:
 - (a) Forms G -invariant with fermions
 - (b) $S \supset (1, 2, \pm\frac{1}{2})$

Model building 6 — Yukawa sector

- Fermions — 3 generations of each (usually flavor \perp GUT):

$$\text{SU}(5) : 10, \bar{5} \quad \text{SO}(10) : 16 \quad \text{E}_6 : 27 \quad (15)$$

- Renormalizable Yukawa terms: FFS

- **Suitable** representations S of scalars to accommodate the Higgs:

(a) Forms G -invariant with fermions

(b) $S \supset (1, 2, \pm \frac{1}{2})$

$$\text{SU}(5) : \quad \bar{5} \otimes \bar{5} = \overline{15}_s \oplus \overline{10}_a \quad (16)$$

$$10 \otimes \bar{5} = 5 \oplus 45 \quad (17)$$

$$10 \otimes 10 = \bar{5}_s \oplus \overline{45}_a \oplus \overline{50}_s \quad (18)$$

$$\text{SO}(10) : \quad 16 \otimes 16 = 10_s \oplus 126_s \oplus 120_a \quad (19)$$

$$\text{E}_6 : \quad 27 \otimes 27 = \overline{27}_s \oplus 351'_s \oplus 351_a \quad (20)$$

Model building 7 — Yukawa sector continued

- Example 1: SU(5) with 5

$$\mathcal{L}_{Yuk} = Y_{10}^{ij} \underbrace{10_i 10_j 5}_{Qu^c H + \dots} + Y_5^{ij} \underbrace{10_i \bar{5}_j 5^*}_{Qd^c H^* + e^c LH^*} \quad (21)$$

→ Yukawa relation $Y_D = Y_E^T$ (at M_{GUT})

NOT GOOD (ENOUGH)!

→ Fix: new operators (45 or non-renormalizable terms)

Model building 7 — Yukawa sector continued

- Example 1: SU(5) with 5

$$\mathcal{L}_{Yuk} = Y_{10}^{ij} \underbrace{10_i 10_j 5}_{Qu^c H + \dots} + Y_5^{ij} \underbrace{10_i \bar{5}_j 5^*}_{Qd^c H^* + e^c L H^*} \quad (21)$$

- Example 2: SO(10) with $10_C \oplus 126$

$$\mathcal{L}_{Yuk} = Y_{10}^{ij} 16_i 16_j 10 + \tilde{Y}_{10}^{ij} 16_i 16_j 10^* + Y_{126}^{ij} 16_i 16_j 126^* \quad (22)$$

Yukawa relations at M_{GUT} :

$$M_U = Y_{10} v_{10}^u + \tilde{Y}_{10} v_{10}^{d*} + Y_{126} v_{126}^u$$

$$M_D = Y_{10} v_{10}^d + \tilde{Y}_{10} v_{10}^{u*} + Y_{126} v_{126}^d$$

$$M_E = Y_{10} v_{10}^d + \tilde{Y}_{10} v_{10}^{u*} - 3 Y_{126} v_{126}^d$$

Fit: works already with Y_{10}, Y_{126}

Model building 7 — Yukawa sector continued

- Example 1: SU(5) with 5

$$\mathcal{L}_{Yuk} = Y_{10}^{ij} \underbrace{10_i 10_j 5}_{Qu^c H + \dots} + Y_5^{ij} \underbrace{10_i \bar{5}_j 5^*}_{Qd^c H^* + e^c L H^*} \quad (21)$$

- Example 2: SO(10) with $10_C \oplus 126$

$$\mathcal{L}_{Yuk} = Y_{10}^{ij} 16_i 16_j 10 + \tilde{Y}_{10}^{ij} 16_i 16_j 10^* + Y_{126}^{ij} 16_i 16_j 126^* \quad (22)$$

- SM Higgs: in $M^2(1, 2, \pm\frac{1}{2})$

- one doublet at **EW** scale
 - fine-tuning or some mechanism
 - issue: doublet-triplet splitting
- an admixture of flavor eigenstates (v's required by fermion fit)

Yukawa relations at M_{GUT} :

$$\begin{aligned} M_U &= Y_{10} v_{10}^u + \tilde{Y}_{10} v_{10}^{d*} + Y_{126} v_{126}^u \\ M_D &= Y_{10} v_{10}^d + \tilde{Y}_{10} v_{10}^{u*} + Y_{126} v_{126}^d \\ M_E &= Y_{10} v_{10}^d + \tilde{Y}_{10} v_{10}^{u*} - 3 Y_{126} v_{126}^d \end{aligned}$$

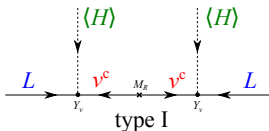
Fit: works already with Y_{10}, Y_{126}

BSM pheno aspects 1 — neutrino physics

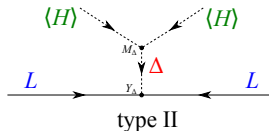
- $\nu^c \in \text{SO}(10), E_6 \rightarrow$ automatically a theory of neutrino mass

BSM pheno aspects 1 — neutrino physics

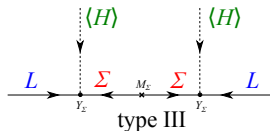
- $\nu^c \in \text{SO}(10), E_6 \rightarrow$ automatically a theory of neutrino mass
- Neutrino masses: $LLHH$ obtained via **see-saw**



$$\nu^c \sim (1, 1, 0)$$



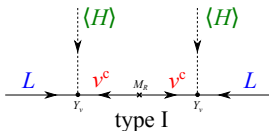
$$\Delta \sim (1, 3, +1)$$



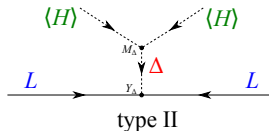
$$\Sigma \sim (1, 3, 0) \quad (23)$$

BSM pheno aspects 1 — neutrino physics

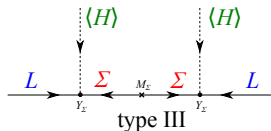
- $\nu^c \in \text{SO}(10), E_6 \rightarrow$ automatically a theory of neutrino mass
- Neutrino masses: $LLHH$ obtained via **see-saw**



$$\nu^c \sim (1, 1, 0)$$



$$\Delta \sim (1, 3, +1)$$



$$\Sigma \sim (1, 3, 0) \quad (23)$$

- Example: type I + II in $\text{SO}(10)$ with $10 \oplus 126$

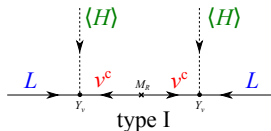
$$Y_\nu v_{EW} = Y_{10} v_{10}^u + \tilde{Y}_{10} v_{10}^{d*} - 3 Y_{126} v_{126}^u \quad M_R = Y_{126} \langle 126 \rangle \quad (24)$$

$$Y_\Delta = Y_{126} \quad M_\Delta : \langle 126 \rangle \cdot 126 \cdot 126^{*2} \quad (25)$$

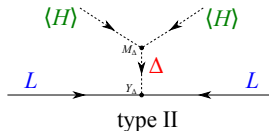
$$\nu^c \subset 16_F, \Delta \subset 126$$

BSM pheno aspects 1 — neutrino physics

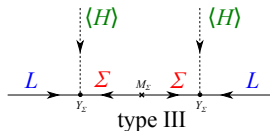
- $\nu^c \in \text{SO}(10), E_6 \rightarrow$ automatically a theory of neutrino mass
- Neutrino masses: $LLHH$ obtained via **see-saw**



$$\nu^c \sim (1, 1, 0)$$



$$\Delta \sim (1, 3, +1)$$



$$\Sigma \sim (1, 3, 0) \quad (23)$$

- Example: type I + II in $\text{SO}(10)$ with $10 \oplus 126$

$$Y_\nu \nu_{EW} = Y_{10} \nu_{10}^u + \tilde{Y}_{10} \nu_{10}^{d*} - 3 Y_{126} \nu_{126}^u \quad M_R = Y_{126} \langle 126 \rangle \quad (24)$$

$$Y_\Delta = Y_{126} \quad M_\Delta : \langle 126 \rangle \cdot 126 \cdot 126^{*2} \quad (25)$$

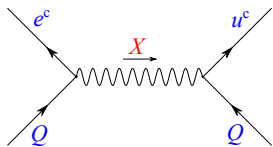
$$\nu^c \subset 16_F, \Delta \subset 126$$

- See-saw scale $\sim 10^{12}$ GeV from intermediate breaking stage

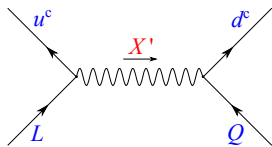
$$M_\nu = -Y_\nu^T M_R^{-1} Y_\nu \nu_{EW}^2 + Y_\Delta \langle \Delta \rangle_{\text{induced}} \quad (26)$$

BSM pheno aspects 2 — proton decay

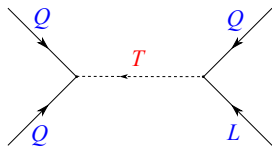
- Operators $qqql$: violate B and L number (preserve $B - L$)



$$X \sim (3, 2, -\frac{5}{6})$$



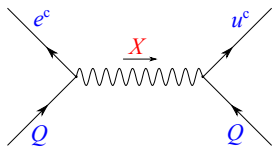
$$X' \sim (3, 2, +\frac{1}{6})$$



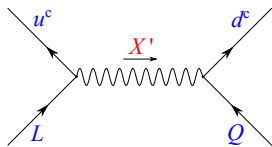
$$T \sim (3, 1, -\frac{1}{3}) \quad (27)$$

BSM pheno aspects 2 — proton decay

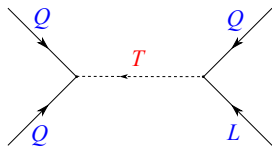
- Operators $qqql$: violate B and L number (preserve $B - L$)



$$X \sim (3, 2, -\frac{5}{6})$$



$$X' \sim (3, 2, +\frac{1}{6})$$

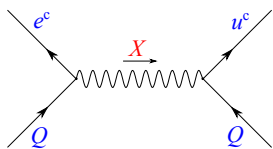


$$T \sim (3, 1, -\frac{1}{3}) \quad (27)$$

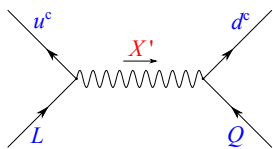
- Gauge-mediated proton decay: **always** found in GUT adjoints
→ $X \subset 24$ of $SU(5)$; $X, X' \subset 45$ of $SO(10)$

BSM pheno aspects 2 — proton decay

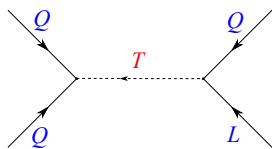
- Operators $qqql$: violate B and L number (preserve $B - L$)



$$X \sim (3, 2, -\frac{5}{6})$$



$$X' \sim (3, 2, +\frac{1}{6})$$

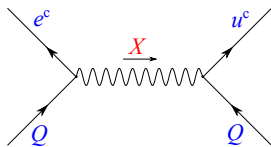


$$T \sim (3, 1, -\frac{1}{3}) \quad (27)$$

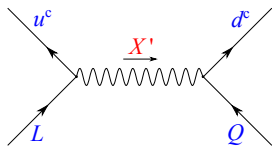
- Gauge-mediated proton decay: **always** found in GUT adjoints
→ $X \subset 24$ of $SU(5)$; $X, X' \subset 45$ of $SO(10)$
- Scalar-mediated proton decay:
→ T usually together with H in GUT [5, 45 of $SU(5)$]

BSM pheno aspects 2 — proton decay

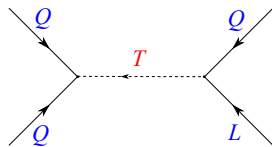
- Operators $qqql$: violate B and L number (preserve $B - L$)



$$X \sim (3, 2, -\frac{5}{6})$$



$$X' \sim (3, 2, +\frac{1}{6})$$



$$T \sim (3, 1, -\frac{1}{3}) \quad (27)$$

- Gauge-mediated proton decay: **always** found in GUT adjoints
→ $X \subset 24$ of $SU(5)$; $X, X' \subset 45$ of $SO(10)$
- Scalar-mediated proton decay:
→ T usually together with H in GUT [5, 45 of $SU(5)$]
- Crude estimate:

$$A_{\text{gauge}} \propto \frac{g^2}{M_X^2}, \quad A_{\text{scalar}} \propto \frac{y^2}{M_T^2}, \quad \Gamma_p \approx \frac{g^4 m_p^5}{M_{\text{GUT}}^4} \quad (28)$$

- SUSY: proton decay faster ($D = 5$ operators)

BSM pheno aspects 3 — proton decay experimentally

- Process: **proton** \rightarrow **meson** + **lepton**

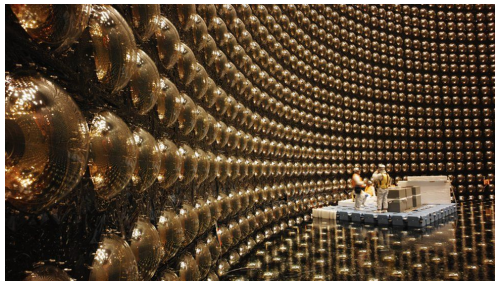
e.g. channels	τ_p limit
$p \rightarrow \pi^0 e^+$	$2.4 \cdot 10^{34}$ yr
$p \rightarrow \pi^0 \mu^+$	$1.6 \cdot 10^{34}$ yr
$p \rightarrow \pi^+ \bar{\nu}$	$2.8 \cdot 10^{32}$ yr
$p \rightarrow K^+ \bar{\nu}$	$6.6 \cdot 10^{33}$ yr

BSM pheno aspects 3 — proton decay experimentally

- Process: **proton** \rightarrow **meson** + **lepton**

e.g. channels	τ_p limit
$p \rightarrow \pi^0 e^+$	$2.4 \cdot 10^{34}$ yr
$p \rightarrow \pi^0 \mu^+$	$1.6 \cdot 10^{34}$ yr
$p \rightarrow \pi^+ \bar{\nu}$	$2.8 \cdot 10^{32}$ yr
$p \rightarrow K^+ \bar{\nu}$	$6.6 \cdot 10^{33}$ yr

experiment	date	volume	exposure
Super-K	1996	27 kt	450 kt yr
Hyper-K	2027	187 kt	/



Super-K: water tank with photomultiplier tubes

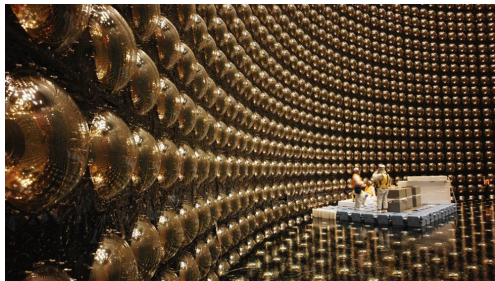
BSM pheno aspects 3 — proton decay experimentally

- Process: **proton** \rightarrow **meson** + **lepton**

e.g. channels	τ_p limit	experiment	date	volume	exposure
$p \rightarrow \pi^0 e^+$	$2.4 \cdot 10^{34}$ yr	Super-K	1996	27 kt	450 kt yr
$p \rightarrow \pi^0 \mu^+$	$1.6 \cdot 10^{34}$ yr	Hyper-K	2027	187 kt	/
$p \rightarrow \pi^+ \bar{\nu}$	$2.8 \cdot 10^{32}$ yr				
$p \rightarrow K^+ \bar{\nu}$	$6.6 \cdot 10^{33}$ yr				

- Intuition:

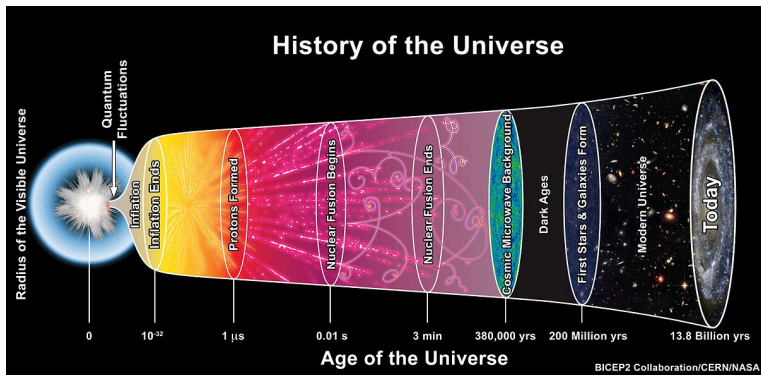
- 1 **proton** for 10^{33} yr
 $\leftrightarrow 10^{33}$ **protons** for 1 yr
- 10^{27} **protons** \sim 1 kg
 $\rightarrow 10^{33}$ **protons** \sim 1 kt
- 1 kt water $\sim (10 \text{ m})^3$
- 10^{34} yr $\rightarrow M_{\text{GUT}} \sim 10^{15.2}$ GeV



Super-K: water tank with photomultiplier tubes

BSM pheno aspects 4 — Cosmology

- Early universe is hot: **GUT era**, cosmology implications



BSM pheno aspects 4 — Cosmology

- Early universe is hot: **GUT era**, cosmology implications

- Generic GUT prediction:

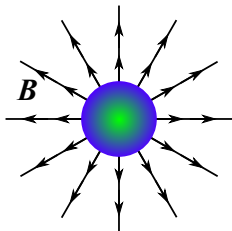
't Hooft–Polyakov monopoles

→ $\pi_2(G/H) \neq \{1\}$

→ core with “knotted” GUT gauge fields

→ size $\sim M_{\text{GUT}}^{-1}$

→ non-observation: # diluted by inflation



BSM pheno aspects 4 — Cosmology

- Early universe is hot: **GUT era**, cosmology implications

- Generic GUT prediction:

't Hooft–Polyakov monopoles

→ $\pi_2(G/H) \neq \{1\}$

→ core with “knotted” GUT gauge fields

→ size $\sim M_{\text{GUT}}^{-1}$

→ non-observation: # diluted by inflation

- Cosmology aspects addressed within GUTs:

1 inflation (usually below M_{GUT})

2 leptogenesis, baryogenesis

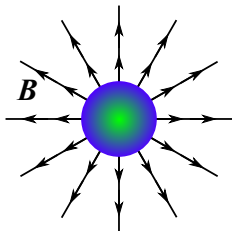
3 dark matter

4 symmetry breaking: phase transitions

→ formation of vacuum bubbles

→ formation of cosmic strings, domain walls, monopoles

→ gravitational wave signals



Intermezzo — a retrospective overview up to now

1 Model choice:

→ unified gauge group G , irreps of fermions \mathcal{F} and scalars \mathcal{S}

Intermezzo — a retrospective overview up to now

1 Model choice:

→ unified gauge group G , irreps of fermions \mathcal{F} and scalars \mathcal{S}

2 A **complete** model:

- Should reproduce (SM+massive neutrinos) at low energy **parameters**: unification, \mathcal{F} masses/mixings, Higgs sector
- BSM observables: **proton decay**, others in specific models
- Will have implications for **cosmology**

Intermezzo — a retrospective overview up to now

1 Model choice:

→ unified gauge group G , irreps of fermions \mathcal{F} and scalars \mathcal{S}

2 A **complete** model:

- Should reproduce (SM+massive neutrinos) at low energy **parameters**: unification, \mathcal{F} masses/mixings, Higgs sector
- BSM observables: **proton decay**, others in specific models
- Will have implications for **cosmology**

3 Many models on the market, no consensus model . . . work in progress

Intermezzo — a retrospective overview up to now

1 Model choice:

→ unified gauge group G , irreps of fermions \mathcal{F} and scalars \mathcal{S}

2 A **complete** model:

- Should reproduce (SM+massive neutrinos) at low energy **parameters**: unification, \mathcal{F} masses/mixings, Higgs sector
- BSM observables: **proton decay**, others in specific models
- Will have implications for **cosmology**

3 Many models on the market, no consensus model . . . work in progress

4 Challenges and limitations for model building:

- Representations too big? Unnecessary DOF, perturbativity issues
- Precision calculations difficult, can have many new parameters

Intermezzo — a retrospective overview up to now

1 Model choice:

→ unified gauge group G , irreps of fermions \mathcal{F} and scalars \mathcal{S}

2 A **complete** model:

- Should reproduce (SM+massive neutrinos) at low energy **parameters**: unification, \mathcal{F} masses/mixings, Higgs sector
- BSM observables: **proton decay**, others in specific models
- Will have implications for **cosmology**

3 Many models on the market, no consensus model . . . work in progress

4 Challenges and limitations for model building:

- Representations too big? Unnecessary DOF, perturbativity issues
- Precision calculations difficult, can have many new parameters

5 Bottom up view [language of SM symmetry]:

- new DOF that complete larger building blocks
- couplings interrelated
- will/can relate BSM phenomena (interesting!)

Axion in GUTs 1 — general considerations

- Axion: a solution to the strong-CP problem, DM candidate

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu}, \quad |\theta_{\text{eff}}| = \left| \frac{\langle a \rangle}{f_a} \right| \lesssim 10^{-10} \quad (29)$$

Axion in GUTs 1 — general considerations

- Axion: a solution to the strong-CP problem, DM candidate

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu}, \quad |\theta_{\text{eff}}| = \left| \frac{\langle a \rangle}{f_a} \right| \lesssim 10^{-10} \quad (29)$$

- Accommodating the axion into GUT:

(a) implement a global Peccei-Quinn $U(1)_{PQ}$

(b) find a home for a complex PQ-charged SM-singlet scalar S :

$$S \sim (1, 1, 0)(q) \quad (30)$$

Axion in GUTs 1 — general considerations

- Axion: a solution to the strong-CP problem, DM candidate

$$\mathcal{L} \supset \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu}, \quad |\theta_{\text{eff}}| = \left| \frac{\langle a \rangle}{f_a} \right| \lesssim 10^{-10} \quad (29)$$

- Accommodating the axion into GUT:

(a) implement a global Peccei-Quinn $U(1)_{PQ}$

(b) find a home for a complex PQ-charged SM-singlet scalar S :

$$S \sim (1, 1, 0)(q) \quad (30)$$

- Axion a : pseudo-Goldstone of $U(1)_{PQ}$, mass from QCD anomaly

$$m_a \approx m_\pi f_\pi / f_a \sim 10 \text{ neV} \cdot (10^{15} \text{ GeV} / f_a) \quad (31)$$

→ $f_a \sim \langle S \rangle$... connection to scales in GUTs

→ relevant for experiment: via $g_{a\gamma\gamma}, g_{aD} \propto 1/f_a$

$$\mathcal{L} \supset g_{a\gamma\gamma} \frac{a}{4} F_{\mu\nu} \tilde{F}^{\mu\nu} - \frac{i}{2} g_{aD} a \bar{\Psi}_n \sigma_{\mu\nu} \gamma_5 \Psi_n F^{\mu\nu} \quad (32)$$

Axion in GUTs 2 — examples in SU(5)

- Example 1A: WGG (Wise-Georgi-Glashow) based on DFS Z
 - In SM language: need 2 **weak doublets** and a **SM-singlet**
 - $SU(5) \times U(1)_{PQ}$

$$3 \times 10_F(-\frac{1}{2}), \quad 3 \times \bar{5}_F(-\frac{1}{2}), \quad 5_1(1), \quad 5_2(-1), \quad 24_C(-1) \quad (33)$$

$$\mathcal{L} \supset Y_{10} 10_F 10_F 5_1 + Y_5 10_F \bar{5}_F 5_2^* - V(5_1, 5_2, 24) \quad (34)$$

- VEVs: $\langle S \rangle \equiv \langle 24 \rangle_{GUT}, \langle 5_1 \rangle_{EW}, \langle 5_2 \rangle_{EW}$
- Axion a : orbital component of $S \subset 24$ (to $\mathcal{O}(M_{EW}/M_{GUT})$)
- nowadays: not realistic (Yukawa sector too simple, no neutrino mass)

Axion in GUTs 2 — examples in SU(5)

- Example 1A: WGG (Wise-Georgi-Glashow) based on DFS Z
 - In SM language: need 2 **weak doublets** and a **SM-singlet**
 - $SU(5) \times U(1)_{PQ}$

$$3 \times 10_F(-\frac{1}{2}), \quad 3 \times \bar{5}_F(-\frac{1}{2}), \quad 5_1(1), \quad 5_2(-1), \quad 24_C(-1) \quad (33)$$

$$\mathcal{L} \supset Y_{10} 10_F 10_F 5_1 + Y_5 10_F \bar{5}_F 5_2^* - V(5_1, 5_2, 24) \quad (34)$$

- VEVs: $\langle S \rangle \equiv \langle 24 \rangle_{GUT}, \langle 5_1 \rangle_{EW}, \langle 5_2 \rangle_{EW}$
- Axion a : orbital component of $S \subset 24$ (to $\mathcal{O}(M_{EW}/M_{GUT})$)
- nowadays: not realistic (Yukawa sector too simple, no neutrino mass)

- Realistic $SU(5) \times U(1)_{PQ}$ examples:

model	additions to 1A	ν mass	m_a [neV]
1B	$24_F(-\frac{1}{2})$	seesaw I+III	[4.8, 6.6]
1C	$15_F(-\frac{1}{2}) \oplus \bar{15}_F(-\frac{1}{2}) \oplus 35(-1)$	at 1-loop	[0.1, 4.7]

Constraints on m_a : unification, proton decay, collider

Axion in GUTs 3 — examples in SO(10)

- Example 2: a realistic $SO(10) \times U(1)_{PQ}$

$$3 \times 16_F(-\frac{1}{2}), \quad 10(+1), \quad 126(-1), \quad 210(-2) \quad (35)$$

Axion in GUTs 3 — examples in SO(10)

- Example 2: a realistic $SO(10) \times U(1)_{PQ}$

$$3 \times 16_F(-\frac{1}{2}), \quad 10(+1), \quad 126(-1), \quad 210(-2) \quad (35)$$

$$SO(10) \xrightarrow{210} 4_C 2_L 2_R \xrightarrow{126} 3_C 2_L 1_Y \xrightarrow{10,126} 3_C 1_{EM} \quad (36)$$

Axion in GUTs 3 — examples in SO(10)

- Example 2: a realistic SO(10) \times U(1)_{PQ}

$$3 \times 16_F(-\frac{1}{2}), \quad 10(+1), \quad 126(-1), \quad 210(-2) \quad (35)$$

$$\text{SO}(10) \xrightarrow{210} 4_C 2_L 2_R \xrightarrow{126} 3_C 2_L 1_Y \xrightarrow{10,126} 3_C 1_{\text{EM}} \quad (36)$$

$$\mathcal{L} \supset Y_{10} 16_F 16_F 10 + Y_{126} 16_F 16_F 126^* - V(10, 126, 210) \quad (37)$$

→ PQ broken by 210 \supset S

→ \tilde{Y}_{10} is now forbidden → more predictive setup

Axion in GUTs 3 — examples in SO(10)

- Example 2: a realistic $SO(10) \times U(1)_{PQ}$

$$3 \times 16_F(-\frac{1}{2}), \quad 10(+1), \quad 126(-1), \quad 210(-2) \quad (35)$$

$$SO(10) \xrightarrow{210} 4_C 2_L 2_R \xrightarrow{126} 3_C 2_L 1_Y \xrightarrow{10,126} 3_C 1_{EM} \quad (36)$$

$$\mathcal{L} \supset Y_{10} 16_F 16_F 10 + Y_{126} 16_F 16_F 126^* - V(10, 126, 210) \quad (37)$$

→ PQ broken by $210 \supset S$

→ \tilde{Y}_{10} is now forbidden → more predictive setup

- Some conceptual alternatives:

- new in SO(10): relate f_a to intermediate scales → heavier axion
→ W is highest PQ-breaking VEV

$$SO(10) \xrightarrow{V} G_1 \xrightarrow{W} G_2 \xrightarrow{Z} 3_C 2_L 1_Y \quad (38)$$

- S is G -singlet: VEV at any scale → not predictive for axion mass

Grand Unified Theories: an attractive framework

- (0) Tentative experimental indication: SM gauge couplings approximately unify in correct high energy window
- (1) Conceptual/philosophical advantages over Standard Model
- (2) Can incorporate and relate different SM or BSM phenomena (particle physics and cosmology)
- (3) Offer possible windows to physics at super-high energies
 - (3a) proton decay
 - (3b) model-specific phenomena . . .
 - (3c) . . . including axions

Grand Unified Theories: an attractive framework

- (0) Tentative experimental indication: SM gauge couplings approximately unify in correct high energy window
- (1) Conceptual/philosophical advantages over Standard Model
- (2) Can incorporate and relate different SM or BSM phenomena (particle physics and cosmology)
- (3) Offer possible windows to physics at super-high energies
 - (3a) proton decay
 - (3b) model-specific phenomena . . .
 - (3c) . . . including axions

Thank you for your attention!