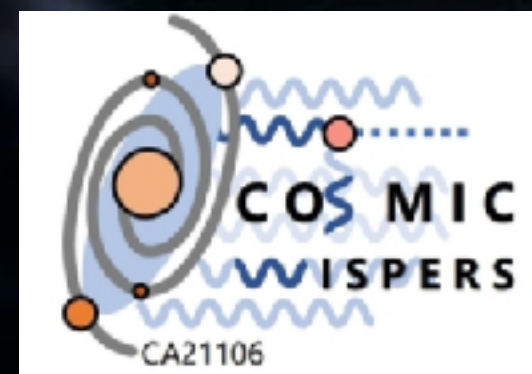


Bosonic Halos: Axion Stars and Dark Matter Capture

Joshua Eby
Oskar Klein Centre
Stockholm University

Barolo Astroparticle Meeting (BAM)
2024/06/13

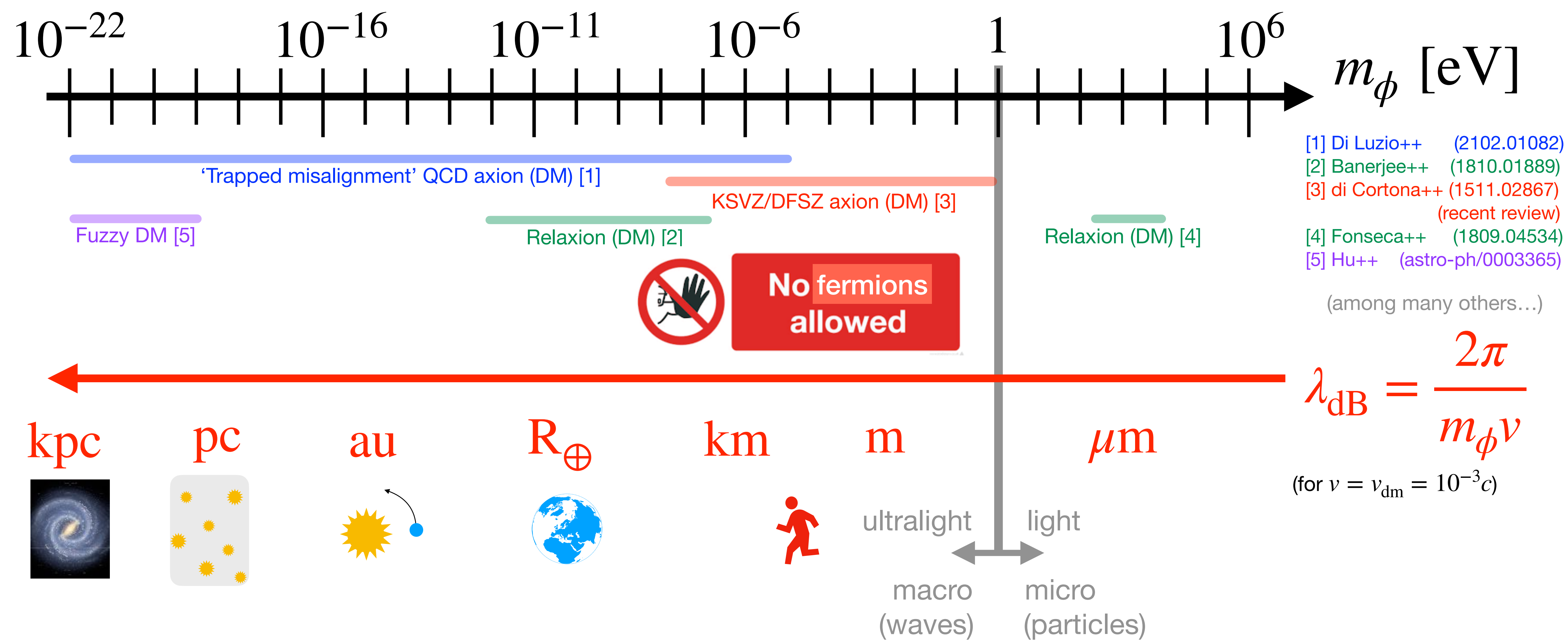


DALL-E 3 illustration
"Bosenova"



Stockholms
universitet

Light and Ultralight Dark Matter



de Broglie wavelength λ_{dB}

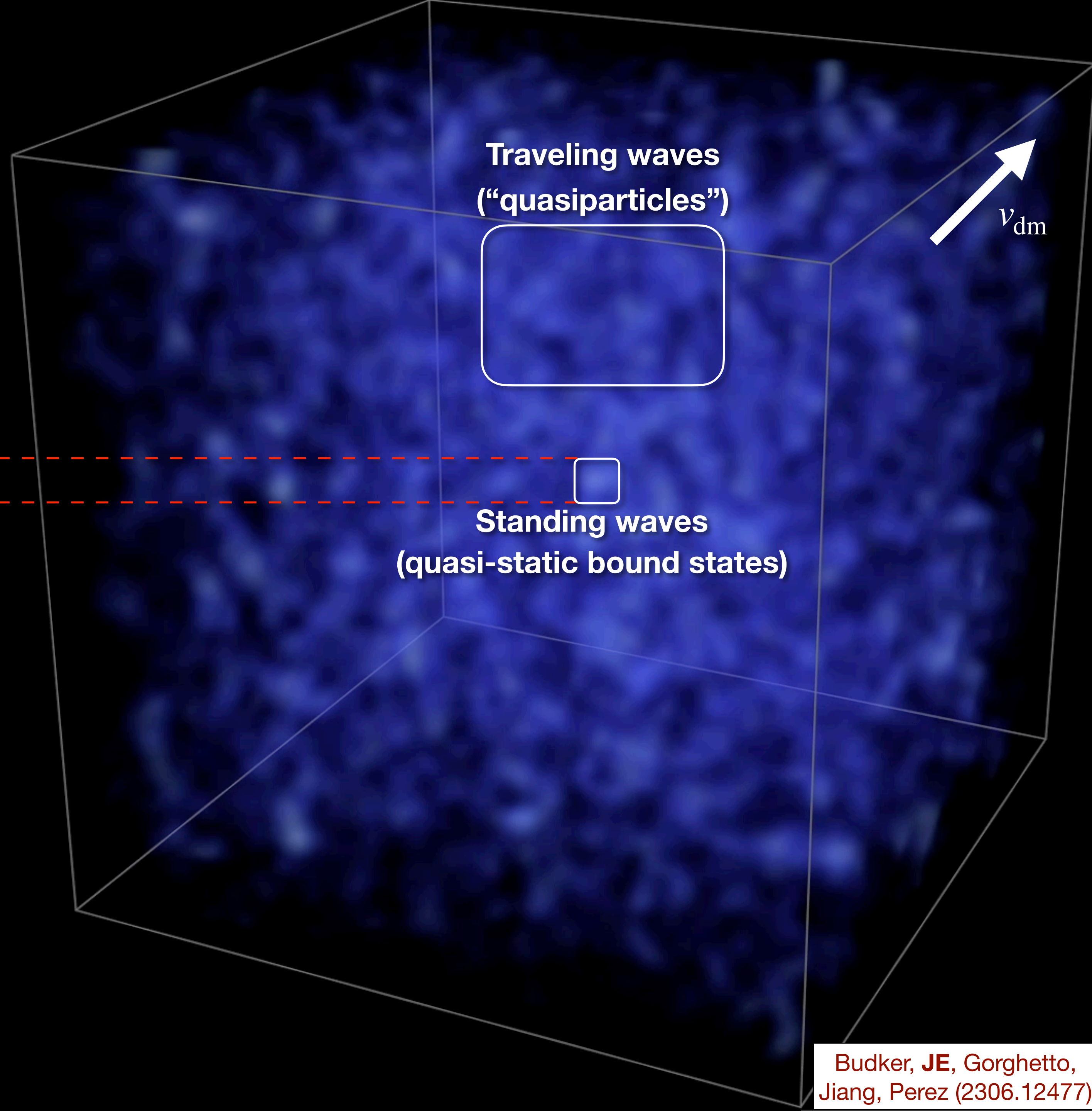


Wave amplitude \iff DM density ρ

Average local density

$$\rho_{\text{dm}} = 0.4 \text{ GeV}/\text{cm}^3$$

with variations on scales of order λ_{dB}



**Traveling waves
("quasiparticles")**



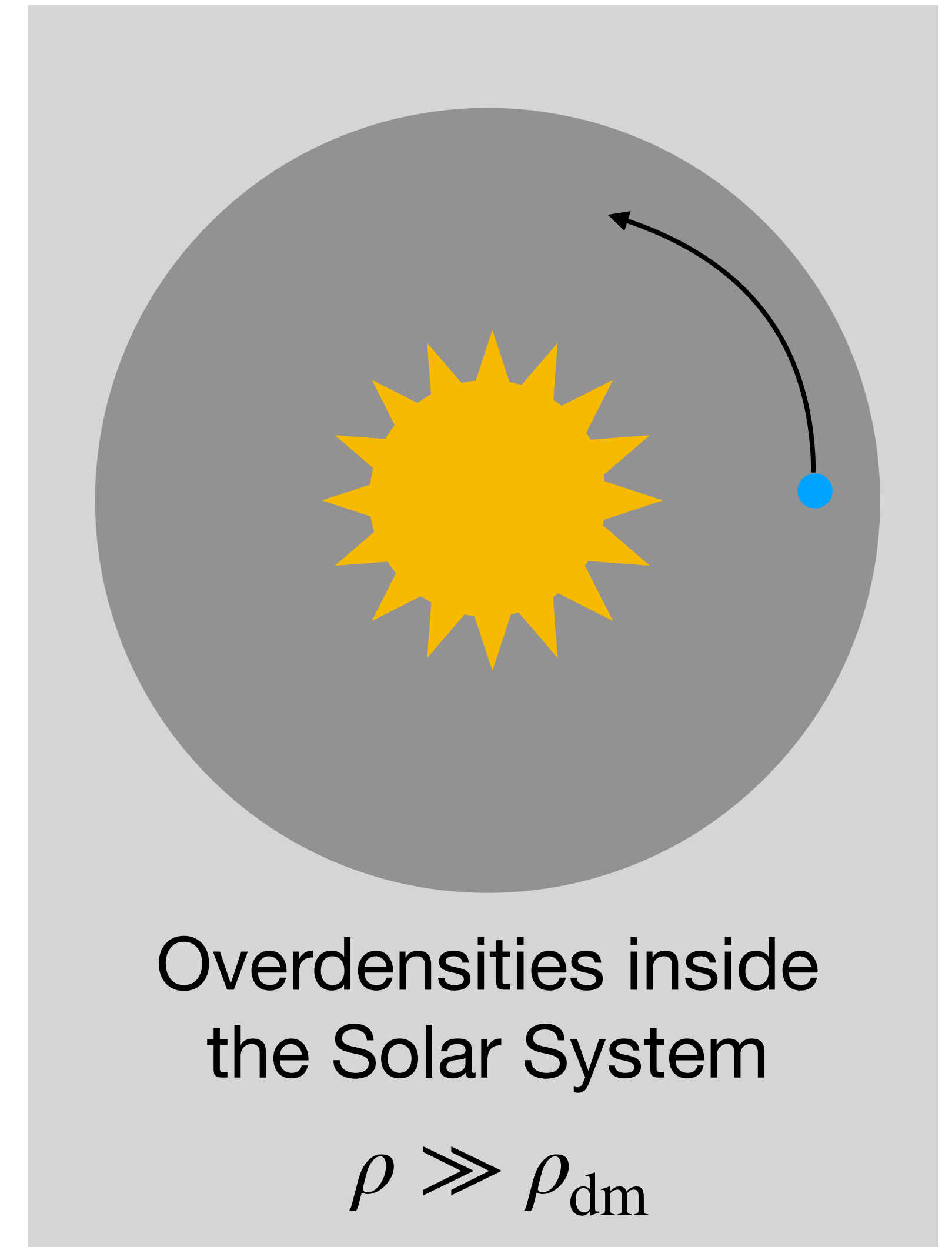
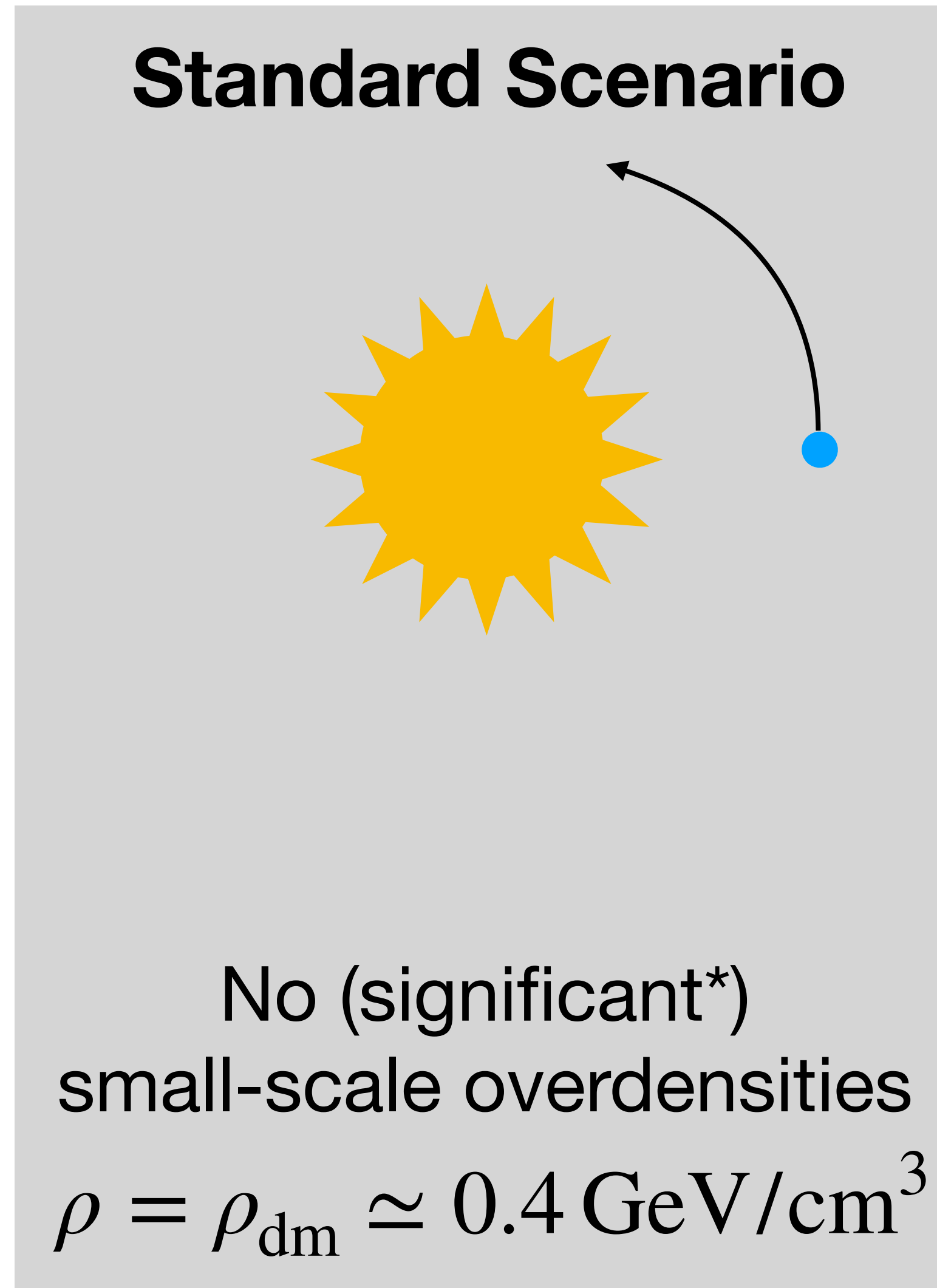
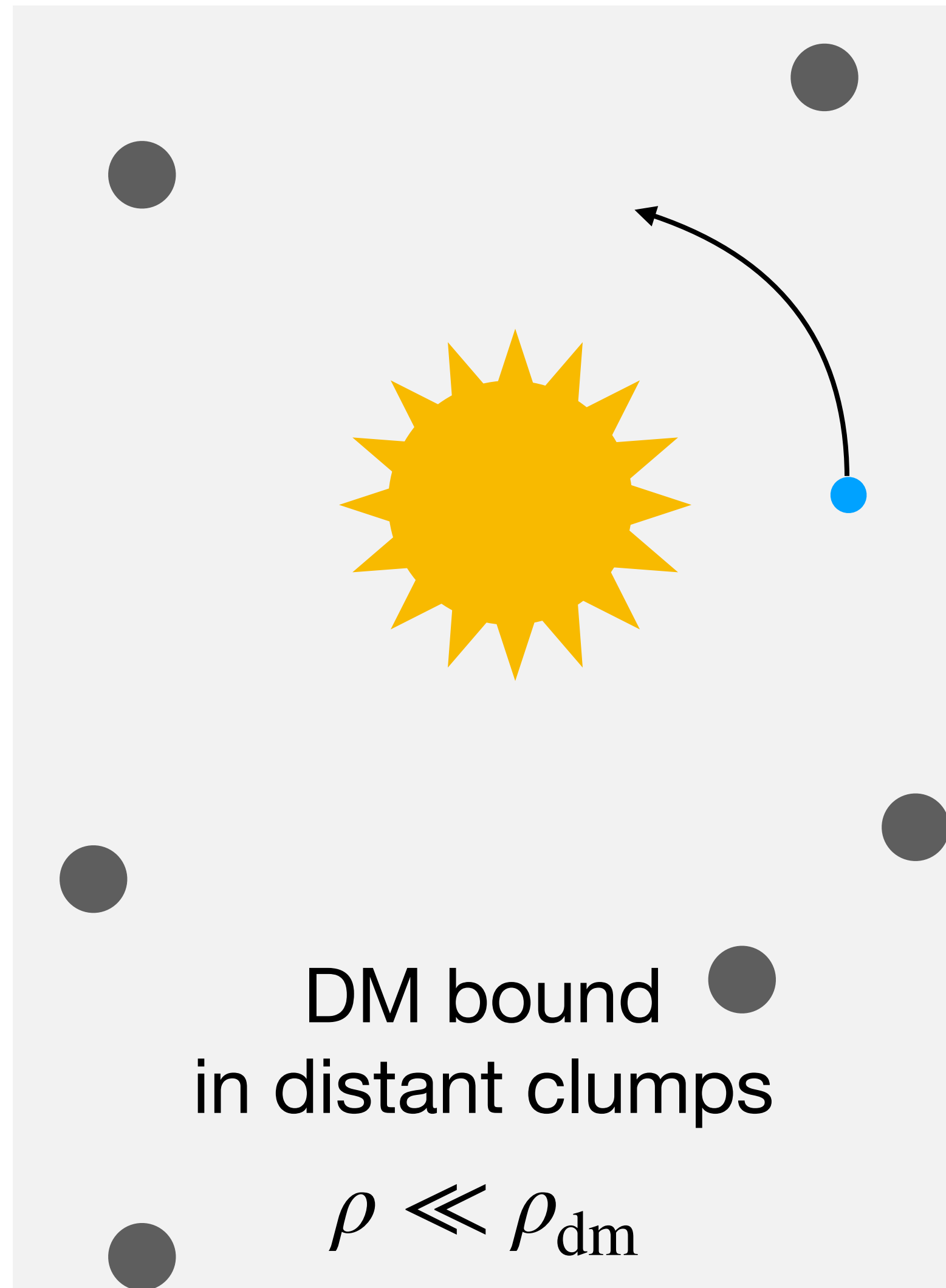
v_{dm}

**Standing waves
(quasi-static bound states)**



The Very Local DM Density

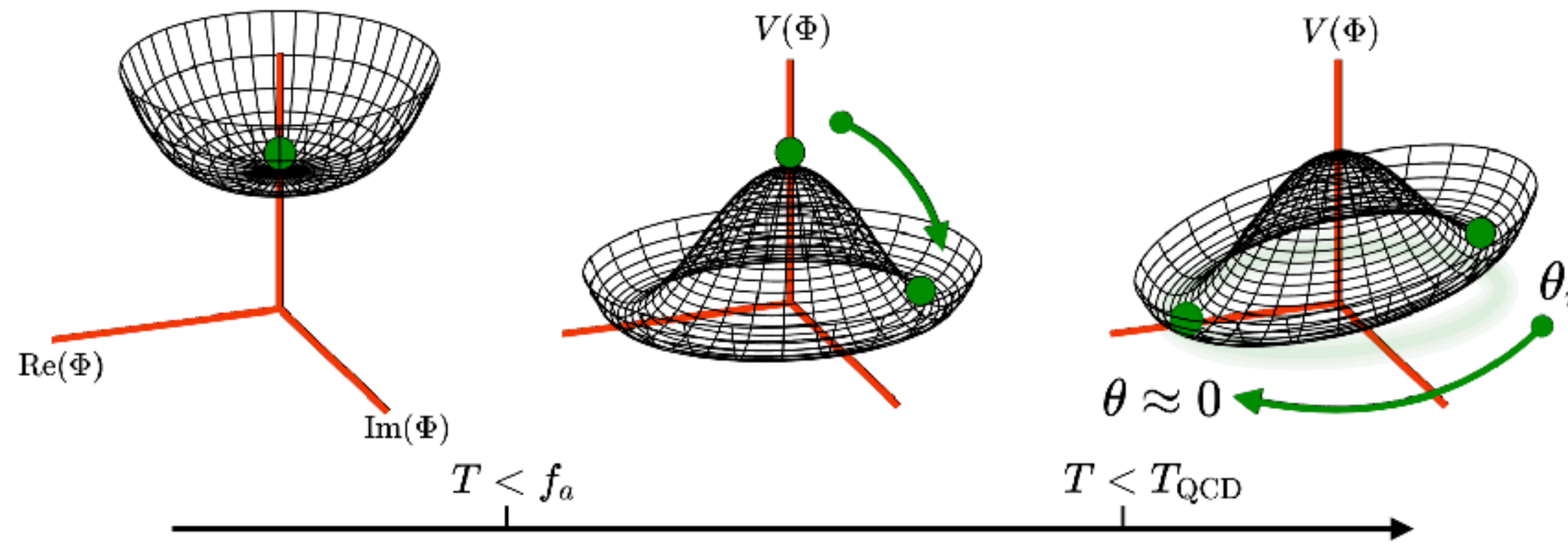
(inside the solar system)



*for ULDM, always $\exists \mathcal{O}(1)$ fluctuations

Axion Cosmology

Recent review:
O'Hare (2403.17697)



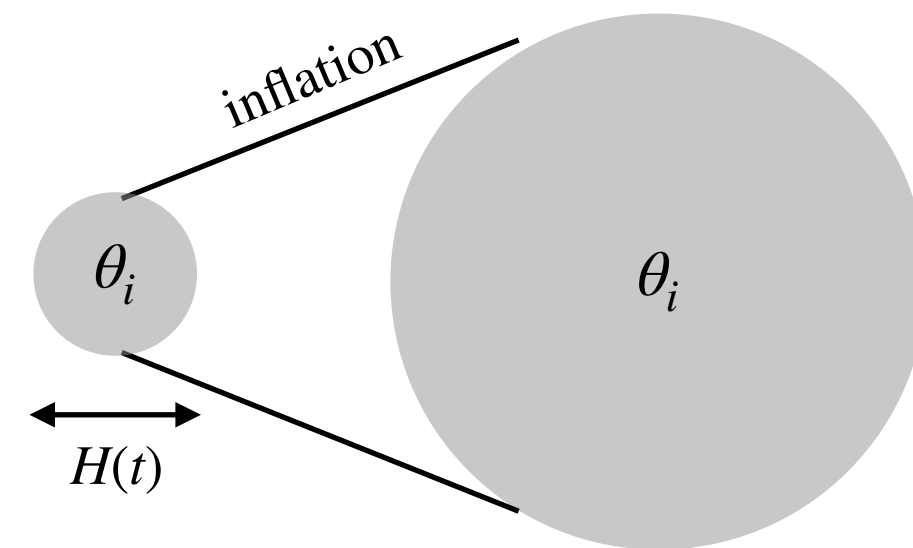
$$V(\Phi) = \lambda_{\Phi} \left(|\Phi|^2 - \frac{f_a^2}{2} \right)^2 + \Lambda^4 \left(1 - \cos \frac{\phi}{f_a} \right)$$

QCD axion:

$$V_{\theta}(\phi) = \left(\theta_{\text{QCD}} + \frac{\phi}{f_a} \right) G^{\mu\nu} \tilde{G}_{\mu\nu} \longrightarrow 0 \quad \Lambda_{\text{QCD}}^4 \simeq m_{\phi}^2 f_a^2$$

Pre-inflationary

$$f_a \gtrsim H_I$$

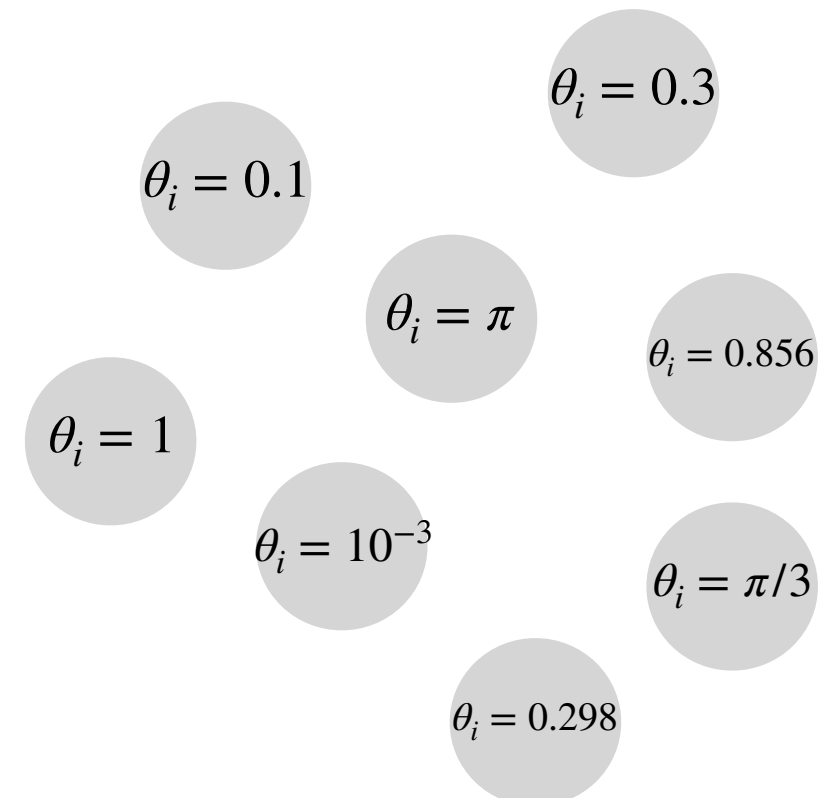


$$\longrightarrow \Omega_a^{\text{misalignment}}(f_a, \theta_i) \simeq 0.1 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \theta_i^2$$

?

Post-inflationary

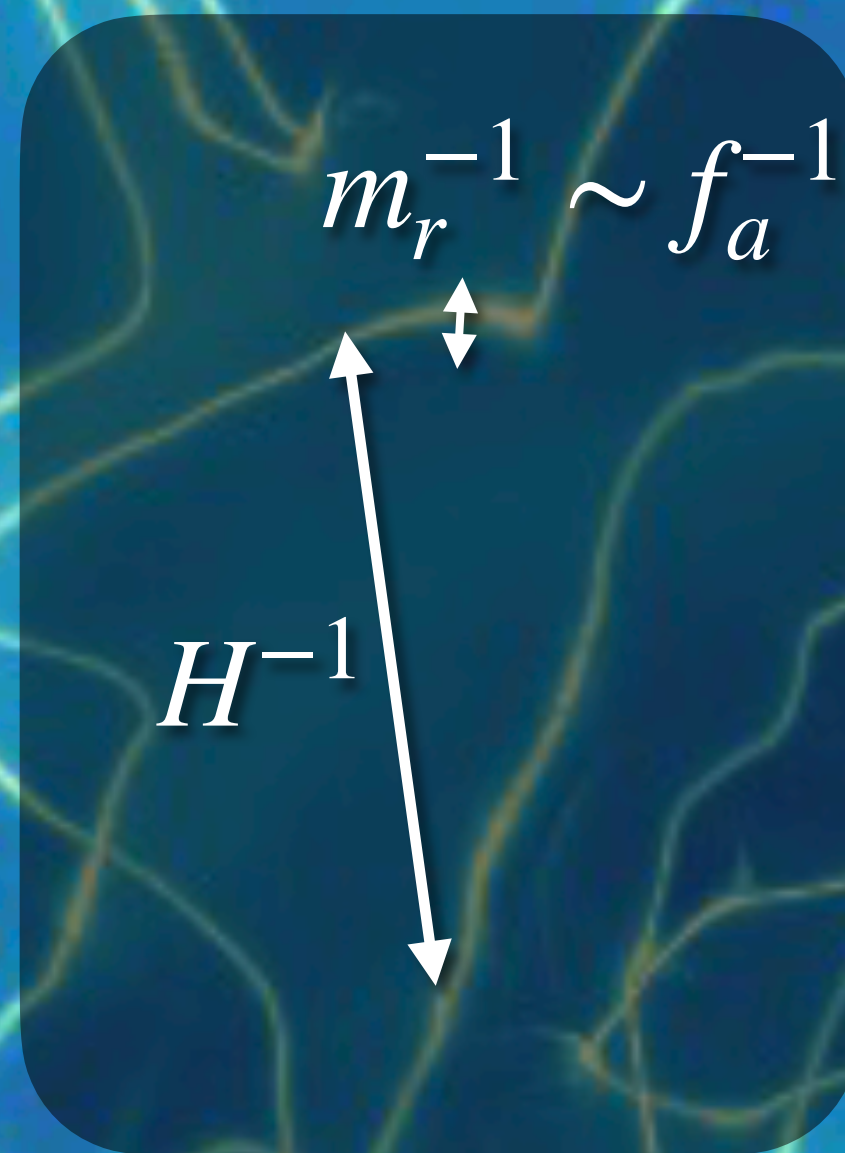
$$f_a \lesssim H_I$$



$$\longrightarrow \Omega_a^{\text{misalignment}}(f_a) \simeq 0.1 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \langle \theta_i^2 \rangle$$

$\pi^2/3$

predictive !

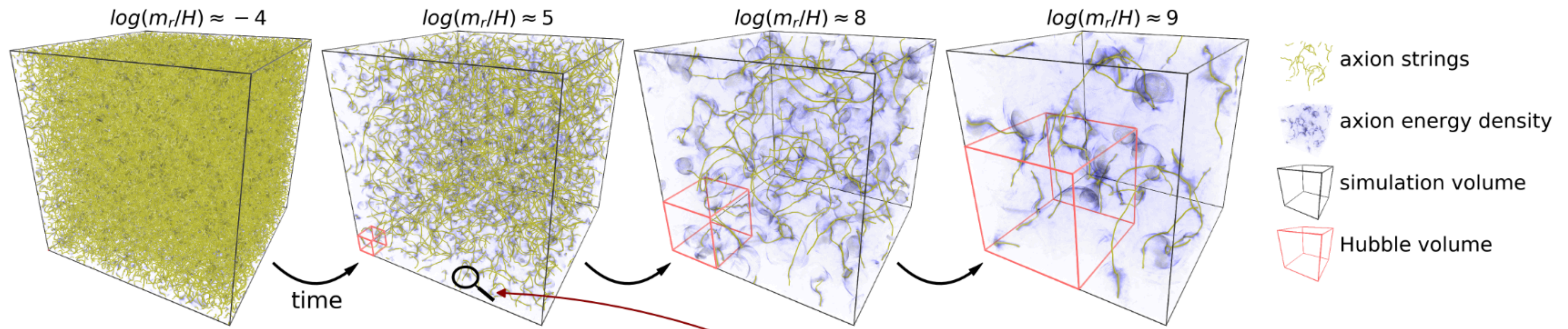


World-leading simulations: $\log \frac{m_r}{H} \simeq 9$

Physical: $\log \frac{m_r}{H} \simeq 70$

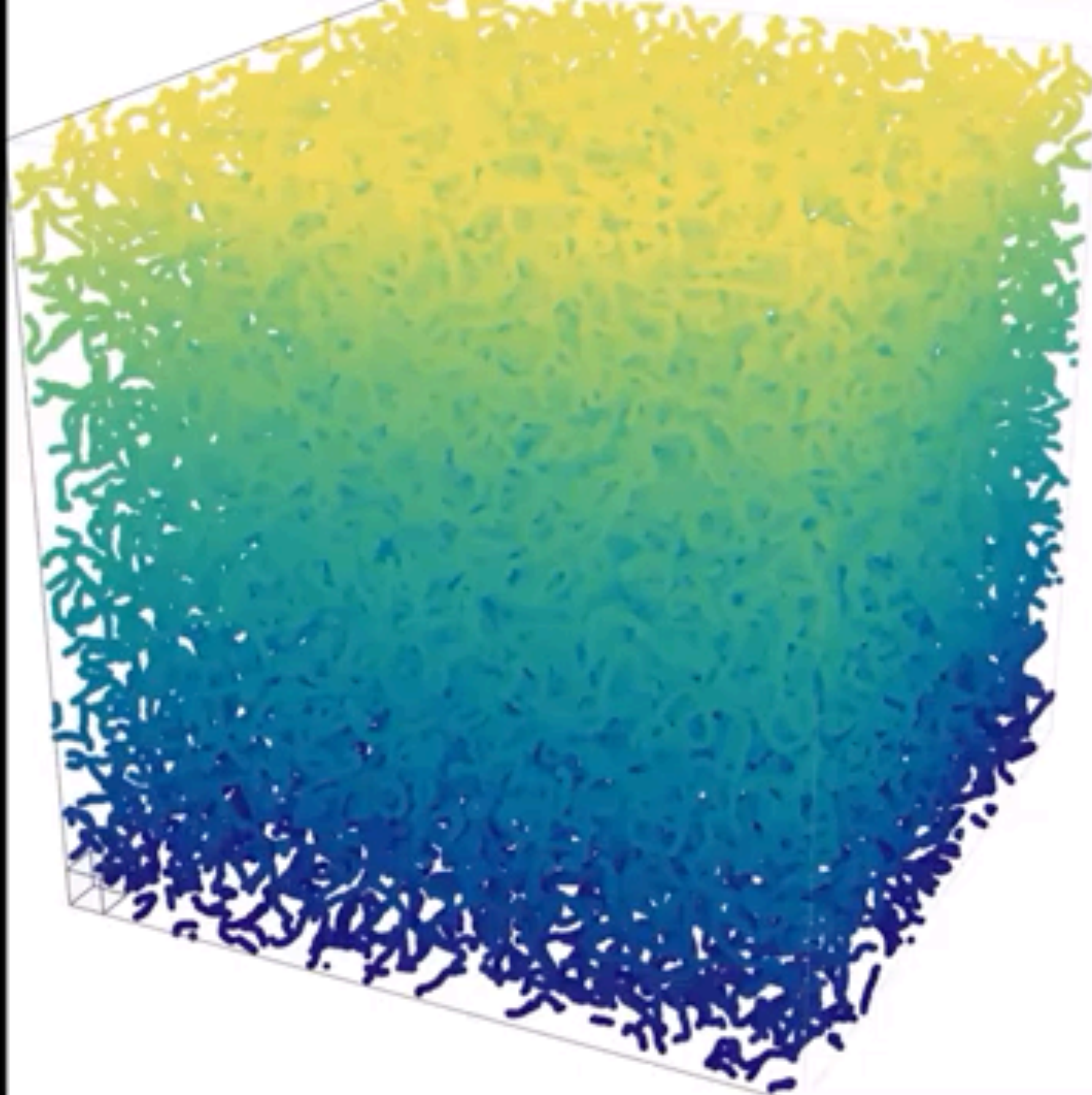
Buschmann, Foster, Hook, Peterson,
Willcox, Zhang, Safdi (2108.05368)

Another view



World-leading simulations: $\log \frac{m_r}{H} \simeq 9$

Physical: $\log \frac{m_r}{H} \simeq 70$



Gorghetto, Hardy, Villadoro
(2007.04990)

Movie via Marco Gorghetto
<https://www.youtube.com/watch?v=DbvM7emtodo>

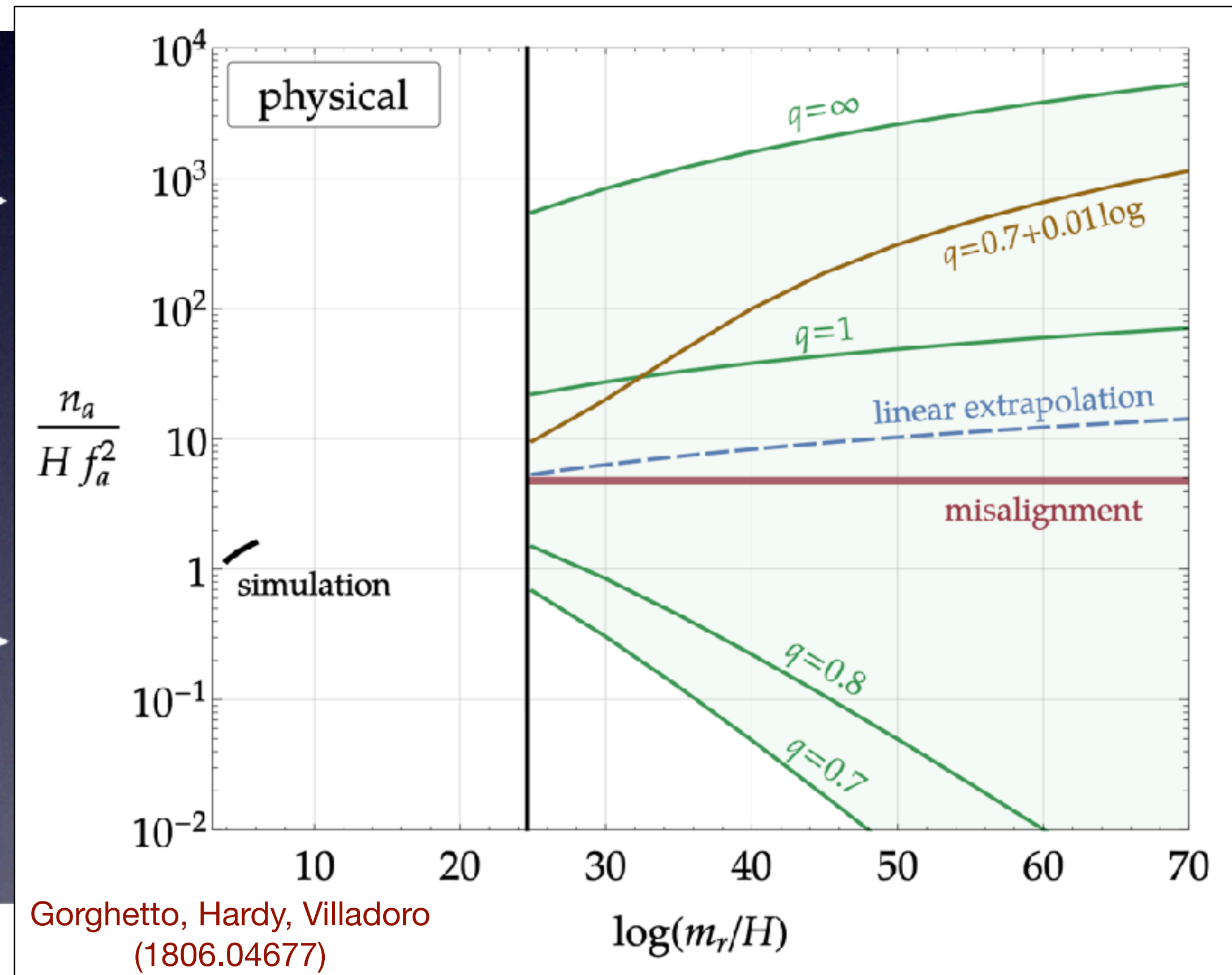
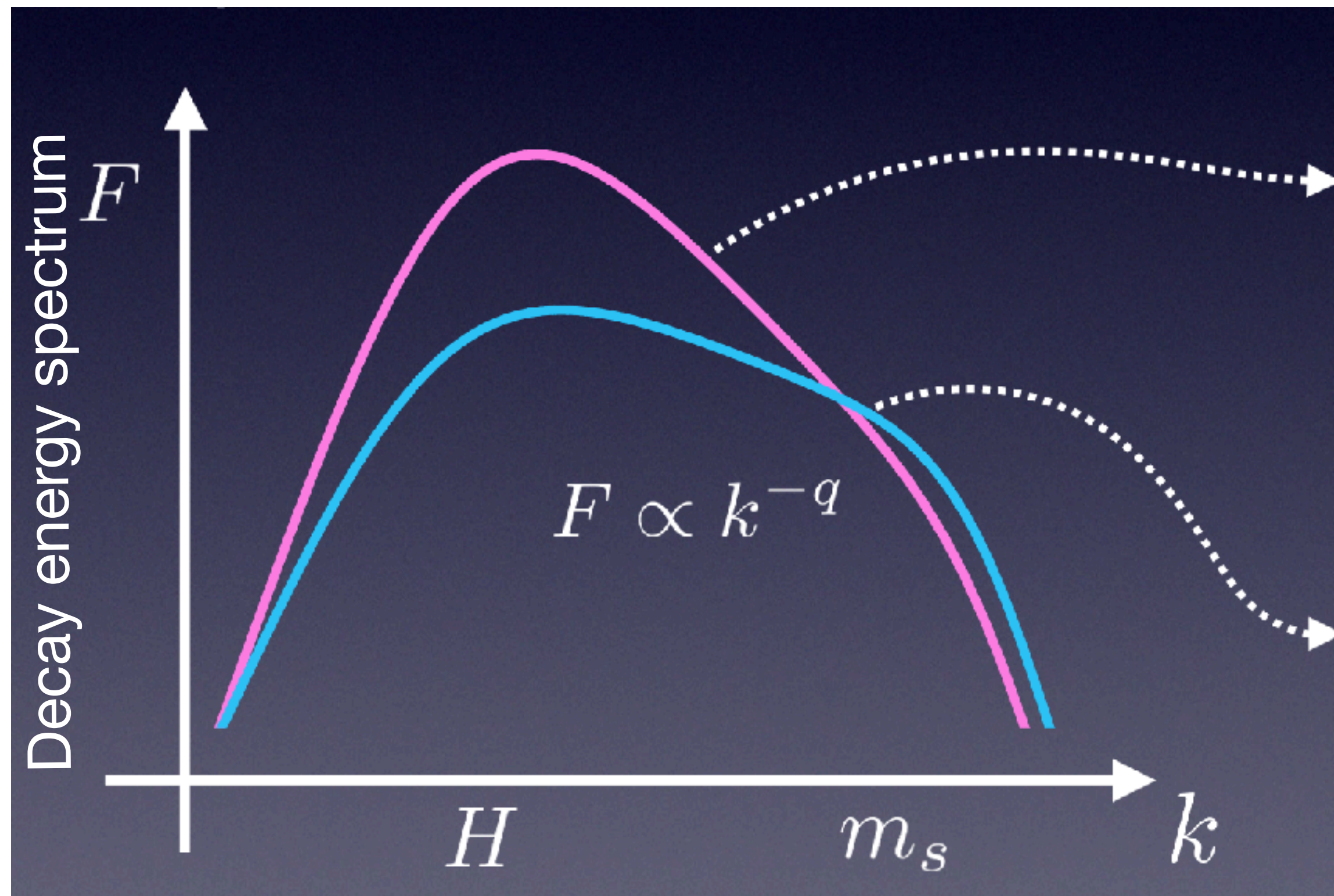
World-leading simulations: $\log \frac{m_r}{H} \simeq 9$

Physical: $\log \frac{m_r}{H} \simeq 70$

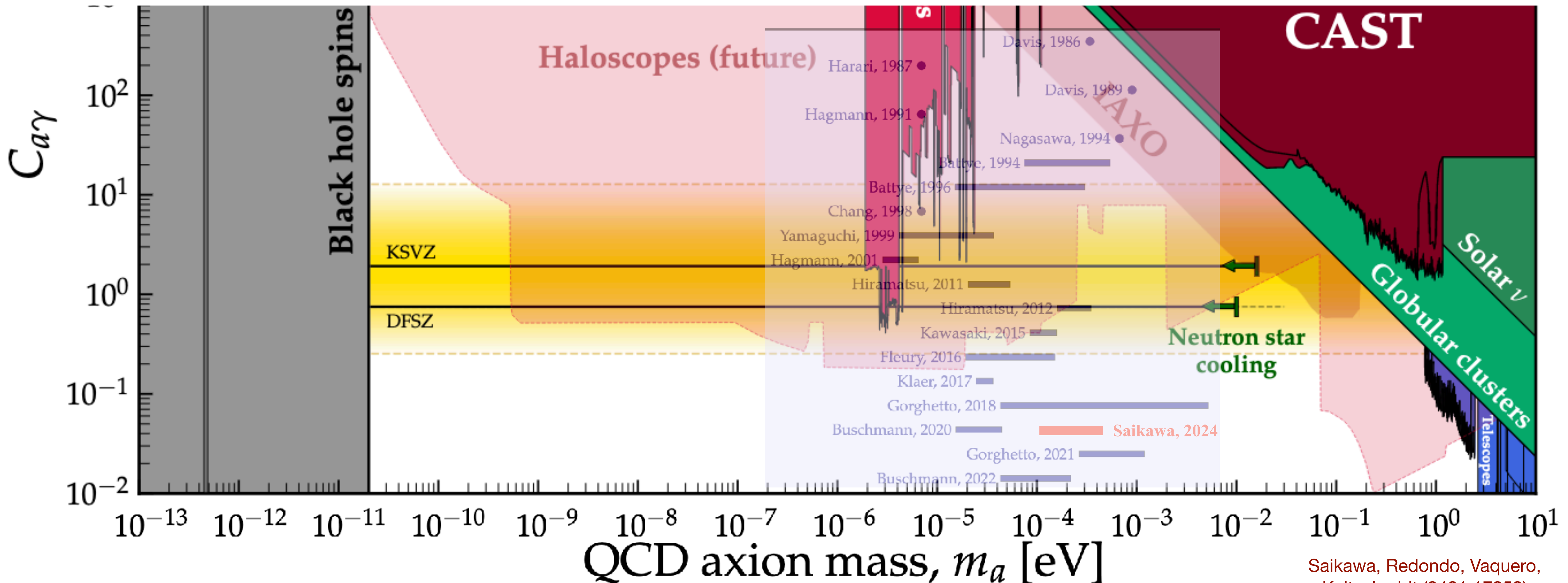
Axion String Decay

Post-inflationary scenario: $\Omega_a = \Omega_a^{\text{misalignment}}(f_a) + \Omega_a^{\text{strings}}(f_a)$

Credit: Ken'ichi Saikawa



The “Predictive” Post-Inflationary QCD Axion

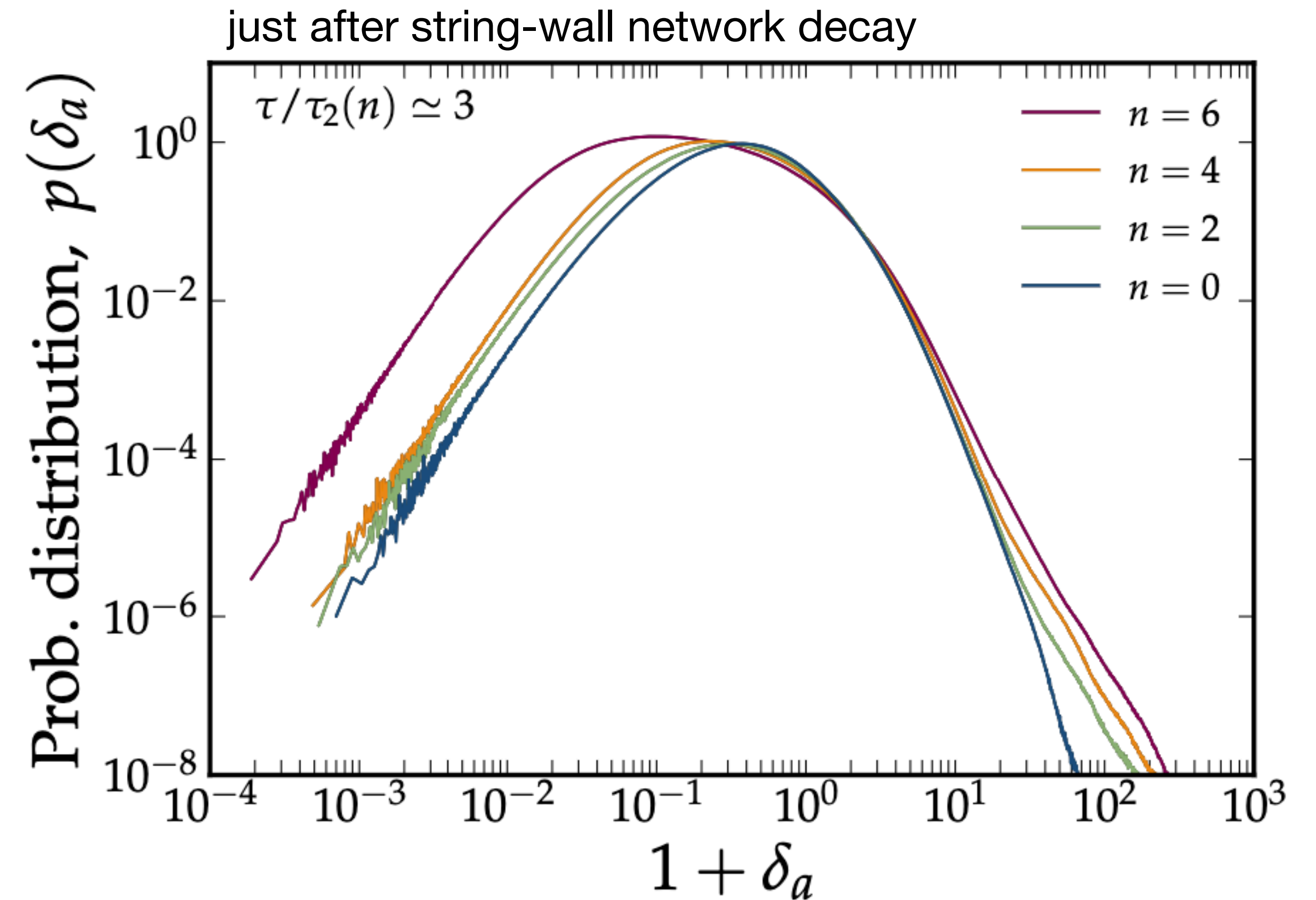
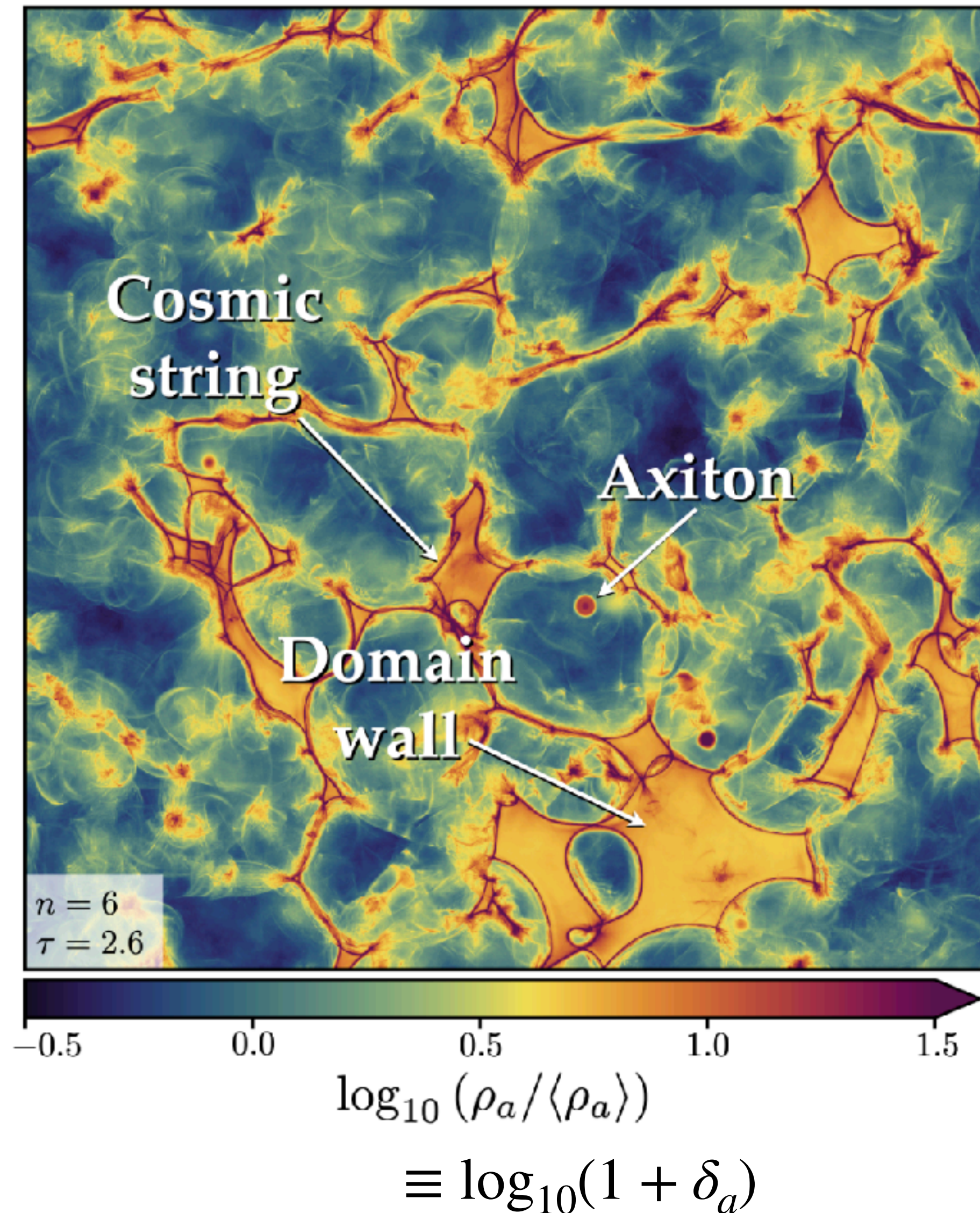


Saikawa, Redondo, Vaquero,
Kaltschmidt (2401.17253)

Rest of this talk: more general because we take m_ϕ, f_a as free parameters
(need not assume QCD axion)

ALPs and Temperature Dependence

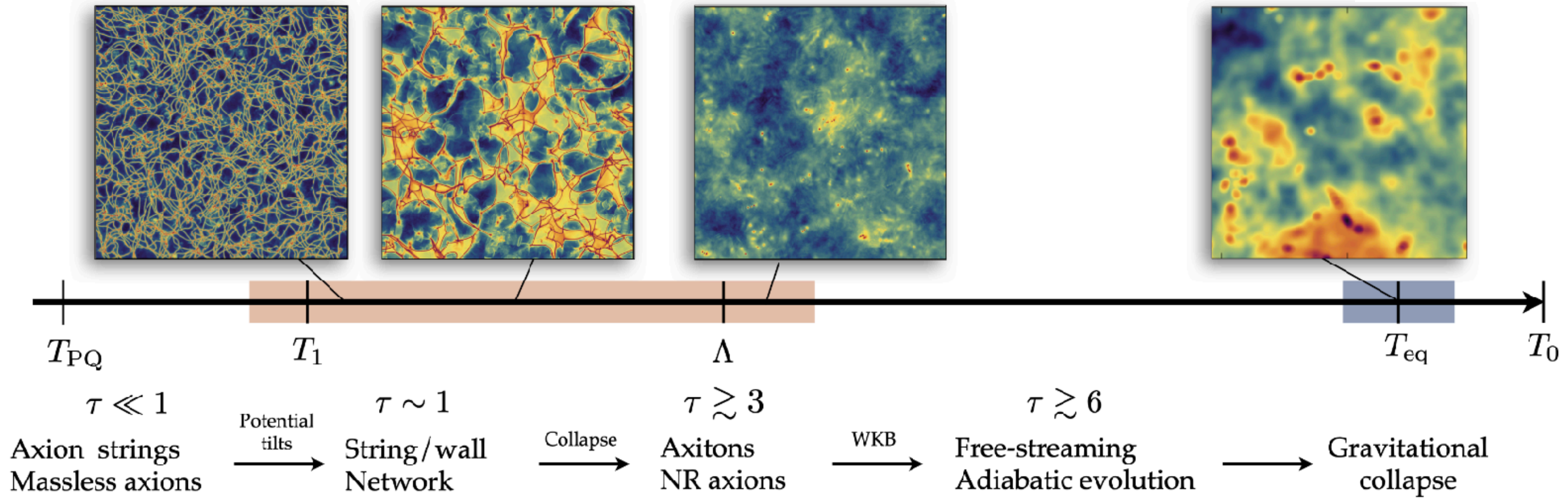
O'Hare, Pierson, Redondo,
Wong (2112.05117)



$$m_\phi(T)^2 \simeq m_\phi^2 \left(\frac{T_\star}{T} \right)^n$$

Axion Miniclusters

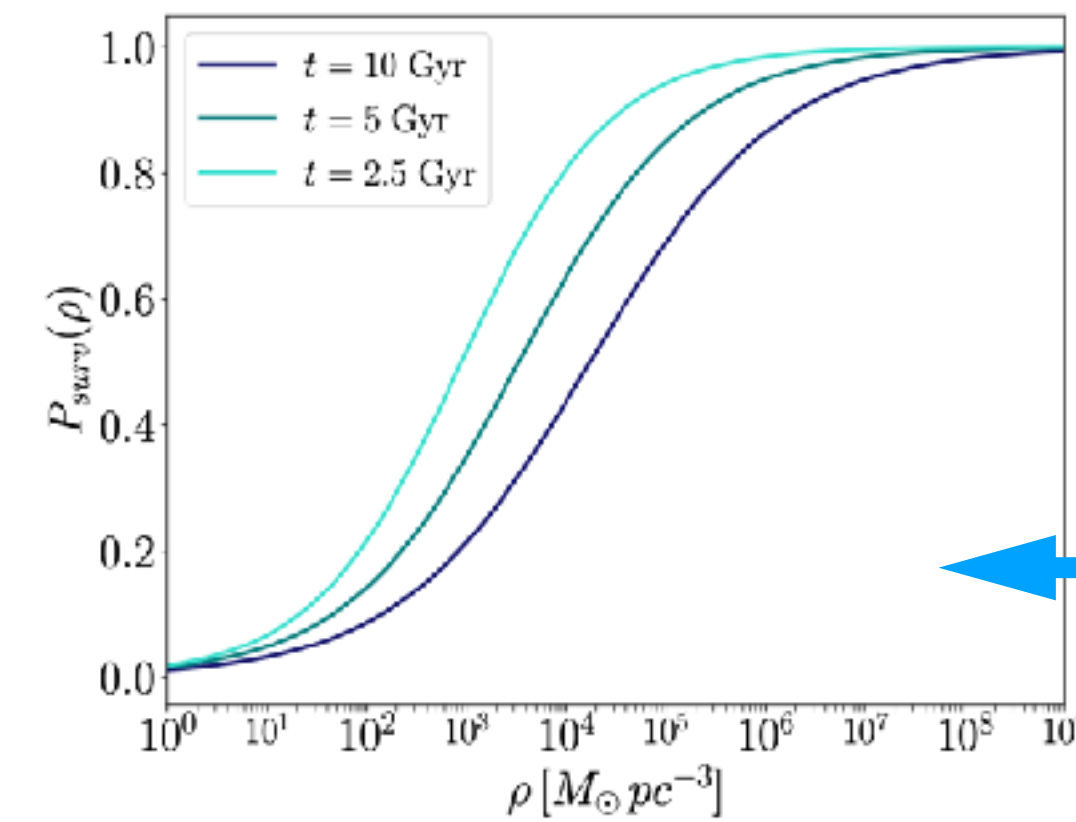
Hogan and Rees (PLB 1988)
 Kolb and Tkachev (hep-ph/9303313)
 ...
 O'Hare, Pierson, Redondo,
 Wong (2112.05117)



'Typical' example with $m(T) = m$, i.e. $n = 0$

$$M_{mc} \sim (1 + \delta_a) M_0 \sim 10^{-10} M_\odot (1 + \delta_a) \left(\frac{f_a}{10^{14} \text{ GeV}} \right)^2 \left(\frac{m_\phi}{10^{-10} \text{ eV}} \right)^2 \left(\frac{\text{GeV}}{T_{osc}} \right)^6$$

$$R_{mc} \sim \frac{L_1}{z_{eq} \delta_a} \sim \frac{10 \text{ au}}{\delta_a} \left(\frac{10^{-10} \text{ eV}}{m_\phi} \right)^{1/2} \quad \rho_{mc} \sim \rho_{eq} \delta_a^3 (1 + \delta_a)$$



Survive tidal disruption?

Dokuchaev, Eroshenko, Tkachev (1710.09586)

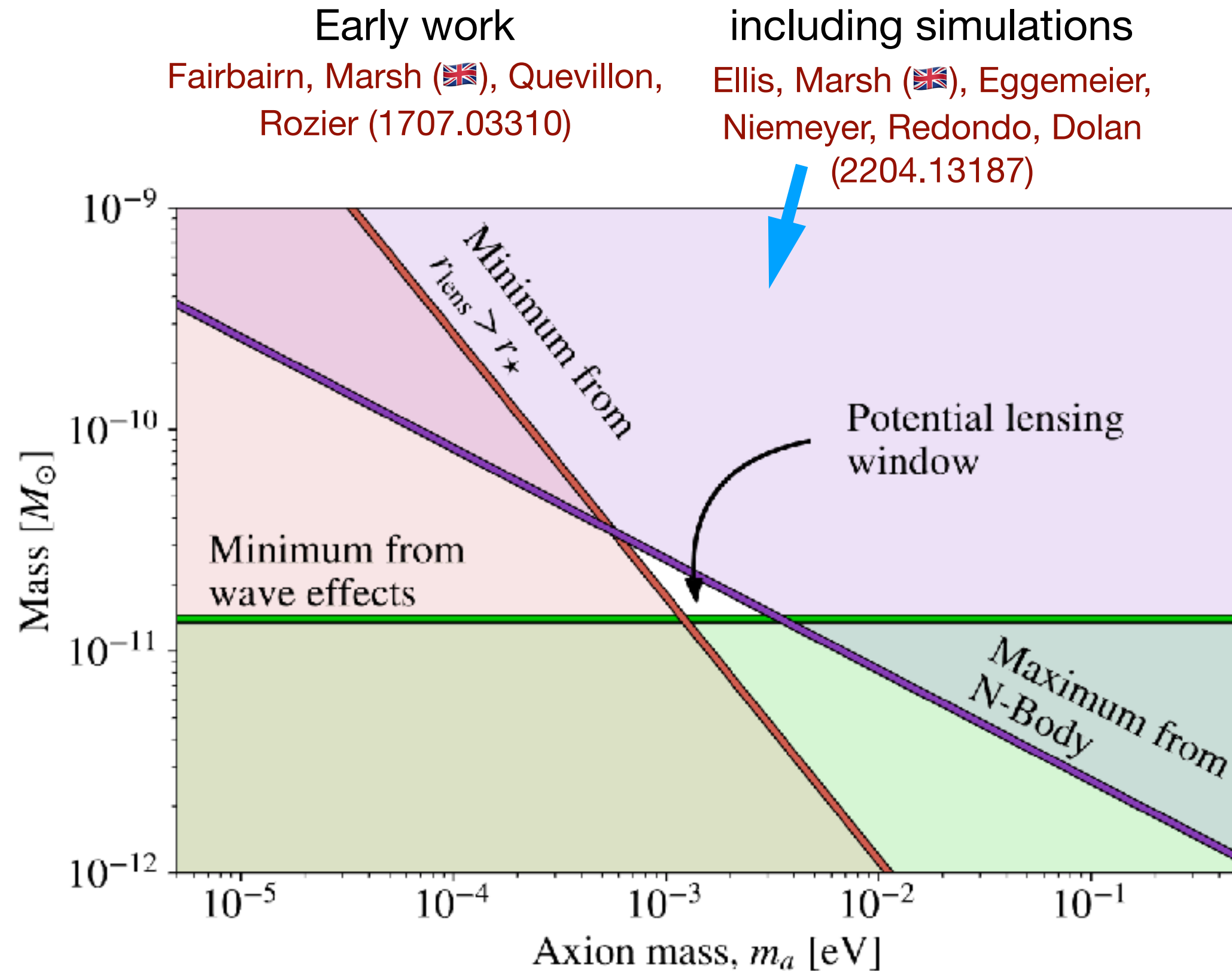
Kavanaugh, Edwards, Visinelli, Weniger (2011.05377)

Dandy, Schwetz, Todarello (2206.04619)

Shen, Xiao, Hopkins, Zurek (2207.11276)

Axion Miniclusters → New Searches 😊

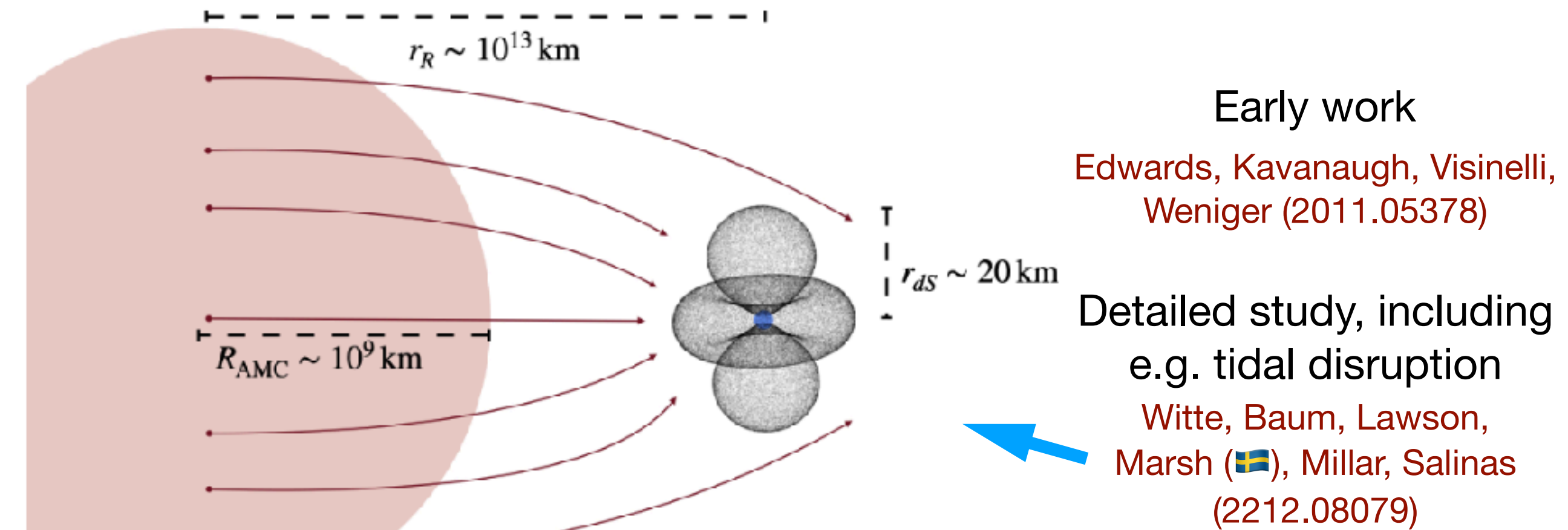
Gravitational microlensing



mc lensing modifies cluster lensing
for highly-magnified stars

Dai and Miralda-Escudé (1908.01773)

Radio signals from neutron star encounters



Gravitational waves from axion miniclusters?

Sun, Zhang (2003.10527)

See also Urrutia (<https://pos.sissa.it/454/046/pdf>)

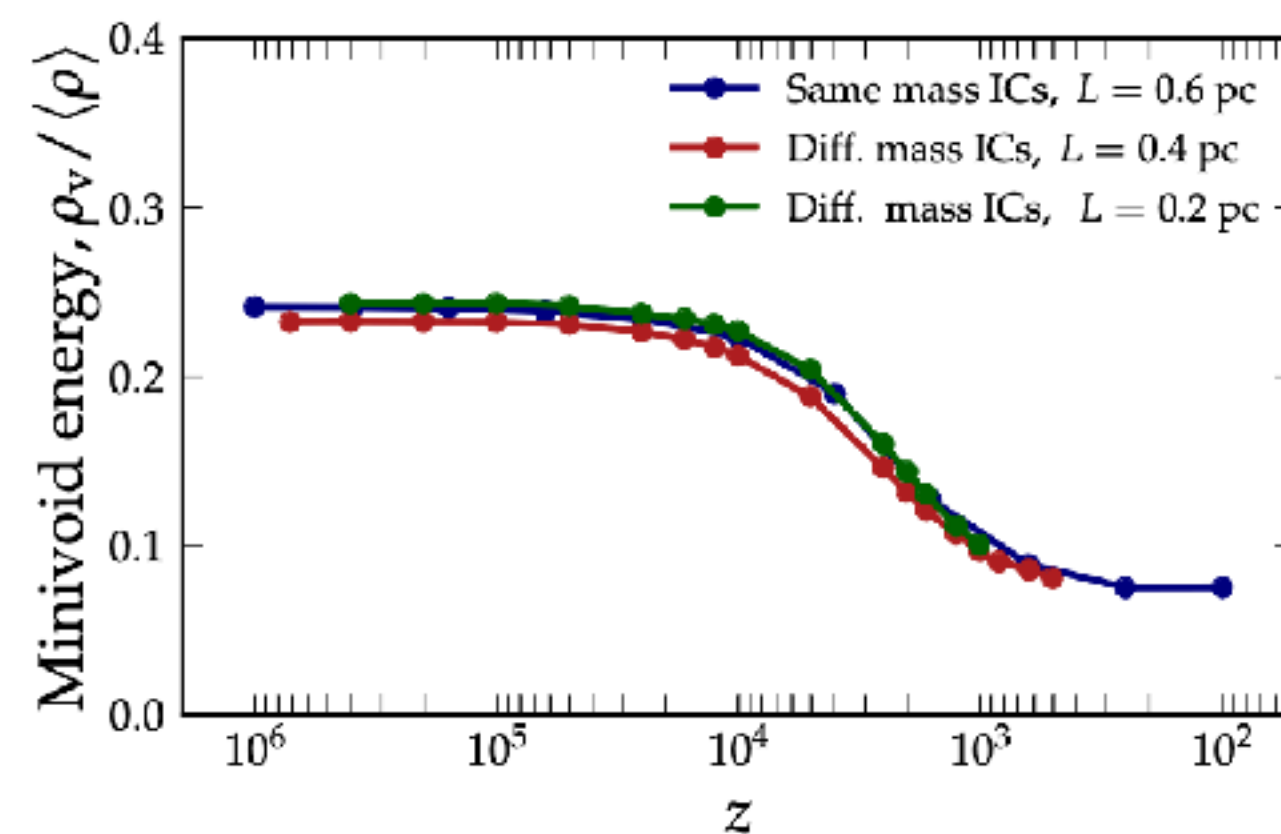
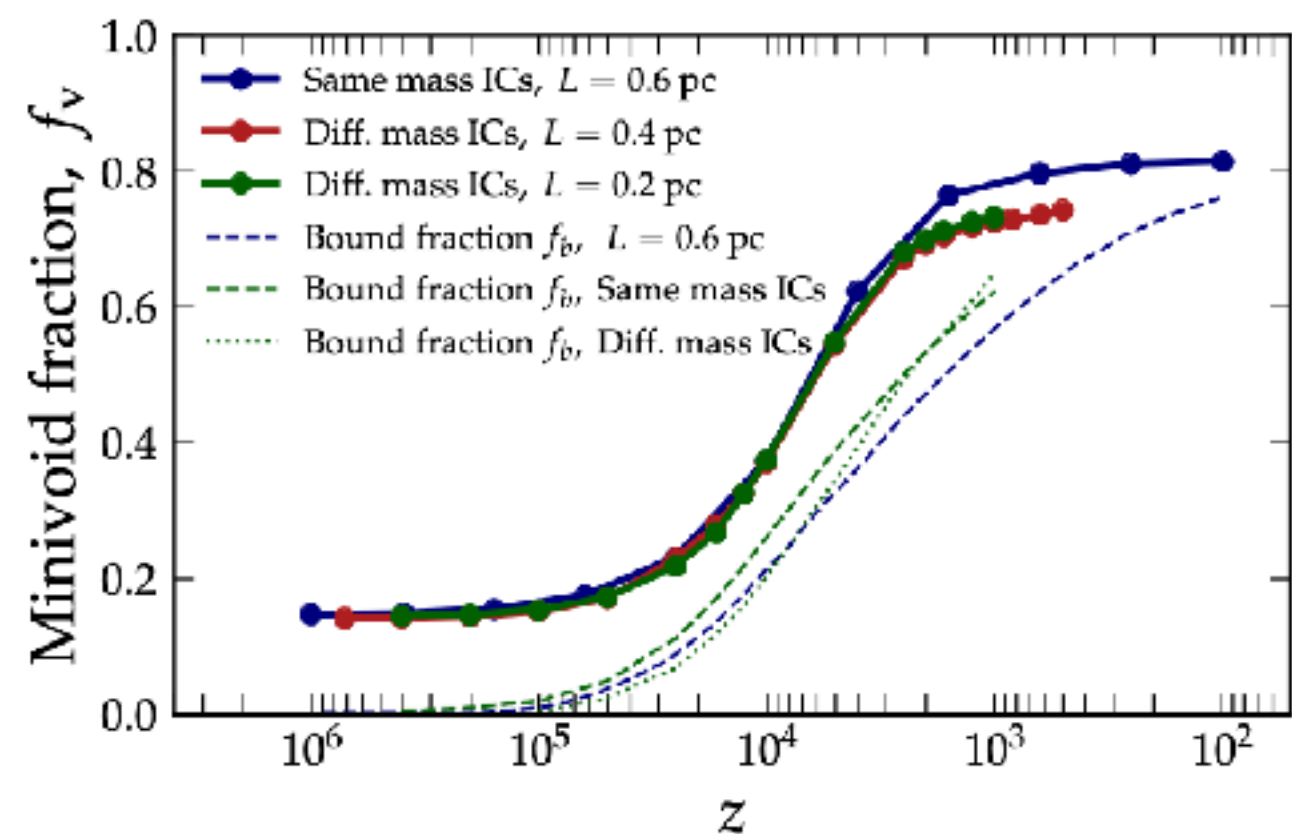
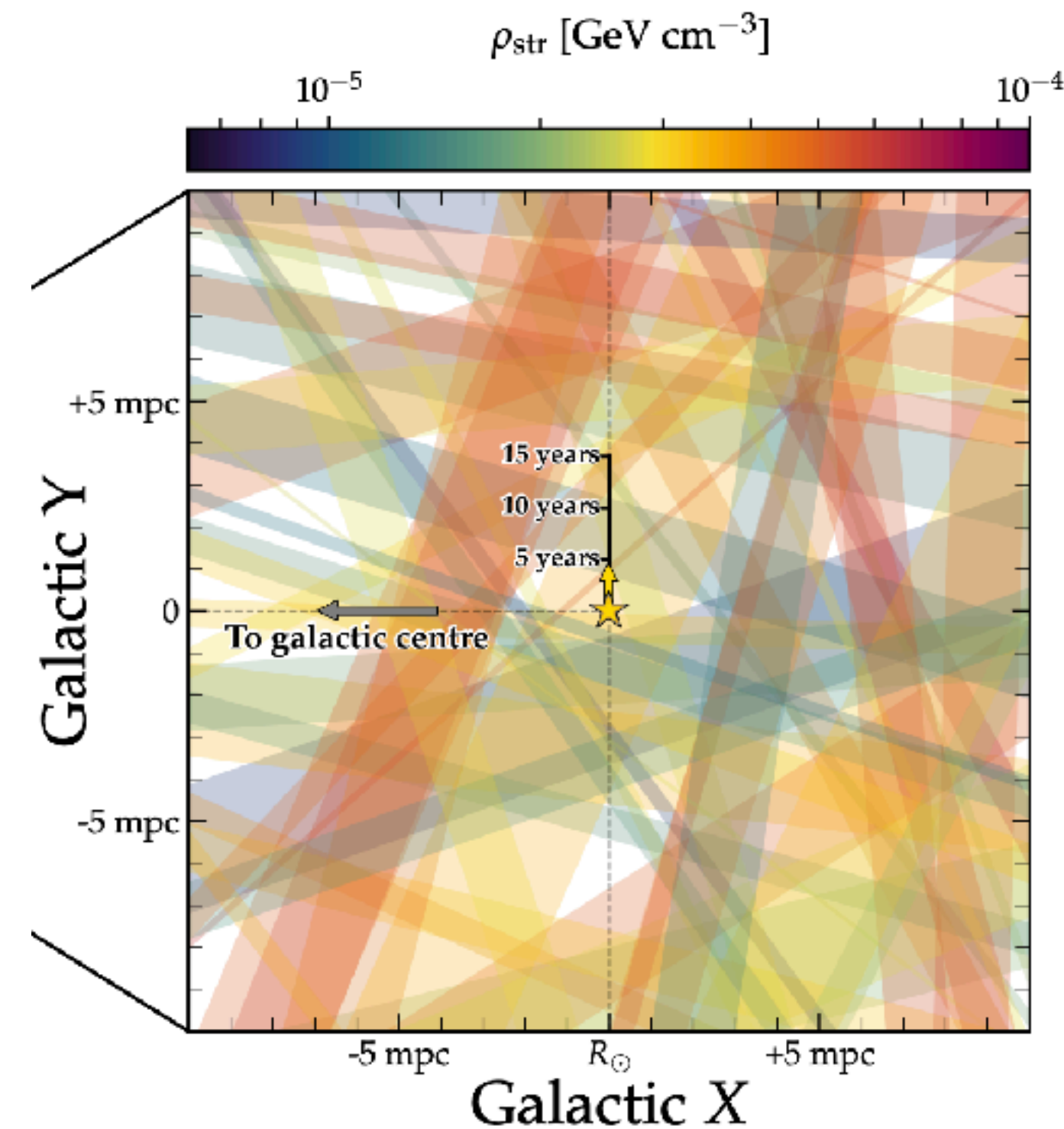
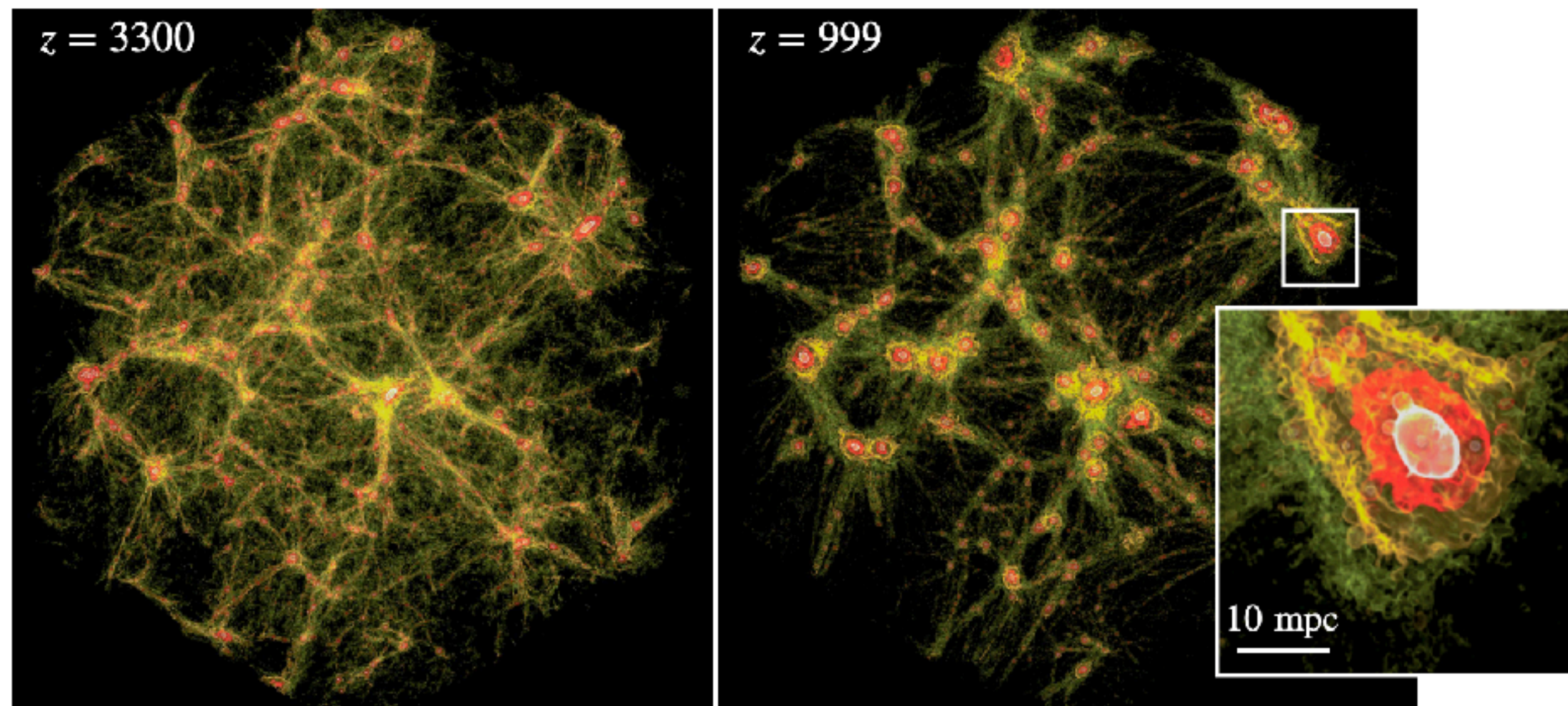
Explain surprisingly luminous early galaxies in JWST?

Hütsi, Raidal, Urrutia, Vaskonen, Veermäe (2211.02651)

Axion Miniclusters → Voids 🙄

Eggemeier, O'Hare, Pierobon, Redondo, Wong (2212.00560)

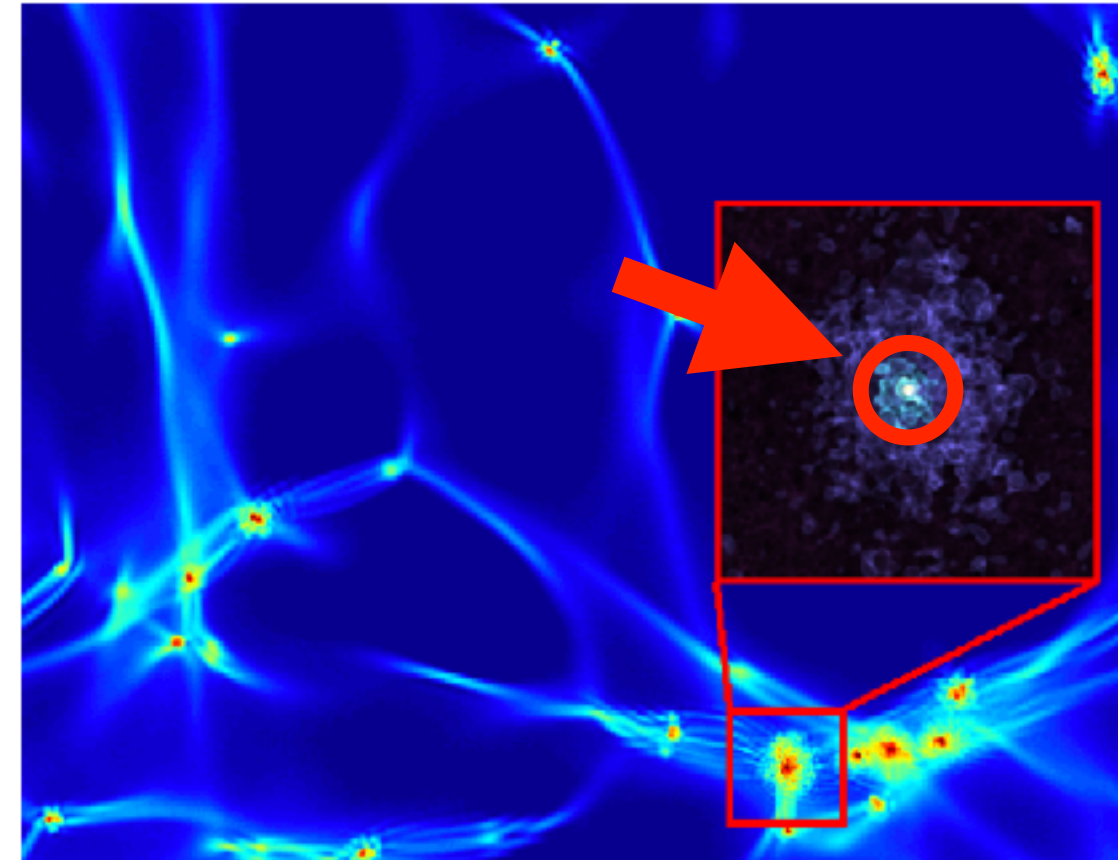
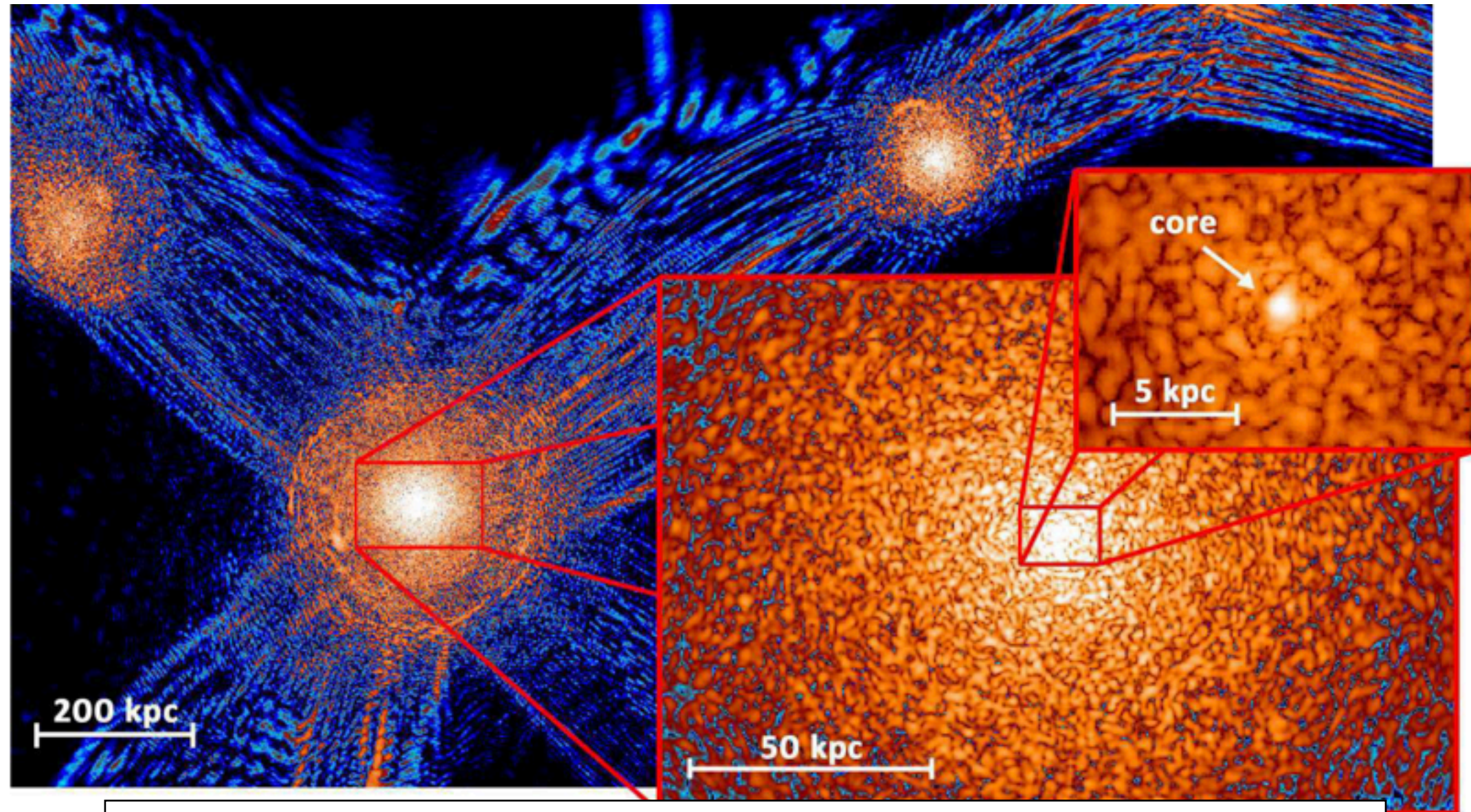
O'Hare, Pierobon, Redondo (2311.17367)



Minicluster tidal disruption

- numerous axion DM streams
- some recovery of local density + nontrivial velocity distribution

Relaxation



Mocz et al. (1705.05845)

Levkov, Panin, Tkachev (1804.05857)
Video via Alexander Panin on YouTube

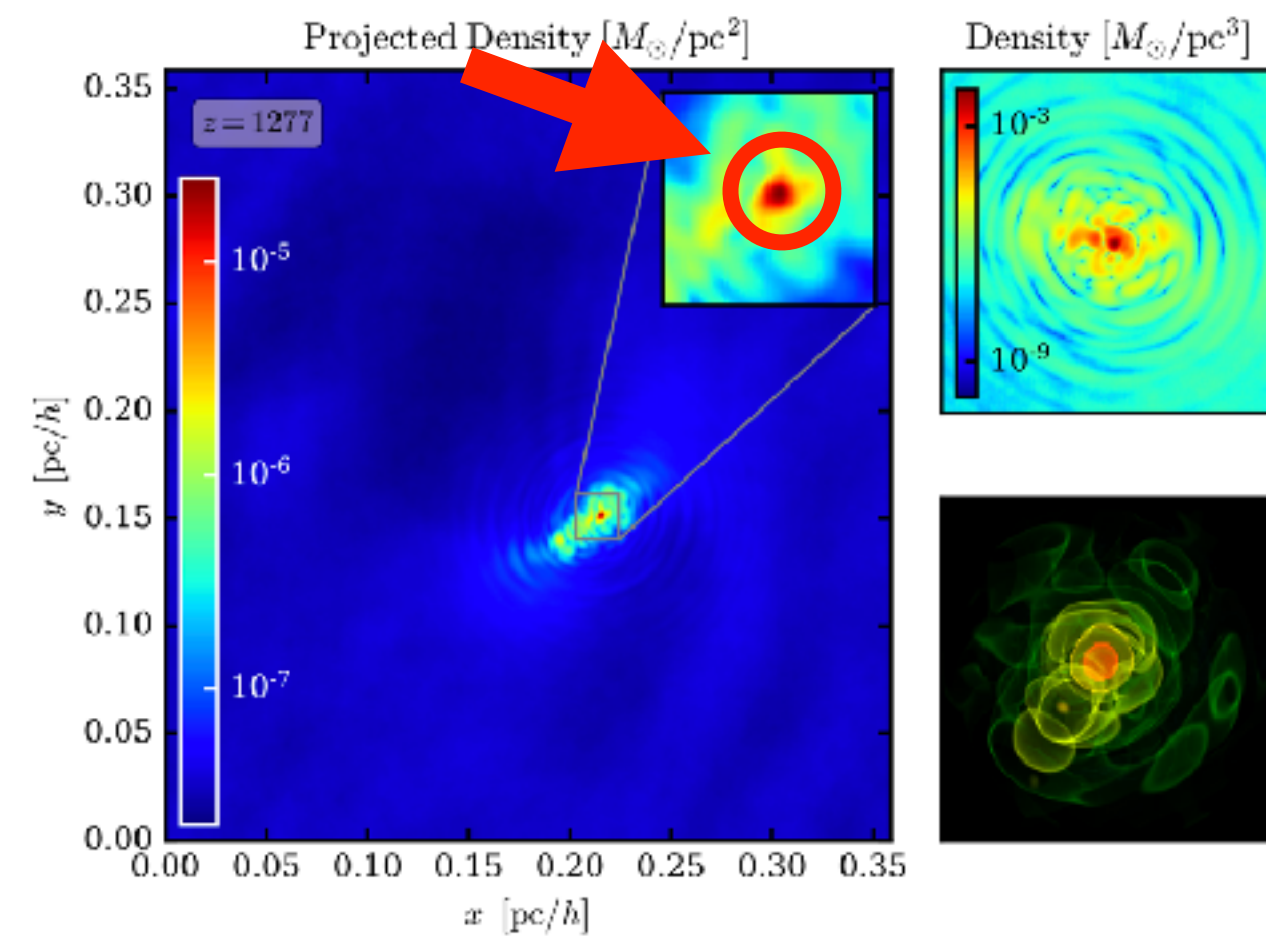
Letter | [Published: 22 June 2014](#)

Cosmic structure as the quantum interference of a coherent dark wave

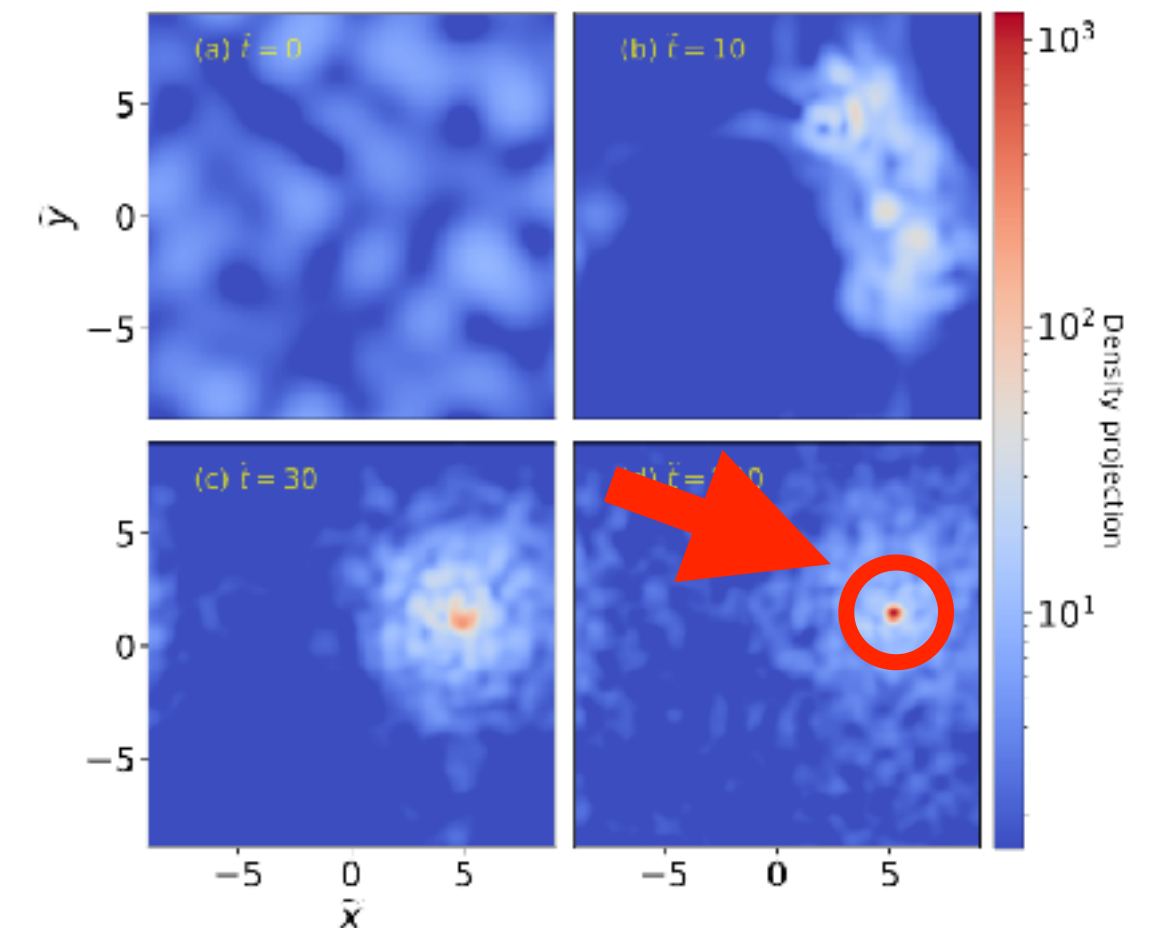
[Hsi-Yu Schive](#), [Tzihong Chiueh](#) & [Tom Broadhurst](#)

Nature Physics **10**, 496–499 (2014) | [Cite this article](#)

35k Accesses | 543 Citations | 145 Altmetric | [Metrics](#)



Eggemeier and Niemeyer (1906.01348)



Chen et al. (2011.01333)

(among others!)

Relaxation Timescale*

Sikivie and Yang (0901.1106)
 Levkov, Panin, Tkachev (1804.05857)
 Kirkpatrick, Mirasola, Prescod-Weinstein (2007.07438)
 Chen, Du, Lentz, Marsh (2109.11474)

$$\tau_{\text{relax}} \sim \left(n_\phi \sigma v_{\text{dm}} \mathcal{F}_\phi \right)^{-1} \sim \frac{1}{n_\phi \sigma v_{\text{dm}}} \frac{1}{n_\phi \lambda_{\text{dB}}^3}$$

$\sigma_g \simeq \frac{8\pi G^2 m_\phi^2}{v_{\text{dm}}^4} \ln \left(m_\phi v_{\text{dm}} R_{\text{halo}} \right)$

$\tau_{\text{relax}}^g \simeq \frac{\sqrt{2}}{12\pi^2} \frac{m_\phi^3 v_{\text{dm}}^6}{\rho^2 G^2}$

$\sim 10 \text{ Gyr} \left(\frac{m_\phi}{10^{-21} \text{ eV}} \right)^3 \left(\frac{10 \rho_{\text{dm}}}{\rho} \right)^2 \left(\frac{v_{\text{dm}}}{10^{-4}} \right)^6$

$\sigma_\lambda \simeq \frac{\lambda^2}{2\pi m_\phi^2} \simeq \frac{m_\phi^2}{2\pi f_a^4}$

$\tau_{\text{relax}}^\lambda \simeq \frac{64 m_\phi^3 f_a^4 v_{\text{dm}}^2}{\rho^2}$

$\sim 10 \text{ Gyr} \left(\frac{m_\phi}{10^{-14} \text{ eV}} \right)^3 \left(\frac{f_a}{10^8 \text{ GeV}} \right)^4 \left(\frac{\rho_{\text{dm}}}{\rho} \right)^2 \left(\frac{v_{\text{dm}}}{10^{-3}} \right)^2$

$$\text{Ratio: } r_{\lambda g} \equiv \frac{\tau_{\text{relax}}^\lambda}{\tau_{\text{relax}}^g} \simeq G^2 f_a^4 v_{\text{dm}}^4 = \left(\frac{f_a v_{\text{dm}}}{M_{\text{Pl}}} \right)^4$$

Note also a possible cross-term $\propto G\lambda$,
 highly relevant when $r_{\lambda g} \sim \mathcal{O}(1)$ or λ repulsive

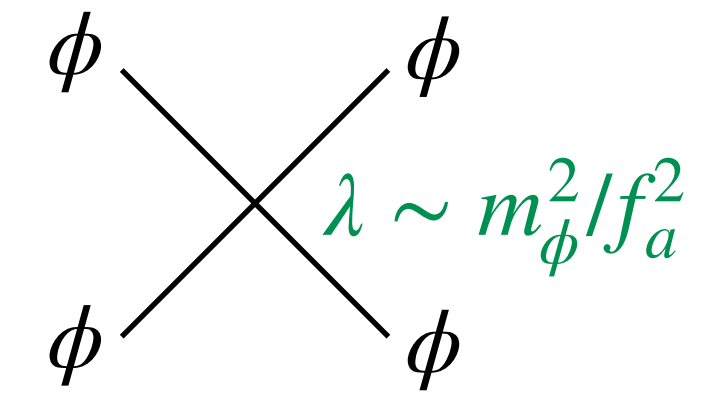
Jain, Wanichwecharungruang, Thomas (2310.00058)

*violent relaxation, e.g. during merger,
 is much faster (basically instantaneous)

ULDM Ground States

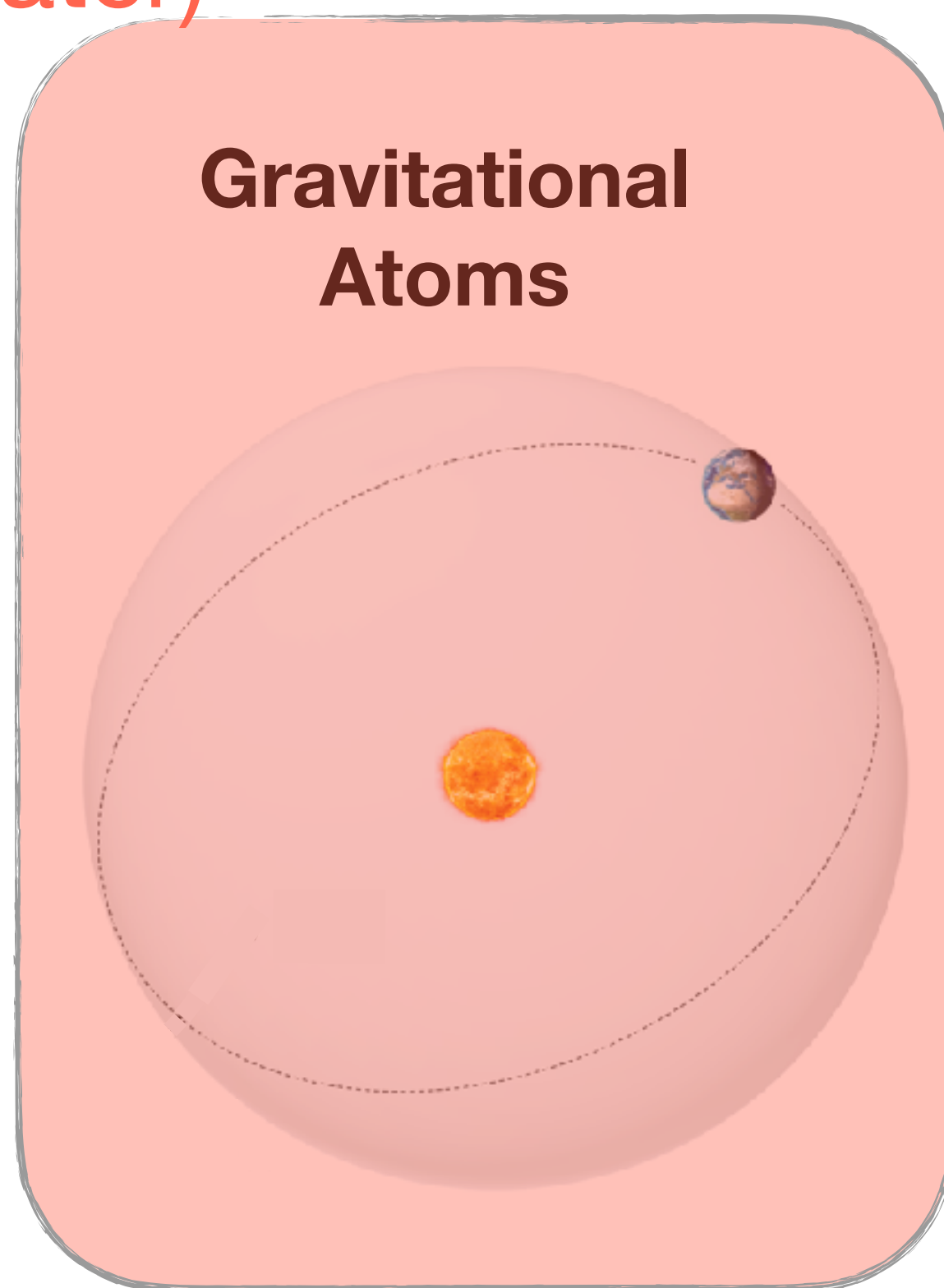
“Quantum” pressure (Repulsive) Gravity (Attractive) Self-interactions (usually attractive)

$$i\psi = \left[-\frac{\nabla^2}{2m_\phi} + V_g(\psi) \right] \psi - \frac{\lambda}{8m_\phi^2} |\psi|^2 \psi$$

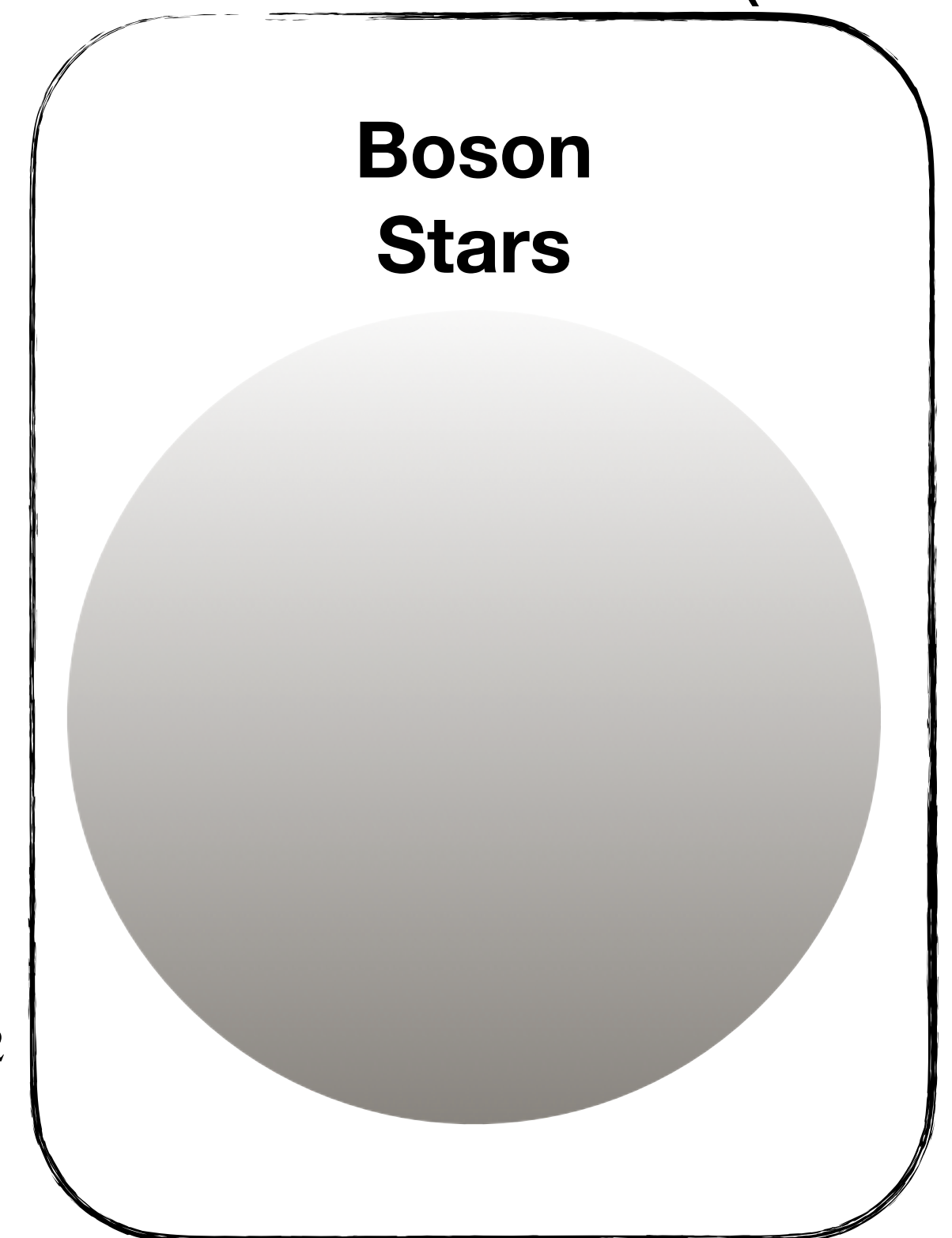


(small when density small, return to this later)

(later)

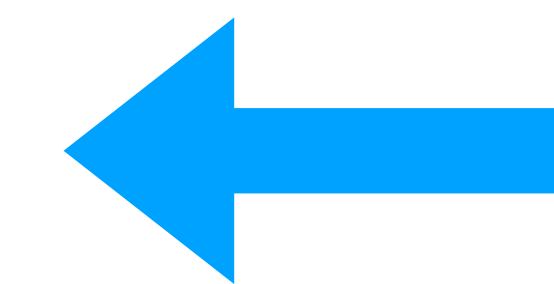


(next)



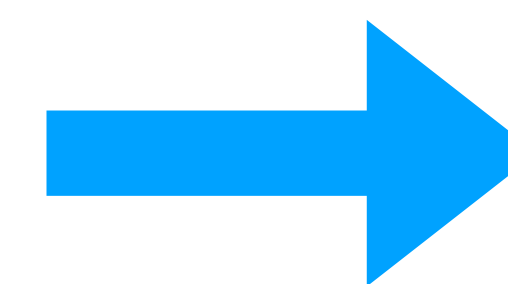
Balance of gradient+gravity in the field

Structure depends on **source of gravity**



External source (bound to other body)

$$V_g(\psi) = \frac{GMm_\phi}{r}$$



ULDM itself (self-gravity)

$$\nabla^2 V_g(\psi) = 4\pi Gm_\phi^2 |\psi|^2$$

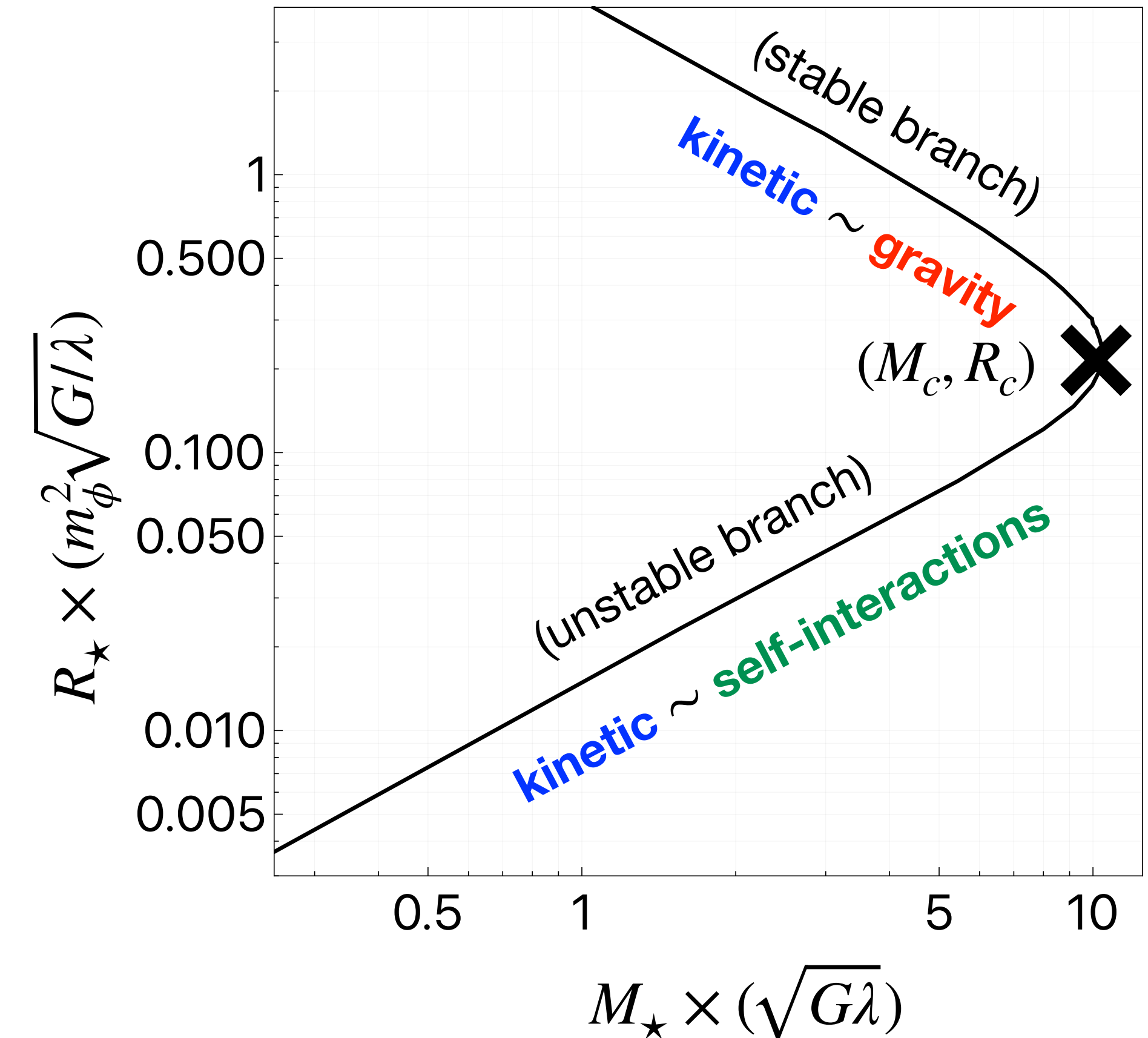
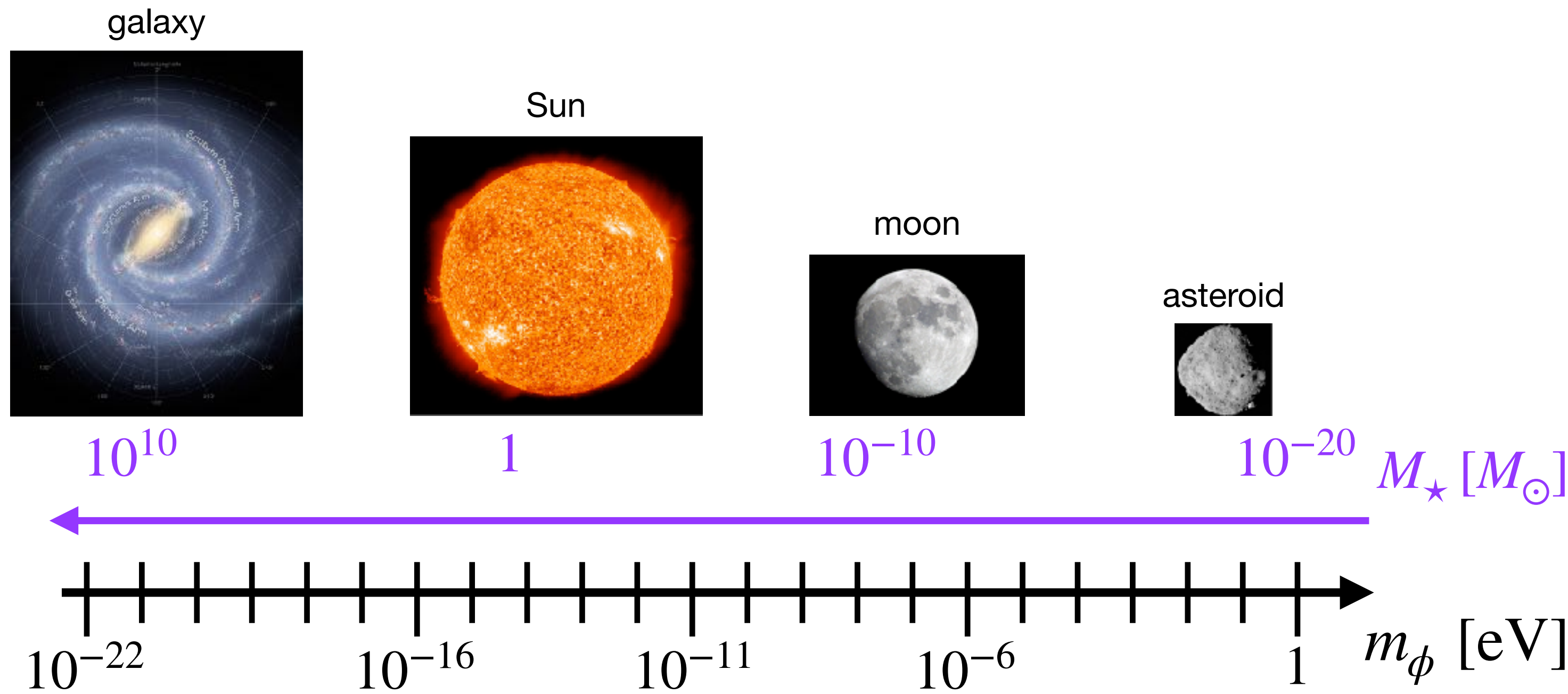
Chavanis (1103.2050)
Chavanis, Delfini (1103.2054)

Size of a Boson Star

The most important fact about a boson star!

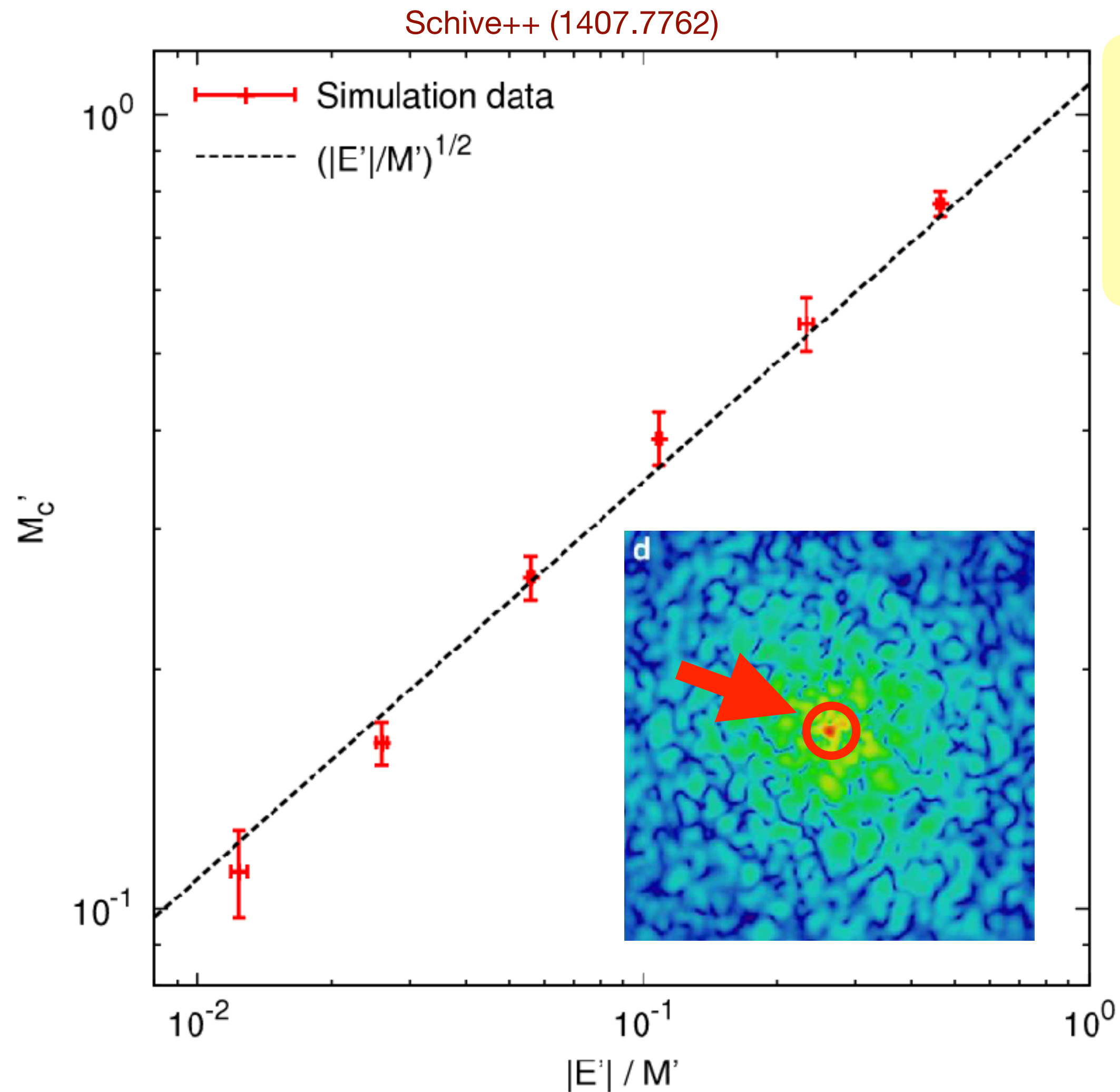
$$M_{\star} \simeq \frac{1}{Gm_{\phi}^2 R_{\star}}$$

$$i\psi = \left[-\frac{\nabla^2}{2m_{\phi}} + V_g(\psi) \right] \psi - \frac{\lambda}{8m_{\phi}^2} |\psi|^2 \psi$$



The Soliton—Host-Halo Relation

(or Core-Halo Relation)



Rule: 1 boson star per halo with

$$M_{ch} \simeq 10^9 M_\odot \left(\frac{10^{-22} \text{ eV}}{m_\phi} \right) \left(\frac{M_{\text{halo}}}{10^{12} M_\odot} \right)^{1/3}$$

Tested in simulations for

- halos with $M_{\text{halo}} \sim (10^8 - 10^{12}) M_\odot$
- ULDM mass $m_\phi \sim (10^{-20} - 10^{-22}) \text{ eV}$
- Other systems with small overdensities (e.g. QCD axion miniclusters)

Reasons to be (at least a little bit) skeptical:

- larger simulation volumes \longrightarrow scatter, $M_{ch} \propto M_{\text{halo}}^{2/5}$? $M_{\text{halo}}^{2/3}$?
- can't be valid for $M_{ch} \rightarrow M_{\text{halo}}$
- Valid when m_ϕ is large? at fixed M_{halo} , predicts very large overdensity

equivalent to

$$\left(\frac{E}{M} \right)_{\text{soliton}} = \left(\frac{E}{M} \right)_{\text{halo}}$$

and therefore

$$\left(\frac{K}{M} \right)_{\text{soliton}} = \left(\frac{K}{M} \right)_{\text{halo}}$$

Bar, Blas, Blum, Sibiryakov (1805.00122)

Bar, Blum, JE, Sato (1903.03402)

Chan, Ferreira, May, Hayashi, Chiba (2110.11882)

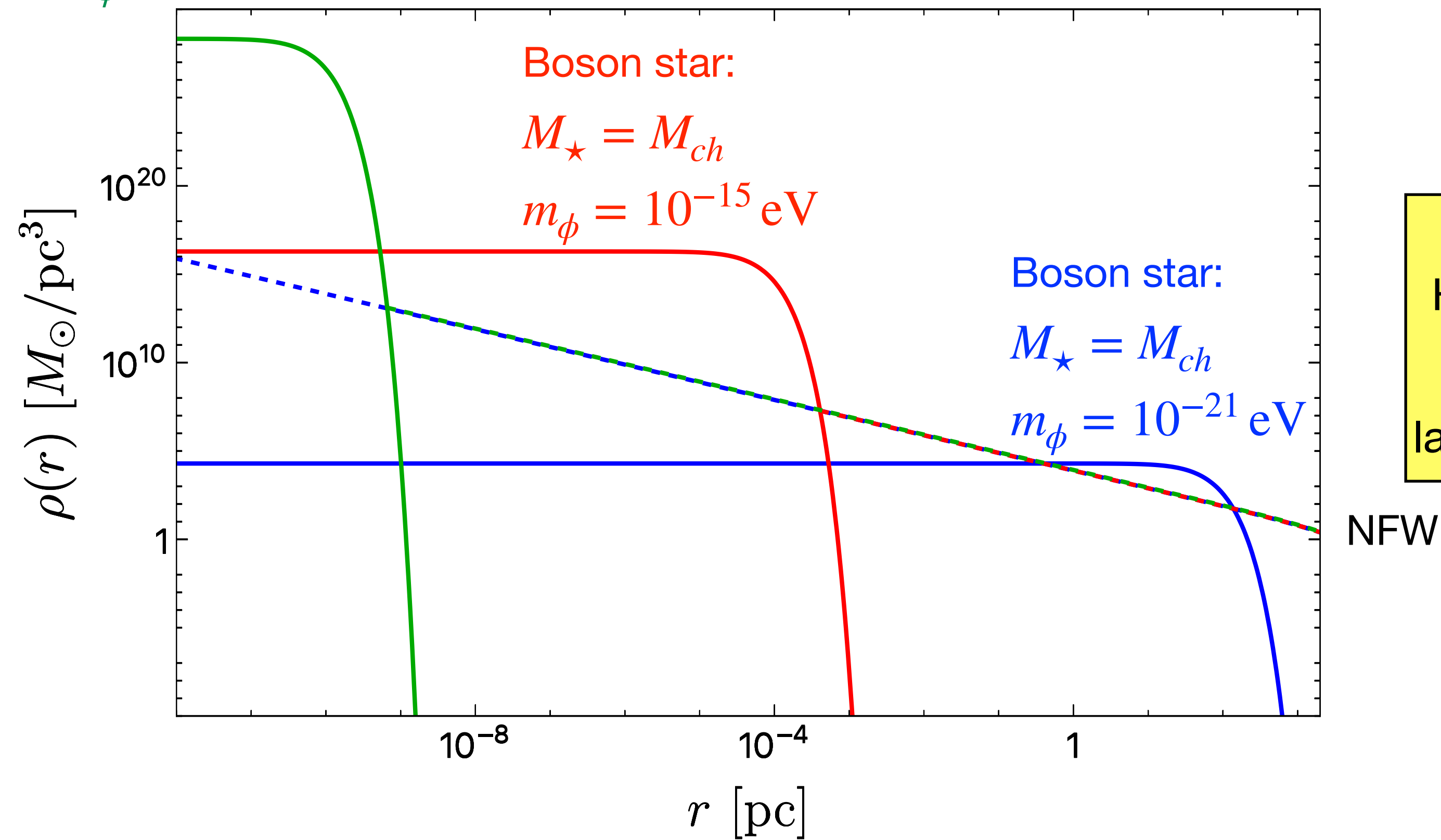
The Soliton—Host-Halo Relation

(or Core-Halo Relation)

Boson star:

$$M_{\star} = M_{ch}$$

$$m_{\phi} = 10^{-9} \text{ eV}$$

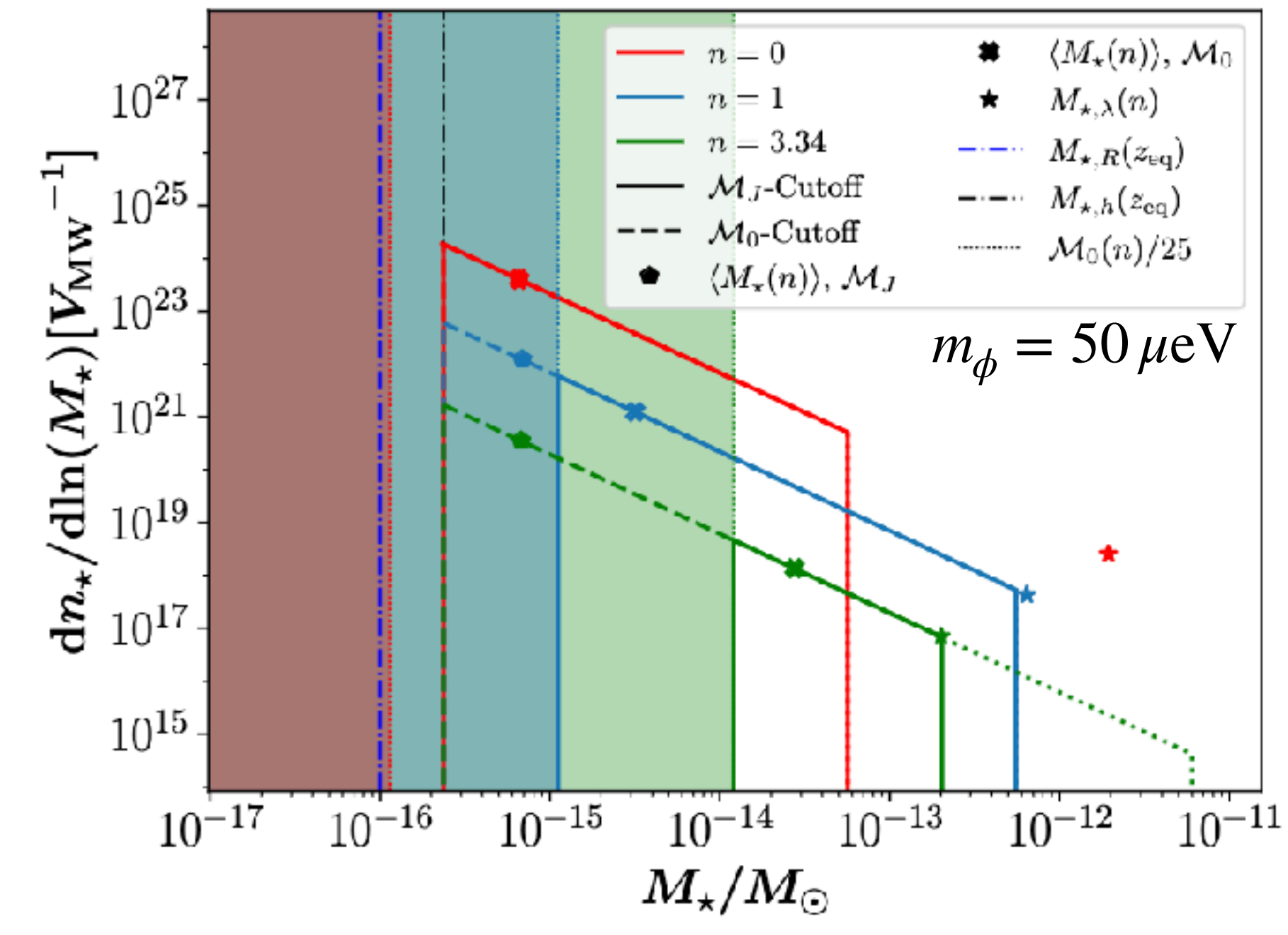
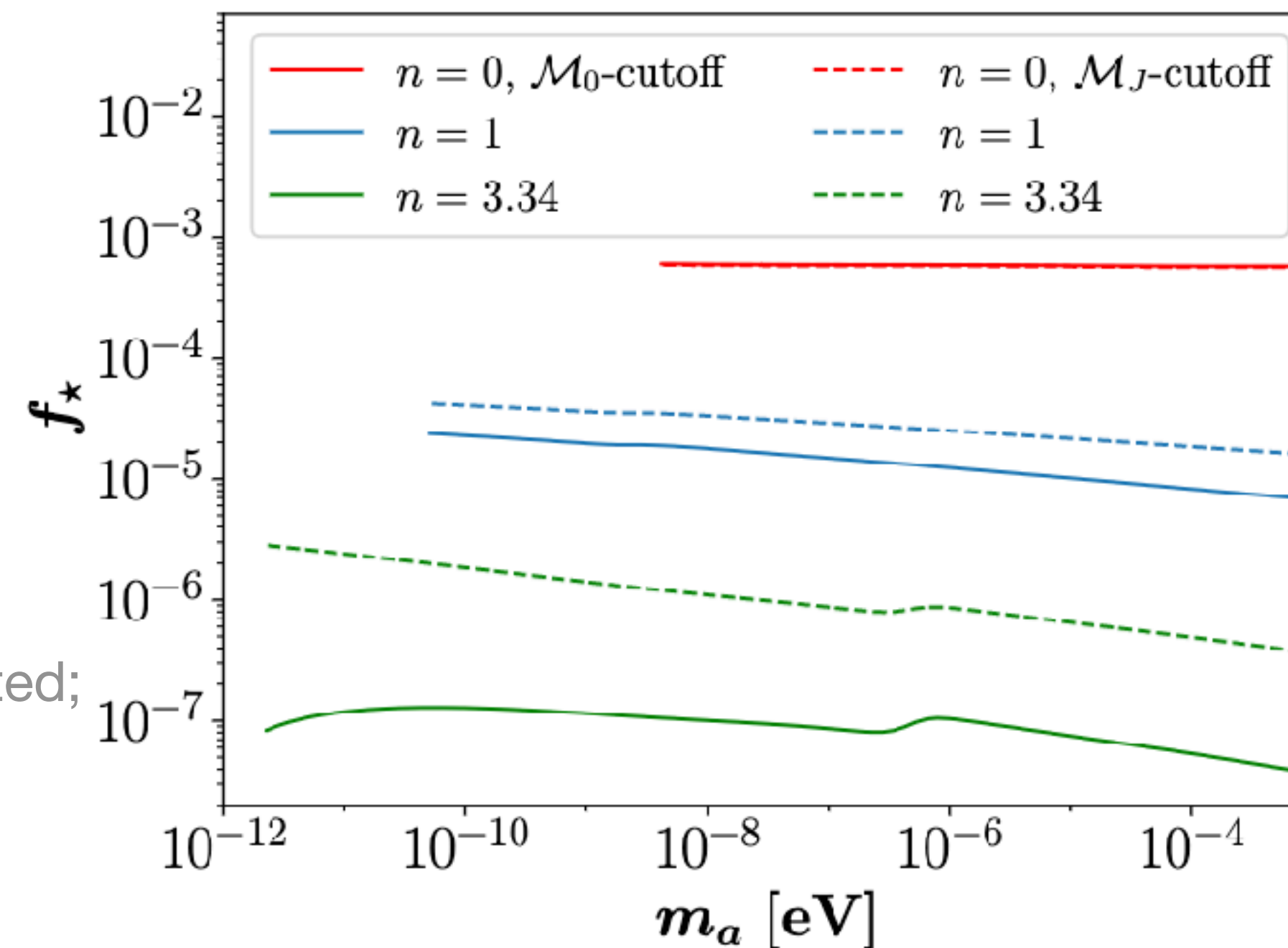
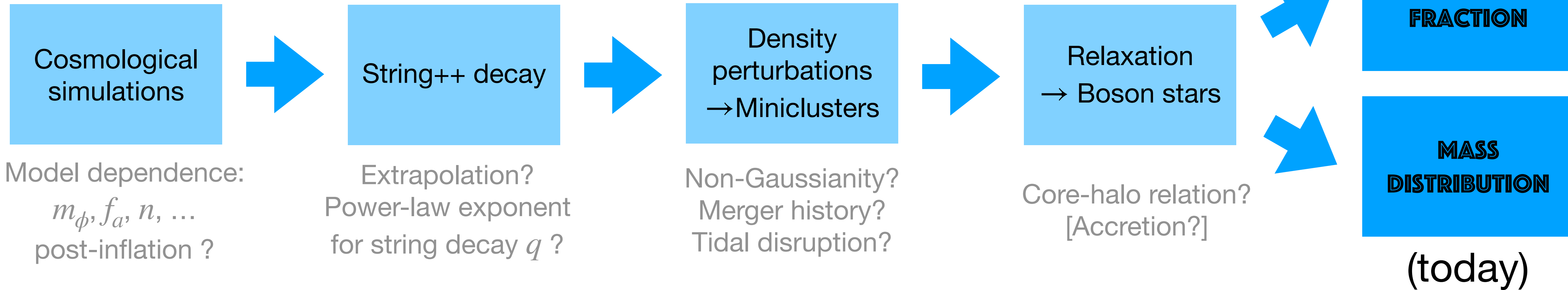


Clarity is still needed:
 How is core-halo relation
 modified in case of
 large M_{halo} with large m_{ϕ} ?

- Valid when m_{ϕ} is small? at fixed M_{halo} , predicts very large overdensity

Boson Star Mass Distribution* Today

Maseizik, Sigl
(2404.07908)

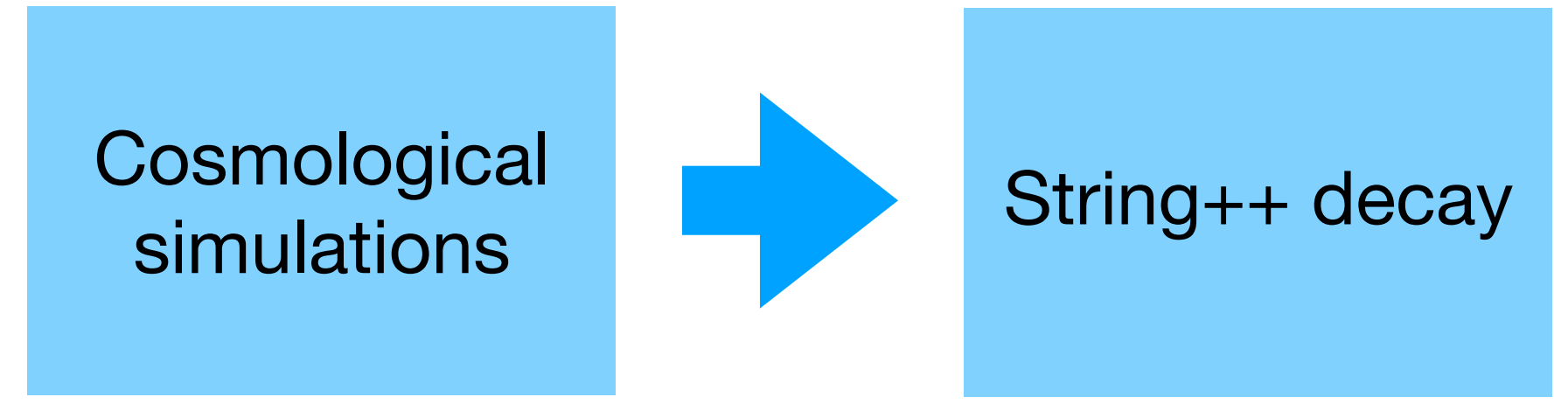


*accretion not accounted for, still contested;

- see
- Levkov, Panin, Tkachev (1804.05857)
 - Eggemeier, Niemeyer (1906.01348)
 - Chen, Du, Lentz, Marsh, Niemeyer (2011.01333)
 - Chan, Sibiryakov, Xue (2207.04057)
 - Dmitriev, Levkov, Panin, Tkachev (2305.01005)

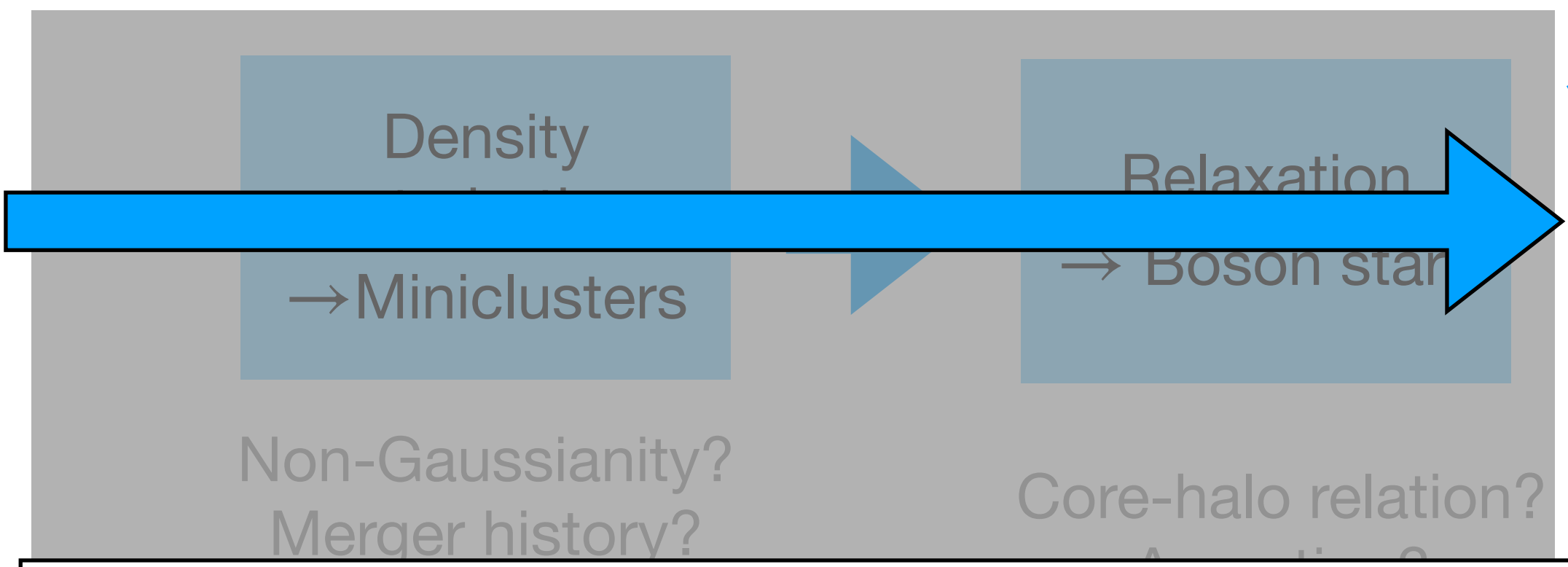
Gorghetto, Hardy, Villadoro (2405.19389)

Early Formation of QCD Axion Stars



Model dependence:
 m_ϕ, f_a, n, \dots
 post-inflation ?

Extrapolation?
 Power-law exponent
 for string decay q ?

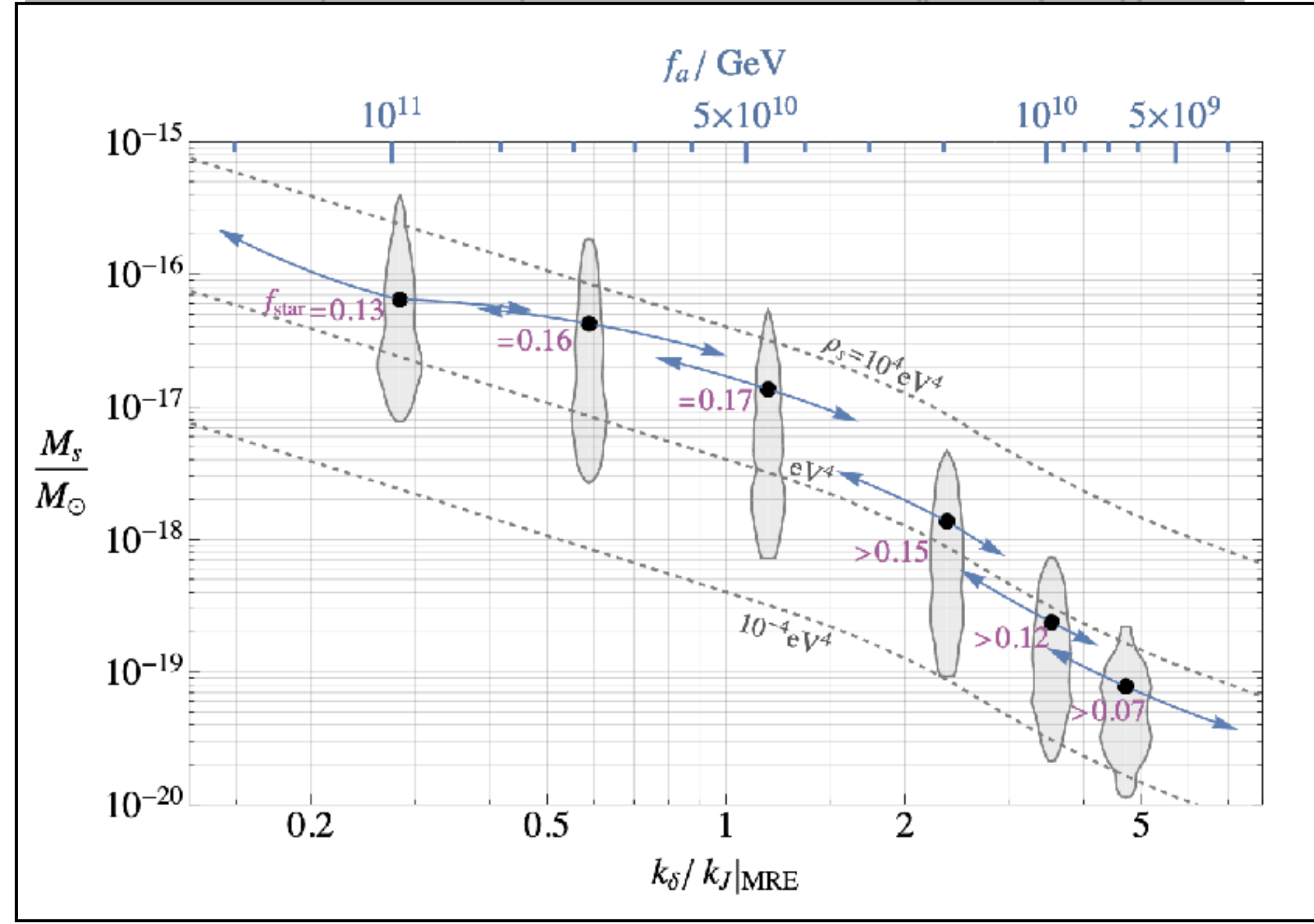


DM MASS FRACTION

MASS DISTRIBUTION

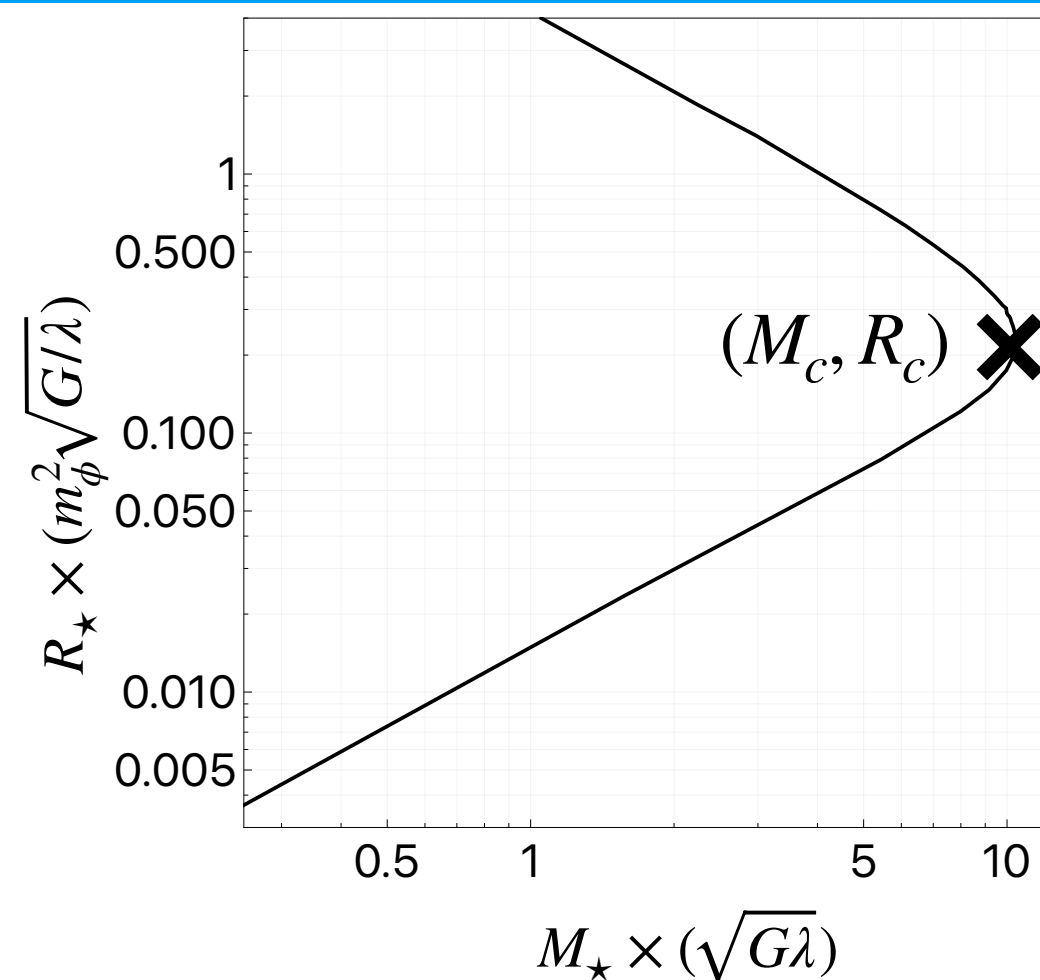
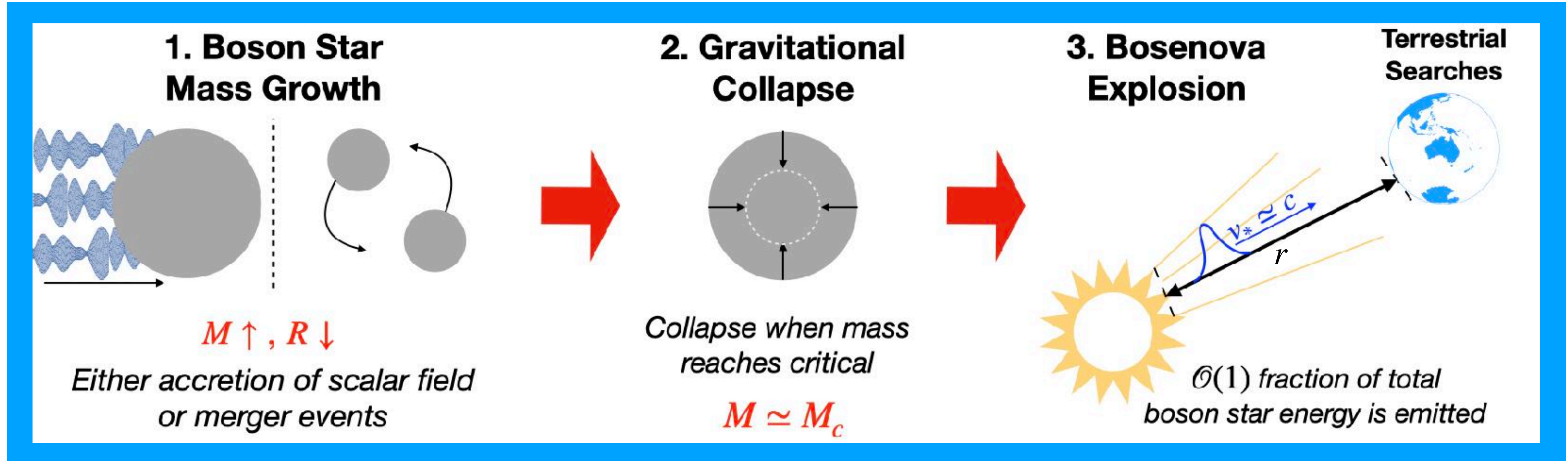
(at MRE)

In preferred QCD axion DM range, collapse at MRE directly leads to $f_{\text{dm}} \sim \text{few} - 20\%$ DM fraction in boson stars!



Boson Star Collapse \longrightarrow Bosenova

Image: Arakawa, **JE**, Safronova, Takhistov, Zaheer (2306.16468)



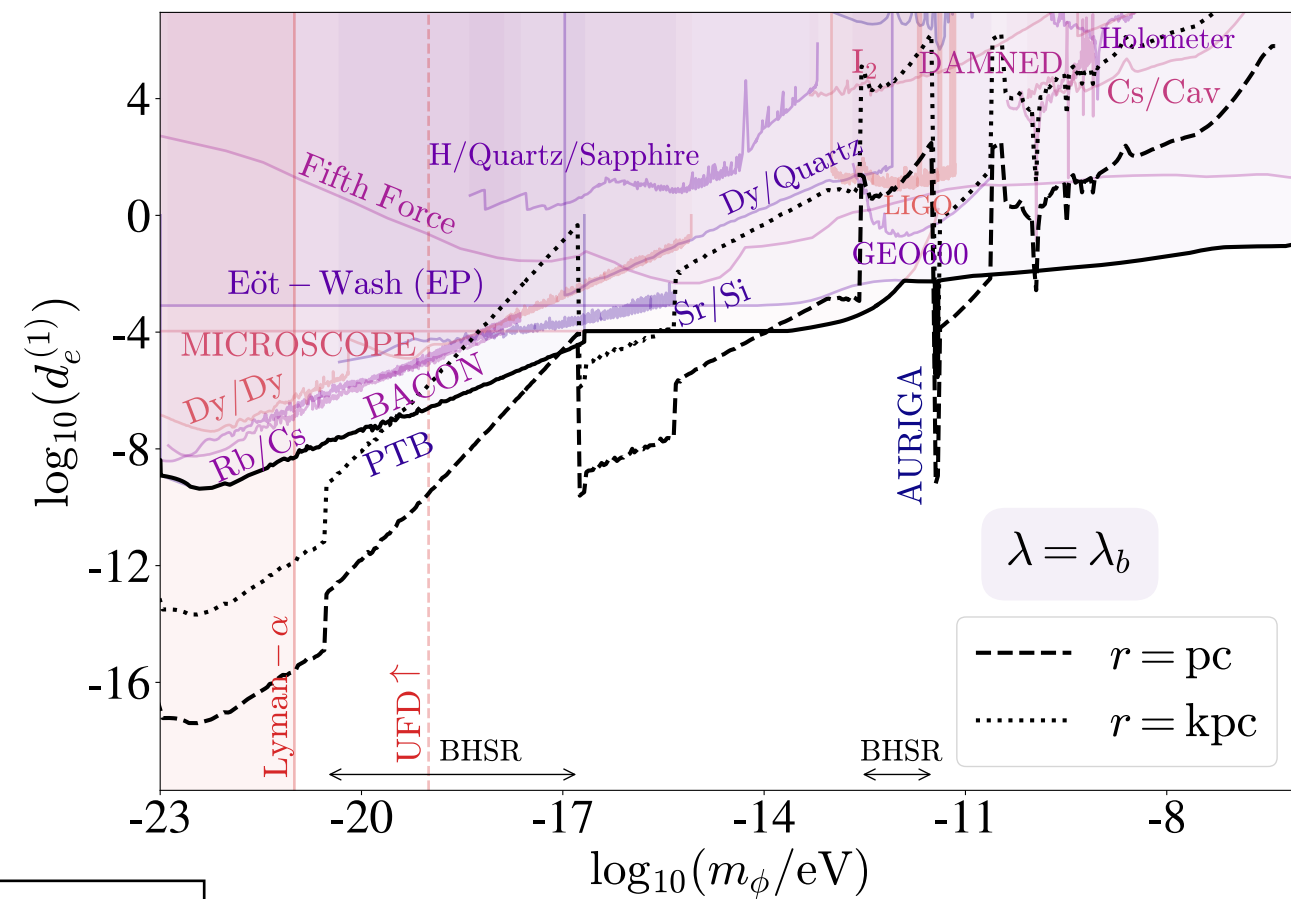
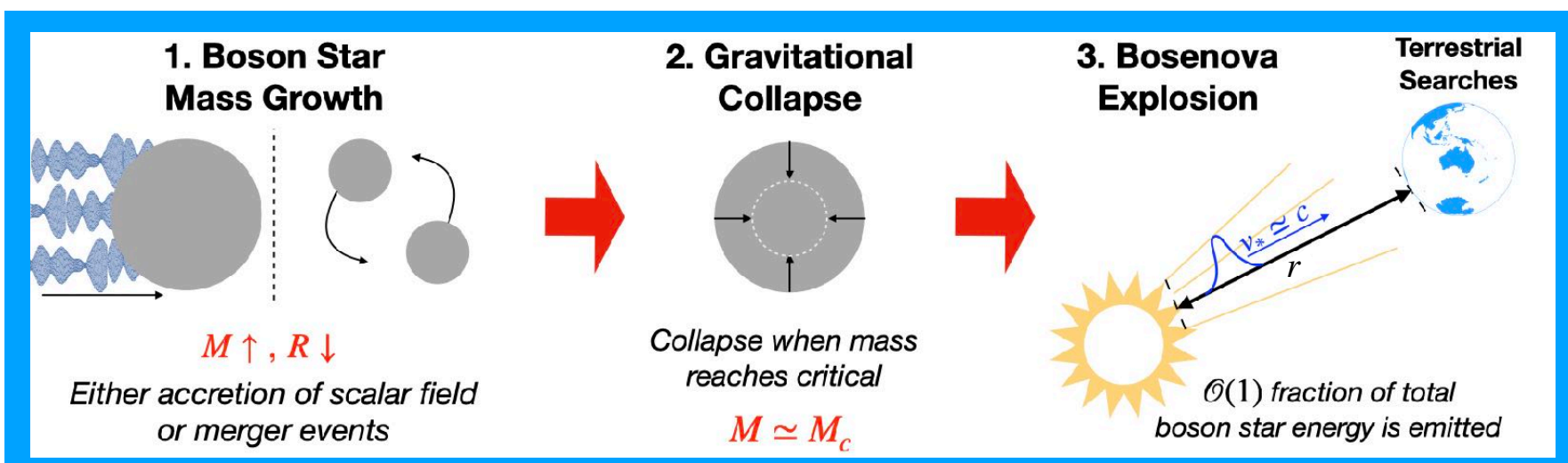
Explicit rate calculation:
 # bosenovae/galaxy today
 can be as large as few/day

Maseizik, Sigl (2404.07908)

Boson Star Collapse \longrightarrow Bosenova \longrightarrow Signals

JE, Shirai, Stadnik, Takhistov (2106.14893)
 Arakawa, JE, Safronova, Takhistov, Zaheer (2306.16468, 2402.06736)

Direct search for relativistic axion waves
 in terrestrial experiments



Caveats

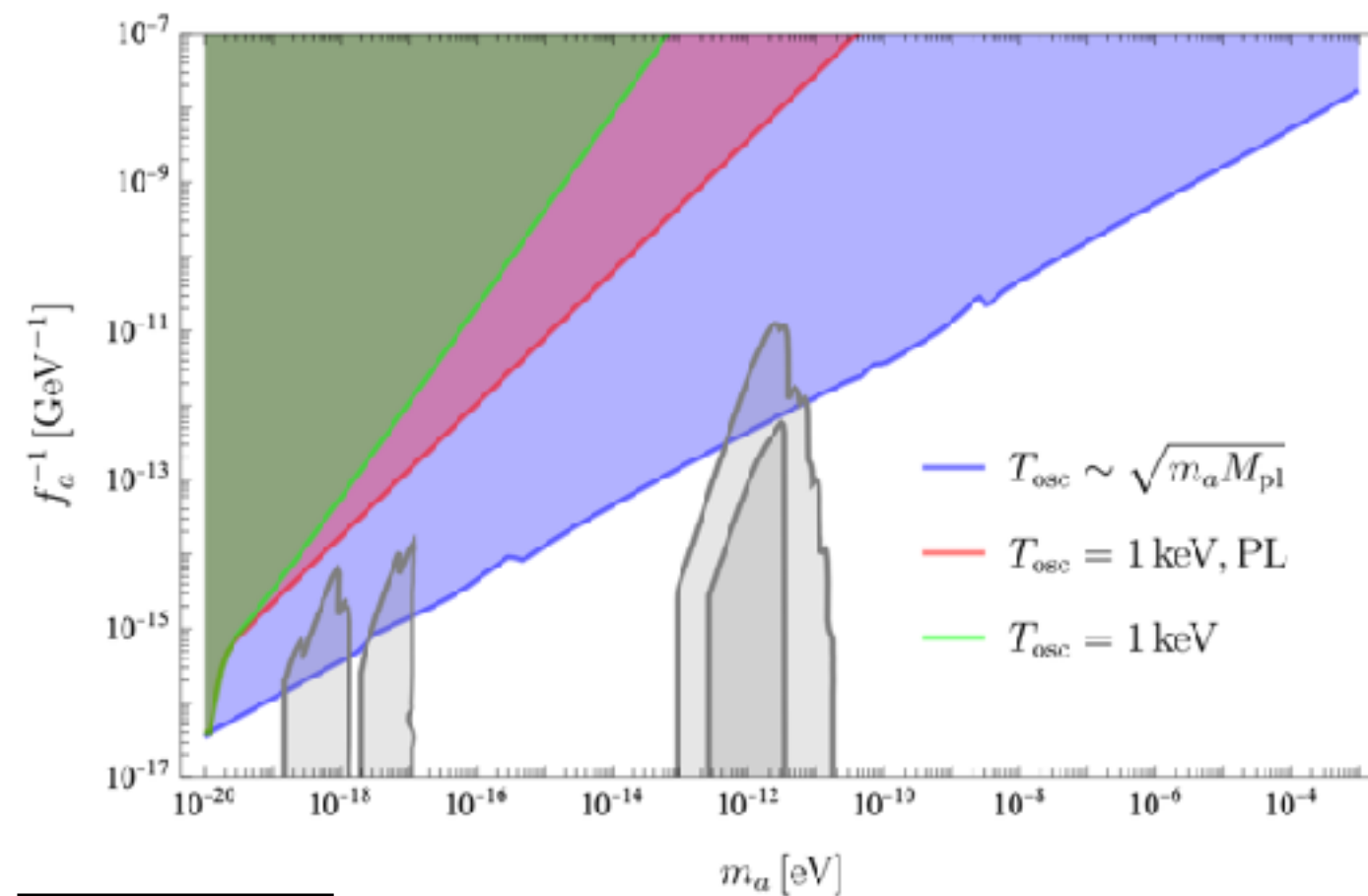
- Need nearby collapse, likely rate-limited
- Rate is *highly* model-dependent

Fox, Weiner, Xiao (2302.00685)

Bosenovae deplete DM at $z \gtrsim 20$,
 constrain total emitted energy

CMB/SDSS limit

$$f_{dDM} \equiv \frac{\Omega_{dDM}}{\Omega_{dDM} + \Omega_{DM}} \leq 2.62\% \text{ (at } 2\sigma\text{)}$$

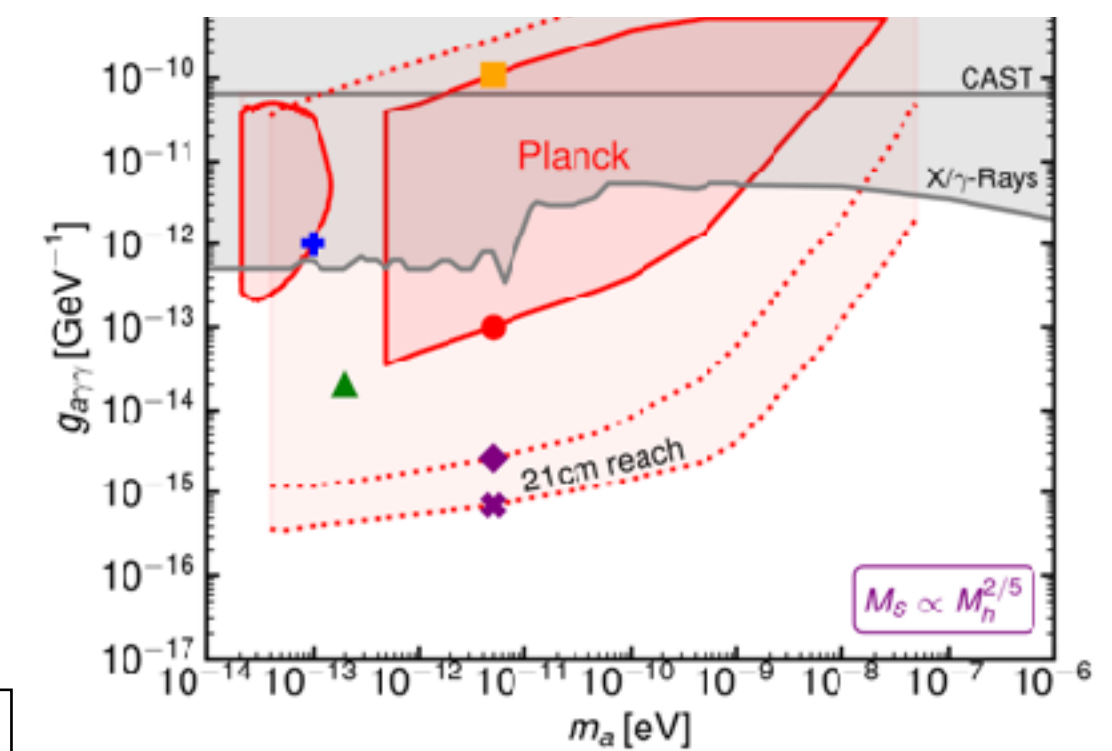
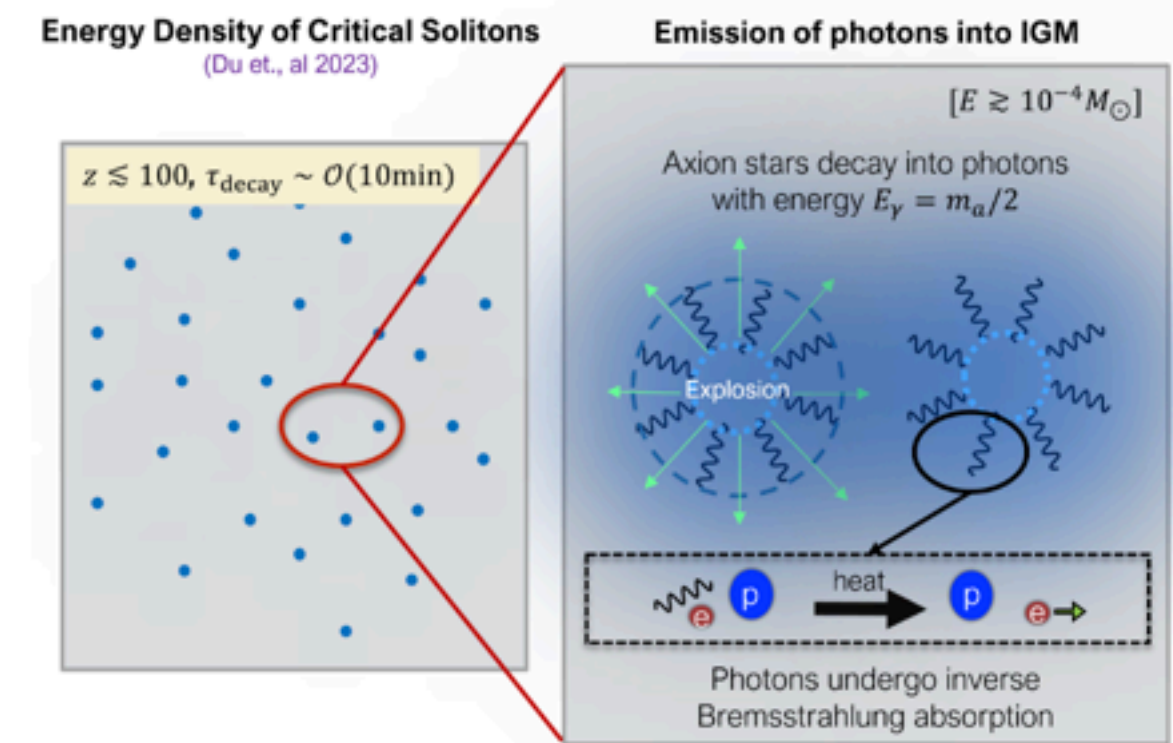


Caveats

- Requires enhancement of small-scale matter power spectrum
- Assumptions about growth rate

Escudero, Pooni, Fairbairn, Blas, Du, Marsh (2302.10206)

For boosted $g_{\phi\gamma}$, bosenova of photons instead,
 look for indirect signals (e.g. IGM heating)

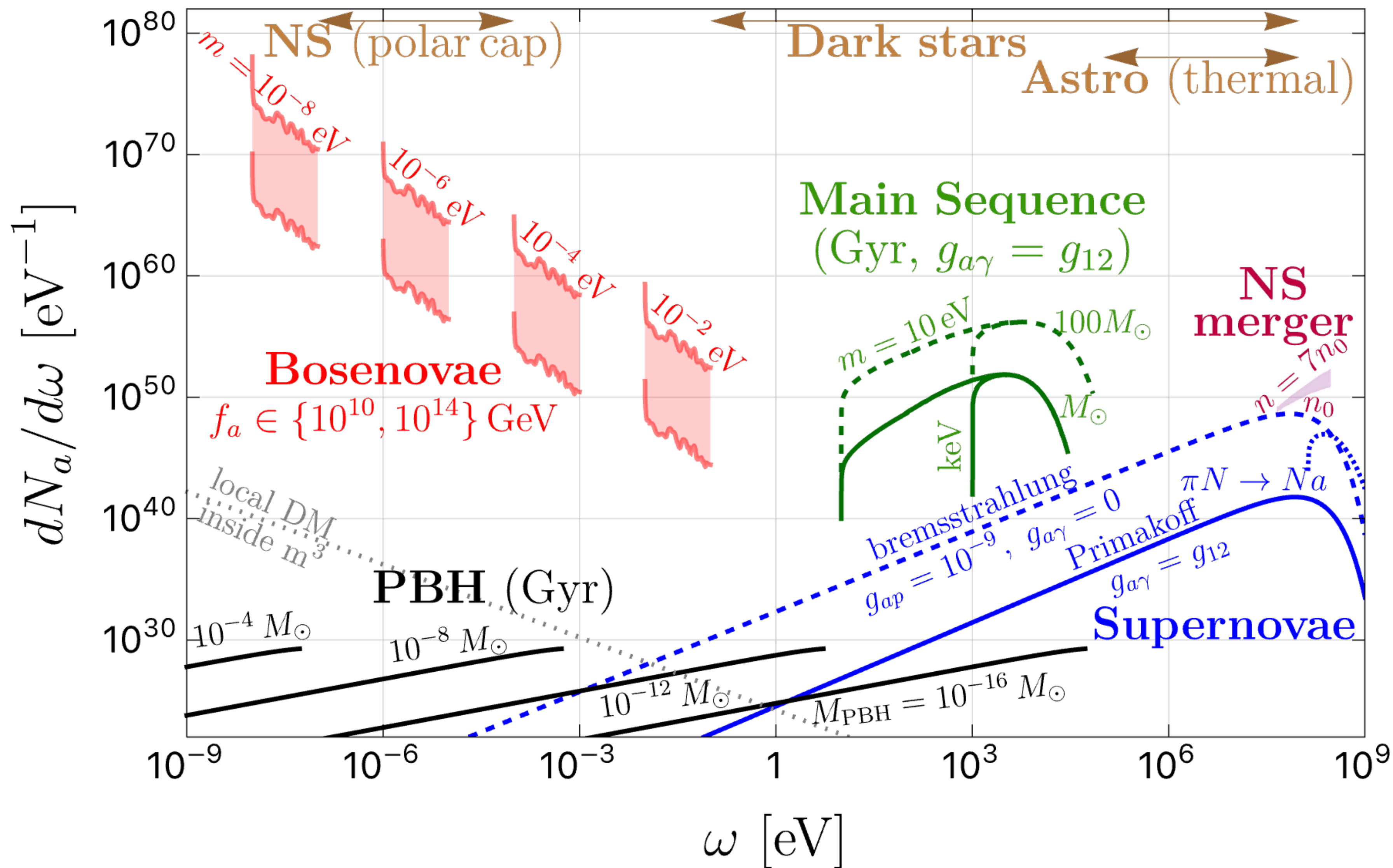


Caveats

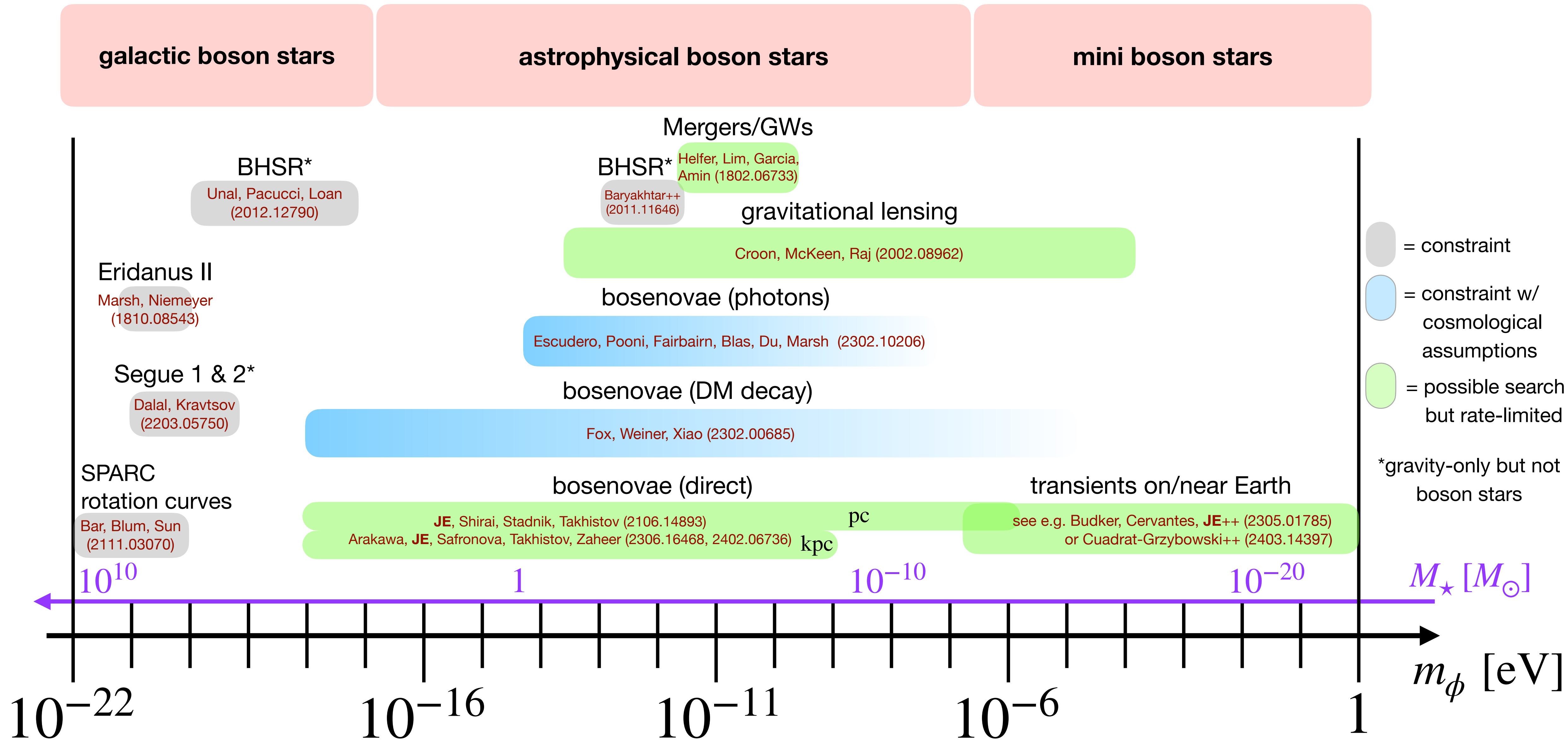
- Requires extrapolation of core-halo relation to large m_ϕ
- Need enhanced $g_{\phi\gamma}$ relative to simplest model

Diffuse Axion Background from Axion Bursts

(including bosenovae)



Searches for Boson Stars



ULDM Ground States

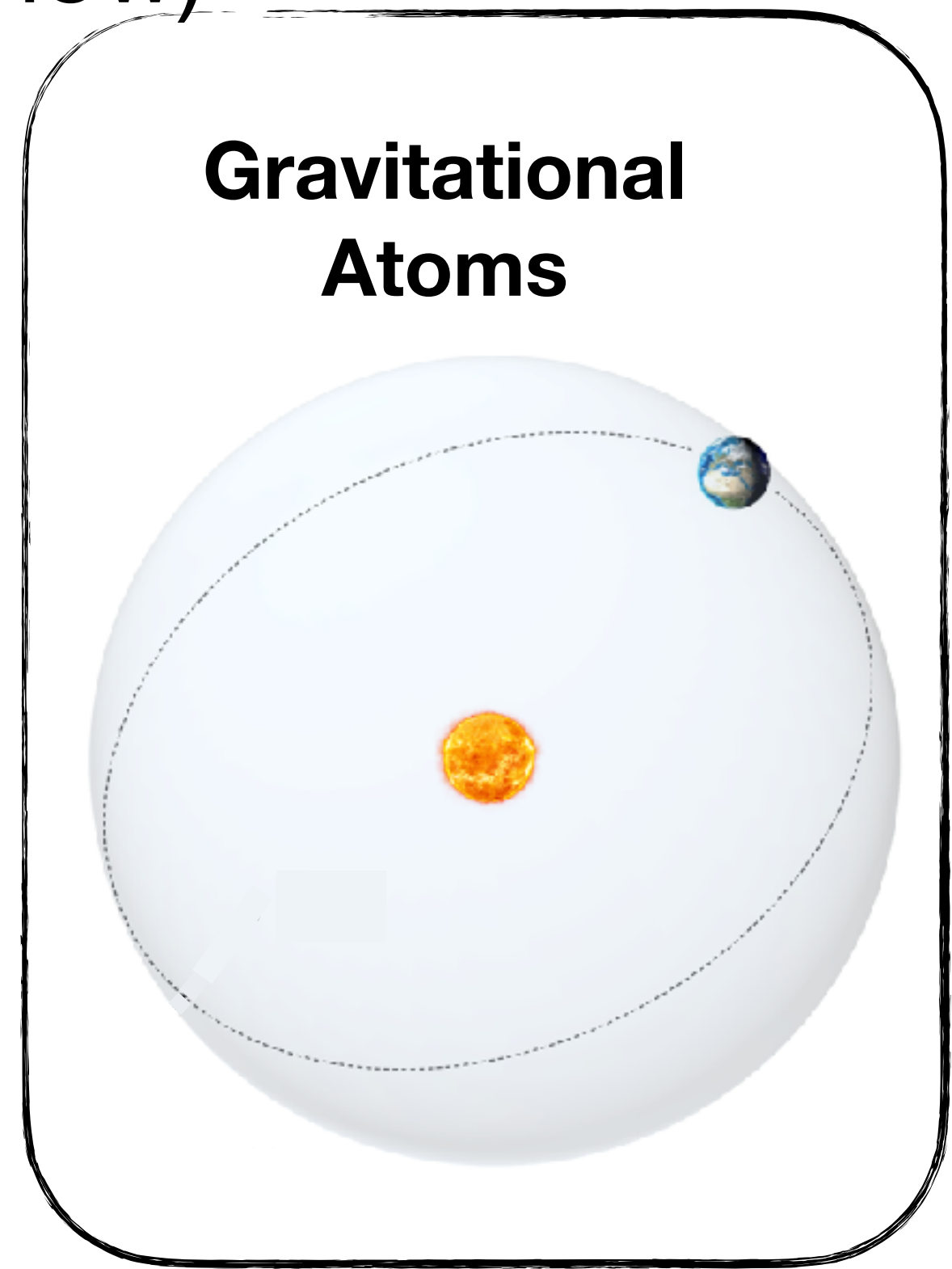
“Quantum” pressure (Repulsive) Gravity (Attractive) Self-interactions (usually attractive)

$$i\psi = \left[-\frac{\nabla^2}{2m_\phi} + V_g(\psi) \right] \psi - \frac{\lambda}{8m_\phi^2} |\psi|^2 \psi$$

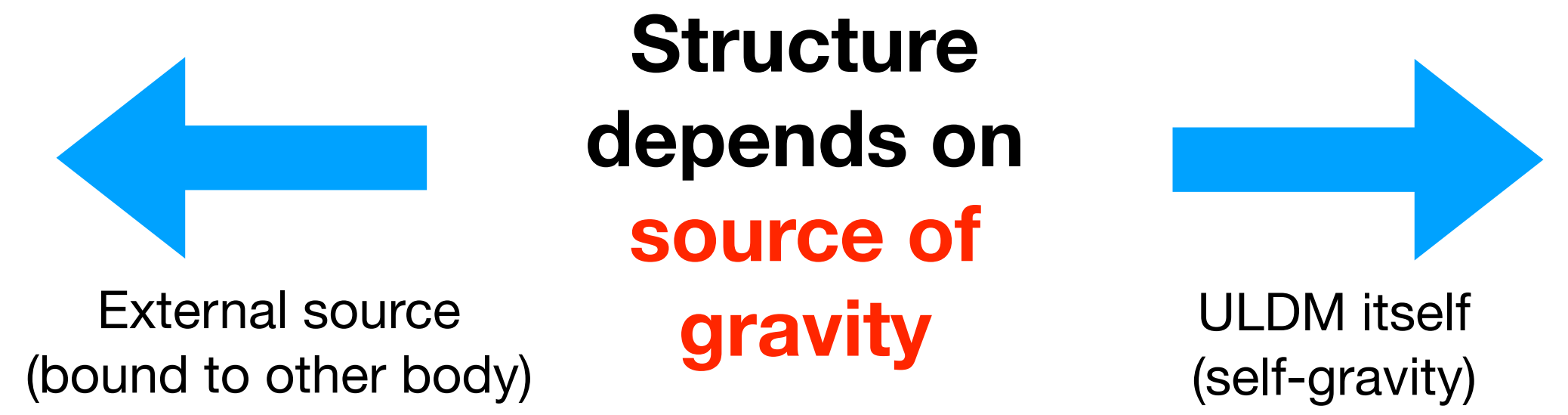
$\lambda \sim m_\phi^2 / f^2$

(small when density small, return to this later)

(now)



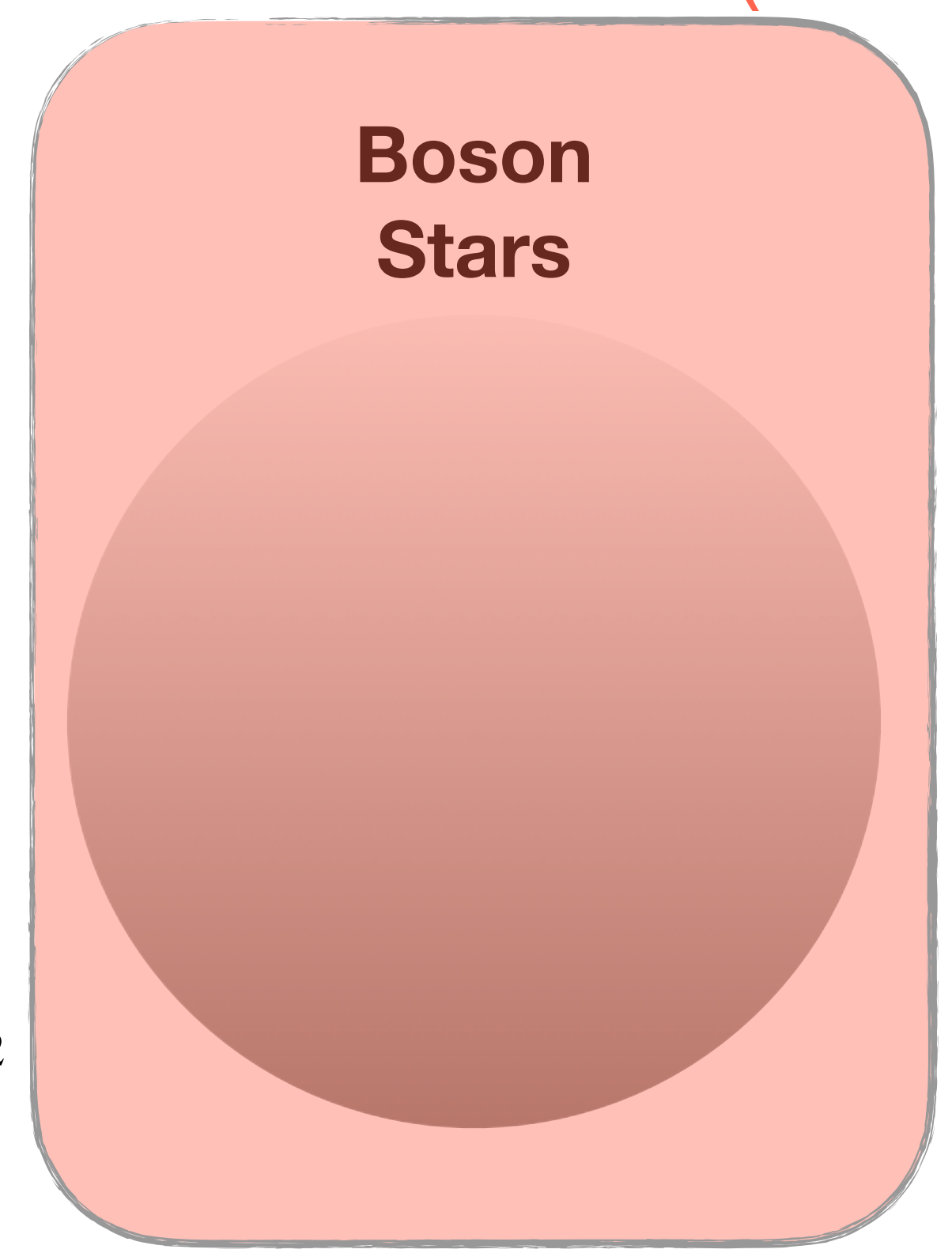
Balance of gradient+gravity in the field



$$V_g(\psi) = \frac{GMm_\phi}{r}$$

$$\nabla^2 V_g(\psi) = 4\pi Gm_\phi^2 |\psi|^2$$

(done)



What is a Gravitational Atom?

Bound states around an external body,
parameterised by $\psi_{n\ell m} = R_{n\ell}(r)Y_{\ell m}(\theta, \phi)$



Gravitational potential:

$$V_g(r) = -\frac{\alpha_g}{r} = -\frac{Gm_\phi M_\odot}{r}$$

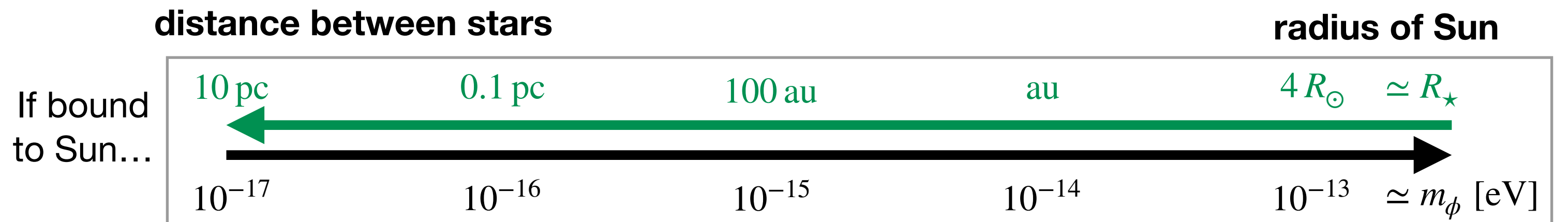
Gravitational “Bohr radius”:

$$R_\star = (m_\phi \alpha_g)^{-1} \simeq \frac{1}{Gm_\phi^2 M_\odot}$$

compare to Hydrogen:

Coulomb potential: $V(r) = -\frac{\alpha}{r}$

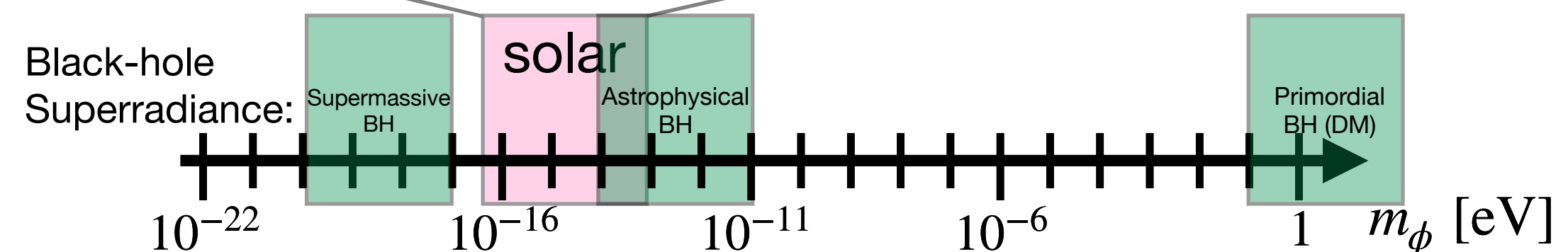
Bohr radius: $a_0 = (m_e \alpha)^{-1}$



“Quantum” pressure
(Repulsive)

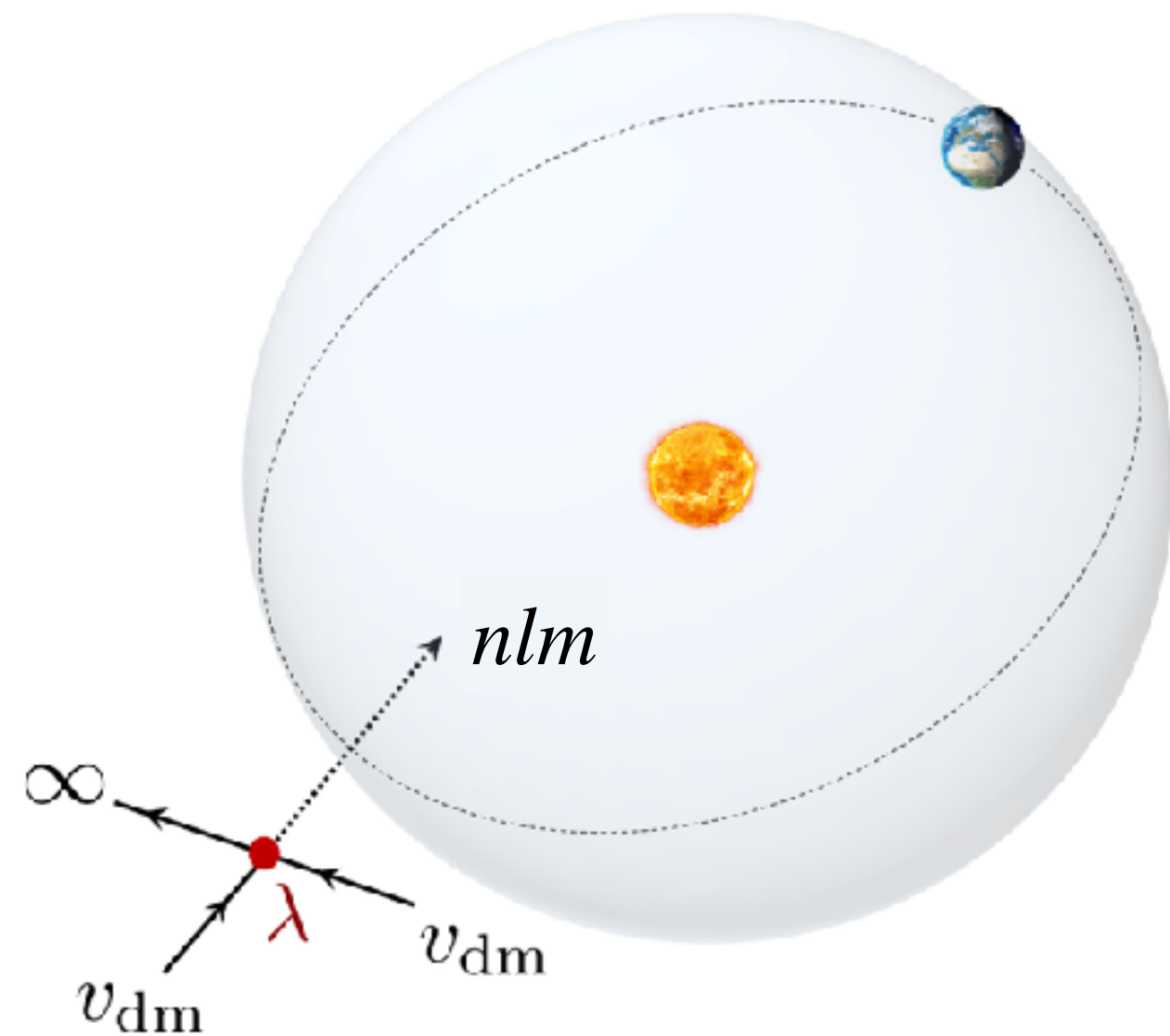
Gravity from source
(Attractive)

$$i\psi = \left[-\frac{\nabla^2}{2m_\phi} + V_g(\psi) \right] \psi$$

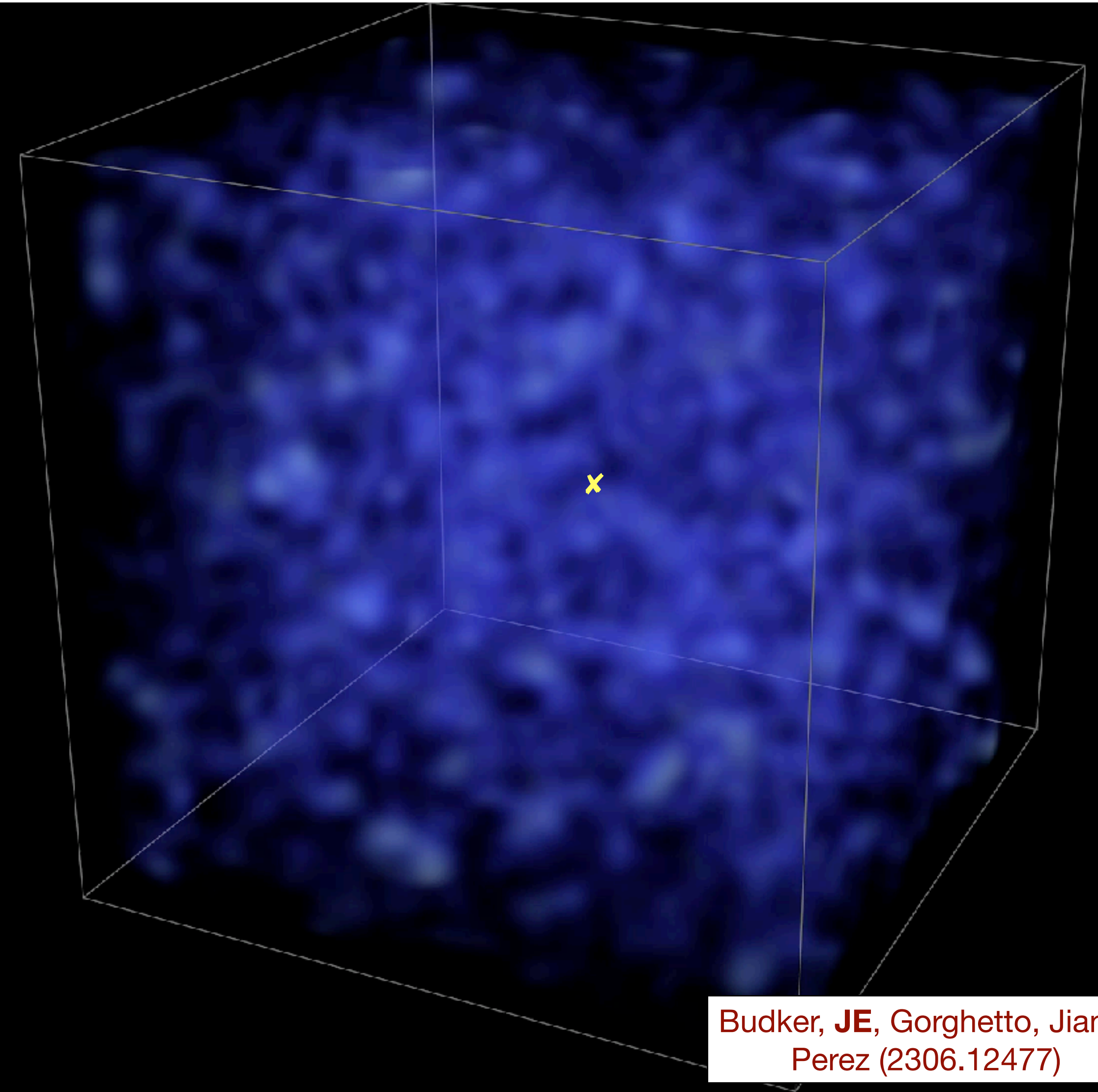


Gravitational Atoms from ULDM Capture

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} - \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$



Self-interactions can move particles from scattering states to bound states (and vice versa)

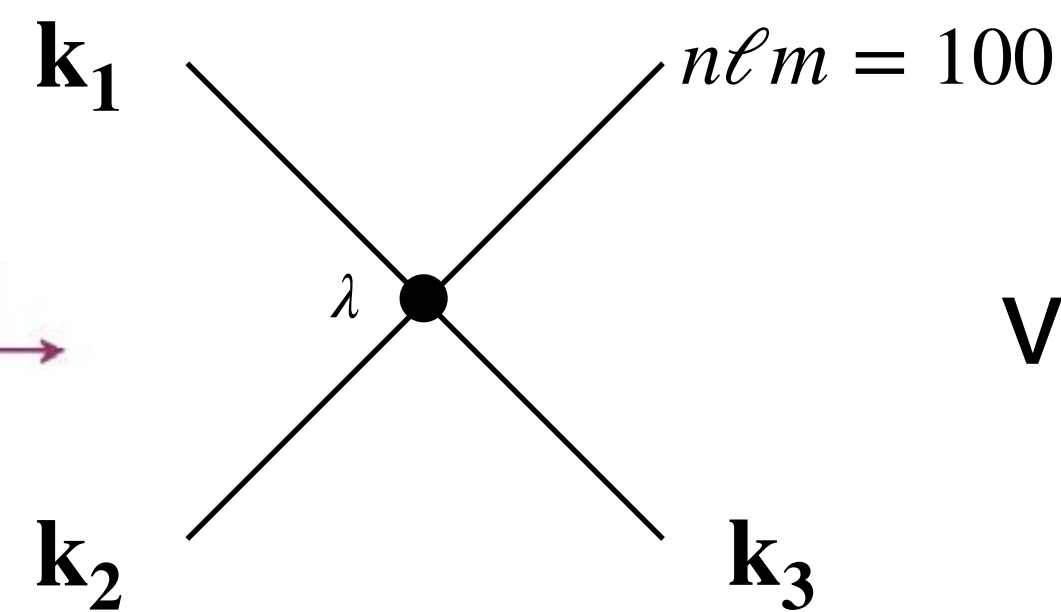


Budker, **JE**, Gorghetto, Jiang, Perez (2306.12477)

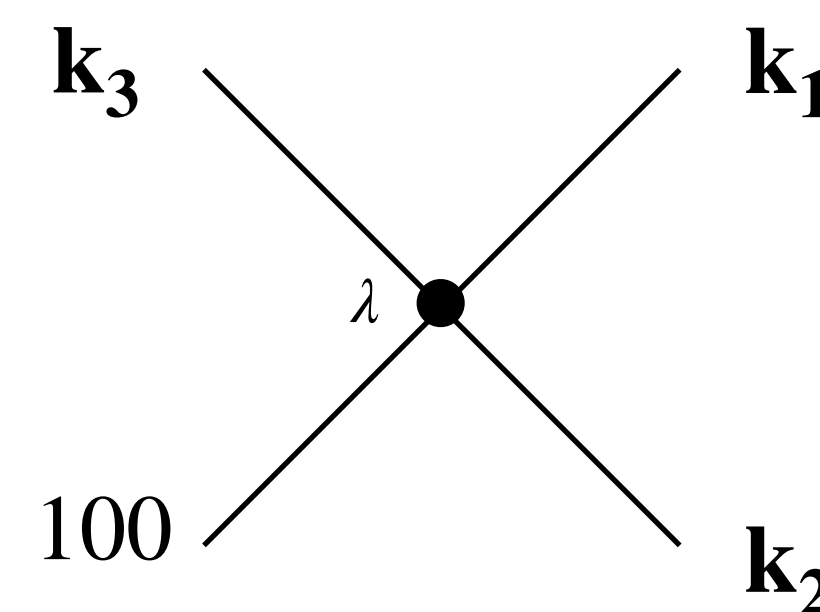
Gravitational Atoms from ULDM Capture

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} - \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$

Stimulated capture



Ionization



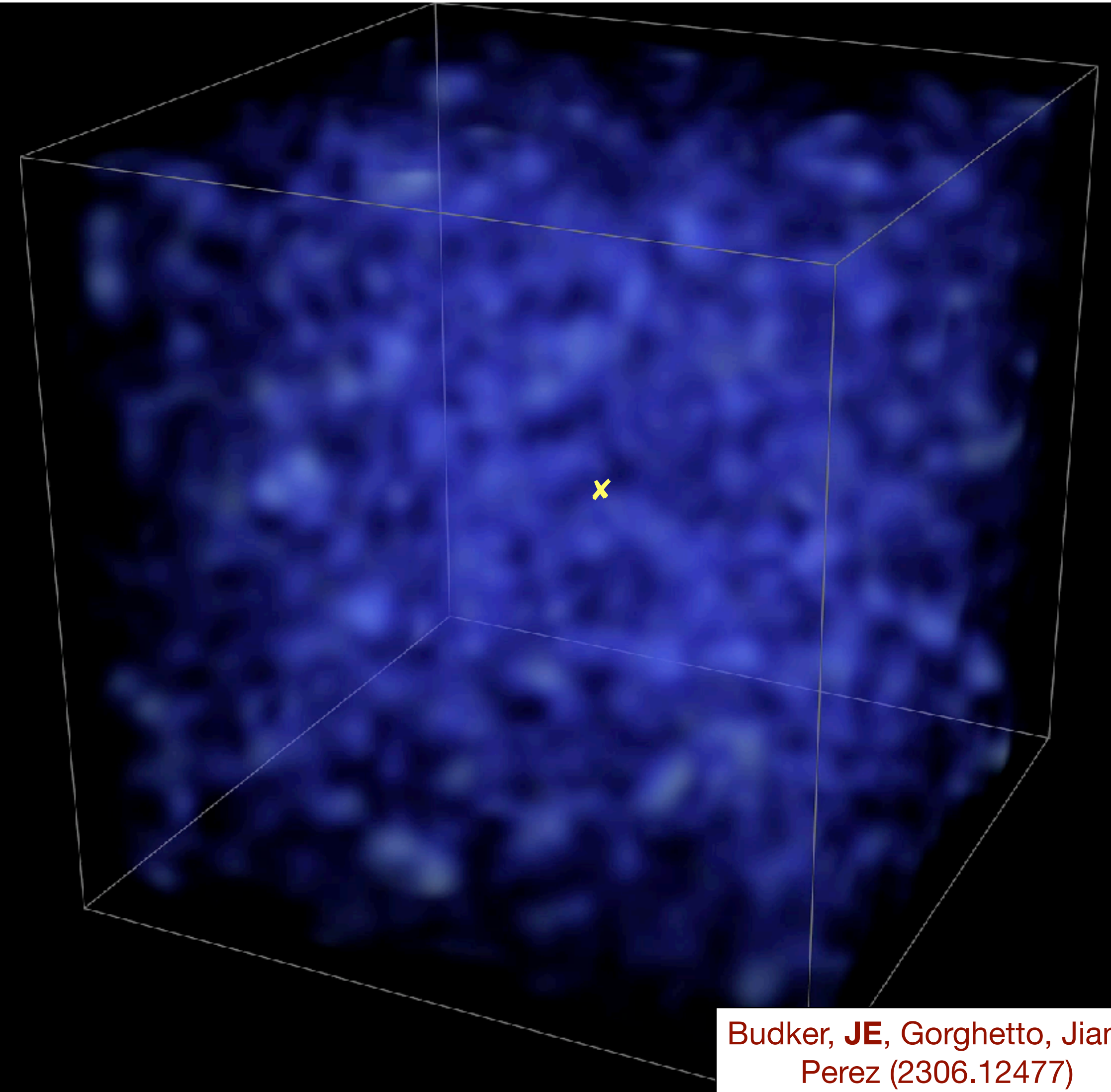
VS

$$\frac{dM_\star}{dt} \simeq C + \Gamma (\Gamma_1 - 2\Gamma_2) M_\star$$

$\Gamma > 0$: Exponential growth

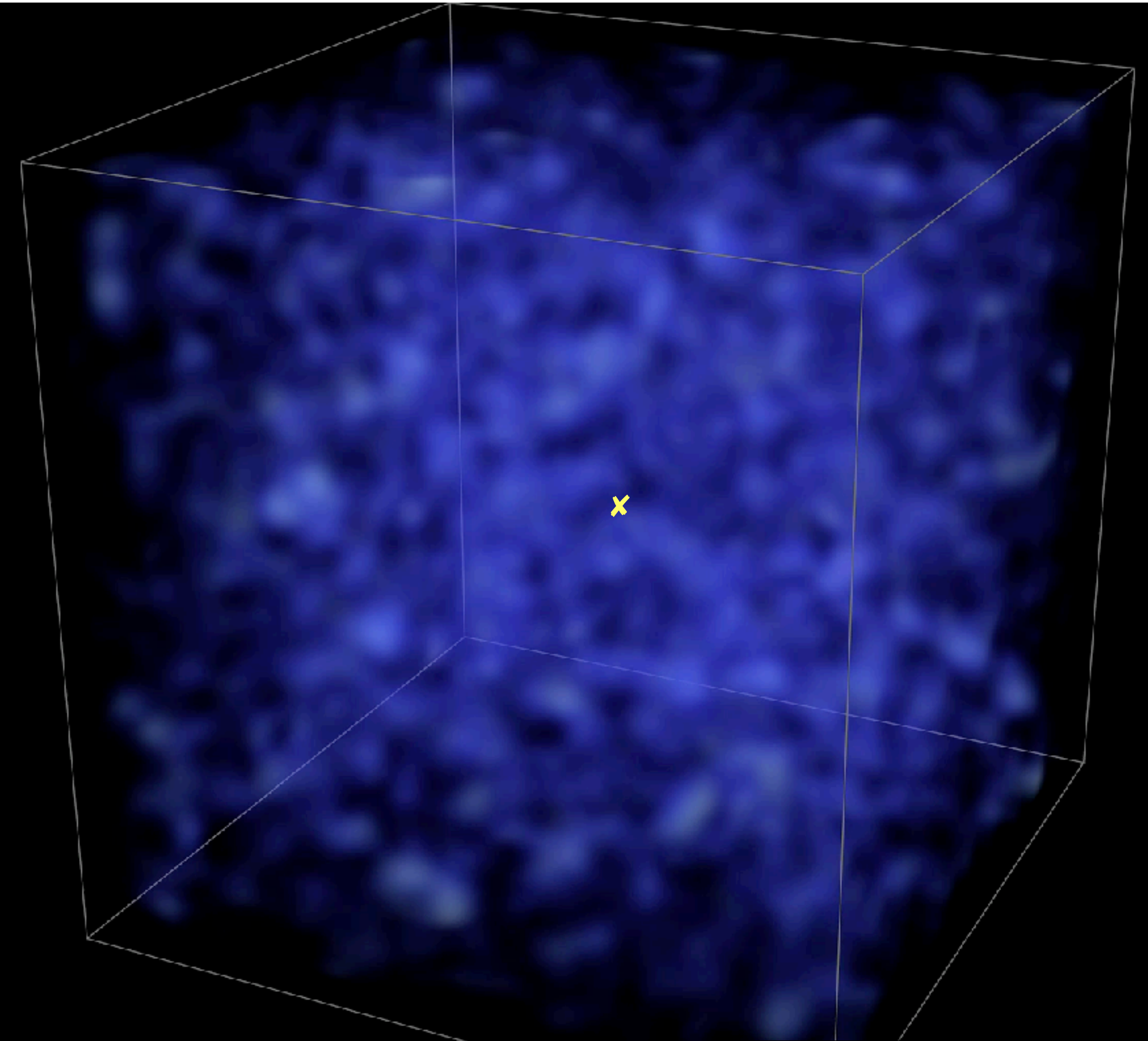
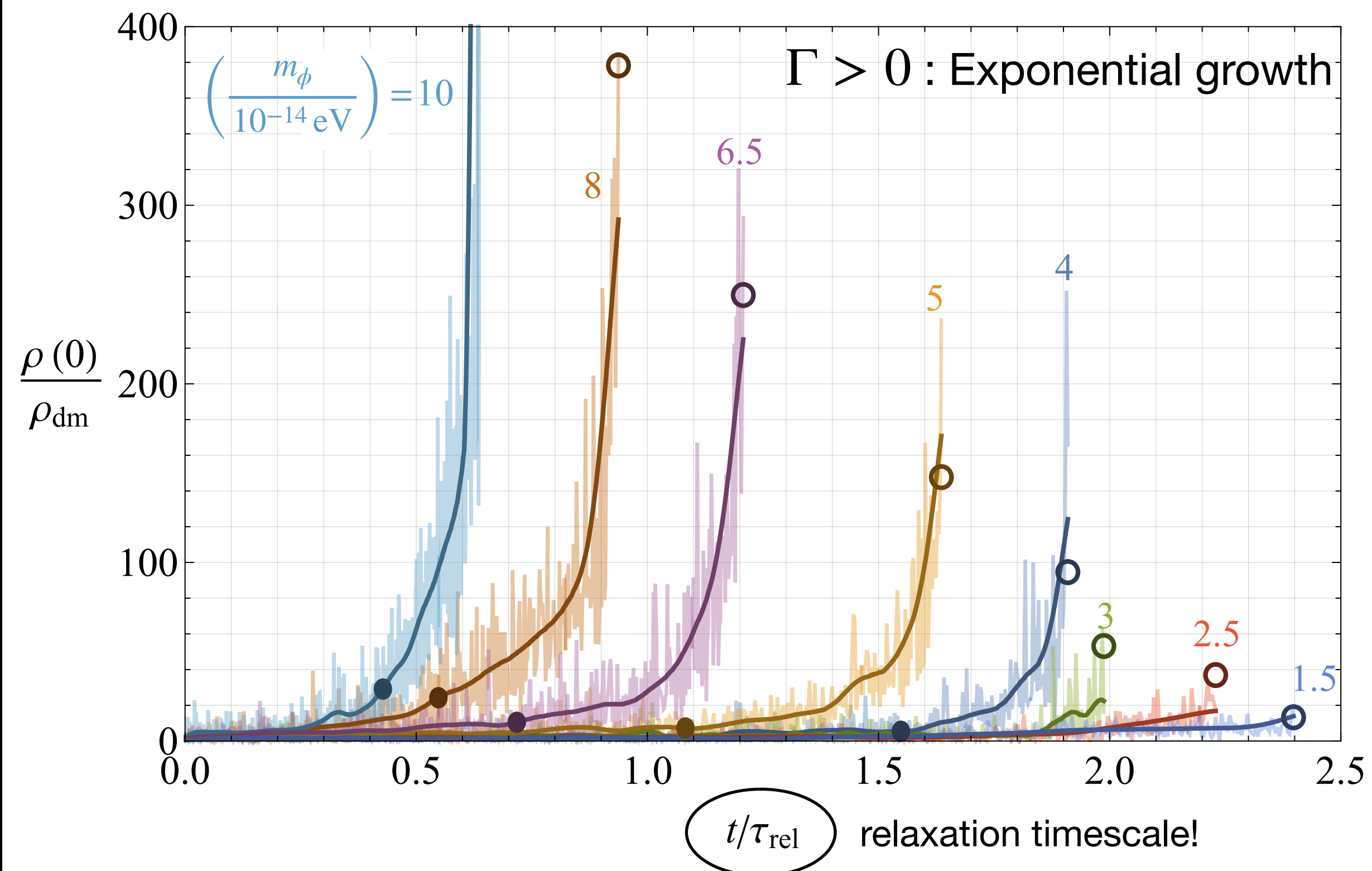
$\Gamma < 0$: Saturation

determines late-time behavior

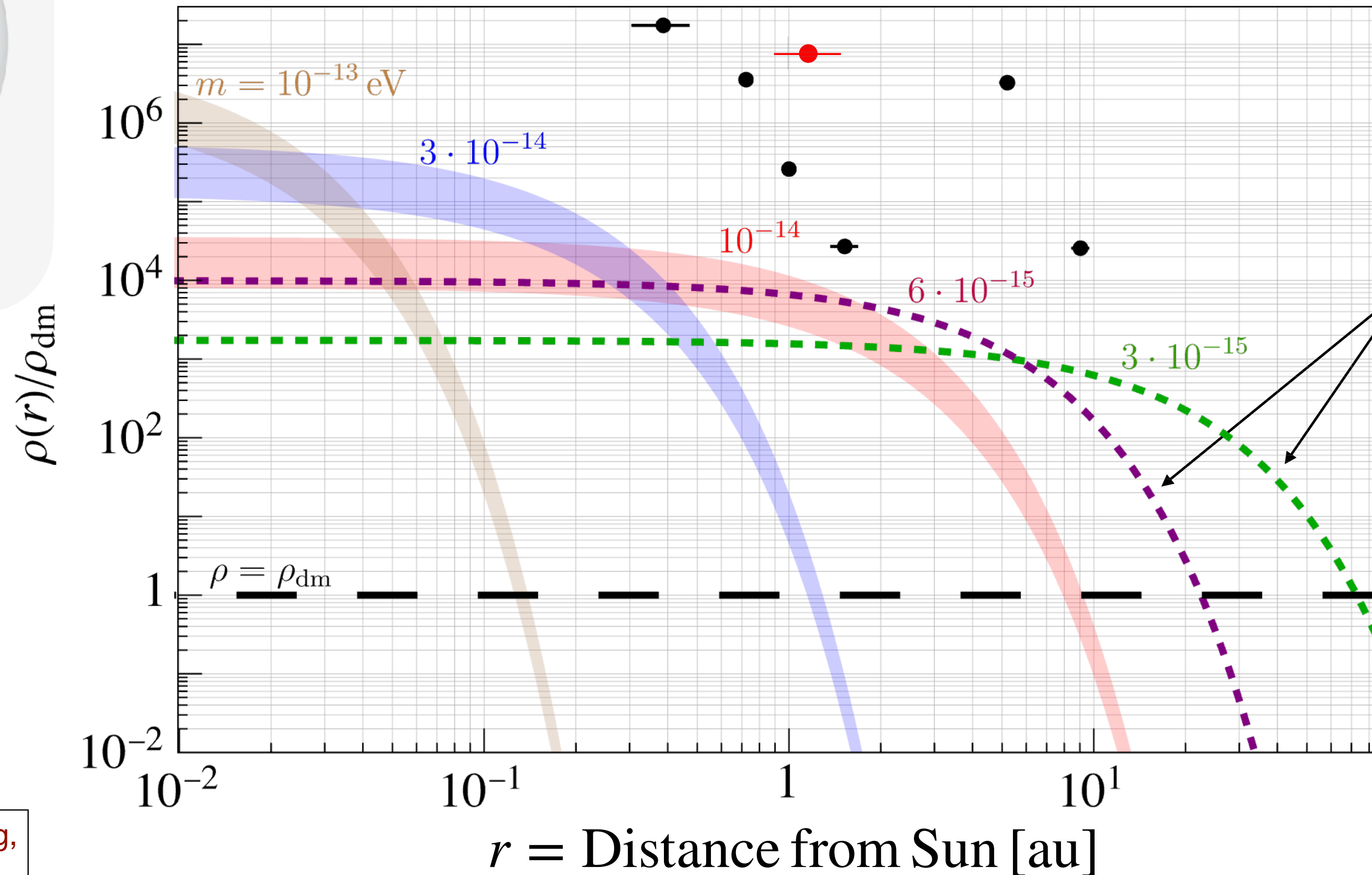


Gravitational Atoms from ULDM Capture

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} - \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$



Gravitational Atom in our Solar System



relevant when DM velocity is very low, $v_{\text{dm}} \ll 200 \text{ km/s}$

Budker, **JE**, Gorghetto, Jiang, Perez (2306.12477)

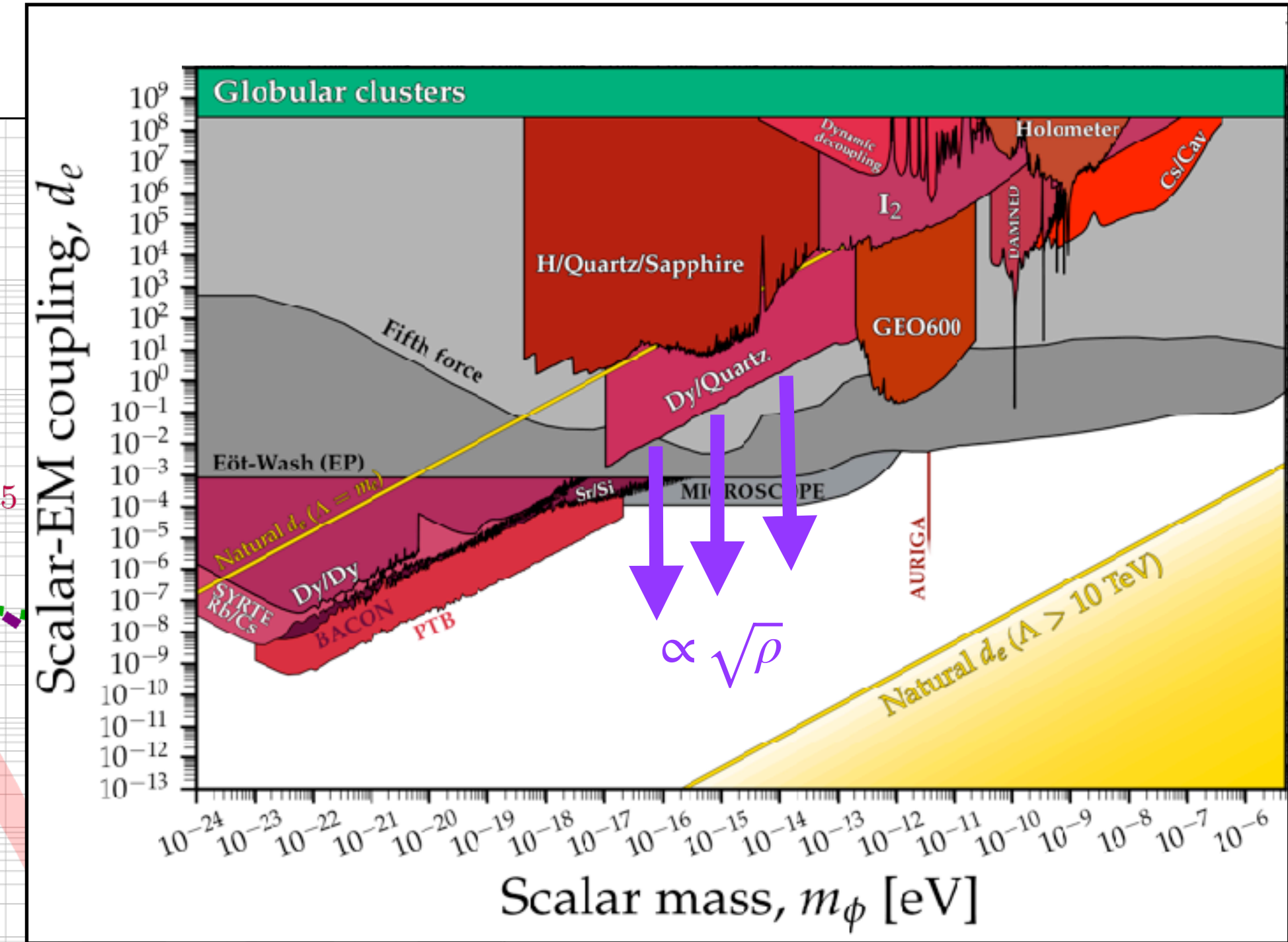
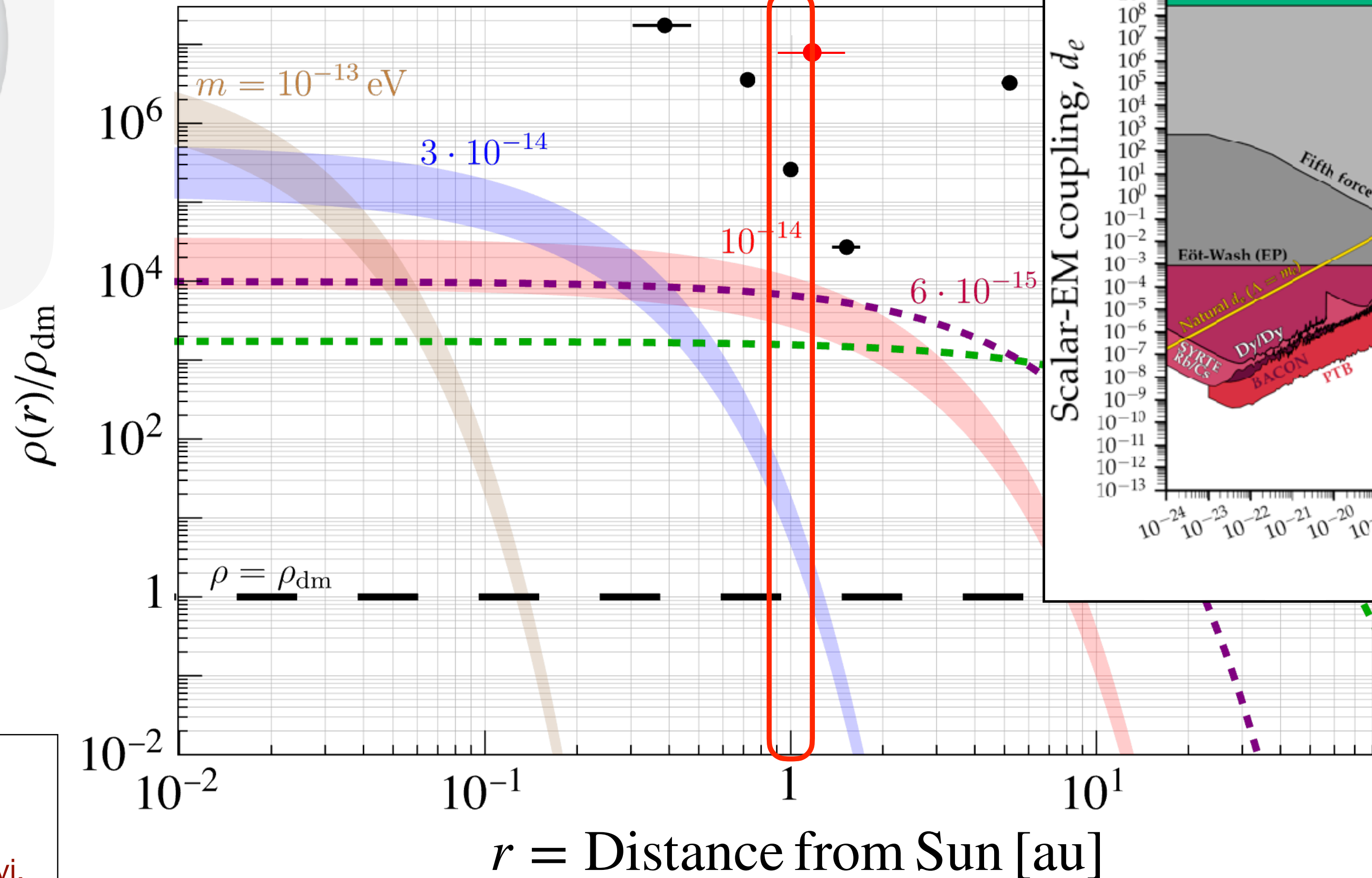
Black points: Constraints from planets
Piitjev and Pitjeva (1306.5534)
Red point: Constraint from Benu asteroid
Tsai, **JE**, Arakawa, Farnocchia, Safronova (2210.03749)

Gravitational Atom in our Solar System



Searches on Earth

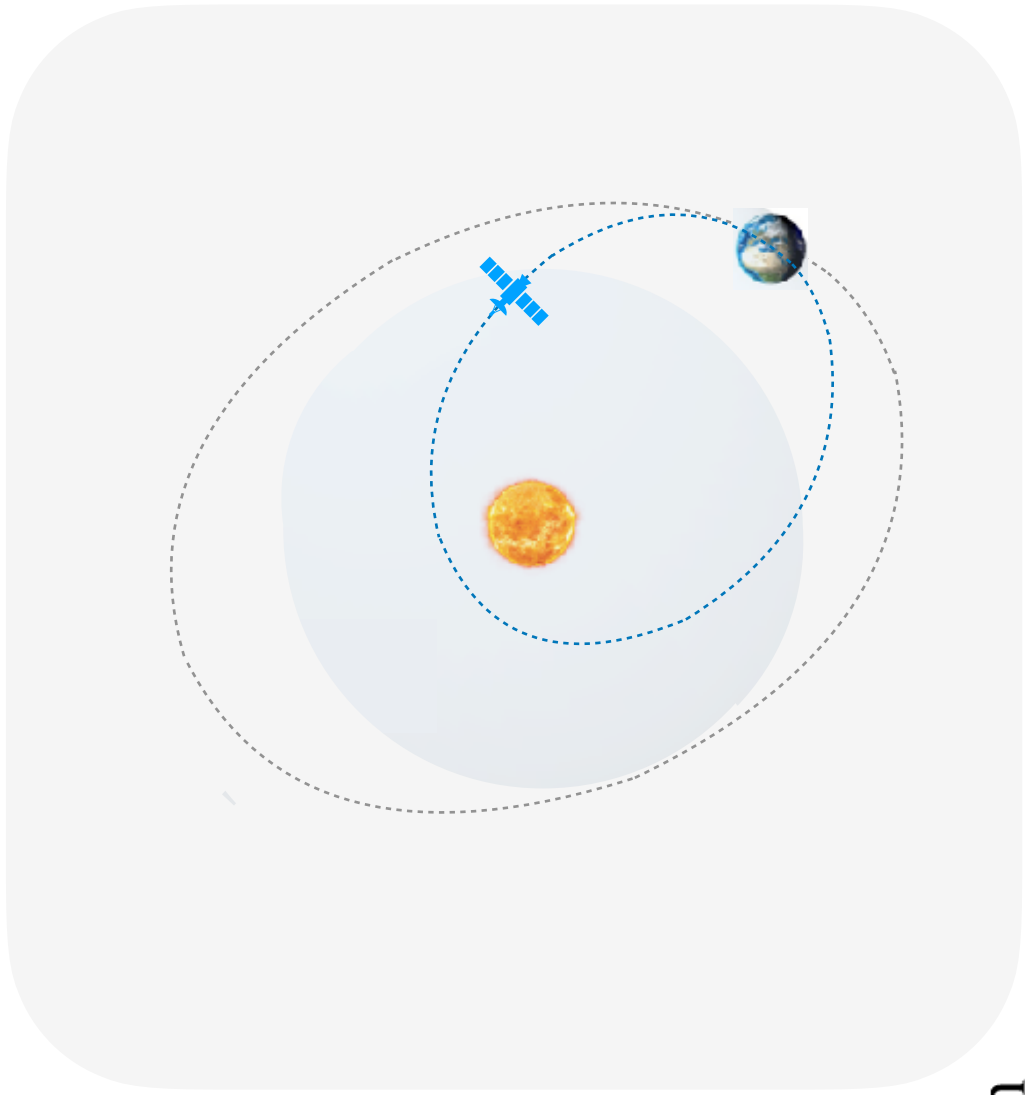
$r = au$



Banerjee, Budker, **JE**, Kim, Perez (1902.08212)
with Flambaum, Matsedonskyi, (1912.04295)

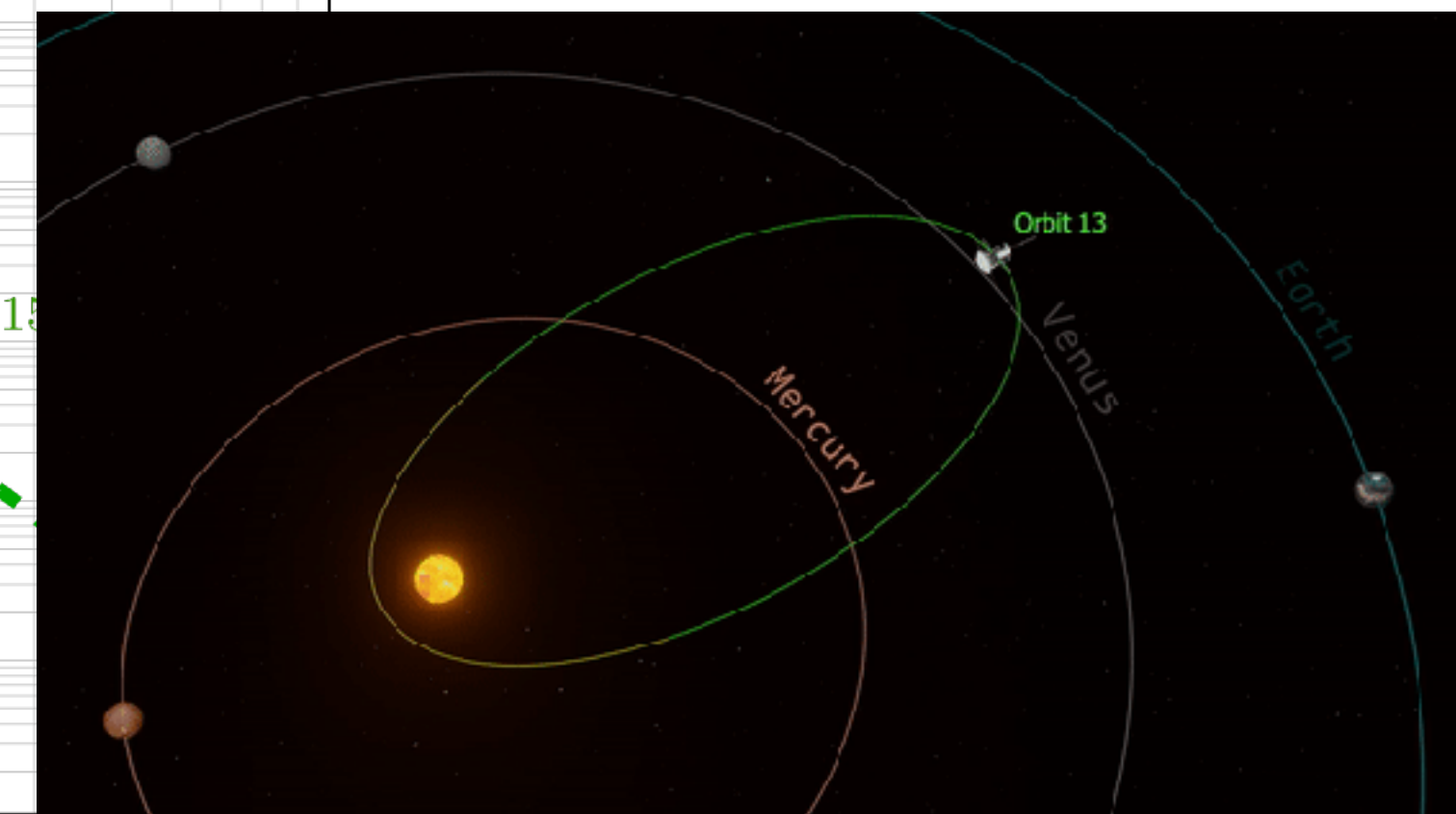
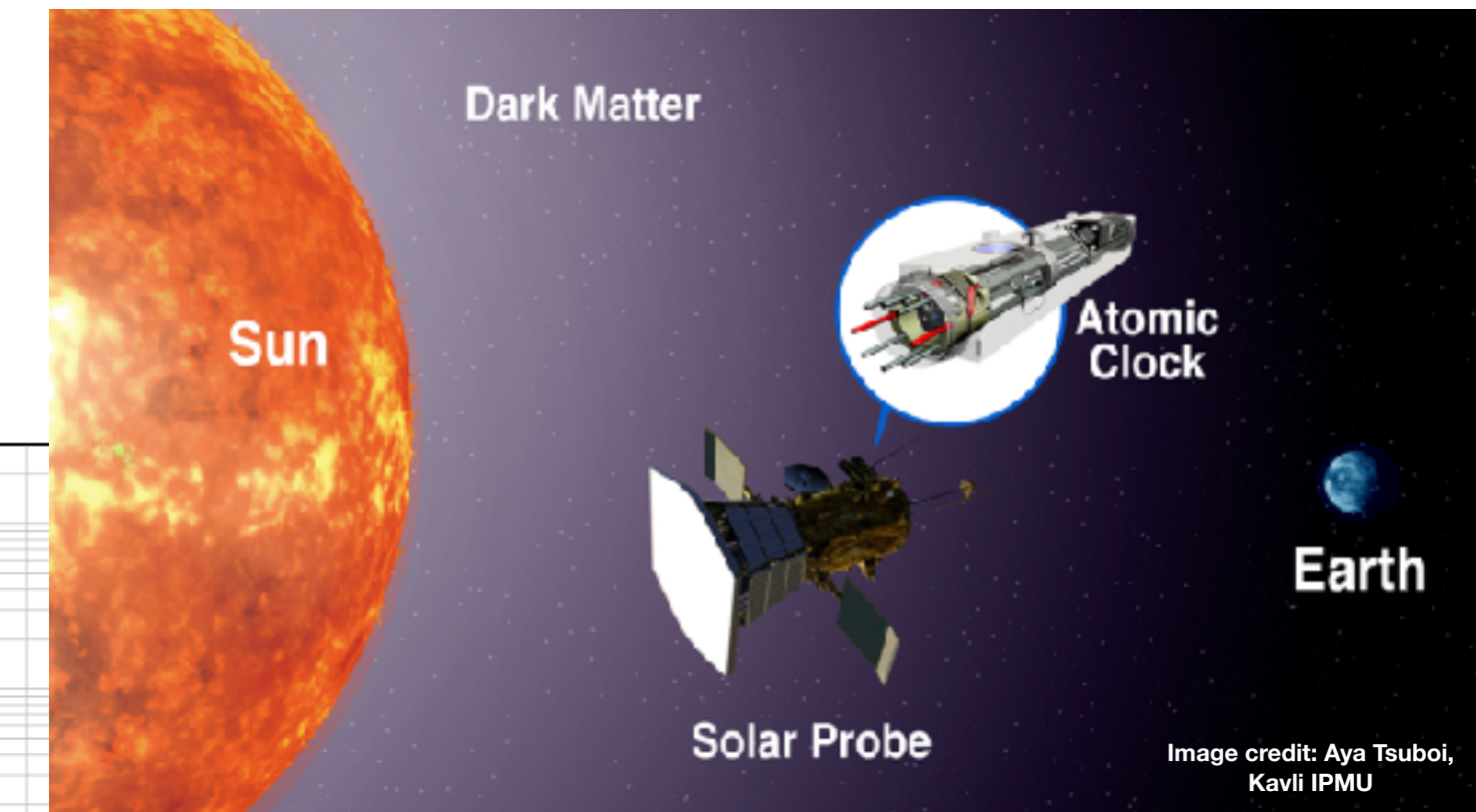
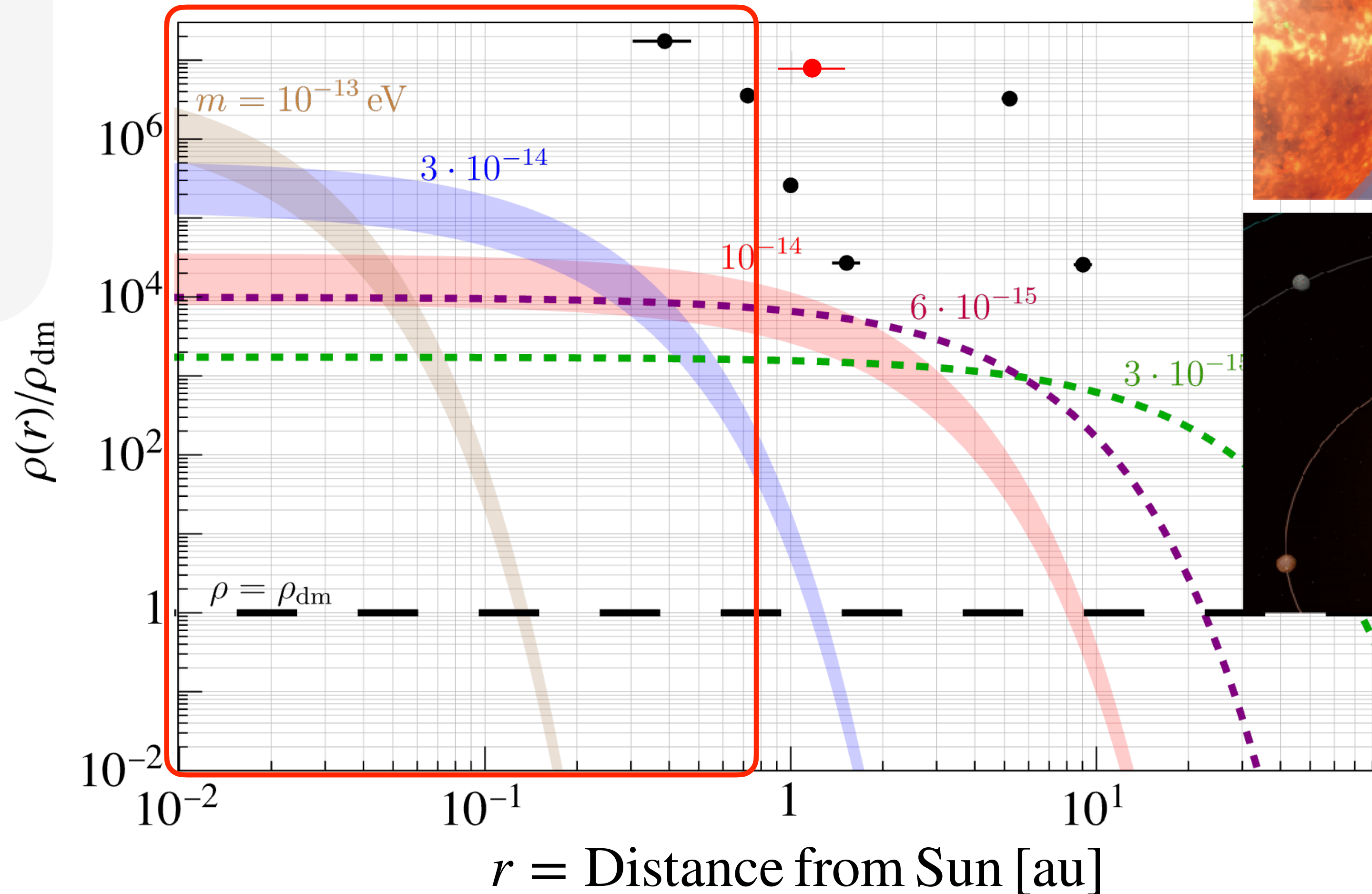
- Experimental Searches:**
- Savalle++, PRL 2020
 - Aharony++, PRD 2021
 - Hanneke++, QST 2020
 - Oswald++, PRL 2022
 - Tretiak++, PRL 2022
 - Manley++, PRD 2023
 - ...

Gravitational Atom in our Solar System



Space missions near the Sun

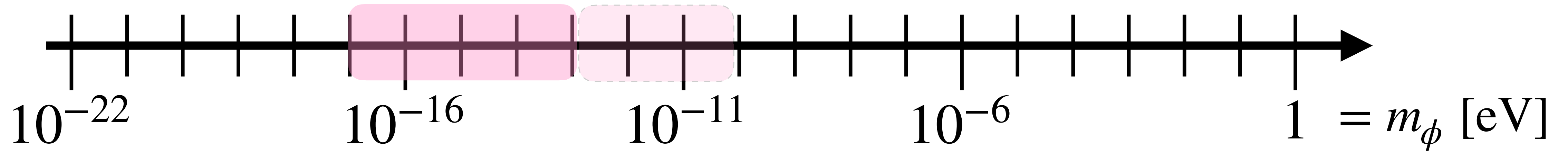
$$r < \text{au}$$



Proposal for future space missions to search near the Sun, where the ULDM density is expected to be largest

Tsai, JE, Safronova
(2112.07674)

ULDM Across the Galaxy



Gravitational atoms bound to...

Sun/stars

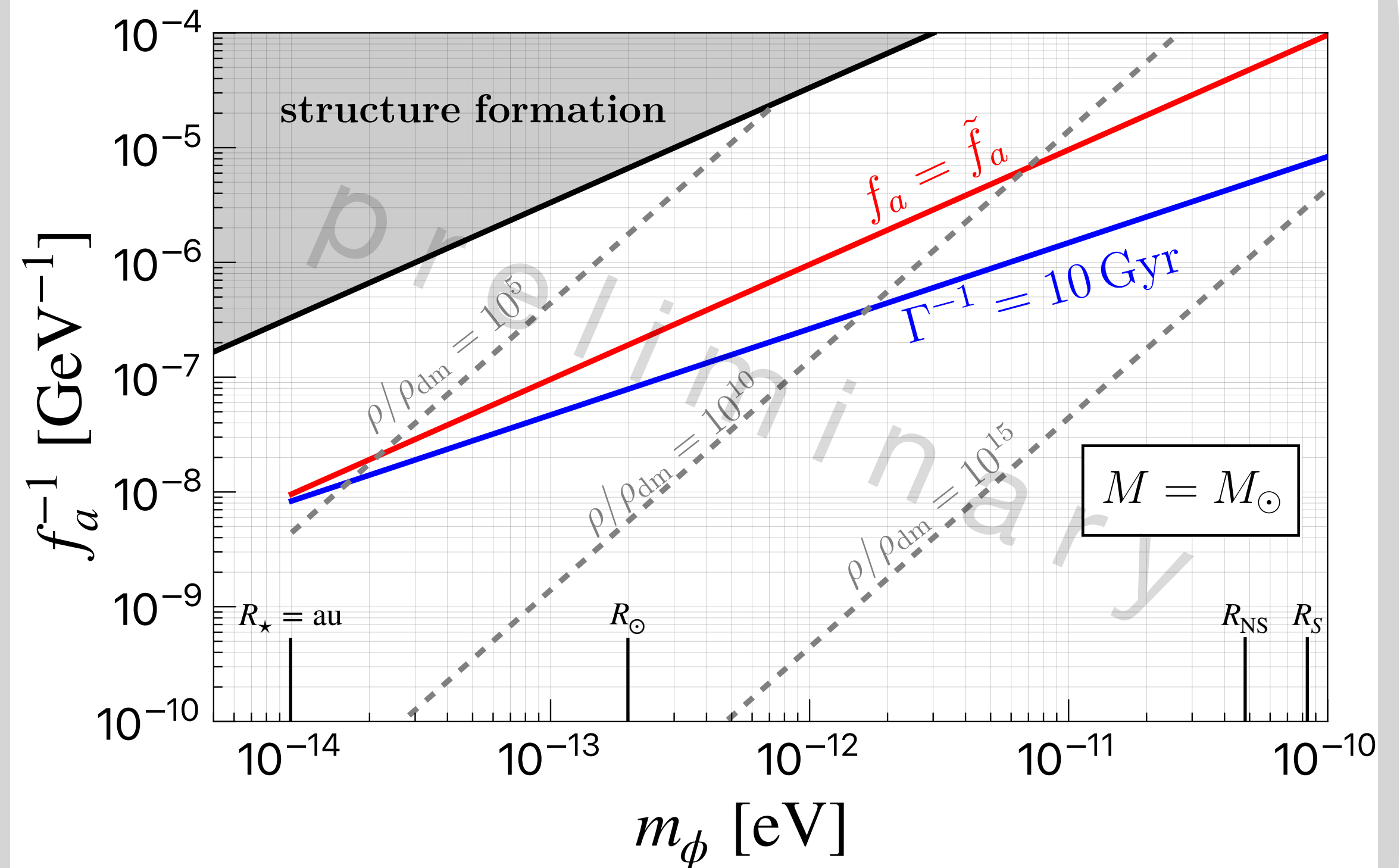
neutron stars / astro BHs

Budker, **JE**, Perez, ...with Banerjee, Kim (1902.08212)
 + Flambaum, Matsedonskyi (1912.04295)
 ...with Gorghetto, Jiang (2306.12477)

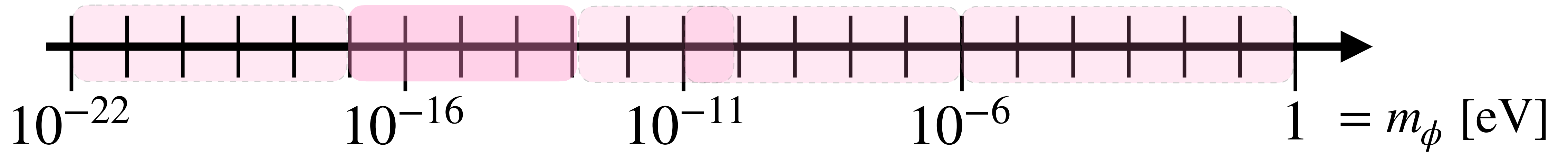
Ongoing work!



Samuel Gómez
 Master's Student
 Uppsala University



ULDM Across the Galaxy



Gravitational atoms bound to...

supermassive black holes

Budker, **JE**, Perez,
 ...with Banerjee, Kim
 + Flambaum, Matsedonskyi
 ...with Gorghetto, Jiang

(1902.08212)
 (1912.04295)
 (2306.12477)

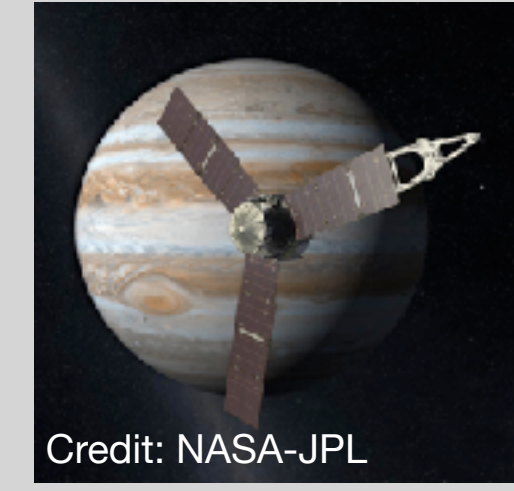
Sun/stars

neutron stars / astro BHs

Ongoing work!

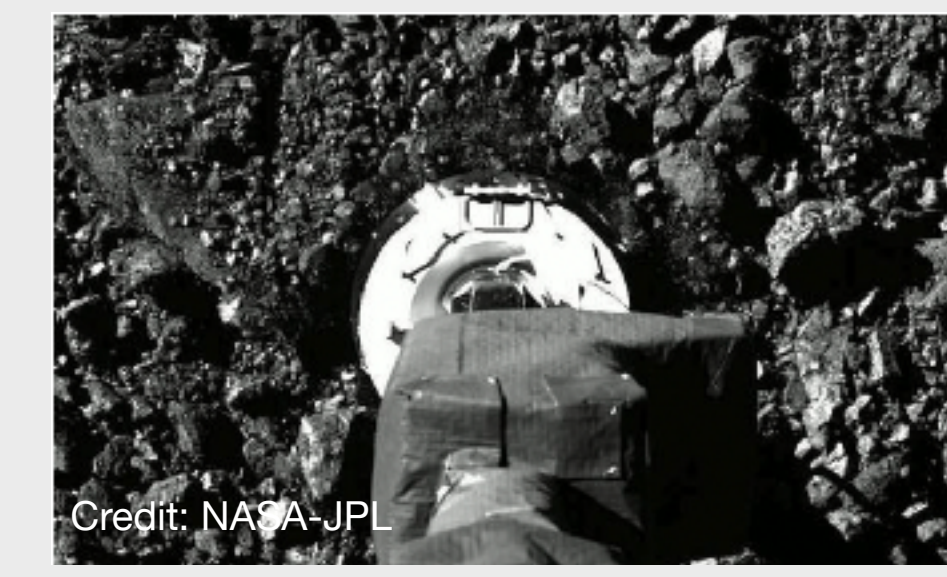
planets

JUNO mission to Jupiter



asteroids

OSIRIS-REx → OSIRIS-APEX



Utilize the rapidly growing subfield of asteroid tracking!

+ future asteroid missions

Birefringence signals:
Gan, Wang, Xiao (2311.02149)



Samuel Gómez
Master's Student
Uppsala University

Earth-bound Gravitational Atoms:

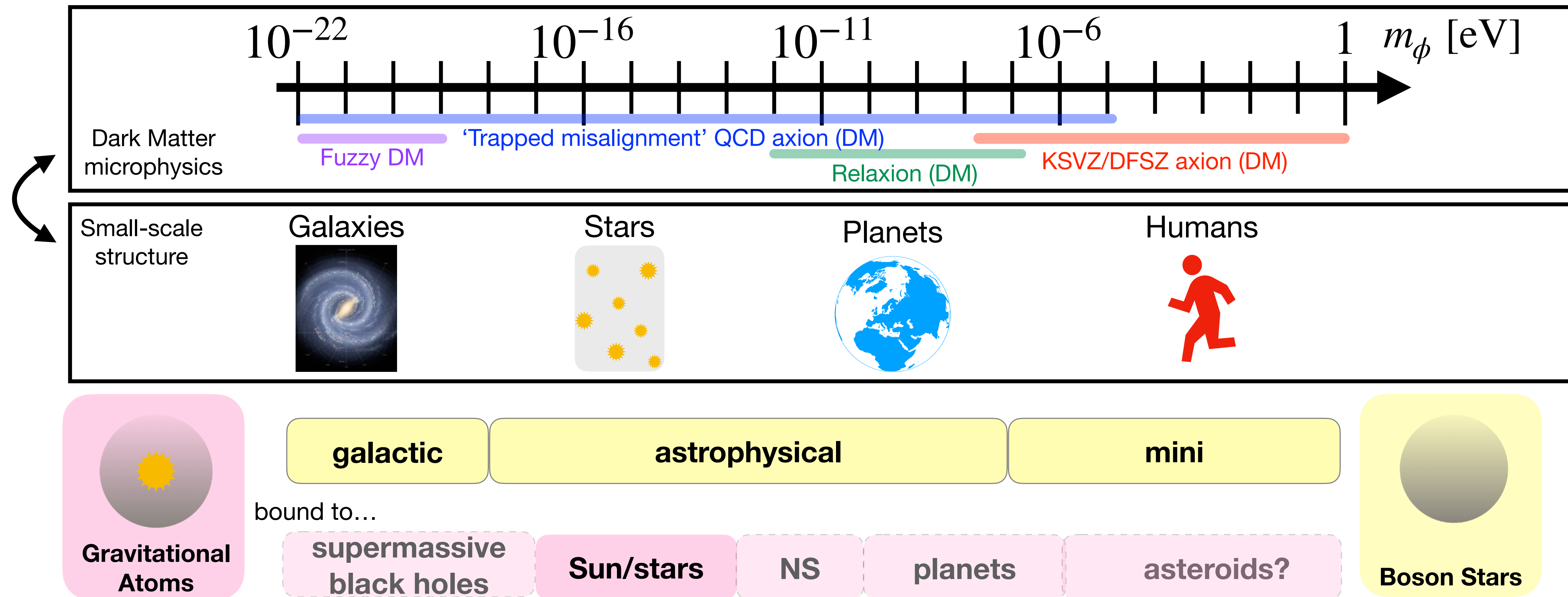
Gravimeters: Hu++ (1912.01900)

Atomic clocks: Kouvaris++ (2106.06023)

Neutrinos: Gherghetta, Shkerin (2305.06441)

Radio telescopes: Gong++ (2308.08477)

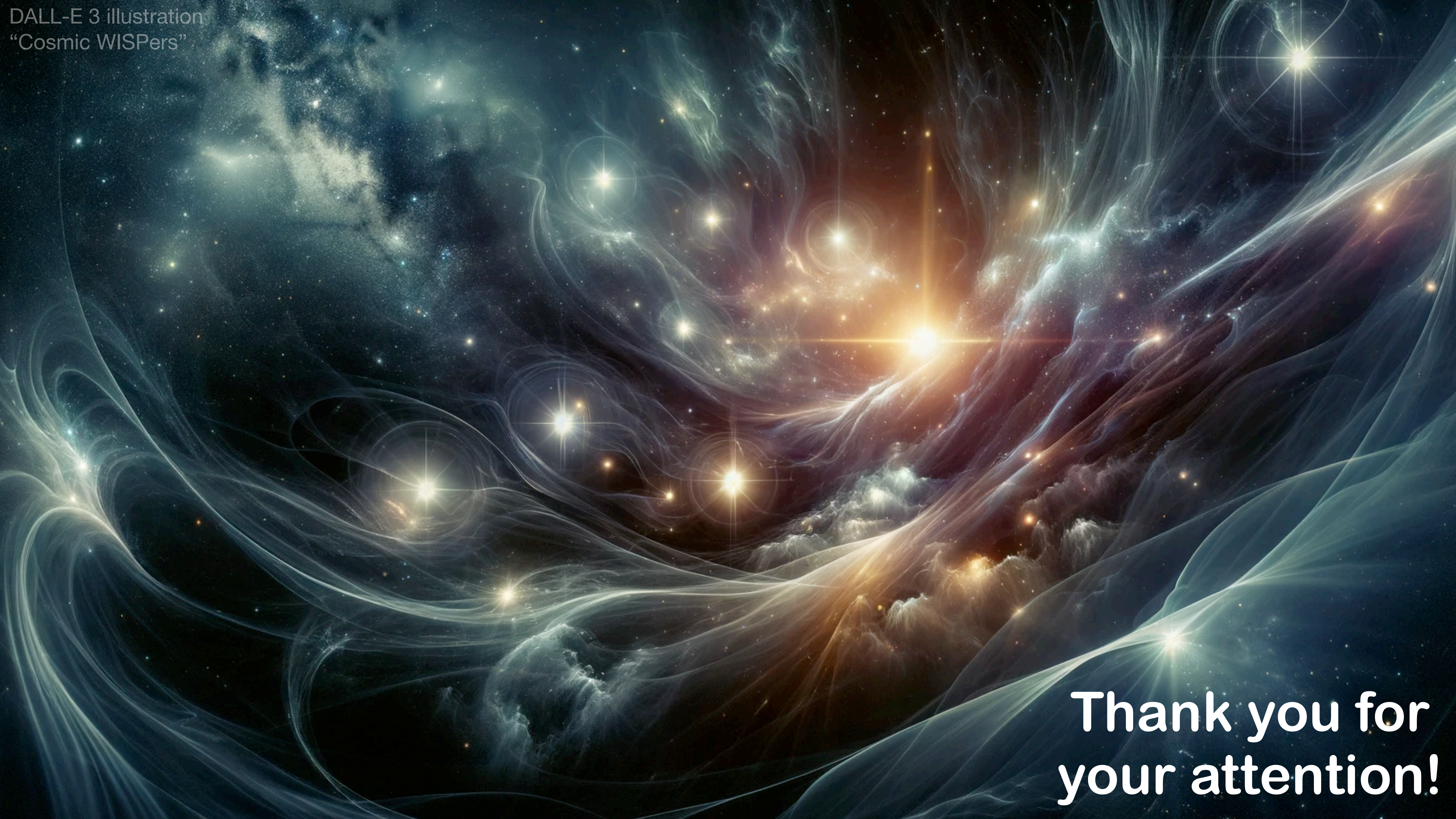
Conclusion: New Paths to Discovery



We are exploring the **unavoidable and unexplored consequences** of the theory, to elucidate the nature of ULDM and find **new paths to discovery**

- Big open questions remain:**
- Core-halo relation? mass growth rate? Mass distribution of boson stars?
 - Signals from GAs across the galaxy?

DALL-E 3 illustration
"Cosmic WISPers"



**Thank you for
your attention!**