

Neutrino cosmology

Maria Archidiacono



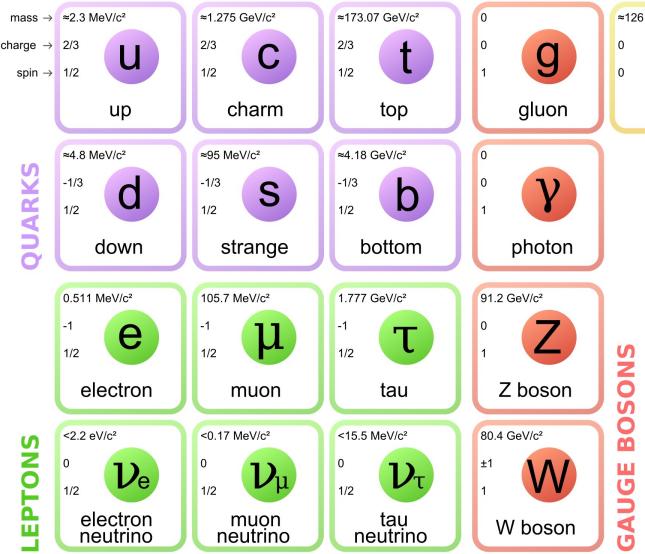
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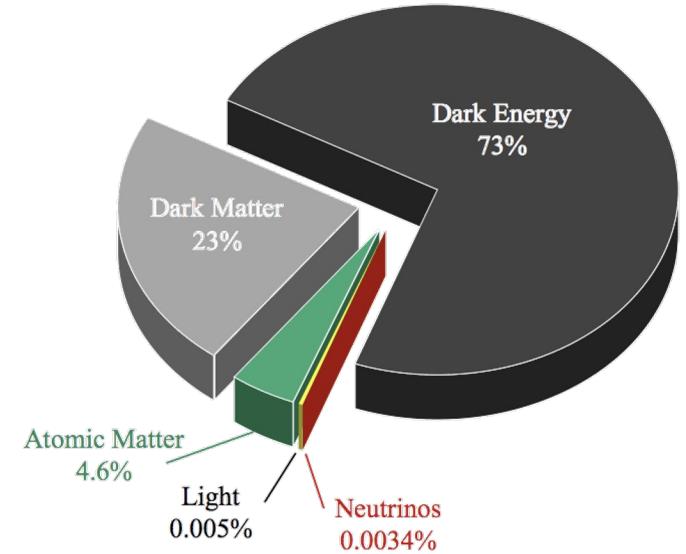
15th International Neutrino Summer School
Bologna, 7.6.2024

Why neutrino cosmology

Standard Model of particle physics

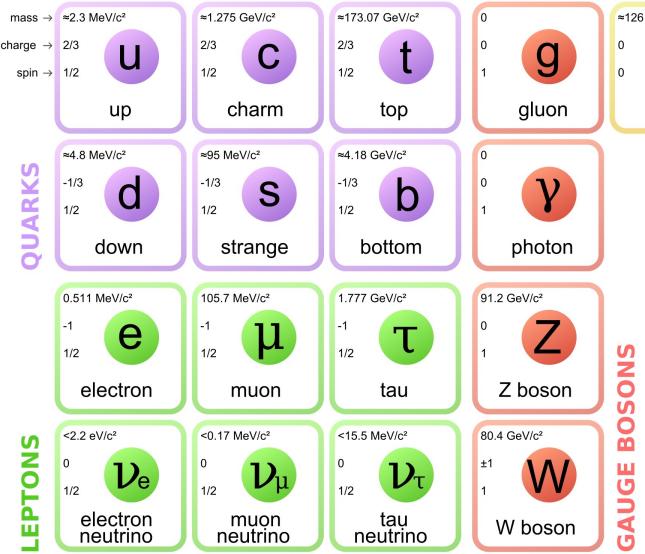


Λ CDM model of cosmology

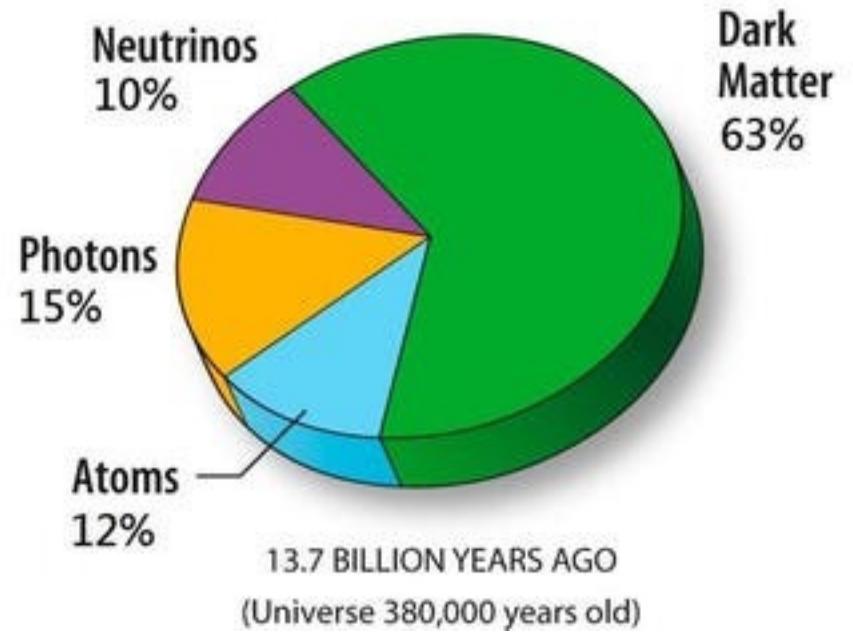


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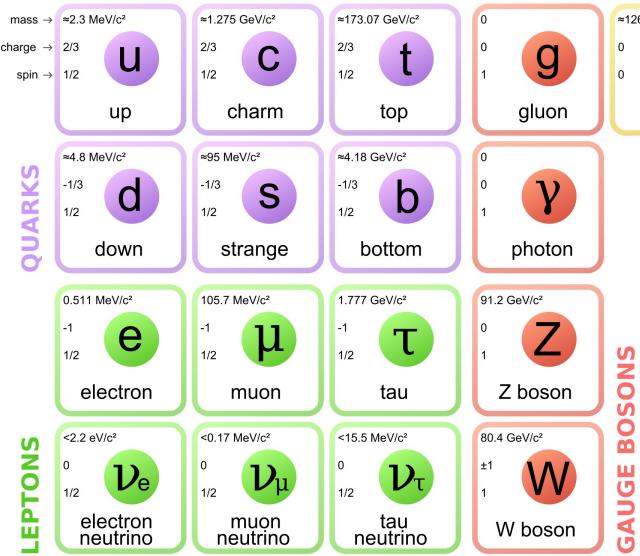


Λ CDM model of cosmology



Why neutrino cosmology

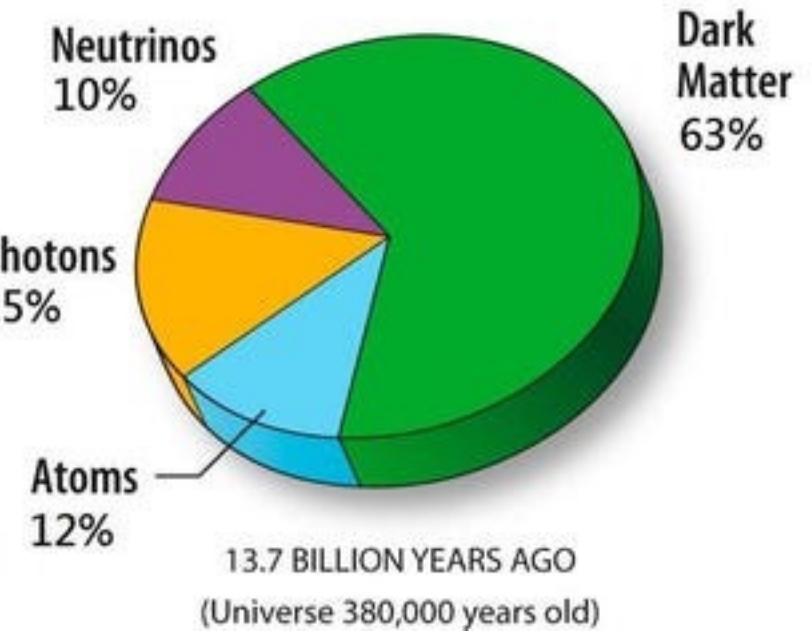
Standard Model of particle physics



Open questions:

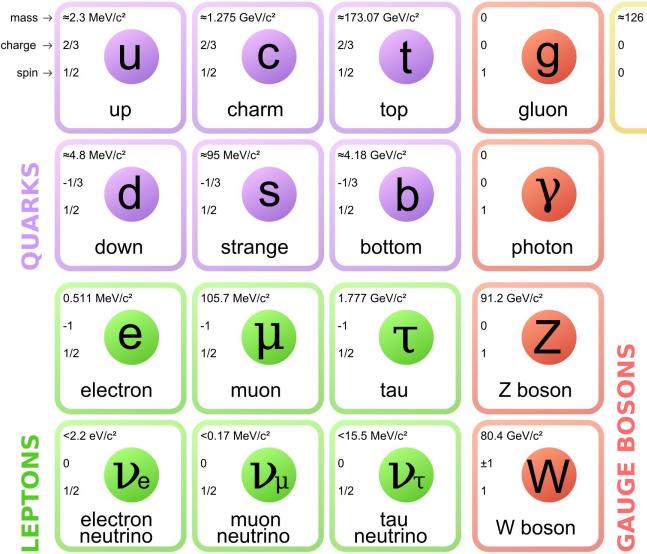
- Neutrino masses
- Dark matter
- Neutrino number
- Dark energy

Λ CDM model of cosmology

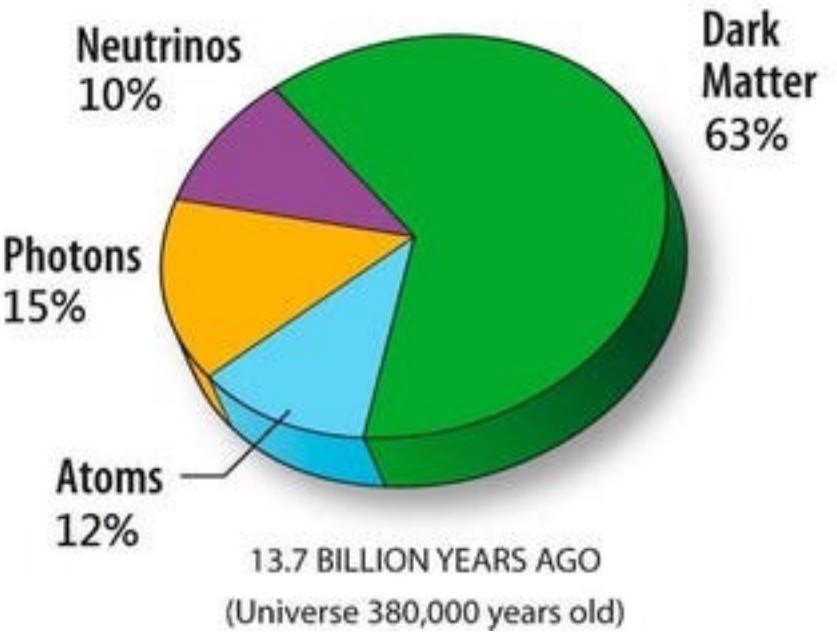


Why neutrino cosmology

Standard Model of particle physics



Λ CDM model of cosmology



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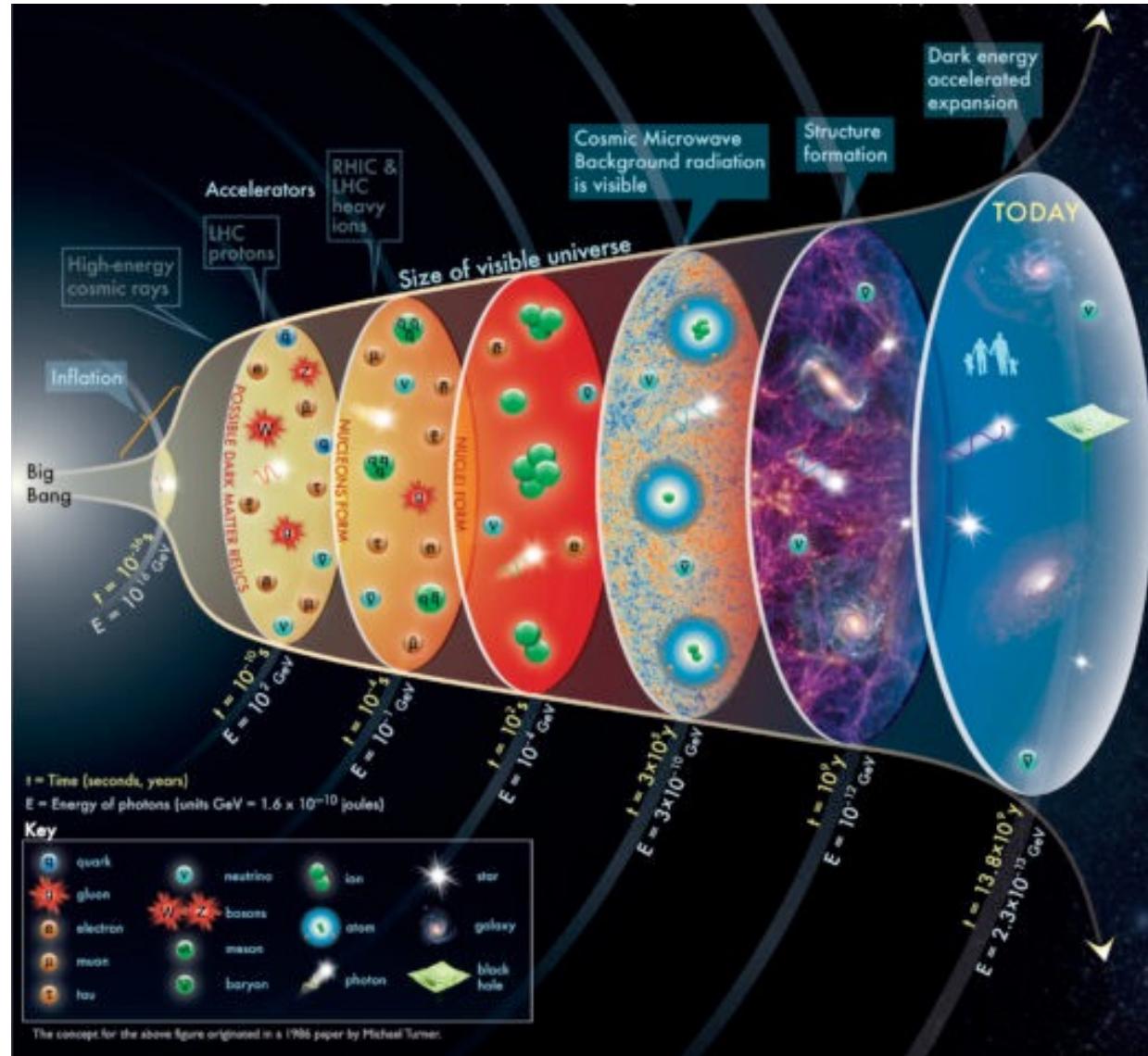
Plan

- A short cosmic history
 - Early Universe: A tale of two backgrounds
 - Cosmic Microwave Background (CMB)
 - Cosmic Neutrino Background (CvB)
 - Late Universe:
 - Large Scale Structure
- Indirect detection of the CvB: “The number of neutrinos...and other light relics”
- Detecting the neutrino mass in the CvB
- Non-standard neutrinos
 - The cosmological neutrino mass problem
 - Sterile neutrinos and new interactions

A short cosmic history

The Universe is expanding from a hot dense and homogeneous state.

Particles decouple from the thermal bath when $\Gamma < H$ leaving behind relics that can survive to present time, and sometimes be detected.

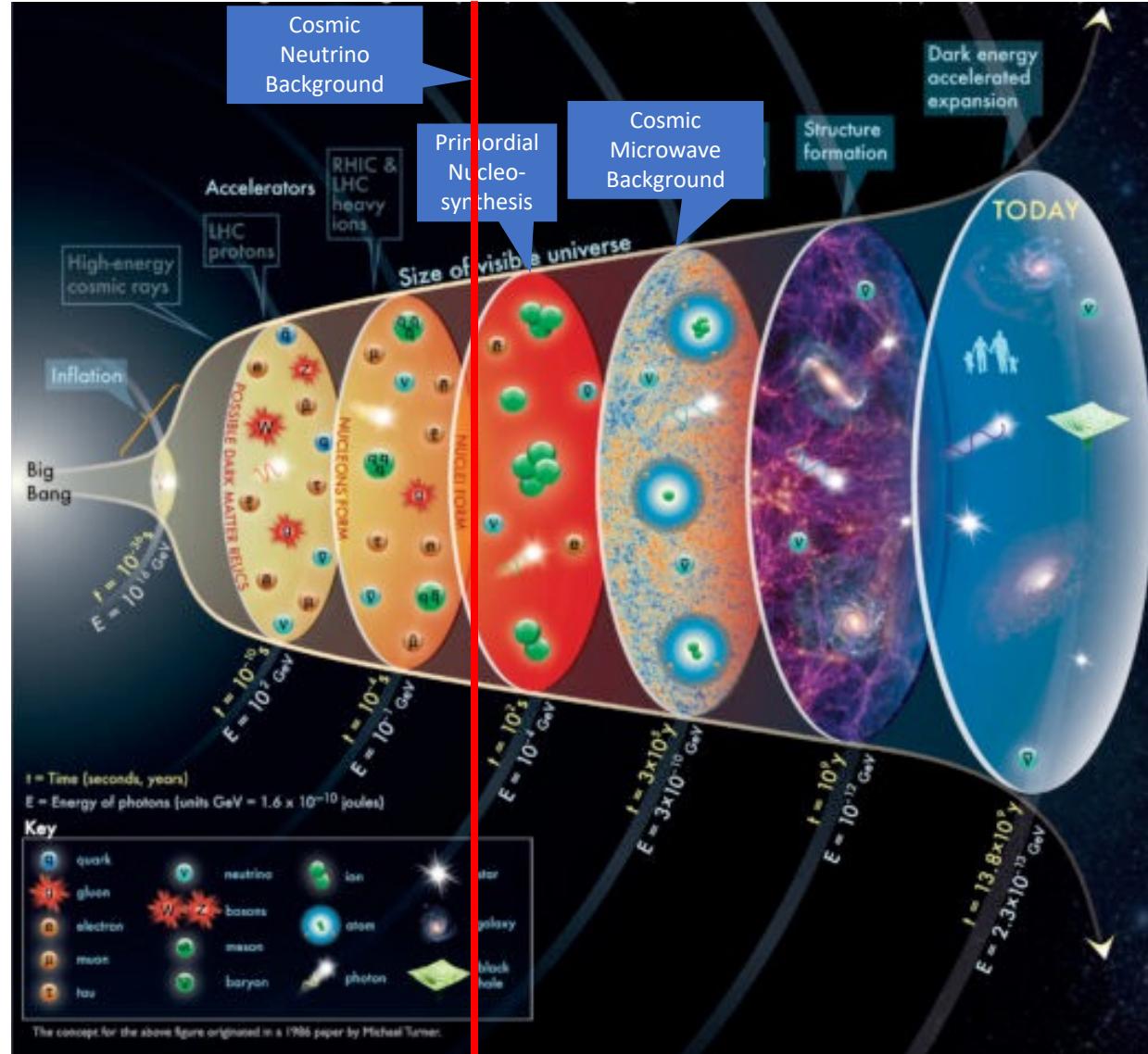


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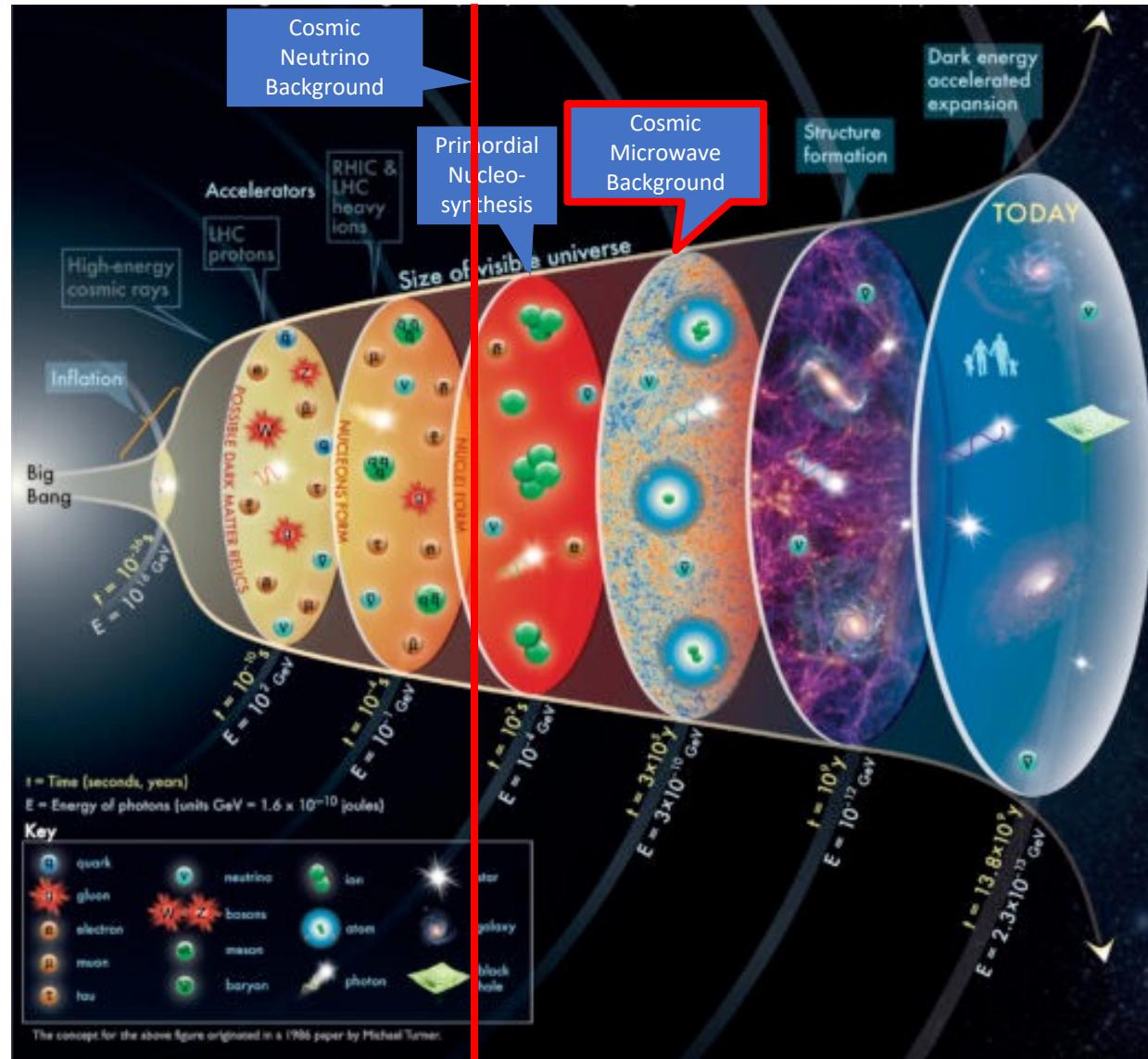


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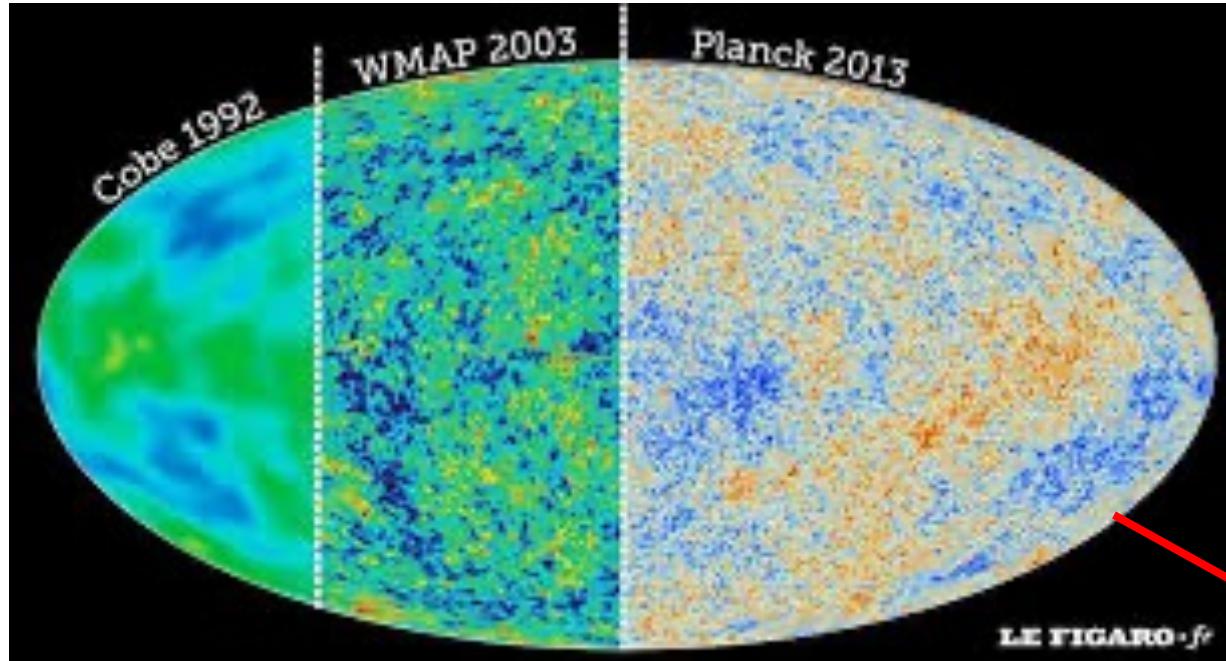
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Cosmic Microwave Background (CMB)

Thermal photon background formed around 300000 years after the big bang ($z \sim 1100$), when the Universe had a temperature of approximately 0.2 eV (~ 3000 K). The Universe becomes neutral and transparent to photons.



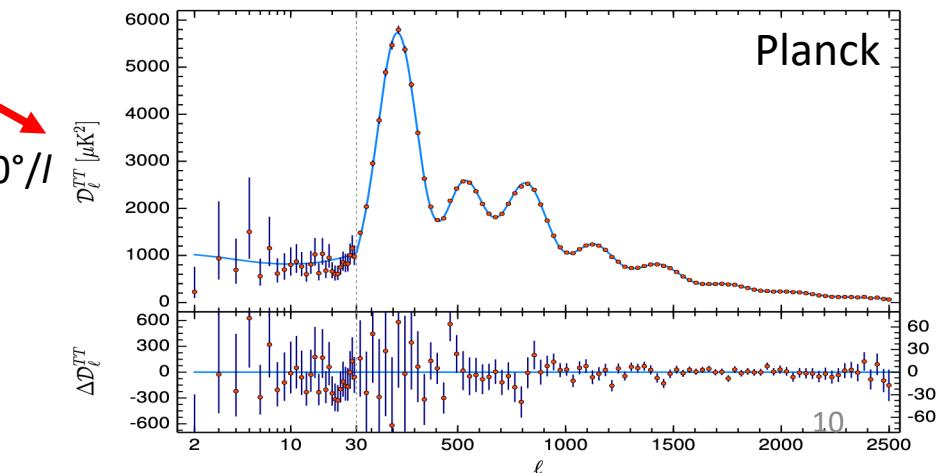
1978 Penzias and Wilson
Discovery of a ~ 2.7 K background



2006 Mather and Smoot (COBE)
Measurement of 10^{-5} fluctuations



2019 Peebles
For predicting these fluctuations in 1970
(among other things)

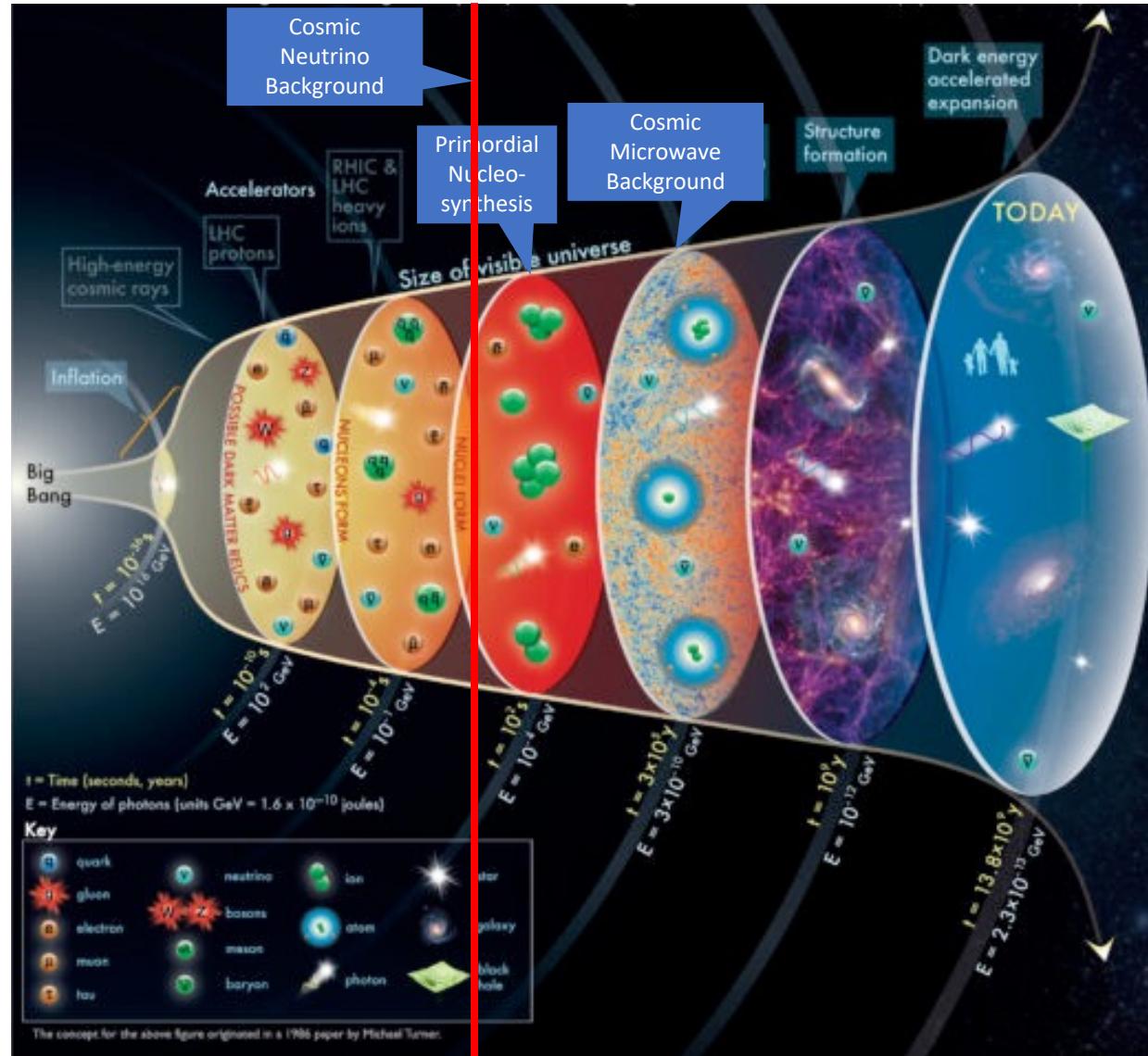


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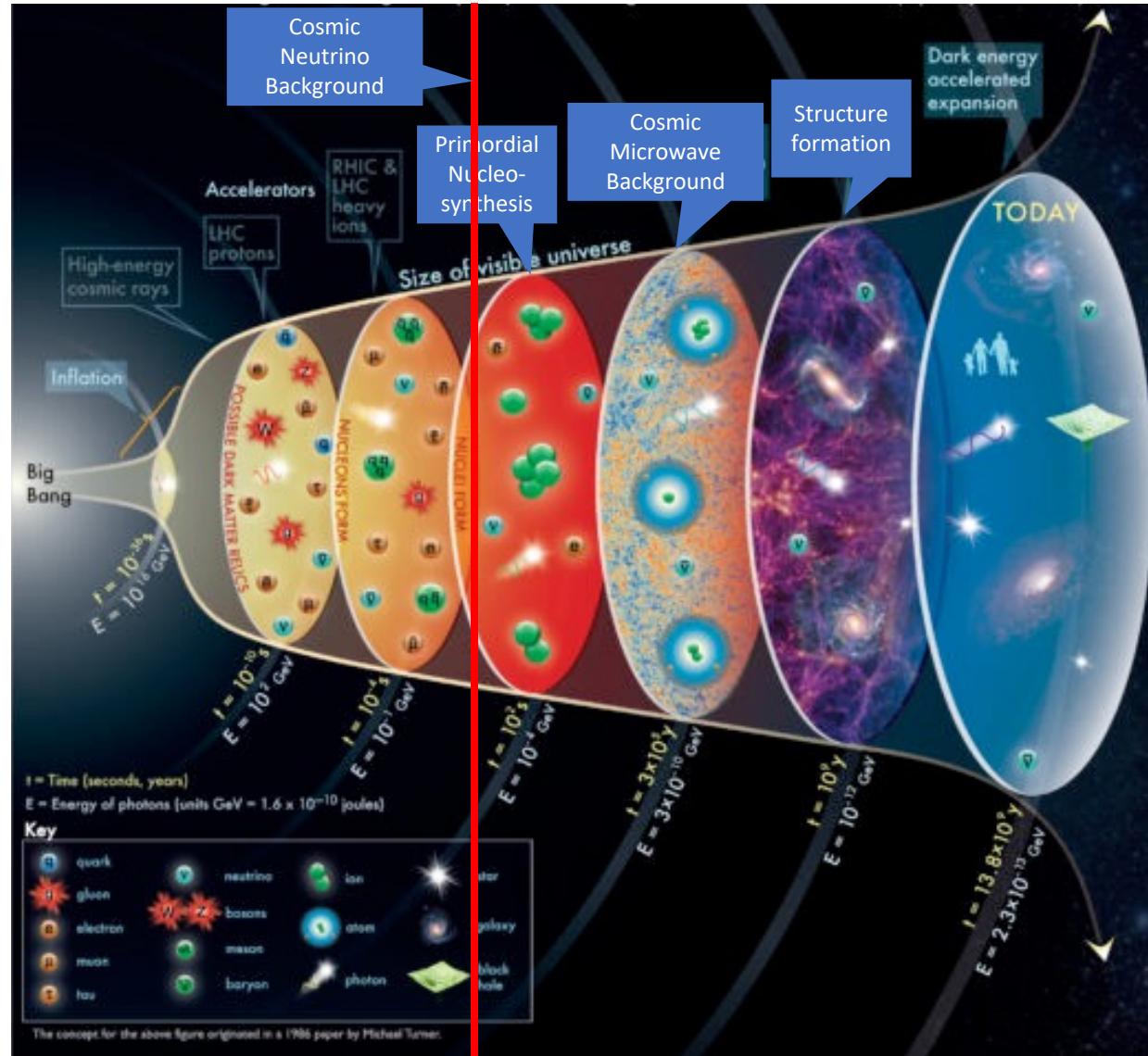


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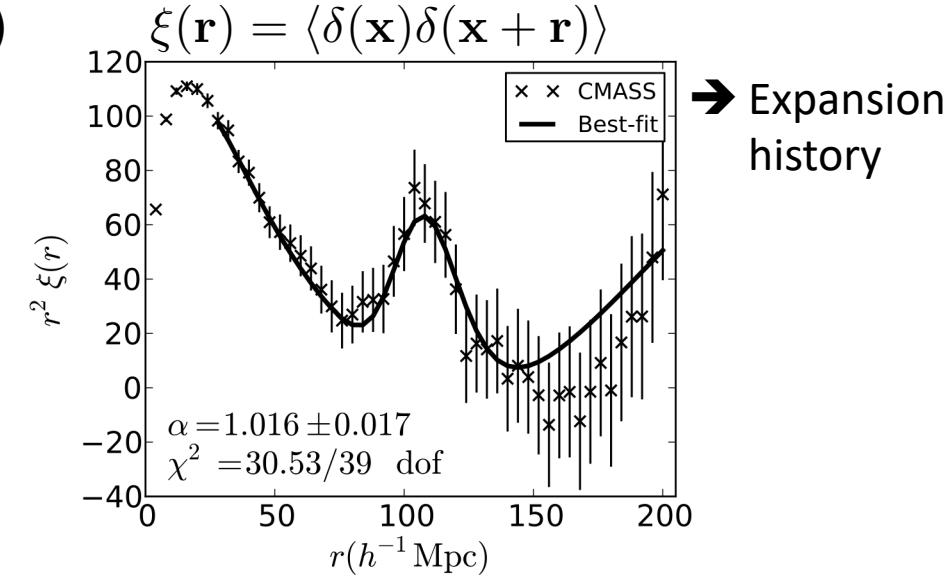
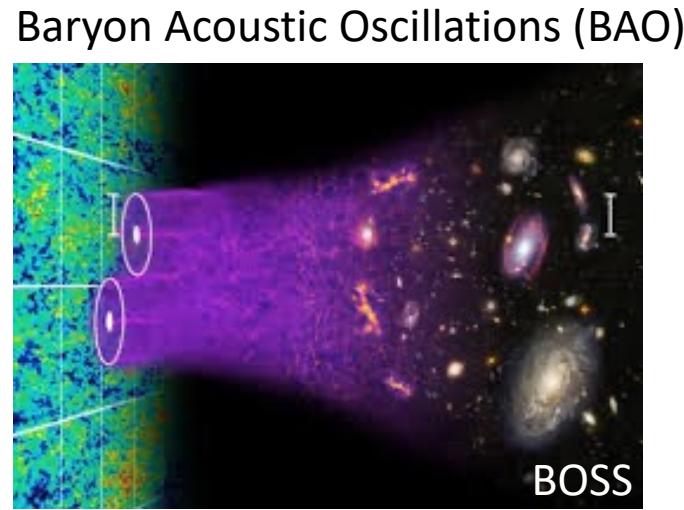
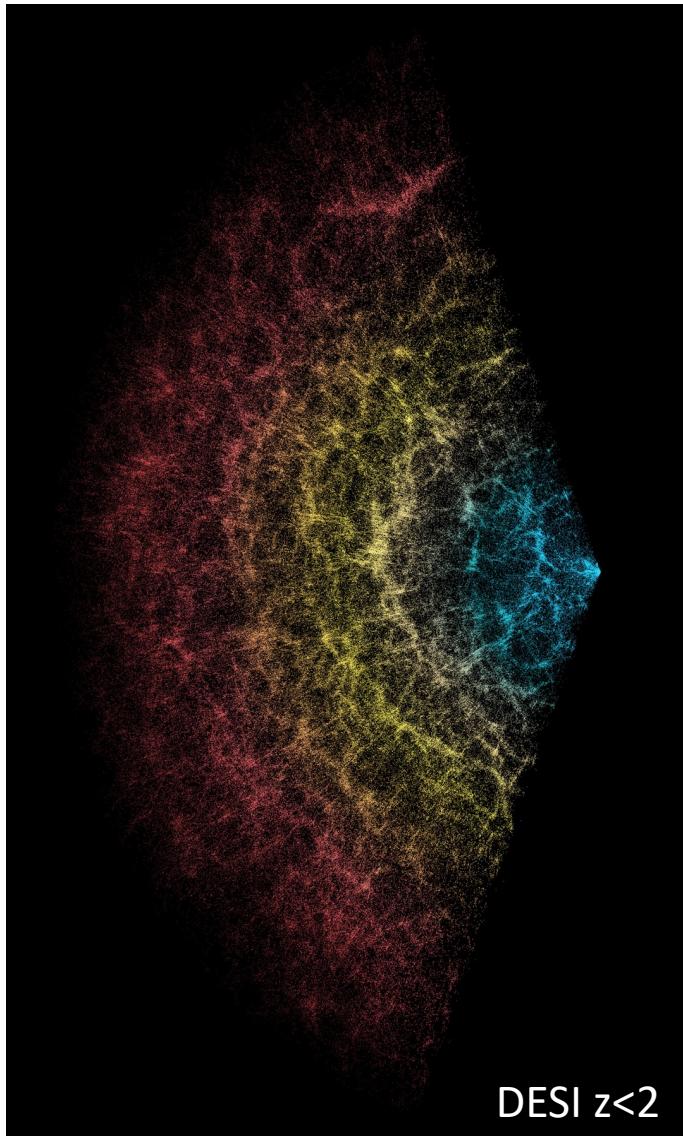
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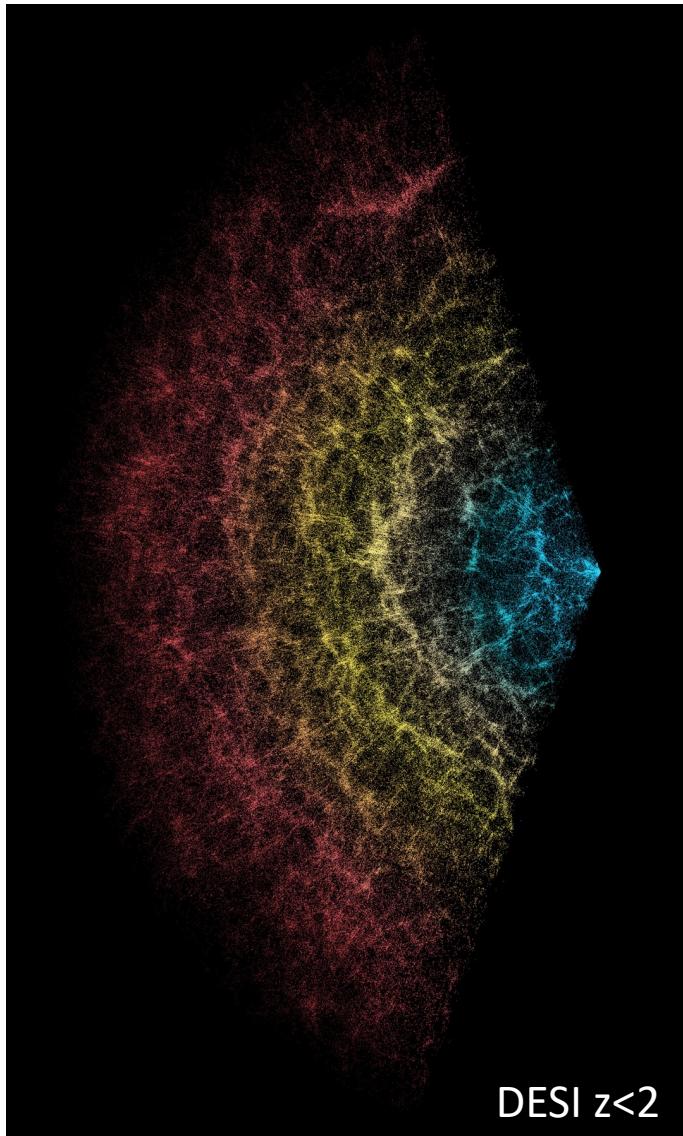
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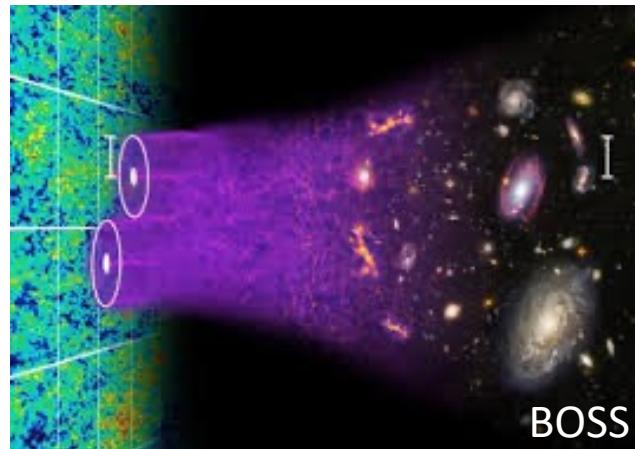
Large scale structure (LSS)



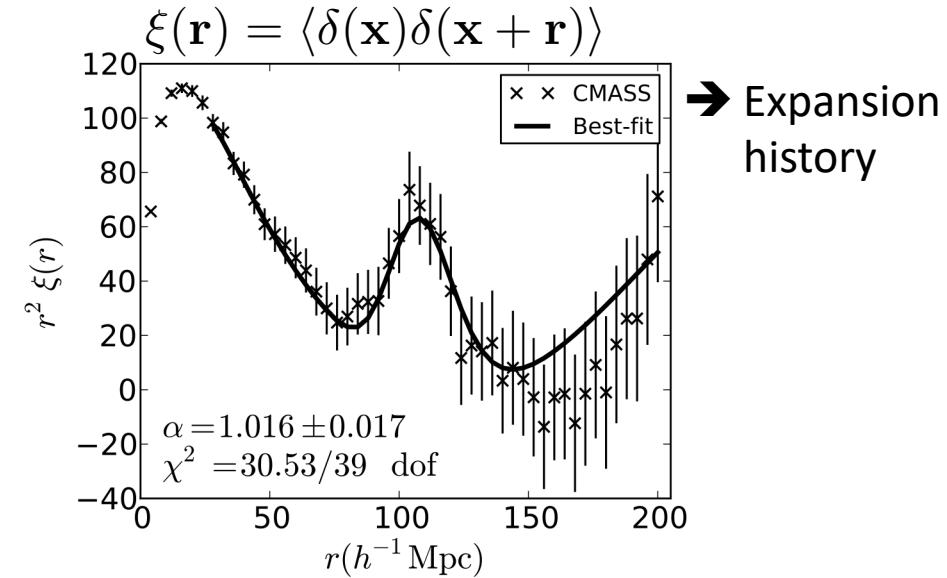
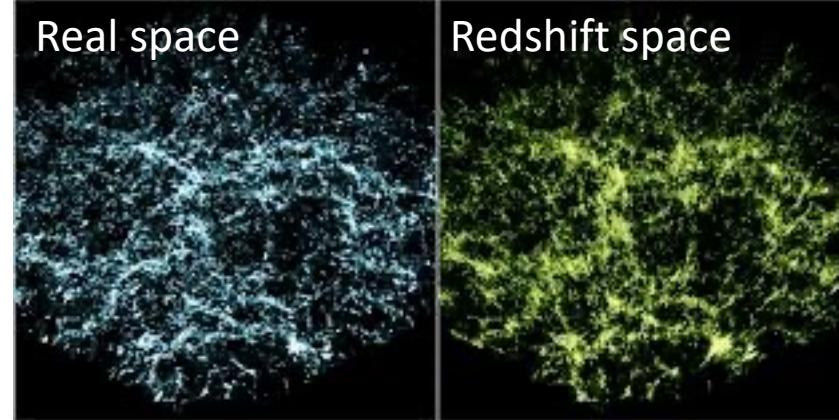
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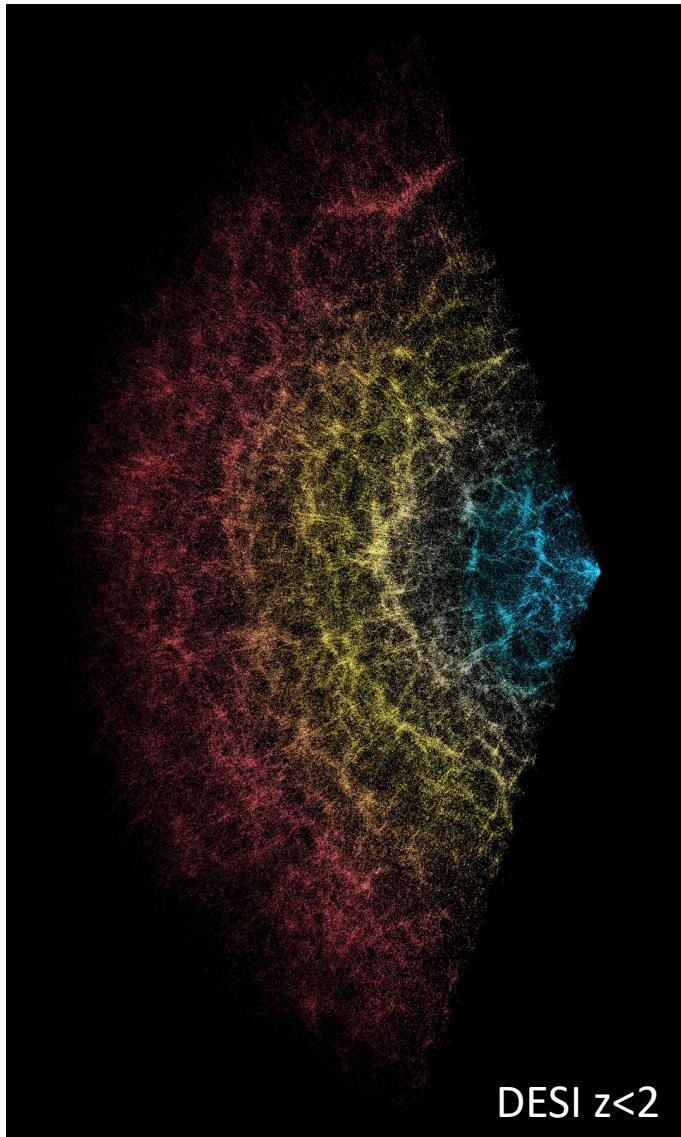
Baryon Acoustic Oscillations (BAO)



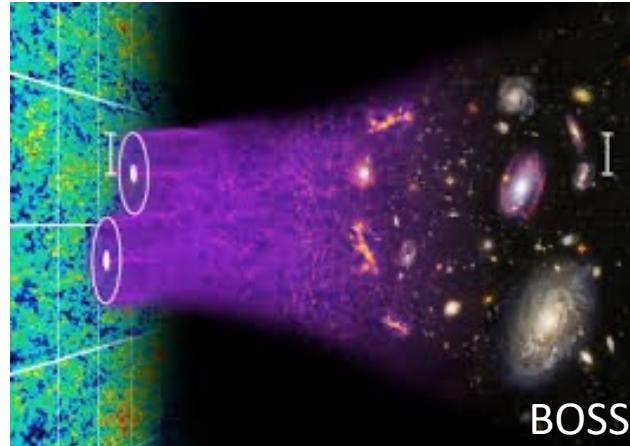
Redshift space distortions



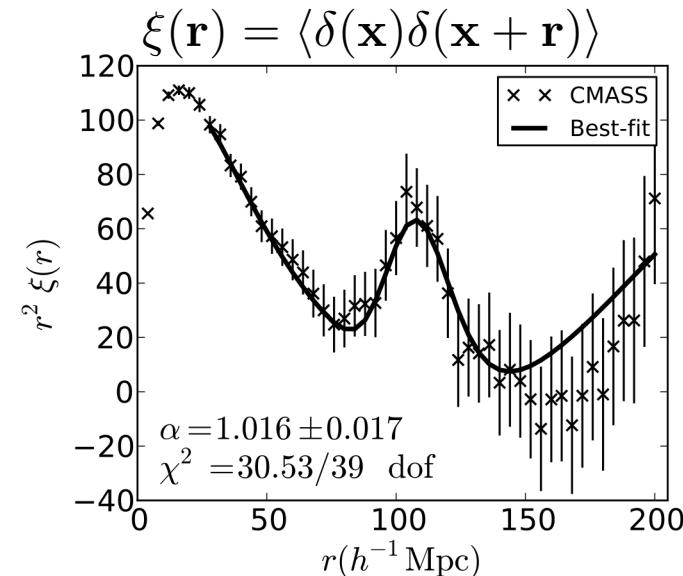
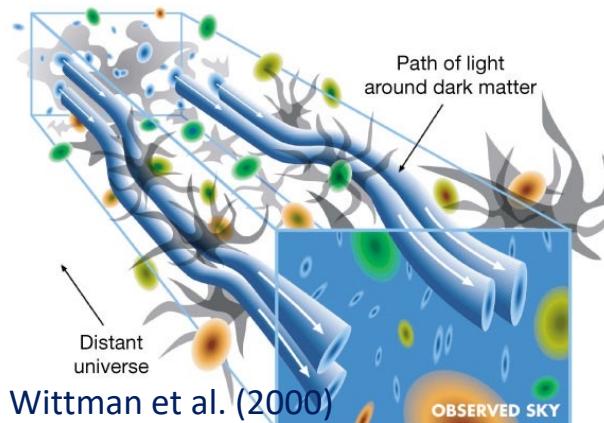
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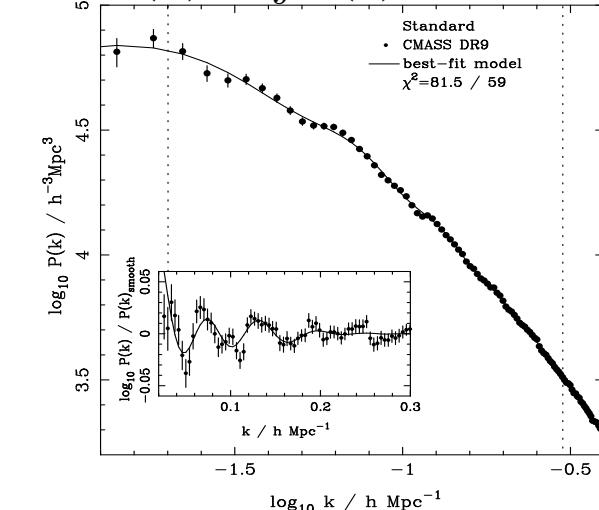
Baryon Acoustic Oscillations (BAO)



With imaging
Weak gravitational lensing



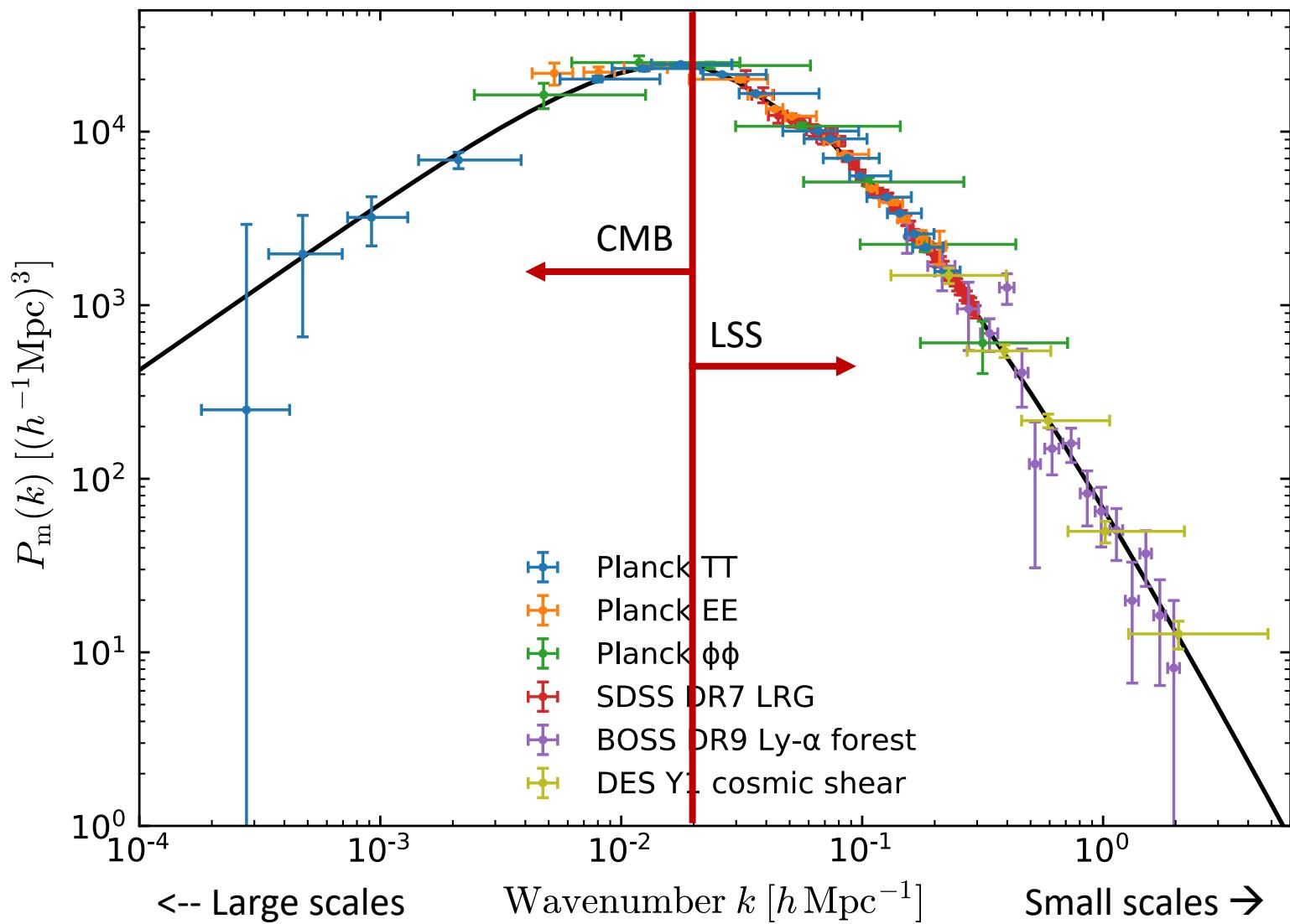
$$P(k) = \int \xi(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}} d^3x$$



→ Expansion history

→ Expansion history and growth of structures

Different observables at different scales



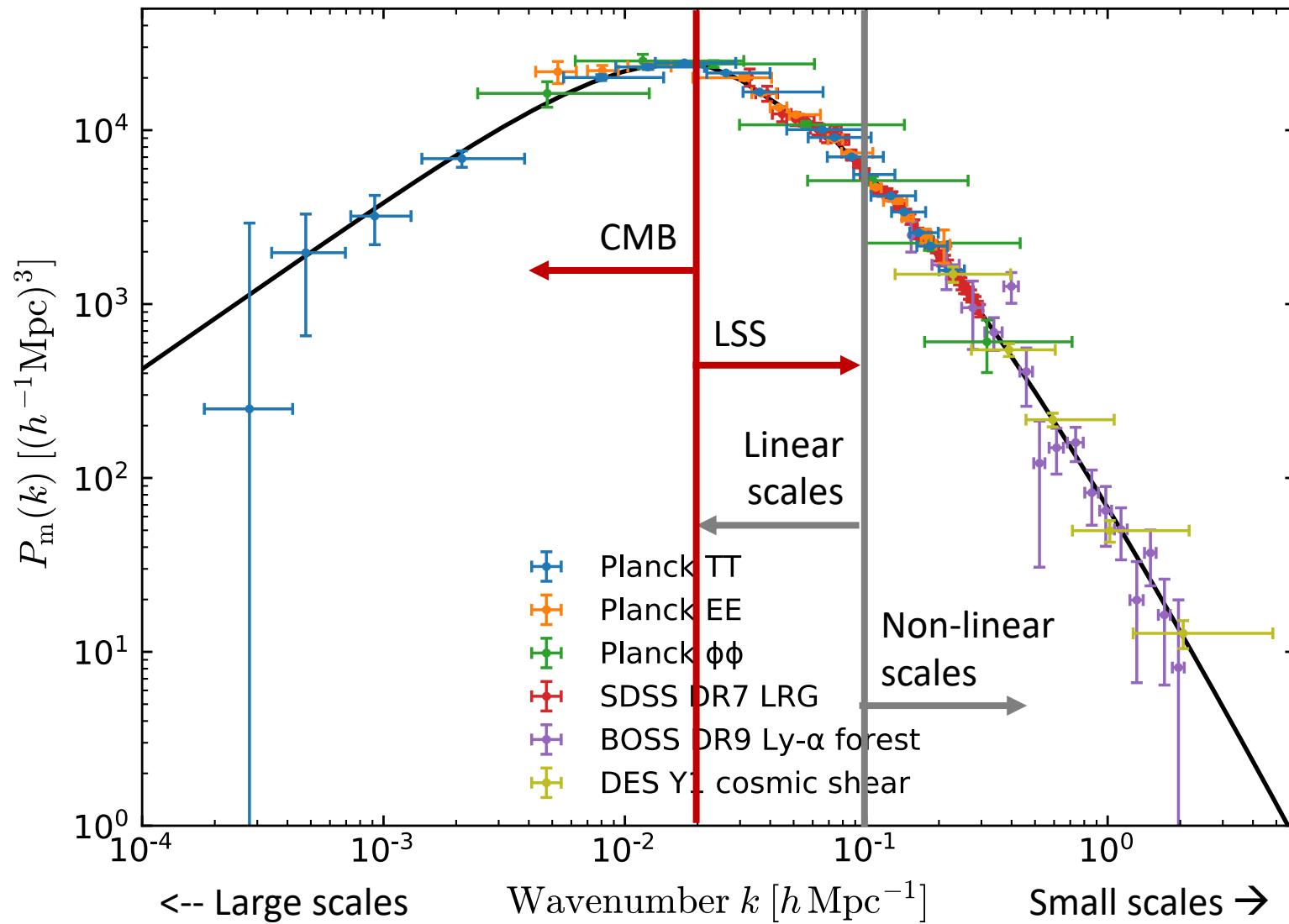
Different observables at different scales

Linear scales:
predictions of
observables from
Einstein-Boltzmann
solvers (CLASS or
CAMB)

CPU time for one
model: << 1 sec

Non-linear scales:
predictions of
observables from N-
body simulations

CPU time for one
model: ~ khrs-Mhrs

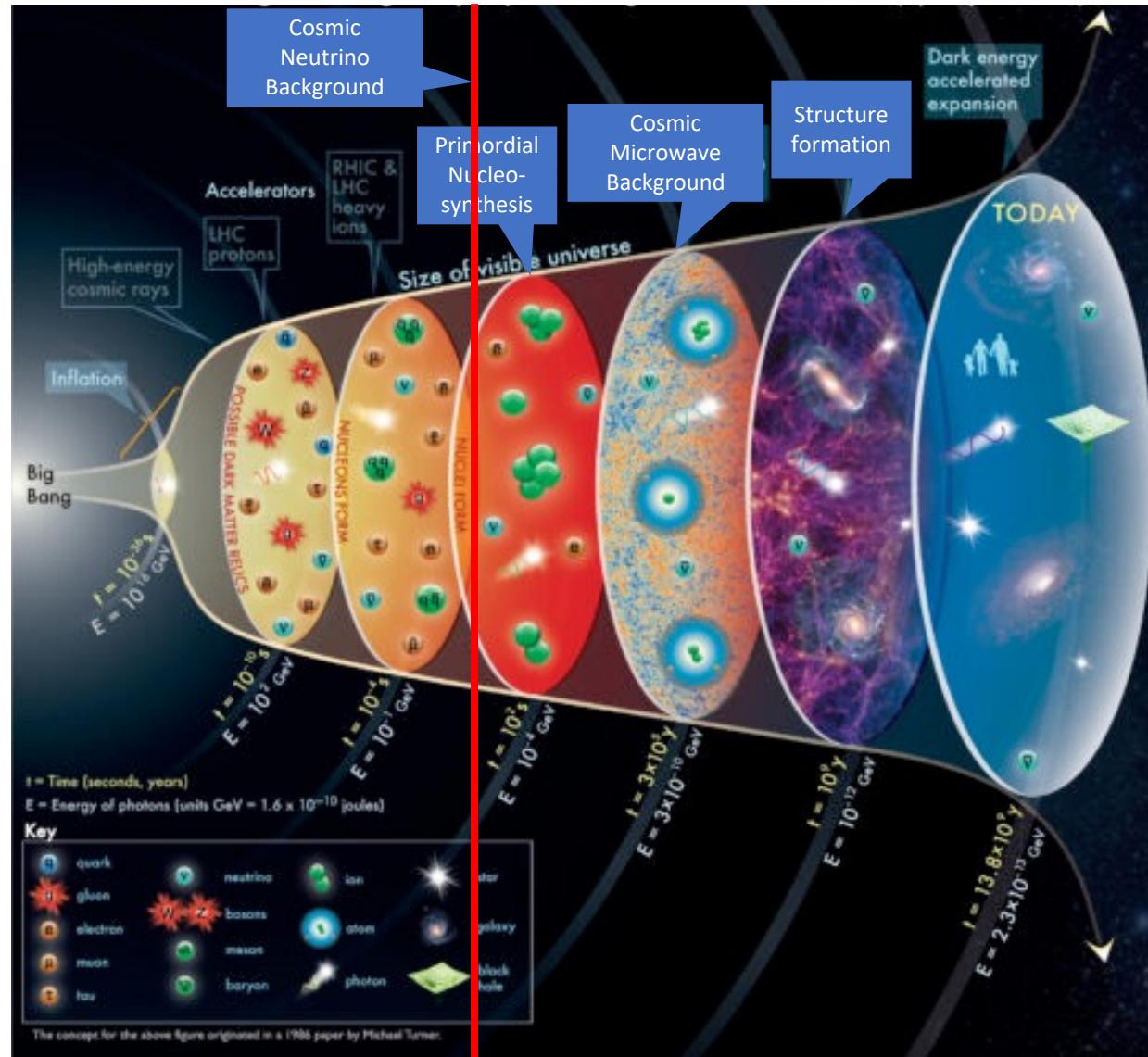


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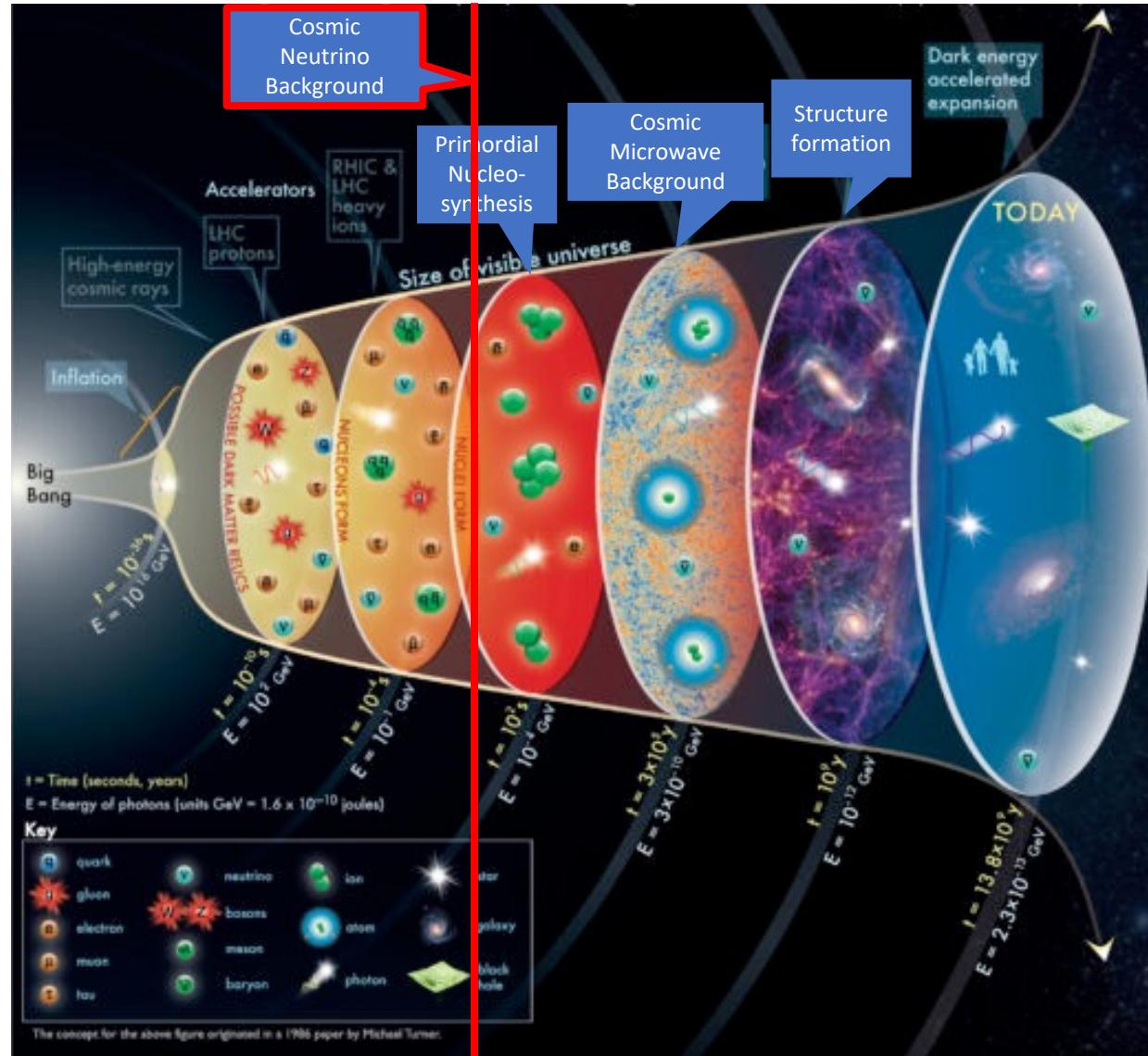


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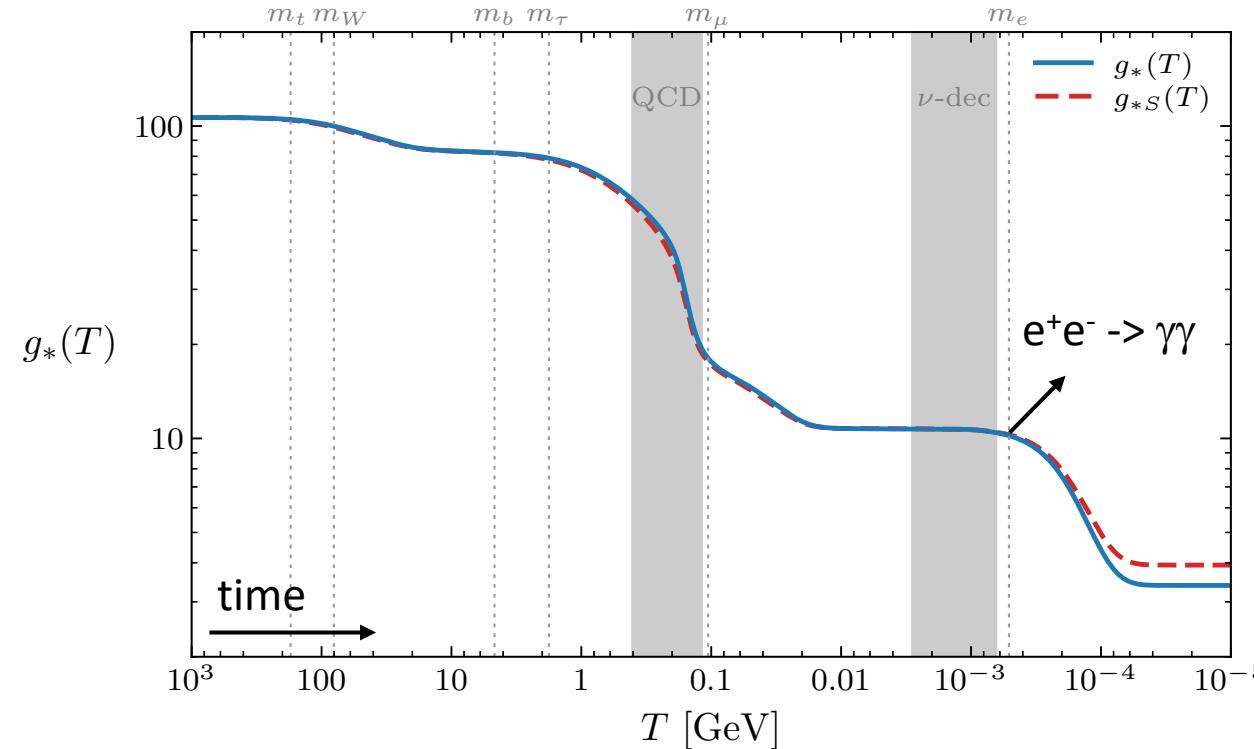
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The cosmic neutrino background (CvB)

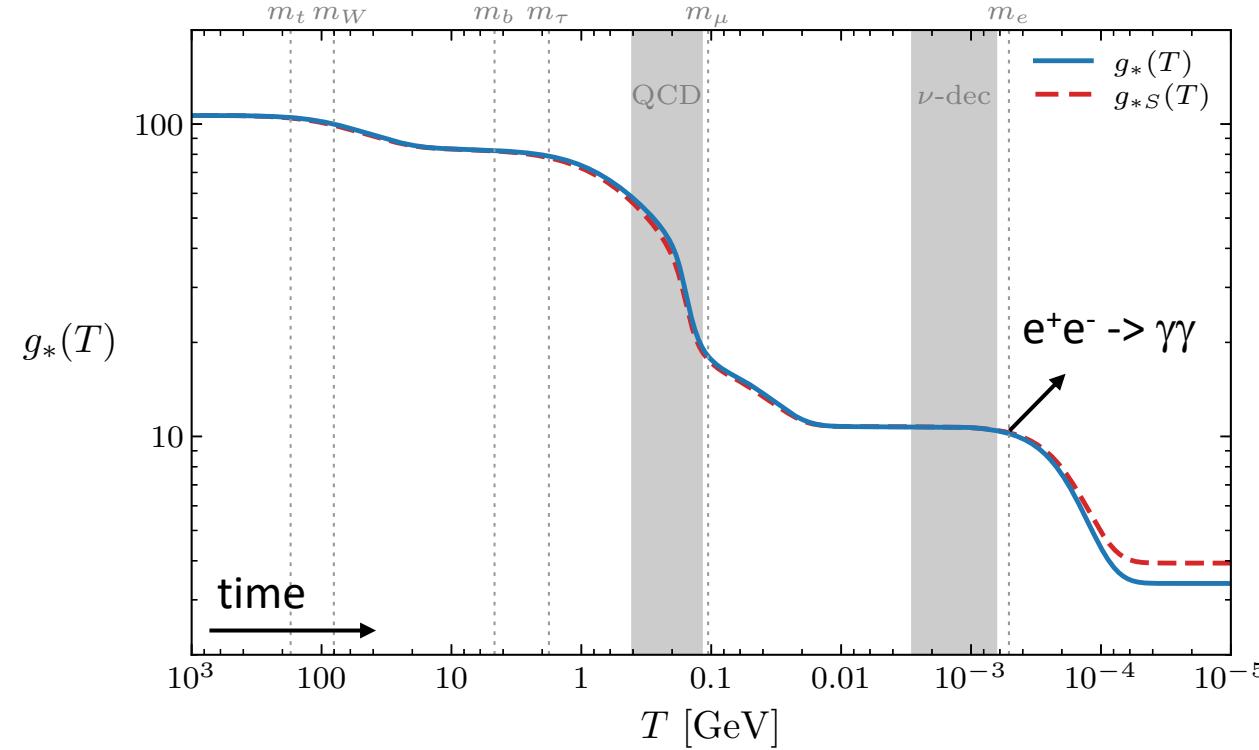


The neutrino background formed 1 sec after the big bang, at $T \sim 1$ MeV, when the weak interactions become inefficient.

Applying entropy density conservation

$$T_{C\nu B} = \left(\frac{4}{11}\right)^{1/3} T_{\gamma, \text{CMB}}$$

The cosmic neutrino background (CvB)



Today:

$$T_{\text{C}\nu\text{B}} \sim 1.9 \text{ K}$$

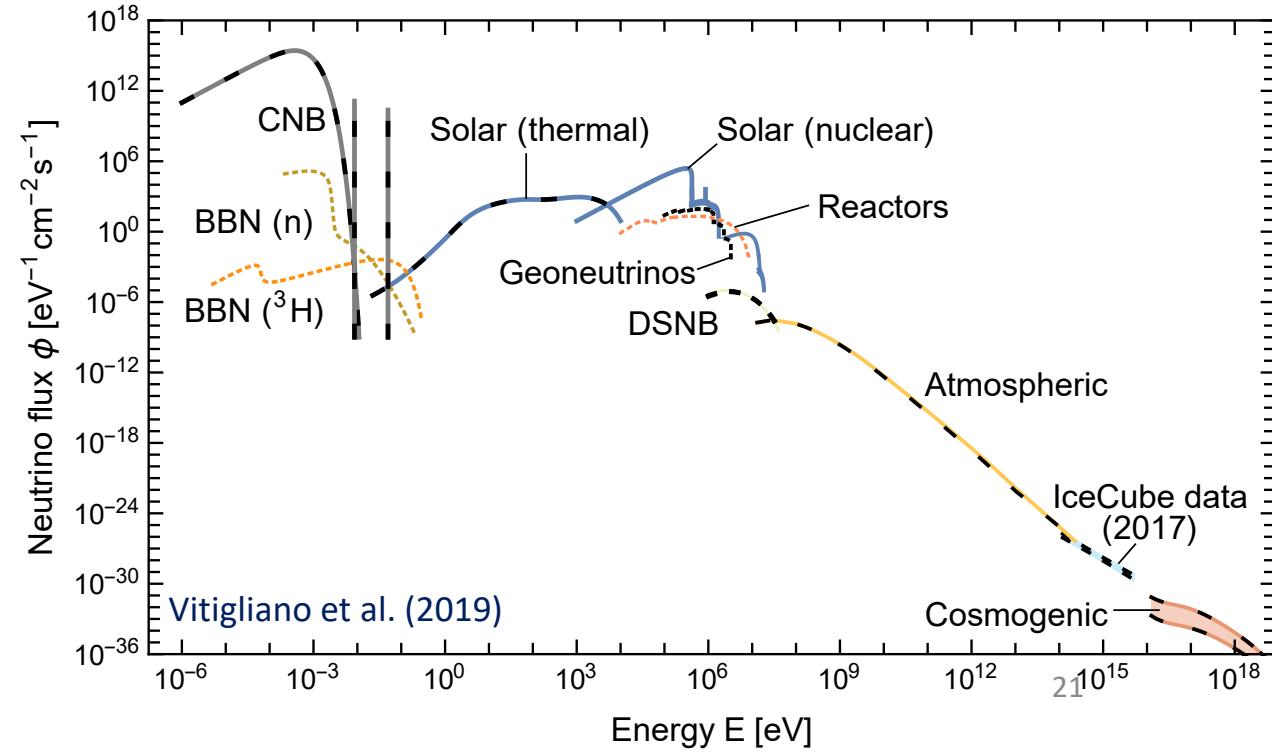
$$n_{\text{C}\nu\text{B}} \sim 110 \text{ cm}^{-3}$$

- ➔ Direct detection of the CvB not in the near future
- ➔ Footprints of the CvB in CMB and LSS

The neutrino background formed 1 sec after the big bang, at $T \sim 1$ MeV, when the weak interactions become inefficient.

Applying entropy density conservation

$$T_{\text{C}\nu\text{B}} = \left(\frac{4}{11}\right)^{1/3} T_{\gamma, \text{CMB}}$$



Indirect detection of the CvB
The number of neutrinos
...and other light relics

The N_{eff} parameter

The expansion rate of the Universe depends on its ingredients

$$H^2(a(t)) = H_0^2 (\Omega_{\text{rad}} a^{-4} + \Omega_{\text{m}} a^{-3} + \Omega_{\Lambda})$$

$$\Omega_i = \frac{\rho_{i,0}}{\rho_{\text{crit}}}$$

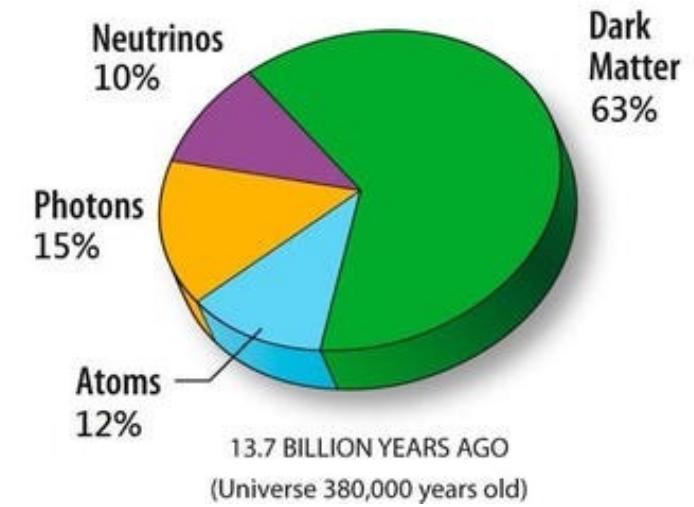
$$\rho_{\text{rad}} = \rho_{\gamma, \text{CMB}} + \rho_{\nu, \text{C}\nu\text{B}} + \rho_{\text{DR}} = \rho_{\gamma, \text{CMB}} \left[1 + N_{\text{eff}} \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \right]$$

DR = Dark Radiation, BSM light relics

$$N_{\text{eff}} = \frac{(\text{energy density of neutrinos} + \text{BSM light particles})}{(\text{energy density of one neutrino species})}$$

For the 3 neutrino families of the SM

$$N_{\text{eff}}^{\text{SM}} = 3.044 \pm 0.001 \text{ [Mangano et al. (2002), Froustey et al. (2020), Bennett et al. (2021), Drewes et al. (2024)]}$$



N_{eff} probes: BBN

Shortly after neutrino decoupling the weak interactions that kept neutrons and protons in statistical equilibrium freeze out.

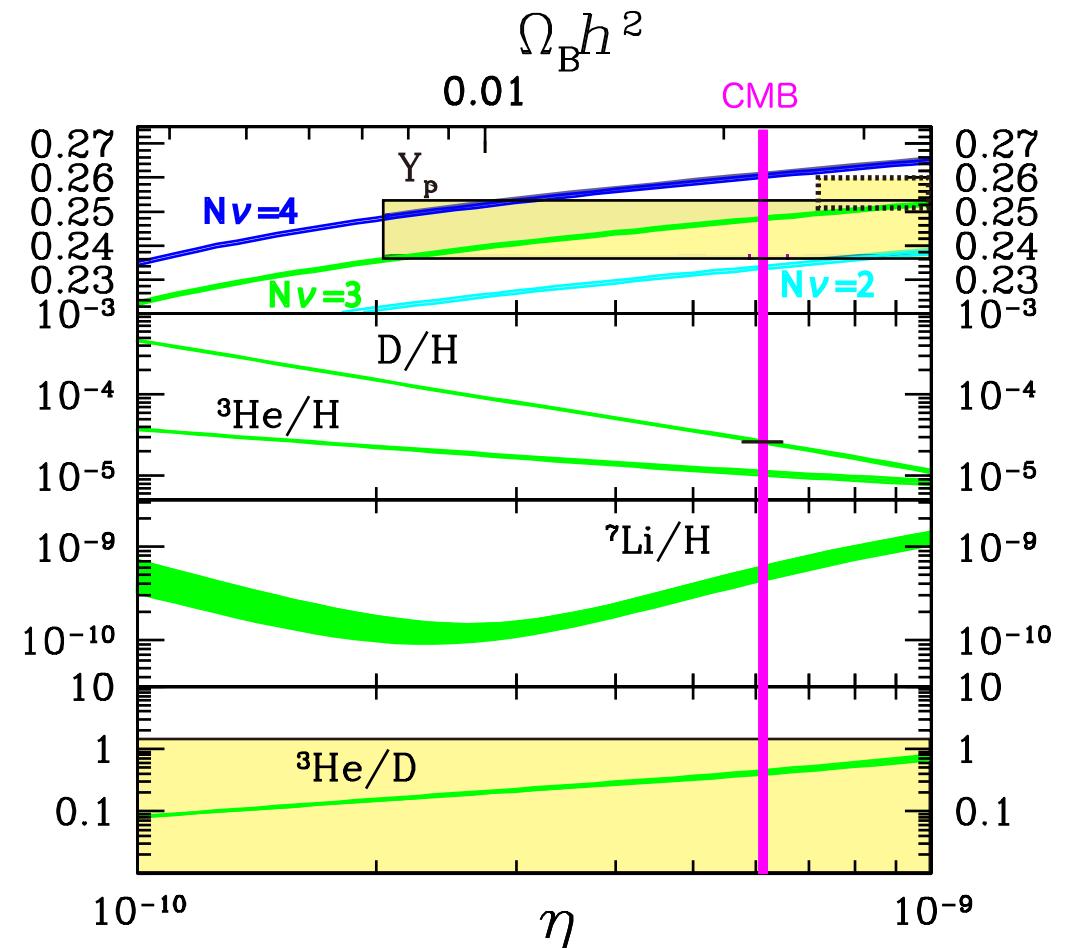
$$H = \Gamma|_{T=T_{\text{freeze}}} \quad T_{\text{freeze}} \approx 0.6 g_*^{1/6} \text{ MeV}$$

$$\left. \frac{n_n}{n_p} \right|_{T=T_{\text{freeze}}} \approx \exp \left(-\frac{(m_n - m_p)}{T_{\text{freeze}}} \right) \approx \frac{1}{6}$$

$$Y_P \approx \left. \frac{2n_n / n_p}{1 + n_n / n_p} \right|_{T \approx 0.2 \text{ MeV}} \propto f(g_*, \Omega_b h^2)$$

$$g_* \rightarrow g_* + \frac{7}{4} \Delta N_{\text{eff}}$$

$$\left| Y_P^{\text{theo}} - Y_P^{\text{obs}} \right|_{\Omega_b} \rightarrow \Delta N_{\text{eff}}|_{\Omega_b}$$



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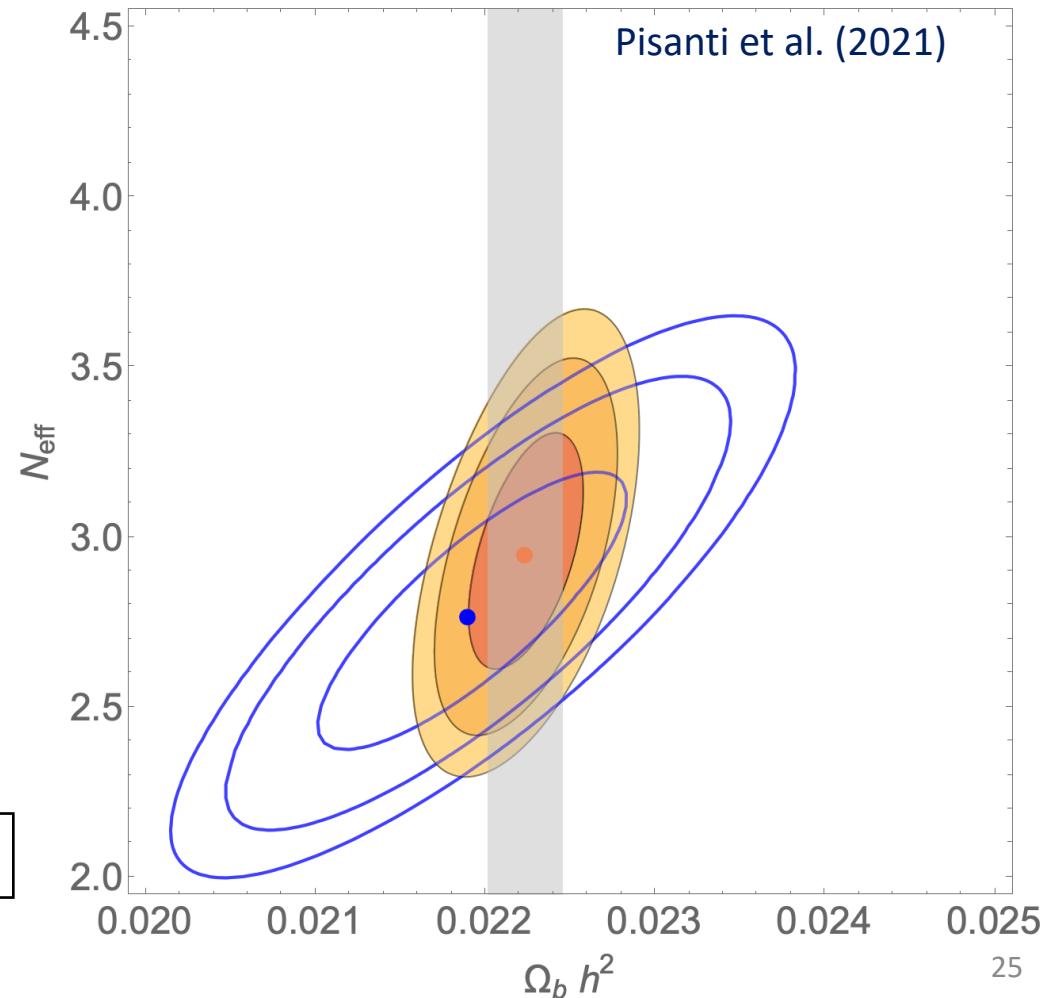
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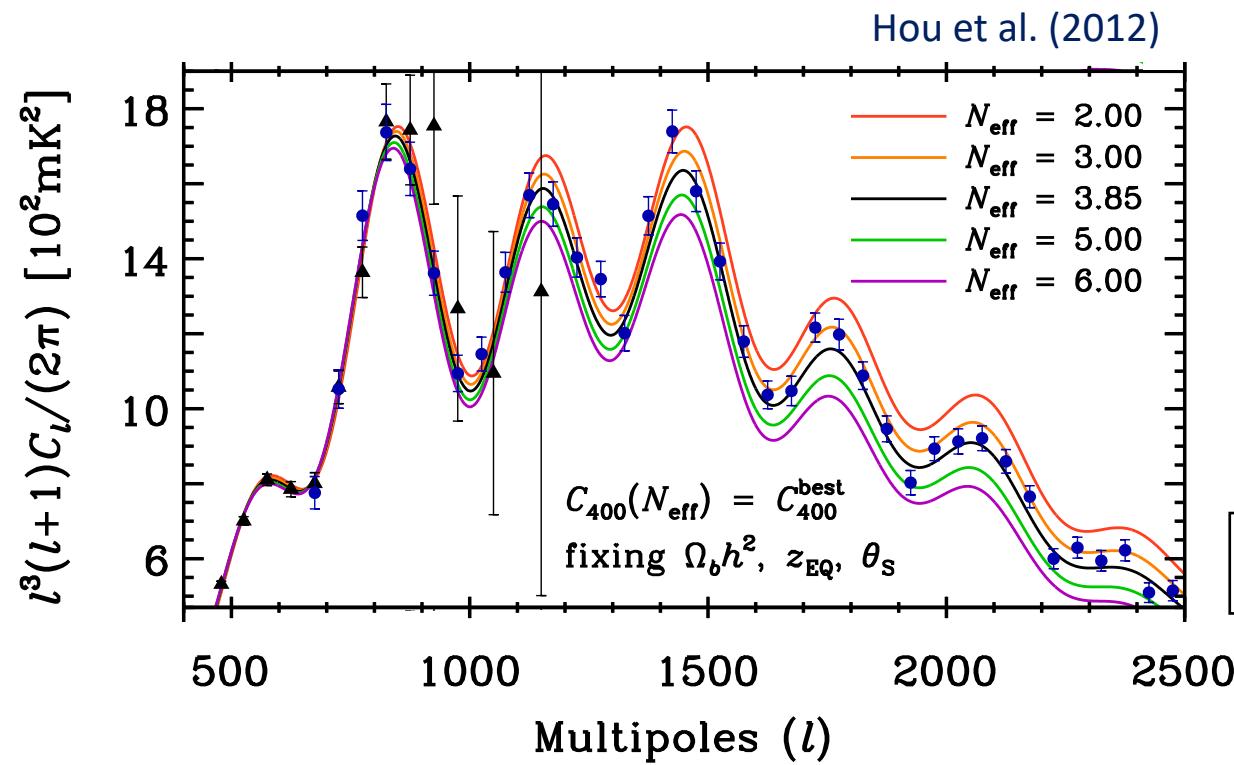
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$$N_{\text{eff}} = 2.78 \pm 0.28 \text{ (68% CL)}$$



N_{eff} probes: CMB



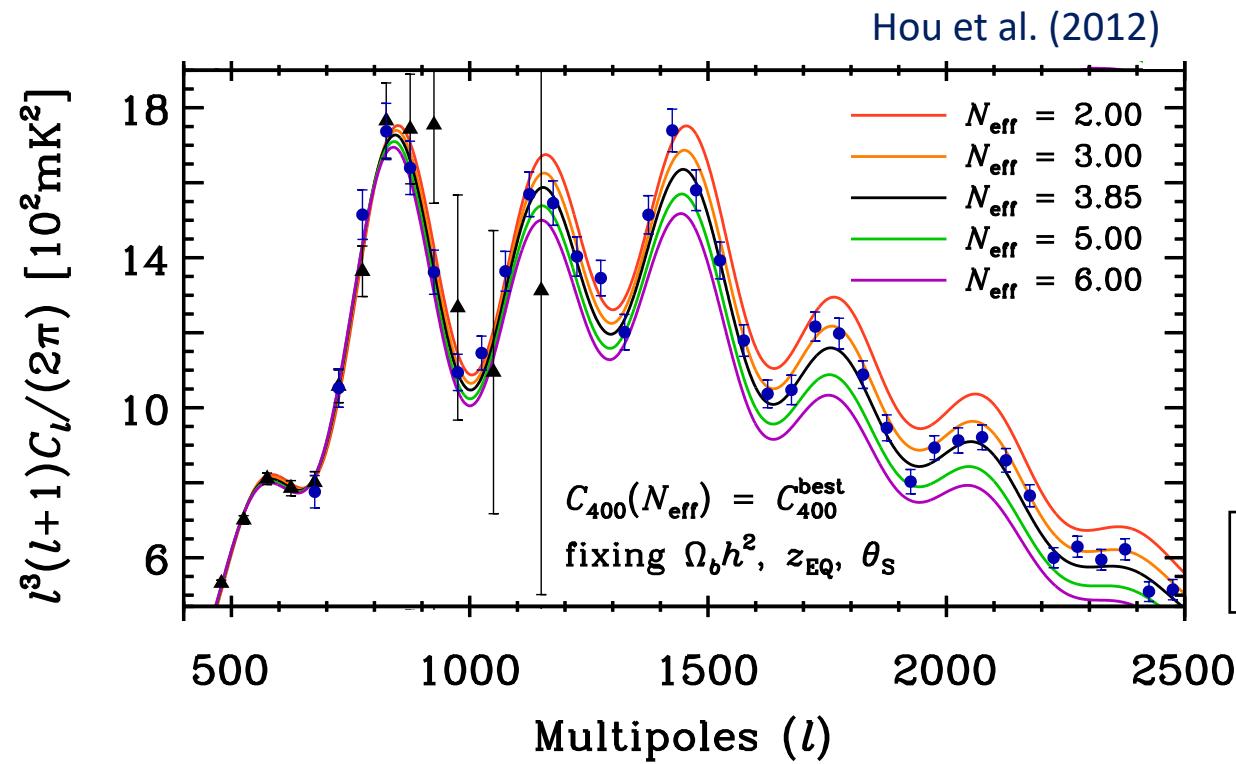
- Increase of diffusion (Silk) damping at large l

$$\theta_d \propto \theta_s H^{0.5}$$

Planck TTTEEE

$N_{\text{eff}} = 2.92 \pm 0.36$ (95% CL)

N_{eff} probes: CMB



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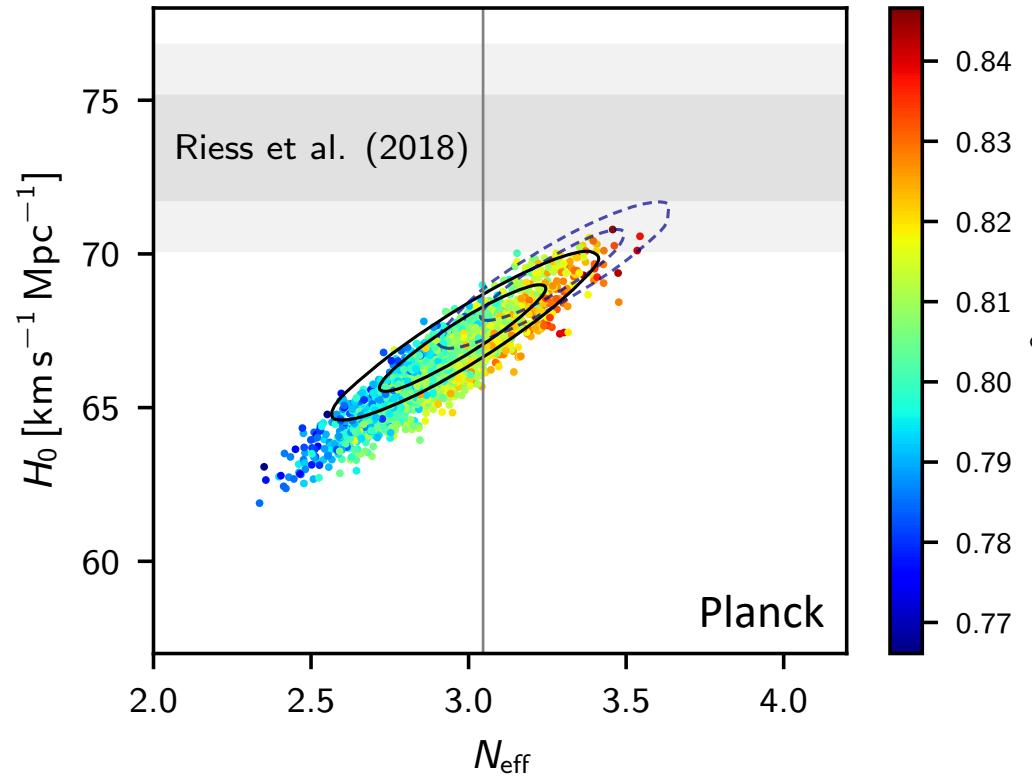
Planck TTTEEE

$N_{\text{eff}} = 2.92 \pm 0.36 \text{ (95% CL)}$

- Phase shift of the CMB acoustic peaks*

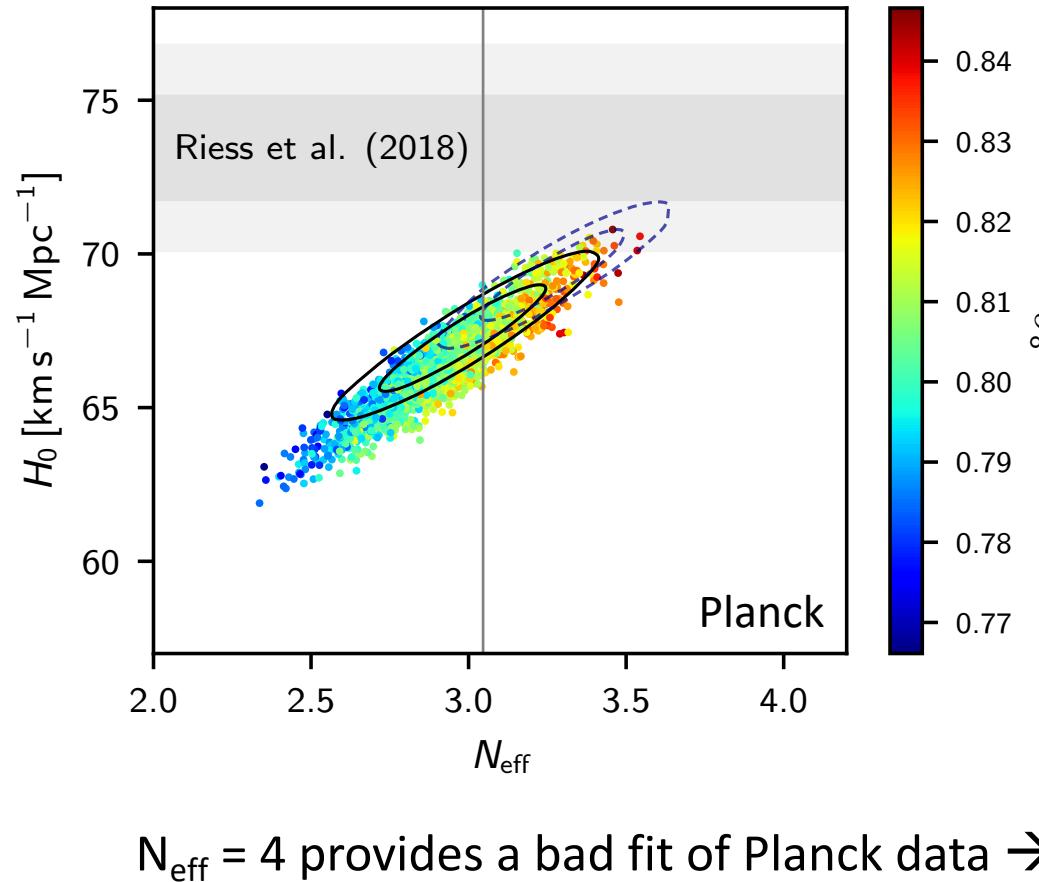
* The angular size of the sound horizon at CMB (θ_s) can be adjusted by increasing H_0

N_{eff} probes: CMB

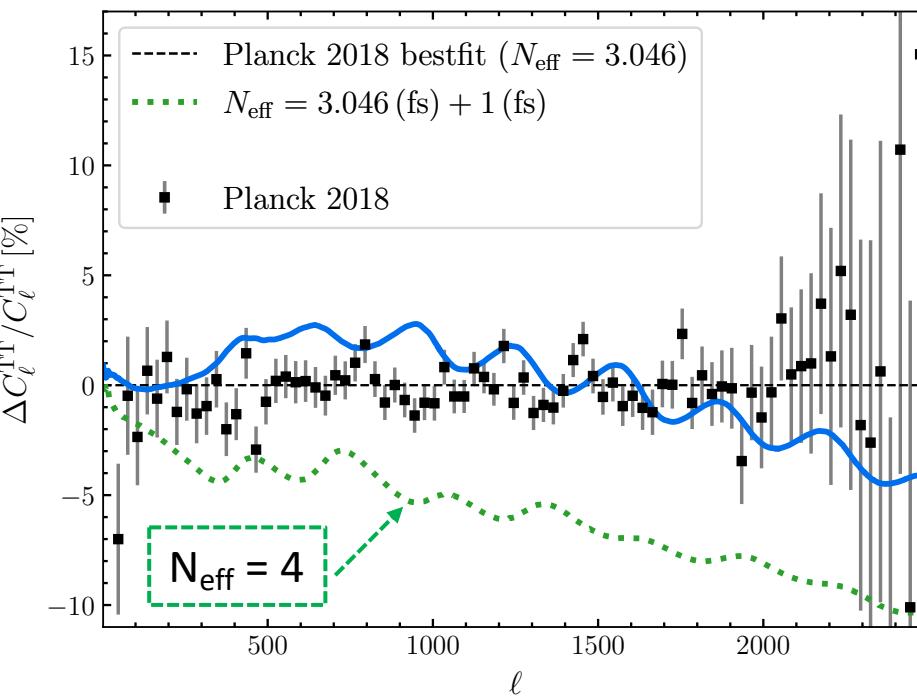


A larger (thermalised) N_{eff} is not a solution to the “ H_0 problem”:
 $H_0 = 67.4 \pm 0.5$ km/s/Mpc (Planck)
 $H_0 = 73.0 \pm 1.0$ km/s/Mpc (SHOES) [Riess et al. (2022)]

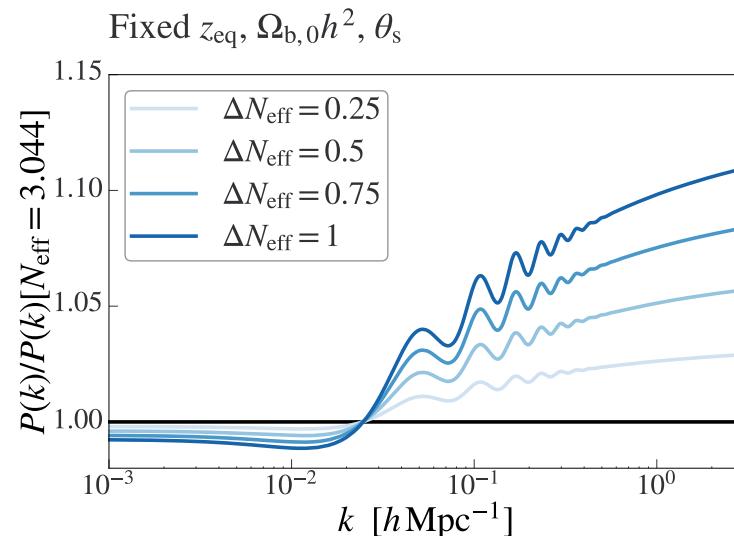
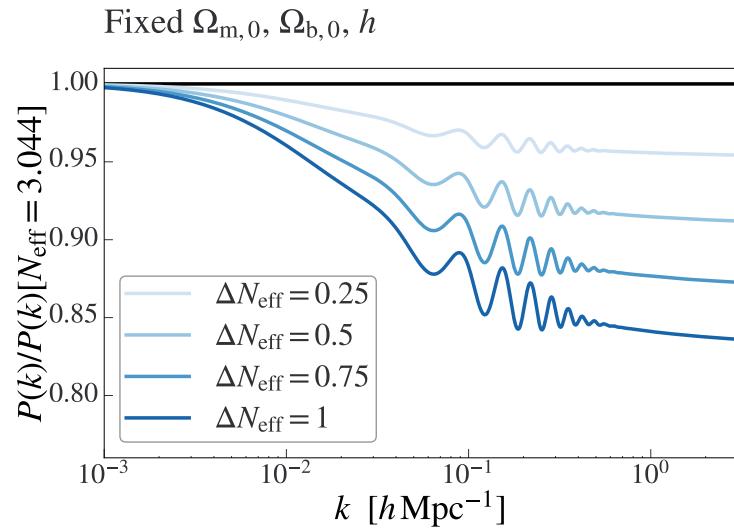
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N_{eff} probes: LSS



- Phase shift of the BAO

CMB (Planck + ACT)

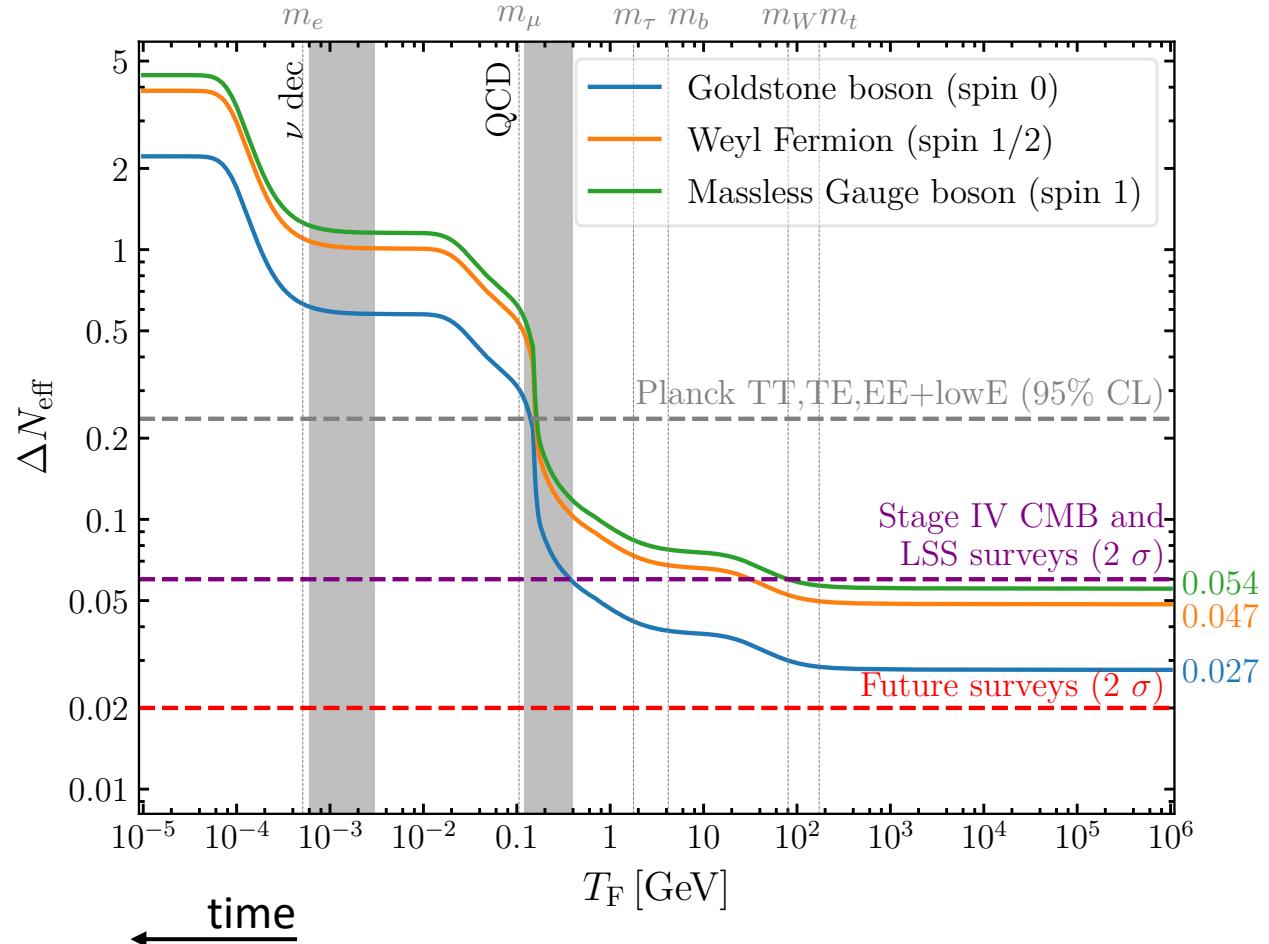
$$N_{\text{eff}} = 2.98 \pm 0.20 \text{ (95% CL)}$$

CMB + DESI BAO [[DESI Collaboration: Adame et al. \(2024\)](#)]

$$N_{\text{eff}} = 3.10 \pm 0.17 \text{ (95% CL)}$$

ΔN_{eff} : other light relics (if any)

$$\Delta N_{\text{eff}} = \frac{\rho_{\text{DR}}}{\rho_\nu} \propto \left(\frac{T_{\text{DR}}}{T_\nu} \right)^4$$



N_{eff} : conclusions

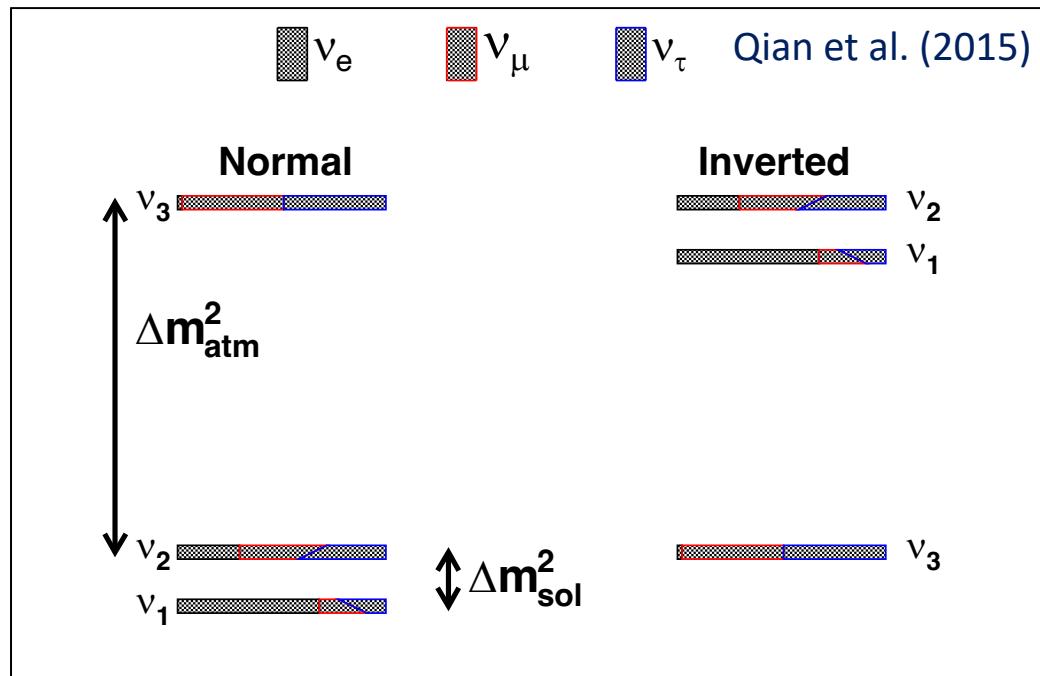
- Cosmology provides an **indirect detection of the CvB**
- Cosmology at different epochs and at different scales is **consistent with Standard Model prediction** $N_{\text{eff}}^{\text{SM}} = 3.044$
- Future CMB and LSS surveys can exclude the existence of **new light particles** decoupling before the onset of QCD phase transition for any spin.

Detecting the neutrino mass in the CvB

Neutrinos from the lab

Neutrino oscillations

Neutrino Mass Hierarchy



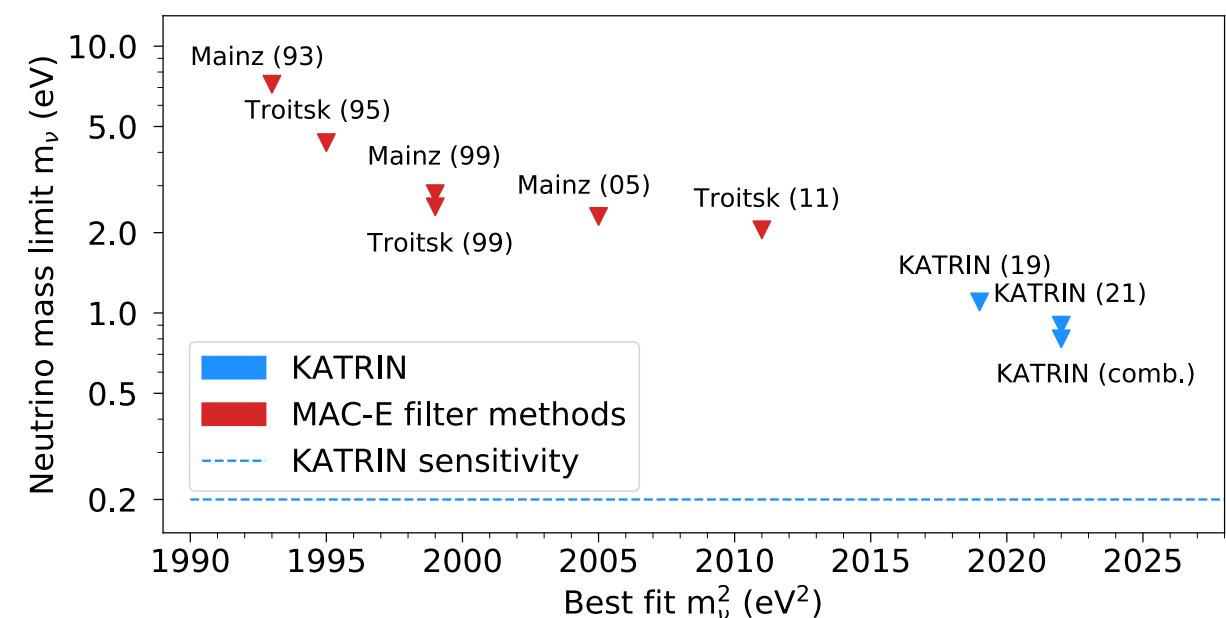
$$\text{Min. } \sum m_\nu (\text{NH}) = 0.058 \text{ eV}; \quad \text{Min. } \sum m_\nu (\text{IH}) = 0.100 \text{ eV}$$

The absolute neutrino mass scale is not yet determined by neutrino oscillation data

Lectures by Di Lodovico, Lisi, Maricic, Petcov

Neutrino β -decay

(Neutrinoless $\beta\beta$ -decay only if neutrinos are Majorana)



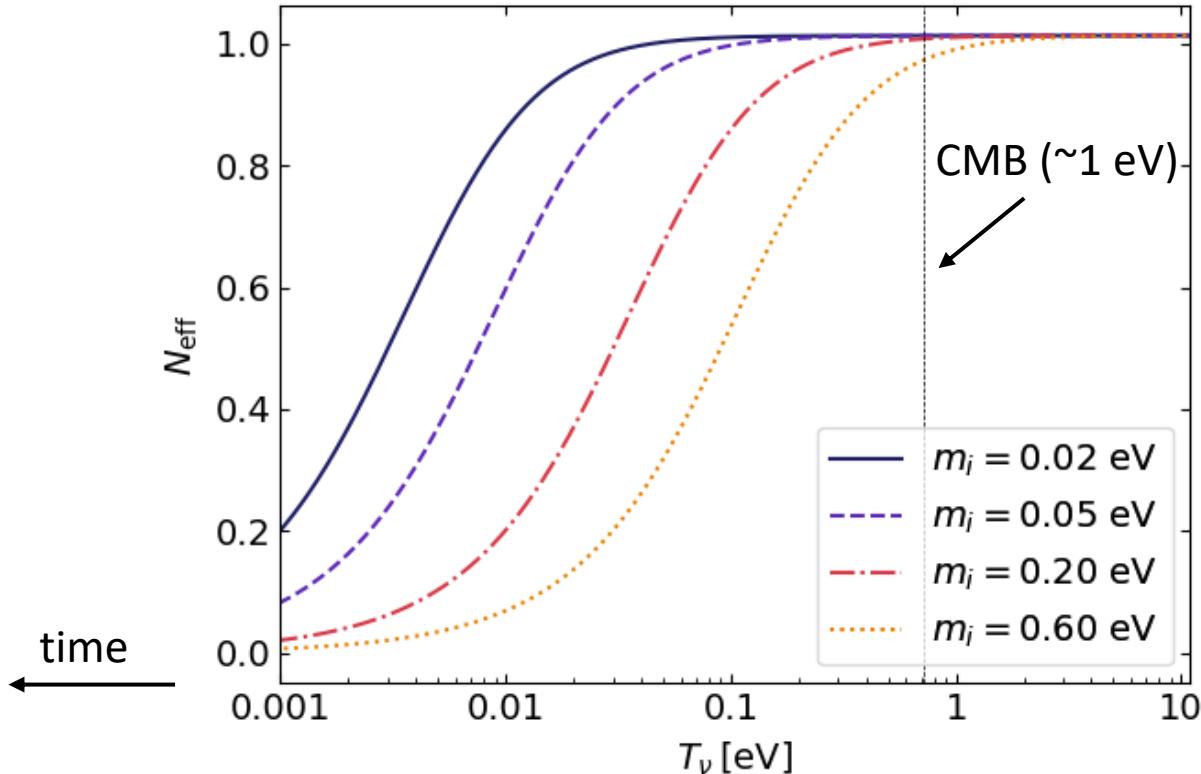
$$m_\beta^2 = \sum_i |U_{ei}|^2 m_i^2 < 0.8 \text{ eV (90% CL)}$$

[Aker et al. (2022)]

Sensitivity: 200 meV (90% CL)

Colloquium by Kathrin Valerius

The duality of the CvB



Early times: neutrinos as **radiation**

$$\rho_{\text{rad}} = \rho_\gamma \left[1 + N_{\text{eff}} \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \right]$$

Neutrino non-relativistic transition

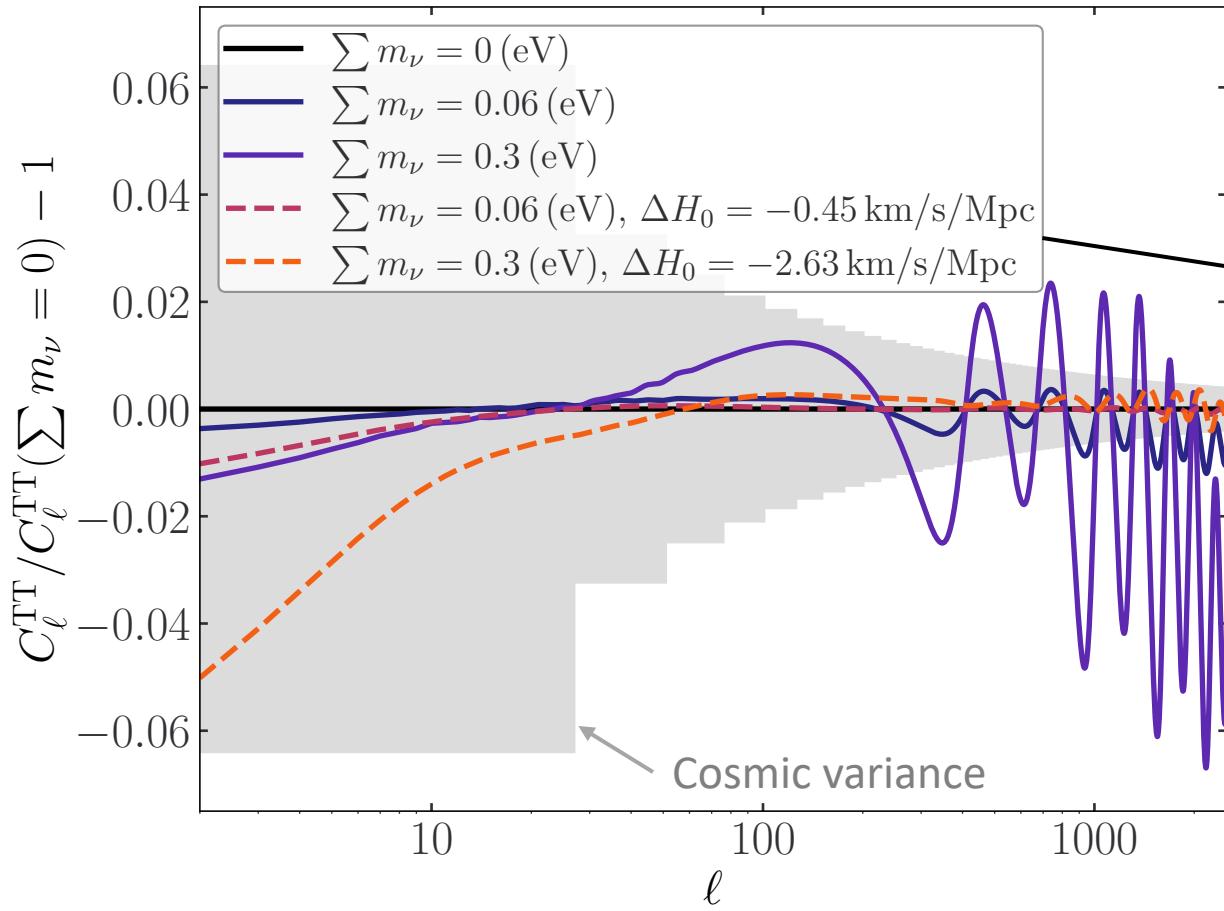
$$z_{\text{nr},i} = 2 \times 10^3 \left(\frac{m_{\nu,i}}{1 \text{ eV}} \right)$$

Late times (after CMB formation): neutrinos as **matter** (contributing to dark matter as hot dark matter)

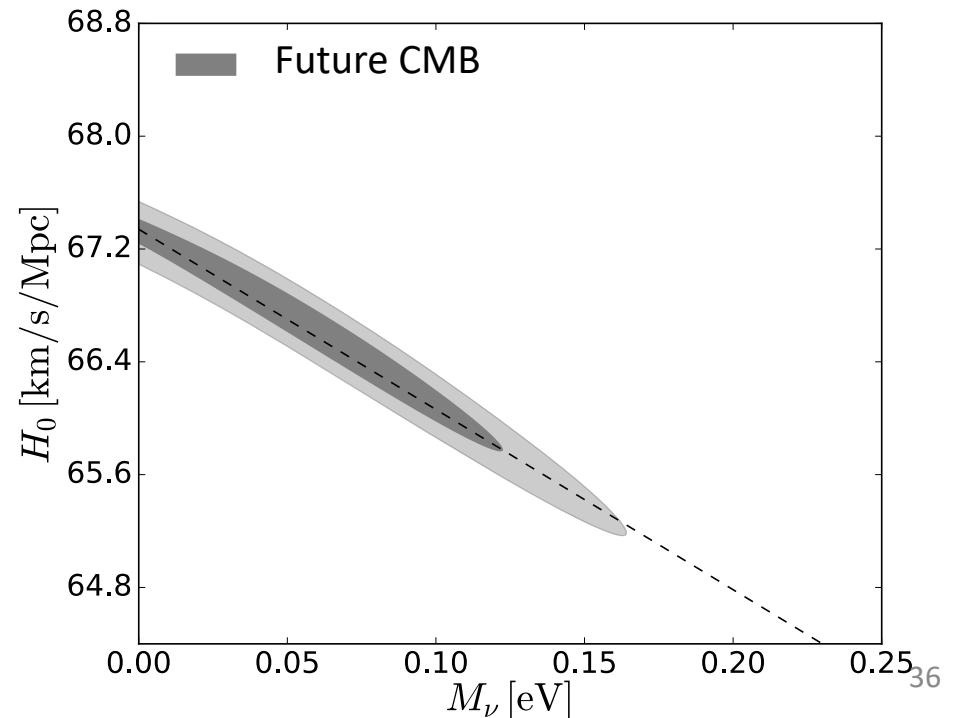
- Neutrino mass $\Omega_\nu h^2 = \frac{\sum m_{\nu,i}}{93.12 \text{ eV}}$ [Mangano et al. 2005, Froustey et al. 2020]

not individual masses

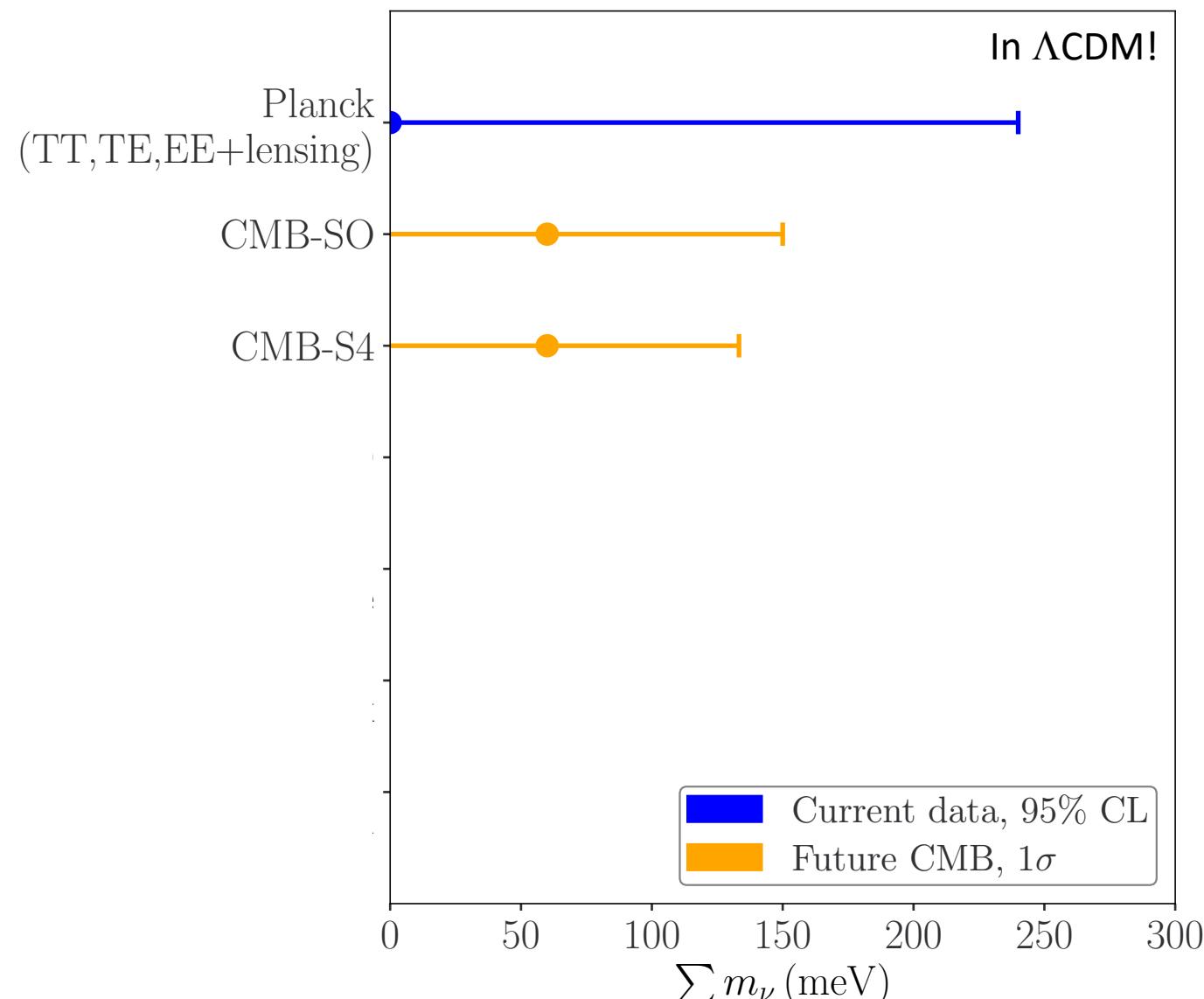
Neutrino mass probes: CMB



- Background effects
 - Perturbation effects
- Varying H_0 to fix θ_s (i.e., the angular size of the sound horizon at recombination)



Neutrino mass constraints: CMB



KATRIN: $\Sigma m_\nu < \sim 2.7$ eV

Fiducial value:

- $\Sigma m_\nu = 58$ meV

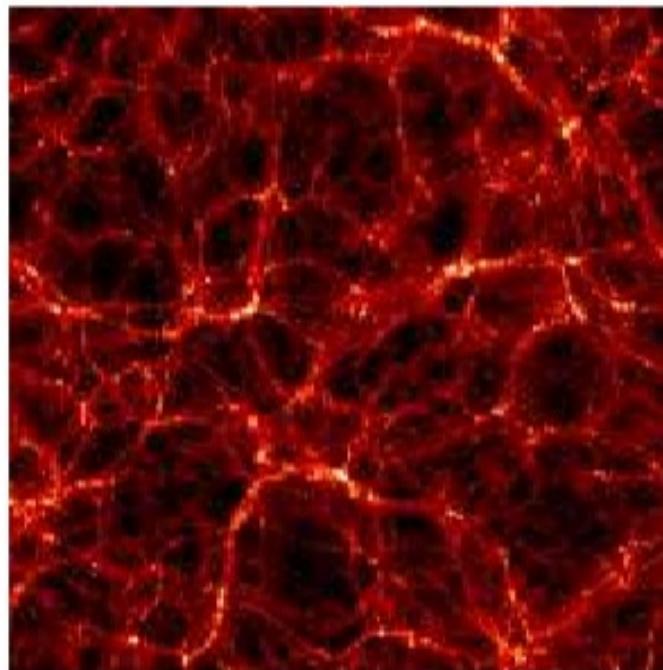
Minimum from neutrino
oscillation experiments

CMB alone will not be able
to detect the neutrino mass
→ Large Scale Structures

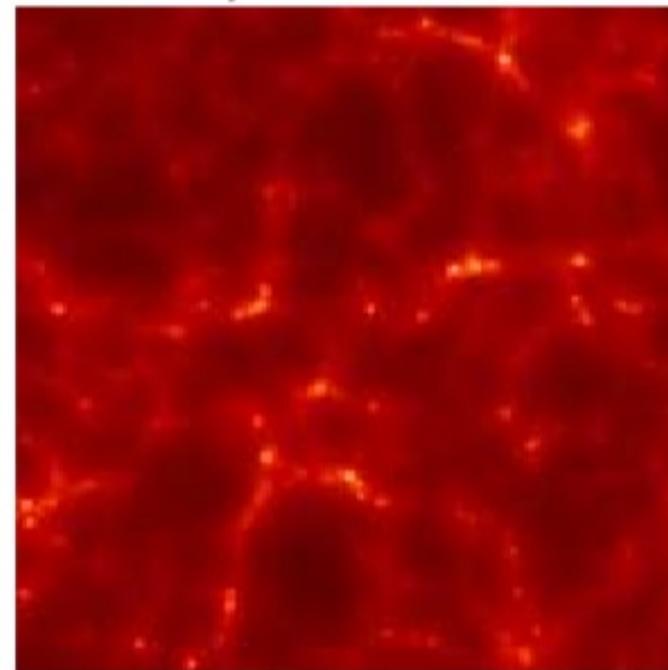
Neutrino mass probes: LSS

- Free-Streaming $d_{\text{FS},i} \sim 1 \text{ Gpc} \frac{eV}{m_{\nu,i}}$

CDM

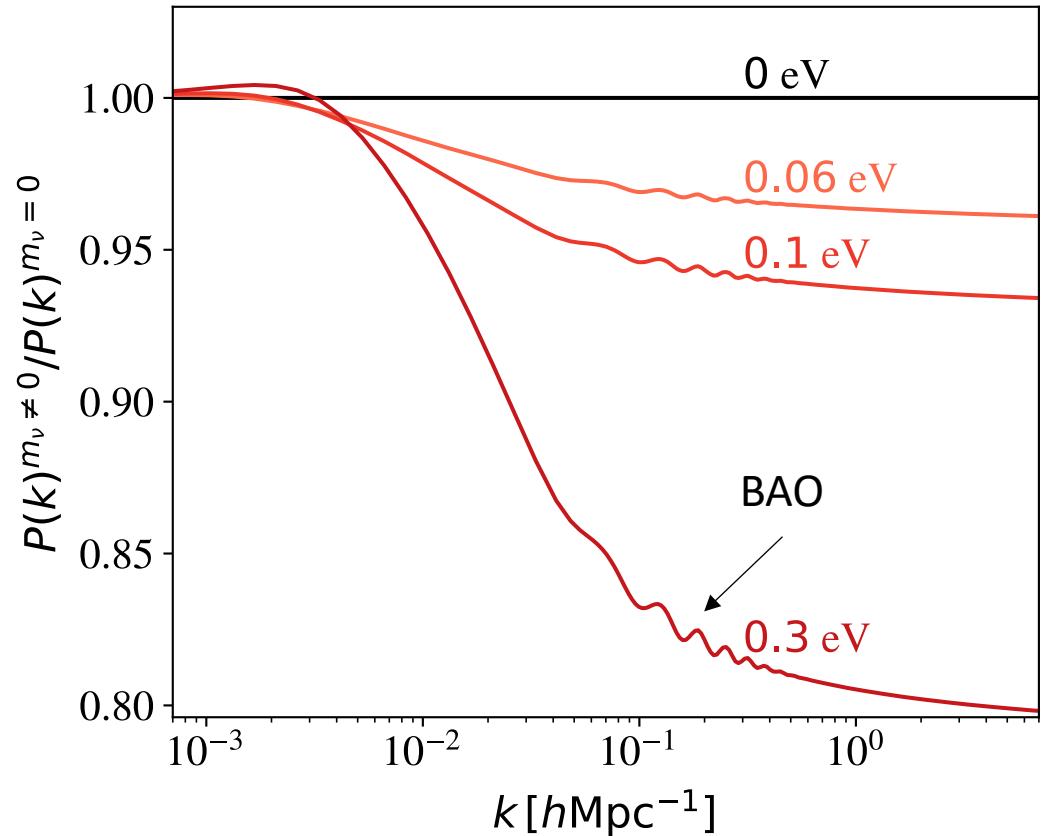


$m_{\nu} = 0.5 \text{ eV}$



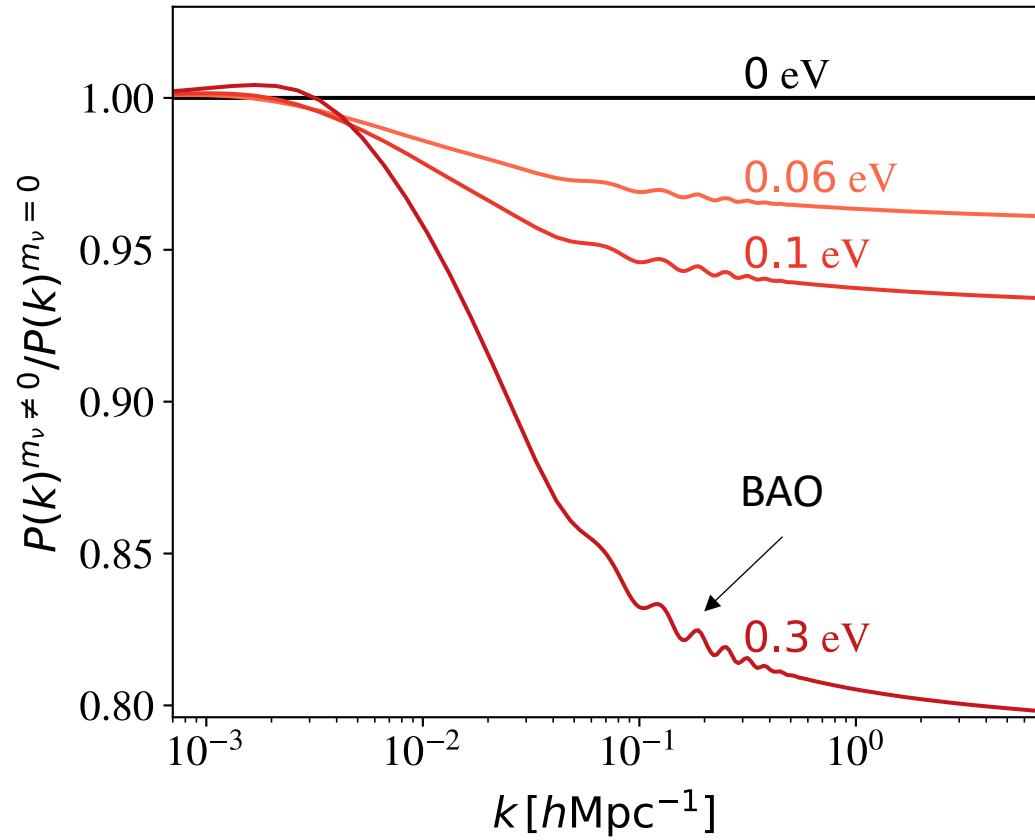
Villaescusa Navarro et al. (2013)

Neutrino mass probes: LSS

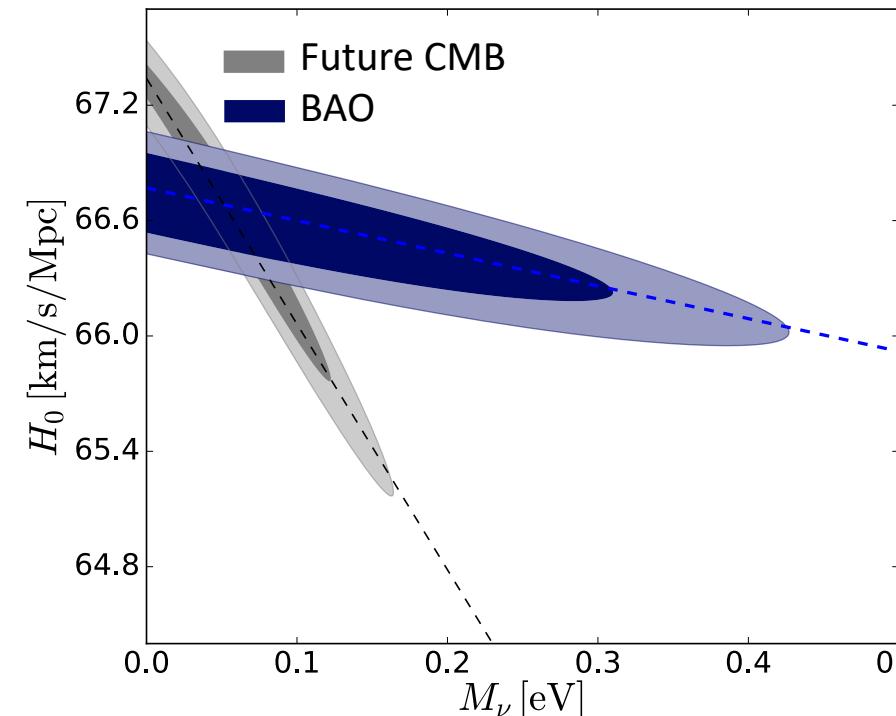


- Massive neutrinos do not cluster
- Massive neutrinos slow down the growth of CDM perturbations
 - Massless neutrino Universe $\delta_{\text{cdm}}^{m_\nu=0} \propto a$
 - Massive neutrino Universe $\delta_{\text{cdm}}^{m_\nu \neq 0} \propto a^{1 - \frac{3}{5} \frac{\Omega_\nu}{\Omega_m}}$

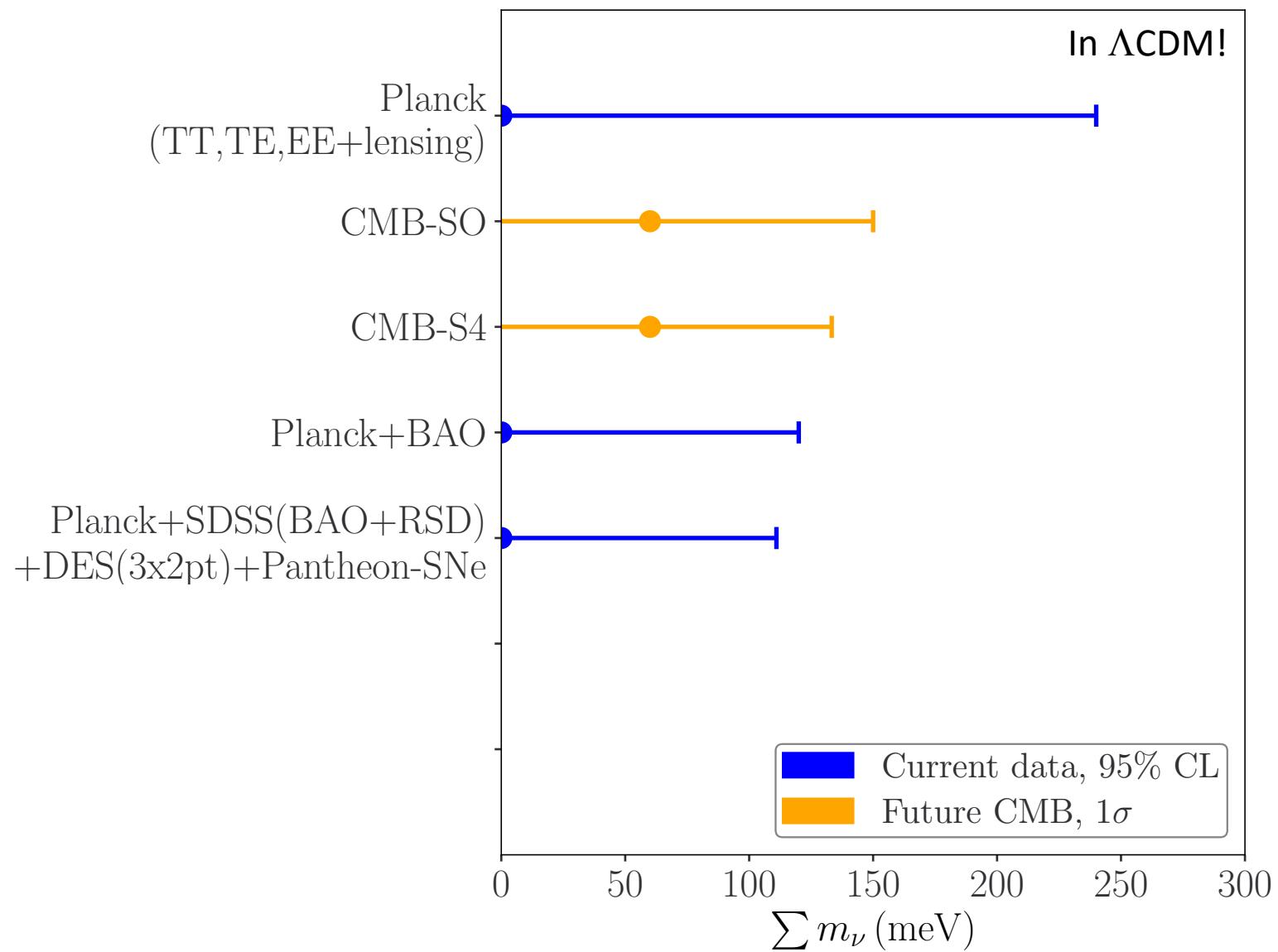
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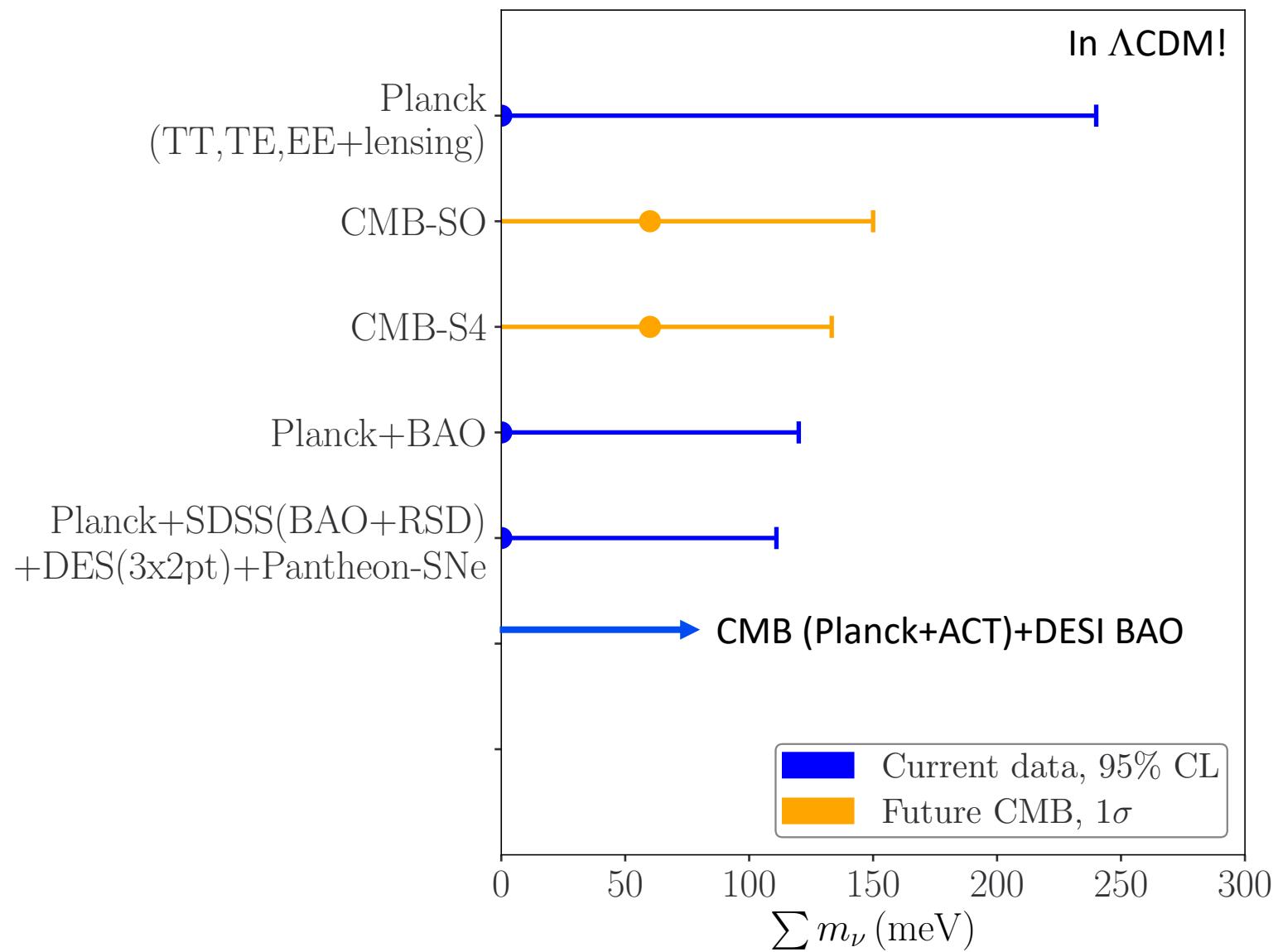
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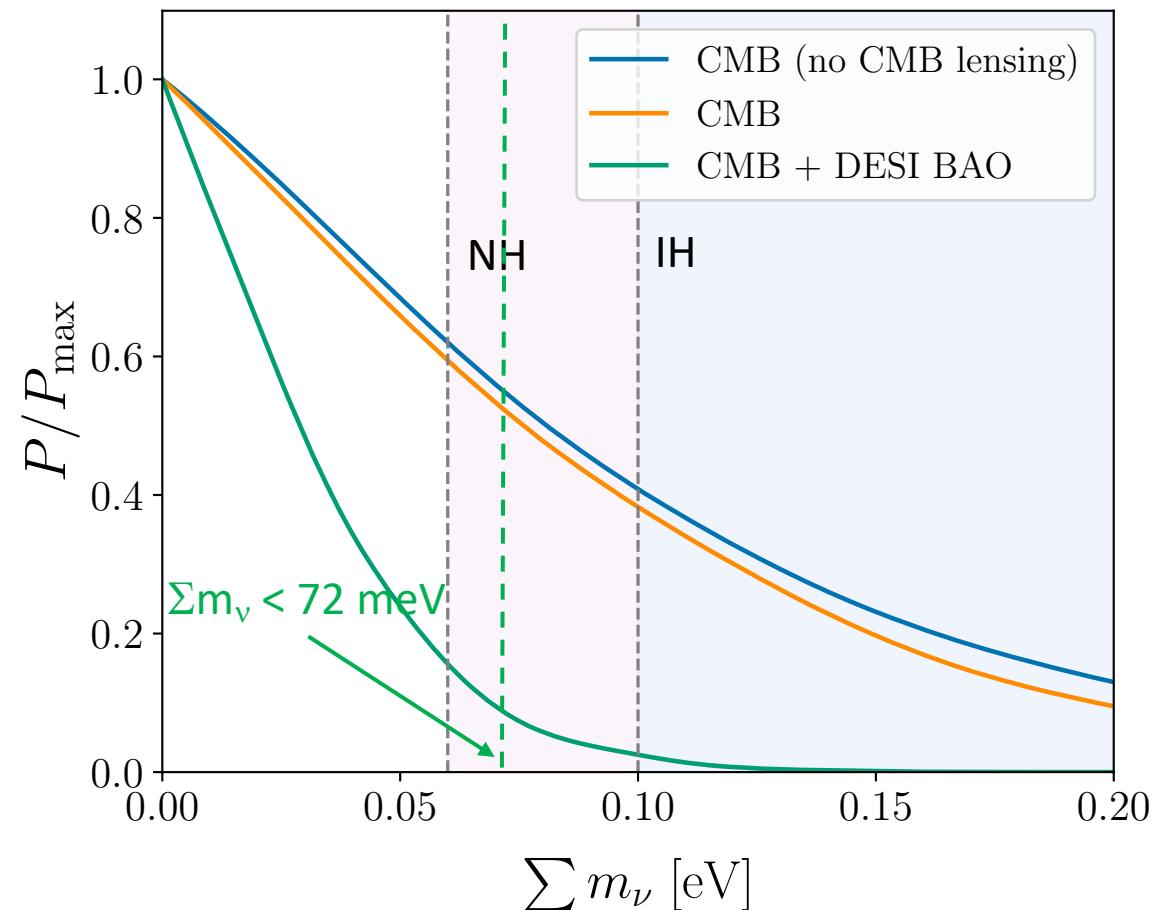
Neutrino mass constraints: LSS



Neutrino mass constraints: LSS



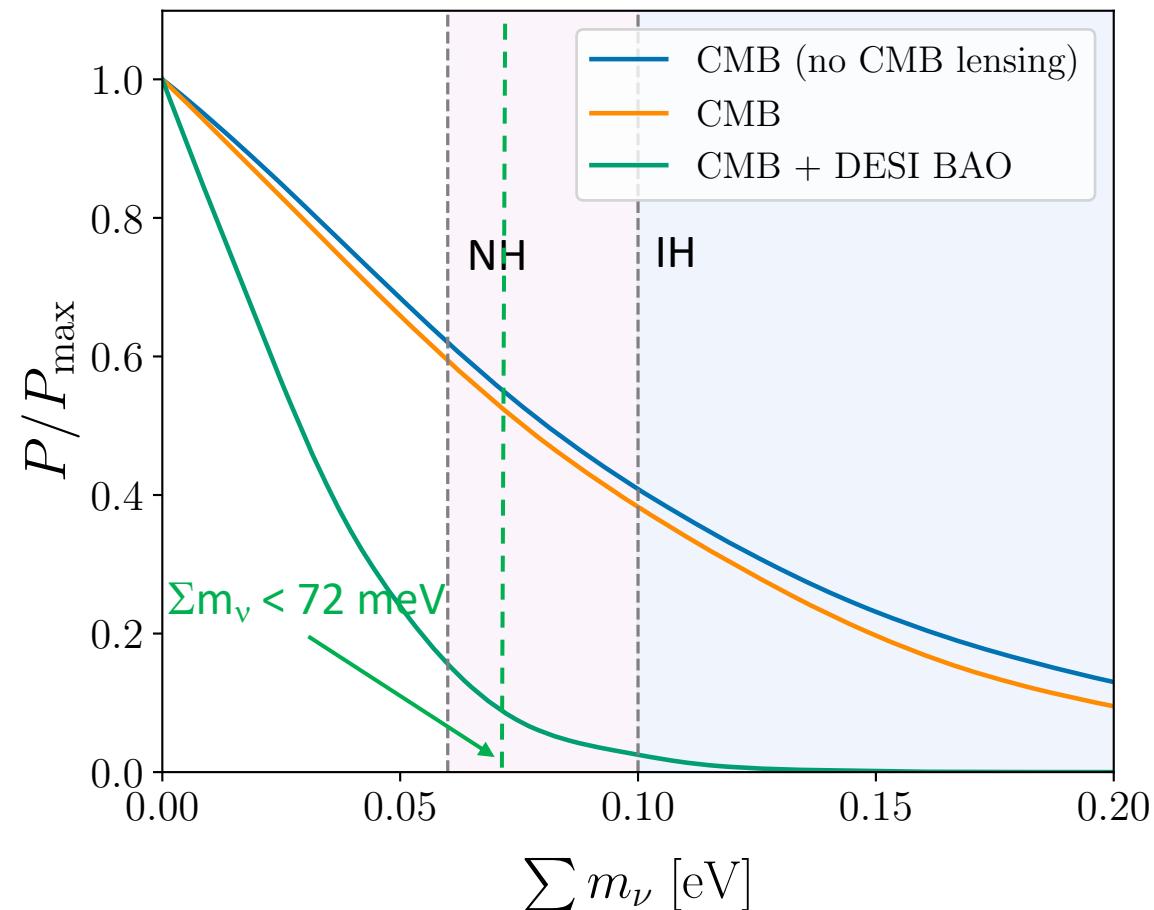
Neutrino mass constraints: DESI



DESI Collaboration: Adame et al. (2024)

CMB (Planck+ACT) + DESI BAO: $\Sigma m_\nu < 72 \text{ meV}$, 95% CL

Neutrino mass constraints: DESI



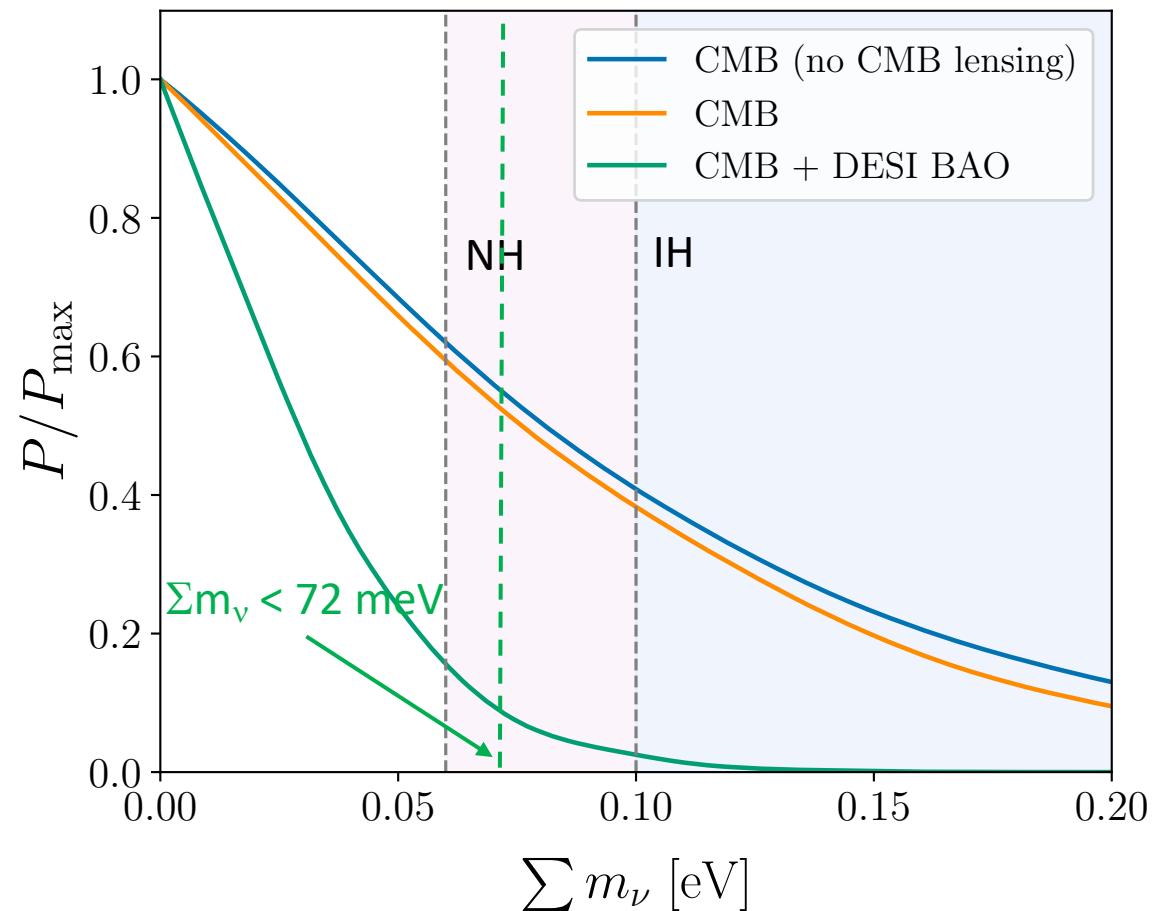
DESI Collaboration: Adame et al. (2024)

CMB (Planck+ACT) + DESI BAO: $\Sigma m_\nu < 72$ meV, 95% CL

Prior: $\Sigma m_\nu > 0$

Prior: $\Sigma m_\nu > 59$ meV $\rightarrow \Sigma m_\nu < 113$ meV, 95% CL

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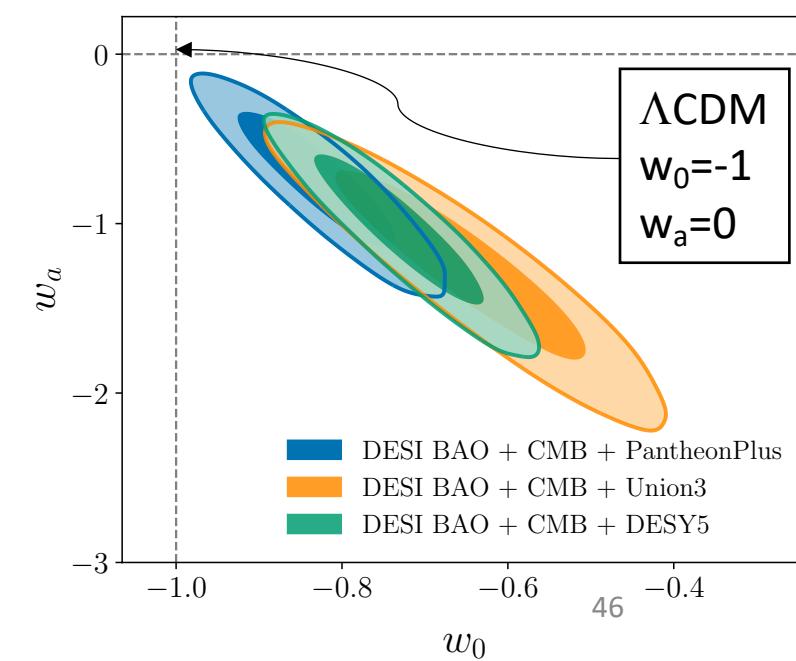
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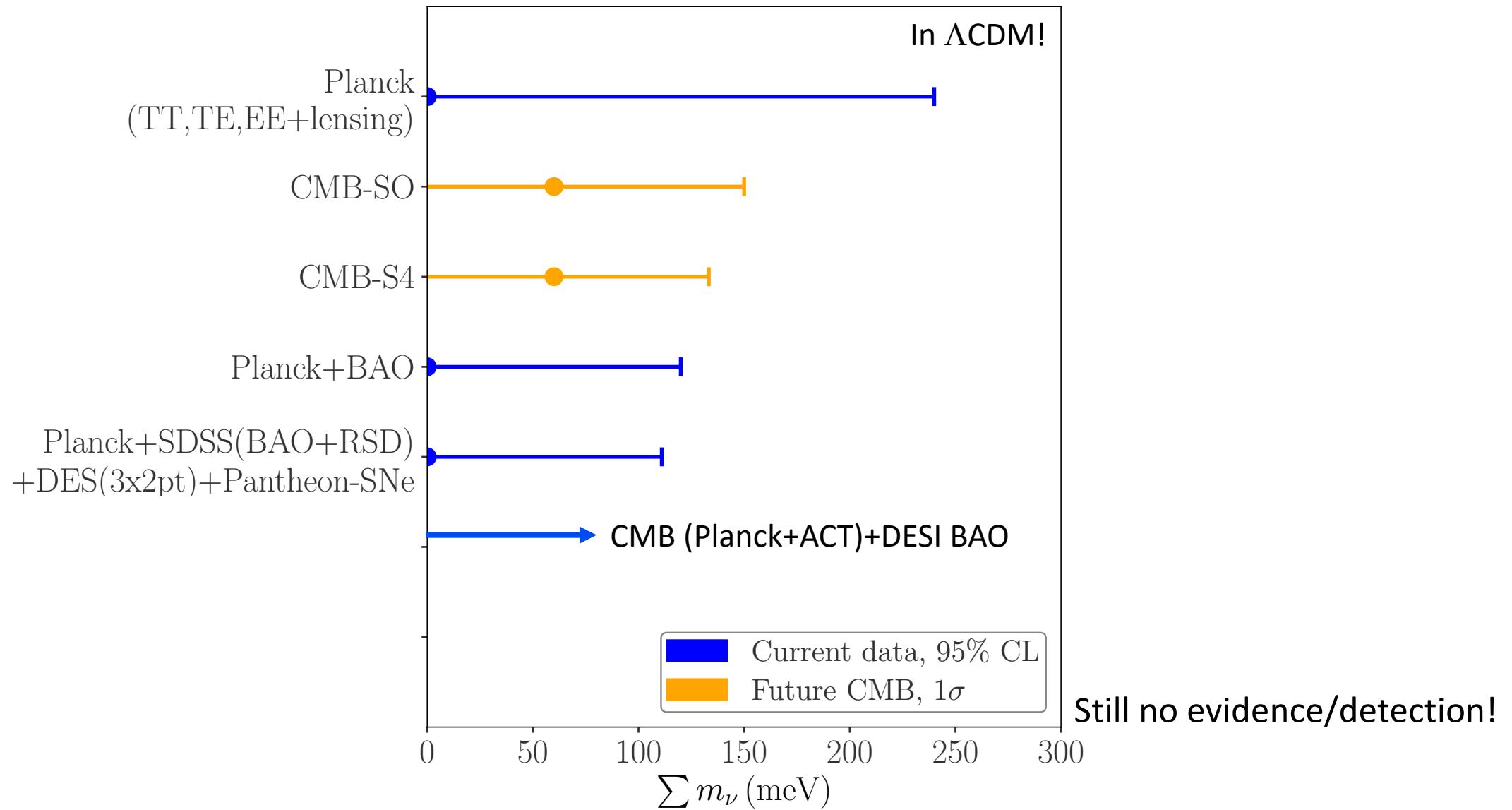
In Λ CDM!

In $w_0 w_a$ CDM:

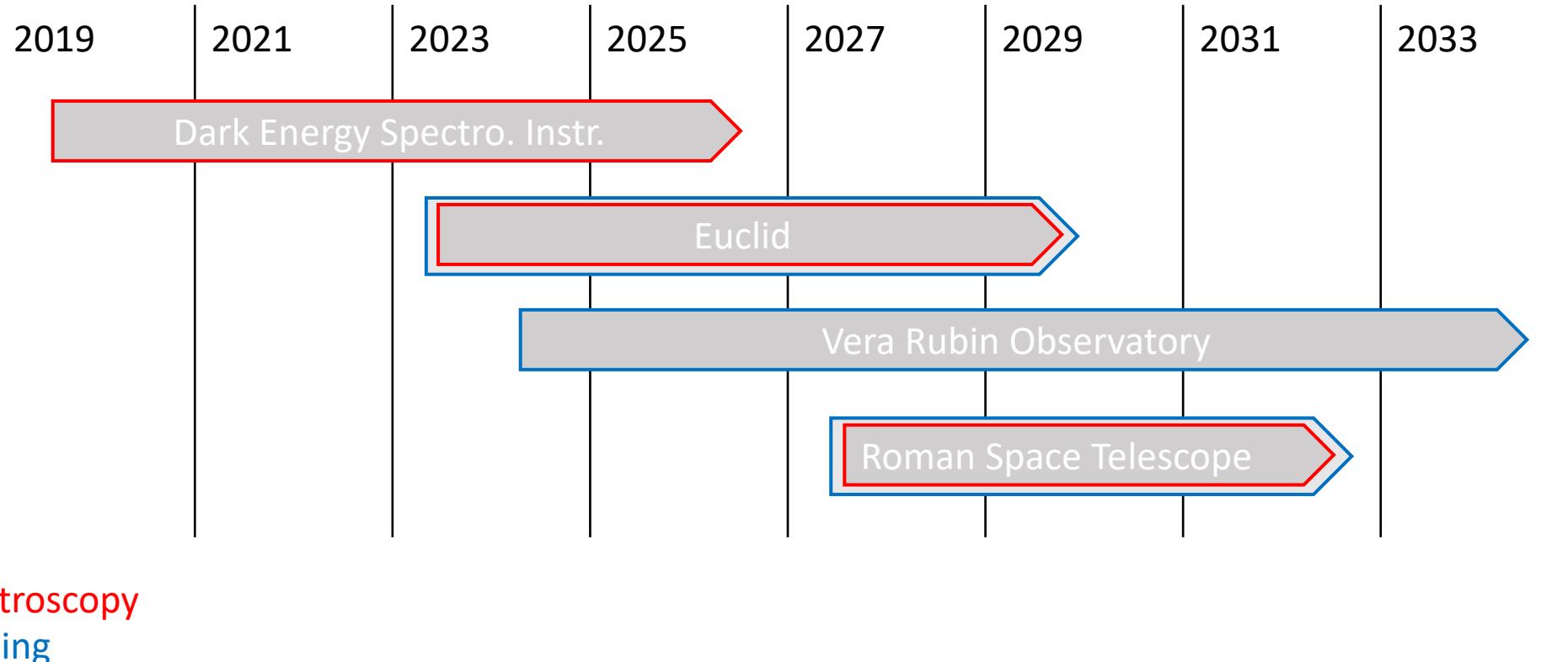
$\Sigma m_\nu < 195$ meV



Neutrino mass constraints: LSS

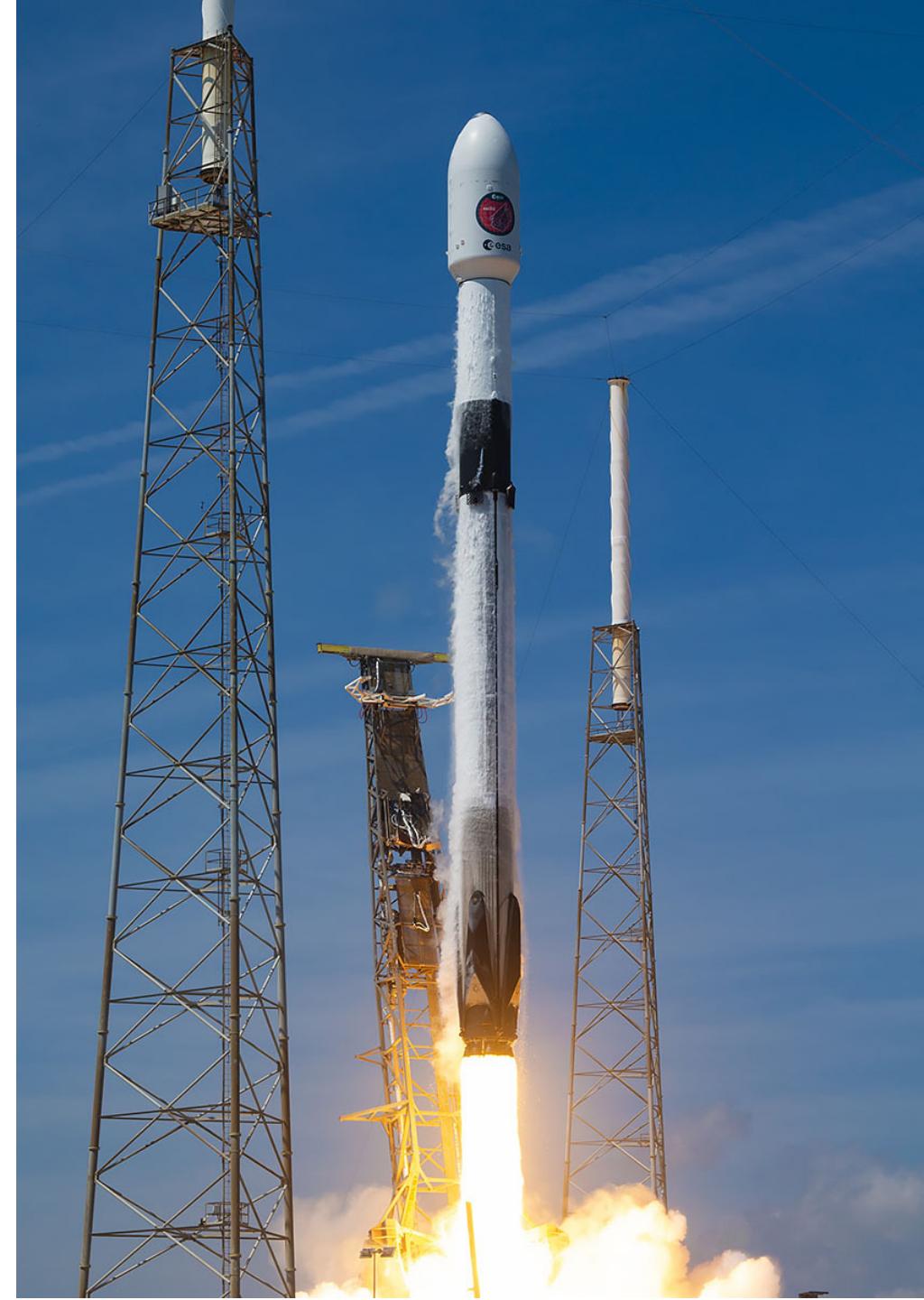


Stage IV Large Scale Surveys

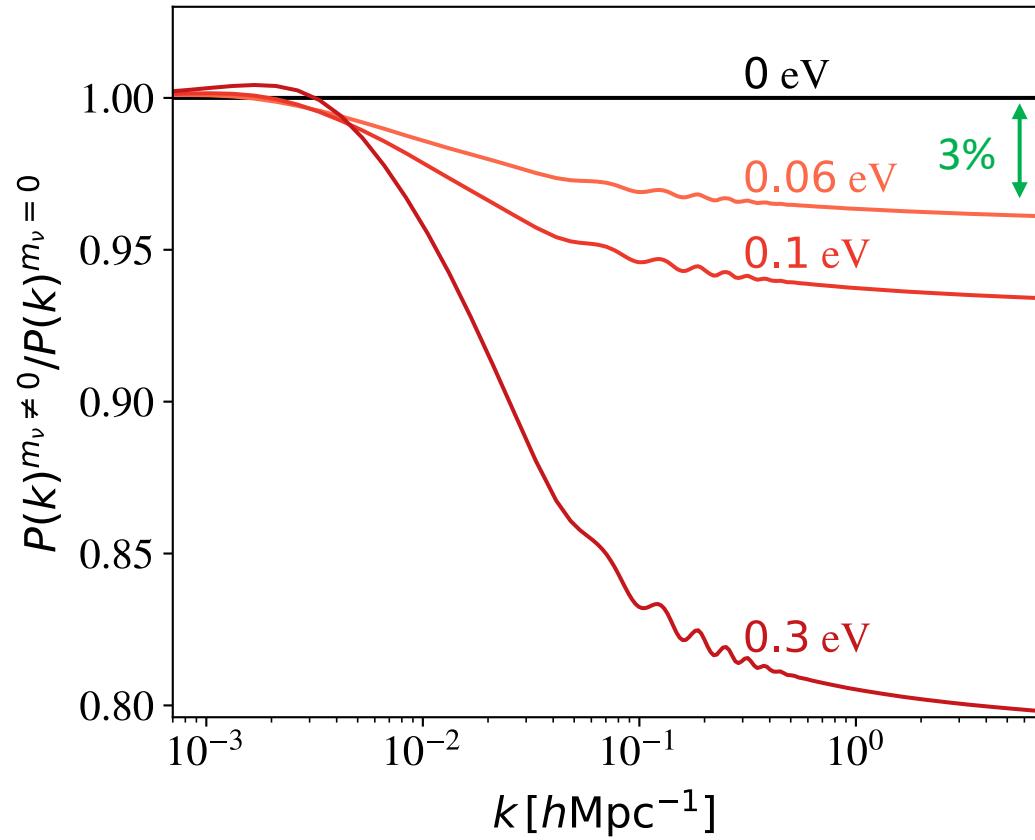


Euclid in a nutshell

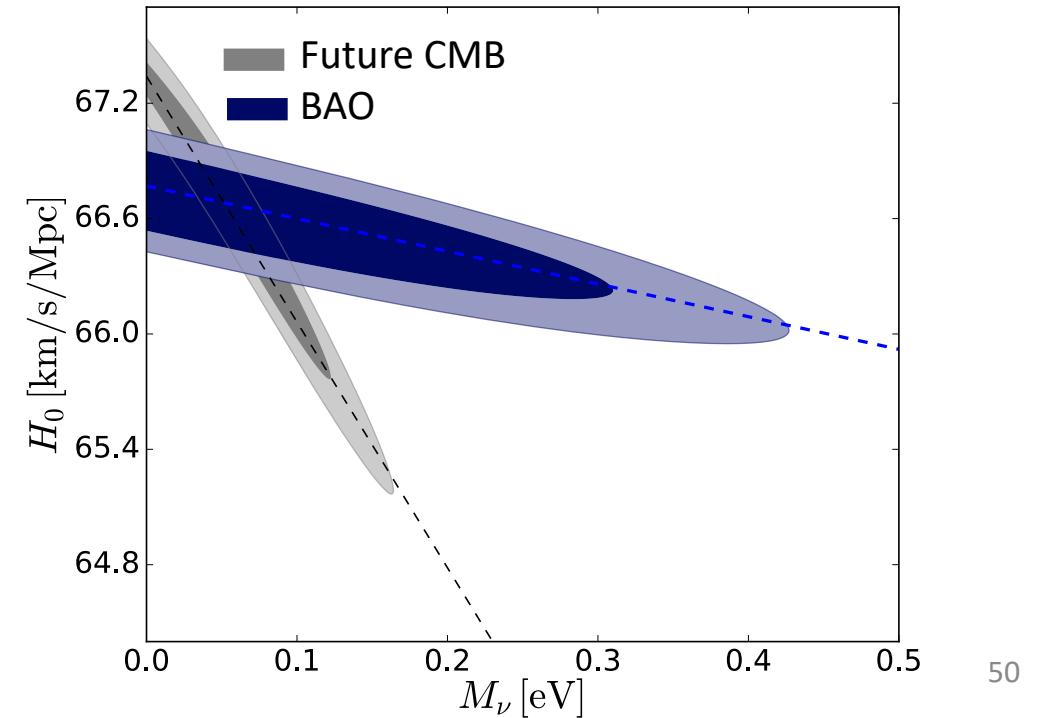
- **ESA M2 space mission** in the framework of the Cosmic Vision program
- Launch **July 1st 2023**. Duration > 6 years
- 1.2m telescope with two instruments: Visible Imager (**VIS**) and Near Infrared Spectrometer and Photometer (**NISP**)
- Wide survey (**14.000 deg²**) and deep survey (40 deg² in 3 different fields)
- Measurements of over **1 billion images** and more than **30 millions spectra** of galaxies out to z>2
- Main scientific objectives: **Dark Energy, Dark Matter, and General Relativity**
- Primary probes: **Galaxy Clustering and Weak Lensing** (**1% accuracy**)



Neutrino mass probes: LSS

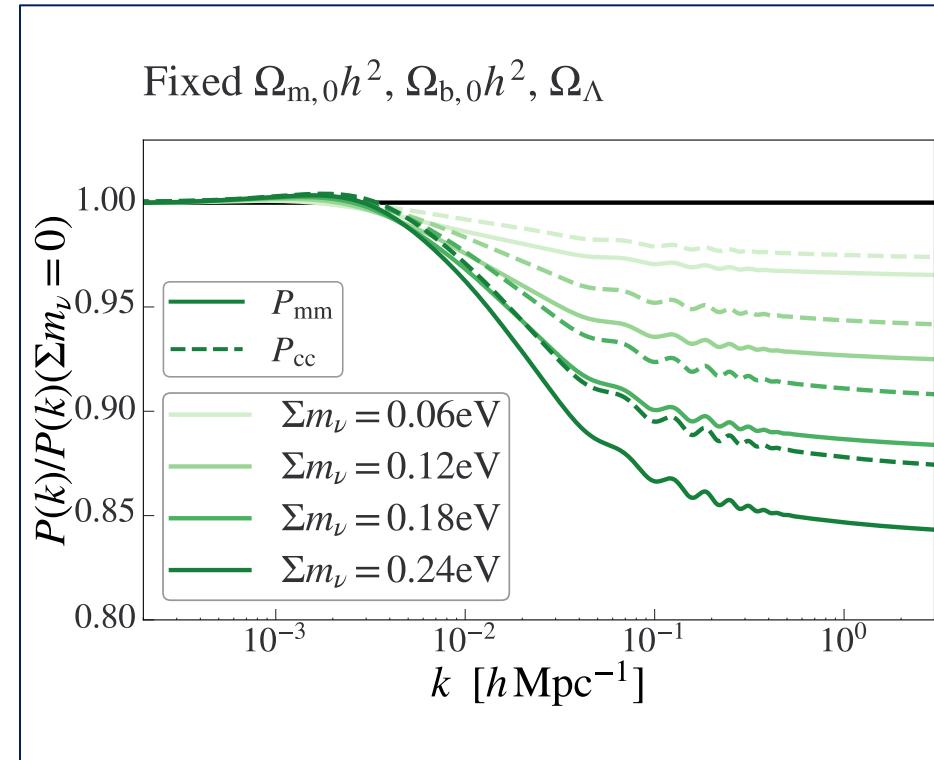


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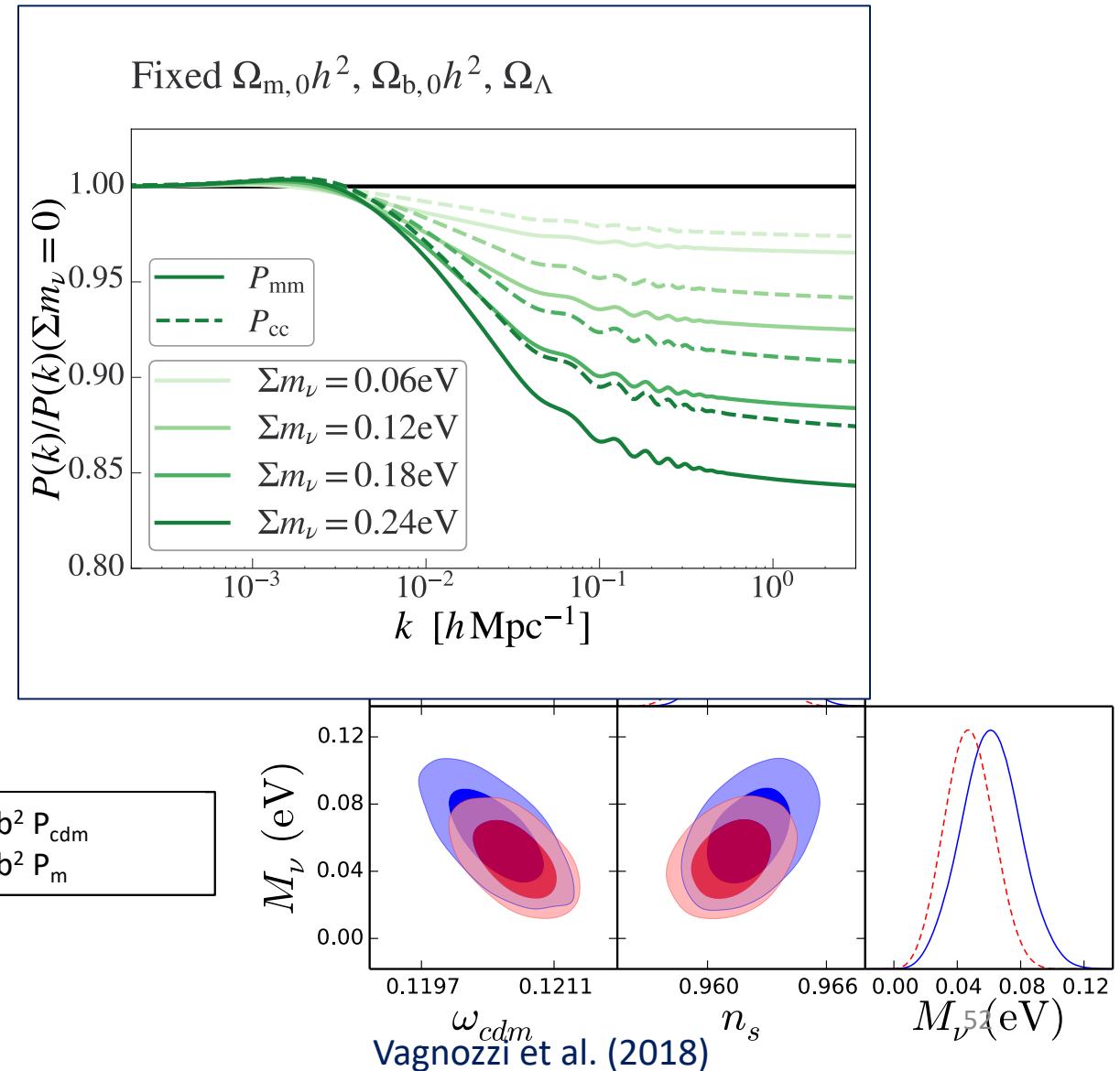
Known unknowns (systematics, etc.)

1. Galaxy bias $P_{\text{galaxy}} = b^2 P_{\text{cdm}}$ [Castorina et al. (2014); Vagnozzi et al. (2018)]



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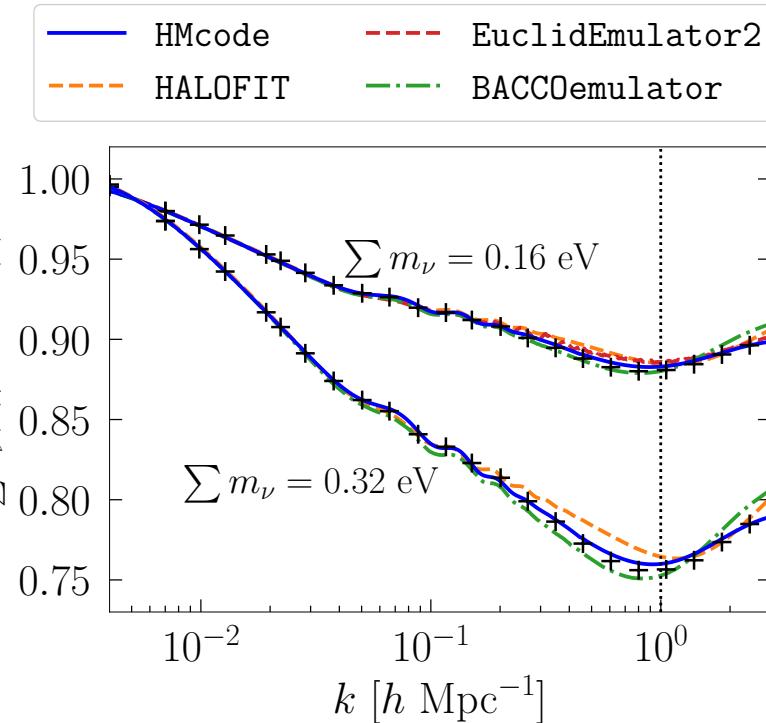
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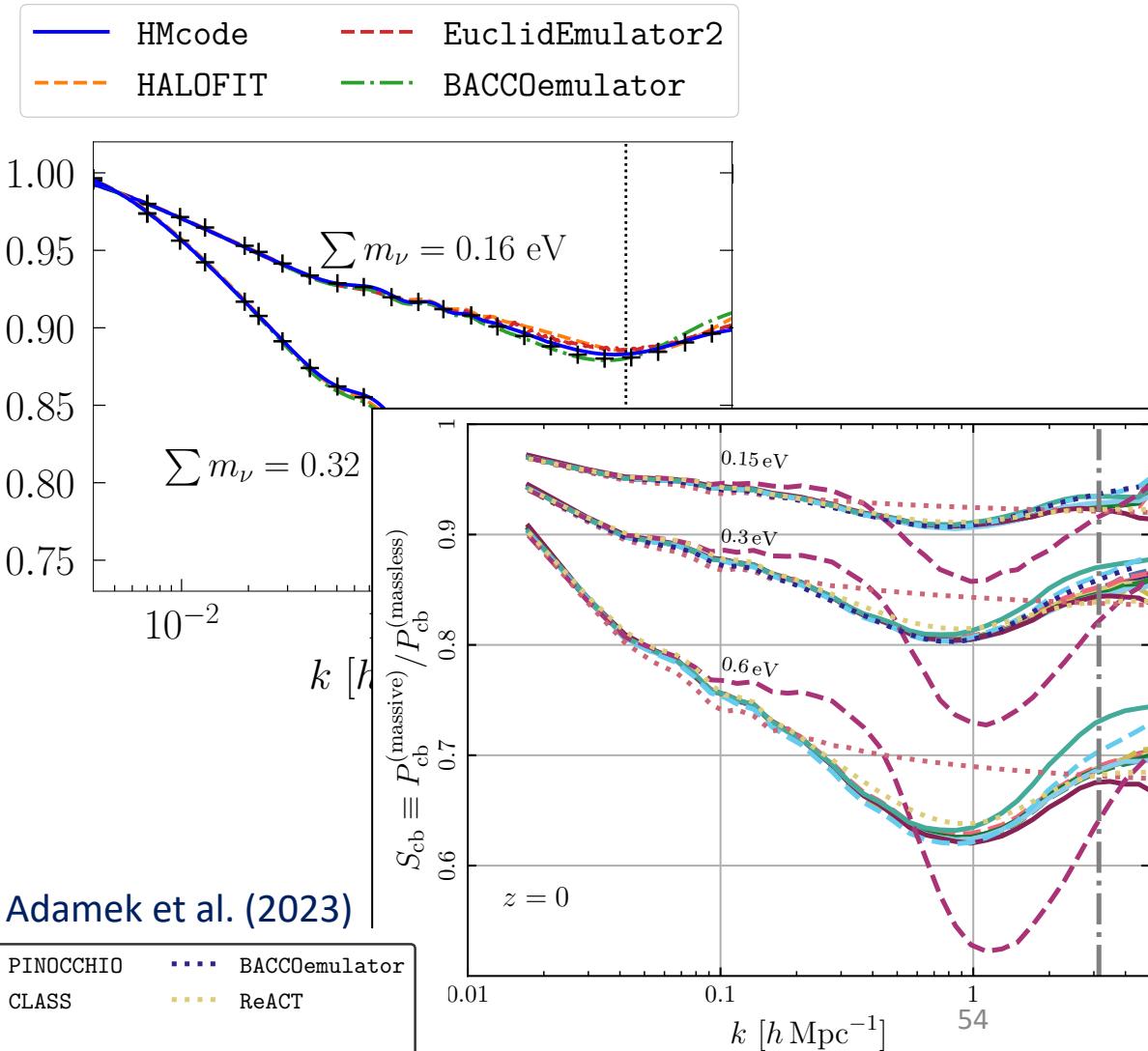
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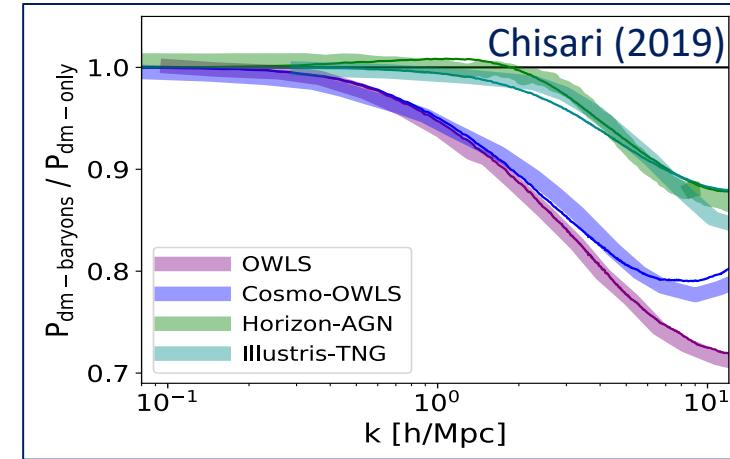
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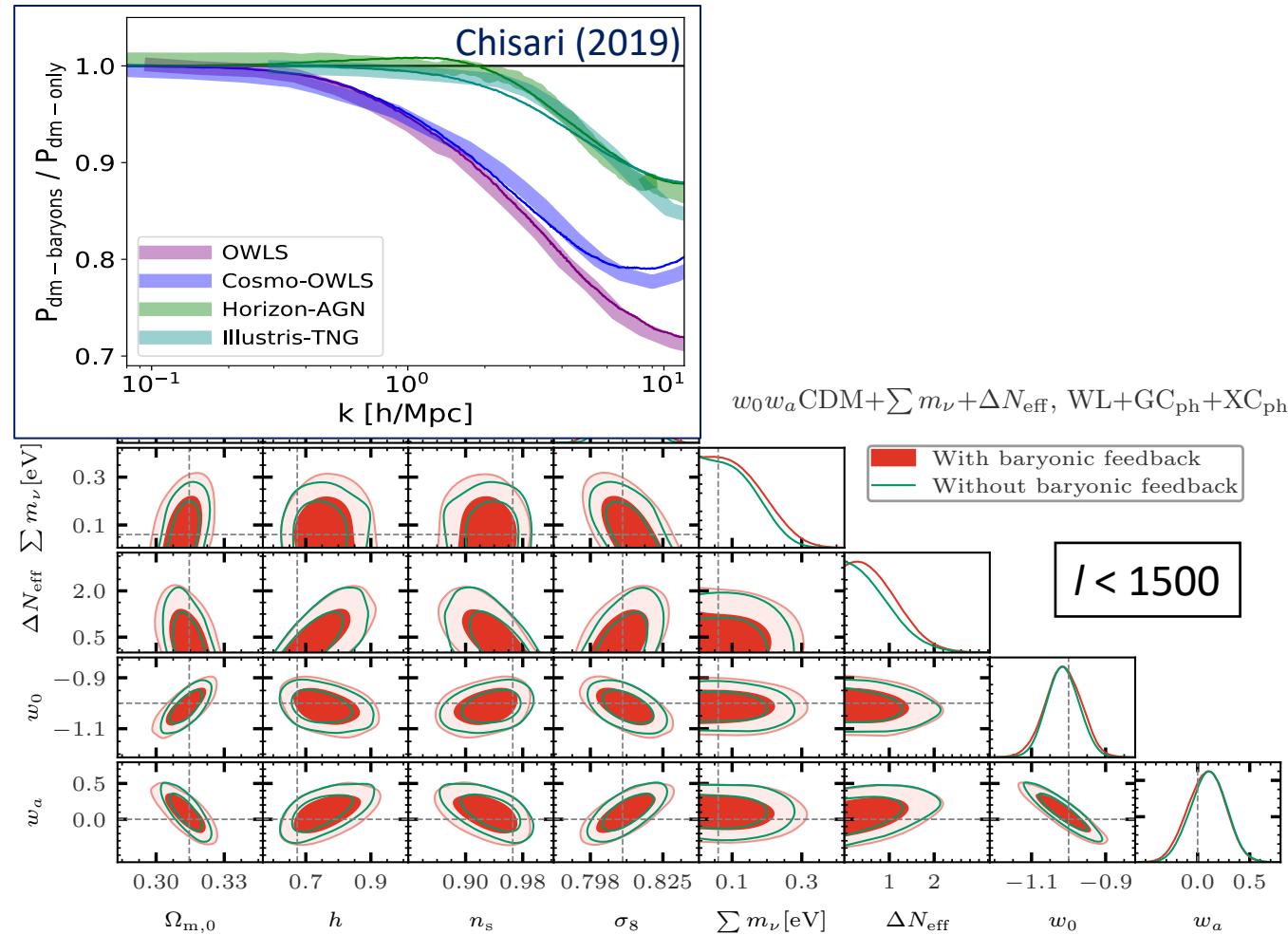
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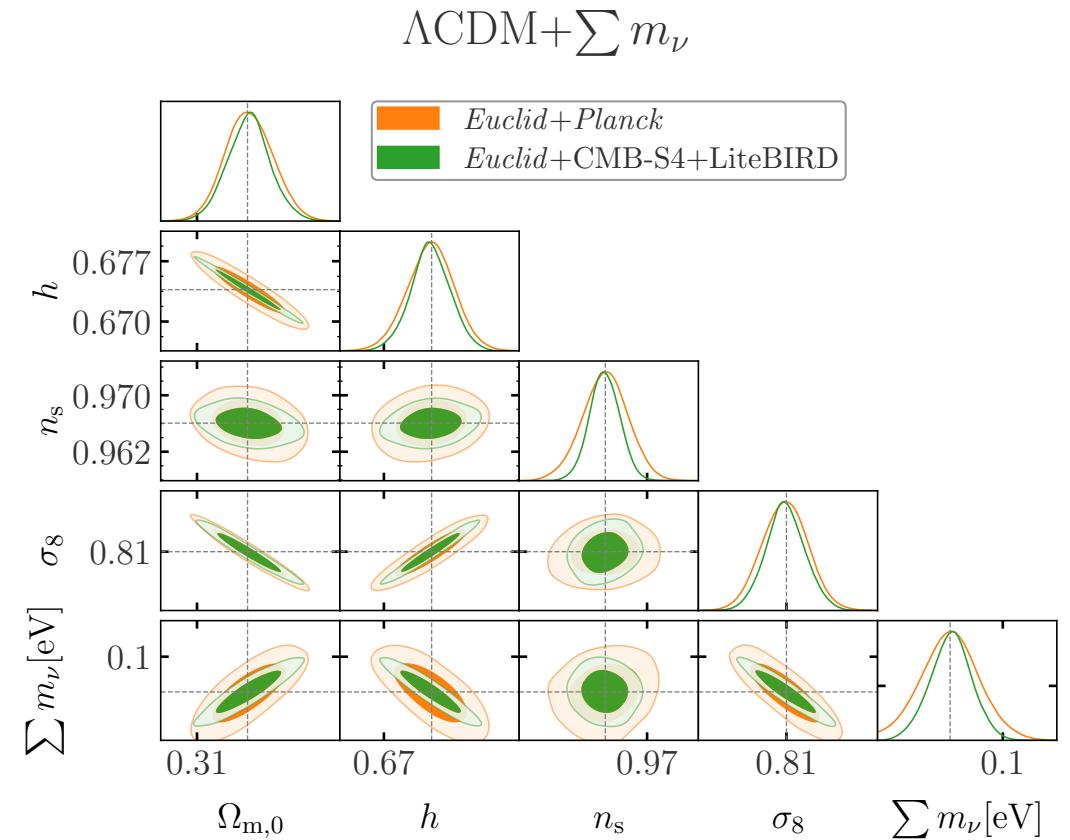
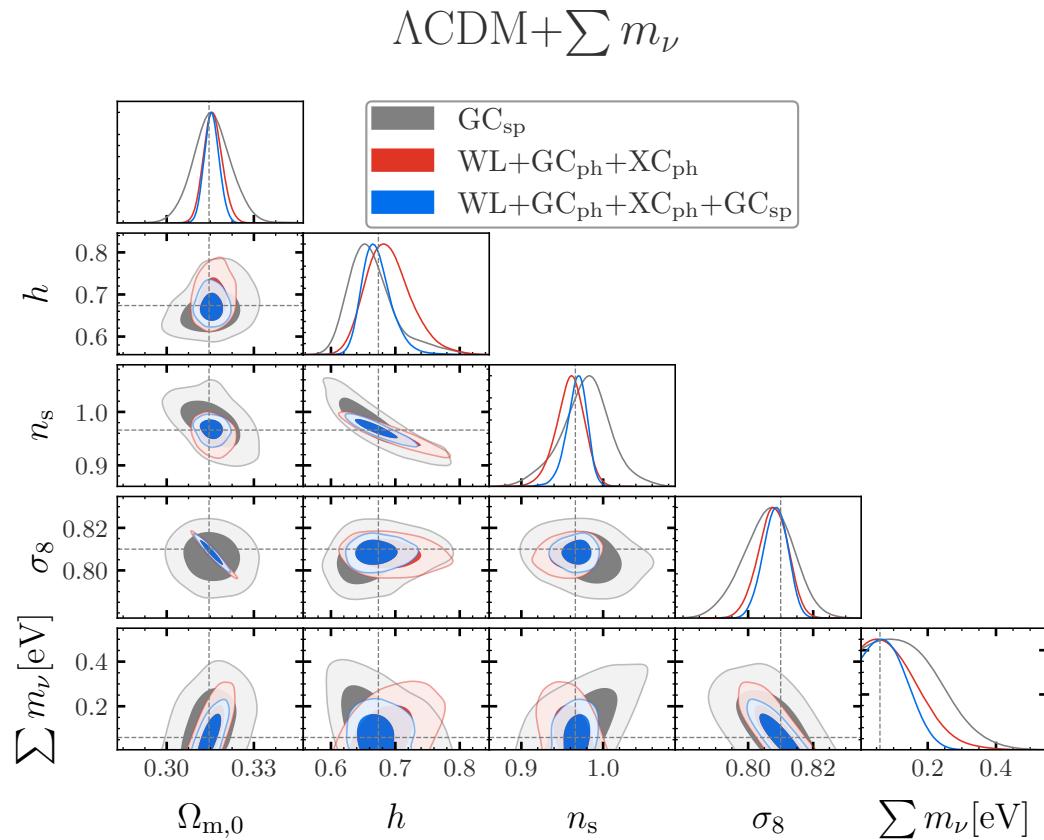


Neutrino mass constraints: the future

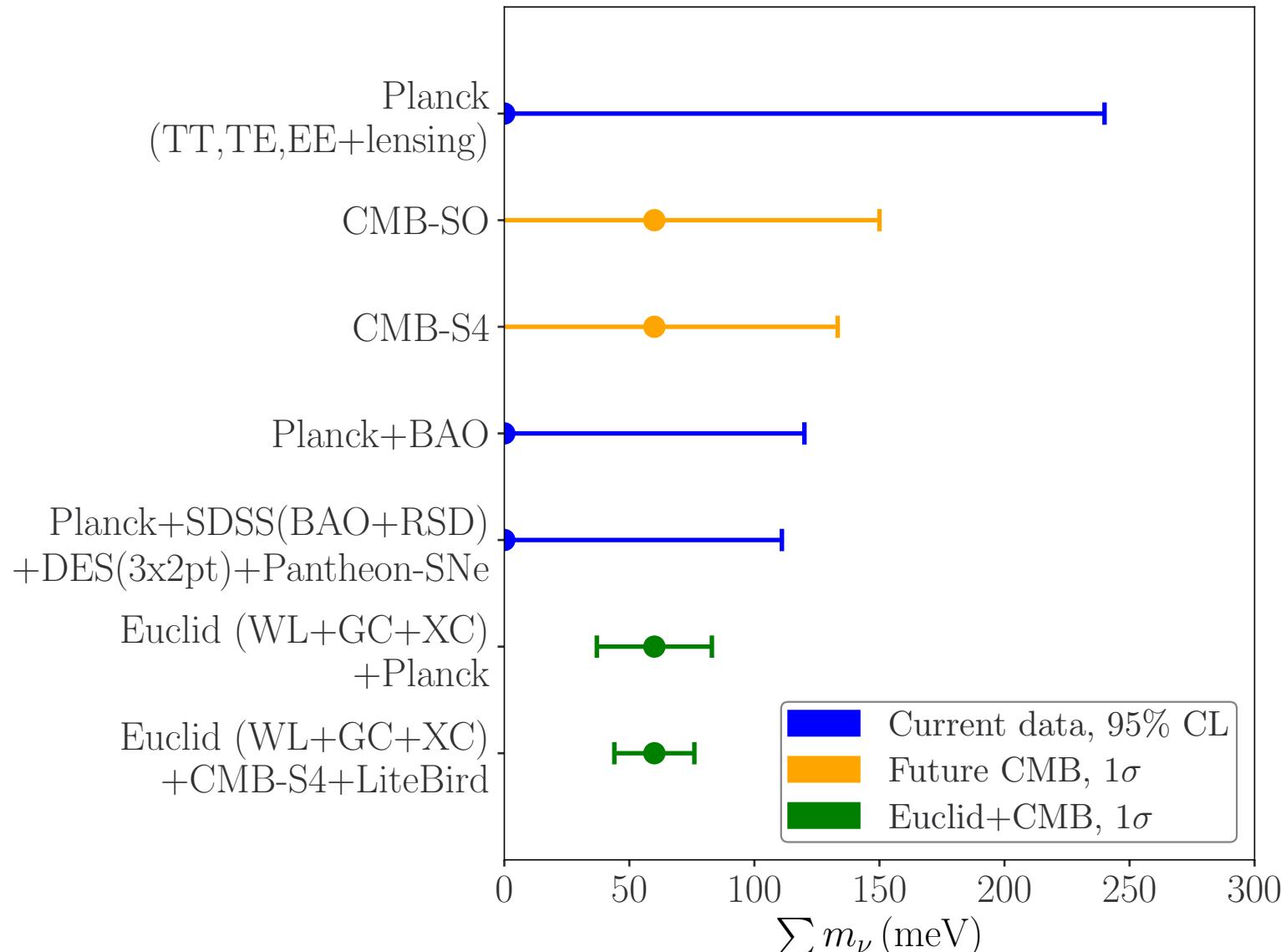
Euclid preparation: Sensitivity to neutrino parameters.

Euclid Collaboration: Archidiacono et al., arXiv:2405.13495

Note: MCMC forecast

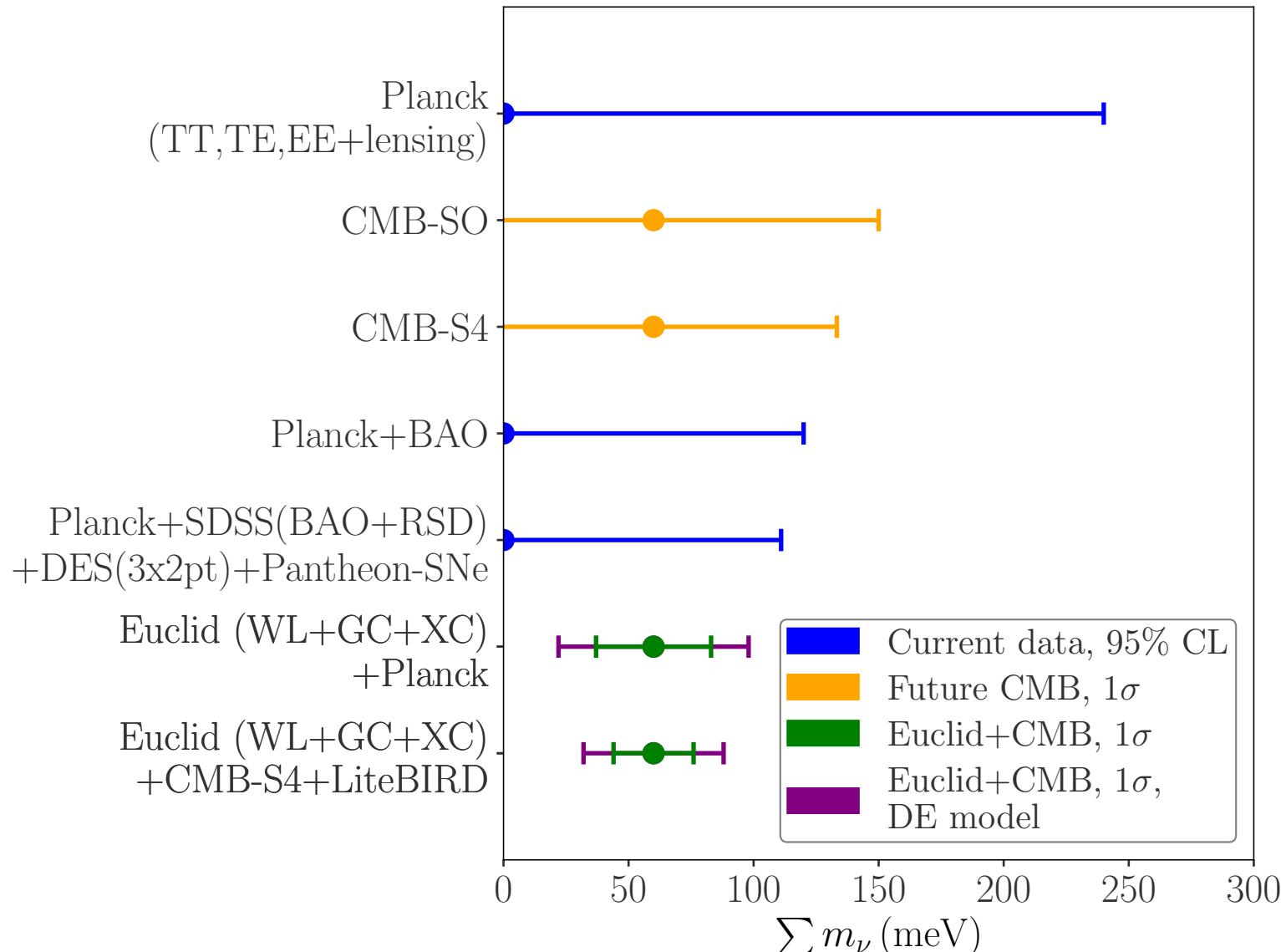


Neutrino mass constraints: the future



Euclid+Planck: >2 σ evidence of a
non-zero neutrino mass sum
Euclid+CMB-S4+LiteBird: >3 σ

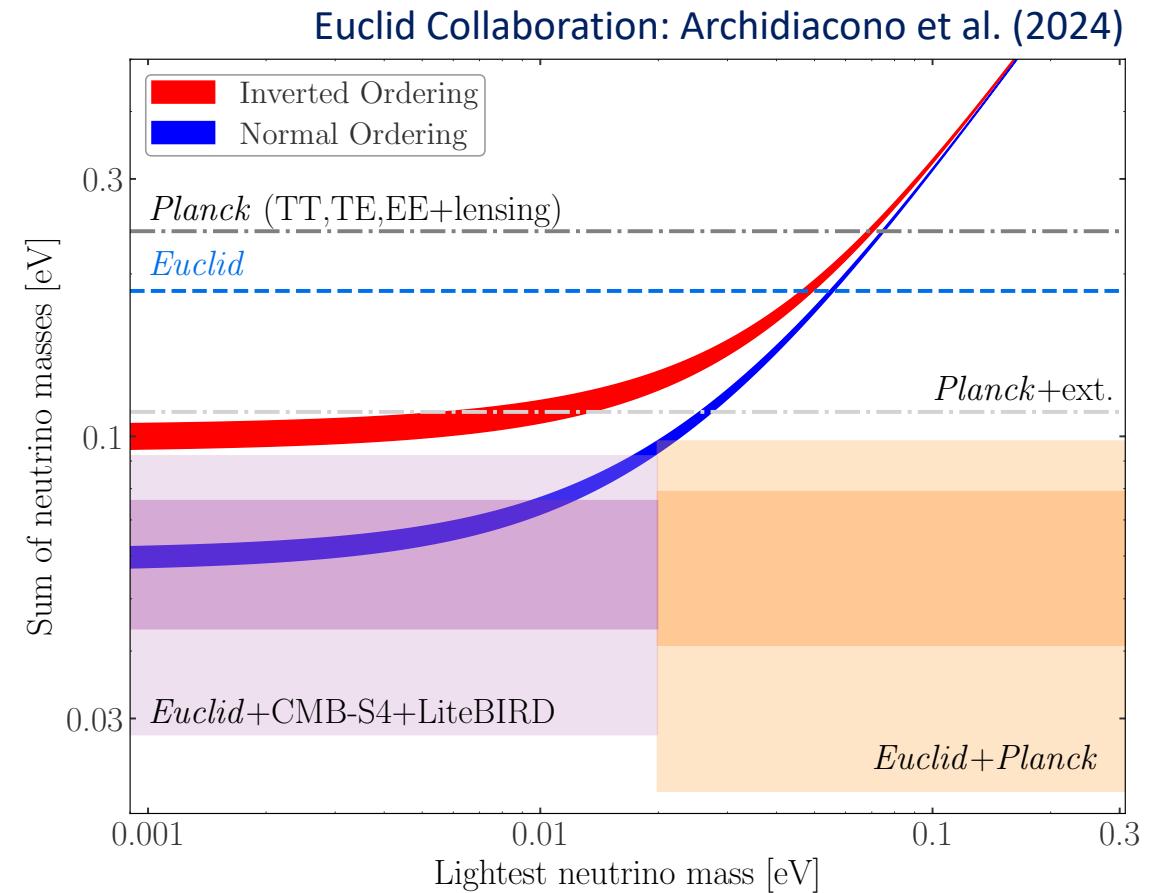
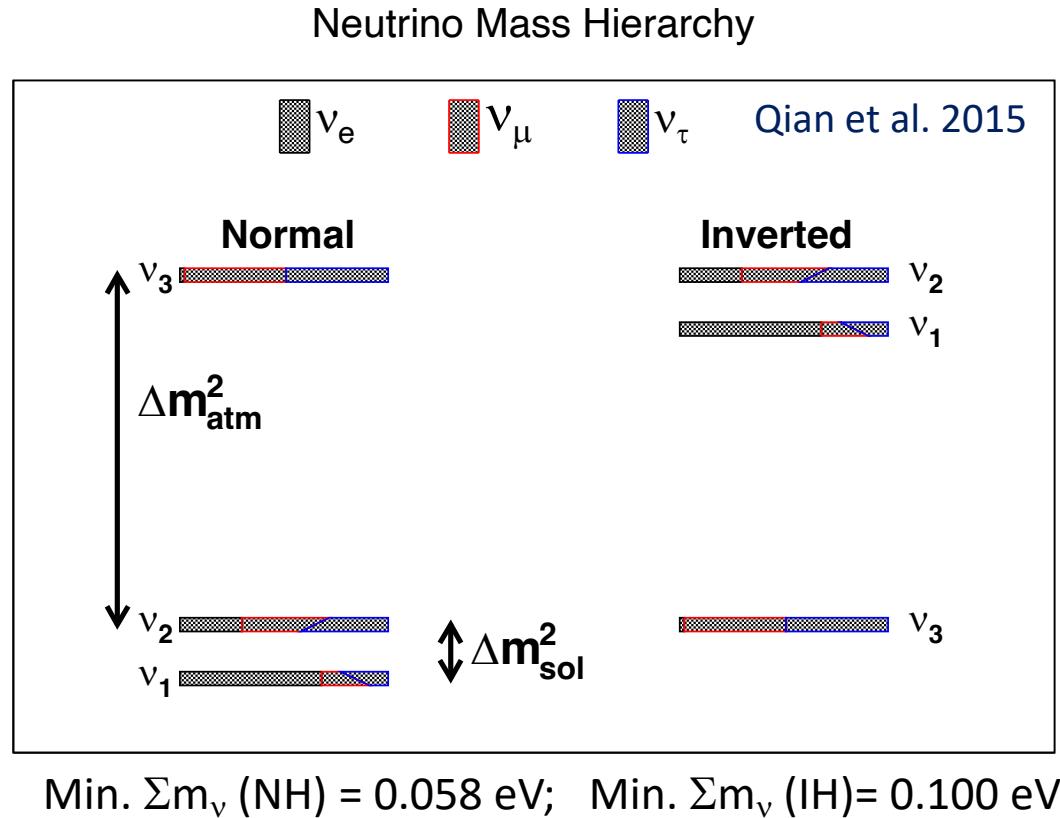
Neutrino mass constraints: the future



Replacing the cosmological constant with dark energy with a time varying equation of state parameter increases the error by a factor 2.

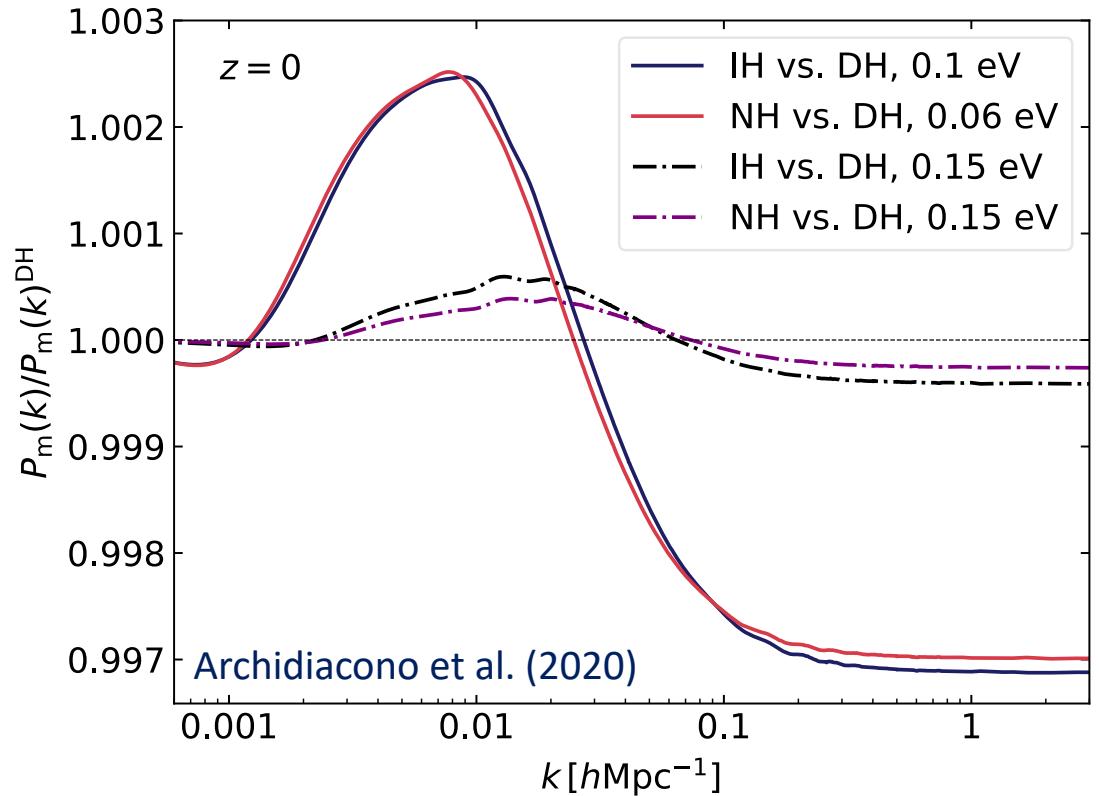
Neutrino mass ordering

Constraints derived under the assumption $m_1=m_2=m_3$ (degenerate hierarchy, DH)



Input fiducial value of the forecast $\Sigma m_\nu = 60$ meV

Neutrino mass ordering



The effect induced by the neutrino mass ordering on the cosmological observables is below the sensitivity of current and planned cosmological surveys.

The DH assumption is valid, and it is more efficient.

See also Gariazzo et al. (2022)

Neutrino mass: conclusions

- Euclid in combination with upcoming CMB surveys can achieve a 5σ **detection** of Σm_ν , even if $\Sigma m_\nu = 0.058$ eV (i.e., min. NH)
- Cosmology is not directly sensitive to the neutrino **mass ordering**, like DUNE, however if $\Sigma m_\nu = 0.058$ eV, then future cosmological constraints can exclude IH at about 2σ
- Cosmology is more sensitive than current and planned β -decay experiments. Caveat: cosmology is **model dependent**, and it requires that **systematic effects** are under control.
Complementarity: cosmology is not sensitive to the Dirac/Majorana nature, mixing angles.
- What if there is a tension between the Cosmos and the Lab?

Non-standard neutrinos – part 1

The cosmological neutrino mass problem

The cosmological neutrino mass problem

What if KATRIN (or Project 8) measures a neutrino mass in disagreement with cosmological bounds?

What if the cosmological bounds cross the minimum value allowed by oscillations?

→ How robust are the cosmological constraints on the neutrino mass? Can they be evaded?

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1. Beyond Λ CDM (replacing Λ with w0wa dark energy)

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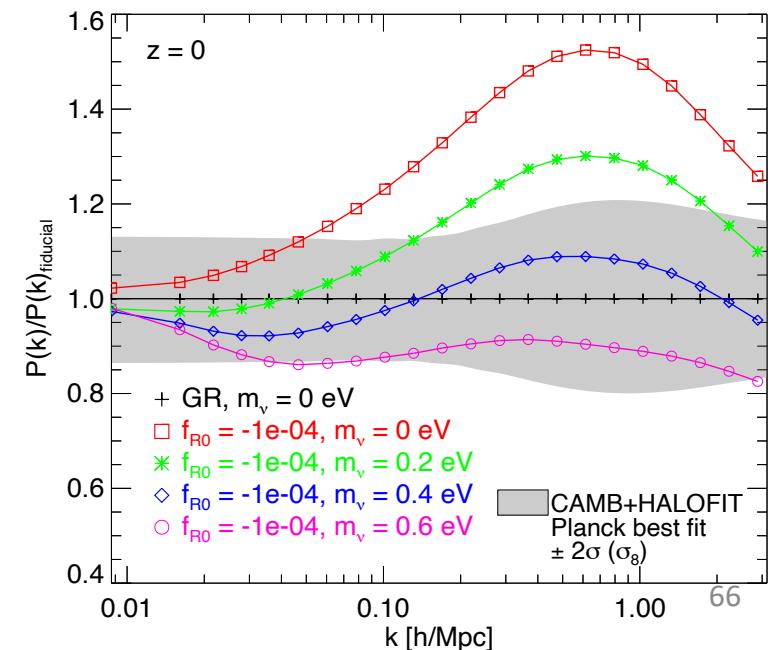
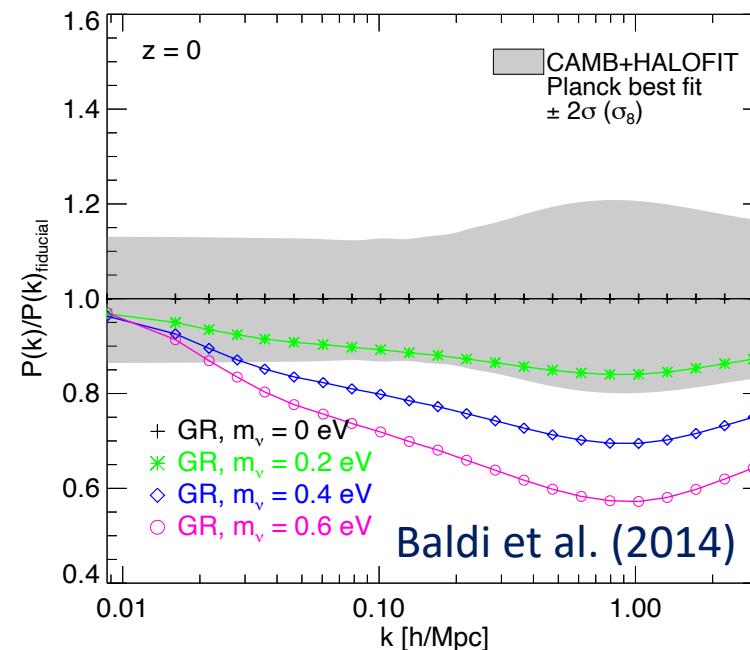
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The cosmological neutrino mass problem

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Yes...

1. Beyond Λ CDM (replacing Λ with w0wa dark energy)
2. Beyond GR
3. Beyond SM
 - Neutrino spectral distortions (from new interactions) [Alvey et al. (2022)] $\sum m_\nu < 3 \text{ eV}$
 - Mass varying neutrinos (late time mass generation) [Lorenz et al. (2021)] $\sum m_\nu < 1.5 \text{ eV}$
 - Invisible neutrino decay into BSM particles [Barenboim et al. (2021)] $\sum m_\nu < 0.2 \text{ eV}$

The cosmological neutrino mass problem

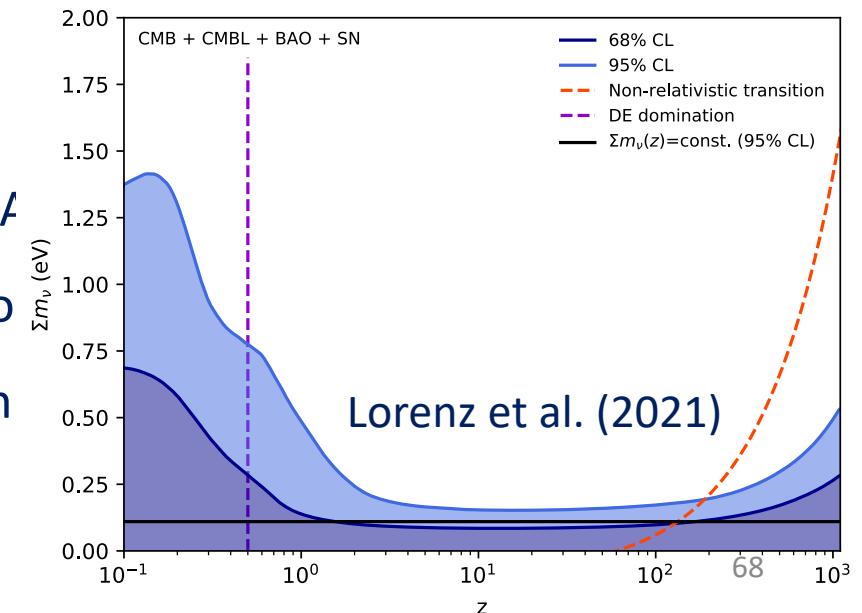
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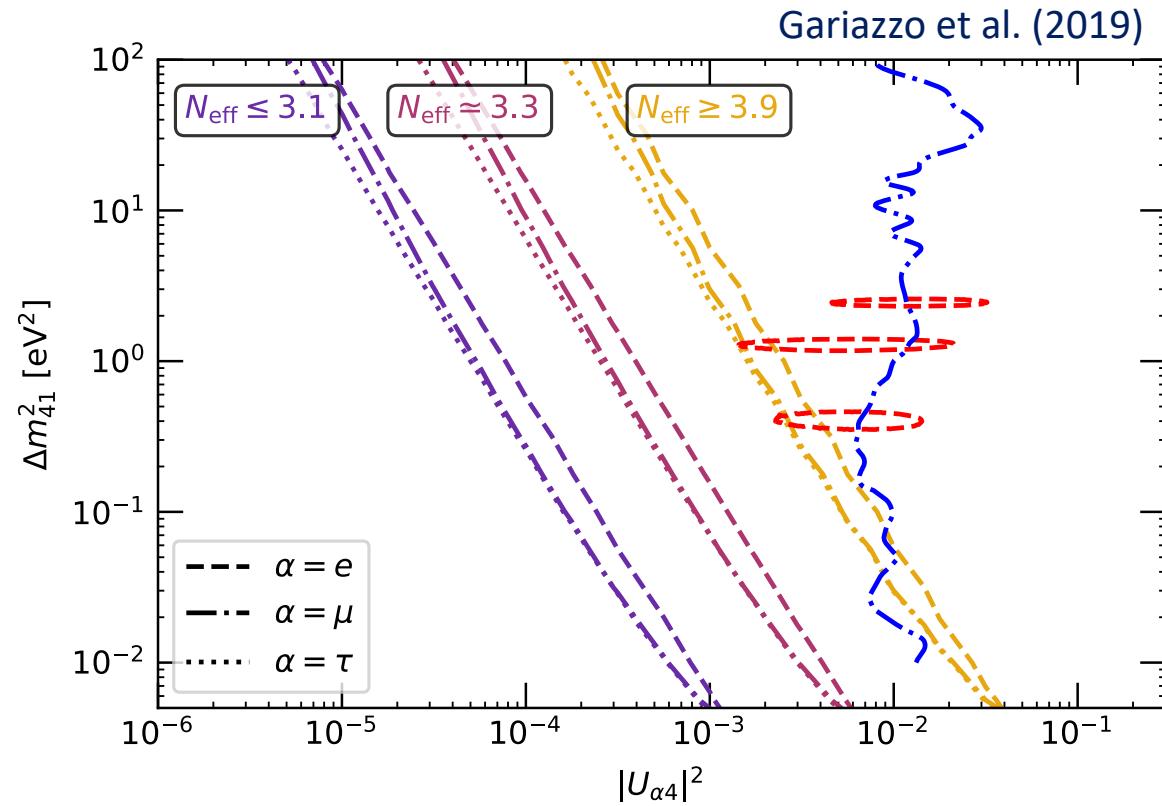
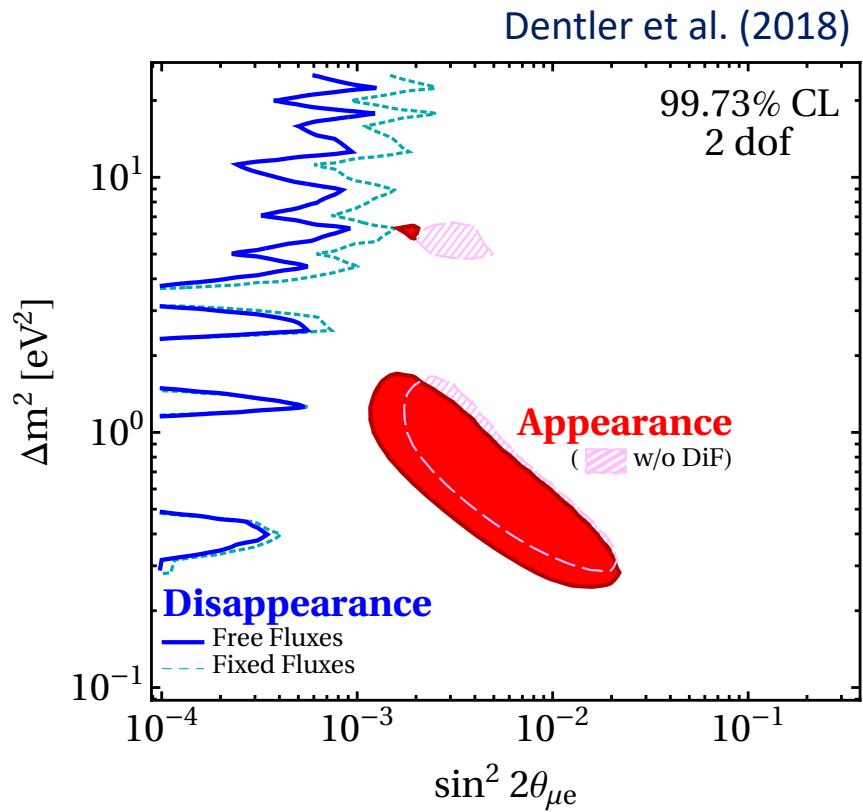
1. Beyond Λ CDM (replacing Λ with w₀w_a dark energy)
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3. Beyond SM
 - Neutrino spectral distortions (from new interactions) [A]
 - Mass varying neutrinos (late time mass generation) [Lo]
 - Invisible neutrino decay into BSM particles [Barenboim]



Non-standard neutrinos – part 2

Sterile neutrinos and new interactions

Light sterile neutrinos



Sterile neutrinos in cosmology are facing two problems:

- They are too many: $N_{\text{eff}} \sim 4$, while current bounds are $N_{\text{eff}} = 2.92 \pm 0.36$ (95% CL)
- They are too massive: $\sum m_\nu \sim 1$ eV, while current bounds are $\sum m_\nu < 0.1\text{-}0.2$ eV

Partial thermalization

$$N_{\text{eff}} = \frac{\rho_{\nu}^{\text{rel}}}{\rho_{\nu, m=0}^{\text{th}}}$$

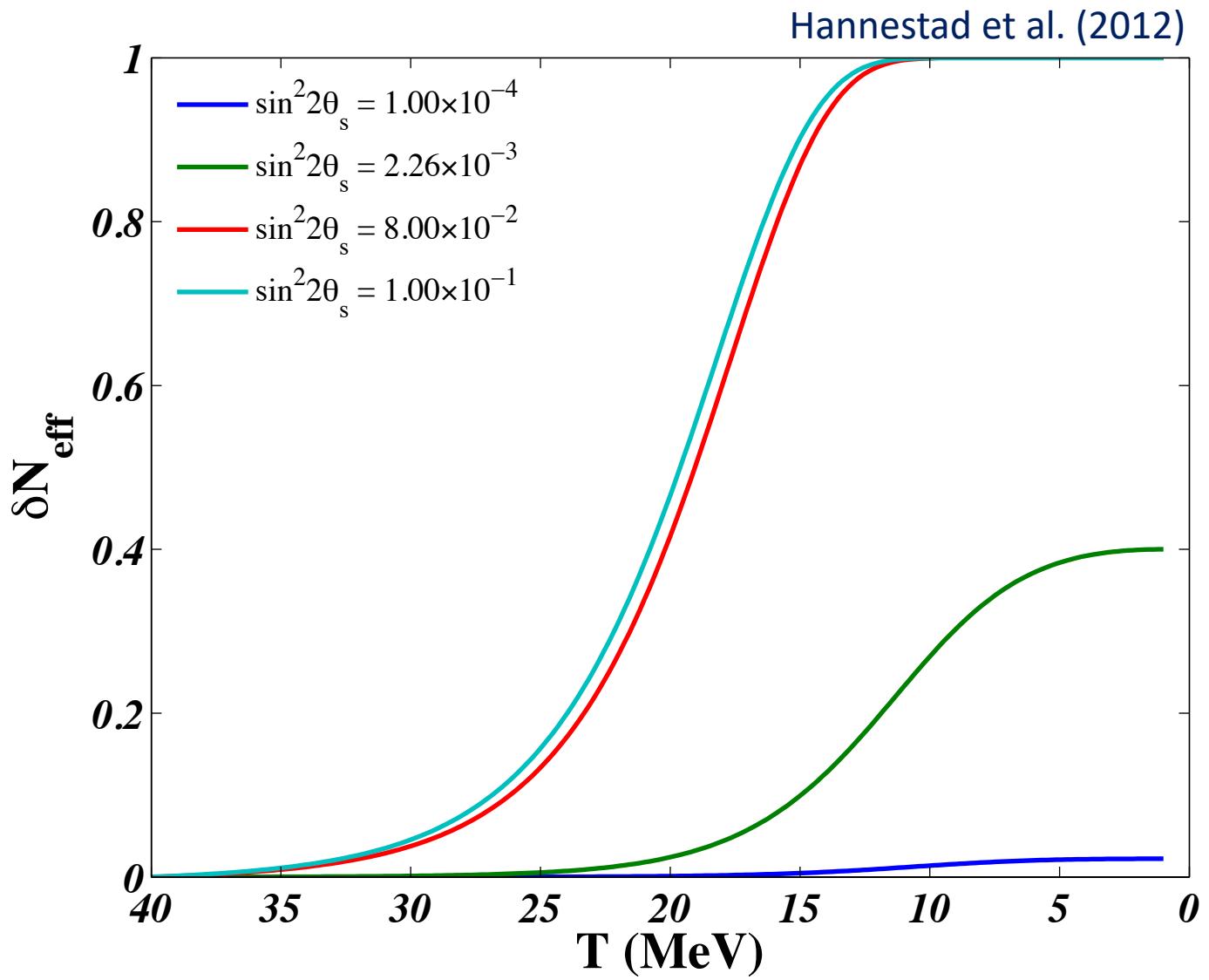
$$\rho_{\nu} = \frac{g}{2\pi^2} \int dp E p^2 f_{\nu}(p)$$

$$f_{\nu}^{\text{th}}(p) = \frac{1}{1 + \exp(E/T)}$$

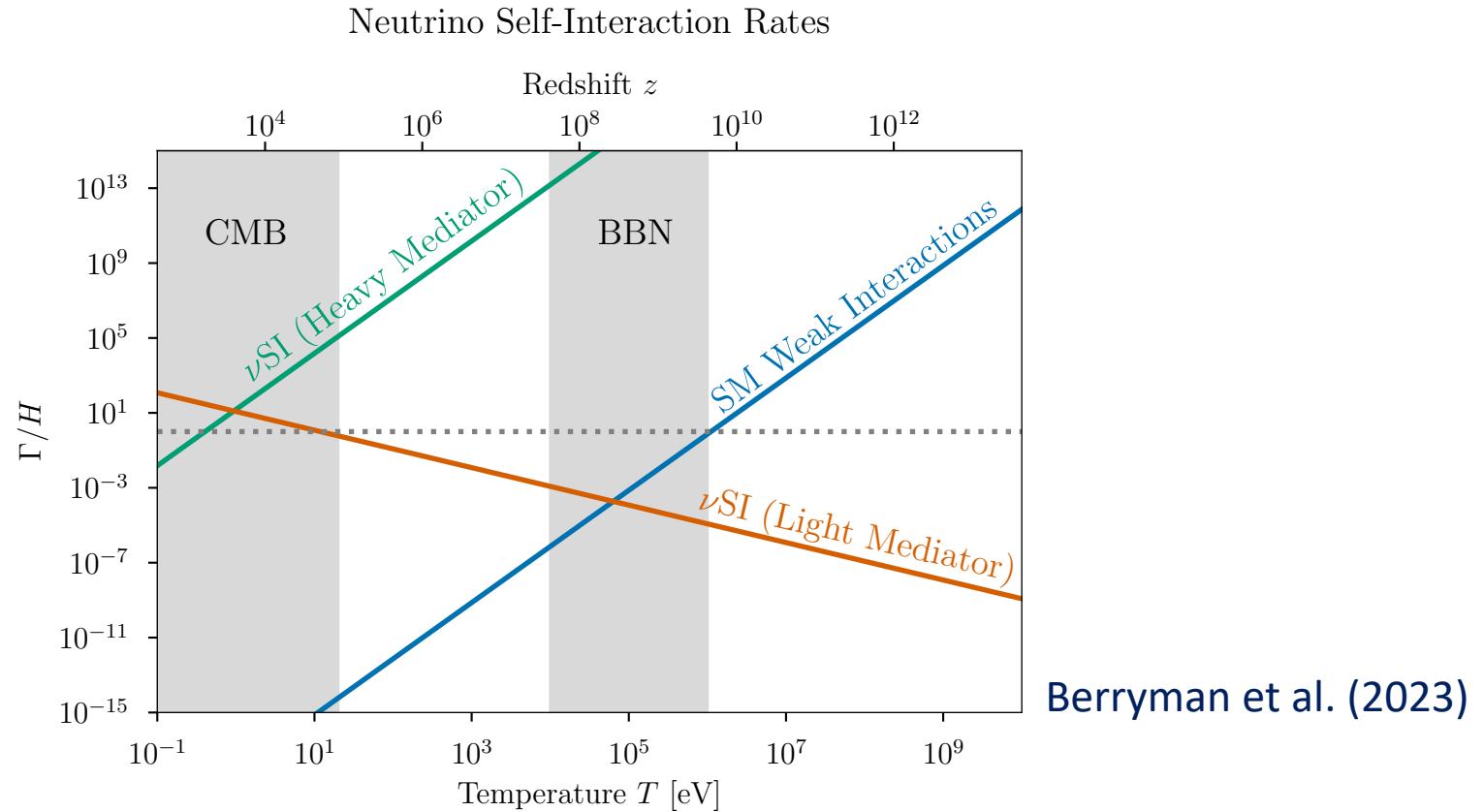
$$\sin^2 2\theta_m = \frac{\sin^2 2\theta_0}{\left(\cos 2\theta_0 + \frac{2E}{\Delta m^2} V_s\right)^2 + \sin^2 2\theta_0}$$

↓

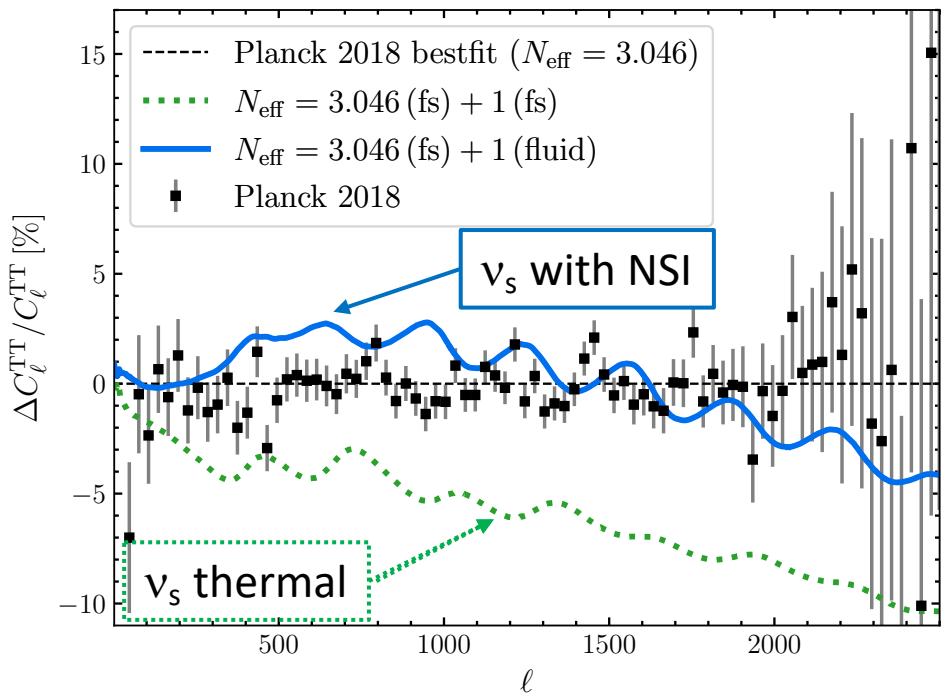
Additional matter-like potential



New interactions

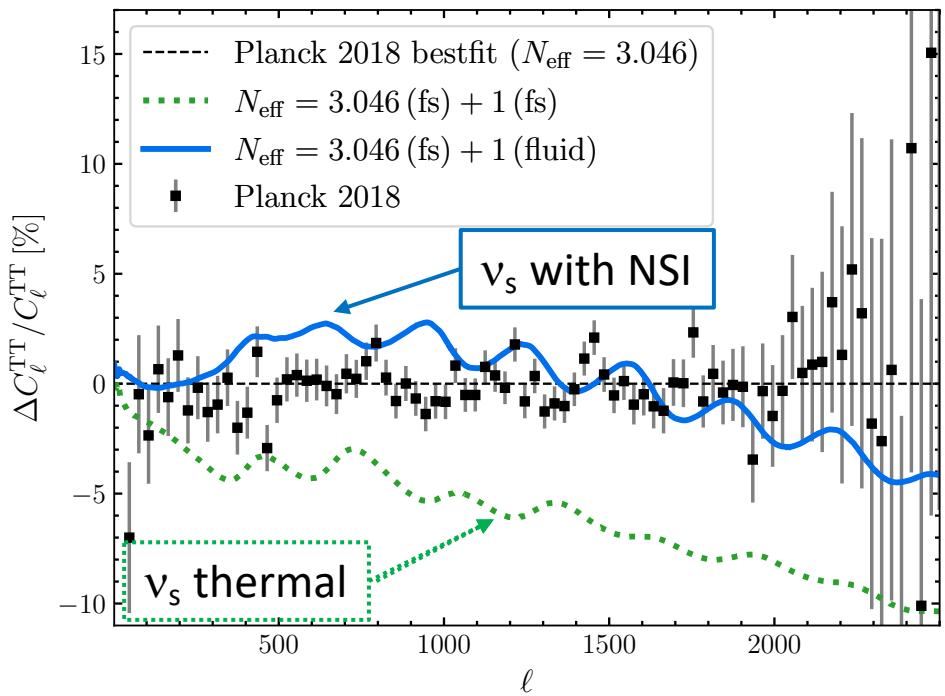


New interactions: CMB



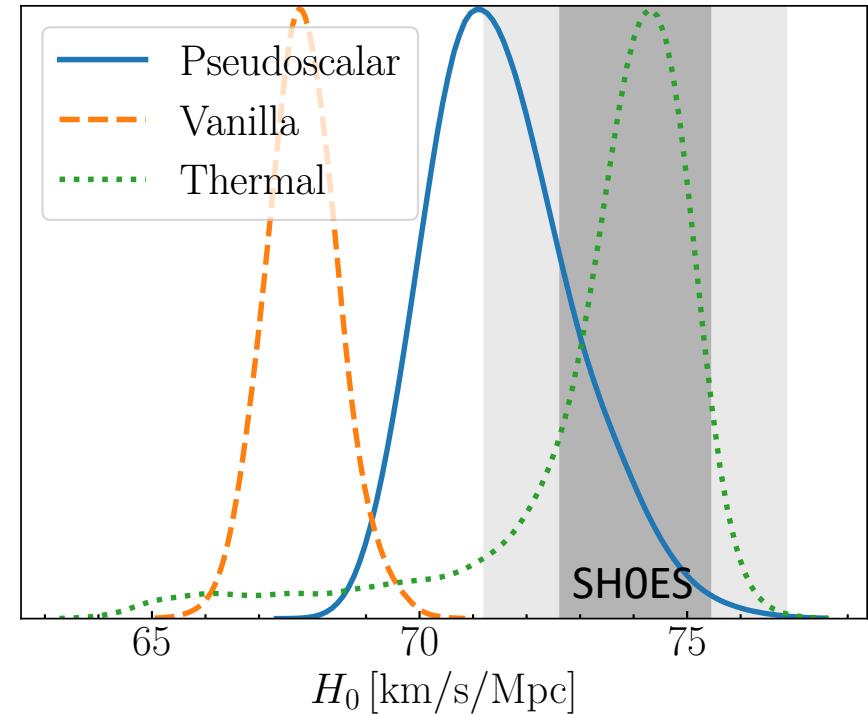
See also Kreisch et al. (2019)

New interactions: CMB

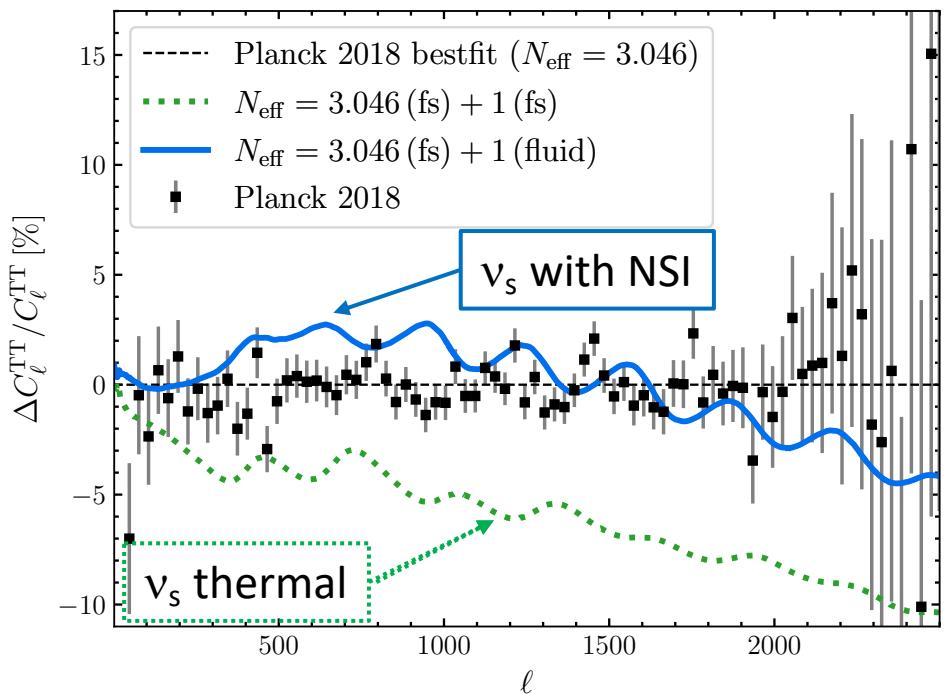


See also Kreisch et al. (2019)

Archidiacono et al. (2021)

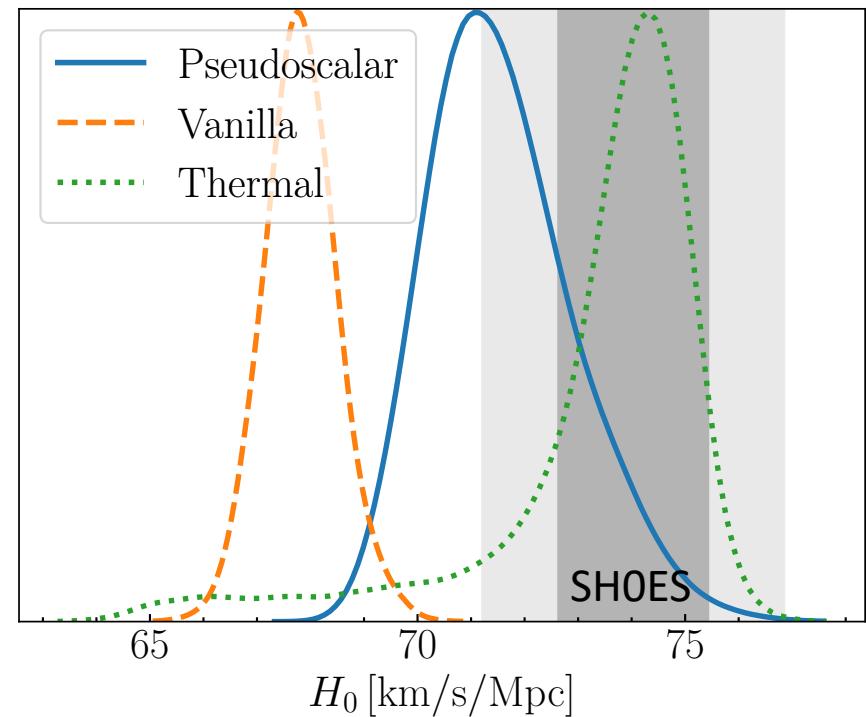


New interactions: CMB

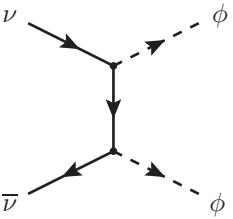


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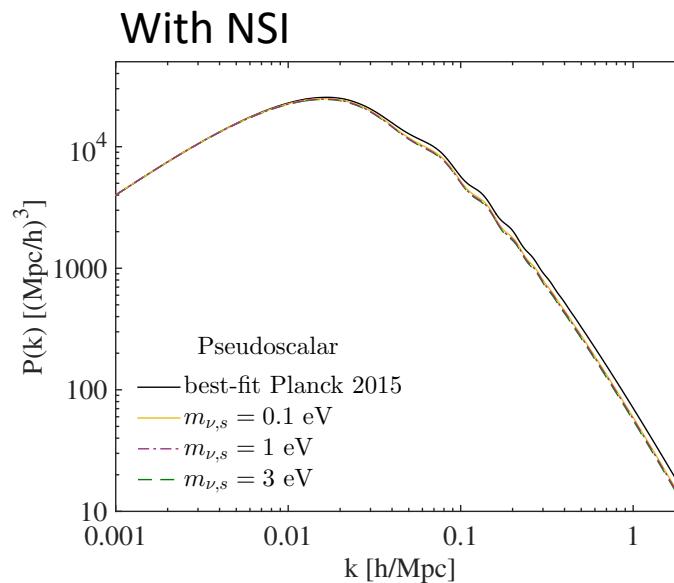
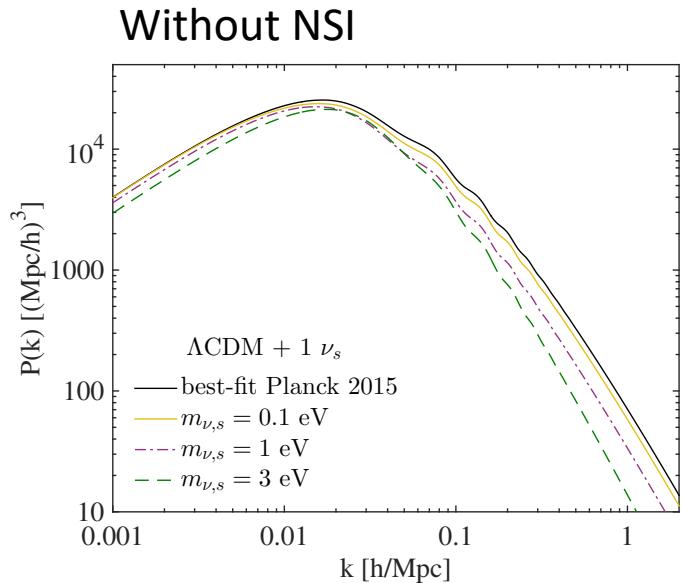
Archidiacono et al. (2021)



New interactions: LSS



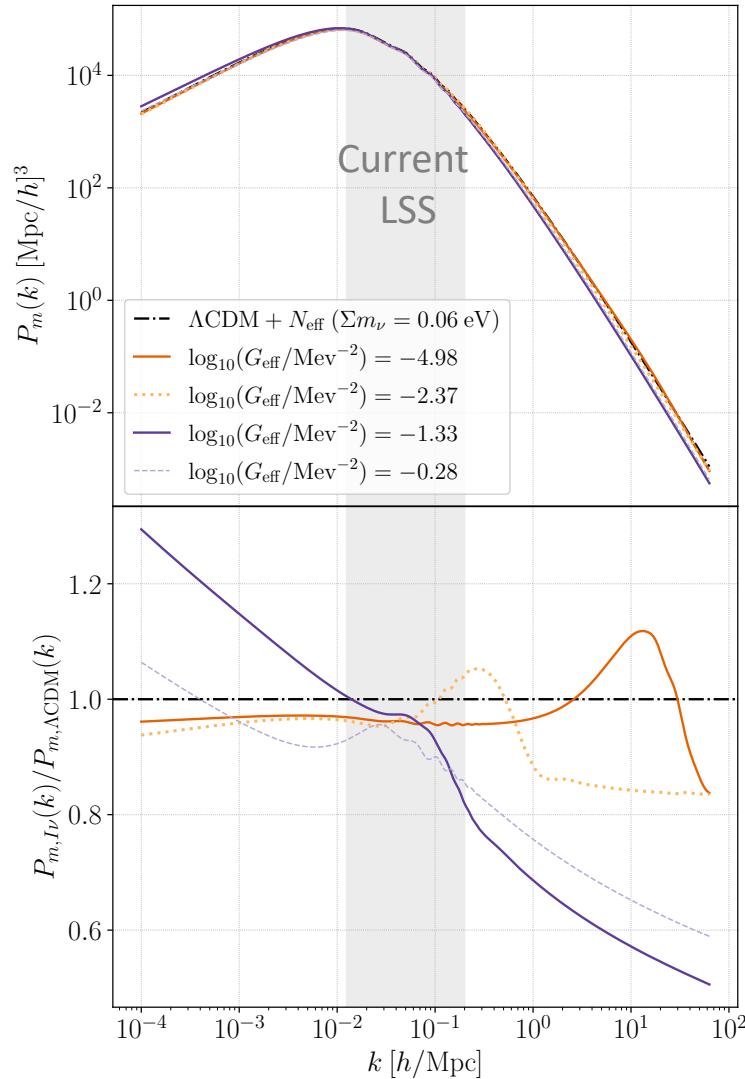
Sterile neutrinos disappear
from the cosmic neutrino
background.



✓ Sterile neutrinos are not too massive anymore

New interactions: LSS

Camarena et al. (2023)

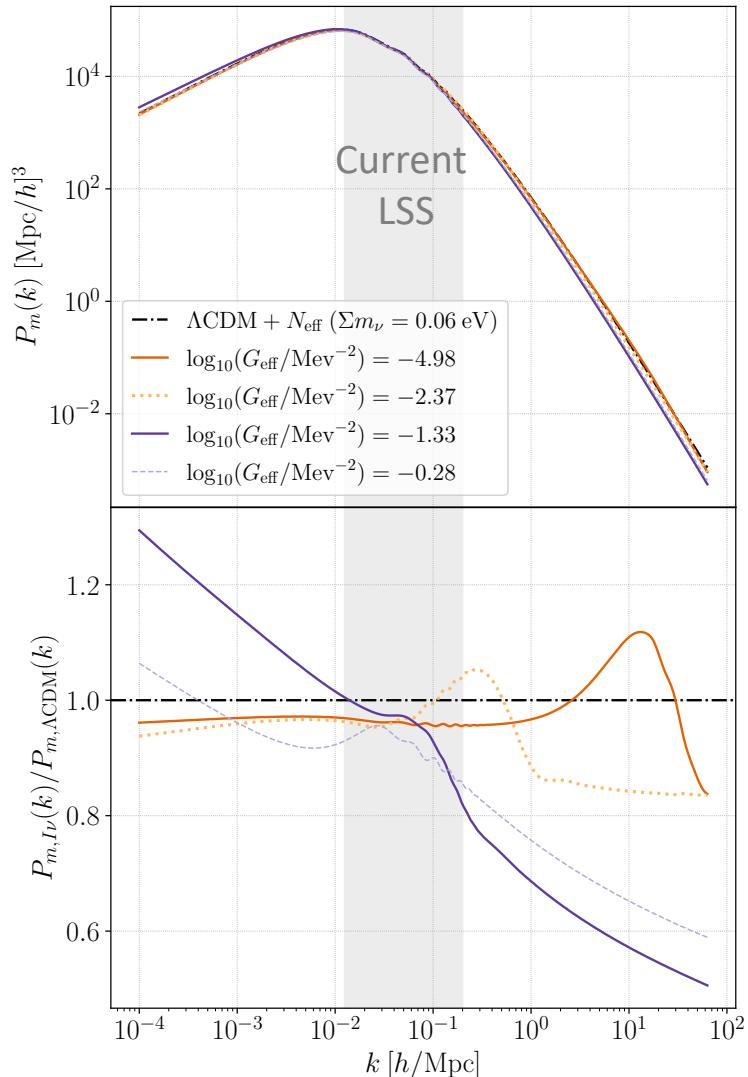


Mild preference for non-vanishing vSI

Stage IV LSS surveys will extend the range where we can test vSI

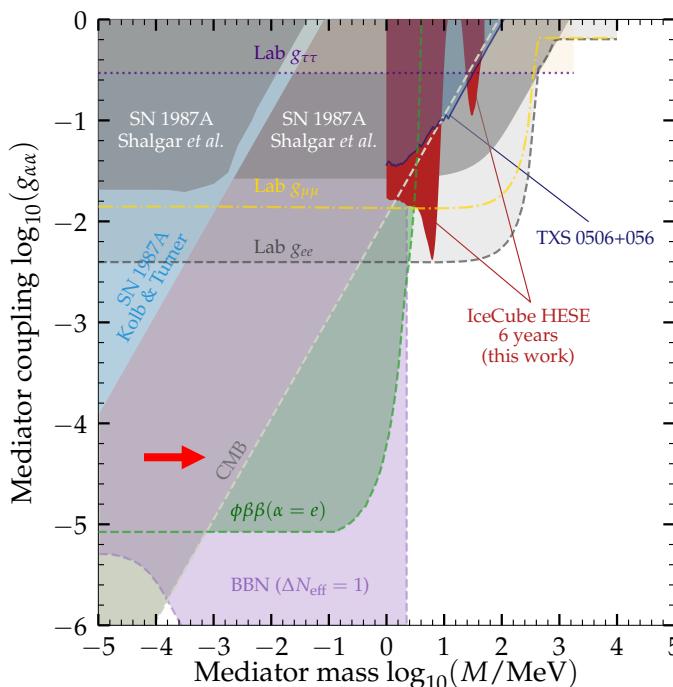
New interactions: LSS

Camarena et al. (2023)



Mild preference for non-vanishing vSI

Stage IV LSS surveys will extend the range where we can test vSI



Astrophysical and laboratory limits might already exclude this model.

Lectures by Coloma, Mirizzi, Sala, Spurio

Bustamante et al. (2020)

Non-standard neutrinos: conclusions

- Neutrino mass: The scenario of no detection, or a tension with ground-based experiments, would require to rethink the cosmological paradigm and/or neutrino physics.

Non-standard neutrinos: conclusions

- Neutrino mass: The scenario of no detection, or a tension with ground-based experiments, would require to rethink the cosmological paradigm and/or neutrino physics.
- Light sterile neutrinos, as hinted at by neutrino oscillation anomalies, are already excluded by Planck with high statistical significance. New neutrino self-interactions provide an elegant way to accommodate light sterile neutrinos in cosmology and to solve the H_0 problem.
- Non-standard interactions can be extended to active neutrinos, but external constraints might rule out their existence.

Useful references

- Textbook
 - J. Lesgourgues, G. Mangano, G. Miele and S. Pastor, *Neutrino Cosmology*
- Reviews
 - A. D. Dolgov, *Neutrinos in cosmology*, Phys. Rept. **370** (2002) 333 [hep-ph/0202122]
 - S. Bashinsky and U. Seljak, *Neutrino perturbations in CMB anisotropy and matter clustering*, Phys. Rev. D **69** (2004) 083002 [astro-ph/0310198]
 - J. Lesgourgues and S. Pastor, *Massive neutrinos and cosmology*, Phys. Rept. **429** (2006) 307 [astro-ph/0603494]

Backup

Neutrino flavour evolution

$$\rho(p,t) = \begin{pmatrix} \rho_{aa} & \rho_{as} \\ \rho_{sa} & \rho_{ss} \end{pmatrix} = \frac{f_0(p)}{2} [P_0(p,t) + \bar{\sigma} \times \bar{P}(p,t)];$$

$$\frac{d\bar{P}}{dt} = \bar{V} \times \bar{P} - D\bar{P}_T + \frac{R}{f_0} \hat{z}$$

$$\bar{V} = \bar{V}_{vacuum} + \bar{V}_{medium} + \bar{V}_s$$

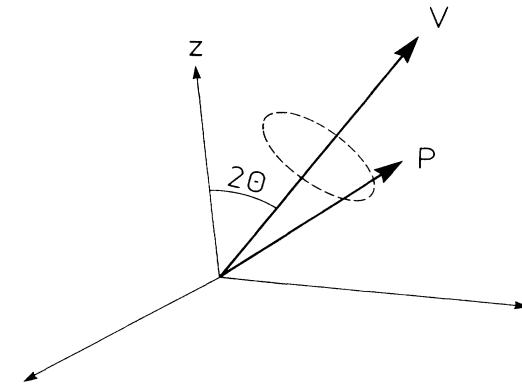
$$V_{vacuum} = \frac{\Delta m^2}{2p}$$

$$V_{medium} \propto \frac{G_F}{M_z^2} n_a p T^4$$

$$D = \frac{1}{2} \Gamma \quad \text{damping}$$

$$R = \Gamma \left(f_0 - \frac{f_0}{2} (P_0 + P_z) \right) \quad \text{repopulation}$$

$$\Gamma_a \propto G_F^2 p T^4$$



Stodolsky PRD (1987)

$$V_s(p_s) = \frac{g_s^2}{8\pi^2 p_s} \int p dp (f_\phi + f_s) \sim 10^{-1} g_s^2 T_s$$

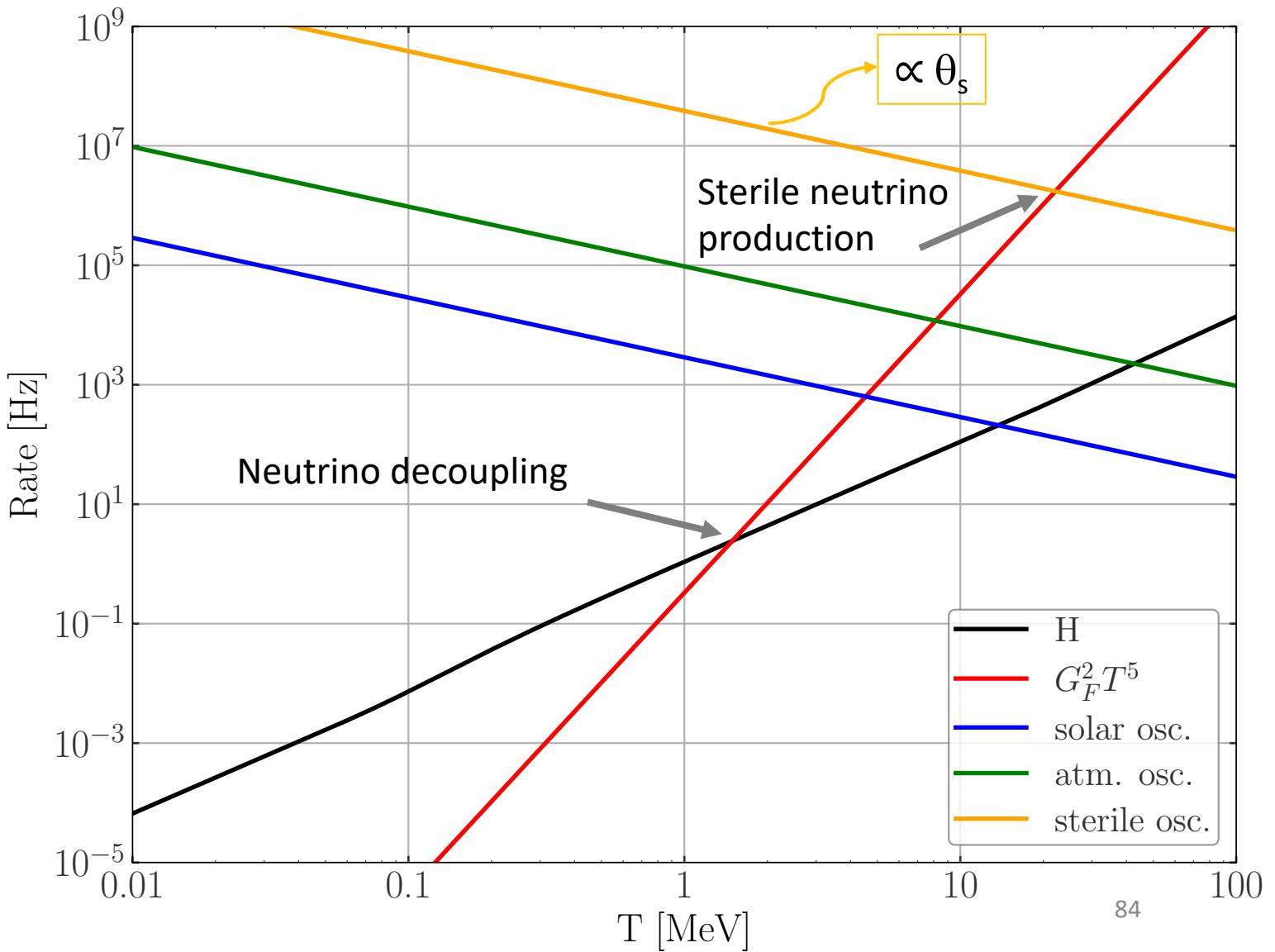
$$\Gamma_s = \frac{g_s^4}{4\pi T_s^2} n_s \quad \text{Babu PLB (1992)}$$

Partial thermalization

$$N_{\text{eff}} = \frac{\rho_{\nu}^{\text{rel}}}{\rho_{\nu, m=0}^{\text{th}}}$$

$$\rho_{\nu} = \frac{g}{2\pi^2} \int dp E p^2 f_{\nu}(p)$$

$$f_{\nu}^{\text{th}}(p) = \frac{1}{1+\exp(E/T)}$$



Light sterile neutrinos with NSI

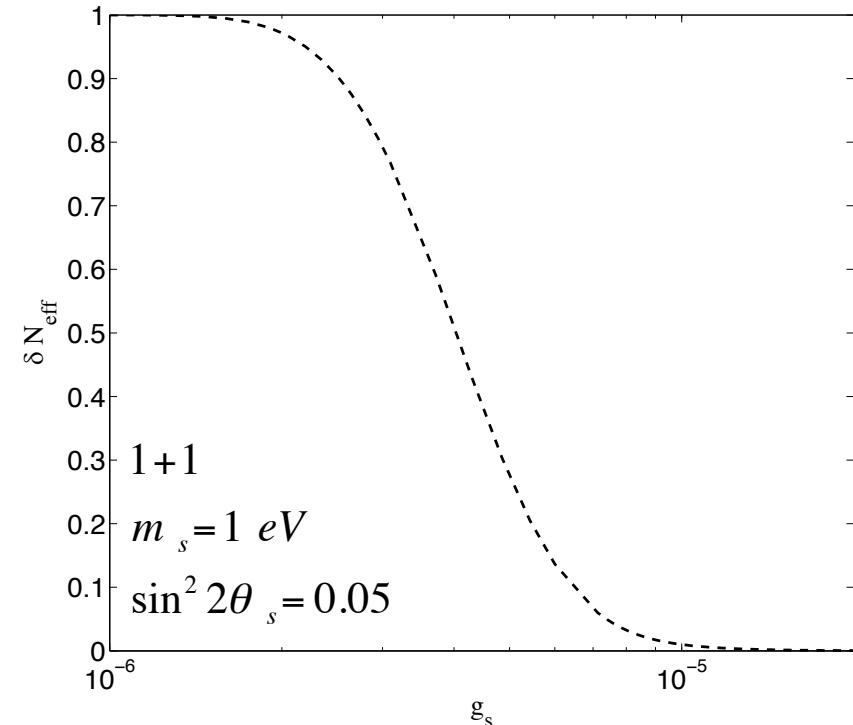
Early time phenomenology

The sterile neutrino is coupled to a new light pseudoscalar ($m_\phi \ll 1\text{eV}$):

$$L_{\text{int}} \sim g_s \phi \nu^{-1}_s \gamma_5 \nu_s$$

If the dimensionless coupling is larger than $g_s \sim 10^{-6}$, the production of sterile neutrinos is delayed until the time of active neutrino decoupling.

One additional sterile neutrino is consistent with the cosmological bounds on N_{eff}



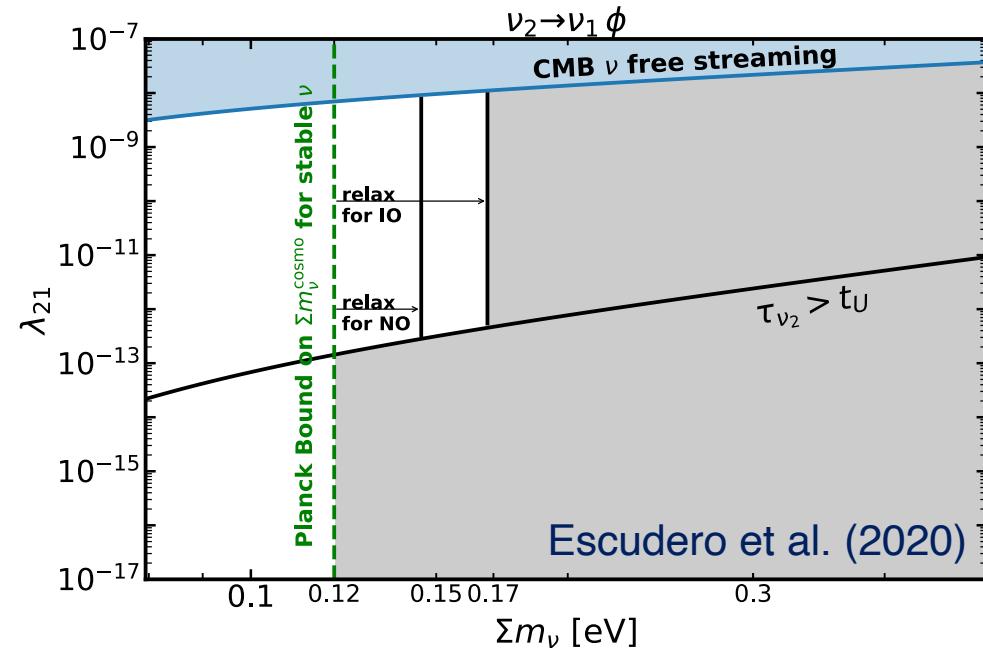
Sterile neutrinos are too many and too massive for cosmology X

BBN bounds

Neutrino mass problem

Cosmological data are more and more pointing towards $\sum m_\nu < 0.06$ eV.

- Extended particle physics models (beyond SM)
 - Invisible neutrino decay into BSM particles, e.g. lighter (sterile) neutrinos plus a massless (pseudo)scalar particle [Barenboim et al. 2021, Escudero et al. 2020]



KeV sterile neutrinos (WDM)

Boyarsky et al. (2019)

