

GW signatures of domain walls

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Outline

- Defects in field theory
- Formation in the early Universe
- Domain walls as seeds for bubble nucleation
- Domain walls as GW sources:
 - Impact of particle friction
 - Improved understanding of the scaling regime

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- Topological defects are non-trivial (space dependent) solutions of the EOM.
- Their classification is based on the topological properties of the vacuum manifold \mathcal{M} :



Fig. from Vilenkin & Shellard

Symmetry group G broken to subgroup $H \Rightarrow \mathcal{M} = G/H$ (coset).

• The type of defects that are supported depends on the non trivial homotopy group of \mathcal{M}



Space equivalent to
$$\mathscr{M} = S^1$$
 with non trivial $\pi_1(S^1) = \mathbb{Z}$

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• Domain walls correspond to disconnected vacuum manifolds: breaking of discrete symmetries



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• Domain walls correspond to disconnected vacuum manifolds: breaking of discrete symmetries



- A domain wall solution is characterized by the tension σ of the wall and its width δ



- Domain wall mass per unit surface:

$$\sigma \approx m_s v_s^2$$

$$+v_s$$

Formation of the network

- Domain walls form according to the Kibble mechanism:
 - Fluctuations of scalar field around T_c have finite correlation length $\xi(T) < d_H$
 - Uncorrelated patches will generally select different points of vacuum manifold \mathcal{M}

Zeldovich et al. 1975 Kibble 1976



Fig. From MIT edu

Walls as impurities

• The presence of domain walls at the time of a first order phase transition can induce exponentially enhanced nucleation on the surface



Seeded critical bubble

SB, Mariotti [2203.16450], PRL

Agrawal, **SB**, Mariotti, Nee [2312.06749], JHEP

> Same catalyzing effect can occur from axion strings: **SB**, Mariotti, [2405.08060]



Walls as impurities

 This scenario can be simulated from the h seeded nucleation (Langevin approach)



SB, Jinno, Konstandin, Rubira, Stomberg, JCAP [2302.06952]

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• This scenario can be simulated from the hydrodynamical point of view, as well as real time



Scaling regime

- regime, with $\mathcal{O}(1)$ domain walls per Hubble volume at any time
- In scaling, the energy density of the network grows compared to the critical density
- annihilates the network



• After formation, the network reaches a dynamical attractor solution known as the scaling

• This "domain wall problem" can be solved by a small energy or population bias that eventually

 $\frac{\rho_{\rm DW}}{\rho_c} \sim G \sigma t$



GWs from the scaling regime

- GWs are radiate by the domain walls during:
 - the scaling the regime (long-lasting source, dominated by later times)
 - the final phase of collapse and annihilation
- Assuming scaling ends at $T = T_*$:

$$\Omega_{\text{peak}} \sim 10^{-6} \alpha_*^2$$
, $f_{\text{peak}} \sim H(T_*)$





Domain wall motion from the Nambu-Goto action

$$S = -\sigma \int d^3 \zeta \sqrt{\gamma}$$

Parameterize a possible friction force by

$$F^{\nu} = \frac{\sigma}{\ell_{\rm f}} (u^{\nu} - x^{\nu}_{,a} \gamma^{ab} x^{\mu}_{,b} g_{\mu\sigma} u^{\sigma}$$

• Equation of motion for the surface:

$$\ddot{x} + \left(3\frac{\dot{a}}{a^2} + \frac{1}{\ell_{\rm f}}\right)\left(1 - \dot{x}^2\right)a\dot{x} = C$$



Curvature

Fig. from Martins, Rybak, Avgoustidis, Shellard [1602.01322]



• Define a total damping length given by:

Scaling regime

 $\overline{\ell_d}$



• The friction length can be evaluated given a particle physics model

$$\Delta P = P_{\rm R} - P_{\rm L} = 2 \int \frac{\mathrm{d}^2 p}{(2\pi)^3} \int_0^\infty \mathrm{d}p_z [f(-v) - f(v)] \frac{p_z^2}{E} \mathcal{R}(p)$$
Right and Left particle distribution in the wall frame
$$\mathcal{R}(p_z)$$

$$\frac{1}{\ell_{\rm f}} = \frac{\Delta P}{\sigma \gamma(v)v} \simeq \frac{\Delta P}{\sigma v}$$
R

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• Consider the case of fermions derivatively coupled to an axion-like-particle (ALP)

$$\mathcal{L}_{a\psi} = \frac{\kappa}{2N_{\rm DW}f_a} \partial_\mu a \,\bar{\psi}\gamma^\mu\gamma_5\psi + \bar{\psi}(i\partial\!\!\!/ - m_f)\psi$$

• Scattering off ALP domain walls:

$$V(a) = \Lambda^4 \left[1 - \cos\left(\frac{aN_{\rm DW}}{v_a}\right) \right]$$

 $a(z) = 4f_a \arctan\left(e^{m_a z}\right)$



Determine the reflection coefficient for the fermion scattering via Dirac equation



• ALPs coupled to **SM leptons** with strength $\sim 1/f_a$





- ALP domain walls destroyed by QCD bias: natural annihilation at $T_* \sim 100 \,\mathrm{MeV}$
 - Introduce a coupling of the ALP to the gluons:

$$\mathcal{L}_a \supset \frac{\alpha_s}{4\pi} \frac{N_c}{v} G\tilde{G}$$

- This will generate a potential which will act as a bias for the pre-existing DW network



- Does this coupling with the QCD sector also induce friction from the ALP domain walls around the time of collapse?
 - Crude approximations:
 - Gluon scattering for $T > 2 \,\text{GeV}$
 - Pion scattering for $T < 60 \,\mathrm{MeV}$

SB, Mariotti, Rase, Sevrin, [2302.06952], JCAP



• Summary plot for the NANOGrav 15yr parameter space of ALP domain walls:



- Friction can be relevant also for the late dynamics of the domain wall network
- When friction is relevant, possible new contribution to the GWs from the plasma?

 The corresponding implications for the GW emission have been studied within the velocityone-scale model implying a suppressed amplitude, no actual numerical simulation so far

GWs from the network collapse

- So far GW spectrum as given by the last moment of scaling prior to collapse
- Additional contribution from the actual phase of annihilation is expected
- This contribution can enhance the GW peak by 1-2 orders of magnitude

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Kitajima, Lee, Murai, Takahashi, Yin, PLB [2306.17146]
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Ferreira, Notari, Pujolàs, Rompineve, JCAP [2401.14331]
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Scaling regime

- - Systematic study in terms of initial fluctuations and $m_s/H(T_i)$
 - Infer the time evolution of ξ during the approach to scaling

Fig. From **SB**, Mariotti, Rase, Vanvlasselaer, in prep.

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• Simulations of domain wall networks support the scaling regime, but no systematic study of initial conditions (many simulations start with less than one wall per Hubble volume)



- A scaling network is characterized by a single scale set by the horizon size



• This argument has been used to prove the flat spectrum from scaling defects during radiation domination, and compared to explicit results from models with $O(N) \rightarrow O(N-1)$







 $\langle \Pi_{ij}^{\mathrm{TT}}(\mathbf{k},t) \Pi_{ij}^{\mathrm{TT}^*}(\mathbf{k}',t') \rangle$

• If the source is scaling Π^2 can only depend on $x_i = \kappa t_i$ up to trivial factors,

 $\Pi^2(k, t_1, t_2)$

• The GW spectrum then reads:

$$\frac{d\rho_{\rm GW}}{d\log k}(x,t) = \Omega_{\rm rad} \frac{4}{\pi} \frac{M_P^2 H_0^2}{a(t)^4} \left(\frac{v}{M_P}\right)^4 \left(\overline{F^T(x)}\right) = \text{const.}$$

• The argument is based on dimensional analysis on the UETC of the stress-energy tensor:

$$\rangle = (2\pi)^3 \Pi^2(k, t, t') \delta_D(\mathbf{k} - \mathbf{k'})$$

$$) = \frac{4v^4}{\sqrt{t_1 t_2}} C^T(x_1, x_2)$$

- Apply the same procedure to domain walls and study the scaling property of the UETC
- A similar (naive) ansatz for domain walls does not work:

$$\Pi^2(k,\tau_1,\tau_2) = \sigma^2(\tau_1\tau_2)^{1/2}C(x_1,x_2), \quad x = k\tau$$

 Need to include powers of the scale factor in the ansatz to account for the scaling in the DW network energy density

See also [2406.17053]

• Extract the scaling properties of the (equal the k^{-1} spectrum



- Extract the scaling properties of the (equal time) Π^2 from the simulation, possibly determine

Vanvlasselaer, in prep.

Summary

- Domain walls can be themselves a powerful source of gravitational waves, and can also act as seeds for catalyzed bubble nucleation
- Friction with the plasma can be relevant at late times. This regime has not been studied with simulations, no clear indication for the GW spectrum so far
- Still room to improve the understanding of the scaling regime: systematic study of the approach to scaling and semi-analytical studies of the UETC.

Thank you!