

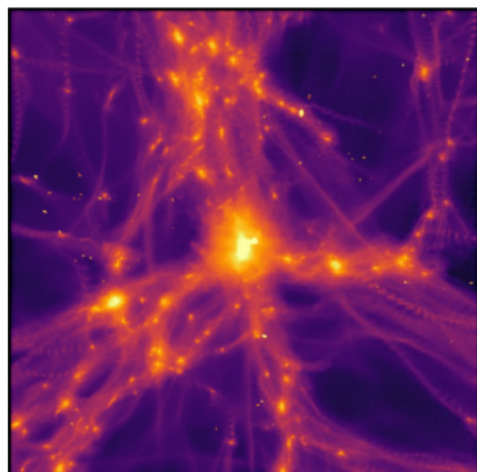
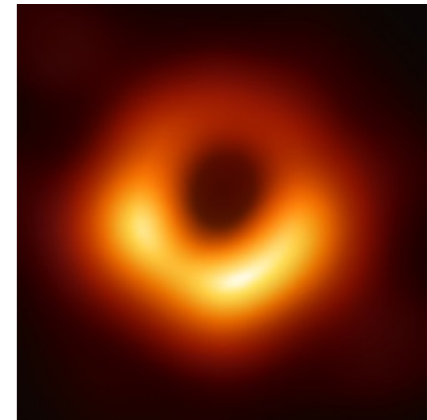
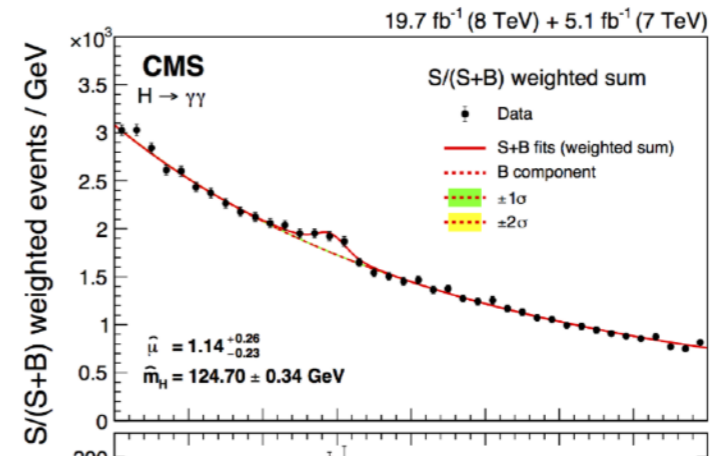
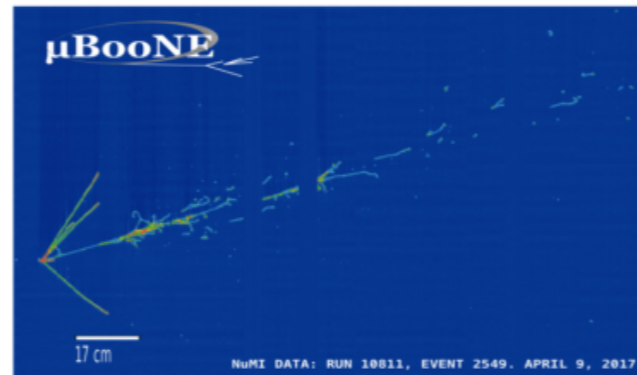
More Axion Stars from Strings

Edward Hardy



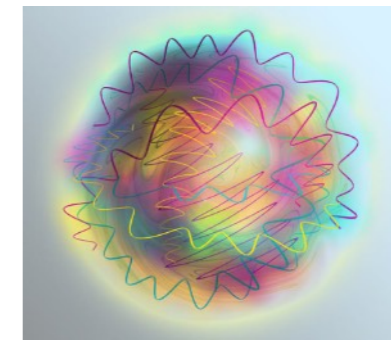
2405.19389 Gorghetto, EH, Villadoro

Dark Matter



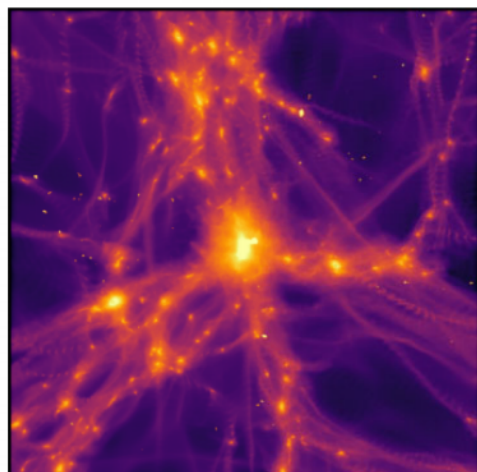
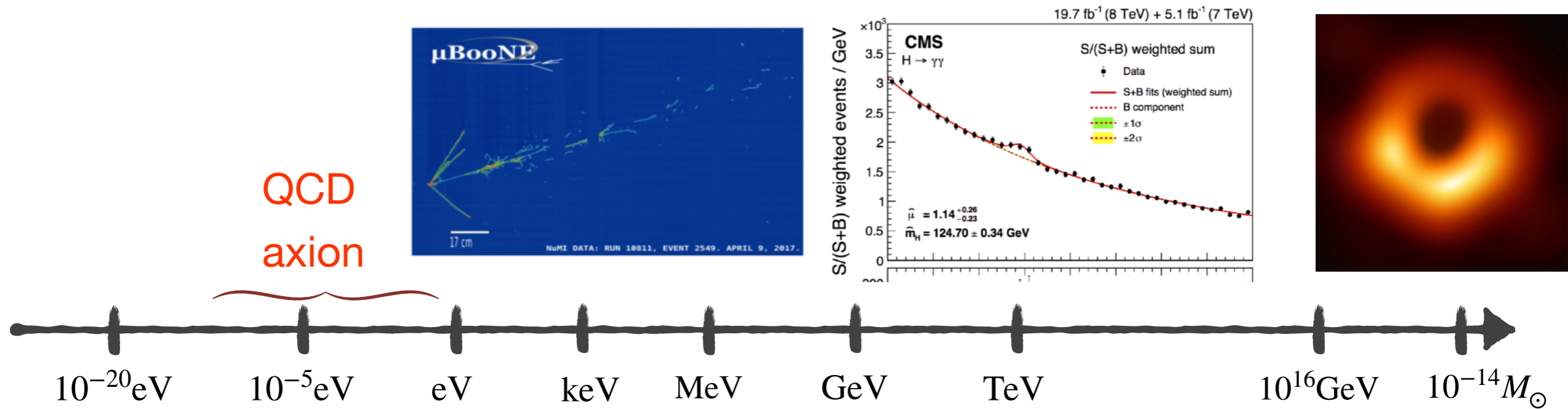
$\lambda_{d.B.}$

$$\frac{n_B - \bar{n}_B}{n_\gamma} \simeq 10^{-10}$$



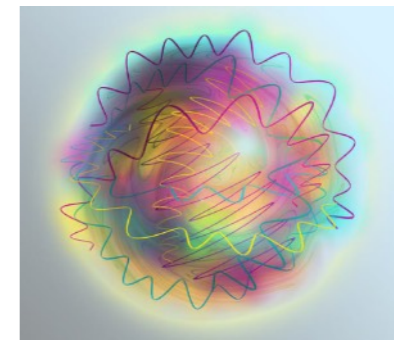
$$\Lambda_{UV} e^{-\frac{1}{2b_0 \alpha(\Lambda_{UV})}}$$

Dark Matter



$\lambda_{\text{d.B.}}$

$$\frac{n_B - \bar{n}_B}{n_\gamma} \simeq 10^{-10}$$



$$\Lambda_{\text{UV}} e^{-\frac{1}{2b_0 \alpha(\Lambda_{\text{UV}})}}$$

The QCD axion

$$\mathcal{L} \supset \theta_0 \frac{\alpha_S}{8\pi} G\tilde{G}$$

$$\theta' = \theta_0 + \arg(\text{Det}M_q) \lesssim 10^{-10}$$

Strong CP problem

The QCD axion

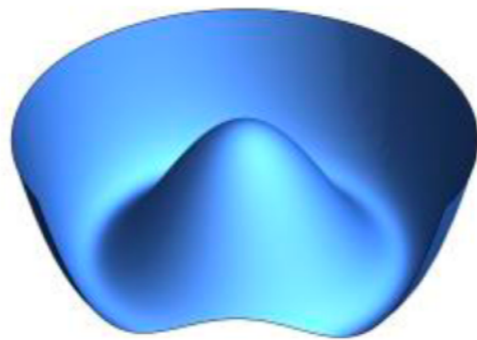
$$\mathcal{L} \supset \theta_0 \frac{\alpha_s}{8\pi} G\tilde{G} \quad \theta' = \theta_0 + \arg(\text{Det}M_q) \lesssim 10^{-10}$$

Strong CP problem

- Axion a , shift symmetry $a \rightarrow a + c$
- Candidate axions generic in high energy theories

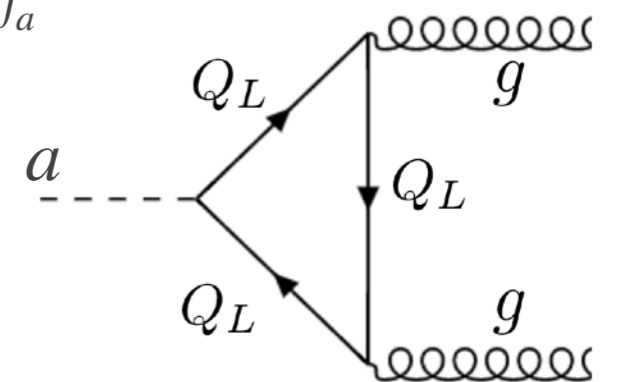
$$\frac{a}{f_a} \frac{\alpha_s}{8\pi} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$\mathcal{L} = (\partial\phi)^2 - \frac{m_r^2}{2f_a^2} \left(|\phi|^2 - \frac{f_a^2}{2} \right)^2$$



$$\phi = \frac{f_a + r}{\sqrt{2}} e^{ia/f_a}$$

$$\theta \equiv a/f_a$$



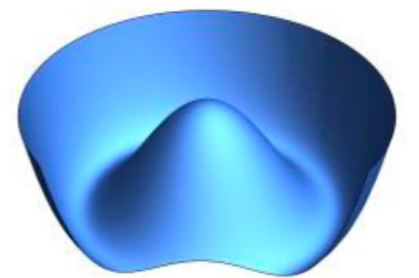
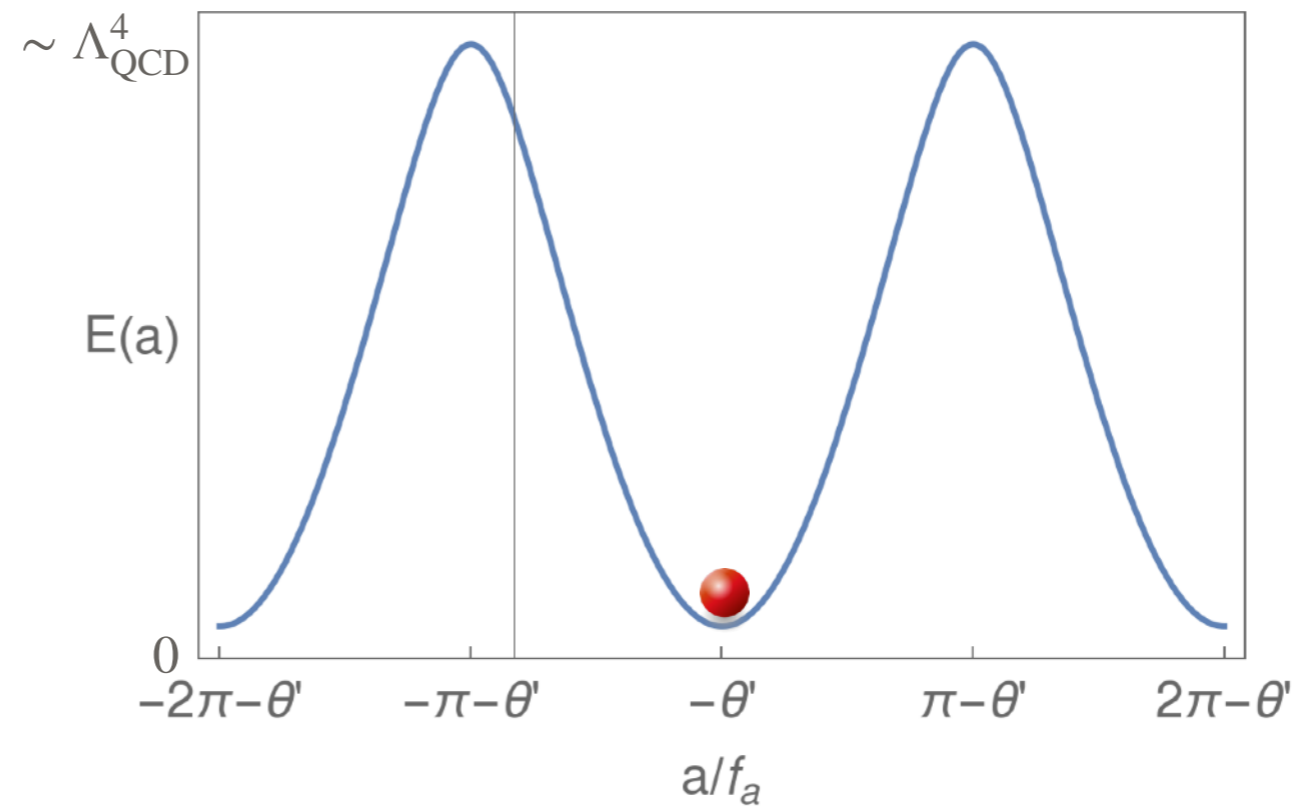
The QCD axion

$$\frac{a}{f_a} \frac{\alpha_s}{8\pi} G_{\mu\nu} \tilde{G}^{\mu\nu}$$



$$E(a) \geq E(a = -\theta')$$

$$\theta_{\text{tot}} = \langle a \rangle + \theta' = 0$$



[Vafa, Witten]

The QCD axion

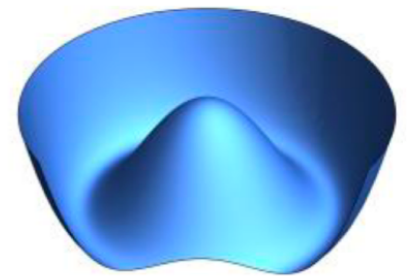
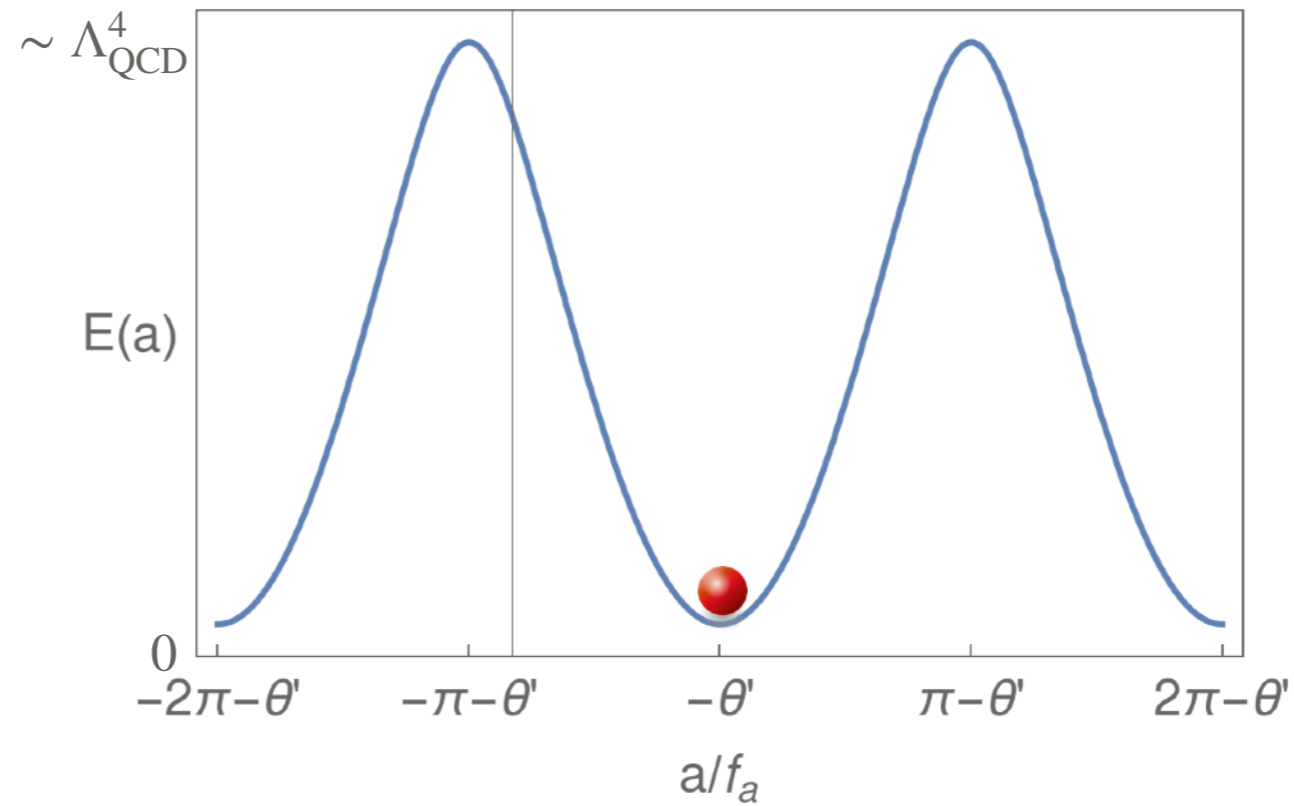
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$$E(a) \geq E(a = -\theta')$$

$$\theta_{\text{tot}} = \langle a \rangle + \theta' = 0$$

$$m_a^2 = \frac{m_u m_d}{(m_u + m_d)^2} \frac{m_\pi^2 f_\pi^2}{f_a^2} + \dots$$

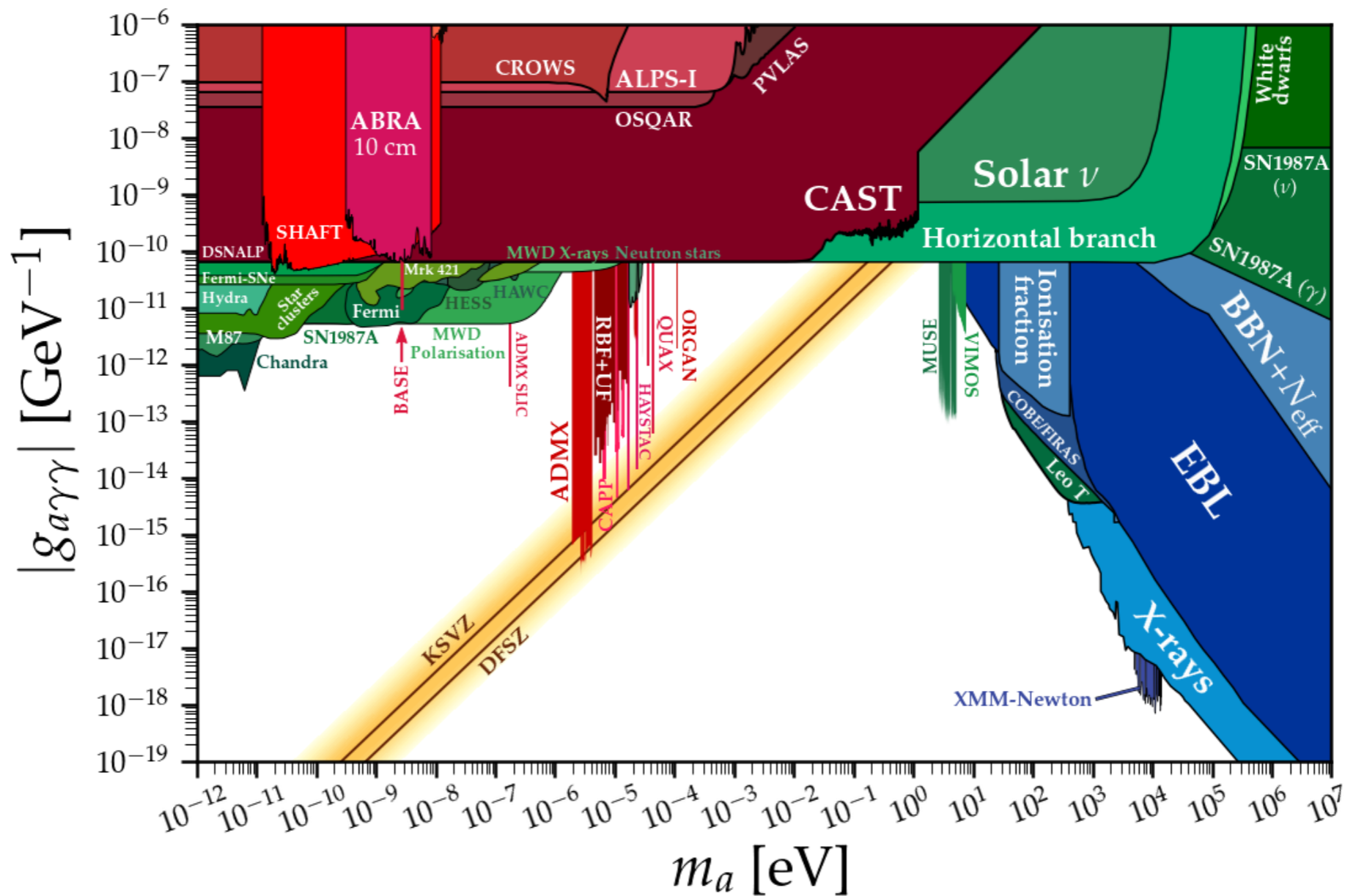


[Vafa, Witten]

$$m_a = 5.70(6)(4) \mu\text{eV} \left(\frac{10^{12} \text{GeV}}{f_a} \right)$$

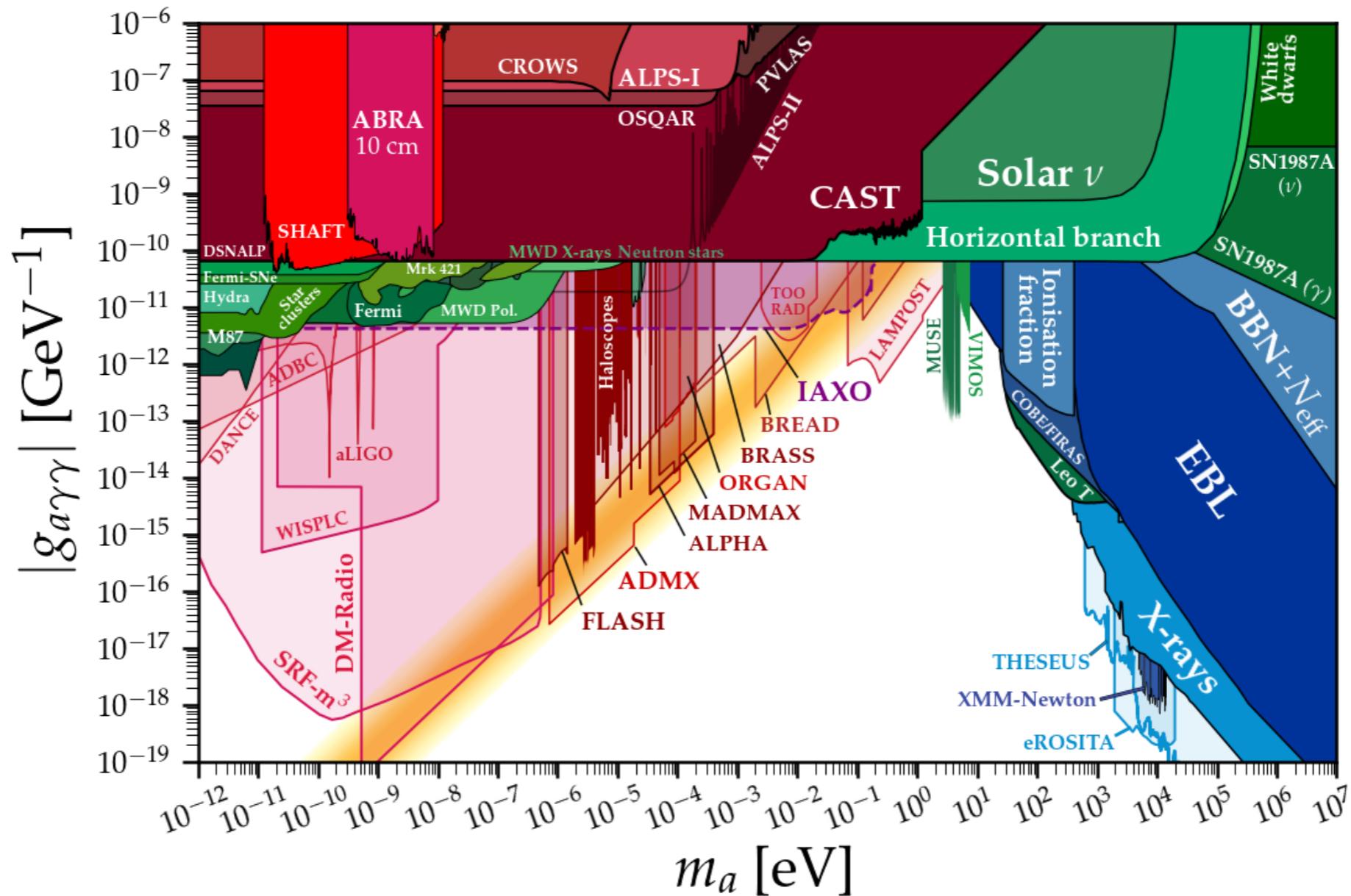
Searches

$$\mathcal{L} \supset -\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



Searches

$$\mathcal{L} \supset -\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

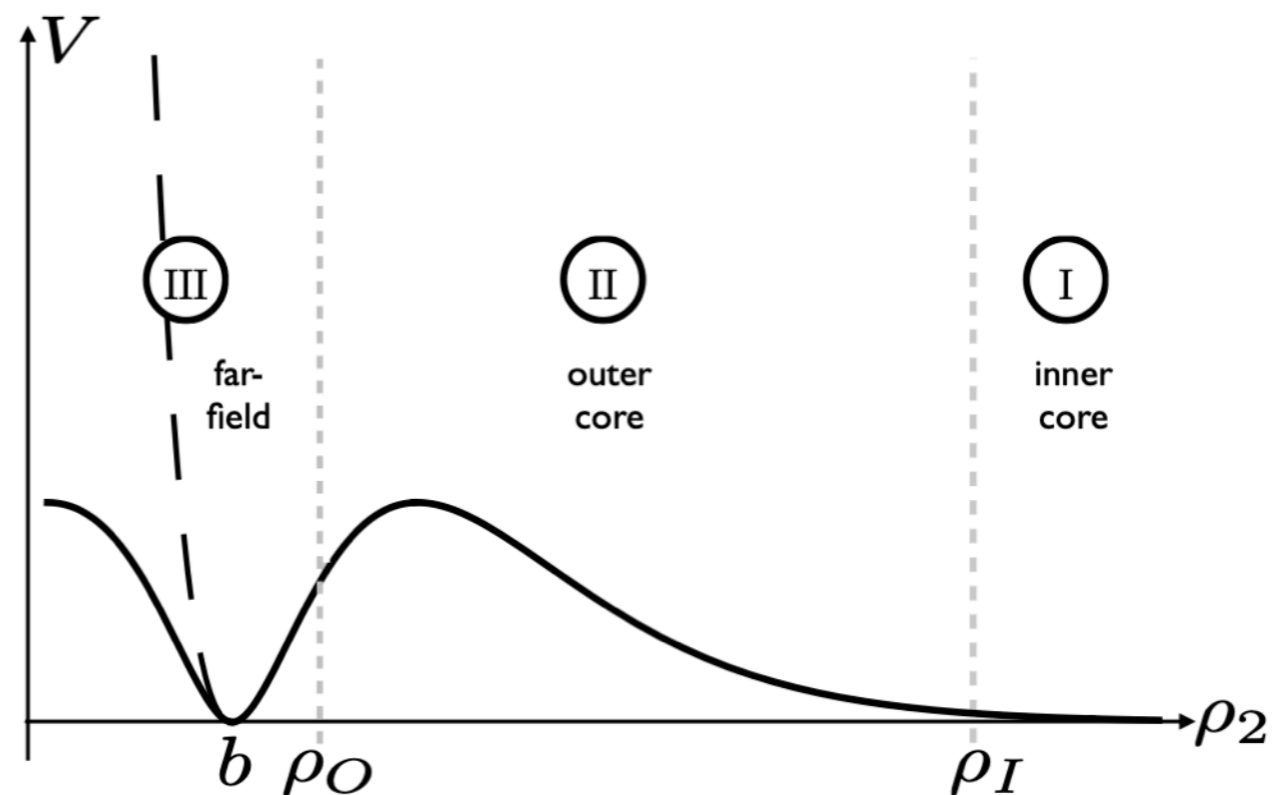


Caution

The dynamics of simple field theory axions, with $N_W = 1$

Possibly important differences for string theory axions, e.g.

- *production of strings*
- *core structure*
- *KK modes*
- *cosmological history*



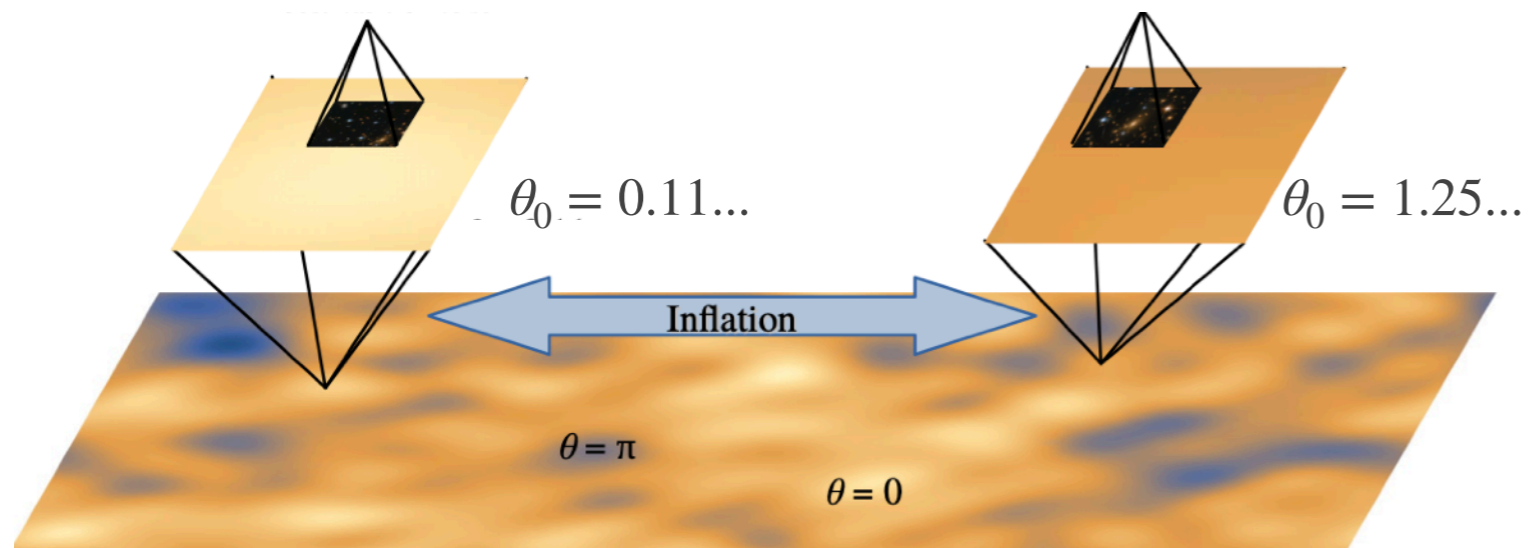
[March-Russell, Tillim]

Initial conditions

Pre-inflationary

Observable universe

Observable universe

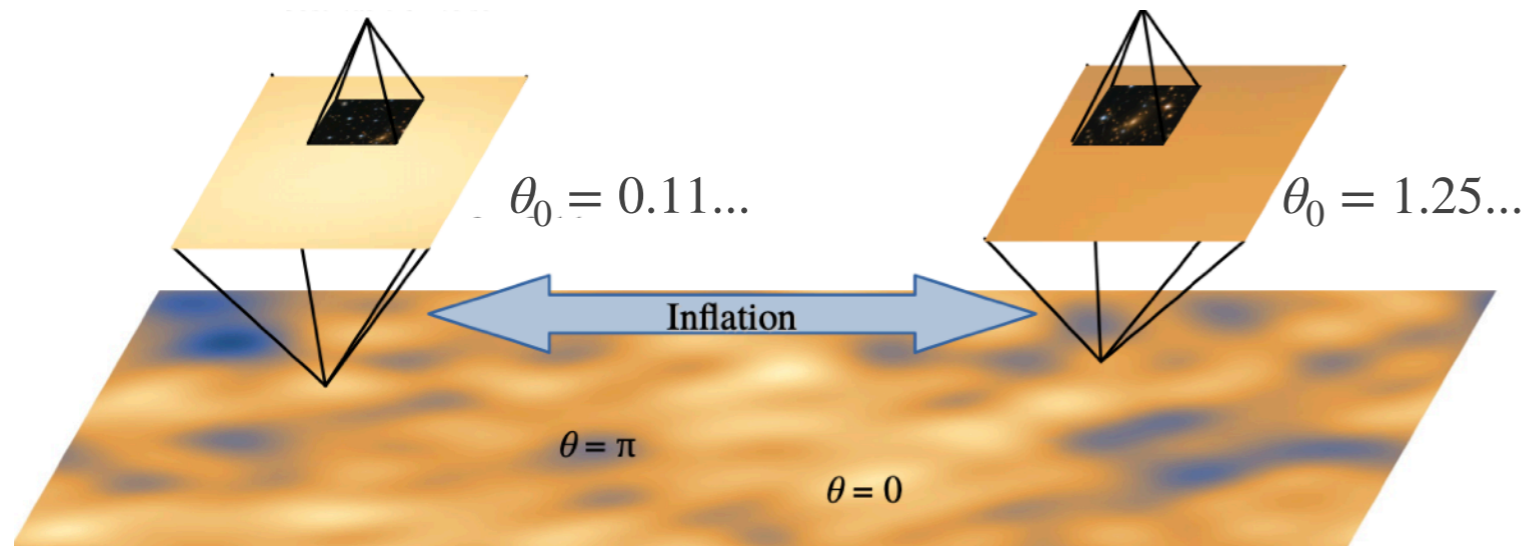


Initial conditions

Pre-inflationary

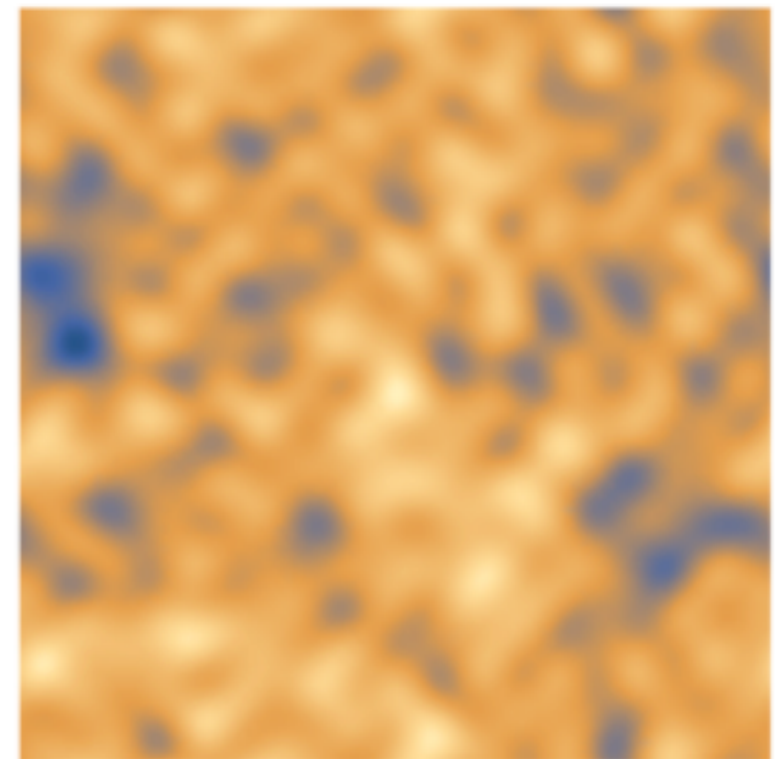
Observable universe

Observable universe

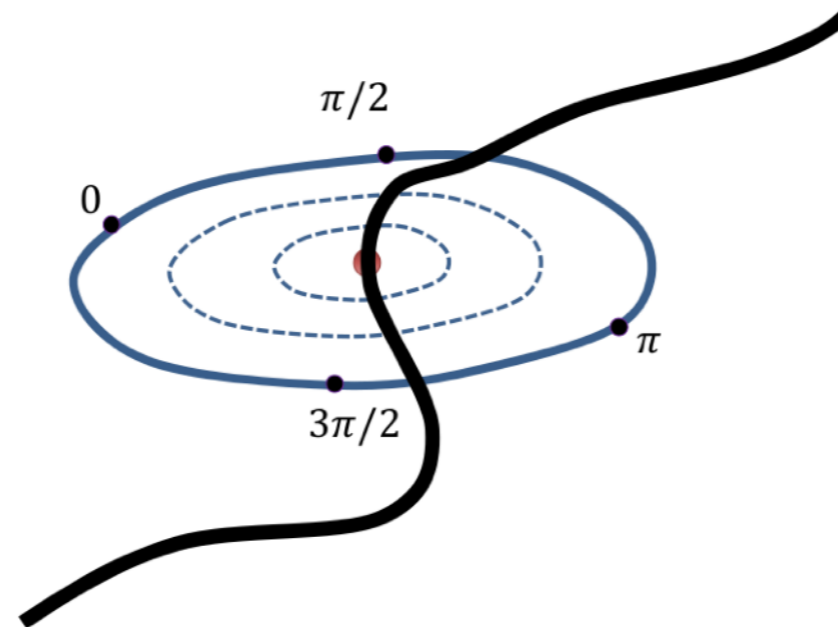
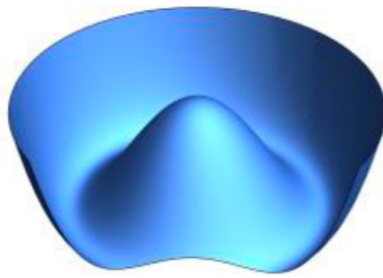
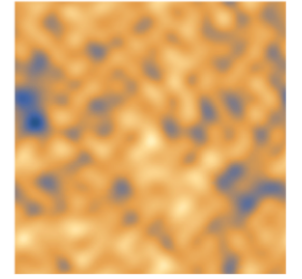


Post-inflationary

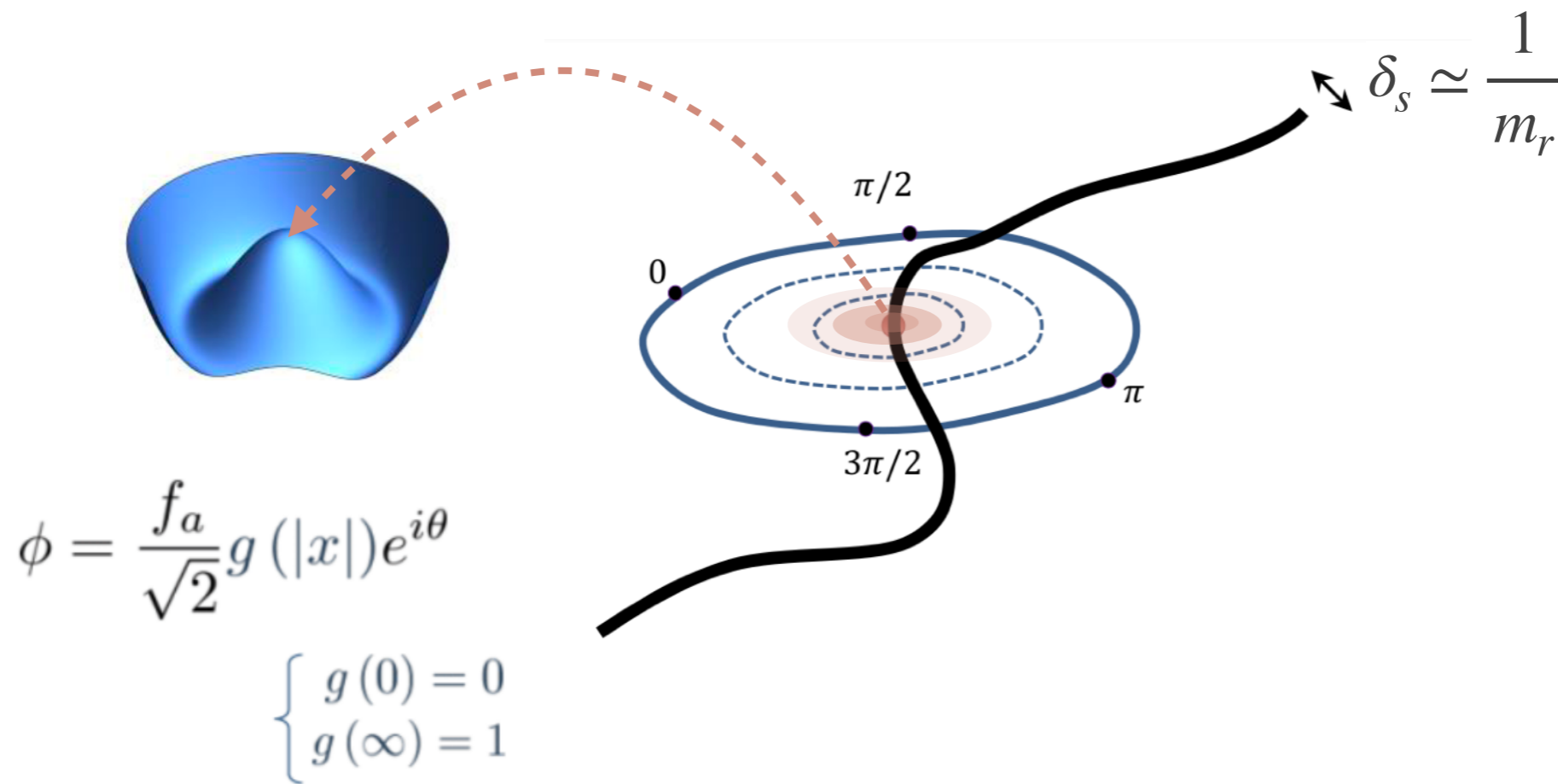
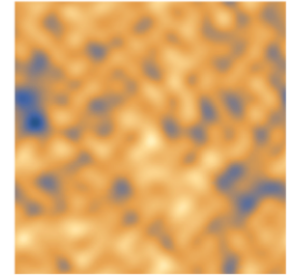
Observable universe



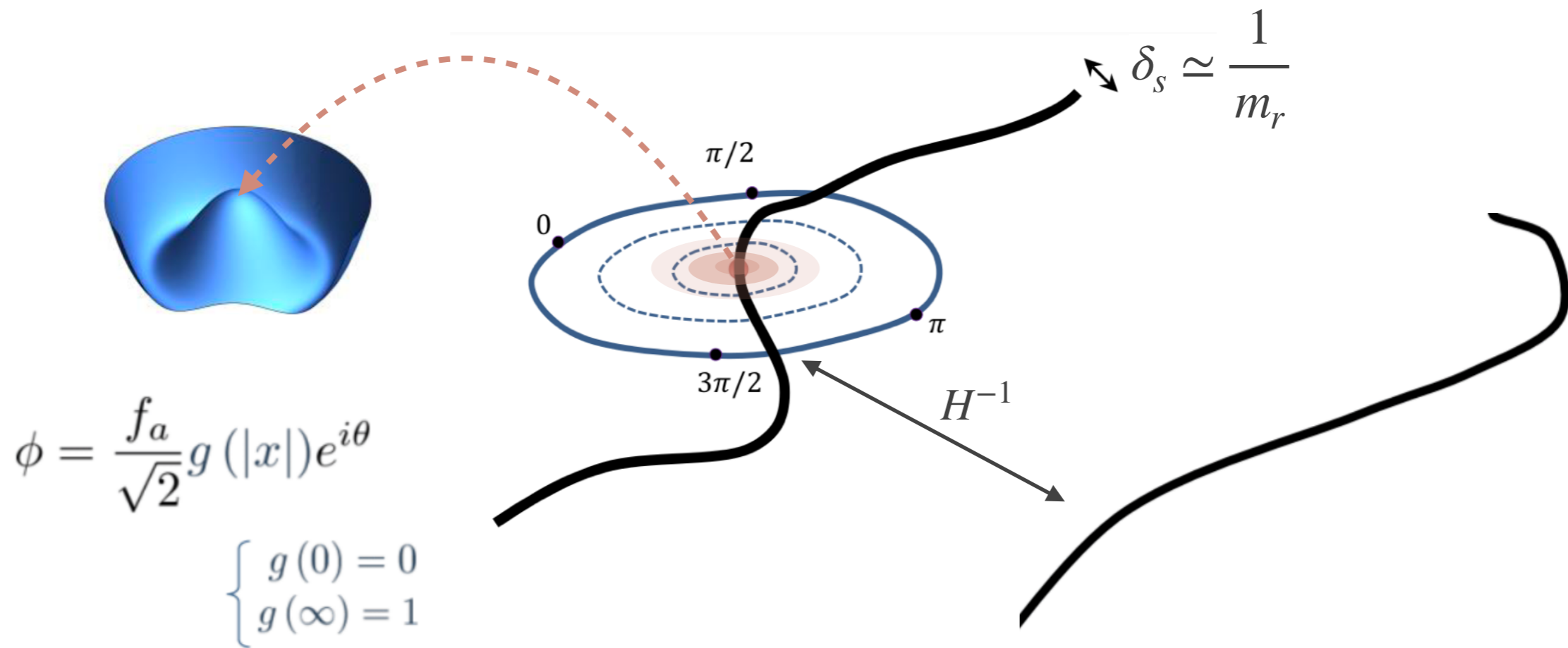
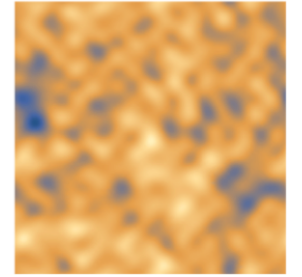
Topological strings



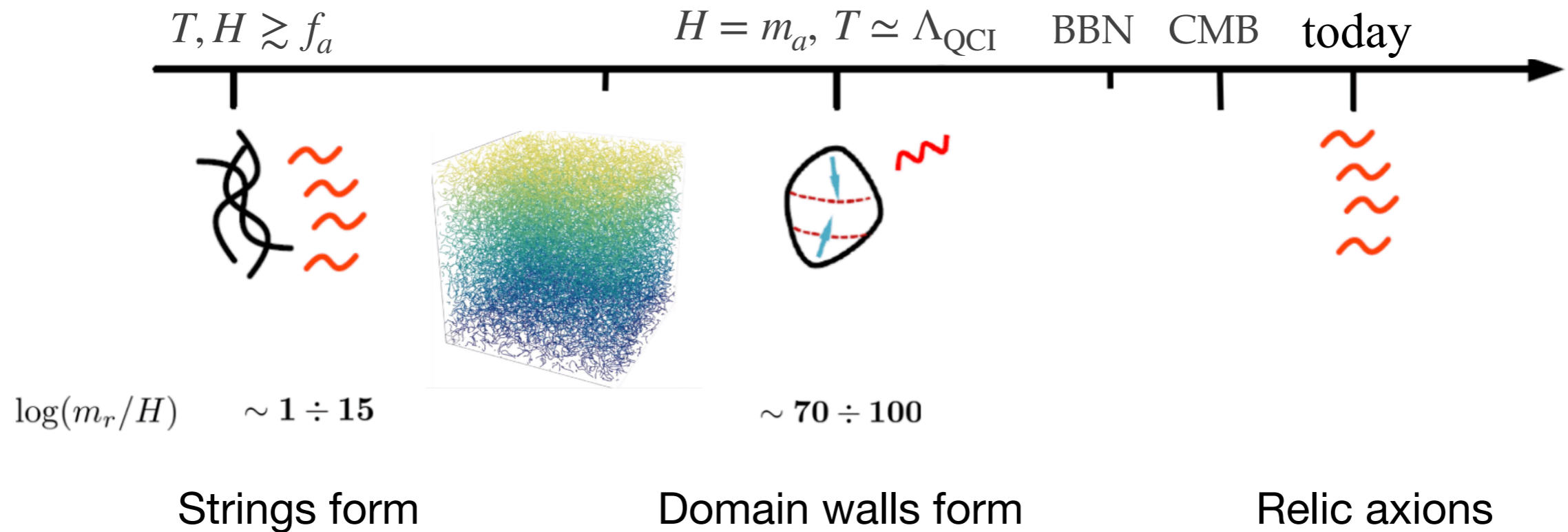
Topological strings



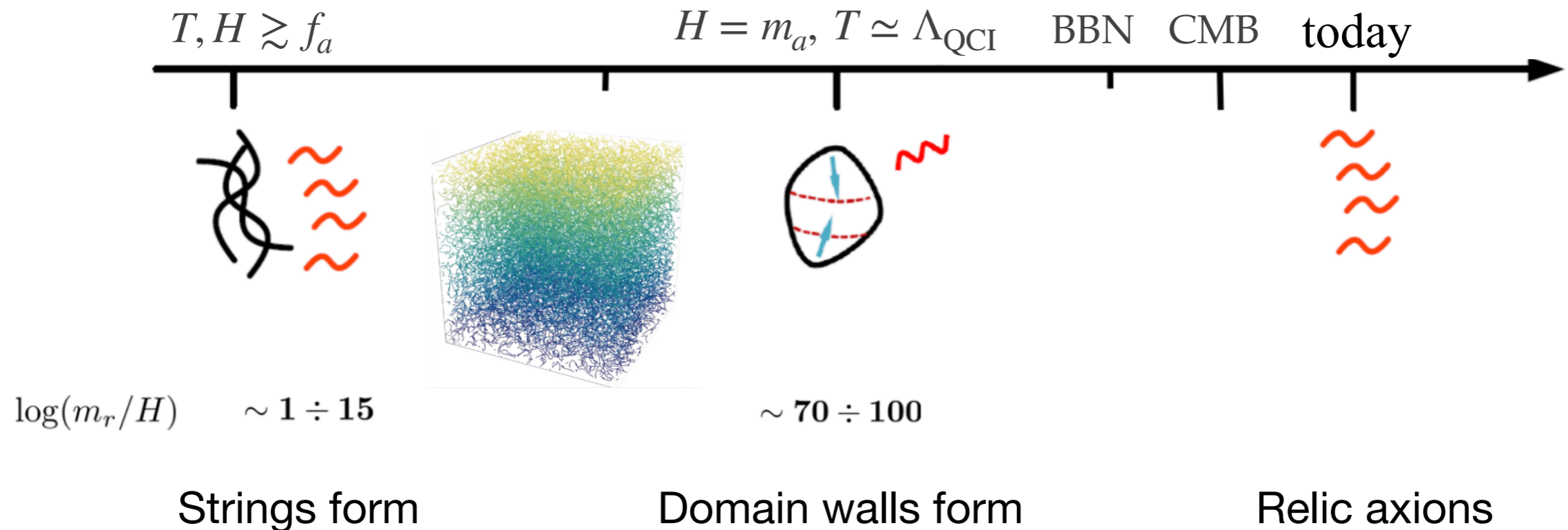
Topological strings



Full evolution



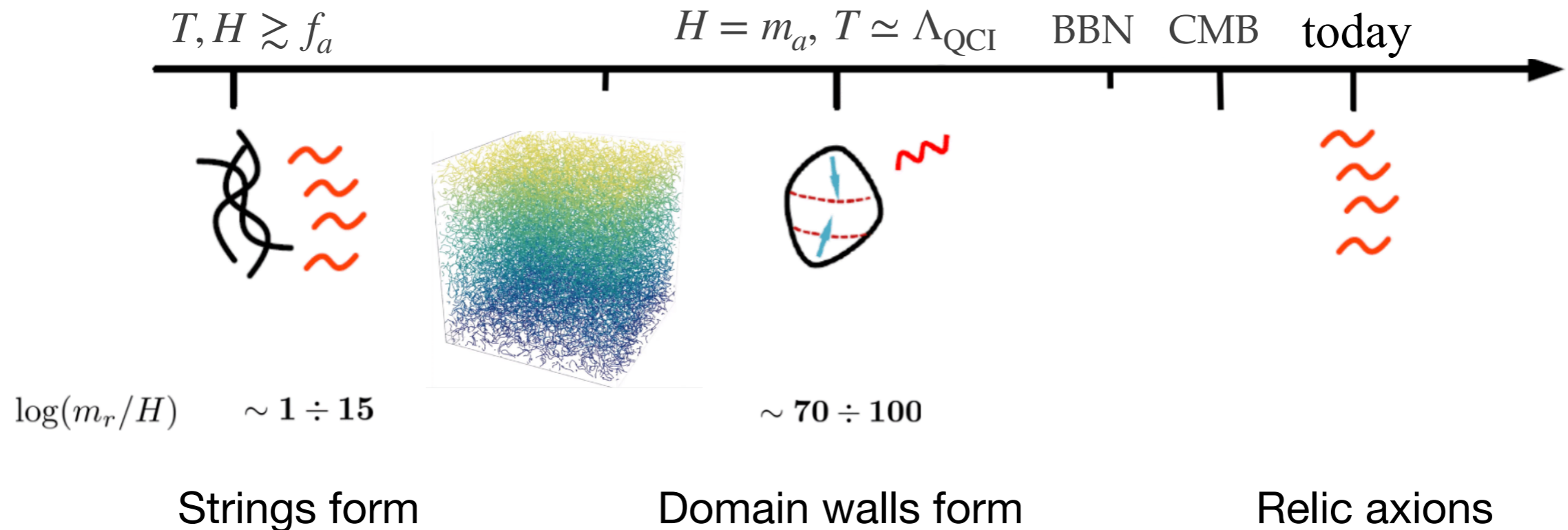
Full evolution



Dynamics:

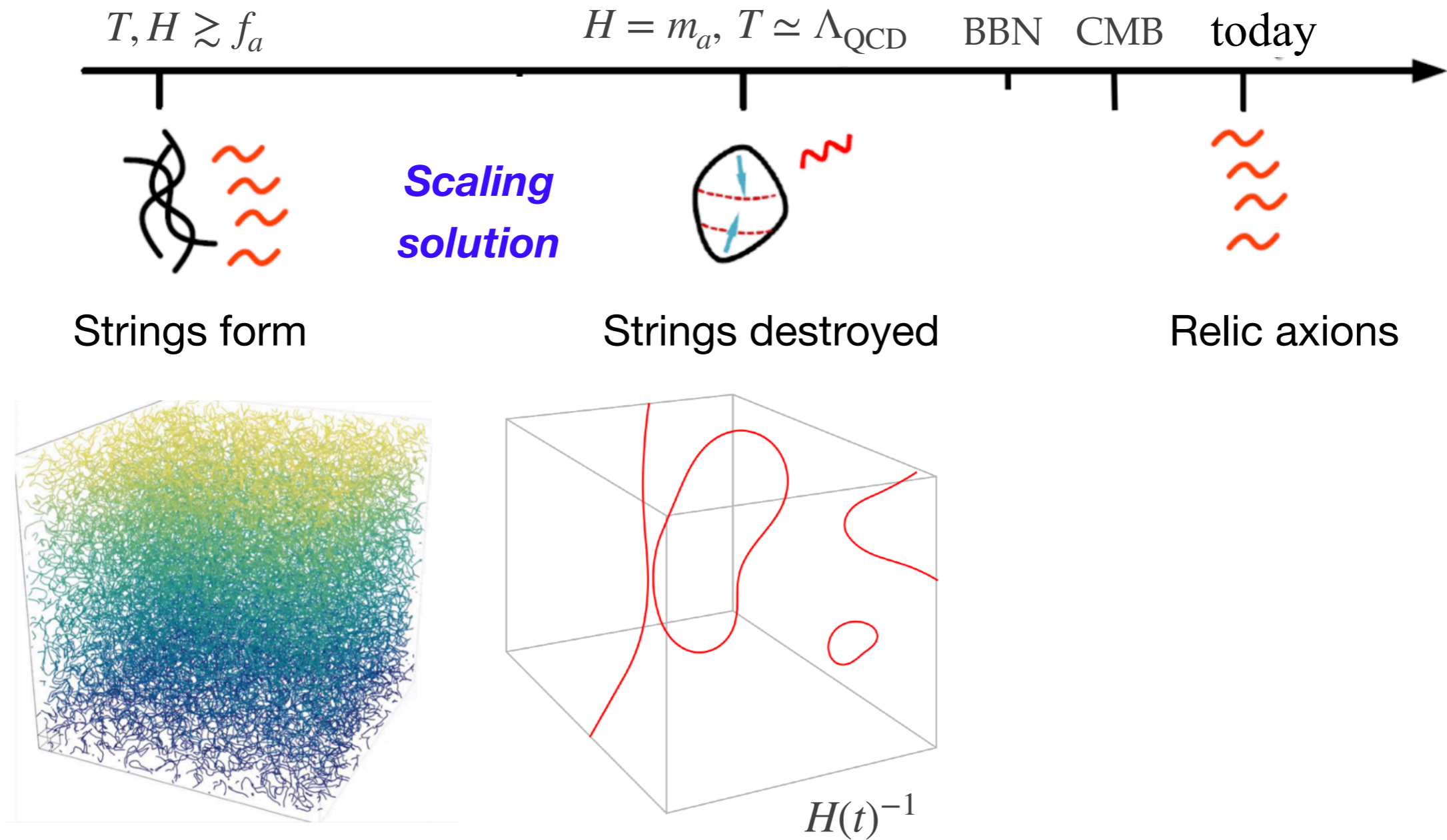
- *nonlinear*
- *large scale separation*

Full evolution

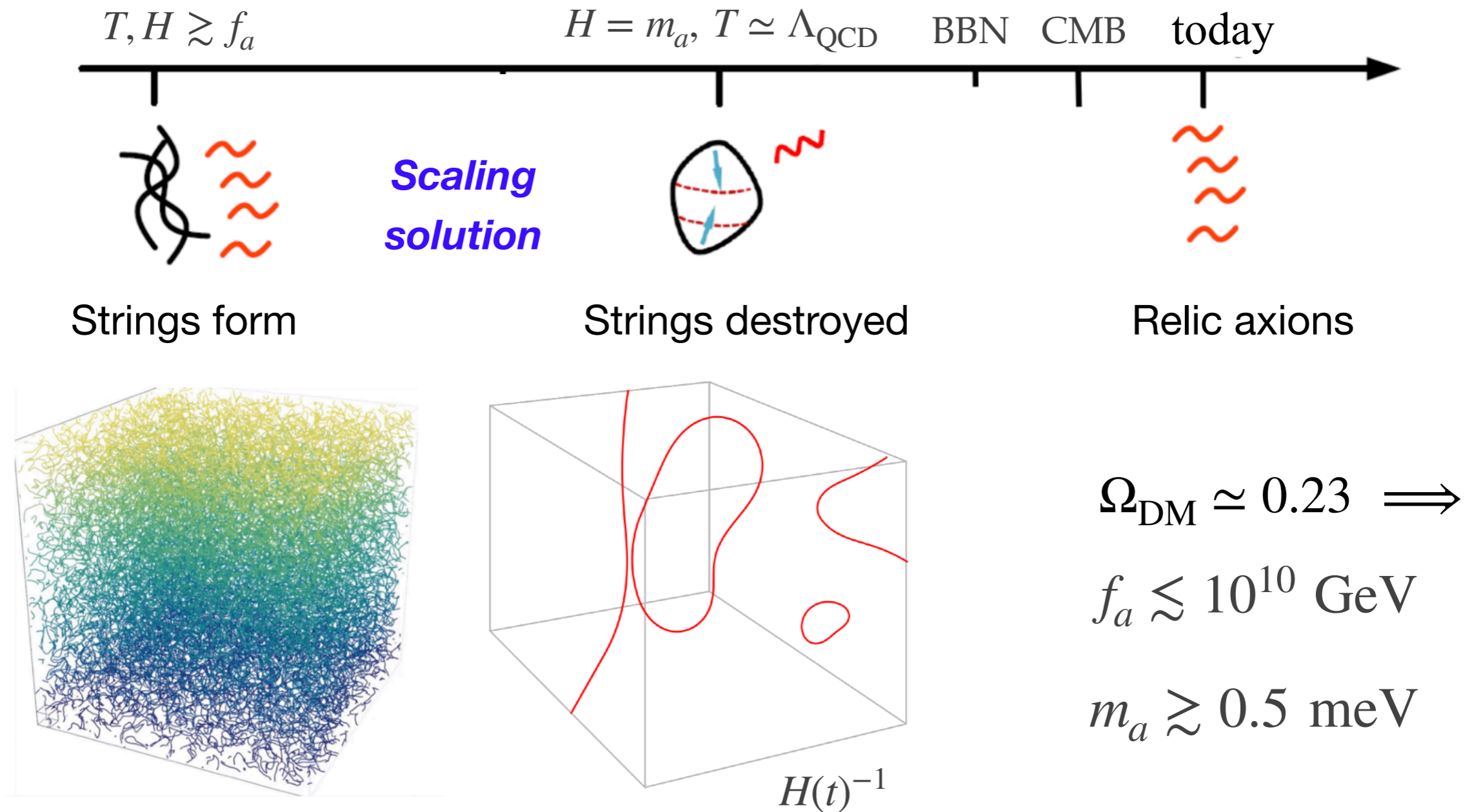


Simulate  *Extrapolate*

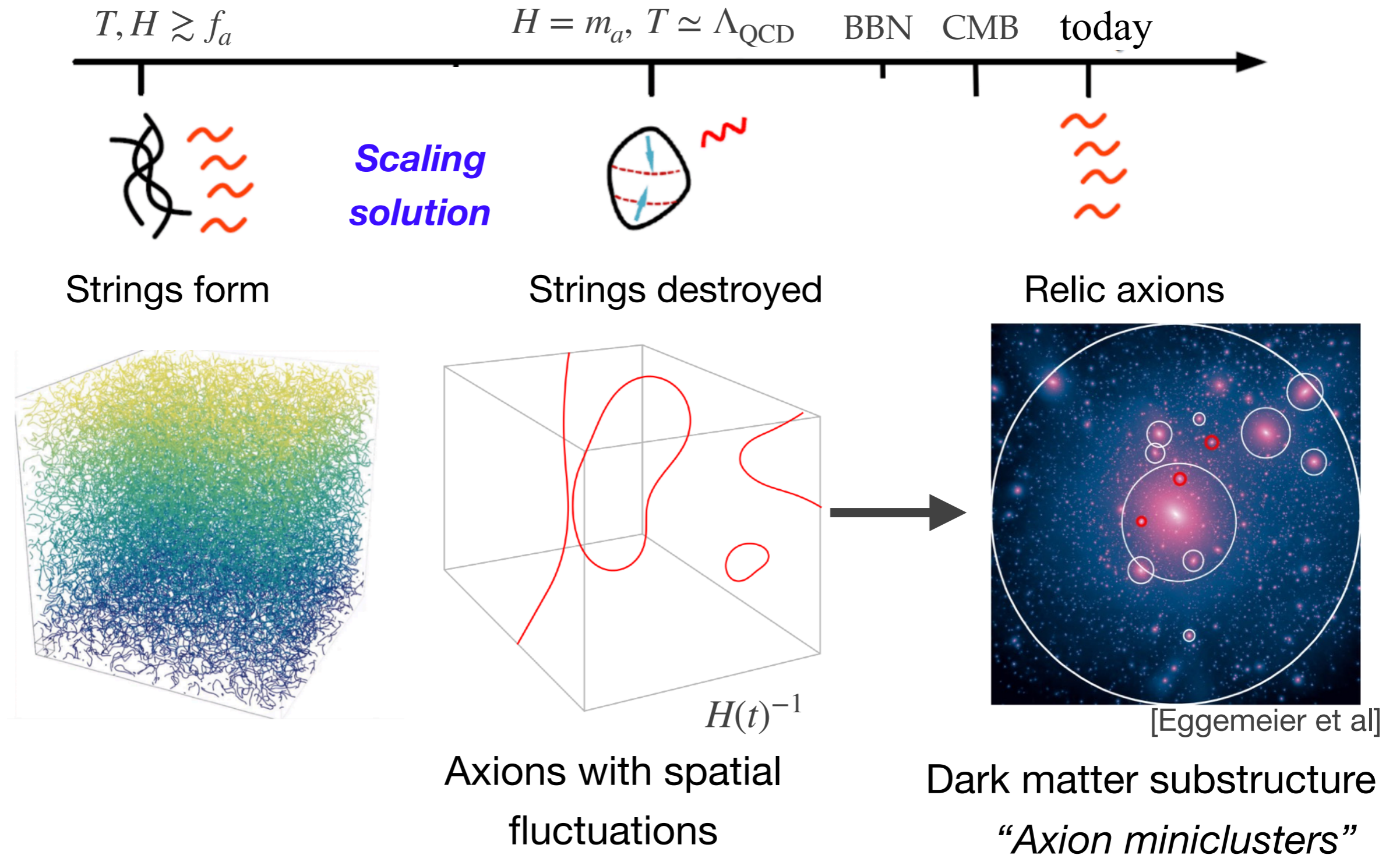
Full evolution



Full evolution

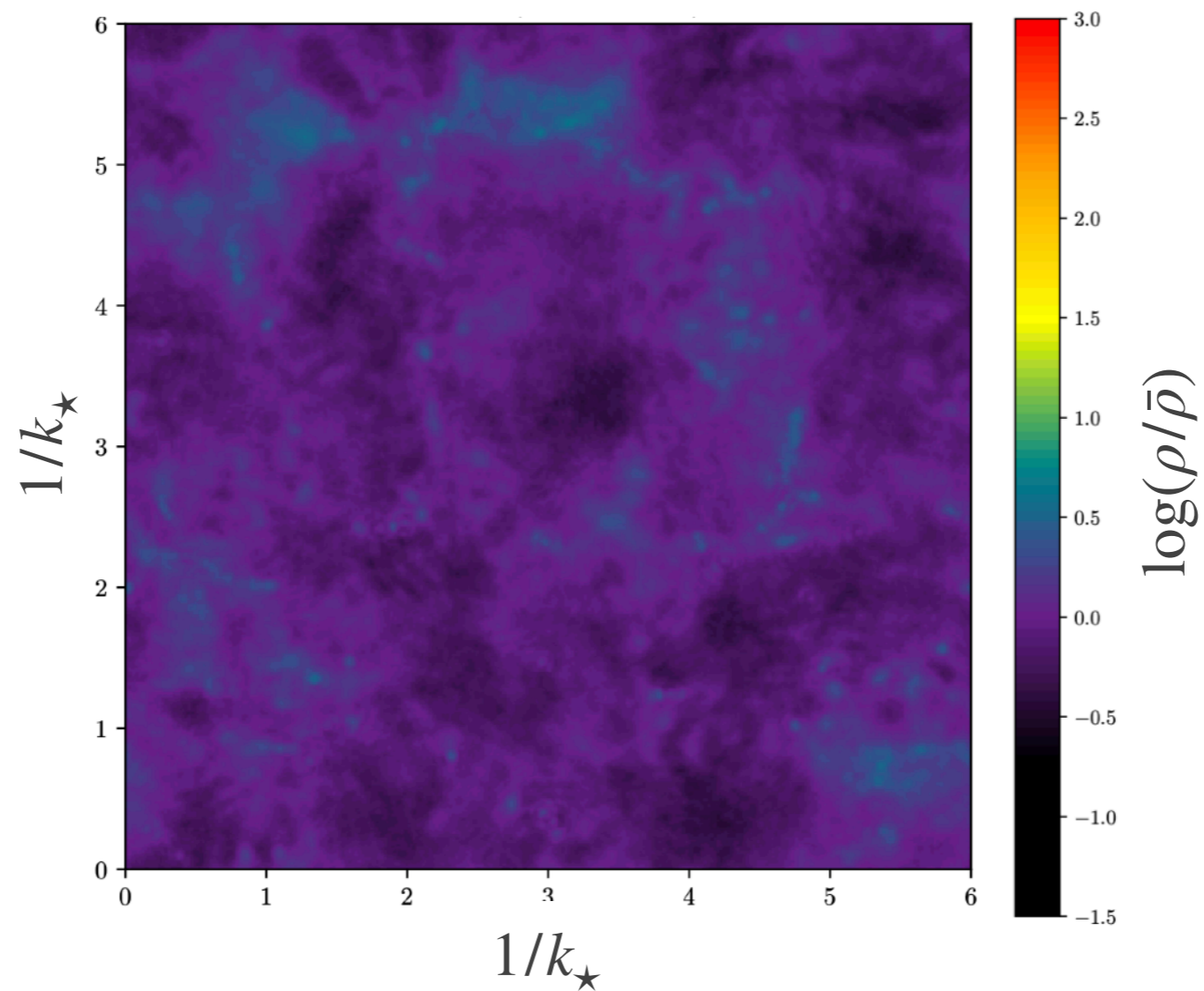


Full evolution



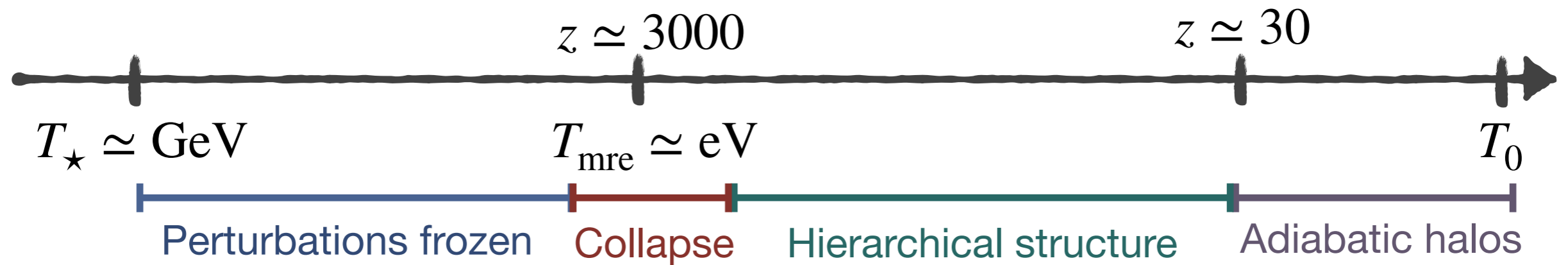
Initial perturbations

Order one fluctuations on co-moving scales $\simeq H_\star$ when $H = m_a(T)$

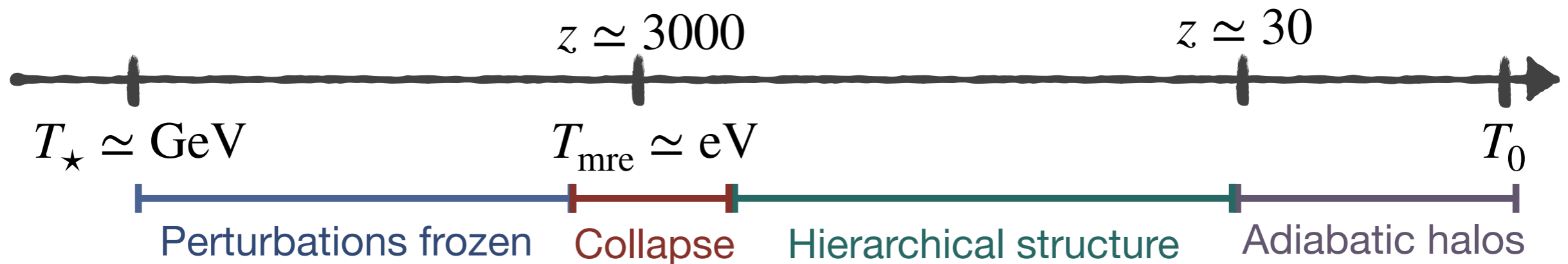


[Eggemeier et al]

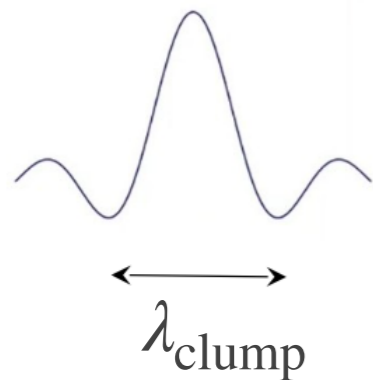
Standard picture



Standard picture



Wave effects at matter radiation equality



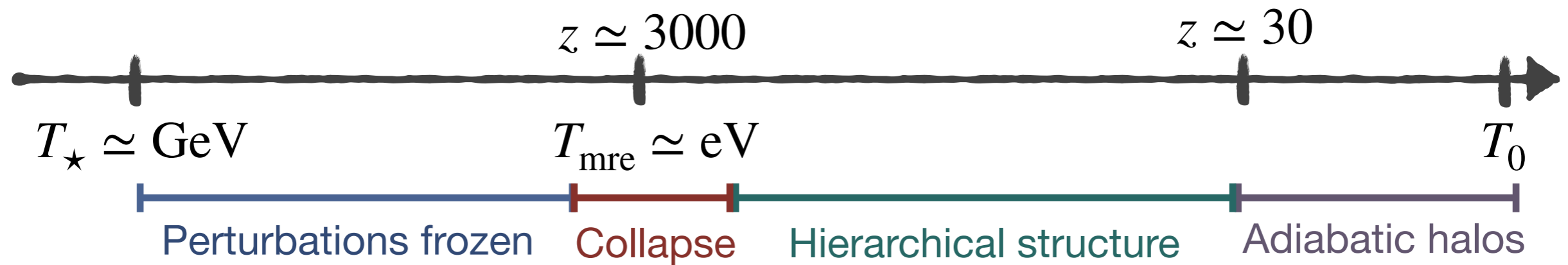
$$\lambda_{\text{dB}} = \frac{1}{m_a v} = \frac{1}{m_a (GM/\lambda_{\text{clump}})^{1/2}} = \frac{1}{\lambda_{\text{clump}} (4\pi G \rho m_a^2)^{1/2}}$$

“Quantum” Jeans scale:

$$\lambda_J \simeq (G \rho m_a^2)^{1/4}$$

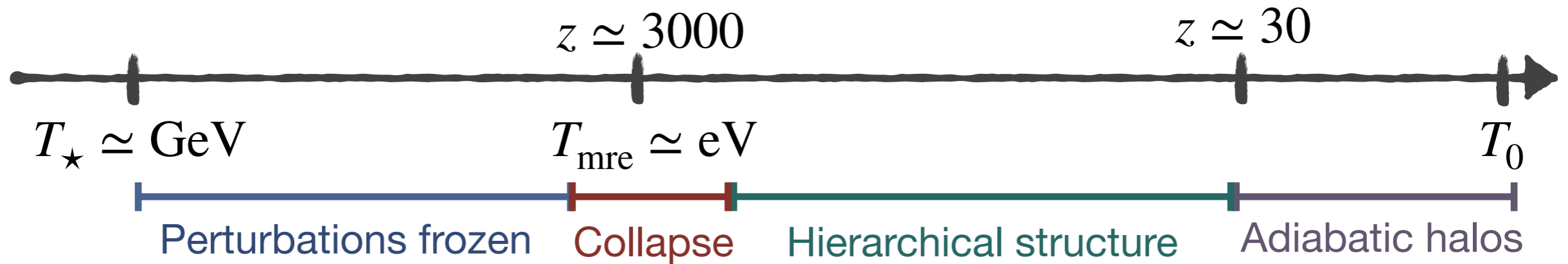
$$k_J/R = (16\pi G \rho m_a^2)^{-1/4}$$

Standard picture



ALP:
$$\left. \frac{\kappa_p}{\kappa_J} \right|_{\text{MRE}} = \frac{\kappa_{p\star} a_\star / a_{\text{MRE}}}{(16\pi G \rho_{\text{MRE}} m^2)^{1/4}} \simeq \frac{\kappa_{p\star}}{H_\star}$$

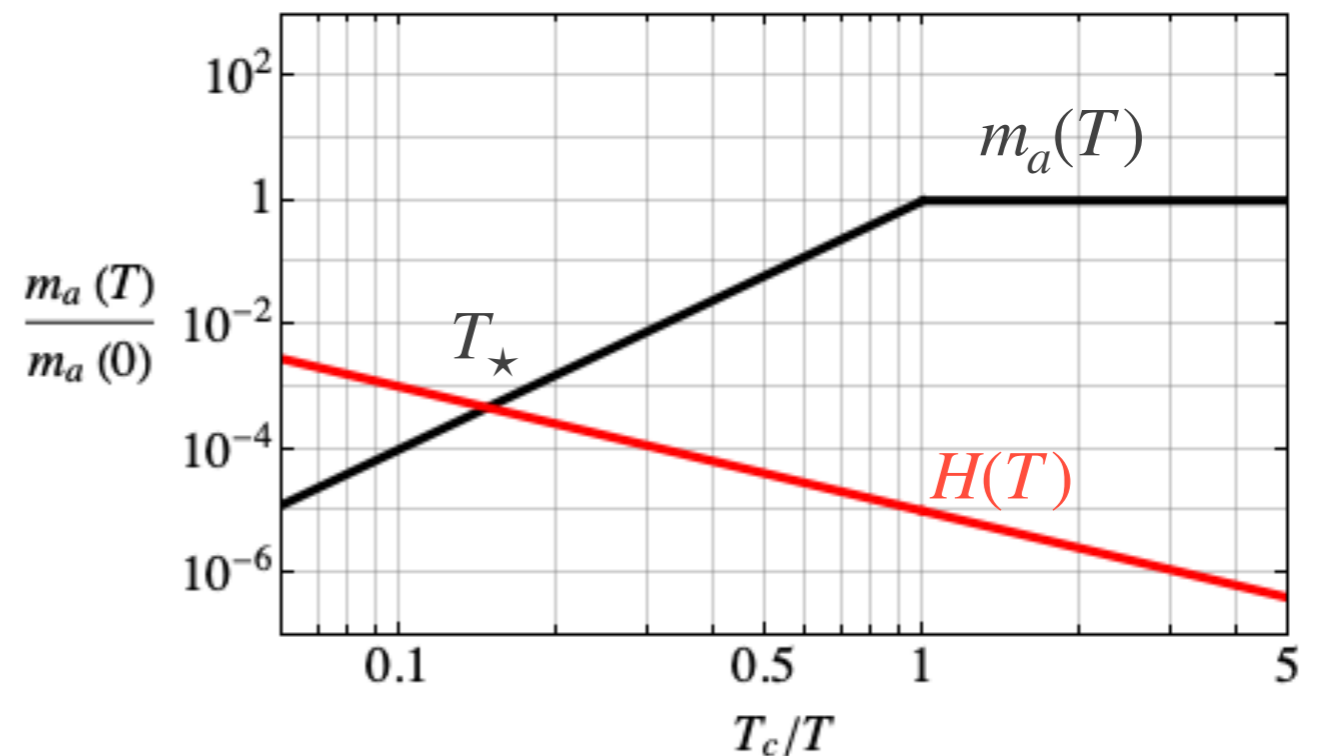
Standard picture



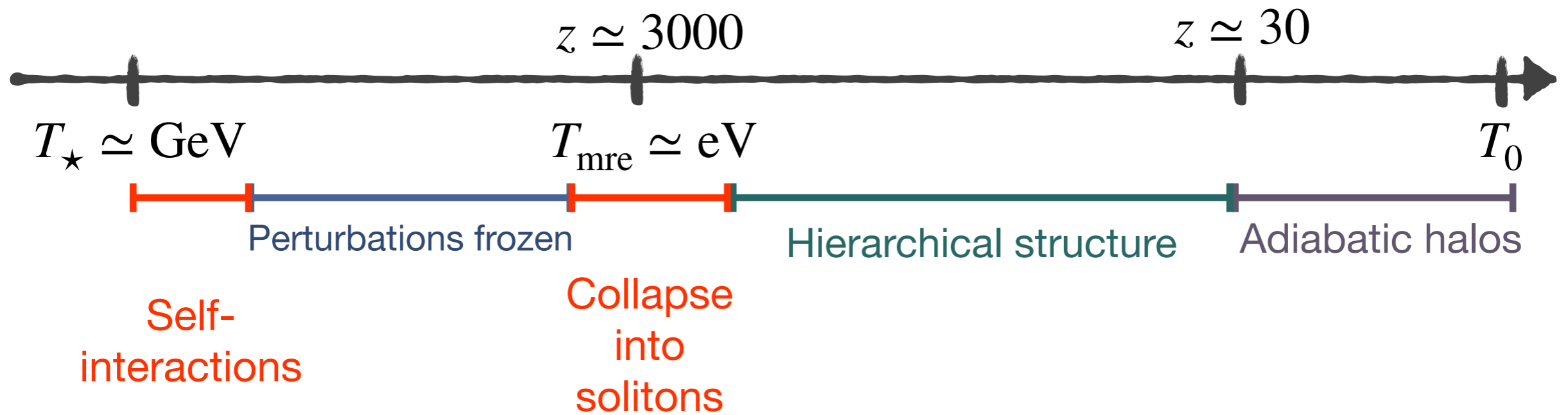
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QCD axion:

$$\left. \frac{\kappa_p}{\kappa_J} \right|_{\text{MRE}} = \frac{\kappa_{p\star}}{6^{1/4} H_\star} \left(\frac{m_\star}{m} \right)^{1/2} \sim 10^{-3} \frac{\kappa_{p\star}}{H_\star}$$



New aspects



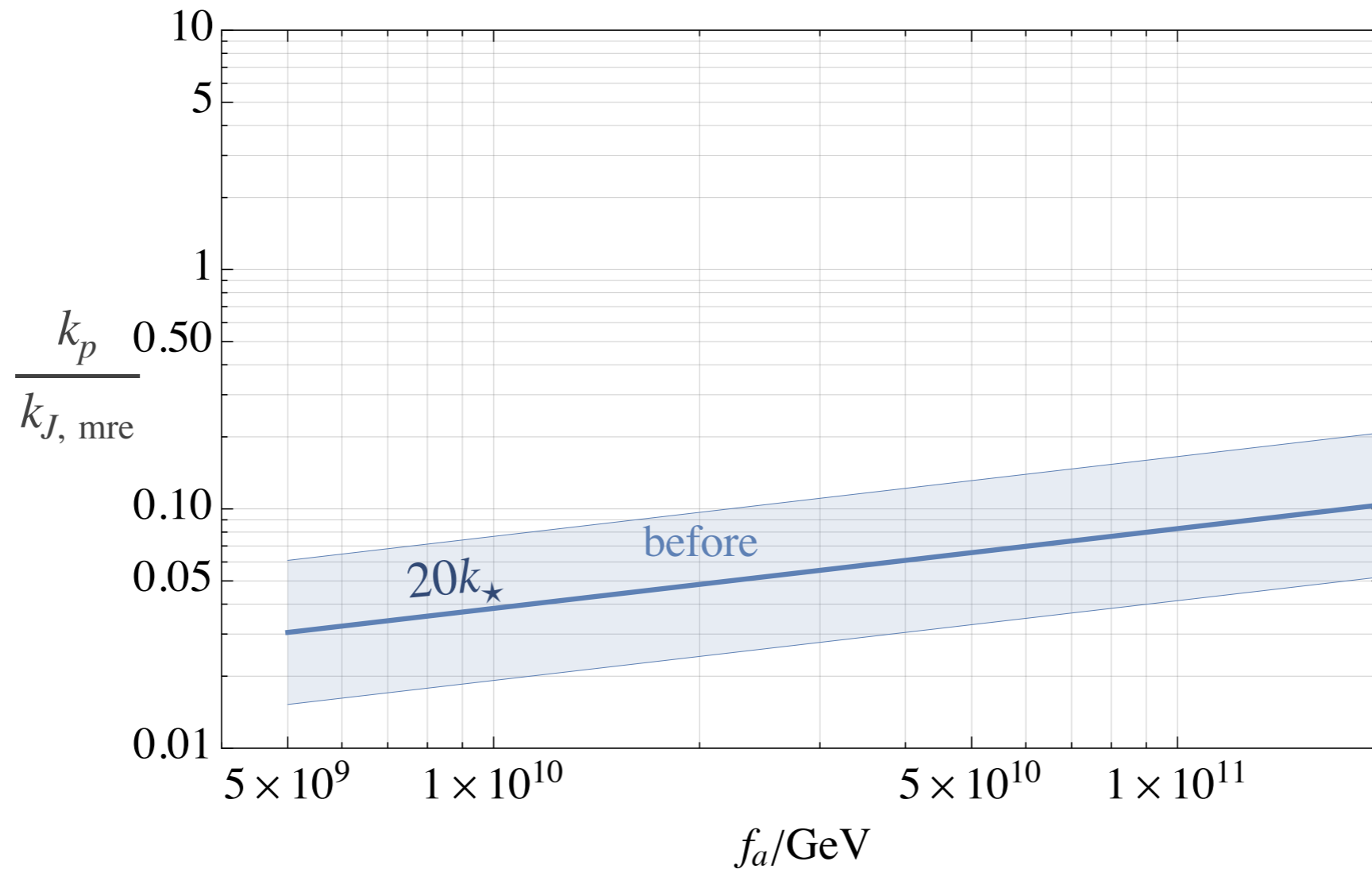
QCD axion:

~~$$\frac{k_p}{k_J} \Big|_{\text{MRE}} = \frac{k_{p\star}}{6^{1/4} H_{\star}} \left(\frac{m_{\star}}{m} \right)^{1/2}$$

$$\sim 10^{-3} \frac{k_{p\star}}{H_{\star}}$$~~

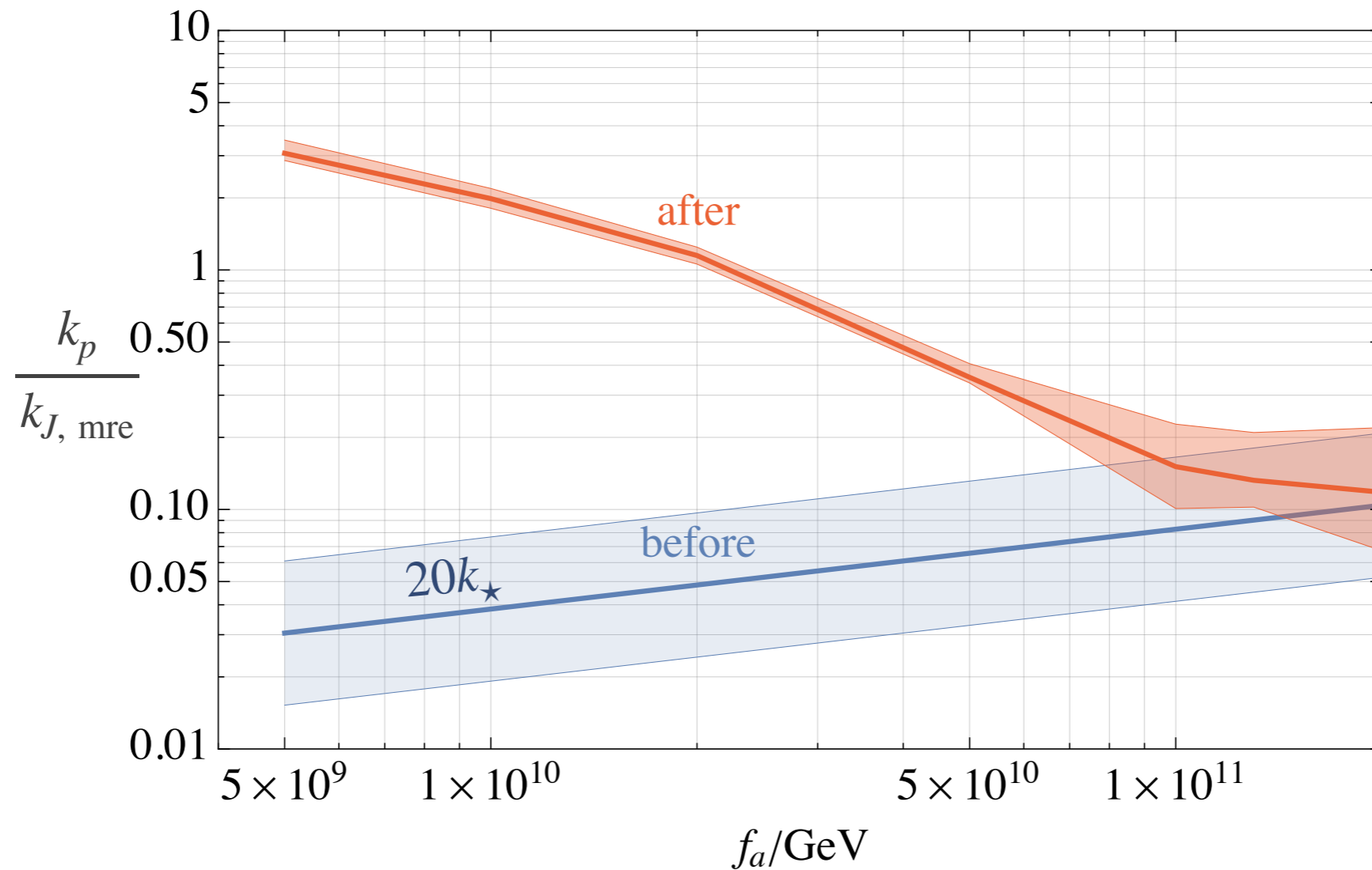
Self-interactions

$$\frac{k_{\star}}{k_{J,\text{eq}}} \simeq 0.002 \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^{1/3}$$

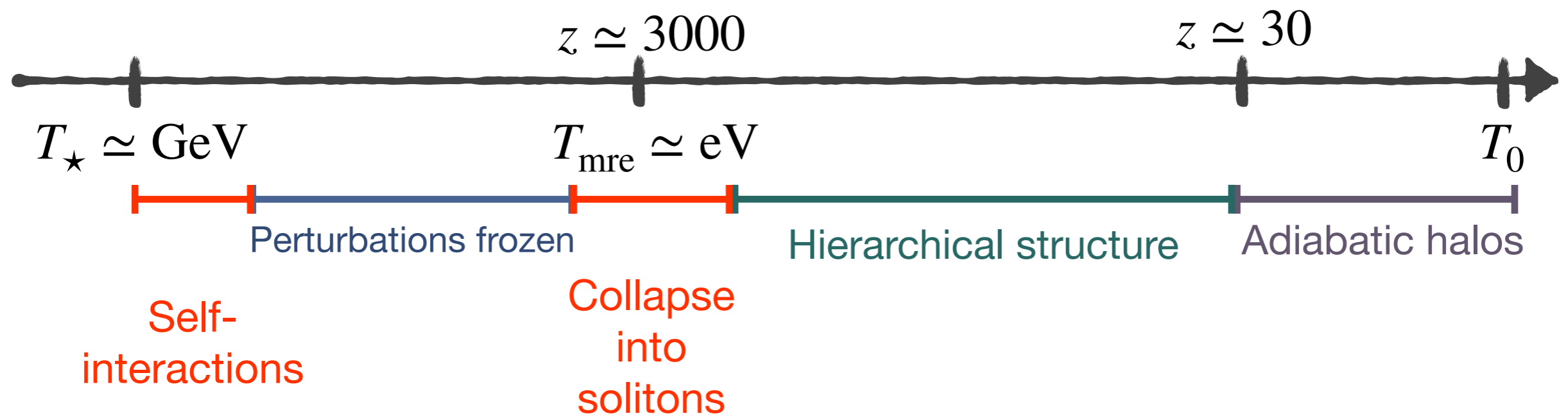


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New aspects

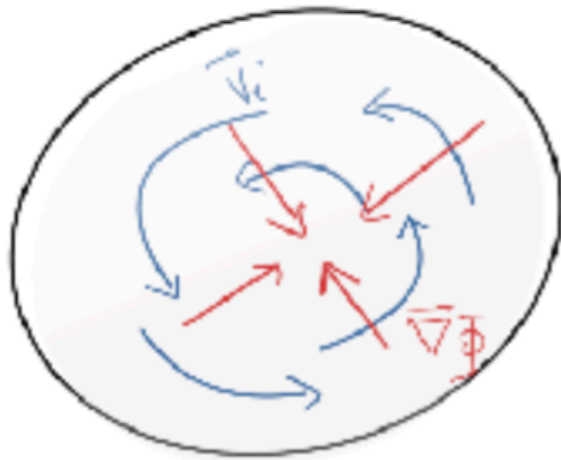


Halos vs solitons

Halos

$$\Phi_Q \simeq 0$$

→ gravitational potential balanced
by velocity term



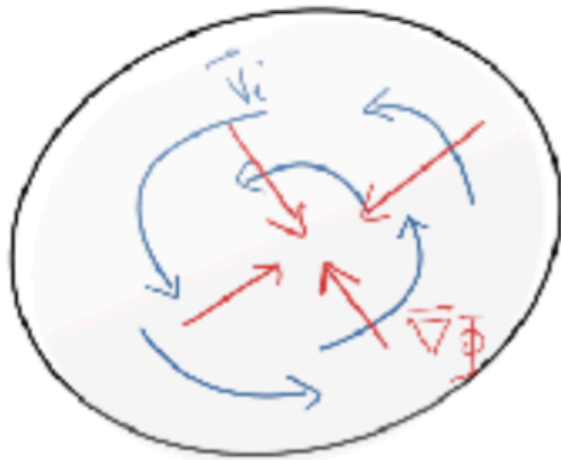
Angular momentum “supports” the
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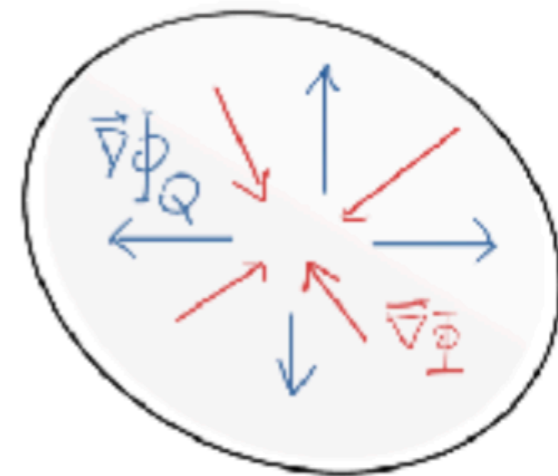


Angular momentum “supports” the
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Soliton

$$\Phi_Q = -\Phi \quad \vec{v} = 0$$

→ gravitational potential balanced by
quantum pressure “*Axion star*”



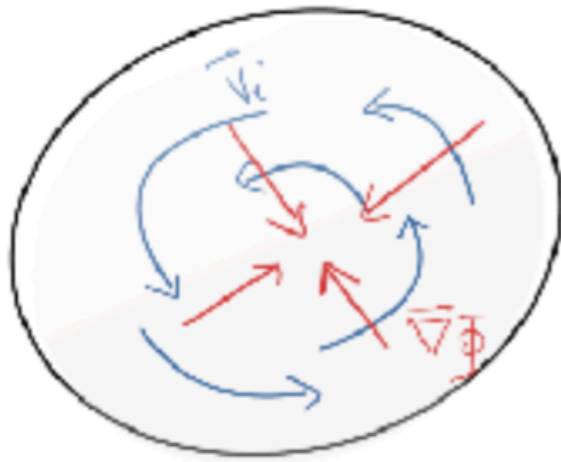
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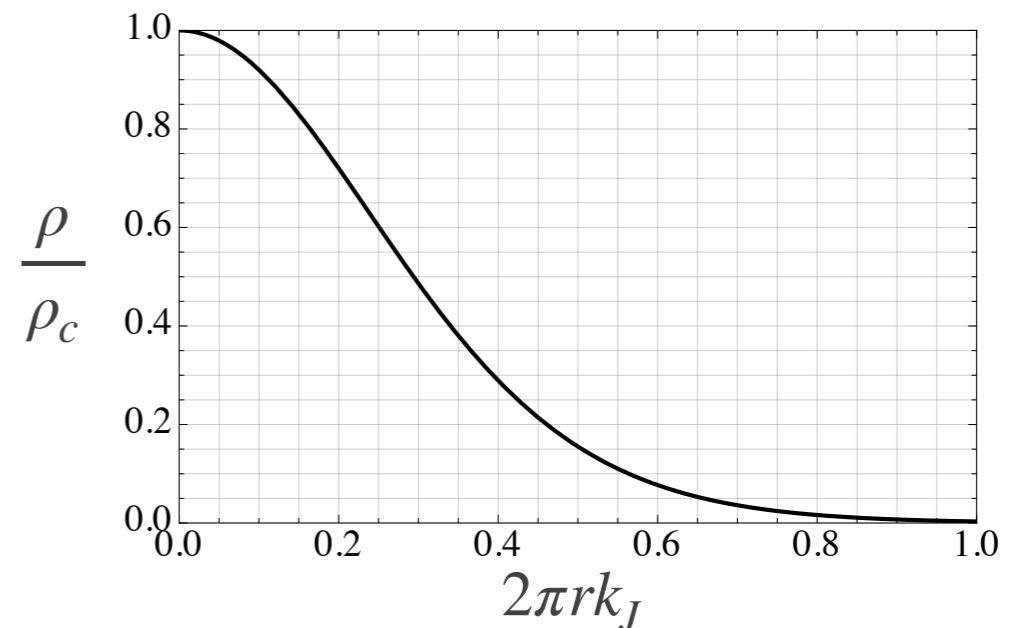


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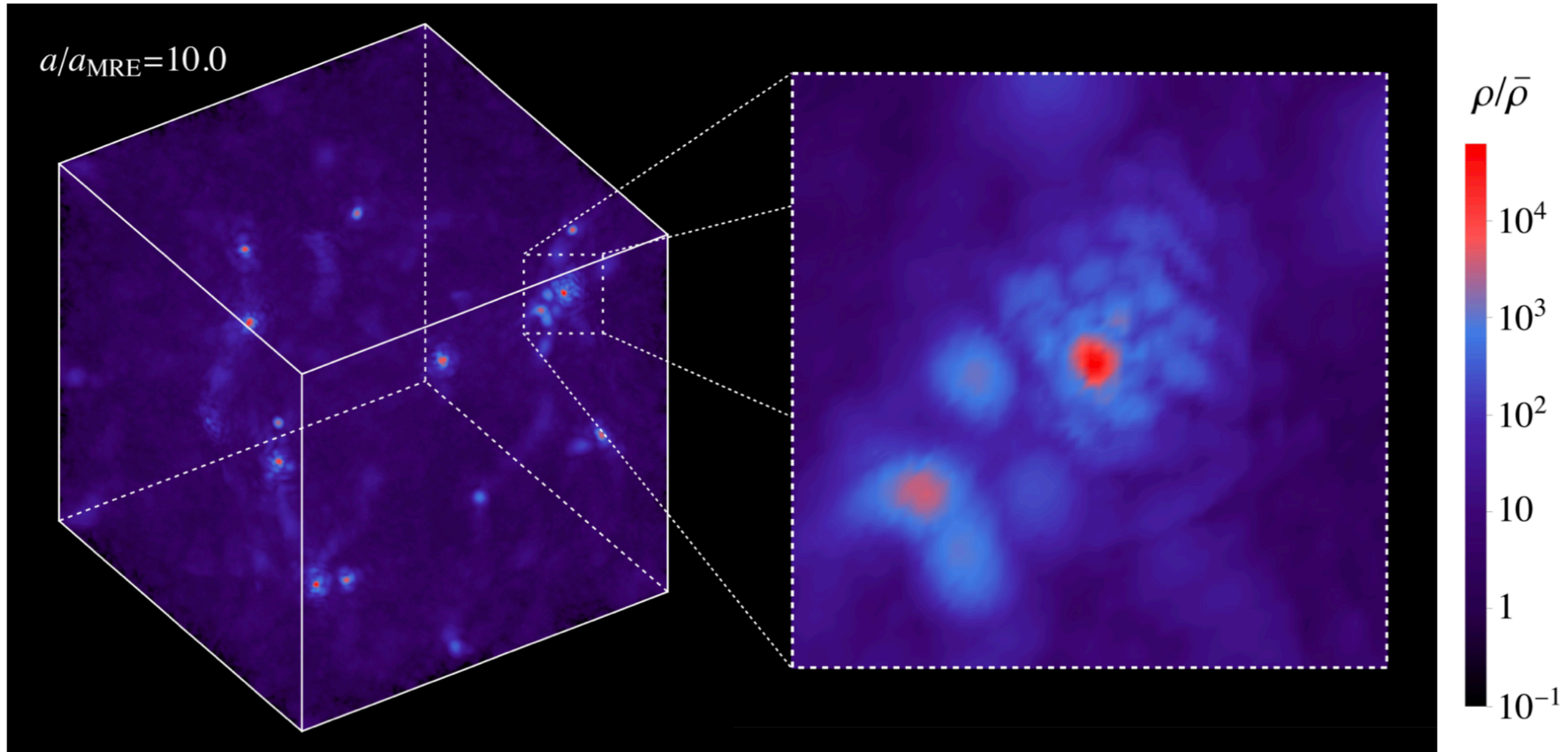
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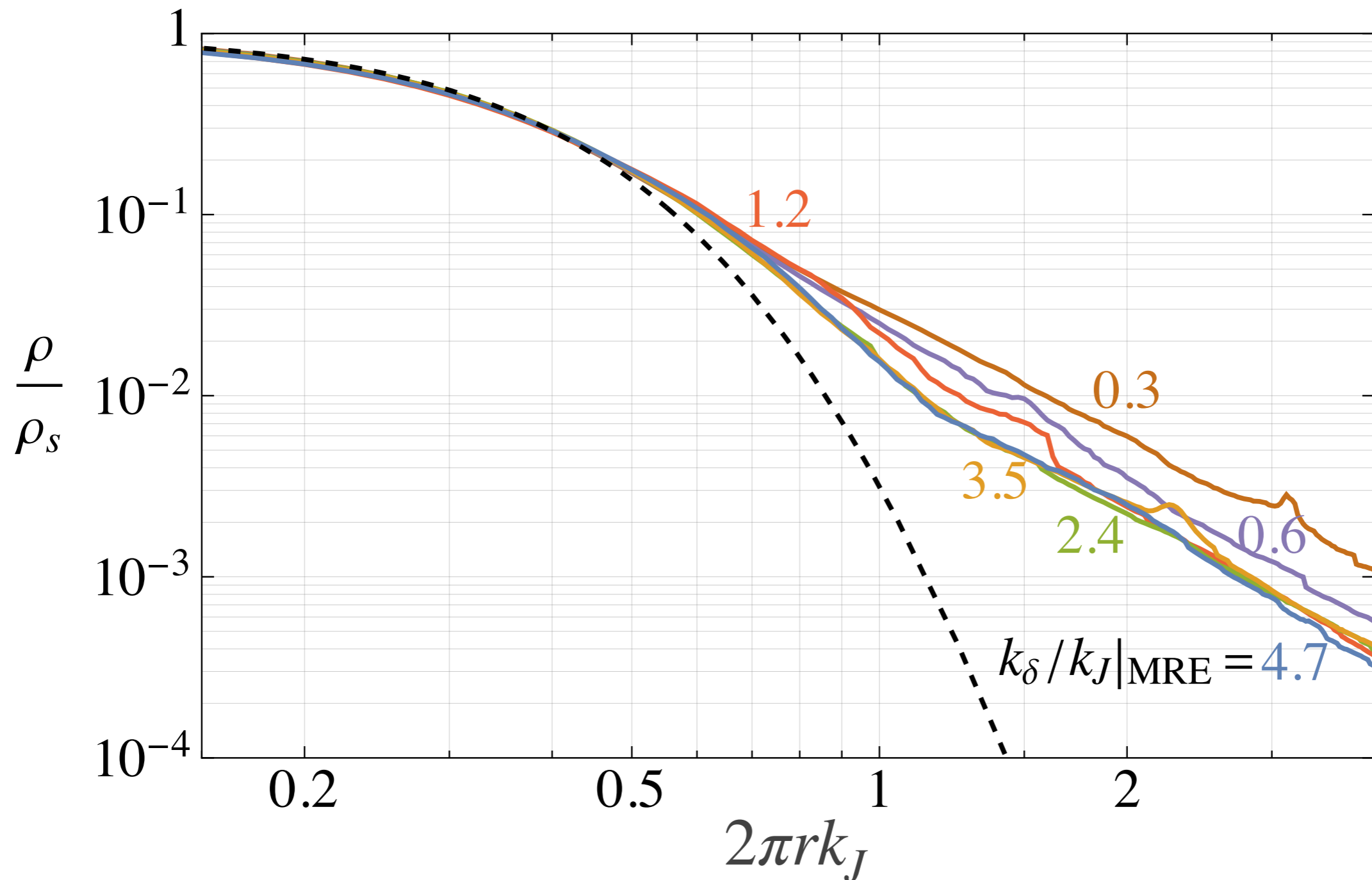


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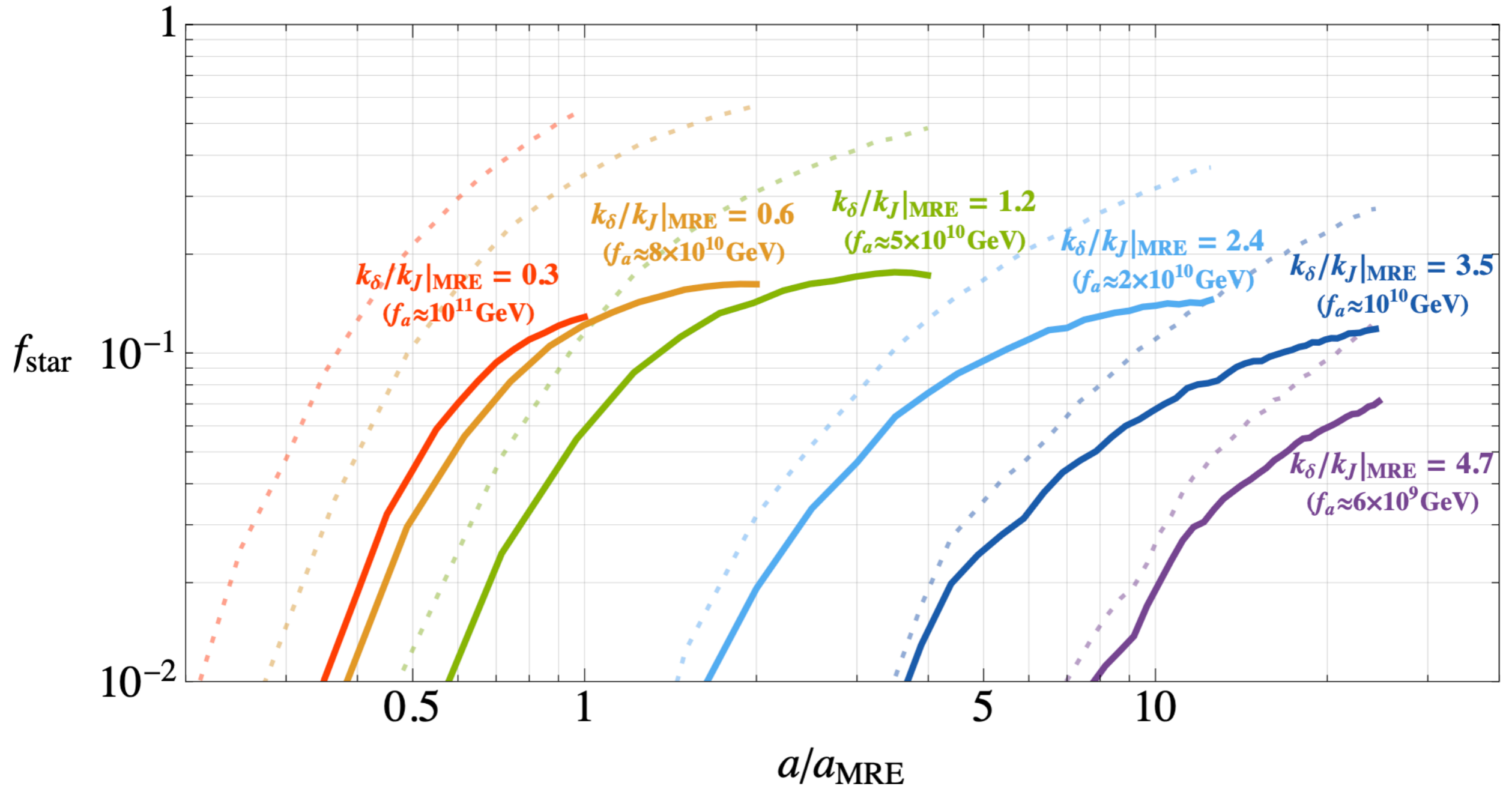
Simulations



Properties of the substructure



Properties of the substructure



Properties of the substructure

$$\bar{M}_s \approx 2 \cdot 10^{-19} M_\odot \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^{\frac{5}{2}}$$

$$R_{0.1} \simeq 4.2 \times 10^6 \text{ km} \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^2 \left(\frac{10^{-19} M_\odot}{M_s} \right)$$

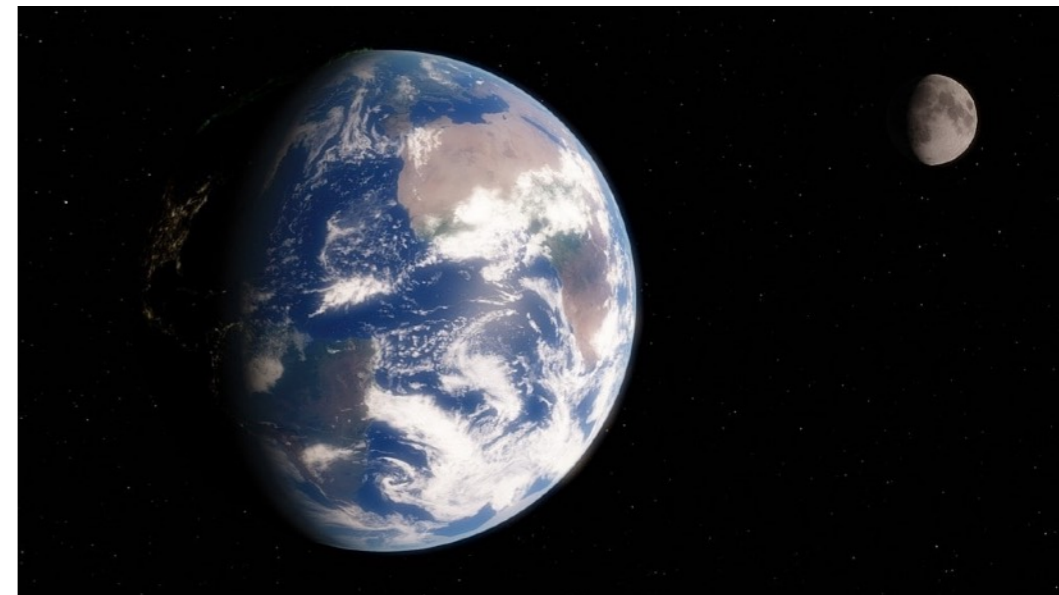
$$\bar{\rho}_s \approx 0.1 \text{ eV}^4 \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^4$$

Properties of the substructure

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Properties of the substructure

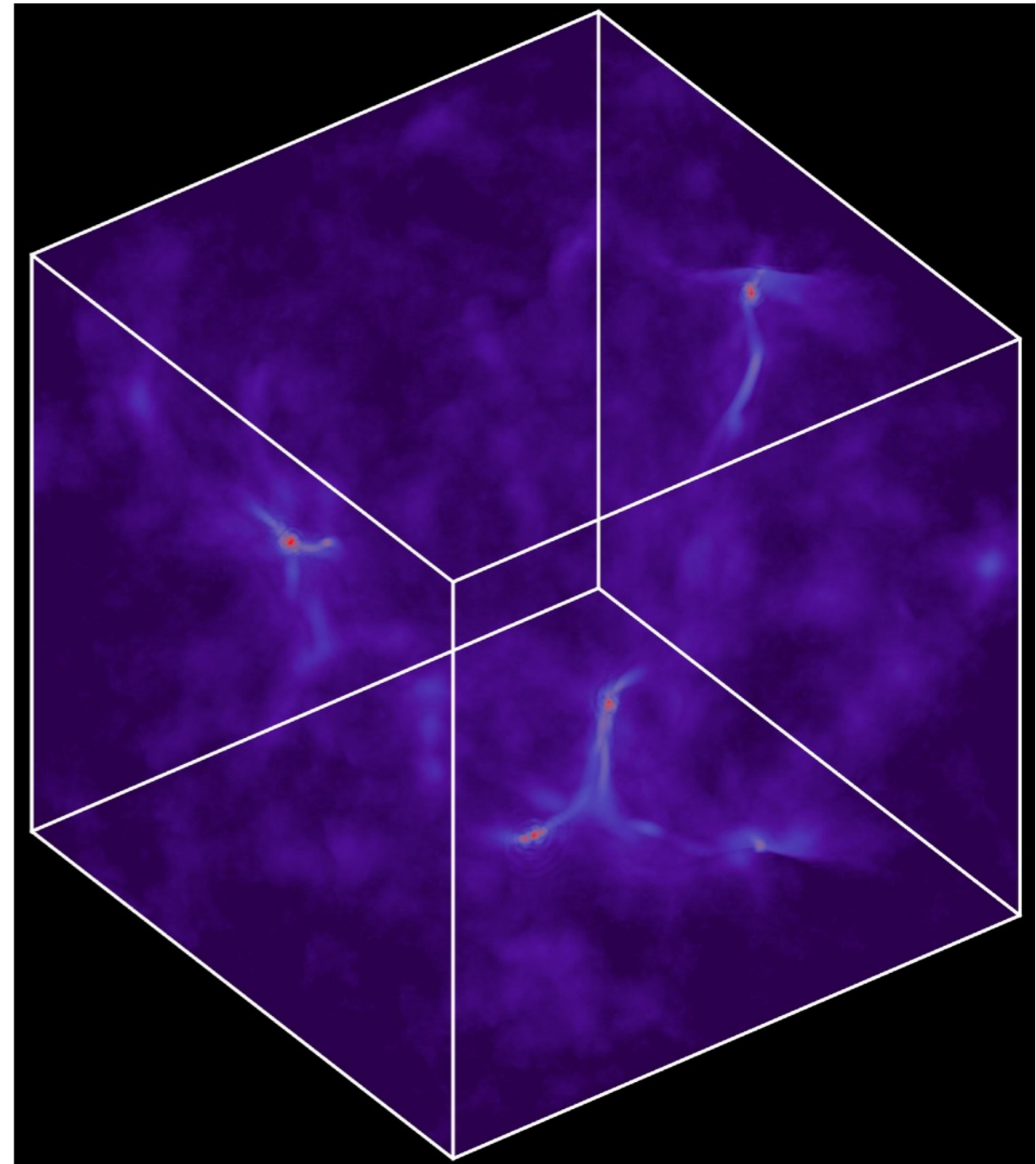
$$\tau_{\oplus} = 5 \text{ yrs} \left(\frac{R_{0.1}}{R} \right)^2 \left(\frac{0.1}{f_{\text{star}}} \right) \left(\frac{\bar{M}_s}{10^{-19} M_{\odot}} \right)^3 \left(\frac{10^{10} \text{ GeV}}{f_a} \right)^4$$

$$\Delta t \simeq \frac{2R_{0.1}}{v_r} \sqrt{1 - \frac{R^2}{R_{0.1}^2}} = 8 \text{ hrs} \left(\frac{10^{-19} M_{\odot}}{\bar{M}_s} \right) \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^2 \sqrt{1 - \frac{R^2}{R_{0.1}^2}}$$

Factor $\gtrsim 10^6$ enhancement compared to background DM density

Summary

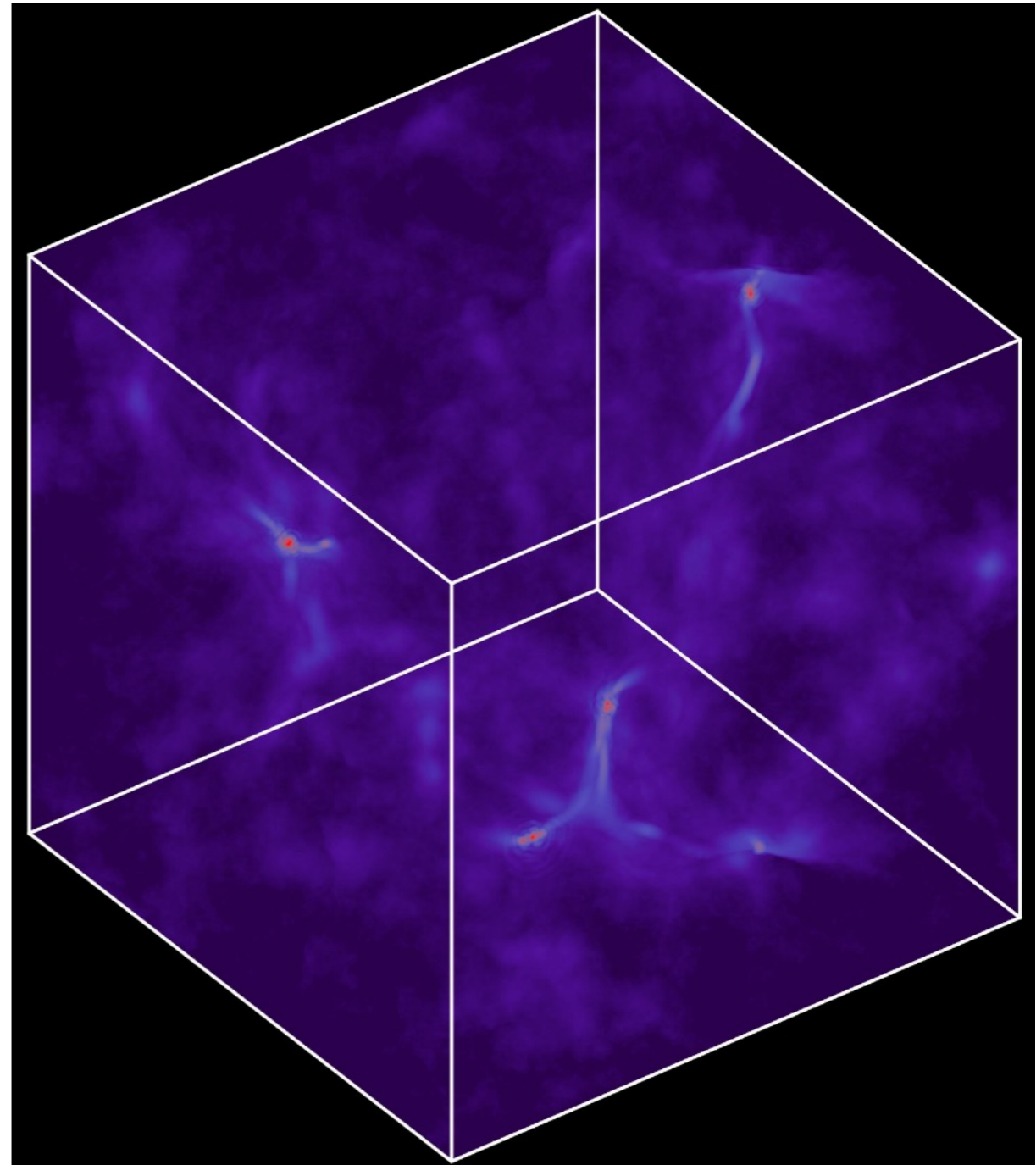
- Previously neglected self-interactions at $T \simeq \Lambda_{\text{QCD}}$ move energy to the UV
- Fluctuations on scales $k_p \simeq k_{J,\text{MRE}}$
- Structures that form around MRE are solitonic “axion stars”
- 20% of DM axions bound



Summary

- Previously neglected self-interactions at $T \simeq \Lambda_{\text{QCD}}$ move energy to the UV
- Fluctuations on scales $k_p \simeq k_{J,\text{MRE}}$
- Structures that form around MRE are solitonic “axion stars”
- 20% of DM axions bound

Thanks



Schrödinger Poisson & QP

$$\left(i\partial_t + \frac{\nabla^2}{2m_a} - m_a\Phi \right) \psi = 0$$

Magdelung Transformation

$$\psi = \sqrt{\rho} e^{i\theta}$$

$$\vec{v} = \frac{1}{m_a} \nabla \theta$$

$$\nabla^2 \Phi = \frac{4\pi G}{R} \left(|\psi|^2 - \langle |\psi|^2 \rangle \right)$$



Continuity: $\partial_t \rho_i + 3H\rho_i + R^{-1} \nabla \cdot (\rho \vec{v}) = 0$

Euler: $\partial_t \vec{v} + H\vec{v} + R^{-1} (\vec{v} \cdot \nabla) \vec{v} = -R^{-1} (\nabla \Phi + \nabla \Phi_Q)$

Perfect fluid with

“quantum pressure”: $\Phi_Q \equiv -\frac{\hbar^2}{2R^2 m_a^2} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}}$

Quantum pressure

negligible if $\nabla \Phi \gg \nabla \Phi_Q$

$$4\pi G R^2 \rho \gg \frac{1}{2R^2 m_a^2} \nabla^2 \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \sim \frac{1}{4R^2 m_a^2} k^4$$

$$k \ll k_J$$

Overdensities dominated by Φ

→ grow and collapse

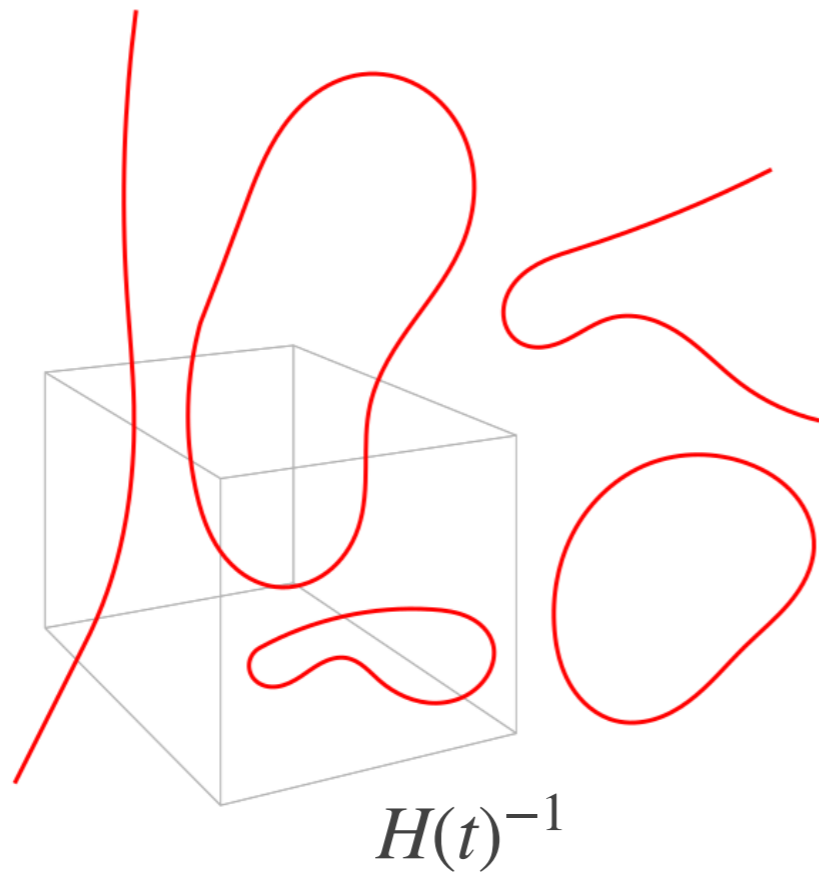
$$\frac{k_J}{R} = (16\pi G \rho m_a^2)^{1/4}$$

$$k \gg k_J$$

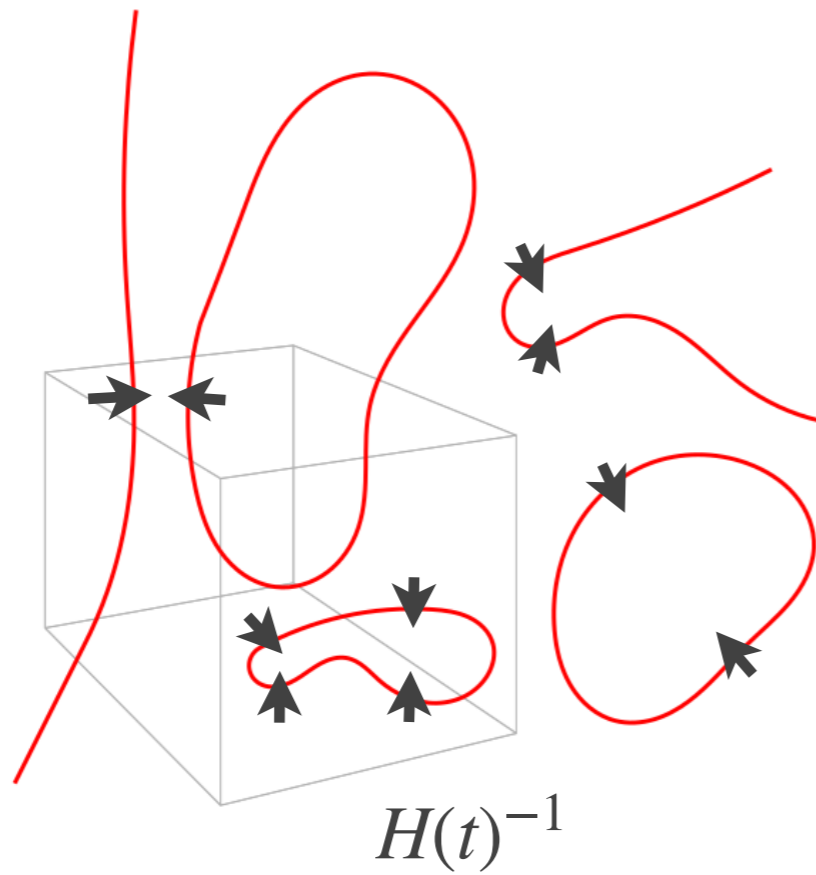
Overdensities dominated by Φ_Q

→ prevented from collapsing and oscillate

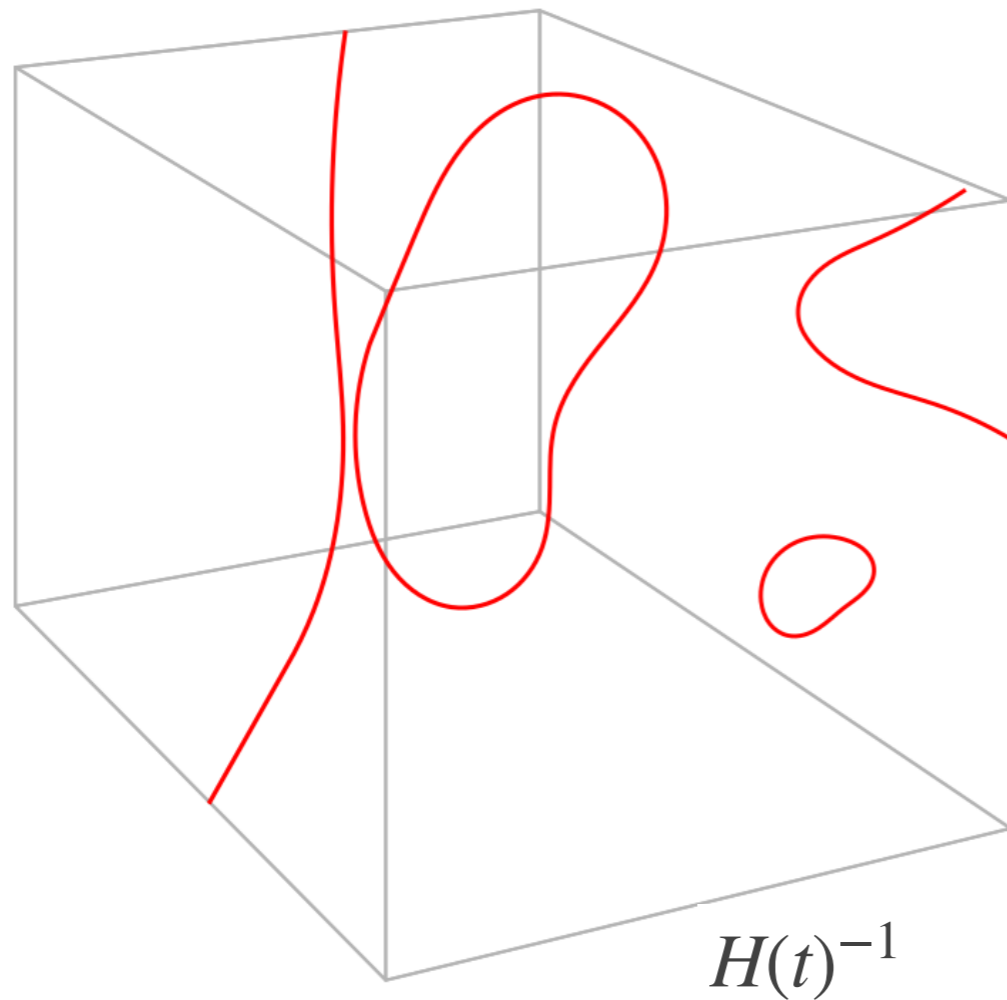
Scaling



Scaling

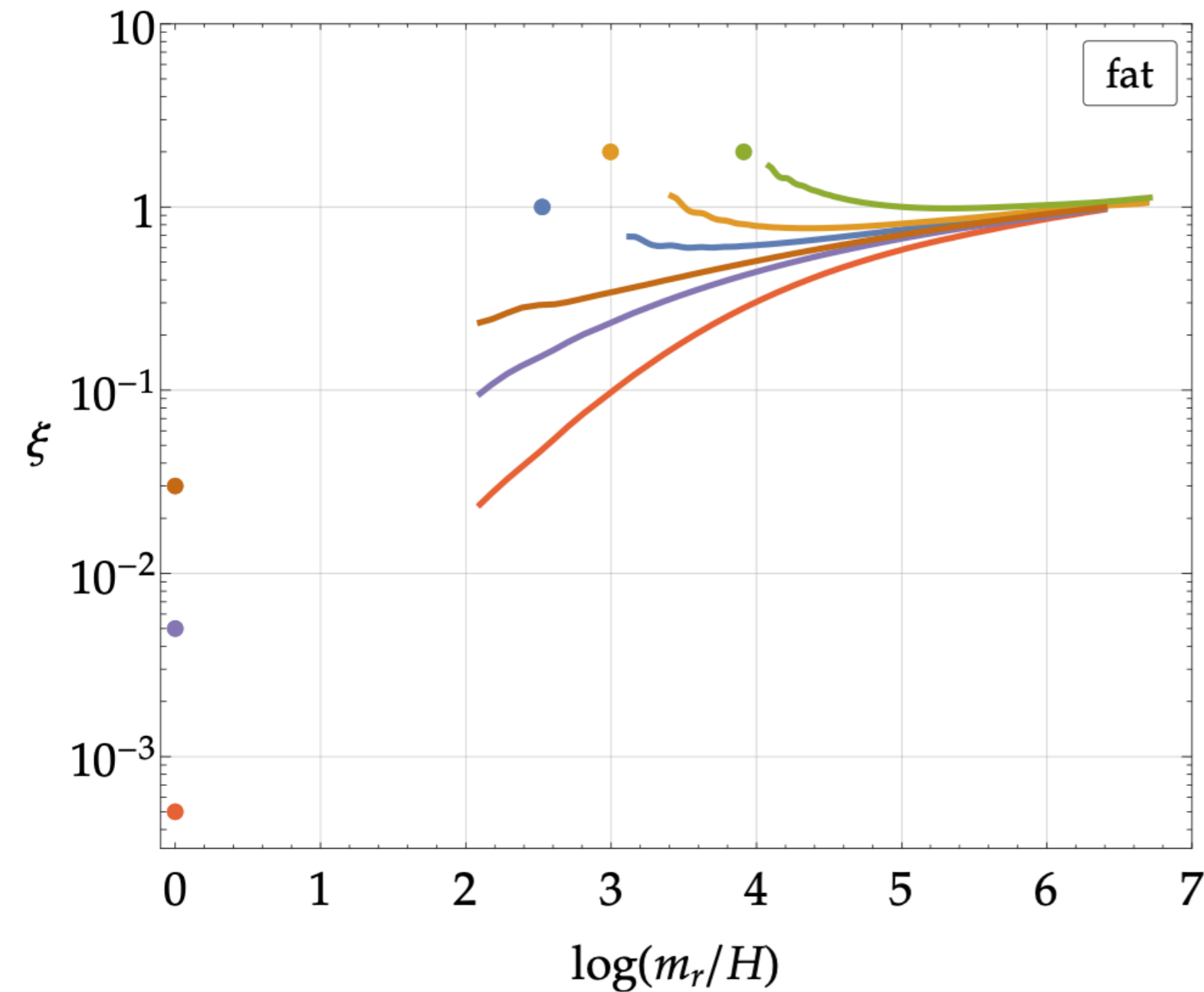


Scaling

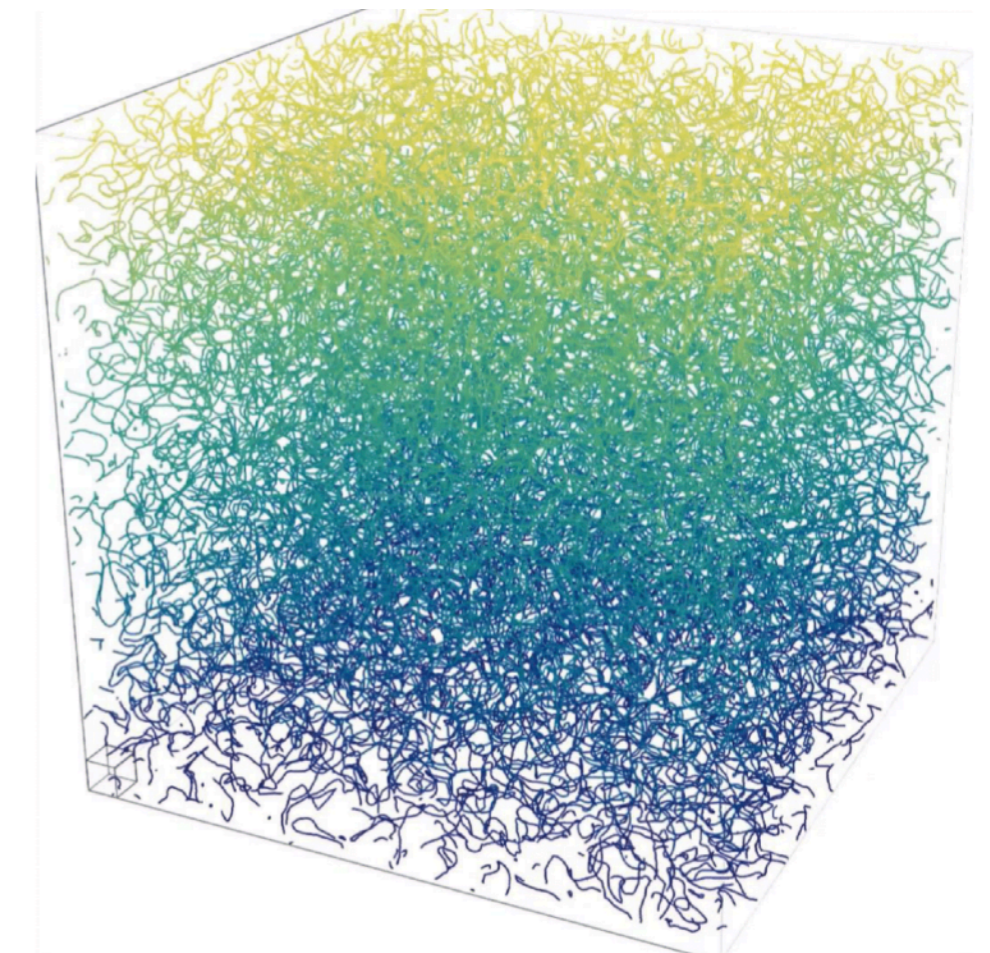


$\xi(t) =$ Length of string in
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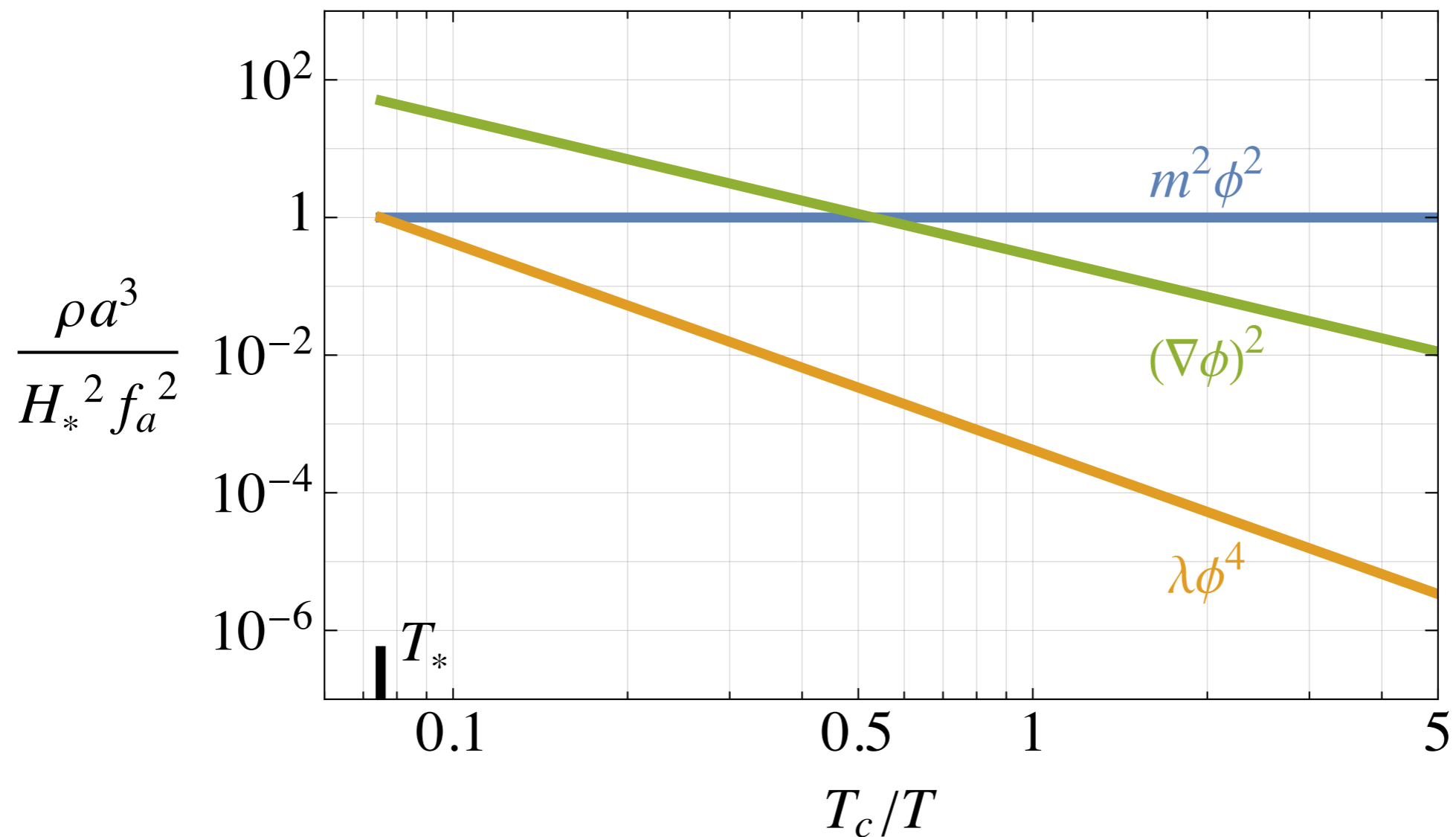


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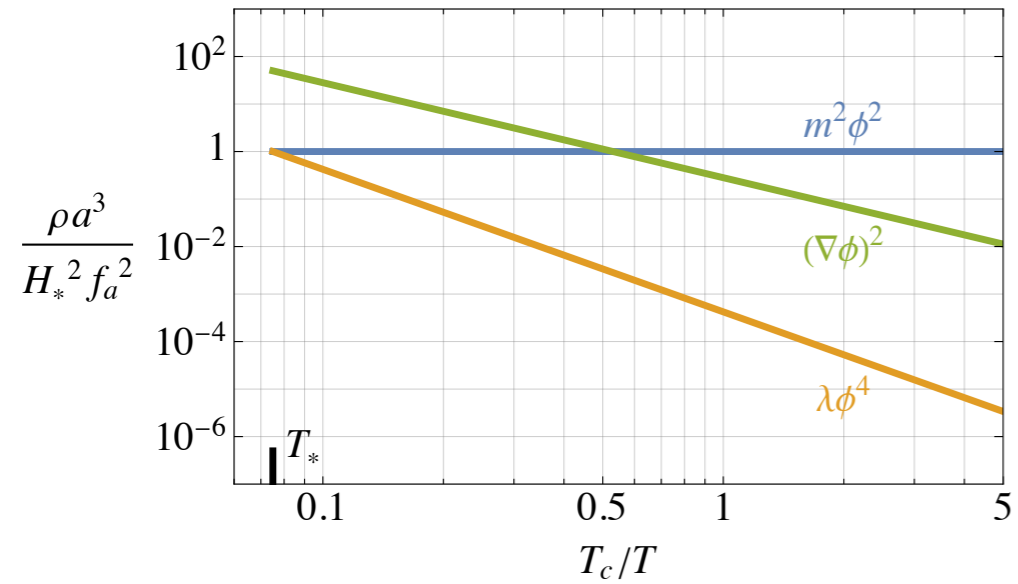
Self-interactions

ALP:



Self-interactions

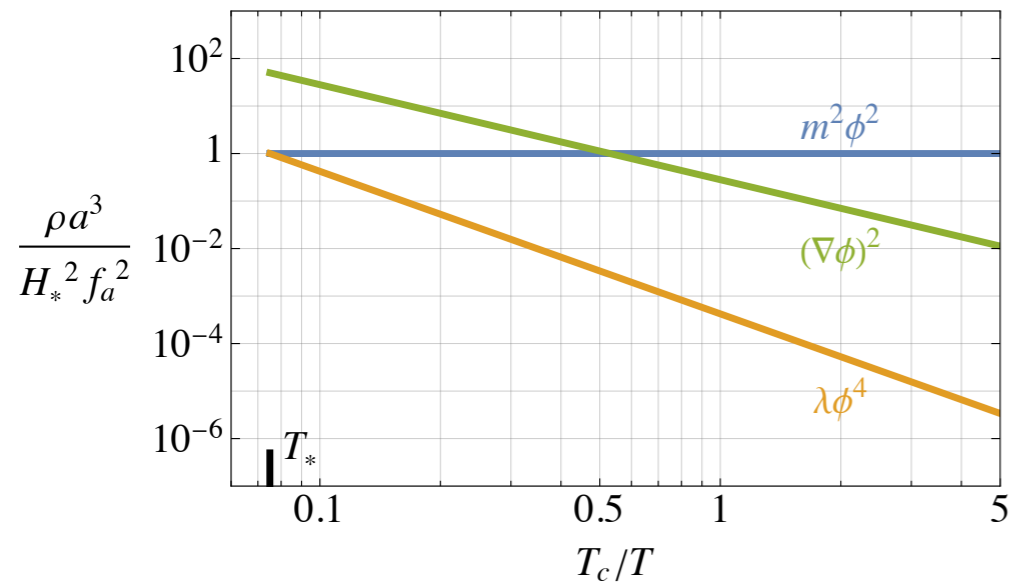
ALP:



QCD axion:

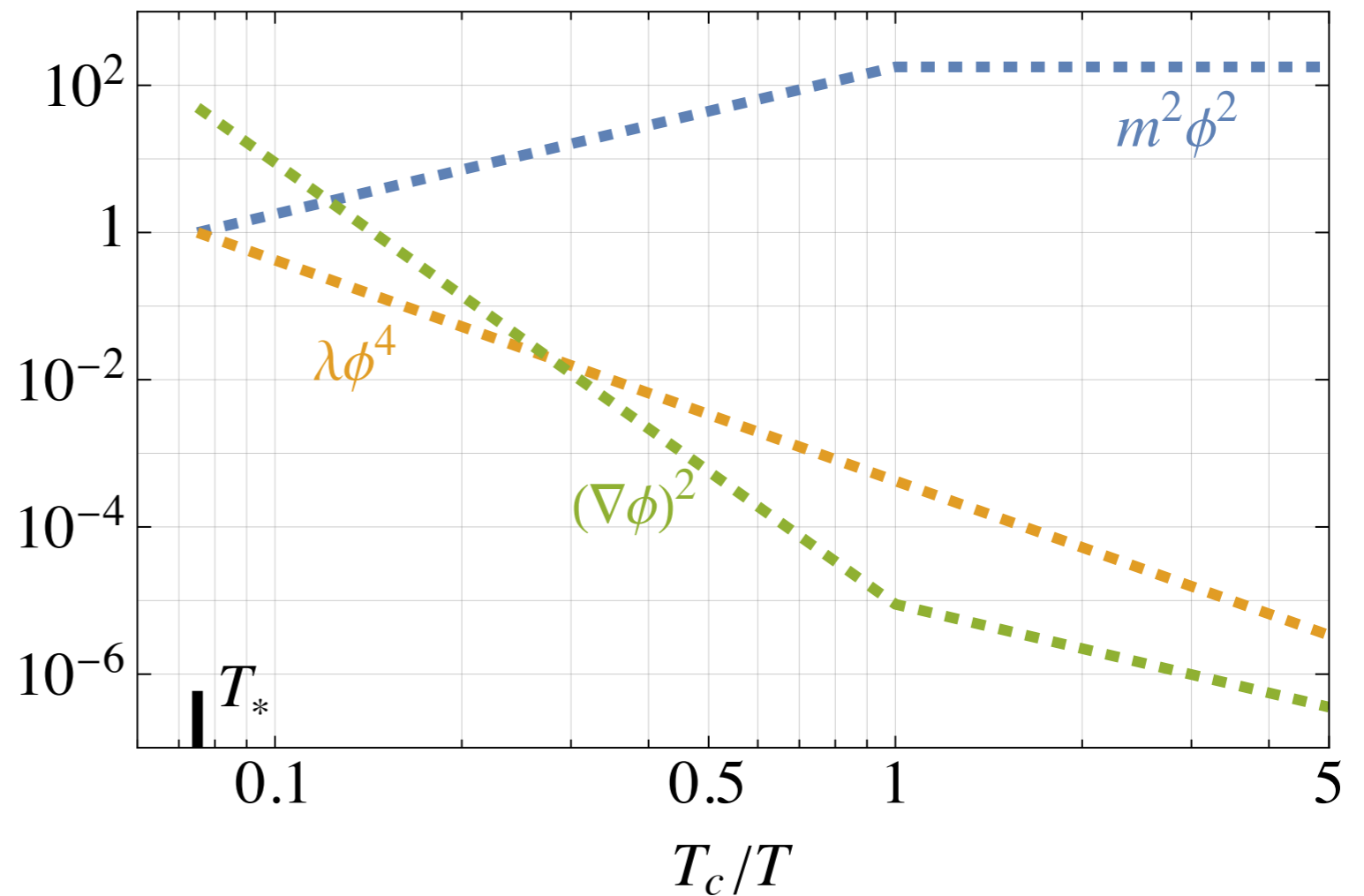
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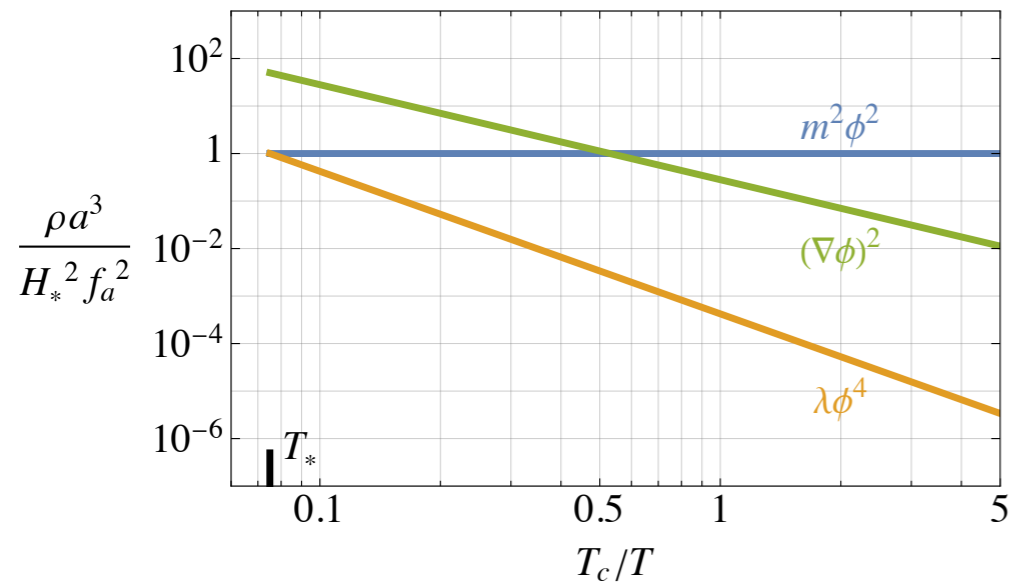
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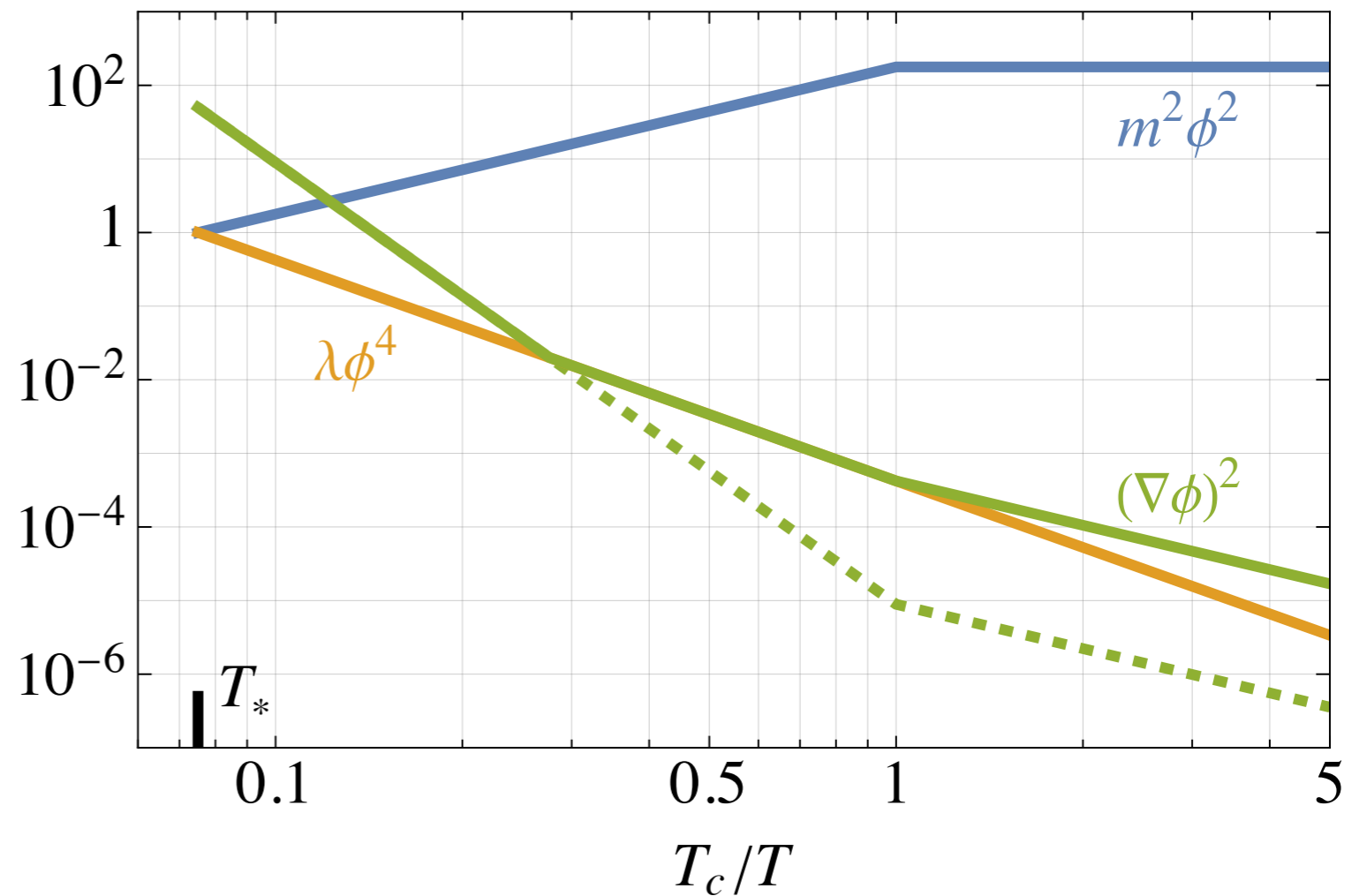
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Self-interactions

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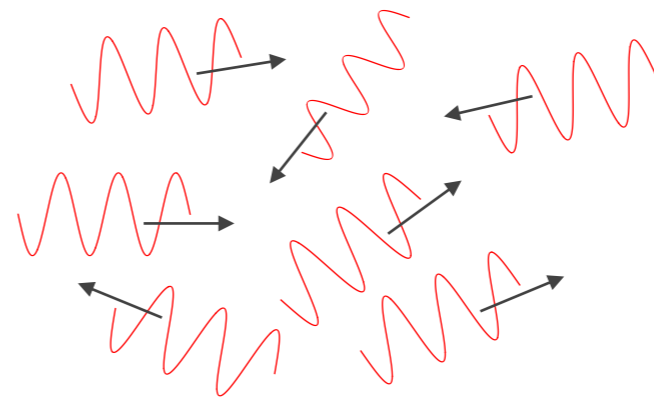
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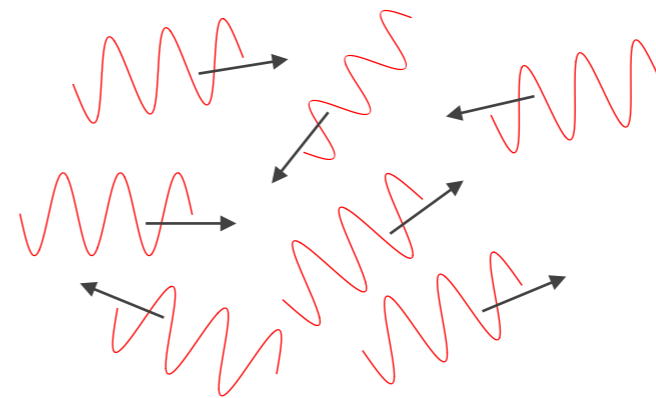
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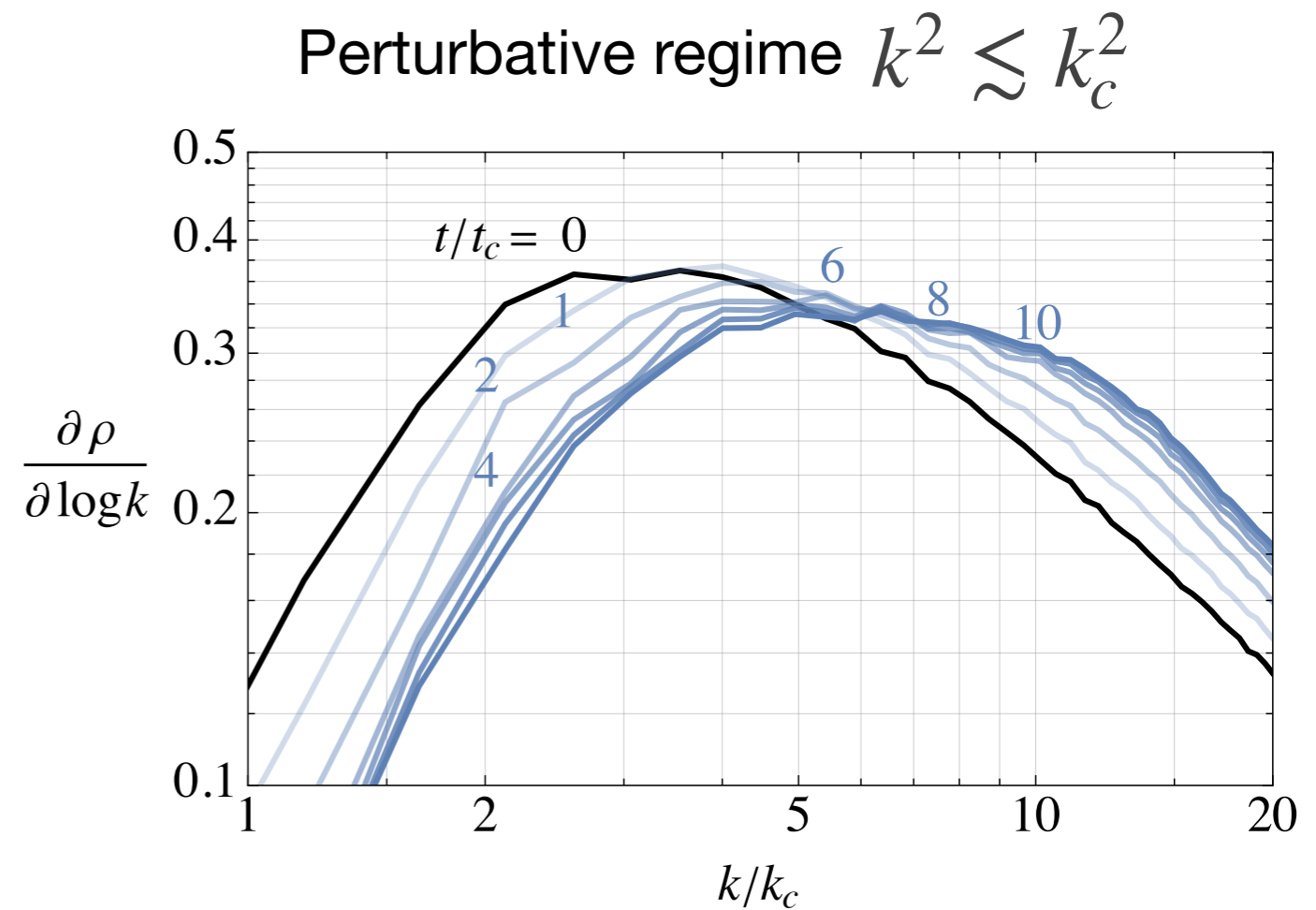
The non-perturbative regime $k^2 \lesssim k_c^2$



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Flat space

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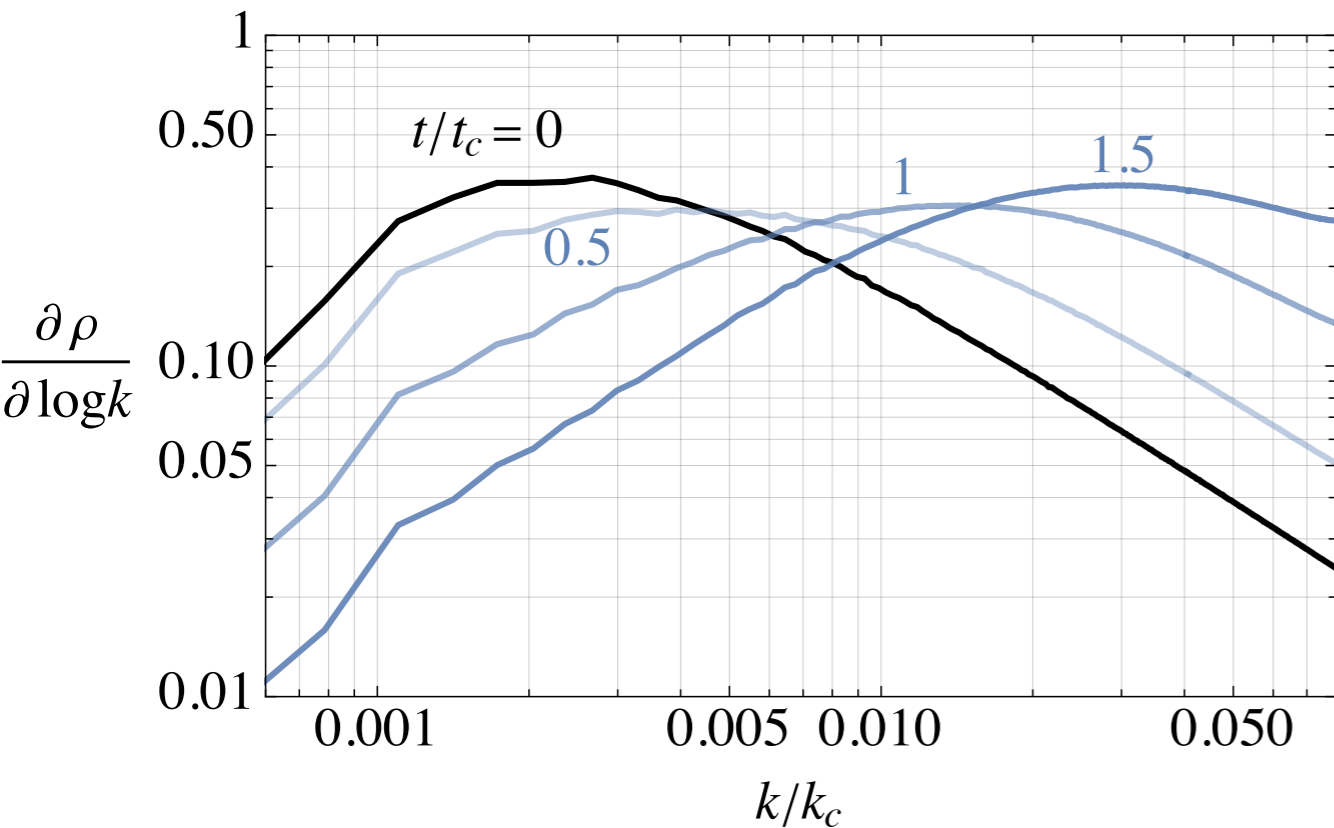


$$t_{\text{rel}} \gg t_c$$

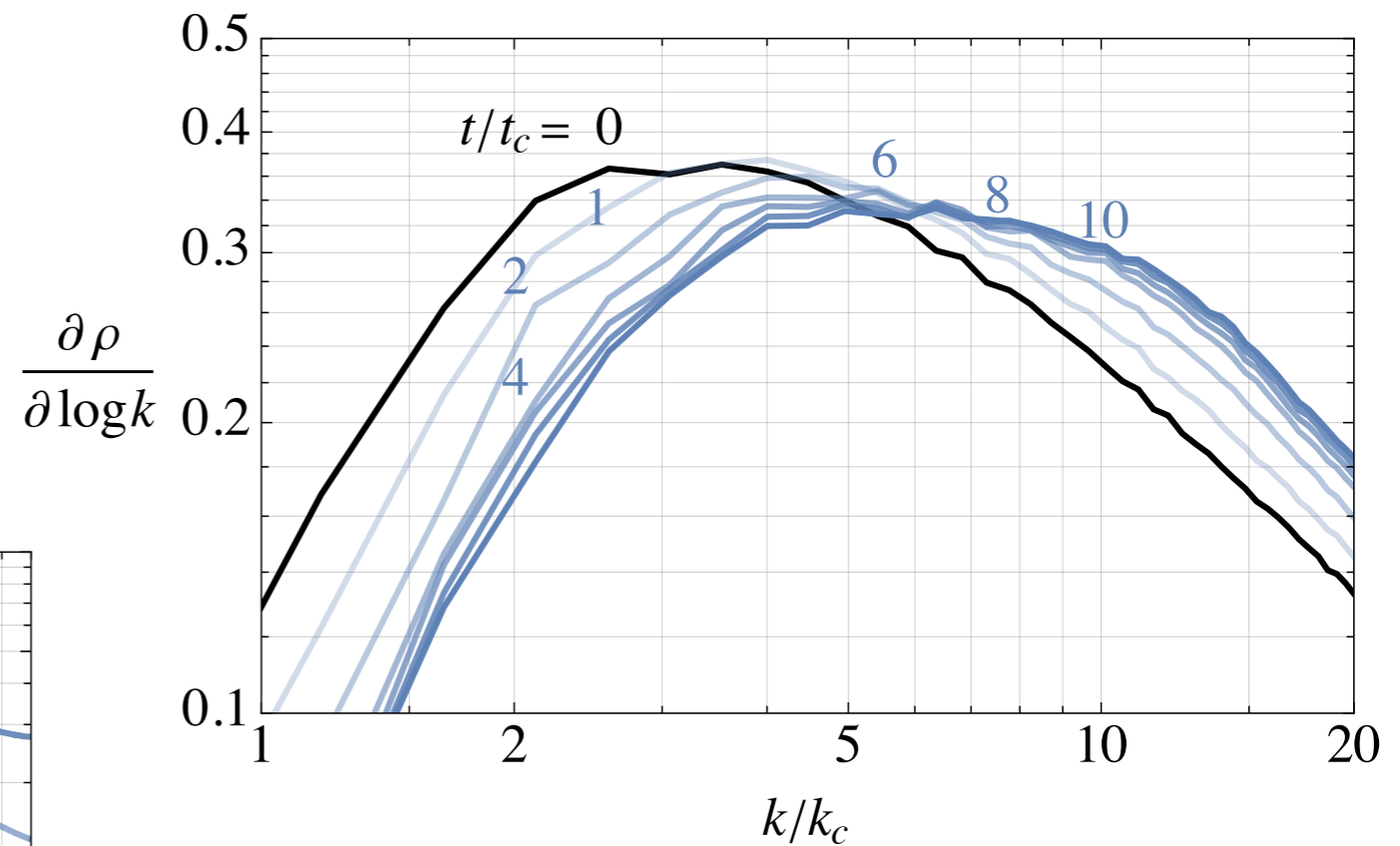
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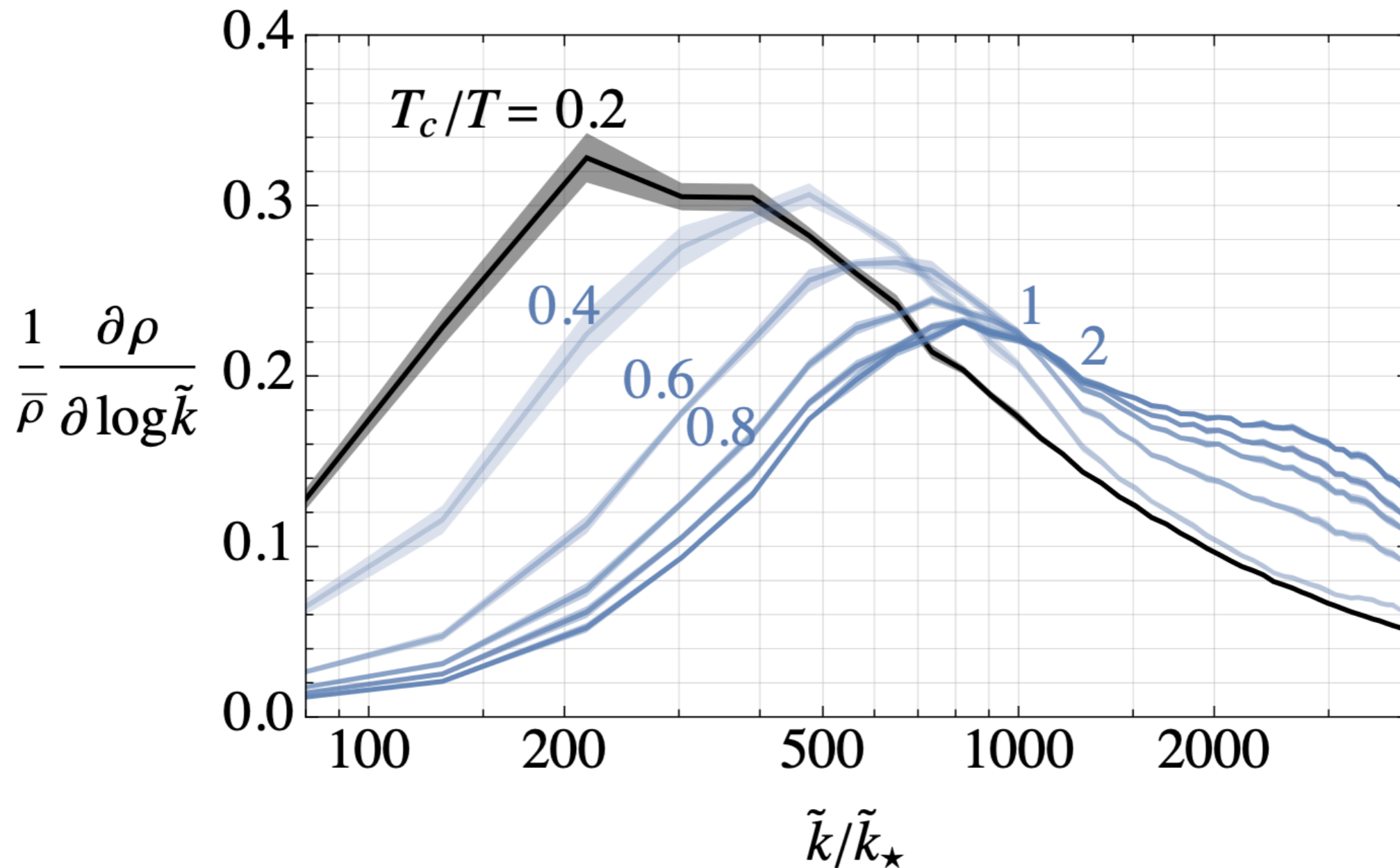
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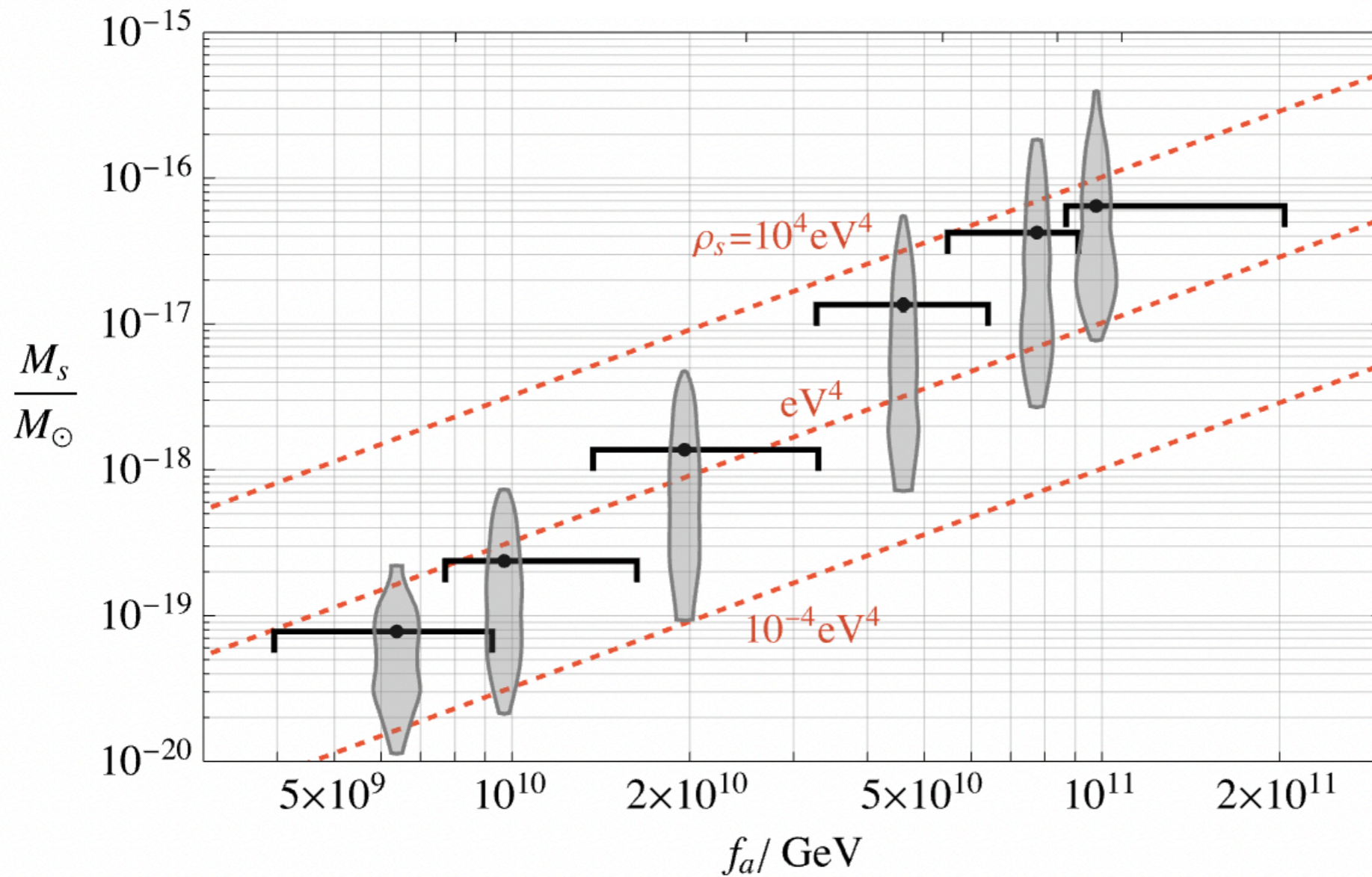
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Results from simulations

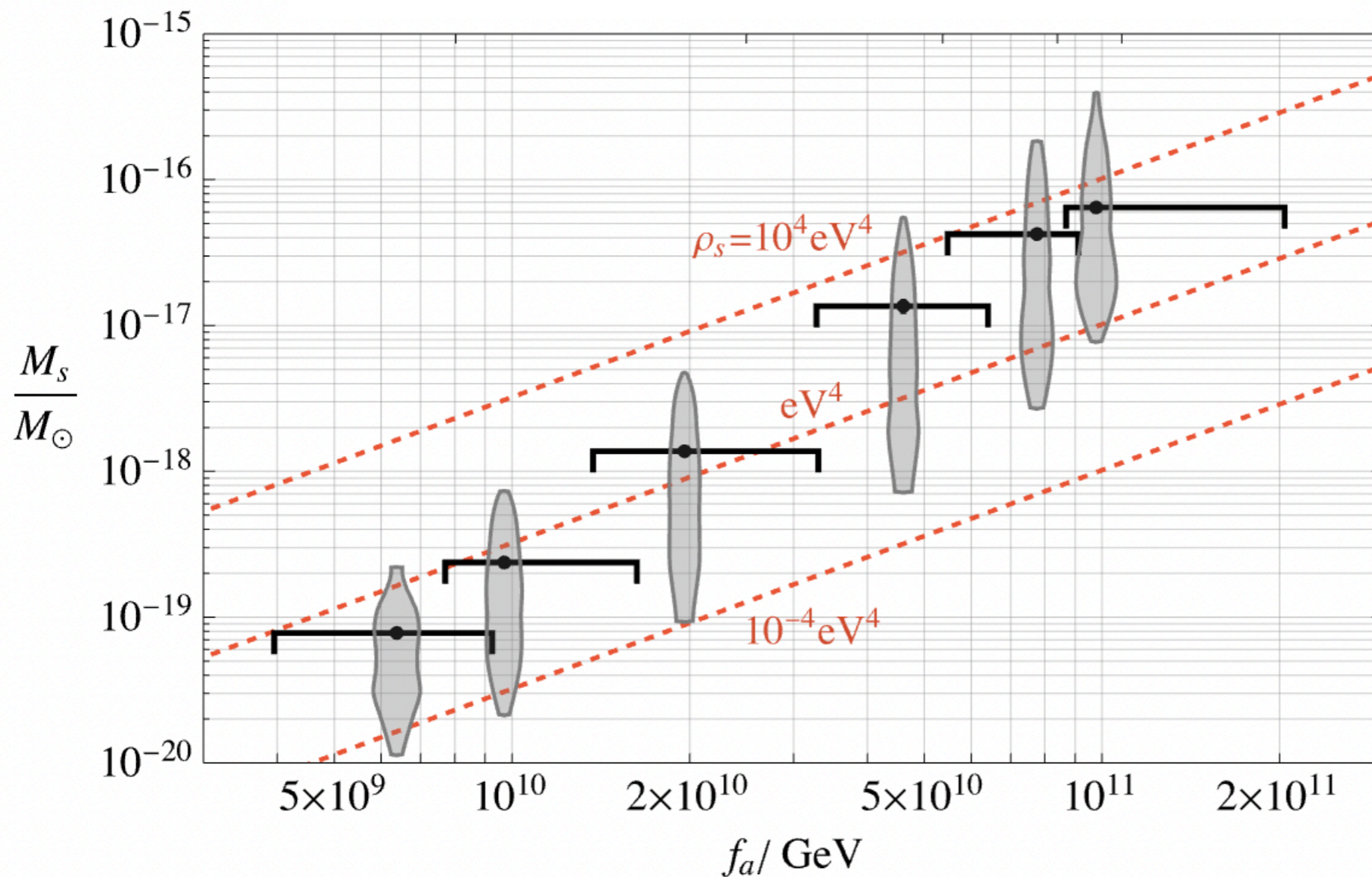
$$f_a = 10^{10} \text{ GeV}$$



Properties of the substructure

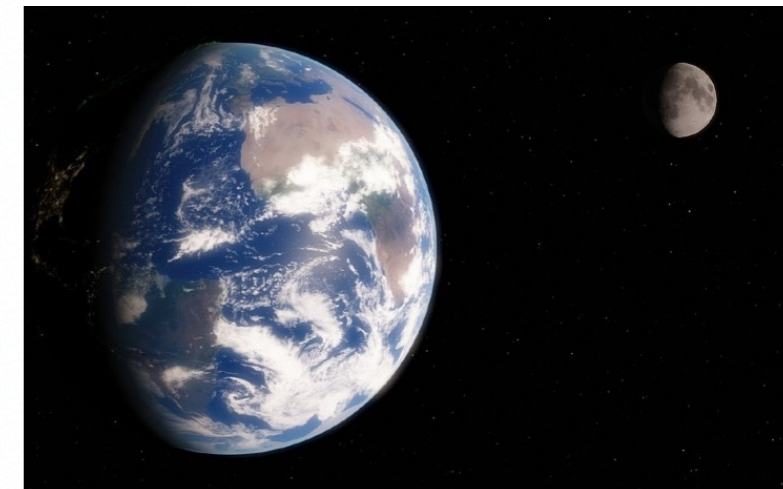
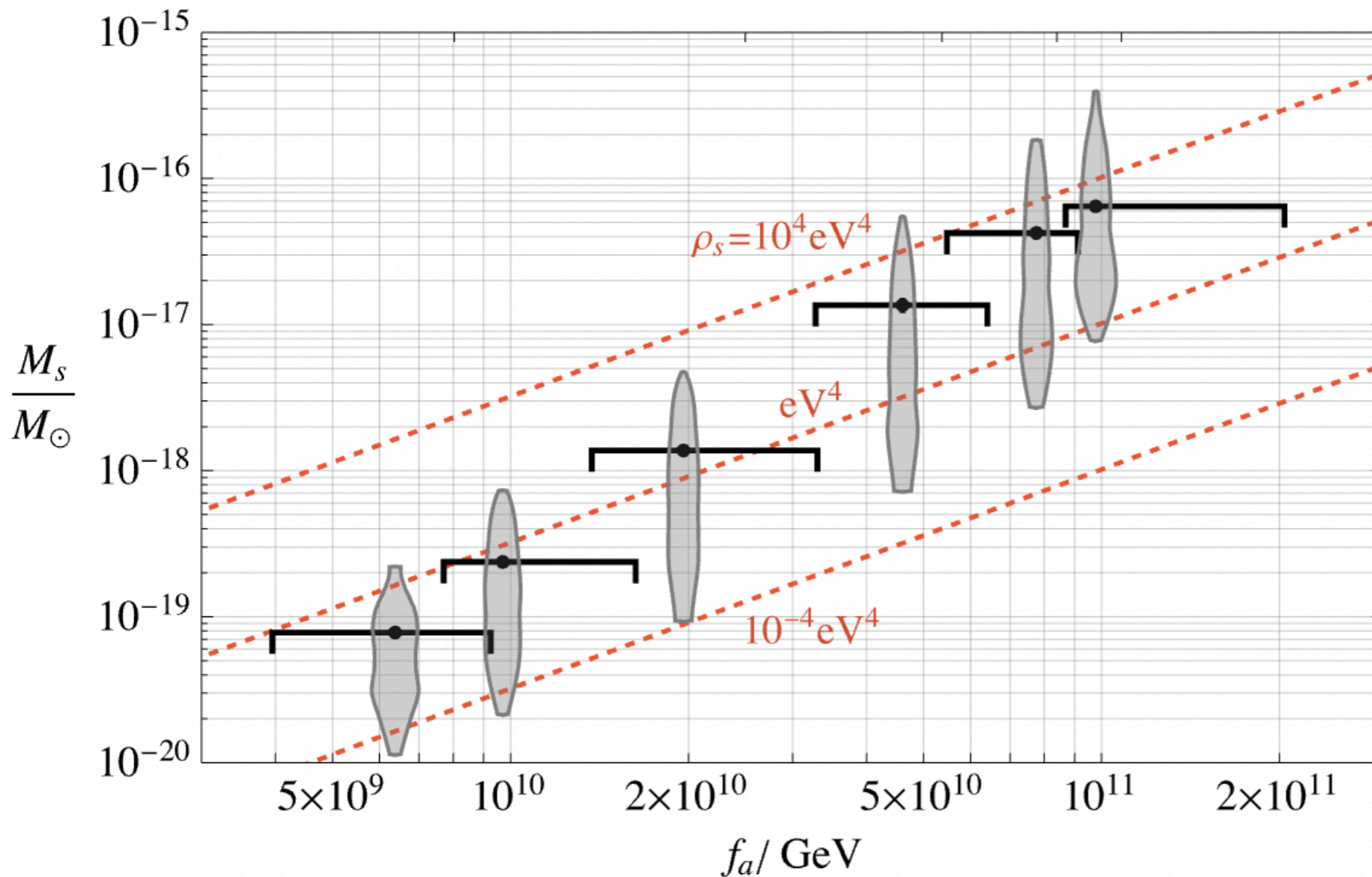


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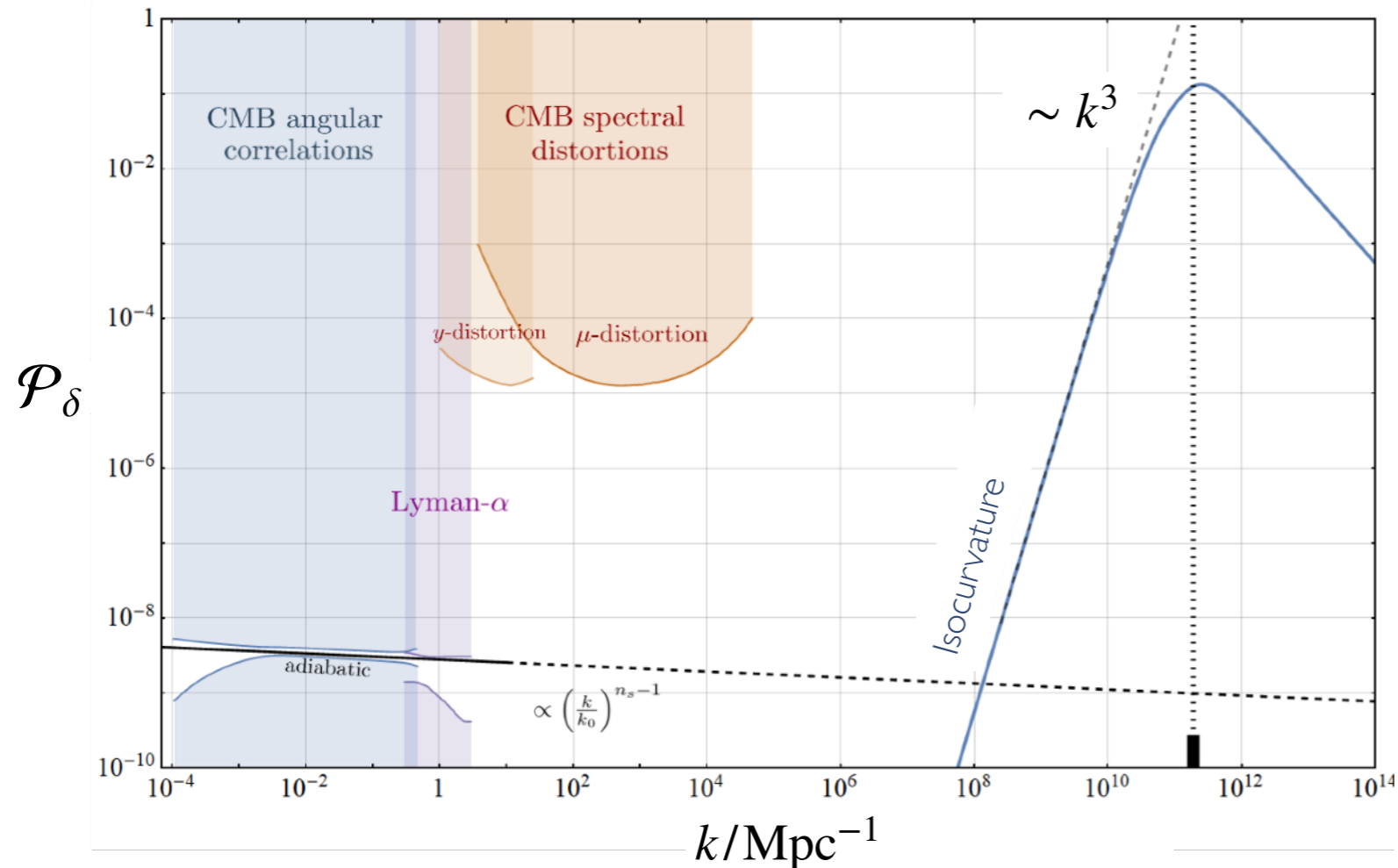
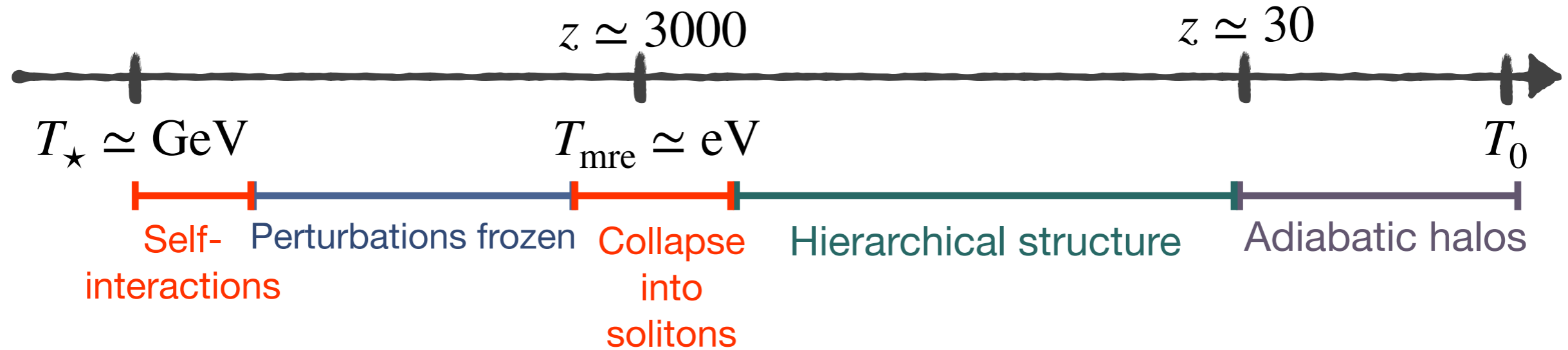
$$R_{0.1} \simeq 4.2 \times 10^6 \text{ km} \left(\frac{f_a}{10^{10} \text{ GeV}} \right)^2 \left(\frac{10^{-19} M_\odot}{M_s} \right)$$

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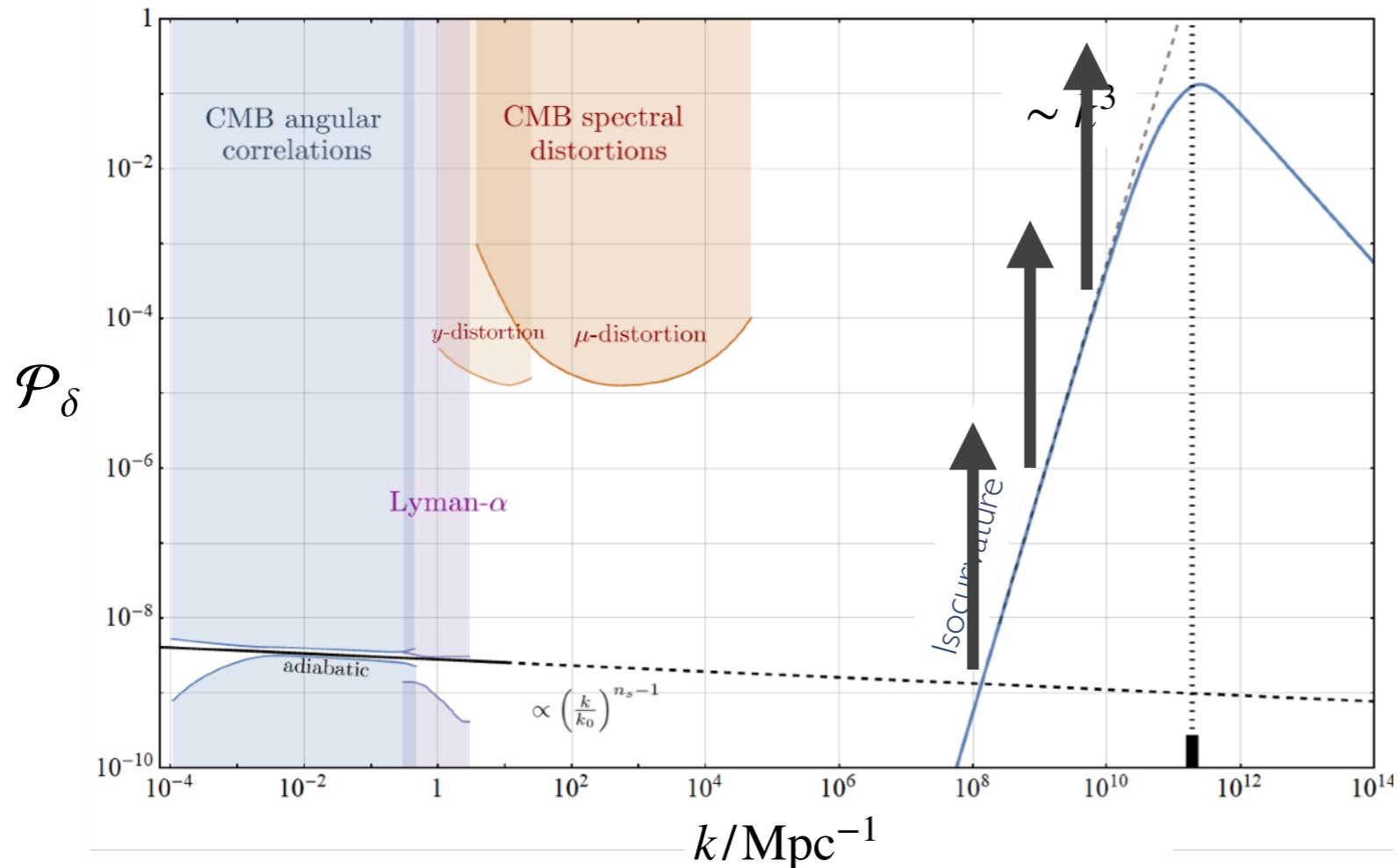
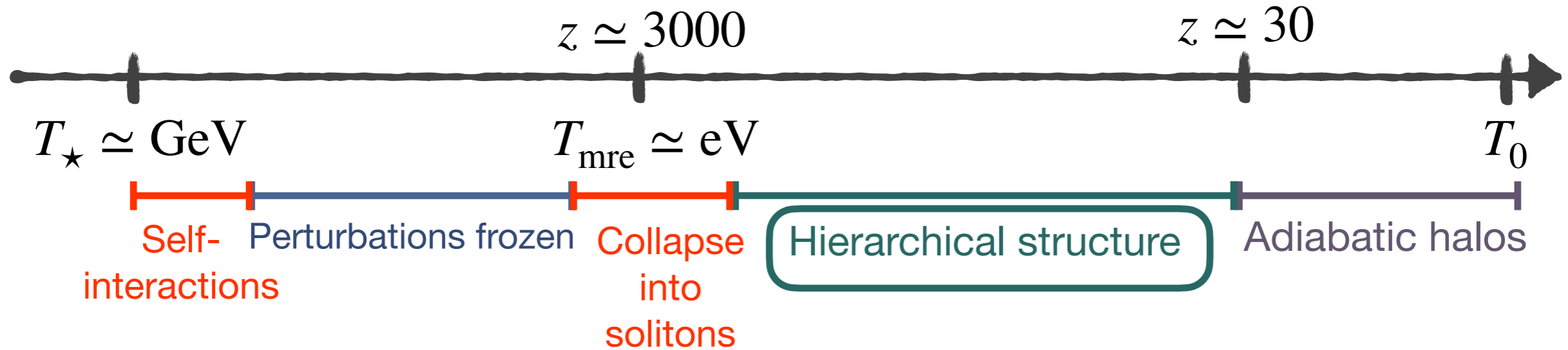


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Subsequent evolution?

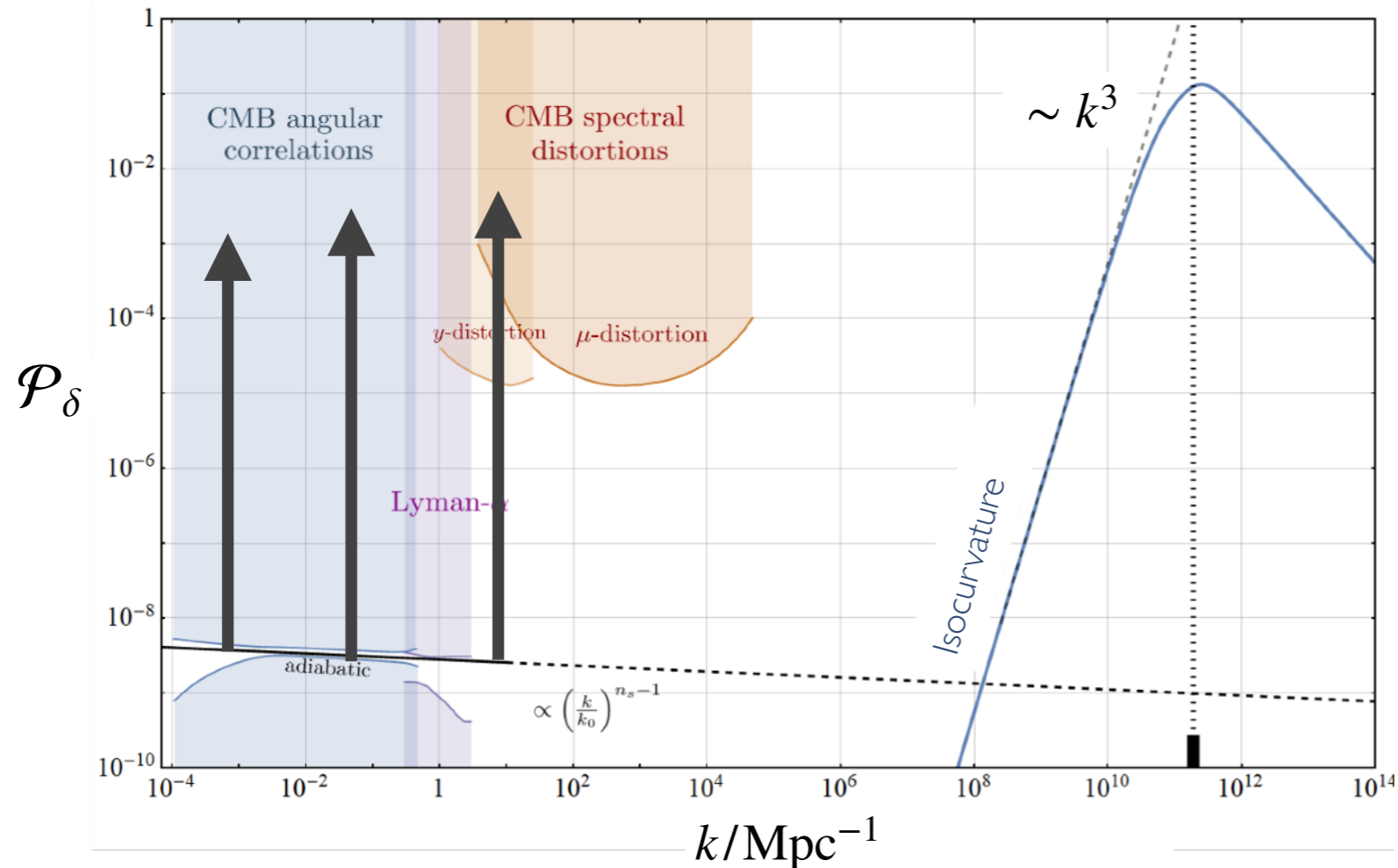
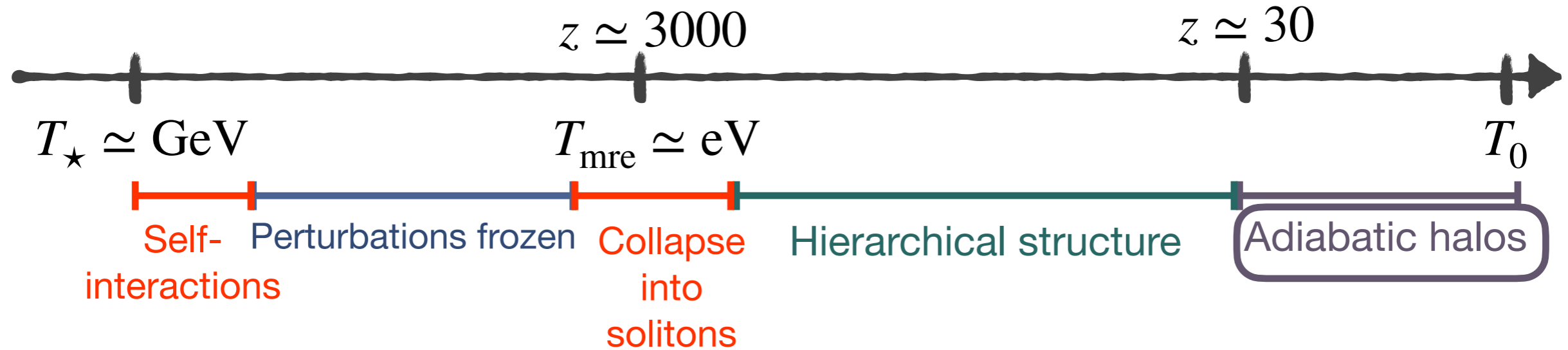


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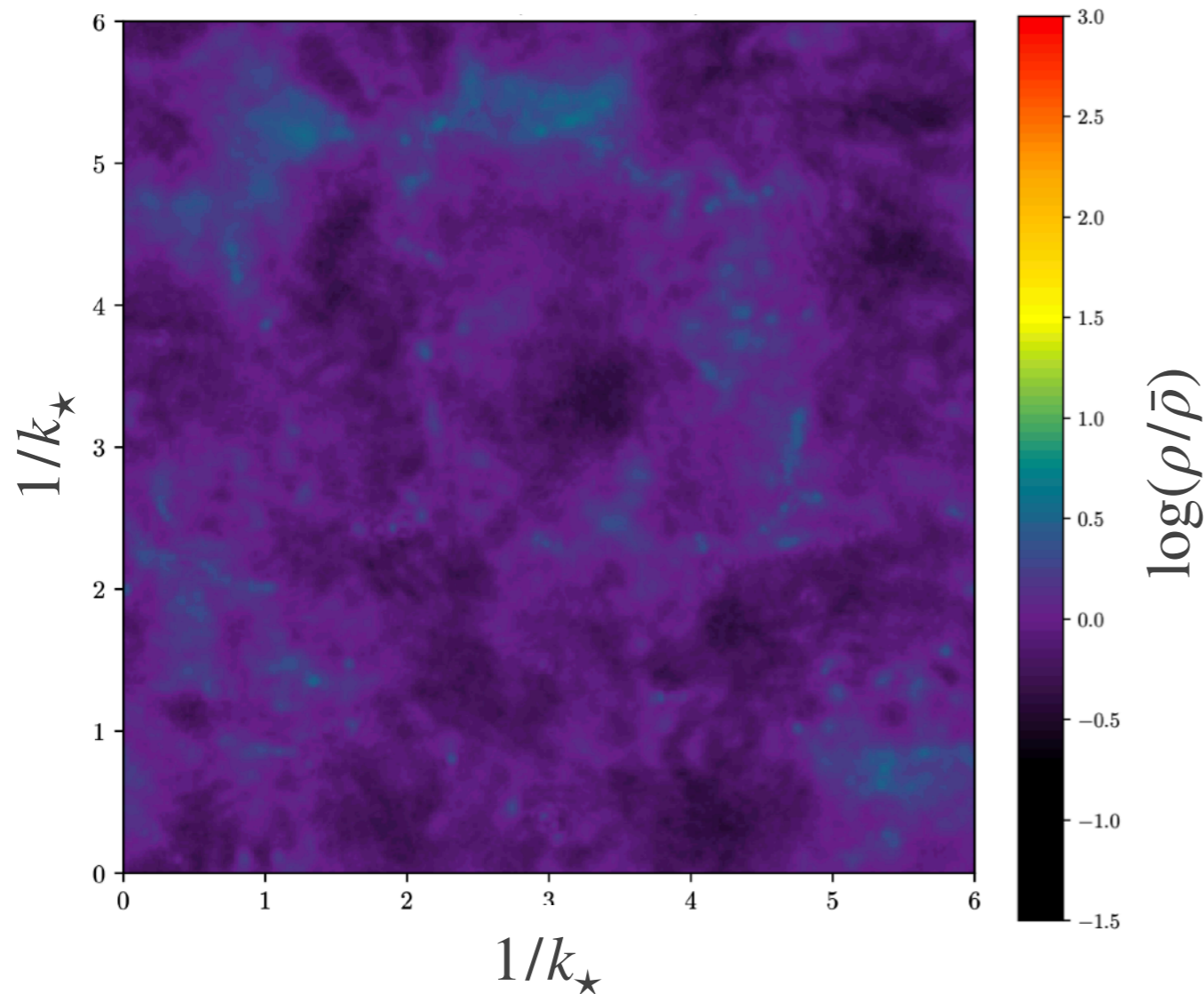
[Eggemeier et al]

Subsequent evolution?



Initial perturbations

Order one fluctuations on co-moving scales $\simeq H_\star$ when $H = m_a(T)$



$$\delta(x) \equiv \frac{\rho(x) - \bar{\rho}}{\bar{\rho}}$$

Dark matter density

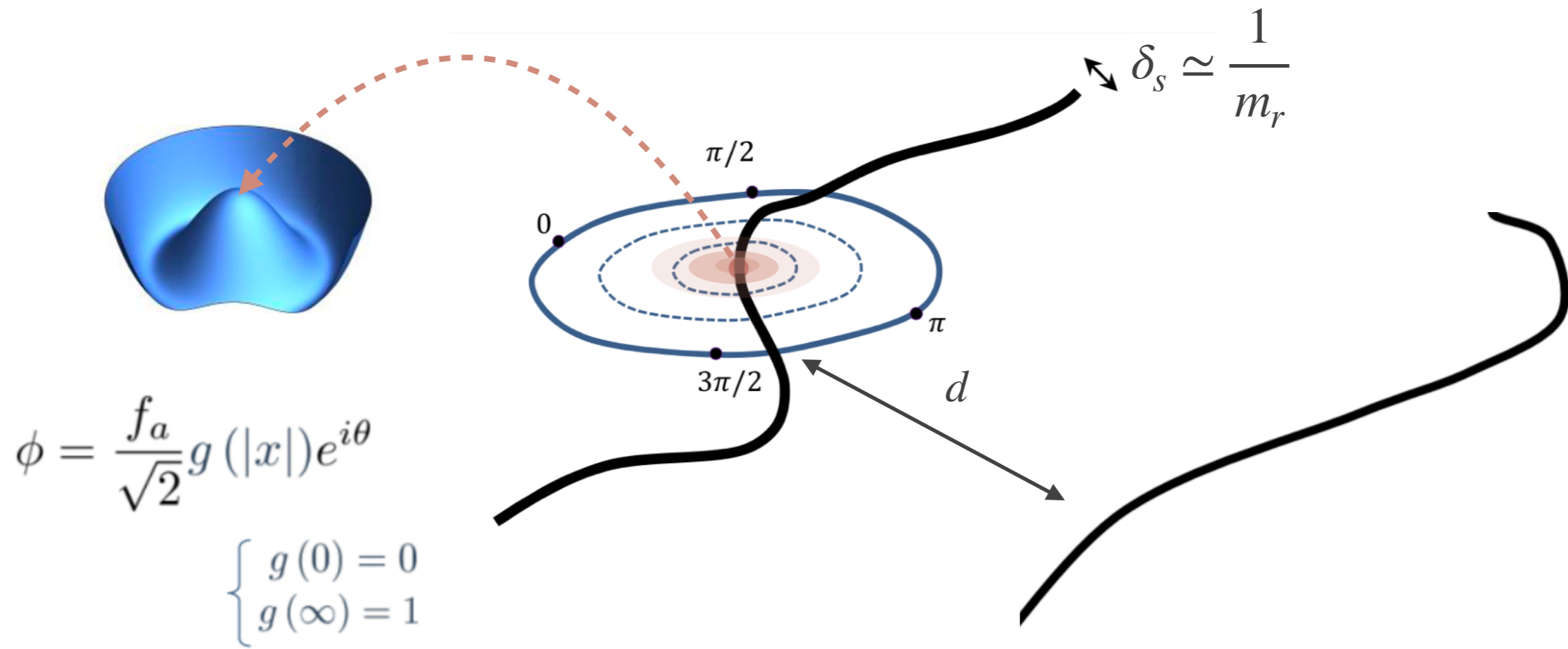
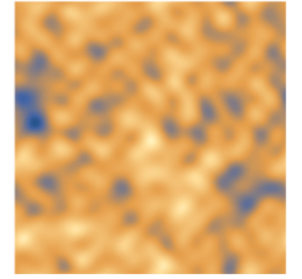
Density power spectrum

$$\langle \tilde{\delta}^*(\vec{k}) \tilde{\delta}(\vec{k}') \rangle \equiv \frac{2\pi^2}{k^3} \delta^3(\vec{k} - \vec{k}') \mathcal{P}_\delta(|\vec{k}|)$$

$$\mathcal{P}_\delta(k) \sim \frac{\partial \rho}{\partial \log k}$$

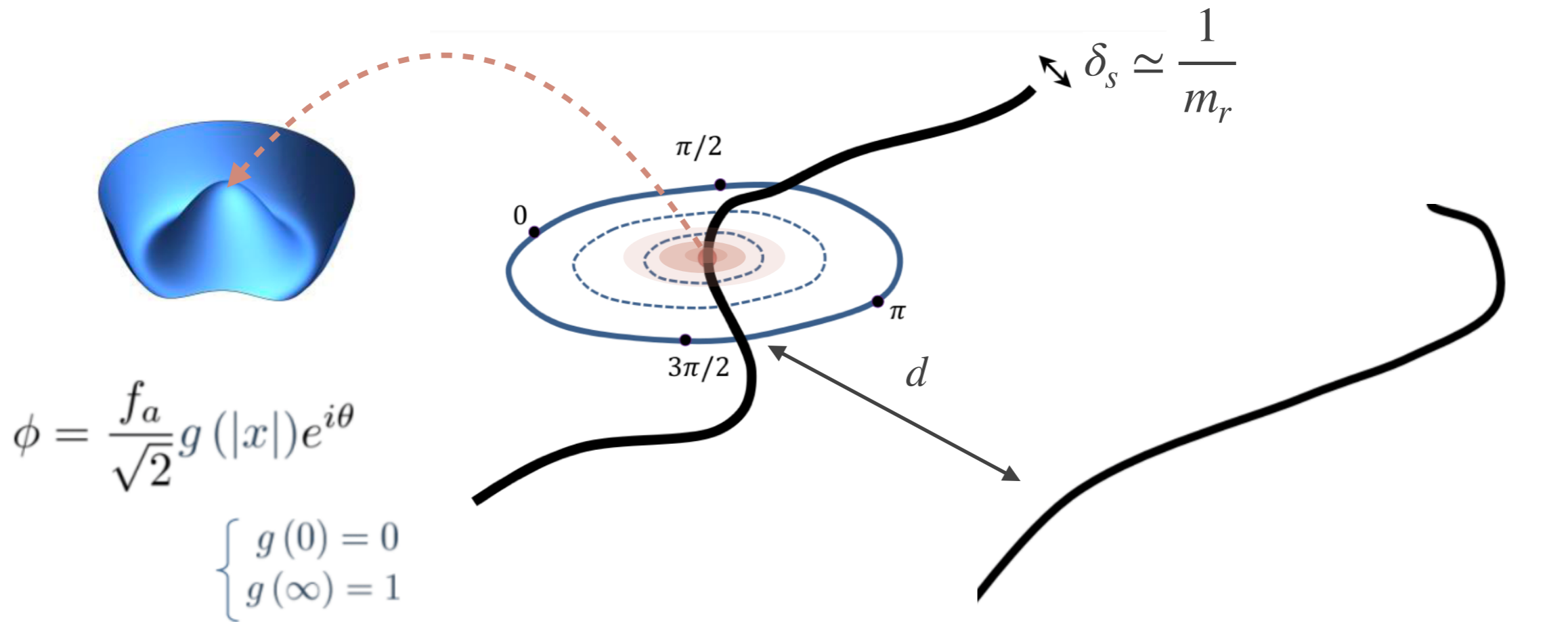
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Topological strings




$$\mu = \frac{E}{L} \sim \underbrace{\pi f_a^2}_{\text{Core}} \underbrace{\log \frac{d}{m_r^{-1}}}_{\text{Gradient}}$$

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$H \sim T^2 / M_{\text{Pl}}$


 Grows logarithmically
 with time