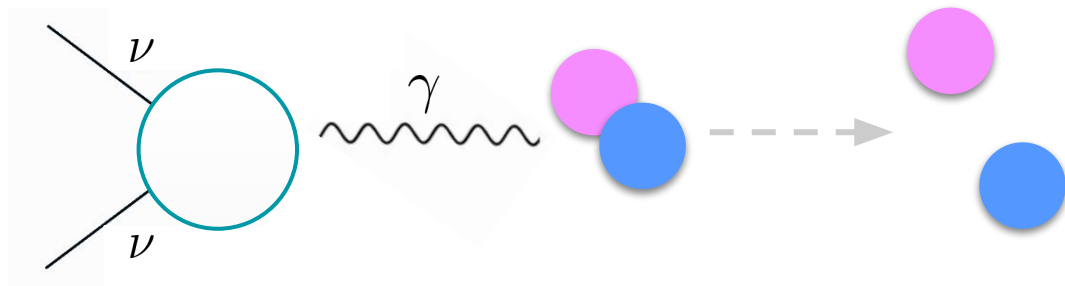


BBN photodisintegration limits from neutrino injections

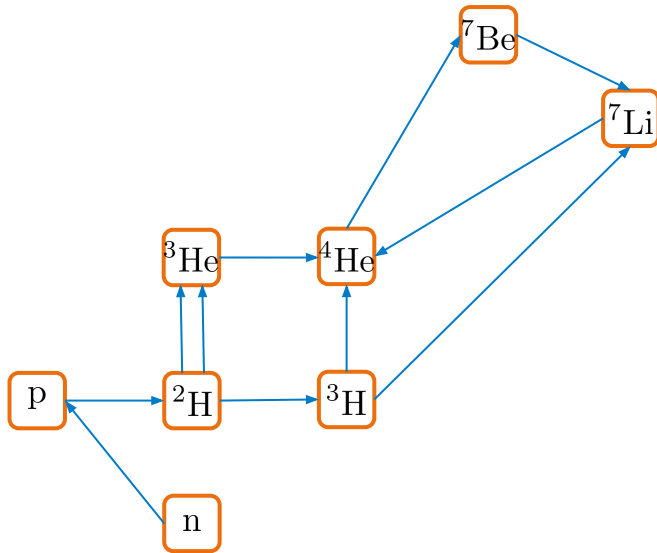
Based on 2407.xxxxx

Sara Bianco
Invisibles2024 Workshop, Bologna
1st July 2024

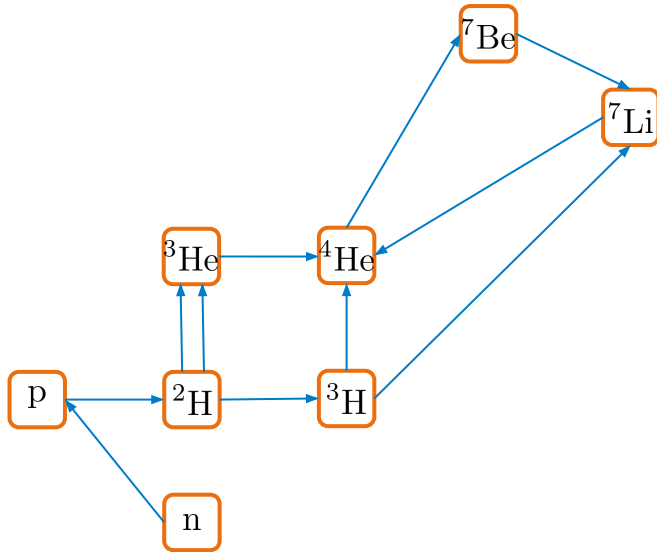


In collaboration with: P. F. Depta, J. Frerick, T. Hambye, M. Hufnagel, and K. Schmidt-Hoberg

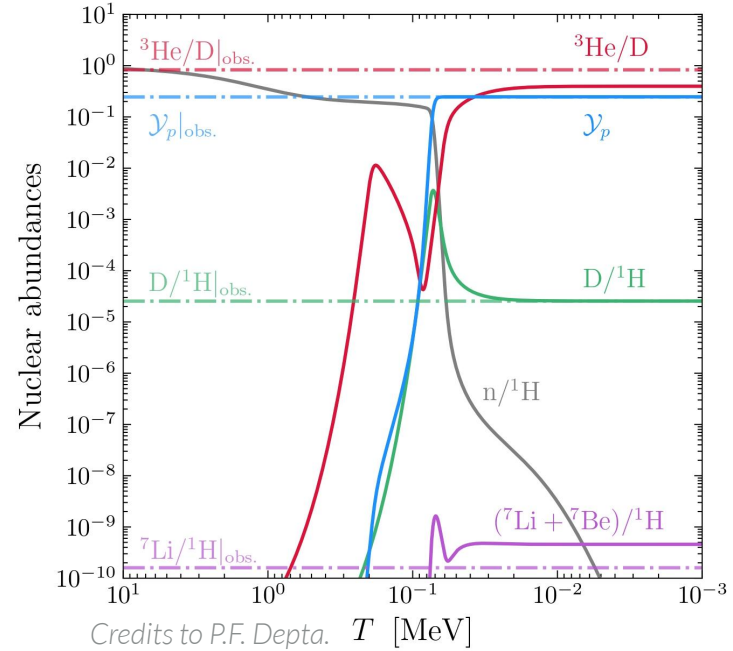
$$t \sim 1\text{s} - 10^3\text{s} \iff T \sim \text{MeV} - \text{keV}$$



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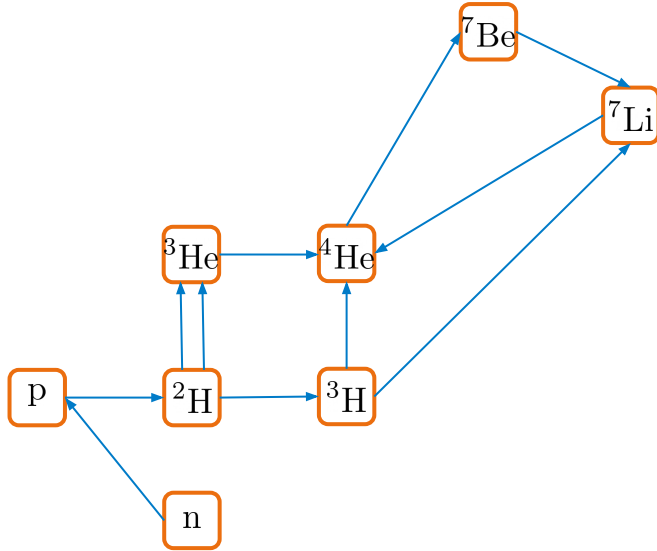


Agreement between predictions and observations:

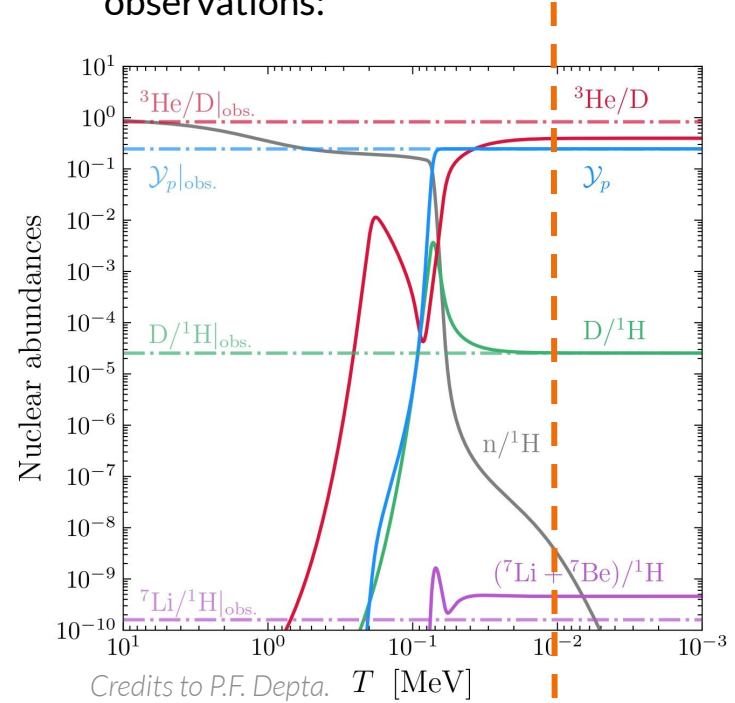


$< 1\sigma$

$$t \sim 1\text{s} - 10^3\text{s} \iff T \sim \text{MeV} - \text{keV}$$

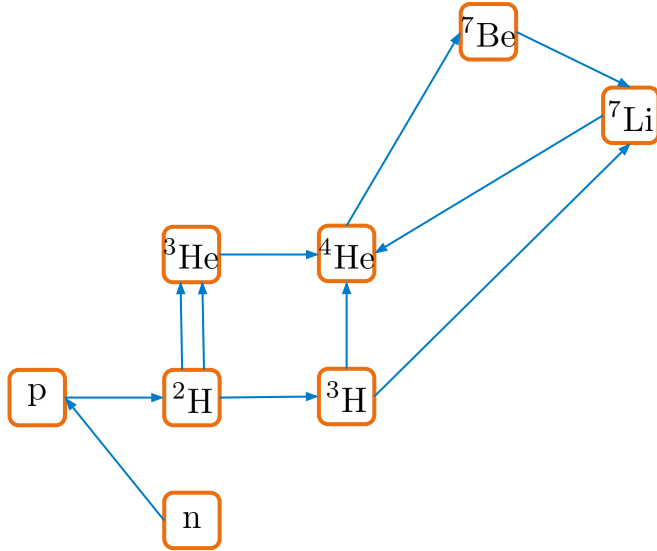


Agreement between predictions and observations:

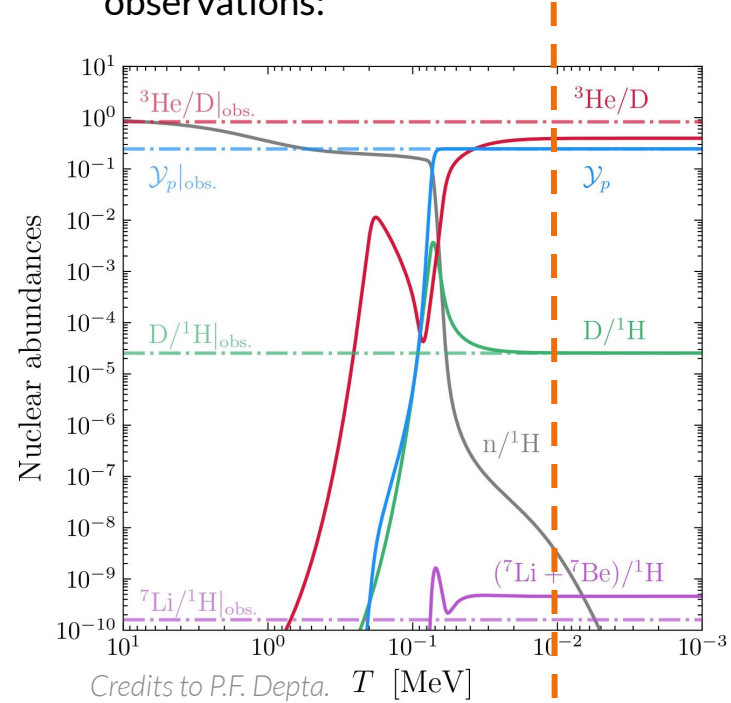


Fixed around 10 keV!

$$t \sim 1\text{s} - 10^3\text{s} \iff T \sim \text{MeV} - \text{keV}$$



Agreement between predictions and observations:



What happens if we inject electromagnetic material at *later times*?

M. Kawasaki and T. Moroi, astro-ph/9412055.

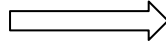
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$$E > 1.5 \text{ MeV}$$

$$t > 10^4 \text{ s}$$

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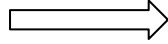
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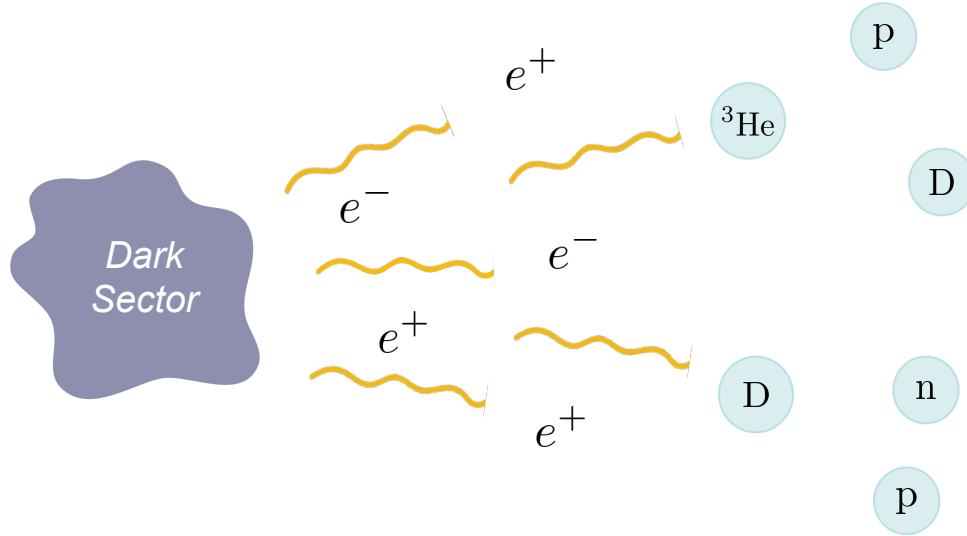
Electromagnetic cascade that leads to **non-thermal** parts of the photon, electron, and positron spectra.

What happens if we inject electromagnetic material at later times?

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Electromagnetic cascade that leads to **non-thermal** parts of the photon, electron, and positron spectra.



What if we inject *neutrinos*?

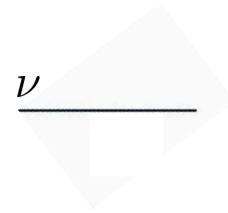
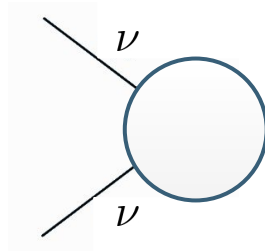
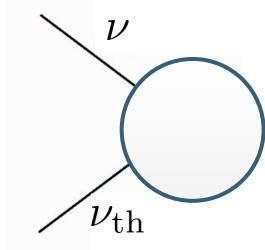
T. Hambye, M. Hufnagel and M. Lucca, 2112.09137.

T. Kanzaki, M. Kawasaki, K. Kohri and T. Moroi, 0705.1200.

$$\phi \rightarrow \nu\nu$$

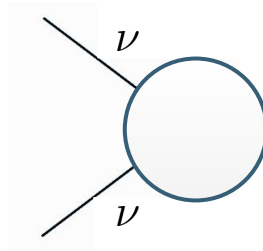
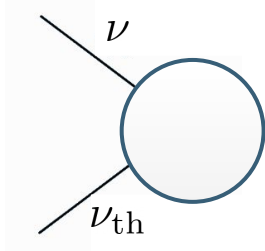
What if we inject *neutrinos*?

$$\phi \rightarrow \nu\nu$$



What if we inject *neutrinos*?

$$\phi \rightarrow \nu\nu$$



$$s_{\max} = 4EE'$$

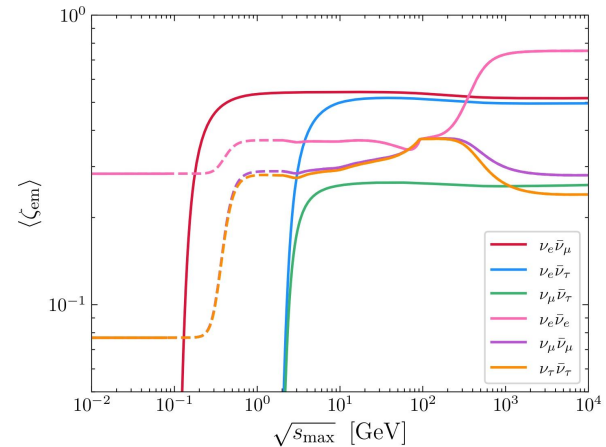
$$\nu\nu_{\text{th}} \rightarrow \nu\nu$$

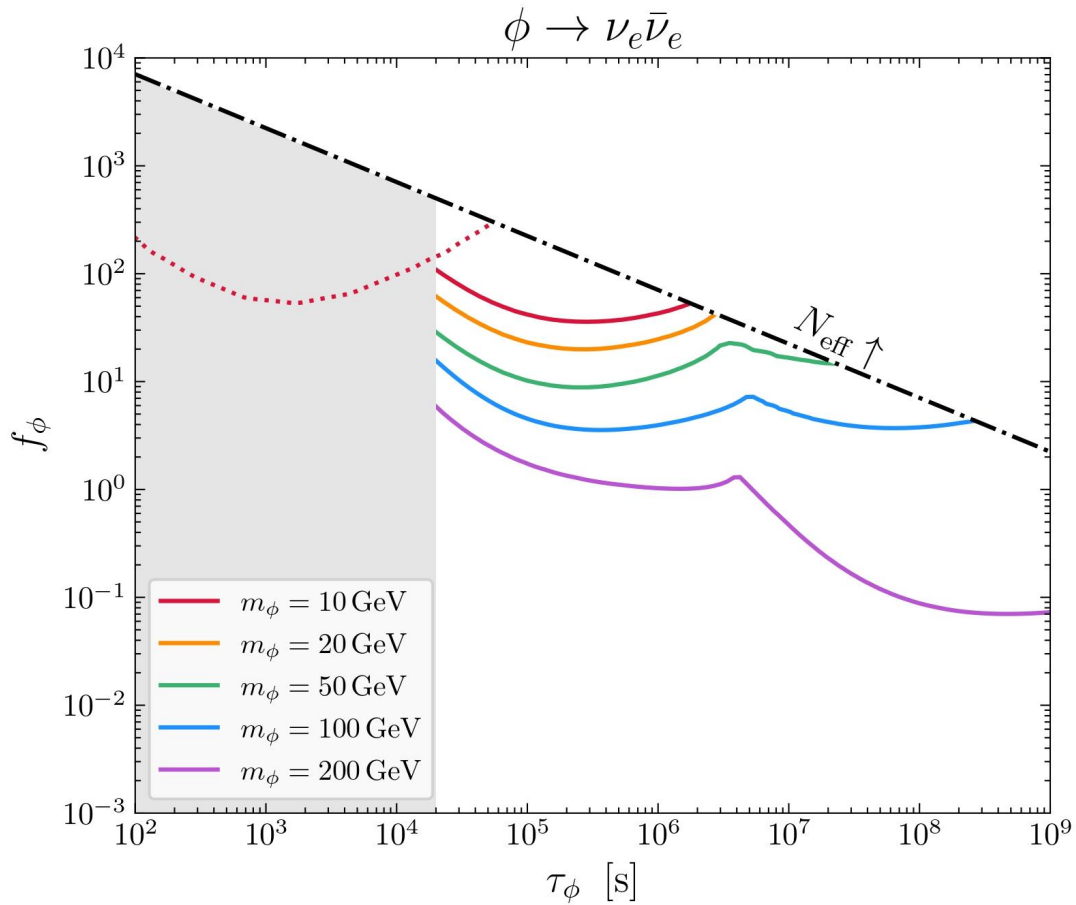
$$\nu\bar{\nu}_{\text{th}} \rightarrow \nu\bar{\nu}$$

$$\nu\bar{\nu}_{\text{th}} \rightarrow e^+e^-$$

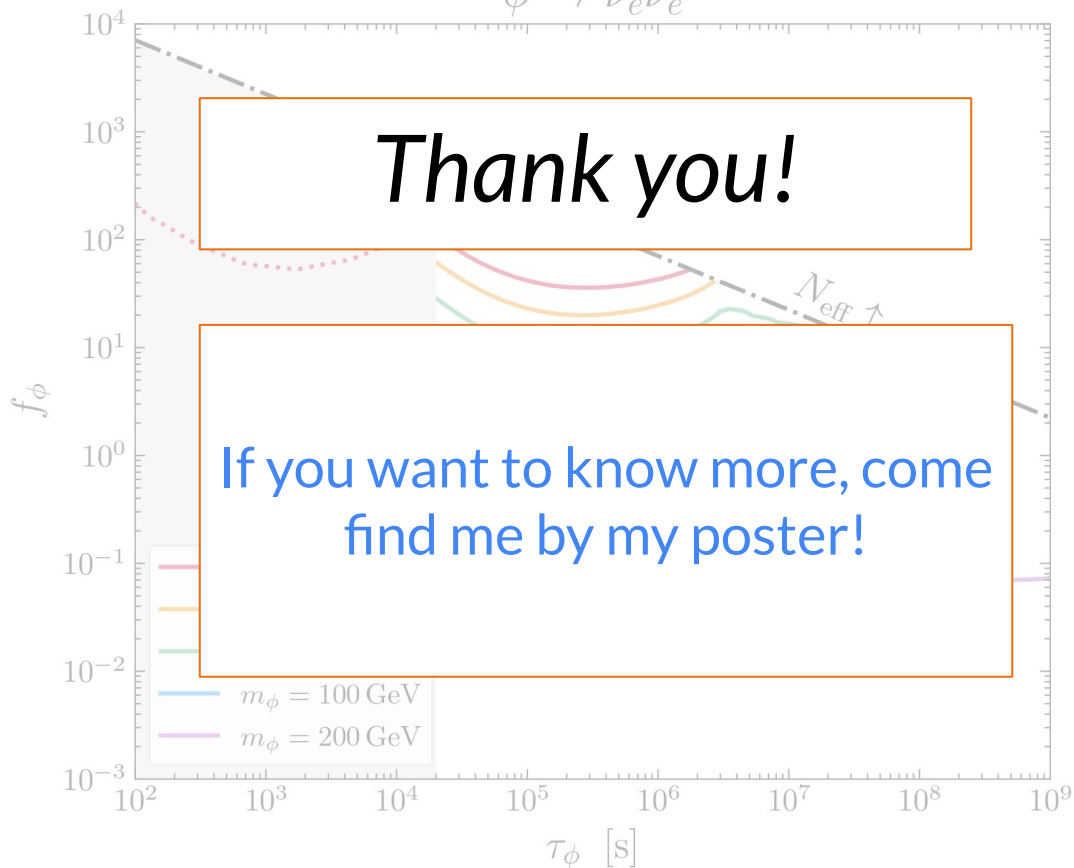
$$\nu\bar{\nu} \rightarrow \pi^-\pi^+$$

...





$$\phi \rightarrow \nu_e \bar{\nu}_e$$



Backup slides

Photodisintegration

Late time decays or residual annihilations of dark sector particles can alter BBN abundances at later times.

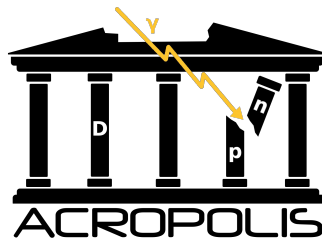
$$\dot{Y}_N(t) = \sum_j Y_j(t) \int_0^\infty dE f_\gamma(t, E) \sigma_{j\gamma \rightarrow N}(E) - Y_N(t) \sum_{j'} \int_0^\infty dE f_\gamma(t, E) \sigma_{N\gamma \rightarrow j'}(E)$$

$$E_{e^\pm}^{\text{th}} \simeq m_e^2 / (22T)$$

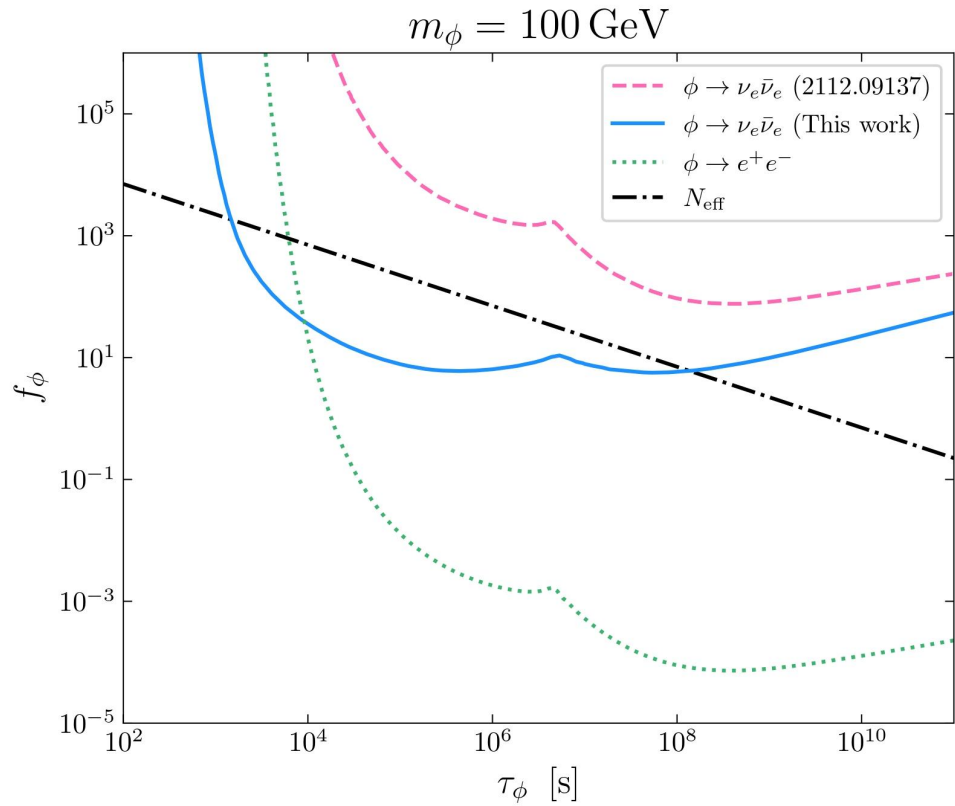
If above pair-production threshold, high-energy photons are rapidly depleted.

	T (keV)	E^{th} (MeV)
D	5.34	2.22
^3H	1.90	6.26
^3He	2.16	5.49
^4He	0.60	19.81
^6Li	3.21	3.70
^7Li	4.81	2.47
^7Be	7.48	1.59

BBN at these T has already finished.
Final abundances of BBN are initial abundances of photodisintegration.



Publicly available code
[2011.06518](#), P. F. Depta, M.
Hufnagel, and K.
Schmidt-Hoberg.



Neutrino oscillations

