

New non-SUSY tachyon-free heterotic theories in 6d and 4d

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C. Angelantonj, I. Florakis, G. L., D. Perugini, arXiv:2407.09597

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In this presentation: construction of heterotic vacua in 6d and 4d enjoying

- SUSY breaking
- Absence of tachyons
- Rigidity
- Rank reduction

SUSY Breaking

Low energy phenomenology requires SUSY breaking, many ways to do it:

[Dixon, Harvey, '86][Alvarez-Gaume, Ginsparg, Moore, Vafa,'86][Seiberg, Witten, '86][Bianchi, Sagnotti, '90]

- Vacua without SUSY (*i.e.* **heterotic**, type 0 (orientifolds),...)

[Bachas,'95][Berkooz,'96][Scherk, Schwarz, '79]

- Compactification (*i.e.* internal manifold, magnetic fields, Scherk-Schwarz,...)

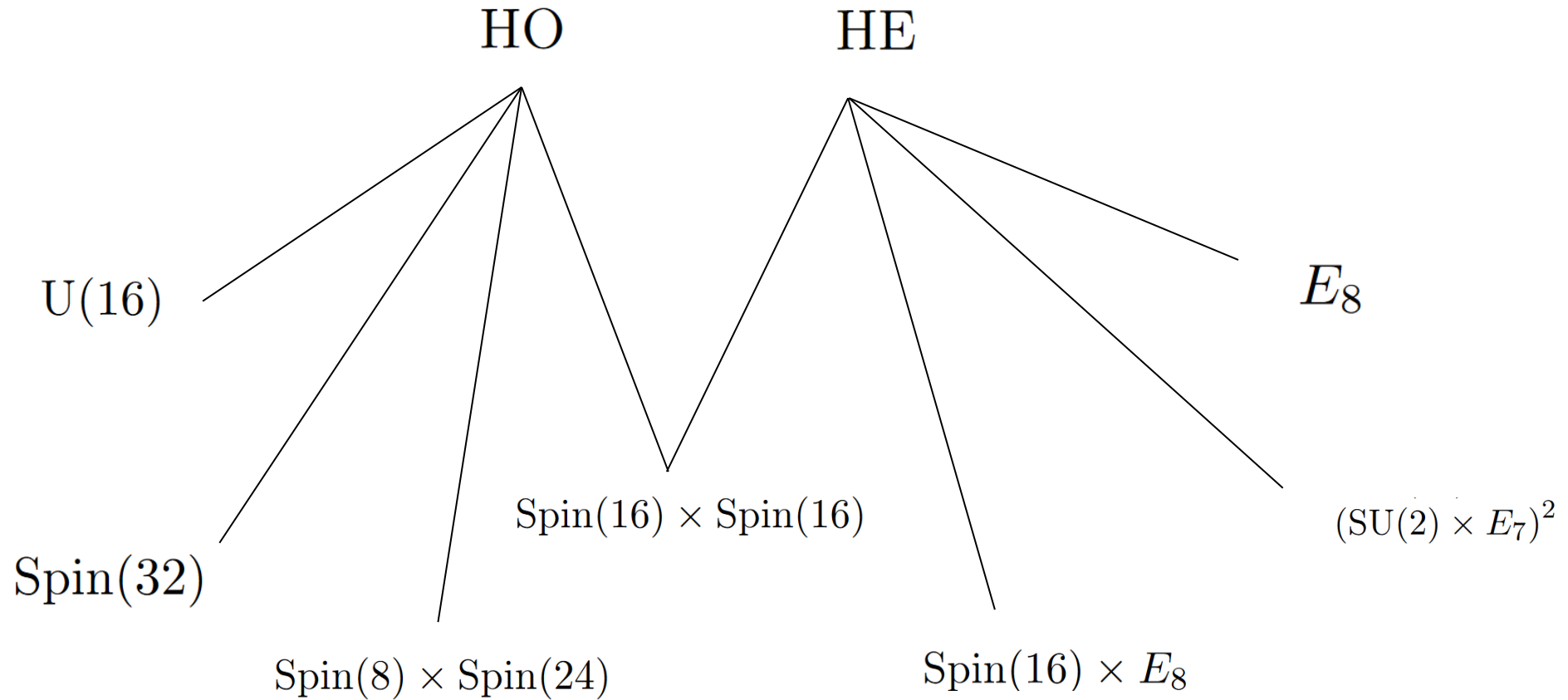
[Sugimoto, '99][Antoniadis, Dudas, Sagnotti, '99][Angelantonj, Condeescu, Dudas, G.L., '24]

- SUSY breaking at the string scale via type II orientifolds (*i.e.* Sugimoto, BSB, ...)





- ...

In 10 d, non-SUSY heterotic vacua

[Dixon, Harvey, '86]
[Alvarez-Gaume, Ginsparg, Moore, Vafa, '86]
[Kawai, Lewellen, Tye, '86]
[Lerche, Luest, Schellekens, '86]



SUSY breaking at the string scale induces instabilities

- Best case scenario  dilaton tadpoles induce shift of the background [Fischler, Susskind, '86]
 -  low-energy EFT analysis [Raucci's talk]
- Worst case scenario  tachyons in the tree level spectrum
 -  classically unstable

Only few models are tachyon-free, in 10d

↙
the $\text{Spin}(16) \times \text{Spin}(16)$
heterotic string

[Dixon, Harvey, '86]

[Alvarez-Gaume, Ginsparg, Moore, Vafa, '86]



GSO-projection breaks
SUSY and does not
admit tachyons

↘
the Sugimoto vacuum

[Sugimoto, '99]



type IIB orientifold
with $O9_+$ planes and
 $\overline{D9}$ branes

↘
the type $0'B$ superstring

[Sagnotti, '95]



type $0B$ orientifold
with closed-string
tachyon projected out

In lower dimensions, situation is different

Sugimoto and BSB constructions cannot be deformed to tachyonic vacua

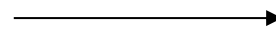
In heterotic vacua, Wilson + deformation moduli may induce tachyons

Ex. on $S^1(R)$

[Ginsparg, Vafa, '86]
[Fraiman, Grana, Parra De Freitas, Sethi, '23]

$\text{Spin}(16) \times \text{Spin}(16)$
no tachyons

$$W = (1, 0^7; (\frac{1}{3})^8), R^2 = \frac{1}{18}$$



$\text{Spin}(16) \times \text{Spin}(18)$

$t : (\mathbf{16}, \mathbf{1})$



tachyon



tree-level instability

Heterotic theories connected to tachyonic vacua

Truly tachyon-free heterotic theories?

One way out: constructions which do not admit deformation moduli



«non-geometric» orbifolds
(asymmetric, quasi-crystalline,...)

[Narain, Sarmadi, Vafa, '87]
[Baykara, Hamada, Tarazi, Vafa, '23]
[Baykara, Tarazi, Vafa, '24]

Rank reduction

Usually String Theory provides gauge groups with high rank

One phenomenological possibility: providing mechanism to reduce the rank

orientifold vacua: turning on a
Kalb-Ramond field

Ex: D9 branes in type I on T^d

$$G_{16} \rightarrow G_{24-b/2}$$

heterotic vacua: modding
out permutation symmetry

$$G_{16} \rightarrow G_8$$

[Bianchi, Pradisi, Sagnotti, '92]
[Angelantonj, '99]
[G.L., TAS]

[Kawai, Lewellen, Tye, '86]
[Chaudhuri, Hockney, Lykken, '95]
[Font, Fraiman, Grana, Nunez, Parra De Freitas, '21]

Appetizer: the E_8 vacuum in 10d

Simplest example: modding out permutation symmetry

[Kawai, Lewellen, Tye, '86]

$$E_8 \times E_8 \xrightarrow{S}$$



$A_\mu :$	248
$g_{\mu\nu}, B_{\mu\nu}, \phi :$	1
$\lambda_L + \lambda_R :$	248
$t :$	248

} E_8 at level 2

↘ tachyon

➔ SUSY broken: $(-1)^F S$



required for the action to be non-trivial

➔ SUSY preserved in lower dimensions: δS



CHL model

[Chaudhuri, Hockney, Lykken, '95]

6d vacuum

Idea: combine permutation symmetry with asymmetric action

Heterotic string $O(16) \times O(16)$ partition function compactified on the $SO(8)$ lattice

\mathbb{Z}_2 asymmetric orbifold action

$$(\mathrm{SO}(4) \times \mathrm{SO}(4) \times \mathrm{SO}(4) \times \mathrm{SO}(4))_L \times (\mathrm{O}(16) \times \mathrm{O}(16) \times \mathrm{SO}(8))_R$$

$$\begin{array}{ccc} \searrow & & \swarrow \\ g : X_L^i(\tau + \sigma) \rightarrow -X_L^i(\tau + \sigma) & & S : \mathrm{O}(16) \times \mathrm{O}(16)' \rightarrow \mathrm{O}(16)' \times \mathrm{O}(16) \\ \chi^i(\tau + \sigma) \rightarrow -\chi^i(\tau + \sigma) & & \end{array}$$

Completion by modular invariance

Action on world-sheet fields Fourier modes in NS sector (untwisted)

Gauge fields

states giving adjoint of SO(8) (symmetry enhancement)

$$b_{-1/2}^{\mu} \otimes [(\lambda_{-1/2}^{A,+} \lambda_{-1/2}^{B,+} \oplus \lambda_{-1/2}^{A,-} \lambda_{-1/2}^{B,-}) |0\rangle \oplus |\psi\rangle] \longrightarrow b_{-1/2}^{\mu} \otimes [(\lambda_{-1/2}^{A,+} \lambda_{-1/2}^{B,+} \oplus (-\lambda_{-1/2}^{A,-} \lambda_{-1/2}^{B,-})) |0\rangle \oplus |\psi\rangle]$$



$$b_{-1/2}^{\mu} \otimes (\lambda_{-1/2}^{A,+} \lambda_{-1/2}^{B,+} |0\rangle \oplus |\psi\rangle)$$

Scalars

$$b_{-1/2}^i \otimes [(\lambda_{-1/2}^{A,+} \lambda_{-1/2}^{B,+} \oplus \lambda_{-1/2}^{A,-} \lambda_{-1/2}^{B,-}) |0\rangle \oplus |\psi\rangle] \longrightarrow -b_{-1/2}^i \otimes [(\lambda_{-1/2}^{A,+} \lambda_{-1/2}^{B,+} \oplus (-\lambda_{-1/2}^{A,-} \lambda_{-1/2}^{B,-})) |0\rangle \oplus |\psi\rangle]$$



$$b_{-1/2}^i \otimes \lambda_{-1/2}^{A,-} \lambda_{-1/2}^{B,-} |0\rangle$$

Light spectrum: gauge group $O(16)_2 \times SO(8)_1$

$$\begin{aligned} A_\mu &: && (\mathbf{120}, \mathbf{1}) + (\mathbf{1}, \mathbf{66}) \\ g_{\mu\nu}, B_{\mu\nu}, \phi &: && (\mathbf{1}, \mathbf{1}) \\ 4\phi &: && (\mathbf{120}, \mathbf{1}) \\ \lambda_L &: && (\mathbf{128}, \mathbf{1}) + (\mathbf{136}, \mathbf{1}) \\ \lambda_R &: && (\mathbf{128}, \mathbf{1}) + (\mathbf{120}, \mathbf{1}) \\ 4\phi + \lambda_R &: && (\mathbf{120}, \mathbf{1}) + (\mathbf{1}, \mathbf{8}_v + \mathbf{8}_s + \mathbf{8}_c) \\ 4\phi + \lambda_L &: && (\mathbf{128}, \mathbf{1}) \end{aligned}$$

Kaehler and complex structure moduli projected away

➡ no deformations allowed*

Similar analysis for the tachyonic $(SU(2) \times E_7)_2 \times SO(8)_1$ heterotic theory

4d vacuum

Heterotic string $O(16) \times O(16)$ compactified on the $SO(12)$ lattice

Straightforward extension: $\mathbb{Z}_2 \times \mathbb{Z}_2$ orbifold action

$$\begin{array}{l|l} g \cdot S & |g : (-, +, -)_L \\ h \cdot S & |h : (+, -, -)_L \end{array}$$

Problem: on last T^2 action is trivialized in the $gh \cdot \mathbf{1}$ twisted sector



deformation moduli reappear in twisted sector

One needs to include asymmetric shifts $\delta : X_R^i \rightarrow X_R^i + \pi$

$$\begin{array}{l|l} g \cdot S & |g : (-, +, -)_L \times (1, \delta, \delta)_R \\ h \cdot S & |h : (+, -, -)_L \times (\delta, 1, \delta)_R \end{array}$$

Light spectrum: gauge group $O(16)_2 \times SO(4)_1 \times SO(4)_1 \times SO(4)_1$

$$\begin{aligned}
 A_\mu &: && (\mathbf{120}, \mathbf{1}, \mathbf{1}, \mathbf{1}) + (\mathbf{1}, \mathbf{6}, \mathbf{1}, \mathbf{1}) + (\mathbf{1}, \mathbf{1}, \mathbf{6}, \mathbf{1}) + (\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{6}) \\
 g_{\mu\nu}, B_{\mu\nu}, \phi &: && (\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{1}) \\
 2\phi &: && (\mathbf{120}, \mathbf{1}, \mathbf{1}, \mathbf{1}) + (\mathbf{1}, \mathbf{4}, \mathbf{4}, \mathbf{1}) + (\mathbf{1}, \mathbf{4}, \mathbf{1}, \mathbf{4}) + (\mathbf{1}, \mathbf{1}, \mathbf{4}, \mathbf{4}) \\
 \psi_D &: && 2(\mathbf{128}, \mathbf{1}, \mathbf{1}, \mathbf{1}) + (\mathbf{136}, \mathbf{1}, \mathbf{1}, \mathbf{1}) + (\mathbf{120}, \mathbf{1}, \mathbf{1}, \mathbf{1}) \\
 2\phi + \psi_D &: && (\mathbf{1}, \mathbf{1}, \mathbf{2}_s + \mathbf{2}_c, \mathbf{2}_s + \mathbf{2}_c) \\
 2\phi + \psi_D &: && (\mathbf{1}, \mathbf{2}_s + \mathbf{2}_c, \mathbf{1}, \mathbf{2}_s + \mathbf{2}_c) \\
 2\phi &: && (\mathbf{1}, \mathbf{2}_s + \mathbf{2}_c, \mathbf{2}_s + \mathbf{2}_c, \mathbf{4})
 \end{aligned}$$

Kaehler and complex structure moduli projected away

➡ no deformations allowed*

Other two similar constructions available

Conclusions and Outlook

Non-SUSY heterotic vacua in 6d and 4d with:

- No tachyons
- No deformation moduli
- Reduced rank

Other similar constructions?

Role of scalars?

THANK YOU FOR THE ATTENTION