

Laser-induced Nuclear Astrophysics @ LNS

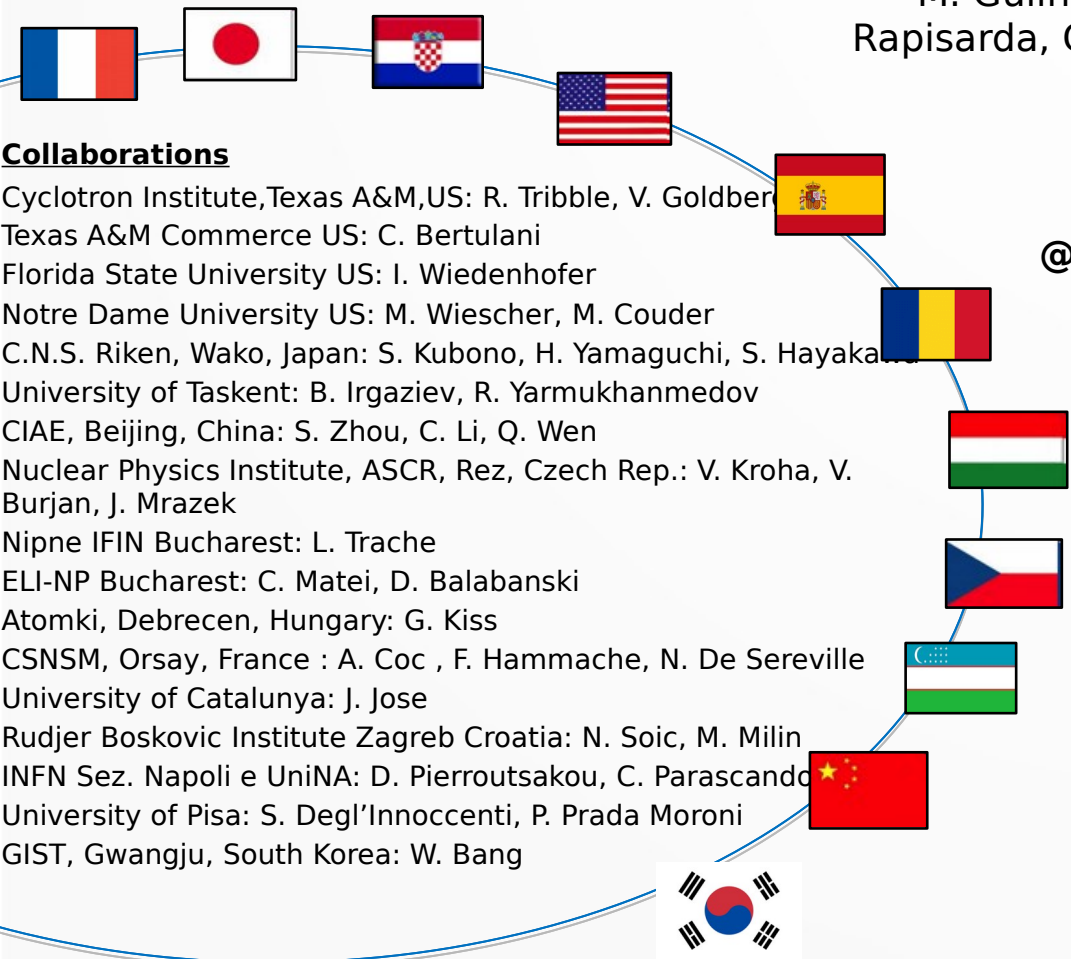
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The AsFiN Collaboration

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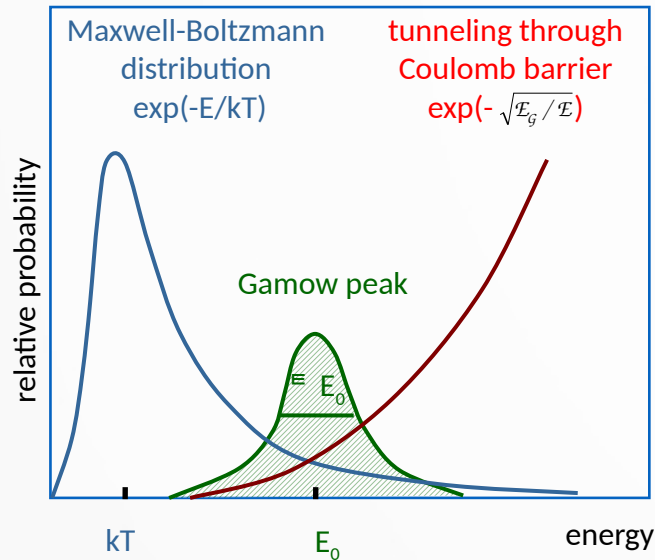
@**Perugia** M. Busso, S. Palmerini, M. Limongi, A. Chieffi, M.C. Nucci



Outline

- Nuclear Astrophysics
- Why lasers?
- How?
- Why at LNS?
- What next?

Nuclear reactions in stars..



Example: $T \sim 15 \times 10^6 \text{ K}$ ($T_6 = 15$)

reaction	C.barrier (MeV)	E_0 (keV)	area under Gamow peak
$p + p$	0.5	5.9	7.0×10^{-6}
$^{12}\text{C} + ^{12}\text{C}$	2.242	56	5.9×10^{-56}
$^{16}\text{O} + ^{16}\text{O}$	10.349	237	2.5×10^{-237}

For $T \sim 200 \times 10^6 \text{ K}$, $E_0 \sim 320 \text{ keV}$
($kT \sim 17 \text{ keV}$)

Gamow peak: most effective energy region for thermonuclear reactions

It is where measurements should be carried out

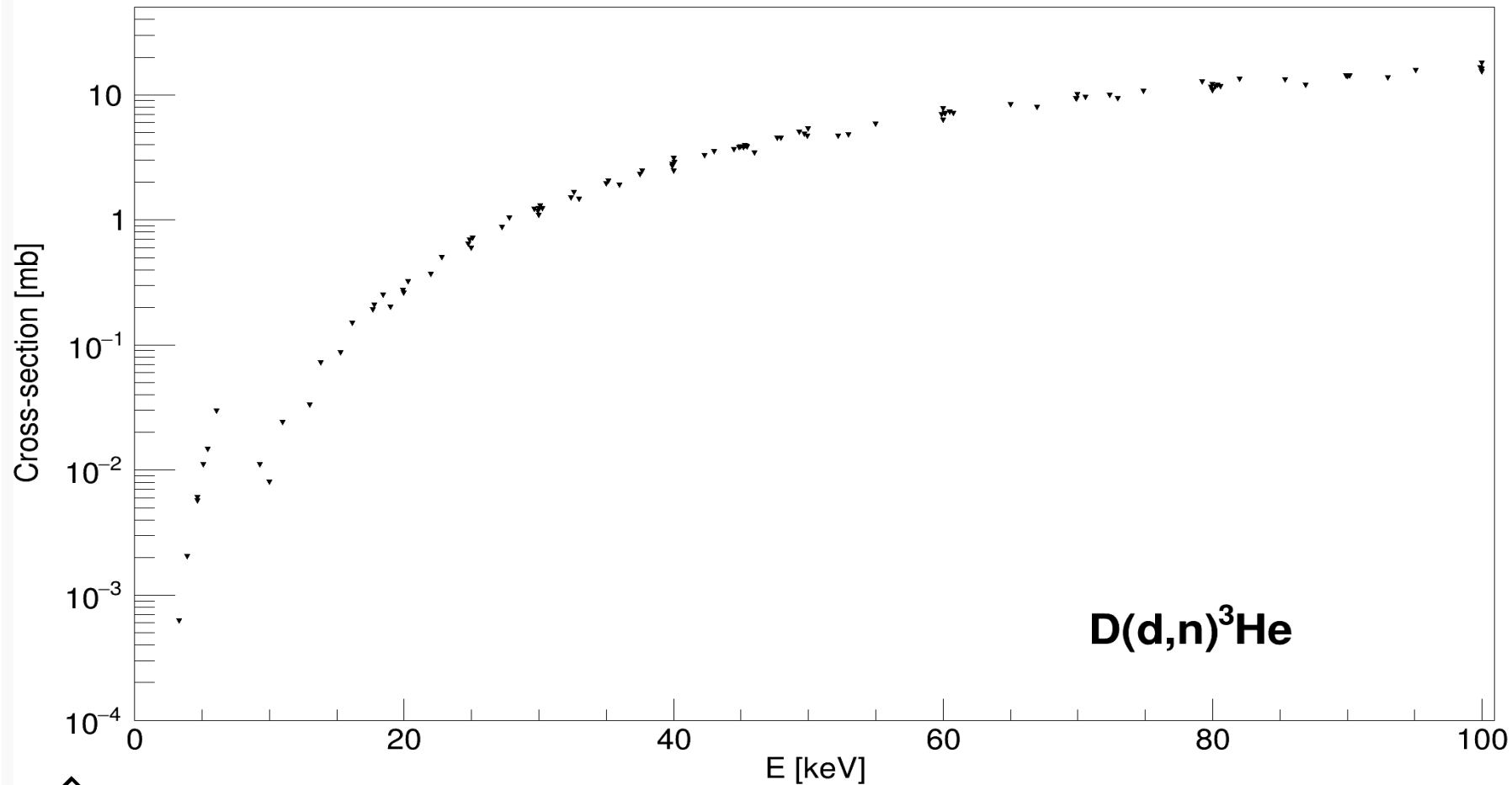
$10^{-18} \text{ barn} < \sigma < 10^{-9} \text{ barn}$



- EXTRAPOLATION
- LUNA, JUNA (background suppression)
- INDIRECT METHOD (THM, ANC)

.. explore directly in a laboratory plasma!

0-day exp: Deuterium burning



GIST23 plan

Temp [keV]	Livetime [h]	Shots
1-2	8	~ 500
2-3	4	~ 150
3-5	1	~ 100
5-10	1	~ 100
> 10	2	~ 200

$D(d,p)t$ also! BR: 50/50

Electron screening

Big Bang Nucleosynthesis

Electron screening

$$\sigma_s(E) = \sigma_b(E+U)$$

$$f(E) = \frac{\langle \sigma_s v \rangle}{\langle \sigma_b v \rangle} \approx e^{\frac{\pi \eta U}{E}}$$

$$U \approx \frac{Z_1 Z_2 e^2}{R_o}$$

$$\langle \sigma v \rangle_{\text{accelerator}} \neq \langle \sigma v \rangle_{\text{plasma}} \approx \langle \sigma v \rangle_{\text{stars}} \neq \langle \sigma v \rangle_{\text{bare}}$$

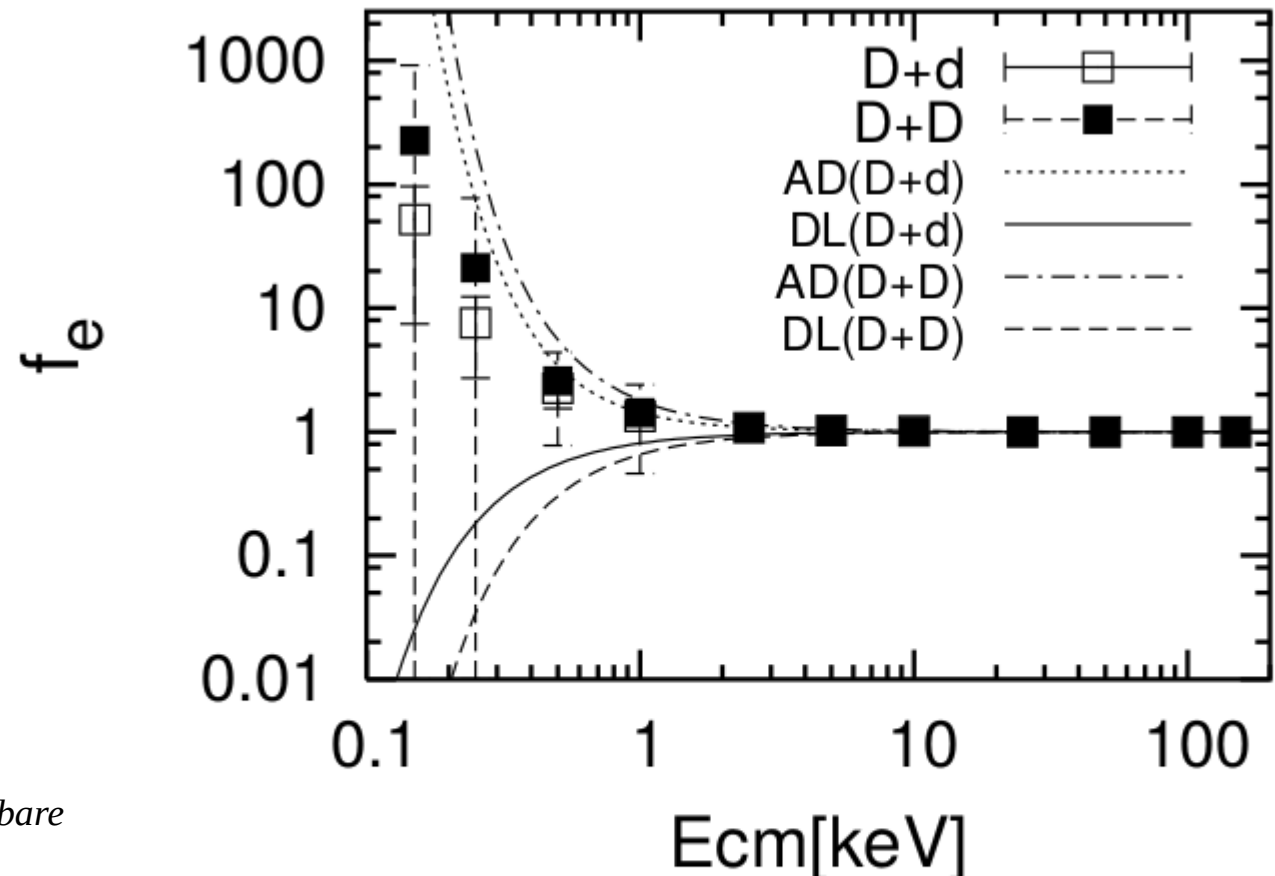


Fig. 1. Enhancement factor as a function of incident center-of-mass energy for the reactions D+d and D+D. Error-bars represent the variances obtained from the events generated for each beam energy.

Electron screening

Relatively small enhancements due to electron screening at energies $E/U = 100$ could cause significant errors in the extrapolation to lower energies, if the cross section curve is forced to follow the trend of the enhanced cross sections without correcting for screening

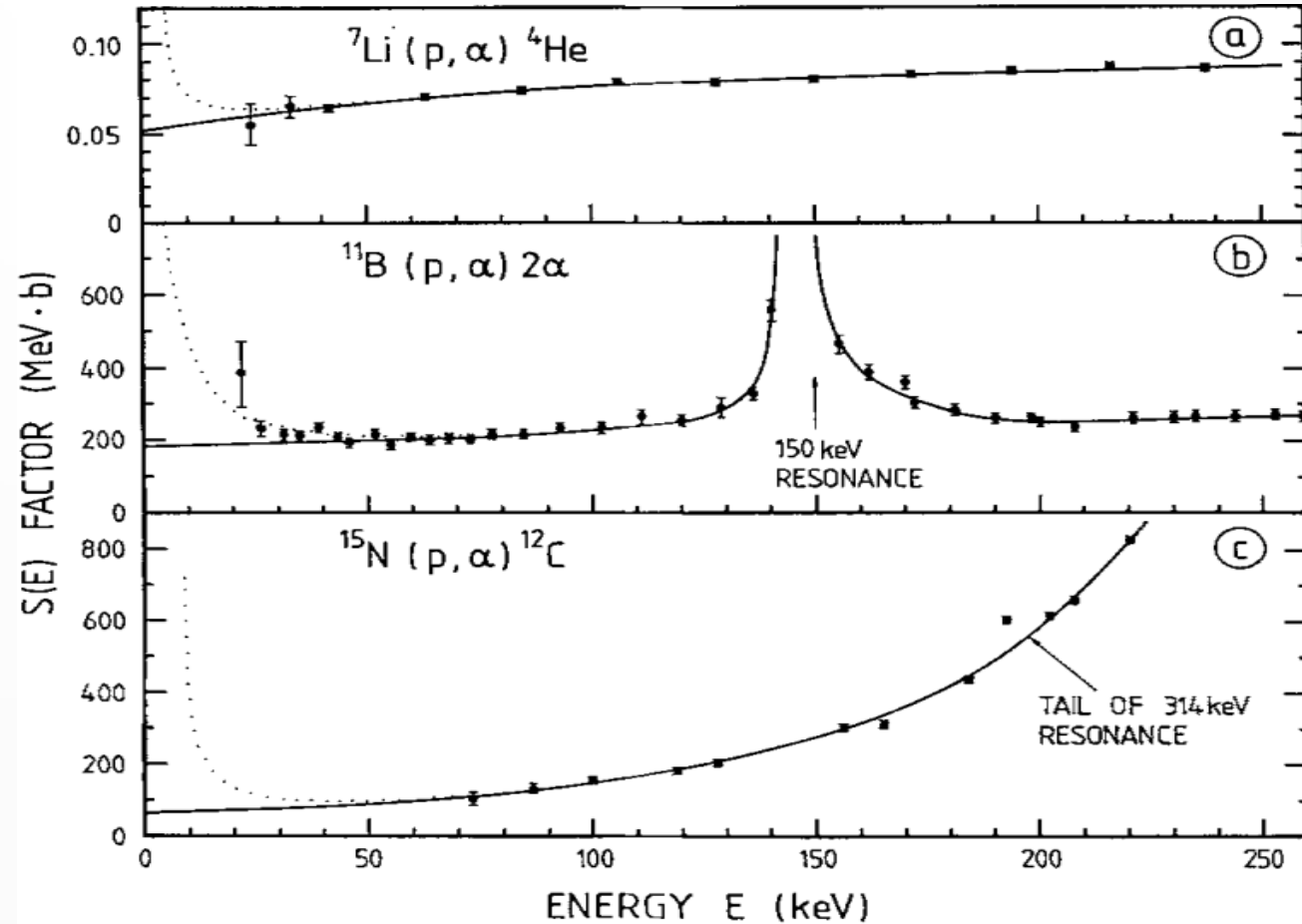
The whole effect of screening in this case ($U_0 \ll E_{\max}$) is then that the reaction rate with screening neglected has to be multiplied by the factor $e^{-\frac{U_0}{KT}}$

$$R_D = 2.812 \times 10^{-7} \rho^{-1/2} T_9^{1/2} \zeta^{-1} (\text{cm})$$

Weak screening: $R_D \gg \text{dist}_{\text{nuclei}} (\text{stars})$

Intermediate screening: $\langle E_C \rangle \approx KT$

Strong screening: $\langle E_C \rangle \gg KT$



Effects of electron screening on low-energy fusion cross sections Assenbaum, H. J., Langanke, K., & Rolfs, C. 1987

ELECTRON SCREENING AND THERMONUCLEAR REACTIONS E. E. SALPETER 1954

How? Coulomb explosion of cryo-cooled molecular clusters

Step 1

Clusters are irradiated by high intensity laser pulse ($\sim 10^{16} \sim 10^{18}$ W/cm²).

Step 2

Laser pulse energy is first absorbed by electrons via heating mechanisms such as rapid collisional heating.

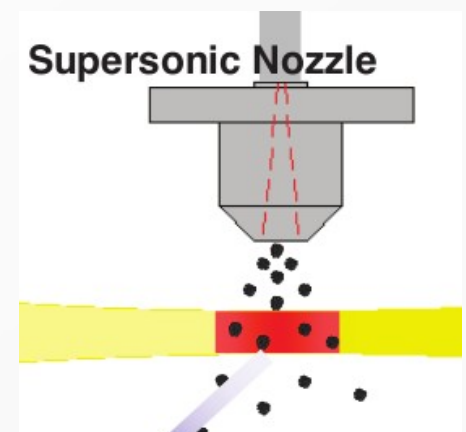
Step 3

Electrons escape from the cluster and leave positive charge build-up on the cluster.

Step 4

The cluster “explodes” and deuterons acquire multi-keV kinetic energy.

- Enhanced laser absorption
- Thermal distribution of accelerated species
- (no ponderomotive force)
- 4π isotropic
- Control over the MB temperature by changing the focus and/or the temperature of the nozzle

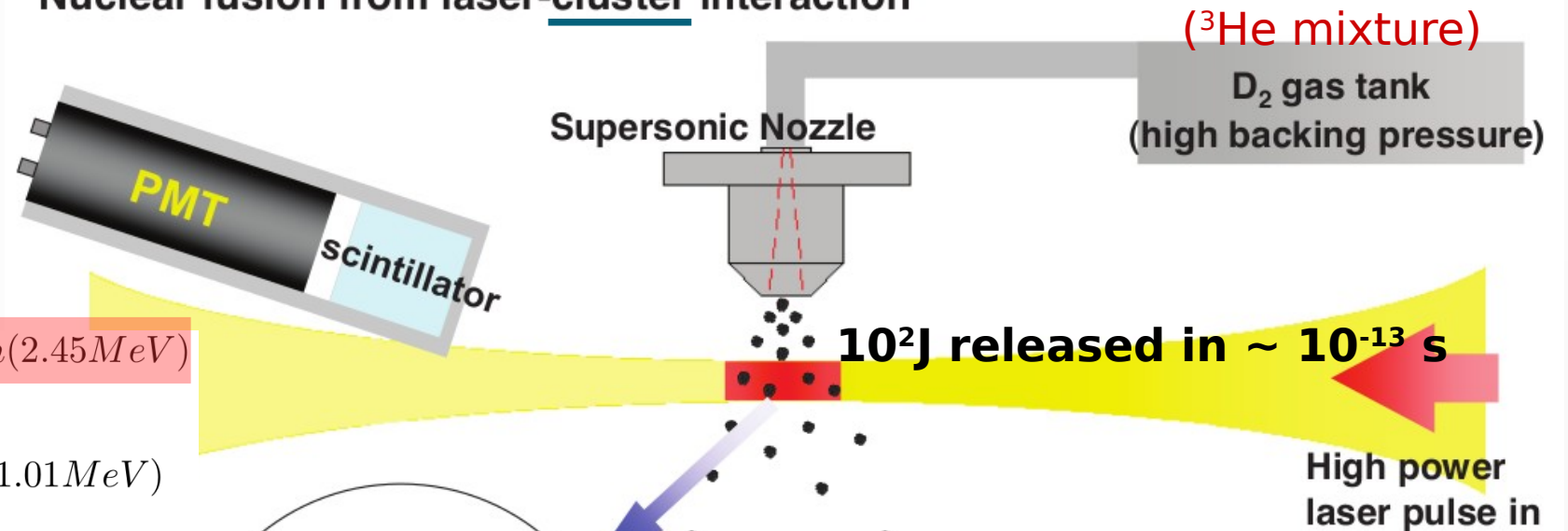


How?

$E = 100-140 \text{ J}$
 Pulse duration: 140 fs
 Rep. Rate: $\sim 1/\text{hour}$
 CW: 1057 nm
 Intensity $\sim 10^{21} \text{ W/cm}^2$

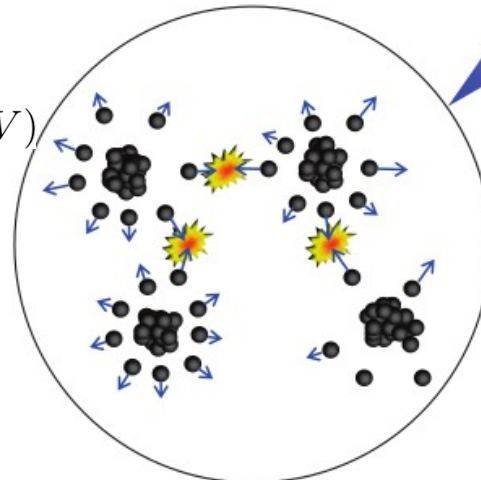
UT Petawatt Laser 2011

Nuclear fusion from laser-cluster interaction



10^2 J released in $\sim 10^{-13} \text{ s}$

High power laser pulse in



- Most of the laser pulse energy is absorbed by the atomic clusters.
- Clusters experience Coulomb explosion after electrons escape.
- DD fusion occurs, and 2.45MeV fusion neutrons are produced.

(and D^3He)

Deuteron Kinetic Energy $< 10^2 \text{ keV}$
 Density $\sim 10^{18} \text{ atoms/cm}^3$

10^5-10^7 neutrons per shot

Ditmire, T. et al. Nuclear fusion from explosions of femtosecond laser-heated deuterium clusters. Nature 398, 489-492 (1999).

How?

D(d,n)³He

D(d,p)t

$$Y = \frac{\rho_1 \int \frac{dN}{dE} S(E) \exp[-2\pi\eta(E)] v \tau dE}{1 + \delta_{12}}$$

$$= l \rho_D \int \frac{dN}{dE} \sigma_{BB}(E) dE$$

$$\sigma(E) = \frac{S(E)}{E} e^{-\sqrt{E_G/E}}$$

$$2\pi\eta = b/\sqrt{E}$$

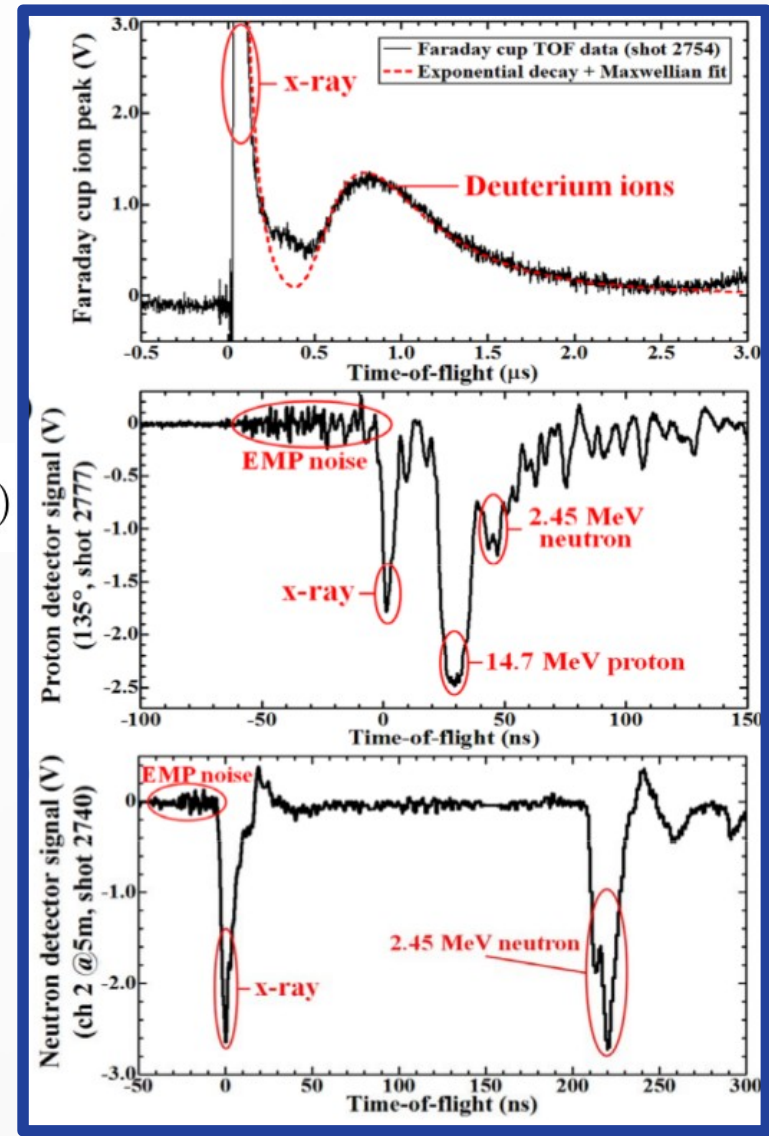
$$b = 0.9898 Z_i Z_j \sqrt{A} \text{ MeV}^{1/2}$$

$$\sigma(E) = \frac{S(E)}{E} \exp(-2\pi\eta(E))$$

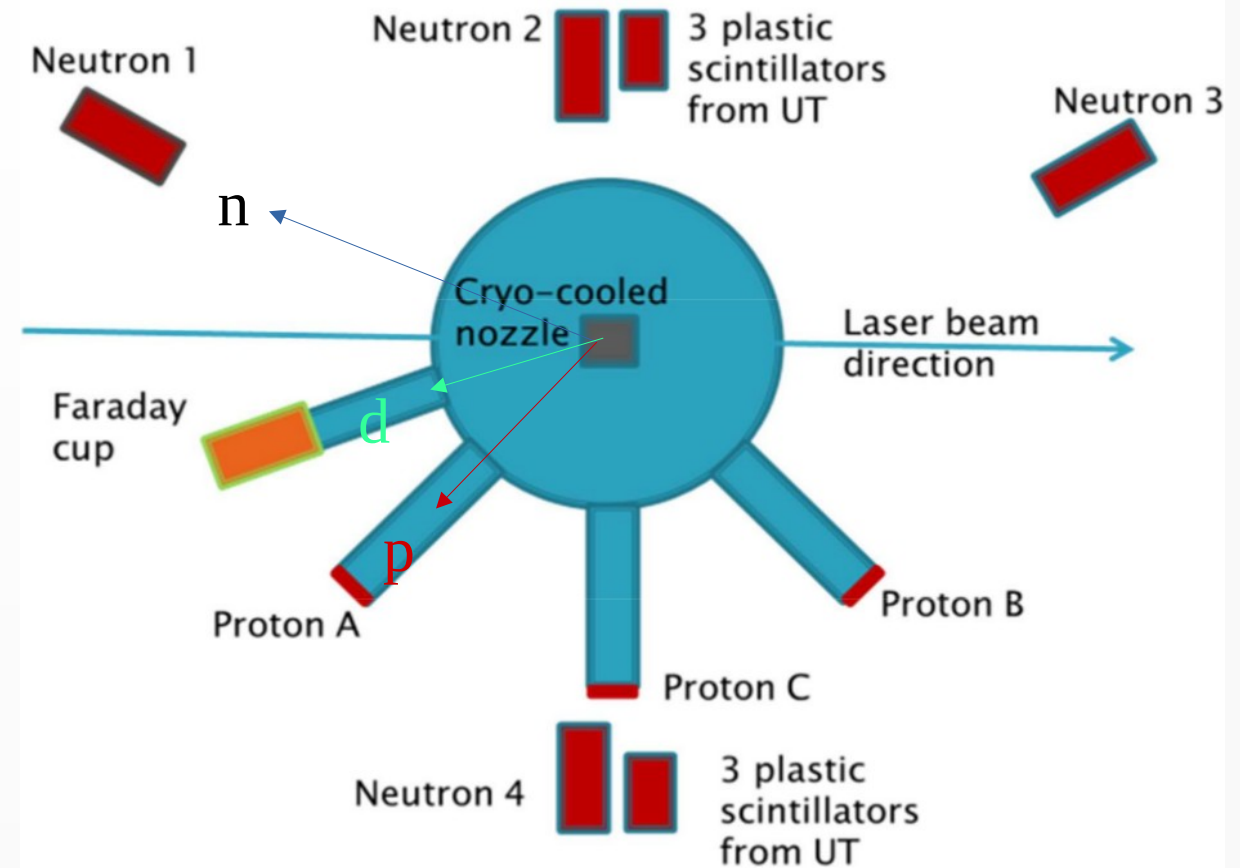
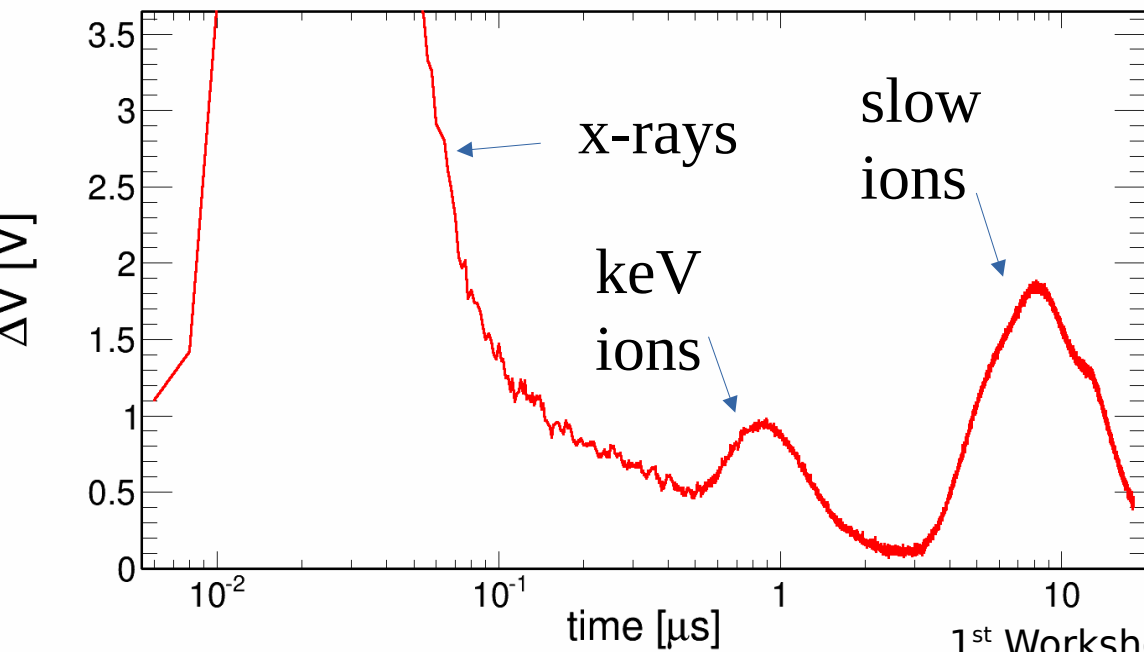
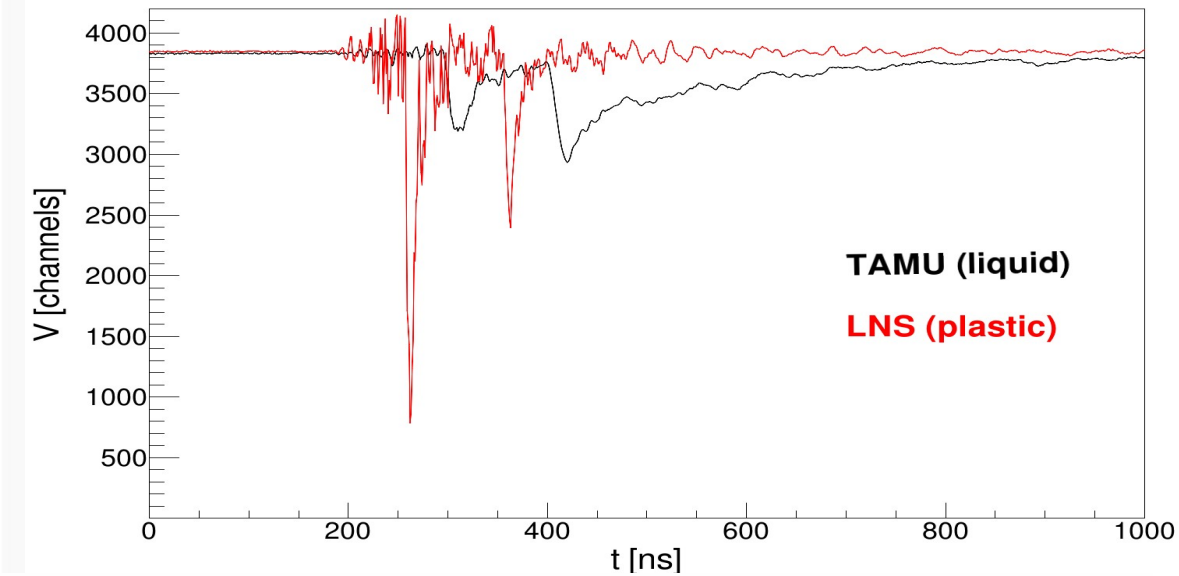
n & p → scintillators, diamond, CR39

Deuteron spectrum → MCP, FC, diamond, CR39

Gas and plasma diagnostics → interferometry, CCD, previous Rayleigh scattering measurements



2011 & 2016 experiments



Model-independent determination of the astrophysical S factor in laser-induced fusion plasmas

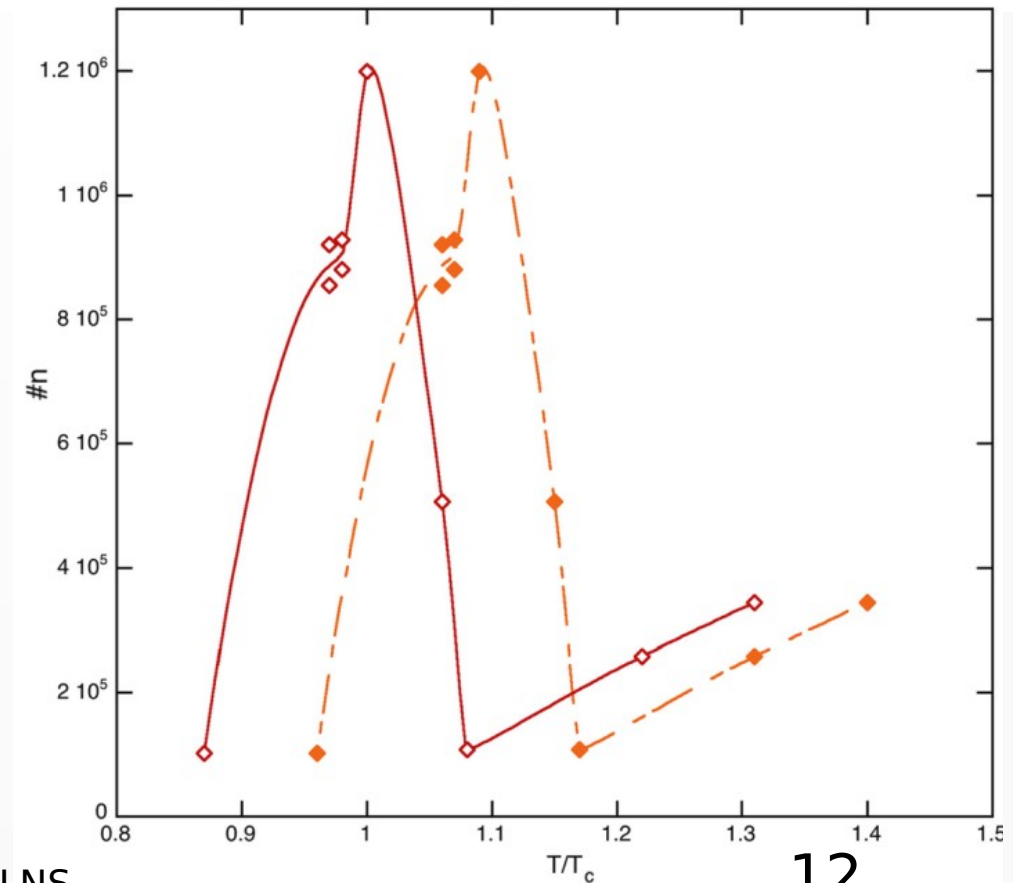
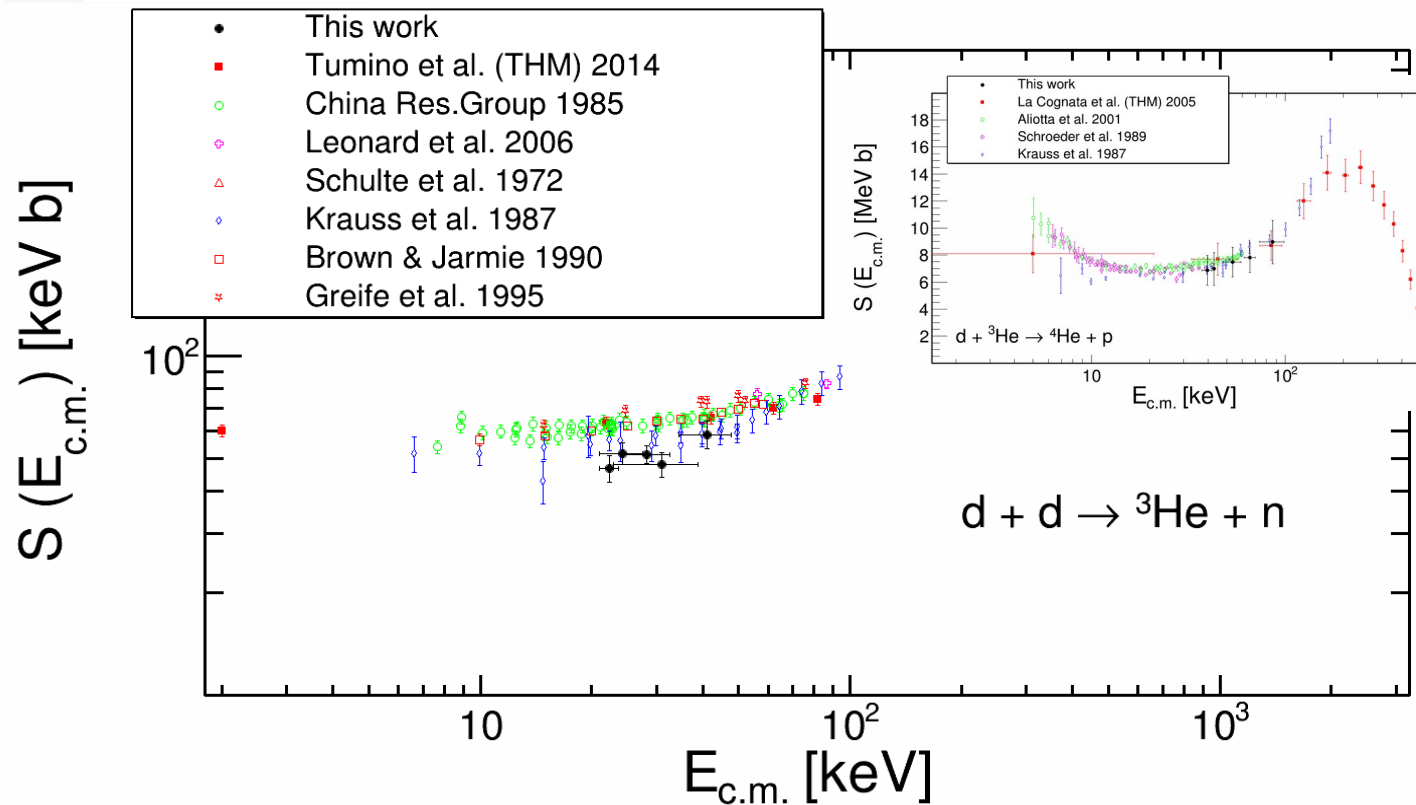
D. Lattuada, M. Barbarino, A. Bonasera, W. Bang, H. J. Quevedo, M. Warren, F. Consoli, R. De Angelis, P. Andreoli, S. Kimura, G. Dyer, A. C. Bernstein, K. Hagel, M. Barbui, K. Schmidt, E. Gaul, M. E. Donovan, J. B. Natowitz, and T. Ditmire
 Phys. Rev. C **93**, 045808 – Published 19 April 2016



Neutron enhancement from laser interaction with a critical fluid



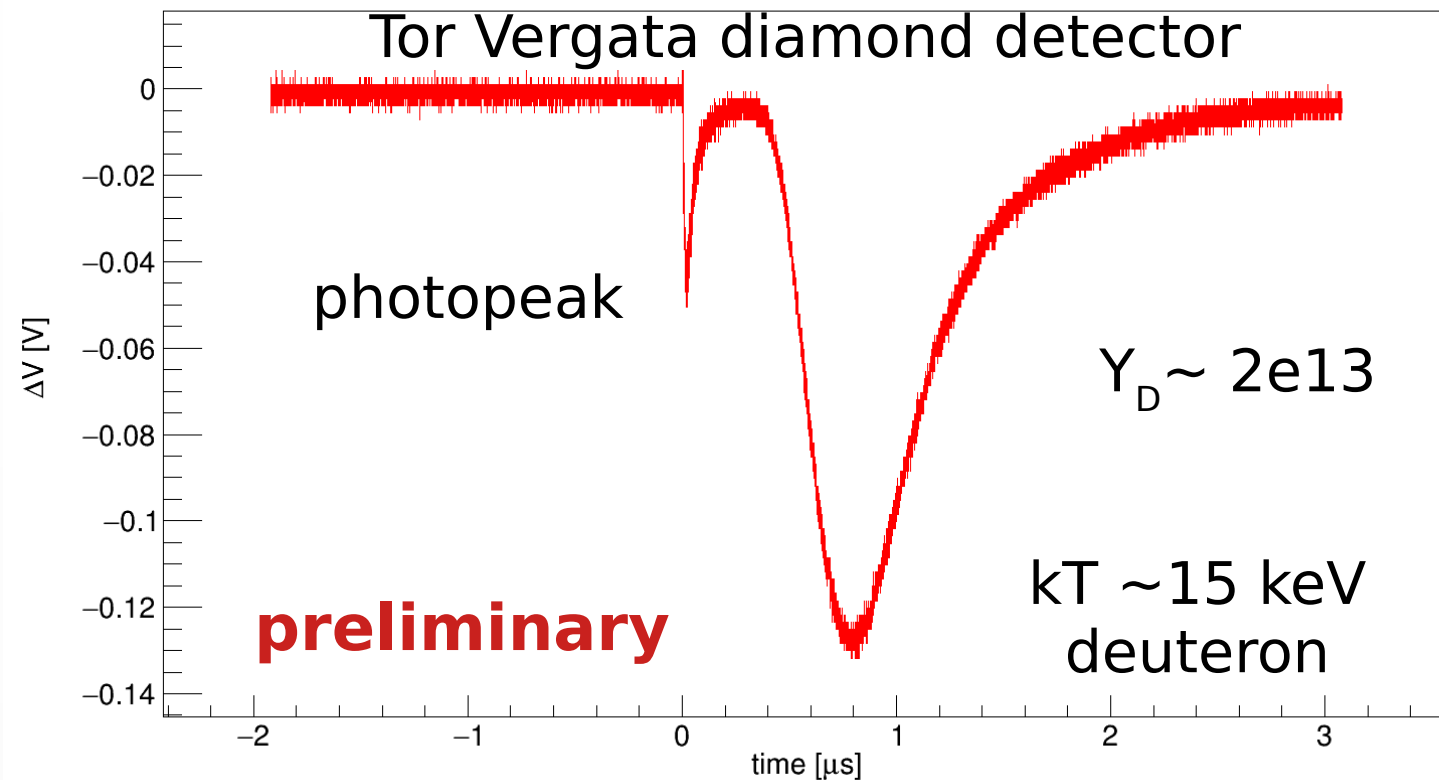
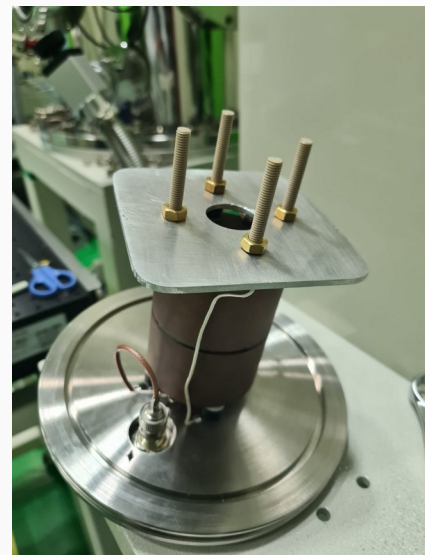
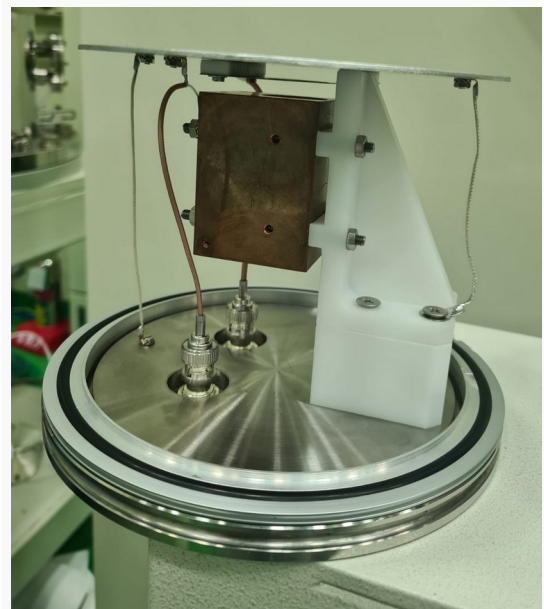
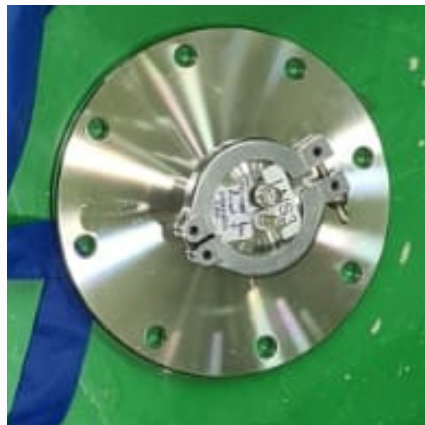
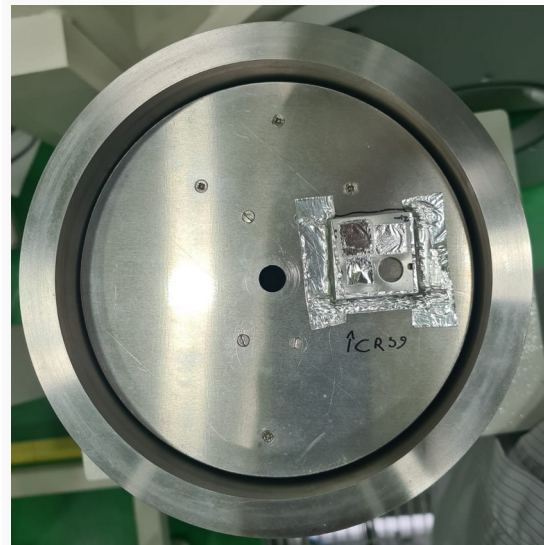
H.J. Quevedo^a, G. Zhang^{b,1}, A. Bonasera^{b,c,*}, M. Donovan^a, G. Dyer^a, E. Gaul^a, G.L. Guardo^c, M. Gulino^{c,d}, M. La Cognata^c, D. Lattuada^c, S. Palmerini^{e,f}, R.G. Pizzone^c, S. Romano^c, H. Smith^a, O. Trippella^{e,f}, A. Anzalone^c, C. Spitaleri^c, T. Ditmire^a



Where

- iLUCE perfect candidate for future high repetition studies
- Long-standing expertise on nuclear physics, accelerators, detectors, electronics, cryogenics, target development, etc.. @LNS (and @LNF)
- Optimization of the setup while under construction (tests @ FLAME)
- Unique possibilities of new paradigm of interaction in synergy with standard accelerators

The Dec23 GIST experiment

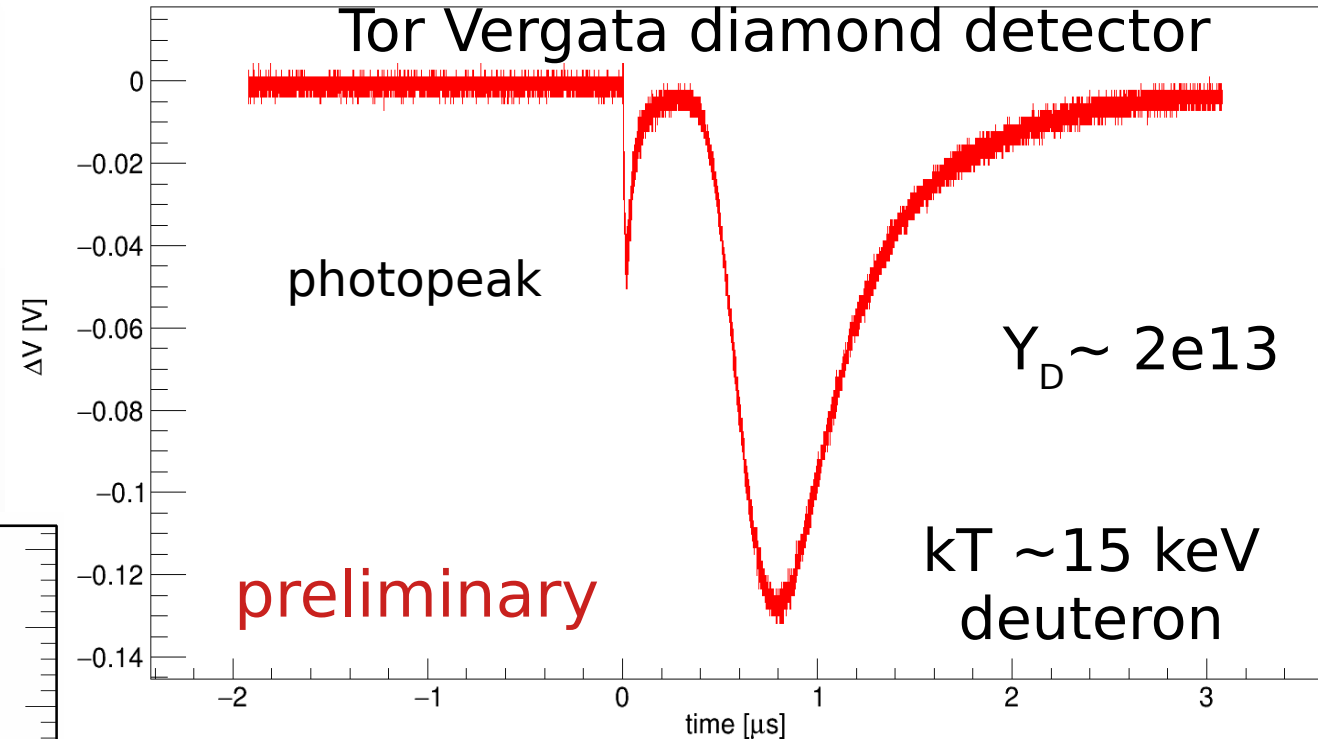
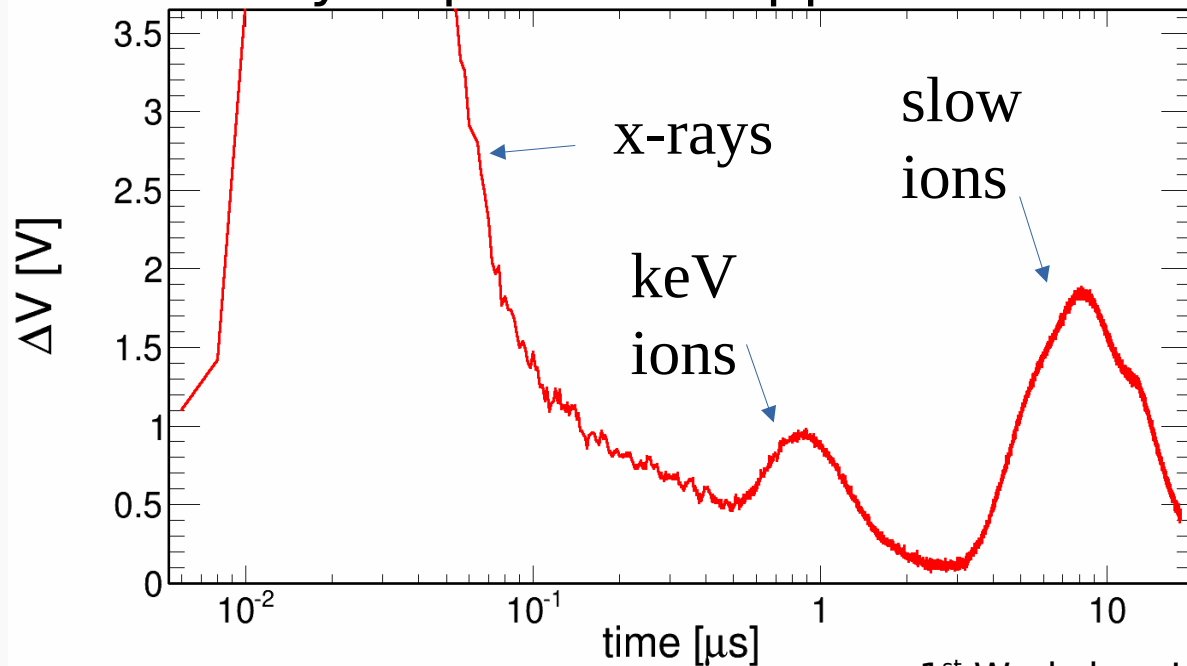


**Over 3000 shots, with T control between
 $kT \sim 1$ keV and $kT \sim 40$ keV**

Experimental Faraday cup with
magnetic suppression (LNS)

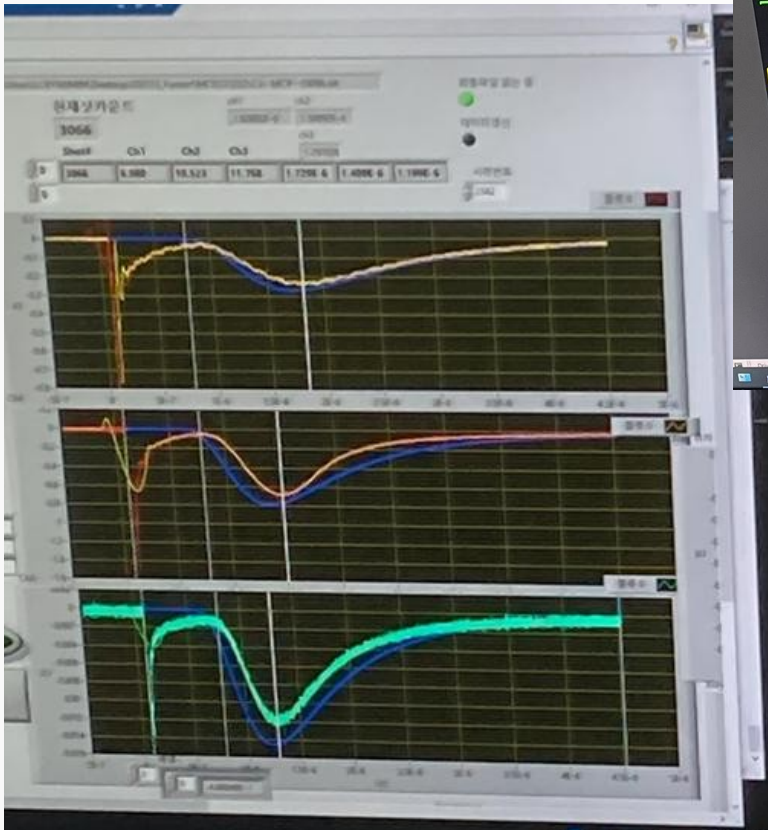
GIST23 & UT2011/16

UT Faraday Cup with V suppression



GIST

The Dec23 GIST experiment



What next?

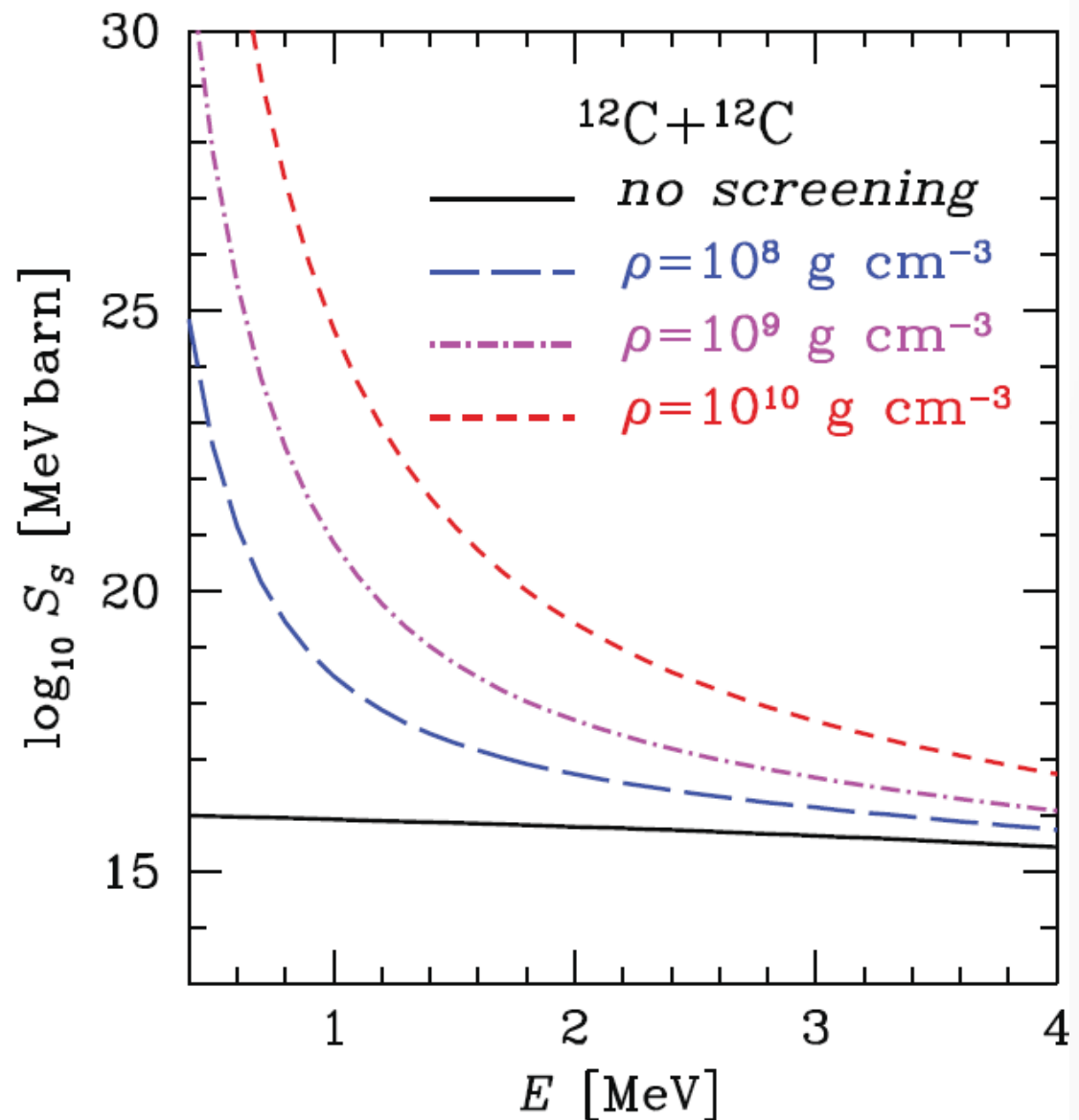
What next?

Can we accelerate ^{12}C ?

Up to which energies?

We can detect them in ToF with FC and diamond detectors, but can we deconvolute the signals?

We would need a Thomson Parabola Spectrometer or..



What else?

- LASER + TANDEM **unique** possibility at LNS
- In-plasma measurement of nuclear reactions relevant for astrophysics with a tandem-like **beam profile** impinging on a laser-induced **plasma** target, with high repetition
- **SETBACKS:** synchronization and reproducibility?
Exploring a different paradigm: continuous beam + 1 laser shot per second/minute and compare → background evaluation can be tricky
- Need for intense R&D and careful planning

Conclusions

- High power lasers technology is steeply advancing
- Huge window of opportunity for Nuclear Astrophysics in-plasma studies
- Laser-cluster experiments provide a perfect paradigm
- In-plasma cross-section measurements have been performed
- In-plasma electron screening studies are possible (synergy with others)
- Need for systematic and multi-diagnostics approach
-
- ..so? Money and people.

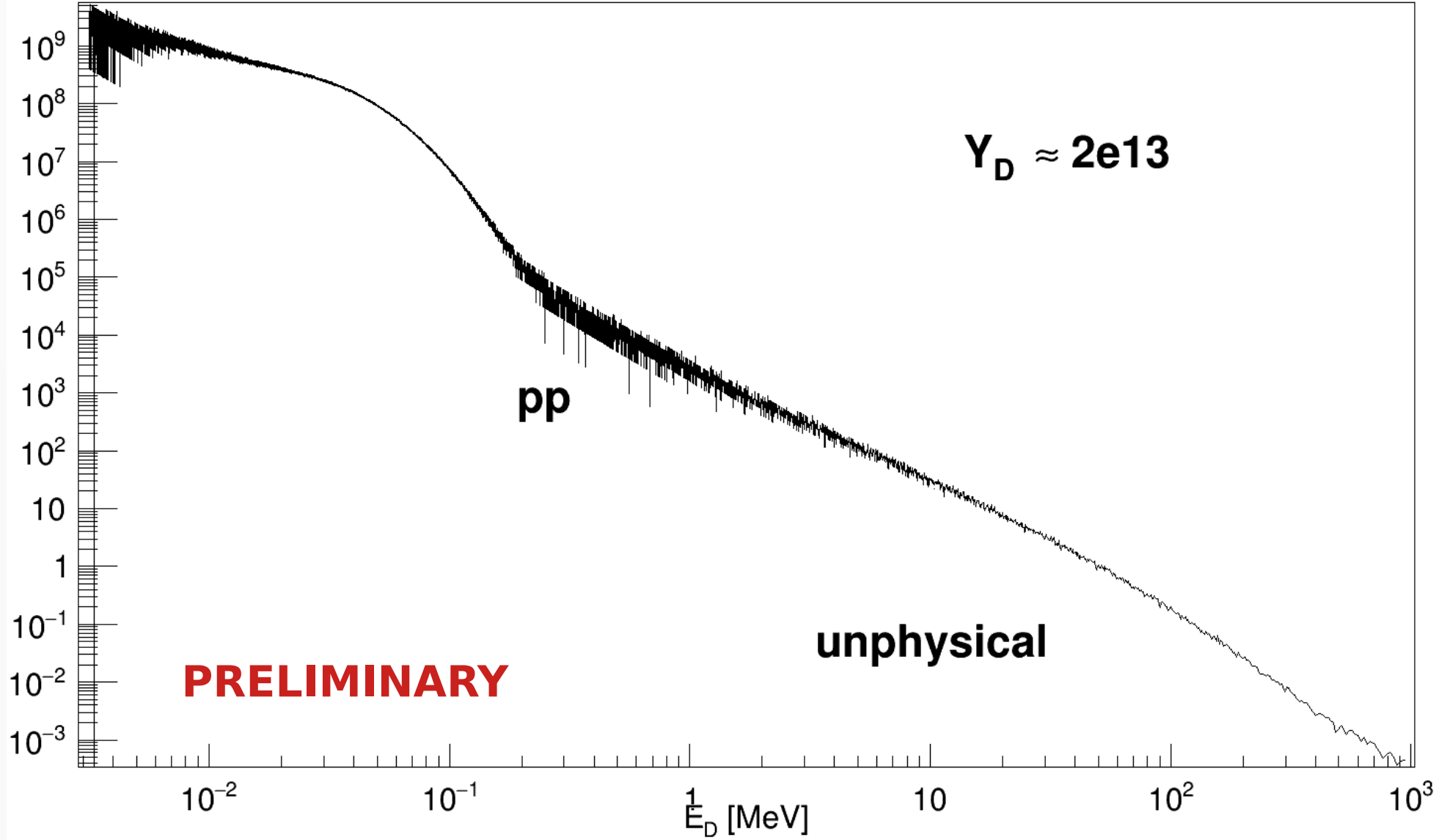


Q&A



BACKUP

GIST 00009



$Y_D \approx 2e13$

pp

unphysical

PRELIMINARY

10^{-2}

10^{-1}

E_D [MeV]

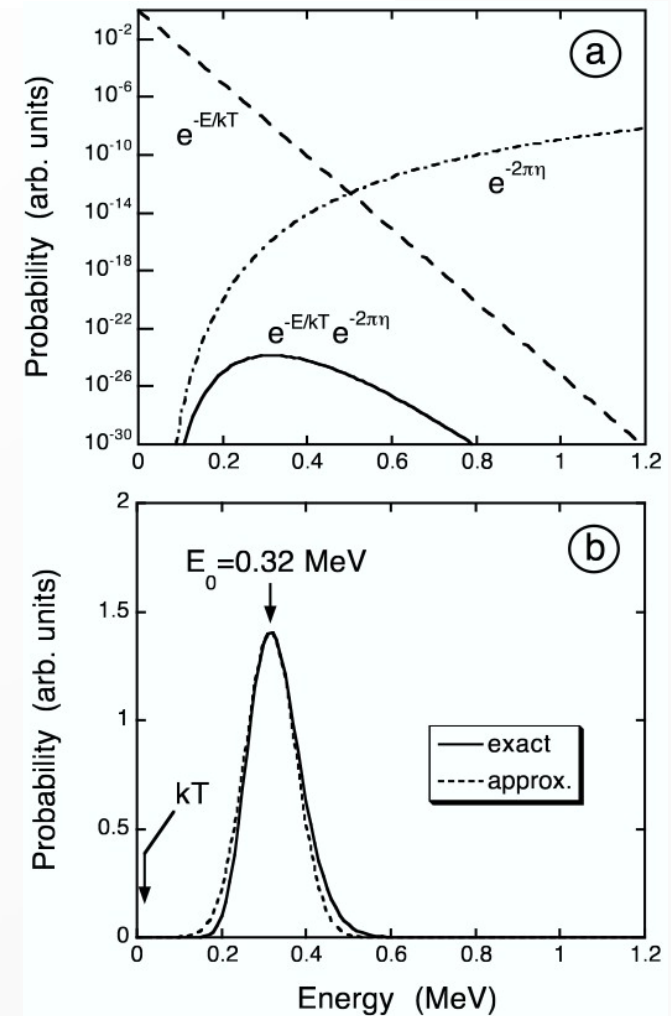
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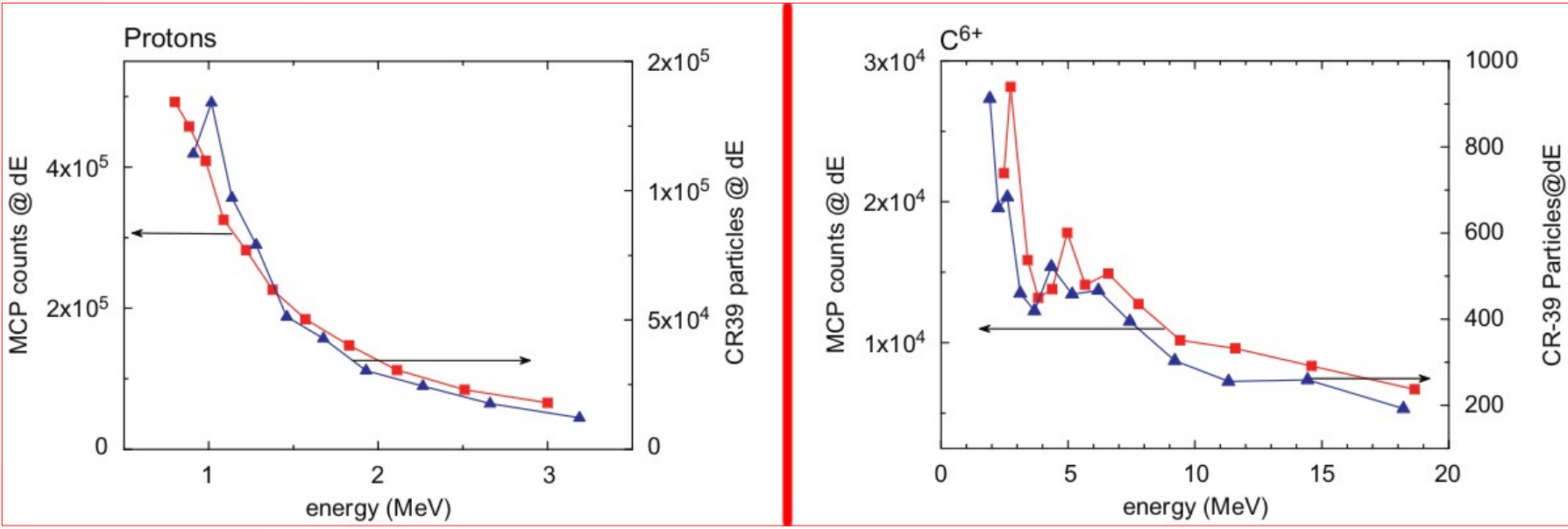
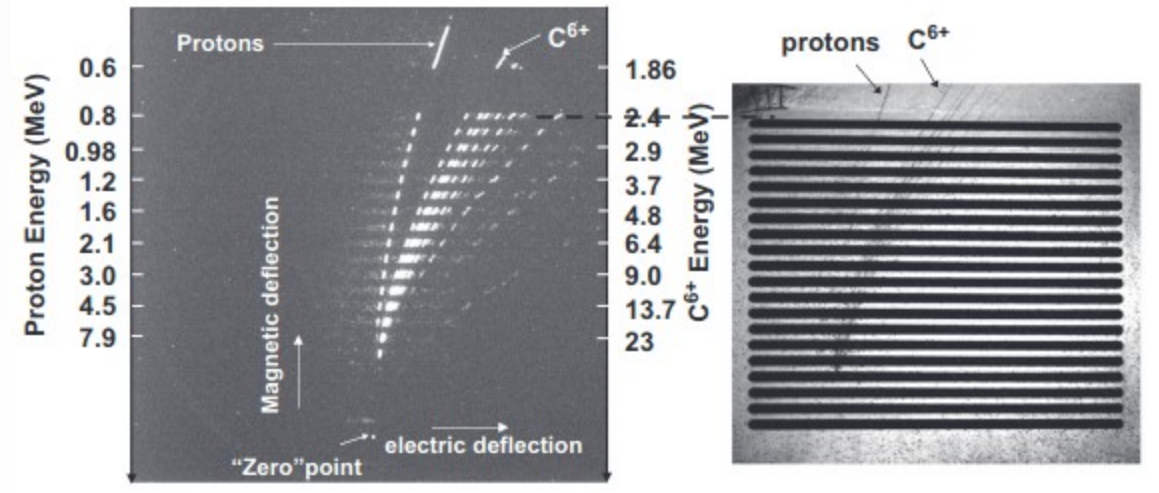
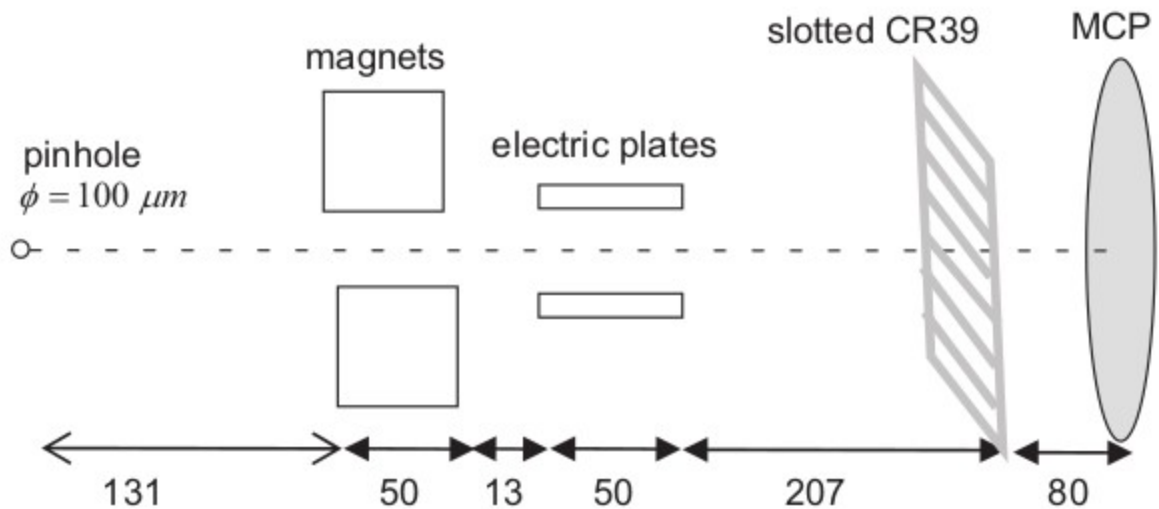
10^2

10^3

$$Y \approx \frac{t_d}{2} \int n_D^2 \langle \sigma_{D(D,n)^3He \nu} \rangle_{kT} dV + N_{ion} \int n_D \langle \sigma_{D(D,n)^3He} \rangle_{kT/2} dl$$

$^{12}\text{C}(\alpha, \text{g})^{16}\text{O}$ - Nuclear
Physics of Stars, C. Iliadis





Read by IP, MCP, LANEX, CMOS ..

R. Prasad et al., Calibration of Thomson parabola—MCP assembly for multi-MeV ion spectroscopy, Nuclear Instruments and Methods in Physics Research A 623 (2010)