Dark matter searches: status and prospects

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Classical (local) DM tests

[Clowe et al., 2006]

(merging) galaxy clusters

(rotationally supported) galaxies

Petac & P.U*., 2020]*

 ms Nowadays not the leading rationale to argue for the *existence* of DM.

labs: these observations provide are mostly done within these DM $\frac{1}{\sqrt{2}}$ a $\frac{1}{\sqrt{2}}$ Direct/indirect searches for DM key ingredients, e.g., DM density and velocity profiles.

Cosmological evidence for DM

cosmic baryon density, while the wider band indicates the BN-The SM for cosmology (ΛCDM model) as a minimal recipe to embed the Universe dynamics and a consistent theory for structure formation, tested against a pletora of cosmological probes, in which the DM term is treated as a classical, cold, pressure-less fluid subject to gravitational interactions only.

Figure 24.1: The primordial abundances of ⁴He, D, ³He, and ⁷Li as predicted by the standard model of Big-Bang nucleosynthesis — the bands show the 95% CL range [44]. Boxes indicate the

> The leading rationale to argue for the existence of DM and the tool to precisely measure it.

"Concordance" cosmology: $\Omega_{\text{DM}}h^2 = 0.1200 \pm 0.0012$ *[Planck Coll., arXiv:1807.06209]*

"Surprises" versus concordance cosmology

The Λ CDM model under extreme scrutiny may show, on top of the astonishing successes, also few discordances. Most relevantly for dark matter/dark sector:

- A small-scale crisis of the CDM paradigm (in the deep non-linear regime, likely where baryonic component modelling do count)? Observational cores versus predicted cusps in the density profile of small dark-matter-dominated galaxies; missing satellites, in particular in the count for the most massive subhalos in the Milky Way and the Local Group - the too-big-to-fail problem; …

- Tensions in cosmological parameters? Indirect versus direct measurements of H_0 - 5σ discrepancy between Planck CMB and SH0ES SNIa + much more; early Universe versus late Universe determination of the normalisation of the matter power spectrum - discrepancy between CMB and weak lensing at $\sim 3\sigma$ level in estimates of S_8 ;

- Surprise from the JWST discovery of the existence of early massive galaxies with stellar masses way larger than what expected within the ΛCDM model;

- DESI baryon acoustic oscillation measurements favouring dynamical dark energy (assuming a time varying equation of state: $w(a) = w_0 + w_a \,(1-a)$ at $\sim 3\sigma$ level the quadrant $w_0 > -1$, $w_a < 0$ is preferred).

Insights on the "true nature" of dark matter and dark energy?

Rephrasing DM as a particle physics problem

In the Λ CDM model the DM term is scale free: there is no insight on how to reformulate the DM puzzle in terms of elementary particles (what mass? what interaction strength with ordinary matter or among themselves?)

The small-scale crisis pointing to an excess of power on small scales (or maybe to baryonic components/baryonic feedback not properly treated in the simulations). Remove power by introducing a new physical scale associated to DM particles: a free-streaming scale (e.g. warm dark matter); a selfinteraction scale; a macroscopic "quantum" scale (e.g. dark matter as a BEC); a large DM-baryon or DM-photon interaction scale; …

Suppressing S_8 at late times, letting dark matter decay or cannibalise itself? Play with subdominant components which again dump power (self-interacting DM, very light axion-like DM, …)?

Steadily moving towards a scenario in which, rather than the SM + a DM particle, you have SM + a multicomponent dark sector in which address the dark matter problem and much more (e.g. the H_0 tension with some early dark energy component???).

The real of (moderately motivated) prejudices

As a starting **assumption**, consider a dark sector in terms of elementary particles, to be possibly treated in the dilute limit (twobody interactions dominating over multi-body interactions).

Disclaimer: this is not the only possible extrapolation! In this talk we will not consider, e.g.: scenarios with "macroscopic granularities", such as primordial black holes - possibly still viable; or scenarios in which gravity is not described by general relativity - no "DM free" variant found so far matching observations on all scales.

Two extra guidelines have been the main model building prejudices:

1) we need a "natural" mechanism to generate dark matter in the early Universe

2) there are some aspects which are not satisfactory in the standard model of particle physics, addressing such open issues will lead to an extension of the standard model embedding dark matter as well.

SM of cosmology & BSM of particle physics $\sum_{i=1}^{n} a_i$ **b**

 A common (particle physicist) roadmap some years ago: $\overline{}$ \sim \sim \sim \sim \sim

i) A (set of) BSM state(s) to be found at colliders; i) A (set of) BSM state(s) $\frac{1}{\sqrt{1}}$ \sim $\frac{1}{\sqrt{2}}$ $\overline{9}$

ii) Direct detection experiments to demonstrate that the (lightest) state is stable and makes the dark matter. detection oului capoilitoillo lu u

% \mathbf{r} $\overline{}$ \mathbf{H} n natura A trigger from naturalness versus the hierarchy problem, and \overline{a} thermal relic **WIMPs** as natural dark matter candidates. (96)

Thermal relics directly coupled to the baryon/photon primordial bath: $\chi \bar{\chi} \leftrightarrow \text{SM} \bar{\text{SM}}$ (with SM is some lighter Standard Model state) <u>"σ Av#T =</u>

$$
\Omega_{\chi} h^2 \simeq \frac{3 \cdot 10^{-27} \text{cm}^{-3} \text{s}^{-1}}{\langle \sigma_A v \rangle_{T=T_f}}
$$

WIMP miracle: "fixed" DM pair annihilation cross section into "visible" particles.

A recipe that can work below about **100 TeV** (unitarity limit *[Griest & Kamionkowski 1990]*; in realistic models up to about 15 TeV) and gets inefficient below about **1 GeV**.

SM of cosmology & BSM of particle physics

A common (particle physicist) roadmap some years ago:

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thermal relic **WIMPs** as natural dark matter candidates.

Thermal relics: the familiar and beloved scheme

ive (and it will alive (and it will be hard to kill the WIMP paradigm is well **fragmentation 1S decay process** trigger" is fading away, *SM* it), however the "naturalness **fragmentation** *species particles species particles* **and/or decay process** *lighter stable* So far, a scheme which has lead only to tentative (and controversial) hints of signals: making to some extent the framework less appealing.

Indirect detection

In principle a straightforward connection between annihilation in the early Universe and in todays DM halos.

Difference Constraints from prompt McDaniel et al., 2023]
P-ray emission (continuum 10^{-21} PH • Population of dwarf spheroidal galaxies have population of dwarf spheroidal galaxies, ideal DM labs: very Strong constraints from prompt McDaniel et al., 2023] γ-ray emission (continuum spectrum) from the local large mass to light ratios, quiet astrophysical environments, relatively close.

In the same plot, models compatible with the excess detected by Fermi from the Galactic Center *[e.g. Fermi collaboration, 2017]*, not an ideal DM lab (is it due to MSPs? *[e.g., latest: Manconi et al.,*

[latest update with 14-yr Fermi data,

Indirect detection, look for signatures

Still a chance to detect smoking gun signals, such as γ-ray lines (arising at loop level): $\chi + \chi \rightarrow \gamma + X$

prominent in some specific models, such as pure Wino DM, and a target for CTAO in galactic center observations:

still to be cross-checked against other detection channels:

[Hryczuk, PU et al., 2014]

 F_{∞} if \Box . Antideuterons as a turther smoking \Box $\mathbb{R}^{S_{\mathbb{Z}/2}}$ $\Big\{$ Antideuterons as a further smoking $\Big\}$ derived using diagonal assumptions. For antiprotions antiprotons (ρ <u>the local density and in the ISRF antarctic</u> containing the ISRF and many latitude and many low latitude and many l $t_{\rm max}$ radiation for the interstellar galaxy and interstellar galaxy $t_{\rm max}$ Tev₁ and the *I-factor constant in the J-factor and foreign*

Direct detection

Measuring the recoil energy in elastic (or inelastic?) scattering of (local Milky Way halo population) dark matter particles off nuclei:

A huge experimental program, originally devised to address "vanilla" (?) WIMPs, with mass close to the weak scale:

Direct detection (2)

Since several years, searches with noble liquid (mainly Xe) detectors have taken the lead in the experimental effort (background rejection + scalability): **CROSS IN THE EXPENIMENTAL ENDIT (DACKGROUND** LAALS SAALCHAS WITH D $m_{\chi} = 50 \,\text{GeV}$

[Snowmass 2021, arXiv:2203.08084, adapted by Baudis 2024]

Direct detection (3)

equilibrium with the SM heat bath, (mostly) because carrying a SM Current limits versus a representative set of dark matter models within the standard WIMP paradigm, namely states in thermal gauge charge (e.g.: SUSY, extended Higgs sectors, …):

 10^{-39} - covering the full expected mass 10^{-41} range; - possibly extending f_{DM} σ_{SI} [cm²] 10^{-43} **Direct Detection** to very small cross sections (crossing 10^{-45} symmetry not Meutrino Fog fulfilled); 10^{-47} ggs **Triplet Majorana** - mostly not NMSSM-1 6
C 工 addressing the **Two** 10^{-49} $NMSSM - 2$ hierarchy problem:100 1000 $10⁴$ 10 M_{DM} [GeV]

> Figure 5: Plots of dark matter-proton cross section σSI times dark matter fraction *f*DM *[Snowmass 2021, arXiv:2203.08084, & refs. therein]*

Enlarging the parameter space

A dark sector containing multiple states offers the possibility of having multiple variants to the standard WIMP paradigm, such as: *[very incomplete lists of models and references]*

- \bullet there is a dark sector thermal bath (with T_D possibly different from T_γ), with thermalisation and freeze-out: $\chi + \chi \leftrightarrow \psi + \bar{\psi}$ led by extra interactions (e.g.: extra $U(1)_D$ with γ_D mediator)
- there is a different process sustaining (dark) thermal populations, such as, e.g.,:
	- coannihilations: $\chi_i + \chi_j \leftrightarrow \psi + \bar{\psi}$ [Griest & Seckel, 1991]
	- semi-annihilations: $\chi + \psi \leftrightarrow \chi + \chi$ [D'Eramo & Thaler, 2010]
- dark matter produced out of equilibrium, because, e.g.:
	- there is a particle-antiparticle asymmetry η_χ (analogous or connected to η_B), asymmetric DM [e.g., Petraki & Volkas, 2013]
	- feebly interacting with the heat bath / never in equilibrium:
	- * super-WIMP freeze-in: $\psi + \phi \rightarrow \chi + \phi$ [Pagel & Primack, 1992]
	- * exponential production: $\psi + \chi \rightarrow \chi + \chi$ [Bringmann et al. 2021]

strong $\overline{}$ TeV - SM inter.
---------------------9 $10s$ TeV SM inter \Join mass $\bm{\times}$ sub keV gravity <

 m

Direct detection away from "vanilla" WIMPs

Below the GeV mass scale alternative techniques and strategies (e.g. electron recoils or small band gaps for electrons excitations):

▸ Experimental Panorama

Filling in the range of possibilities in the EU

At some early cosmological epoch (temperature much larger than the particle mass) the abundance of the DM candidates relative to SM particles also spans huge ranges, e.g.:

- It is order 1 for **WIMPs** (since the sizeable interaction ensures thermal equilibrium)

- It is very small for **super-WIMPs** (never in thermal equilibrium because of their tiny interactions, e.g. they leak out the thermal bath through the freeze-in mechanism)

- It is very large for **super-cold DM** (very light bosons, almost non interacting, with huge occupation numbers of their lowest momentum state, e.g.: axion DM)

Natural matching $\Omega_\chi \sim \Omega_{\rm CDM}$? Several of the scenarios mentioned above simply do envisage fine-tuning.

Reintroducing a (another) particle physics motivated framework?

The axion solution to the strong CP problem $\frac{1}{2}$ **The axion solution to the strong CP problem** where *G* is the gluon field strangth. This lagrangian term violates *P* and *T*, and therefore, by **CPTC** also consider to the otrong of proprom. *CPT* theorem, also *CP*. The operator *GG*˜ is a total derivative, but because of the non-trivial the axion solution to the strong of problem

In the Lagrangian of QCD is not forbidden to have: s_{max} In the Lagrangian of QCD is not forbidden to have: $tho!$ \overline{a} and \overline{b} axion \overline{a}

$$
\mathcal{L}_{\text{QCD}} \supset \theta \frac{g_s^2}{32\pi^2} G \tilde{G}
$$
 (gluon field strength contracted to its dual)

which violates T, P & CP and give rise to dangerous operators such as an electric dipole moment for the neutron: nich violates 1, P & CP and give rise to dangerous operators such as an
ectric dipole moment for the neutron: ⇠ ¹⁰¹⁶✓ *^e ·* cm (3.3) *f f f f f******f f******f f f f f f******f f*

$$
d_n \sim \frac{\theta}{(4\pi)^2} e \frac{m_\pi}{m_N^2} \sim 10^{-16} \theta e \cdot \text{cm}
$$

 μ_{N}
 μ_{N} μ_{N from a loop diagram where α and β run in the loop and β run in the loop and β emits a photon and β Experimentally $|d_n| \lesssim 10^{-26} e\mathrm{~cm} \implies \theta \lesssim 10^{-10}$ The strong CP problem: why is this parameter so small? Experimentally $|u_n| \geq 10$ *e* chiral \rightarrow $v \geq 10$ is this param $T_{\rm vnormal}$ by σ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ Experimentally $|d_n| \lesssim 10^{-26}e$ cm $\implies \theta \lesssim 10^{-10}$ interstivily

The strong CP problem: why is this parameter so small?

The axion solt **Fandal** if to a dynamical field and "redurity The axion column promote it to a crynannoar noid and notax it to of so-called "*strong CP problem*". The axion solution: promote it to a dynamical field and "relax" it to 0: The axion solution: promote it to a dynamical field and "relax" The axion s

$$
\mathcal{L} \supset \left(\theta + \frac{a(x)}{f_a}\right) \frac{g_s^2}{32\pi^2} G\tilde{G}
$$

initially with flat potential but at $T \sim \Lambda_{QCD}$ acquiring via i tially with flat potential but at Δ initially with flat potential but at $T \sim \Lambda_{QCD}$ acquiring via instantons: initially with flat potential but at $T \sim \Lambda_{QCD}$ acquiring via in:

$$
V(a) \simeq \Lambda_{QCD}^4 \left(1 - \cos\frac{a}{f_a}\right) \simeq \frac{1}{2}\Lambda_{QCD}^4 \frac{a^2}{f_a^2} + \dots \simeq \frac{1}{2}m_a^2 a^2 + \dots
$$

The QCD axion as a dark matter candidate

Misalignment production: very light scalars trapped in modes with coherent oscillations, which behave cosmologically as CDM:

 $\sim 10^{-6} - 10^{-4}$ eV *(also ferm from sf* Follow the eq. of motion: $\ddot{a} + 3H\dot{a} +$ dV *da* $\simeq \ddot{a} + 3H\dot{a} + m_a^2 a = 0$ when $3H < m_a$, coherent oscillations with frequency m_a start, i.e. $\rho_a =$ 1 2 $\rho_a = \frac{1}{2} \left(\dot{a}^2 + m_a^2 a^2 \right) \;$ evolves as: $\dot{\rho}_a \simeq$ $\int \dot{m}_a$ *m^a* 3*H* ◆ $\rho_a \Rightarrow \rho_a \propto$ *m^a* $(1 + z)^{-3}$

i.e. as matter. $\Omega_a\sim \Omega_{\rm CDM}\ \Rightarrow\ m_a\sim 10^{-6}-10^{-4}\,{\rm eV}$ (also term from string decay)

The phenomenology mostly based on the axion coupling with photons:

$$
\mathcal{L}_{a\gamma\gamma} = -\frac{g_{a\gamma}}{4} F_{\mu\nu} \tilde{F}_{\mu\nu} a = g_{a\gamma} \mathbf{E} \cdot \mathbf{B} a
$$

within a specific model: α $\left(E\right)$

$$
g_{a\gamma} = \frac{\alpha}{2\pi f_a} \left(\frac{E}{N} - 1.92(4) \right)
$$

= $\left(0.203(3) \frac{E}{N} - 0.39(1) \right) \frac{m_a}{\text{GeV}^2}$

Generalising to axion-like particles (ALPs)

Misalignment production as a generic scheme for light scalar CDM. Remove the link to the strong-CP problem and find much lighter or more strongly coupled cases. Models arising, e.g., in stringy frameworks; possible links with inflation, dark energy, …

haloscopes to helioscopes, to cooling of stellar systems, to modifying the γ-ray 10^{-6} A more generic parameter space and a variety of detection techniques: from horizon, … *[Snowmass 2021, arXiv:2203.14923]*

A set of new ideas and a set of new search techniques being explored and looking promising:

A few windows of opportunity

why he is **porture the marky he wink certisty strip where lamppost cartoon...** *minutes the policeman asks if he is sure he lost them here, and the drunk replies, no, and that he lost them in the park. The policeman the light is."*

Waiting for super-precision cosmology to solve it all (but on small scales and the difficulties in modelling baryons, it is not expected to happen very soon), the dark matter phenomenologist faces hard times, running the risk of getting trapped by the infamous "streetlight effect".

ham Kaplan who referred to this as the "principle of the 20 is called the Can Gall it lampnost searches, still the scientific program is very d Freedman In 2010 Book wrong. Also alle we few gaty and ample cases: ept has been used by social scientists way back in 1964 like

Dark matter boosted by galactic cosmic rays

Reconsider the "light mass wall" in nuclear recoil direct detection experiments, namely the maximum recoil energy does **not** fulfil:

$$
E_R^{\text{max}} \sim \frac{m_\chi^2 v^2}{M_N} > E_{\text{th}}
$$

for non-relativistic galactic DM halo particles ($v \sim 10^{-3}c$), if the DM particle mass m_χ is lighter than $\sim 1 \, \text{GeV}$.

However the same coupling DM-ordinary matter being tested in the experiment, may be relevant in the up-scattering by galactic cosmic rays (mainly $\frac{2}{5}$
protons) of a fraction of the DM $\frac{1}{5}$ protons) of a fraction of the DM galactic population to high energies, making sub-GeV dark matter candidates potentially detectable.

Blazer boosted dark matter

Is there in Nature a potentially more powerful and/or more efficient dark matter booster?

Blazers are the ideal case:

Extremely powerful flux of protons (electrons) through an extremely dense dark matter environment (dark matter spike accreted around the blazer black hole engine), potentially generating a sizable DM flux towards us. our simplified method are in good agreement with those protono (oloolono) unough could be avoided by averaging over the full production of the full product over the full product of α DM-proton scattering is too weak to leave any recoil in engine), potentially t_{scat}

T^N the nuclear recoil energy is the BBDM induced target that the second target target target the second target target target target target target target ta Tightest limits/best discovery a iu iigit dan i potential for light dark matter

 $See also:$

[Granelli, P.U., Wang, JCAP 2022]

Multi-wavelength signals from dwarf galaxies?

On top of prompt γ-ray emission, in dwarf galaxies there can be radiative emission connected to leptonic components from DM annihilations/decays:

Early analysis predicting such signal for Draco:

If you trust these predictions (as done in some later analyses), you would conclude that from radio surveys which did not detect such a signal can put constraints in the (WINEP) parameter space at a leyel competing with γπειν telescopes **10** Way cosmic ray halo**?** (as assumed **10 -17** ν **I(** $($ reqų ired to contine $e^+\!f\!\!\!\!/e^-$) in dwarfs? **10** caveat: what is the level of turbulence Piugiun
100¹ **-14** $\sqrt{2}$ **an M_u = 100 GeV, σv at EGRET level**
an Dut constraints in the **7.5 10 12.5 12.5 12.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 20 22.5 β** 30-14-15 cm-2 stription \overrightarrow{BA} in \overrightarrow{A} fight 0.5 0.1 These predictions are how sver model dependent, with especially a serious Comparable to the one in the Milky in the plot)

23 **log (** ν **[Hz])**

log (ν **[Hz])**

Self-confinement of DM-induced cosmic rays (e.g., a spherical source versus turbulence flowing along lines of a regular magnetic field being large scale dipole or **Self-confinement of DM-induced cosmic rays** 5 @*W* = DI *Dkk*(*W*) luced c (*r*²*vAW*) + CR(*ne, k*)*W .* (4)

DM would naturally induced a charged particle density gradient, in turn sourcing turbulence: solve in a dwarf the two coupled eqs.: Divi would indicturity inducted a charged parties which may depend on the radial distance *r* and energy *E*, *v^A* is the Alfv´en velocity that transports the electrons sourcing turbulence: solve in a dwarf the two coupled eqs.: @*t* @*k* @*k r*2 @*r* rcing turbulence: solve in a dwarf the two coupled eqs.:

i)
$$
\frac{\partial n_e}{\partial t} = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 D \frac{\partial n_e}{\partial r} - r^2 v_A n_e \right] + \frac{2v_A}{r} \frac{\partial}{\partial E} \left[\frac{p}{3} \beta c n_e \right] - \frac{\partial}{\partial E} \left[\dot{E} n_e \right] + q_{\text{CR}}
$$
\nwith: $q_{\text{CR}}(r, E) = \langle \sigma v \rangle_f \frac{\rho_{\text{DM}}^2(r)}{2m_{\text{DM}}^2 dE}$
\ni)
$$
\frac{\partial W}{\partial t} = \frac{\partial}{\partial k} \left[D_{kk}(W) \frac{\partial W}{\partial k} \right] - \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_A W) + \Gamma_{\text{CR}}(n_e, k) W
$$
\nwith: $D_{kk}(W) = c_k v_A k^{7/2} \sqrt{W}$ and $\Gamma_{\text{CR}} = \frac{4\pi c v_A}{3k W(k) B_0^2/(8\pi)} \left[\beta(p) p^4 \left| \frac{\partial f_e}{\partial r} \right| \right]_{p=p_{\text{res}}}$
\nWithin a simplified
\ngeometry and at
\nsteady state a
\nrelevant diffusion
\naxis:
\n
$$
D(r, p, t) = \frac{D_B(p) 4/\pi}{k W(r, k, t)}
$$
\n
$$
= \frac{p_{\text{max}} - p_{\text{max}} \beta}{\frac{p_{\text{max}} - p_{\text{max}} \
$$

24

 Z

 $\frac{L \left[\alpha v \right]_1}{24}$

Pegis, P.U. et al, 2023

[Regis, P.U. et al, 2023]

Self-confinement of DM-induced cosmic rays (e.g., a spherical source versus turbulence flowing along lines of a regular magnetic field being large scale dipole or **Self-confinement of DM-induced cosmic rays** 5 @*W* = DI *Dkk*(*W*) luced c (*r*²*vAW*) + CR(*ne, k*)*W .* (4) The value of **B0, instead, see chain** and the synchrotron power and the synchrotron power and the confinement time the signal go in the opposite directions, partially canceling each other. Indeed, the larger is *B*⁰ the larger is *Psyn* (thus enhancing the signal), but also the larger is *D* (which depletes the signal), as we already discussed in the previous

DM would naturally induced a charged particle density gradient, in turn sourcing turbulence: solve in a dwarf the two coupled eqs.: Divi would indicturity inducted a charged parties which may depend on the radial distance *r* and energy *E*, *v^A* is the Alfv´en velocity that transports the electrons sourcing turbulence: solve in a dwarf the two coupled eqs.: @*t* @*k* @*k r*2 @*r* bound. rcing turbulence: solve in a dwarf the two coupled eqs.: The only exception to this argument is for scenarios with low *B*0, low *m*DM and soft channels of annihilation. e: solve in a dwarf the two coupled egs.: $\qquad \qquad \bullet$ s. Solve in a awar the two ocapied equ..

i)
$$
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$$
\nwith:
$$
q_{\text{CR}}(r, E) = \langle \sigma v \rangle_f \frac{\rho_{\text{DM}}^2(r)}{2m_{\text{DM}}^2} \frac{dN_e^f}{dE}
$$
\nii)
$$
\frac{\partial W}{\partial t} = \frac{\partial}{\partial k} \left[D_{kk}(W) \frac{\partial W}{\partial k} \right] - \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_A W) + \Gamma_{\text{CR}}(n_e, k) W
$$
\nwith:
$$
D_{kk}(W) = c_k v_A k^{7/2} \sqrt{W} \text{ and } \Gamma_{\text{CR}} = \frac{4\pi c v_A}{3k W(k) B_0^2/(8\pi)} \left[\beta(p) p^4 \left| \frac{\partial f_e}{\partial r} \right| \right]_{p=p_{\text{res}}}
$$
\n
$$
\text{Leading to}
$$
\n
$$
\frac{10^{30}}
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\n
$$
\frac{10^{30}}
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\n
$$
\text{consecutive but}
$$
\n
$$
\frac{10^{30}}
$$

25

nice overview is given for example in Ref. [31]. For the systems under consideration, i.e., with low stellar turbulence,

A minimal DM scheme and $(g$ -2) $_{\mu}$

Account for the muon (g-2) anomaly within the most minimal BSM recipe embedding also a DM candidate: a thermal relic pure Bino + 2 scale muon partners (this is NOT the MSSM).

It works up to the TeV scale and beyond:

[Acuña, Stengel, P.U., PRD 2022]

A minimal DM scheme and $(g$ -2) $_{\mu}$

Account for the muon (g-2) anomaly within the most minimal BSM recipe embedding also a DM candidate: a thermal relic pure Bino + 2 scale muon partners (this is NOT the MSSM).

No "traditional" WIMP detection method working in this case; kinetic heating of neutron stars would be instead extremely efficient and future infrared surveys of old neutron star populations should probe the entire parameter space!

[Acuña, Stengel, P.U., arXiv:2209.12552]

Conclusions

Is cosmology shaping a dark sector which will guide us to the solution of the dark matter problem?

New prejudice-free paths to address the dark matter problem from a particle physics perspective; are there efficient ways of walking through them and discriminating among each other?

Several windows of opportunities for dark matter detection still open; will we enter the stage in which models are confirmed or rejected?