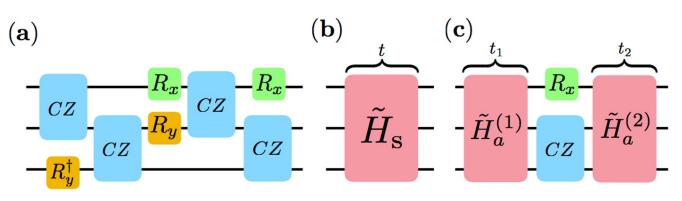
Quantum Architectures for Analogues and Theory Applications

Outline



- Digital Simulations
 - Spin Systems (PRL 112, 200501 (2014), Phys Rev X 5 021027 (2015))
 - Fermionic Systems (10.1038/ncomms8654)
 - QG (PRL 119, 040501 (2017), Nature volume 612, pages 51–55 (2022))
- Analog Simulations
 - GAAH Model (arXiv:2206.13107)
 - Hawking Radiation (Nature Communications | (2023) 14:3263)
 - Dirac Equation (New Journal of Physics 15 (2013) 055008)
- Digital Analog Simulations
 - Fermion Fermion scattering (PRL 114, 070502 (2015))
- Optimal Control (PHYSICAL REVIEW A 101, 062307 (2020))
- Analog Circuits (REVIEWS OF MODERN PHYSICS 84 2012, EPL 128, 24002 (2020)

Digital Simulations

(a) $\begin{bmatrix} CZ & R_x \\ -R_y \end{bmatrix} \begin{bmatrix} CZ & -R_x \\ -R_y \end{bmatrix} \begin{bmatrix} CZ & CZ \\ -R_y \end{bmatrix} \begin{bmatrix} CZ & CZ \\ -R_y \end{bmatrix}$

Spin Systems

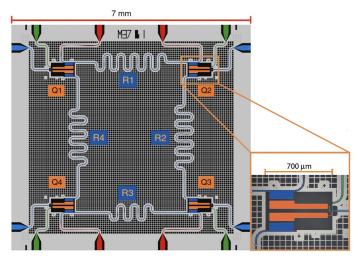
Heisenberg Model

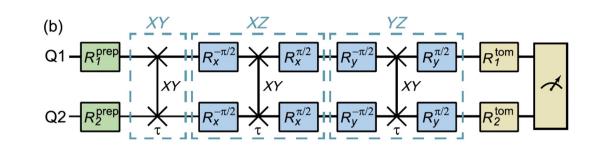
$$H_{xyz} = \sum_{(i,j)} (J_x \sigma_i^x \sigma_j^x + J_y \sigma_i^y \sigma_j^y + J_z \sigma_i^z \sigma_j^z)$$

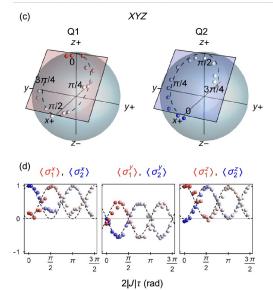
Suzuki-Lie-Trotter

expansion

$$e^{-iHt} \simeq \left(e^{-iH_1 t/l} \cdots e^{-iH_M t/l} \right)^l + \sum_{i < j} \frac{[H_i, H_j]t^2}{2l}$$







Phys Rev X 5 021027 (2015)

Fermionic Systems

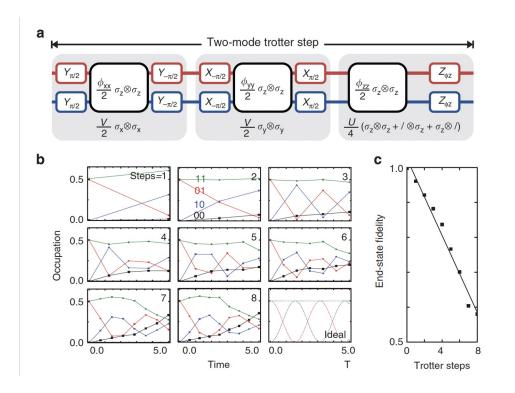
Fermi-Hubbard model

$$H=-Vig(b_1^\dagger b_2+b_2^\dagger b_1ig)+Ub_1^\dagger b_1b_2^\dagger b_2$$

Jordan-Wigner mapping

$$\begin{array}{rcl} b_1^{\dagger} & = & \mathbb{I} \otimes \mathbb{I} \otimes \sigma^+, \\ b_2^{\dagger} & = & \mathbb{I} \otimes \sigma^+ \otimes \sigma^z, \\ b_3^{\dagger} & = & \sigma^+ \otimes \sigma^z \otimes \sigma^z. \end{array}$$

a U_{14} U_{23} U_{23} U

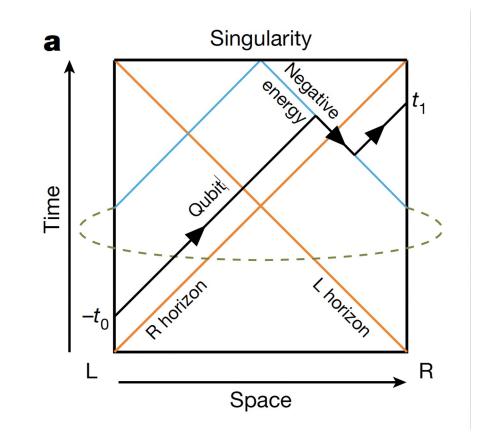


Quantum Gravity

Sachdev-Ye-Kitaev (SYK) model

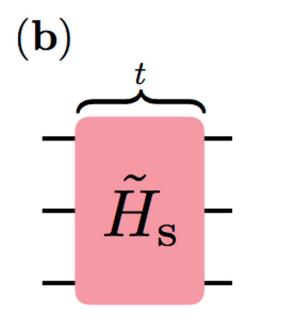
$$H = \frac{1}{4 \times 4!} \sum_{i,j,k,l=1}^{N} J_{ijkl} \chi_i \chi_j \chi_k \chi_l,$$

Majorana fermionic operators are codified as $\chi_{2n-1} = (\prod_{j=1}^{n-1} \sigma_j^z) \sigma_n^x$ and $\chi_{2n} = (\prod_{j=1}^{n-1} \sigma_j^z) \sigma_n^y$, with $\{\chi_i, \chi_j\} = 2\delta_{ij}$.



Simulation on the Google Sycamore superconducting qubit array with a nine-qubit circuit of 164 controlled-Z gates and 295 single-qubit gates

Analog Simulations



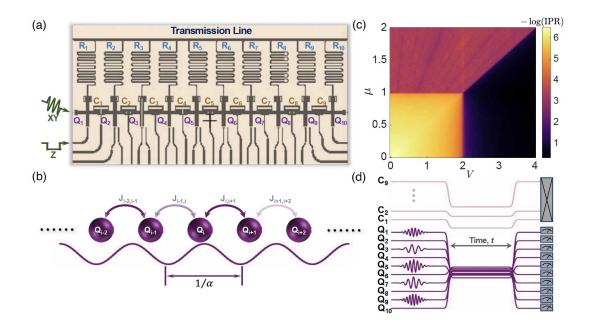
GAAH model

Generalized Aubry André Harper model

$$\frac{\hat{H}}{\hbar} = \lambda \sum_{j=1}^{9} \left(1 + \mu \cos \left[2\pi \left(j + \frac{1}{2} \right) \alpha + \delta \right] \right) \hat{a}_j^{\dagger} \hat{a}_{j+1} + \text{H.c.}$$

$$+ \lambda \sum_{j=1}^{10} V \cos (2\pi j \alpha + \delta) \hat{a}_j^{\dagger} \hat{a}_j.$$
(2)

We initially excite the leftmost qubit Q_1 , i.e., the system is initialized as $|\psi(0)\rangle = |100000000\rangle$, where $|0\rangle$ ($|1\rangle$) denotes the ground (excited) state of a qubit. Then we apply the fast Z pulse on each qubit and coupler, and the system will evolve under the Hamiltonian Eq.(2), satisfying Schrödinger equation $|\psi(t)\rangle = e^{-i\hat{H}t} |\psi(0)\rangle$. We monitor its dynamics from t = 0 to 500 ns, by measuring the photon occupancy probabilities of each qubit $P_j(t) = \langle \psi(t) | \hat{a}_j^{\dagger} \hat{a}_j | \psi(t) \rangle$. For each time point, we perform 5000 repeated single-shot measurements.



10 qubits with tunable couplers

Hawking Radiation

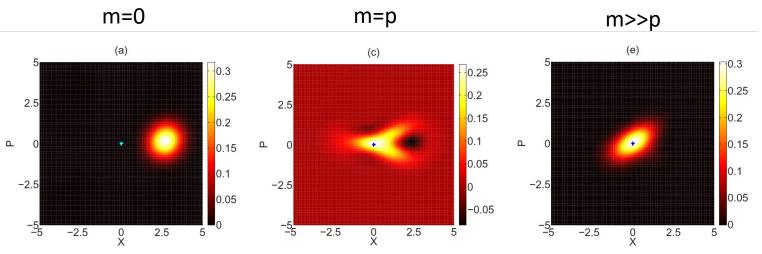
(1+1) Dira field on curved space

$$\mathrm{i}\gamma^{a}e^{\mu}_{(a)}\partial_{\mu}\psi+\frac{\mathrm{i}}{2}\gamma^{a}\frac{1}{\sqrt{-g}}\partial_{\mu}\left(\sqrt{-g}e^{\mu}_{(a)}\right)\psi-m\psi=0,$$

$$\hat{H} = -\sum_{j} \kappa_{j} \left(\hat{\sigma}_{j}^{\dagger} \hat{\sigma}_{j+1}^{-} + \hat{\sigma}_{j}^{-} \hat{\sigma}_{j+1}^{+} \right) - \sum_{j} \mu_{j} \hat{\sigma}_{j}^{\dagger} \hat{\sigma}_{j}^{-}$$

Dirac Equation

By using three classical microwave drives, superconducting qubit strongly coupled to a resonator field mode can be used to simulate the dynamics of the 1+1 Dirac equation:

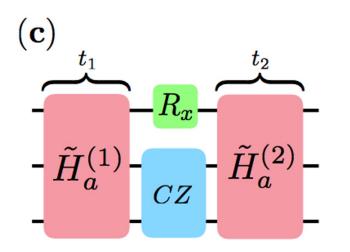


Wigner function W(x, p) of the field mode state inside the resonator

$$\mathcal{H} = \frac{\hbar\omega_q}{2}\sigma_z + \hbar\omega a^{\dagger}a - \hbar g\left(\sigma^{\dagger}a + \sigma a^{\dagger}\right) - \hbar\Omega\left(e^{i(\omega t + \varphi)}\sigma + e^{-i(\omega t + \varphi)}\sigma^{\dagger}\right) \\ -\hbar\lambda\left(e^{i(\nu t + \varphi)}\sigma + e^{-i(\nu t + \varphi)}\sigma^{\dagger}\right) + \hbar\xi\left(e^{i\omega t}a + e^{-i\omega t}a^{\dagger}\right),$$

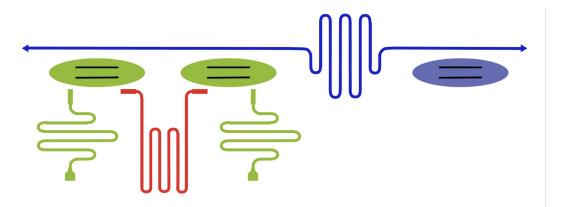
$$\mathcal{H}_{\text{eff}} = \frac{\hbar\lambda}{2} \sigma_z + \frac{\hbar g}{\sqrt{2}} \sigma_y \hat{p} + \hbar\xi \sqrt{2} \, \hat{x}, \qquad \hat{x} = (a + a^{\dagger})/\sqrt{2}, \ \hat{p} = -i(a - a^{\dagger})/\sqrt{2}, \ \text{mass} \quad \text{kinetic} \quad \text{potential}$$

Digital-Analog Simulations

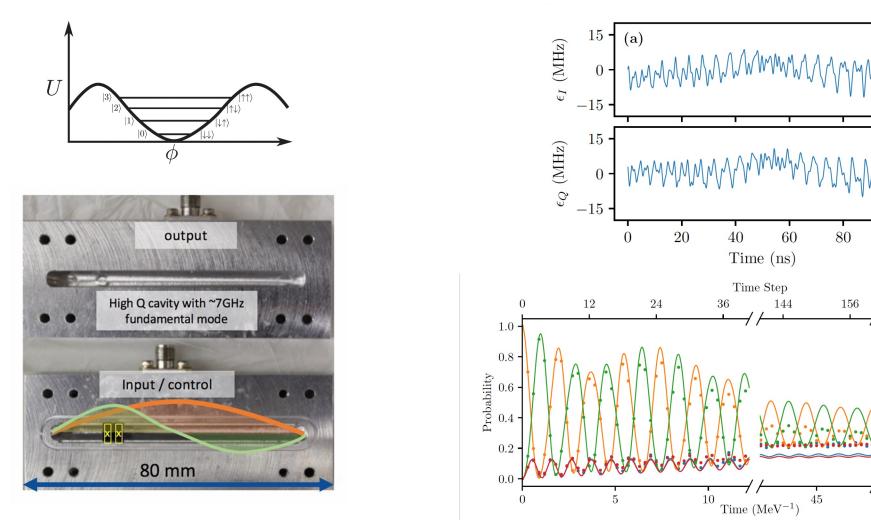


Fermion Fermion Scattering

Analog-digital quantum simulation of fermionfermion scattering mediated by a continuum of bosonic modes. Qubits simulate fermionic modes via digital techniques, while the continuum of an open transmission line simulates bosonic modes in quantum field theory.



Optimal Control on a qudit



100

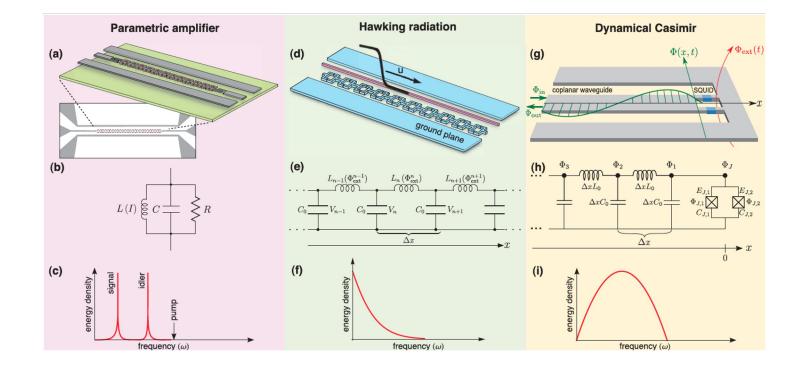
324

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95

Pederiva et al. PHYSICAL REVIEW A 101, 062307 (2020)

Analog Circuits

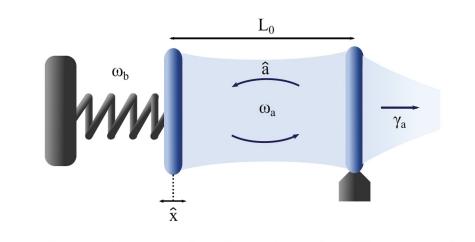


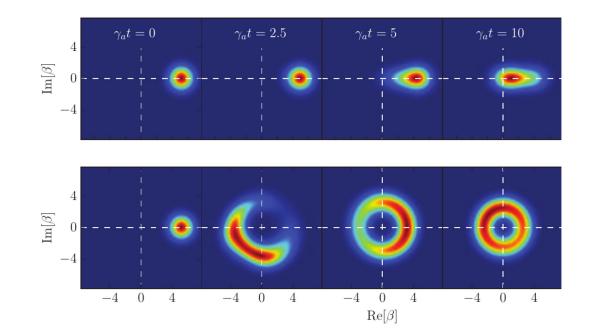
REVIEWS OF MODERN PHYSICS, VOLUME 84, JANUARY–MARCH 2012

Possibility of transforming the virtual photons corresponding to the zero-point fluctuations into an observable quantum vacuum radiation when the background on which the quantum field lives is modulated in either space or time

Baby toy model: friction force on dynamical Casimir emission

CPW





Superconducting Devices at INFN





Consiglio Nazionale delle Ricerche



CNR, National Research Council







PIEMONTE QUANTUM ENABLING TECHNOLOGY

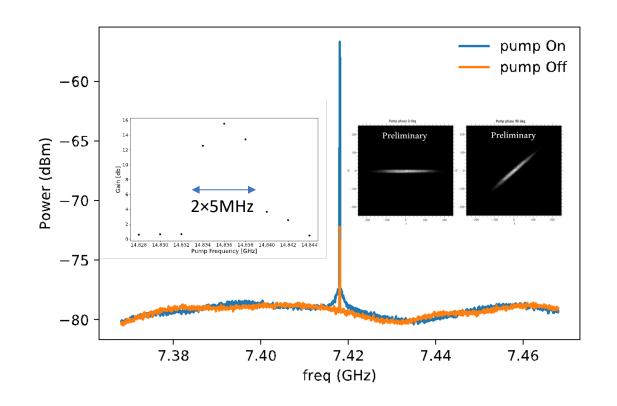
INRiM, Italian Metrology Institute

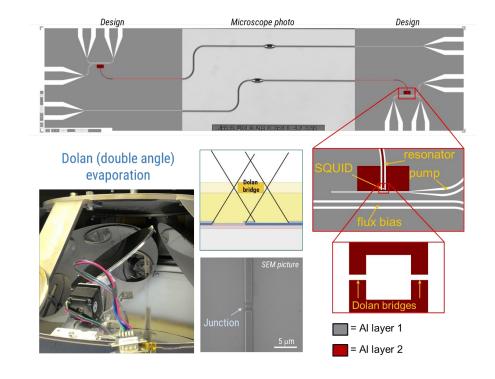


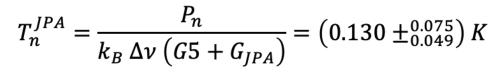
FBK, Fondazione Bruno Kessler



Realization of a Flux JPA



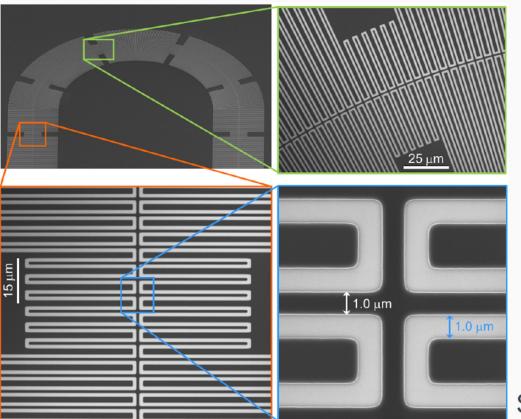


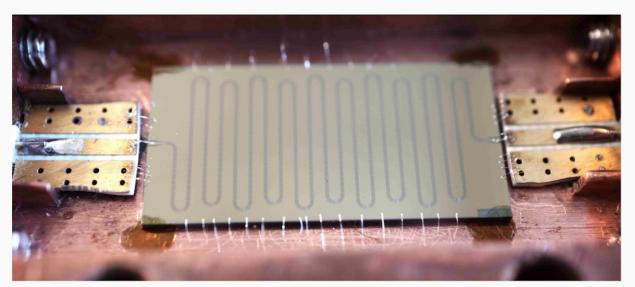


FONDAZIONE BRUNO KESSLER

KITWPA: production of the half-size amplifier





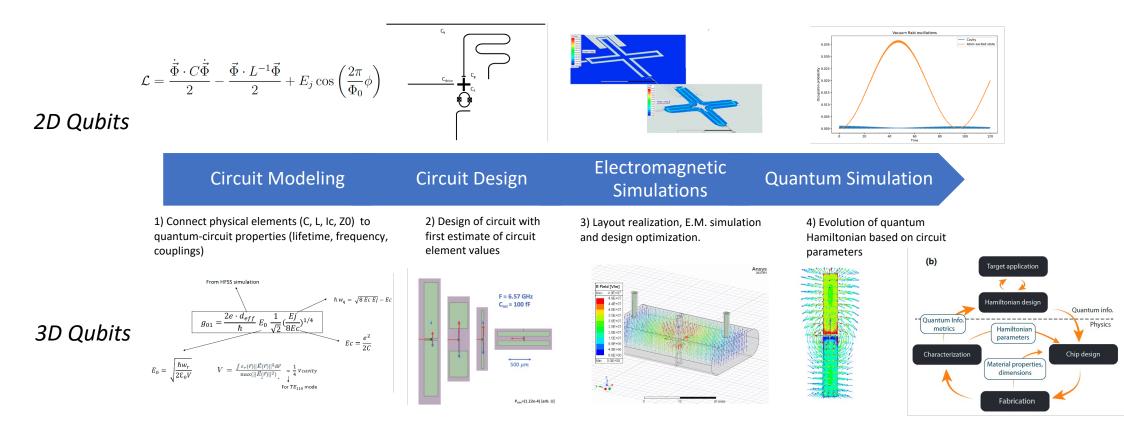


KITWPA in the copper box

SEM pictures of the fabricated KI-TWPA.

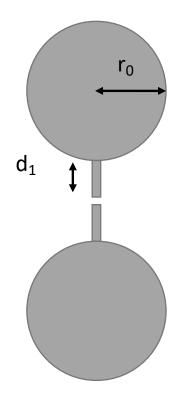
- The first KI-TWPA prototype based on the preliminary half-size layout L1 has been produced early in 2023;
- Device composed of 523 super-cells with a length of about 17 cm. Gain expected in the (7–11) dB range;
- Characterization results from these preliminary amplifiers are be crucial in refining the final design;

Design of 2D and 3D Superconducting Qubits

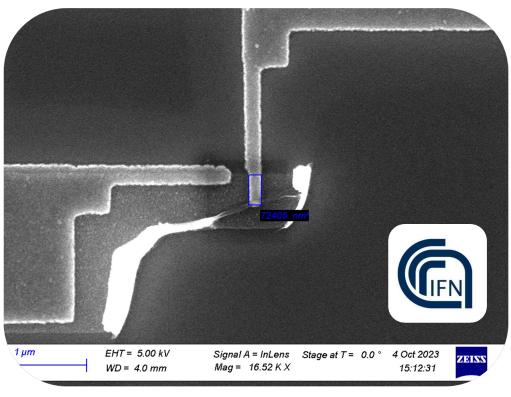


Y. Gao, PRX QUANTUM 2, 040202 (2021)

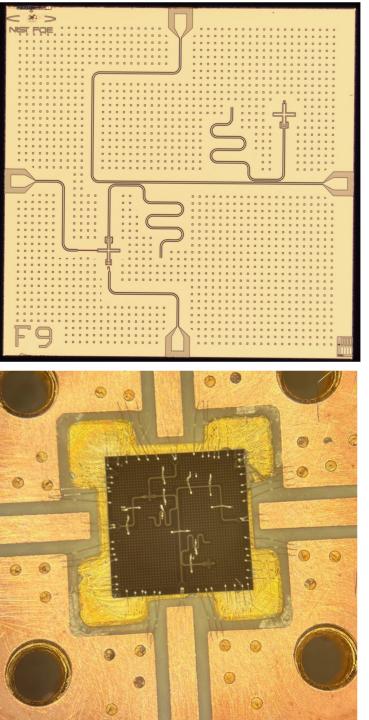
Manufactoring of 3D qubits with circular pads at CNR







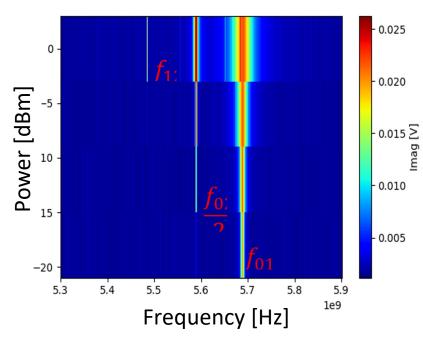
Aluminum JJ with area approx. 200 x 350 nm

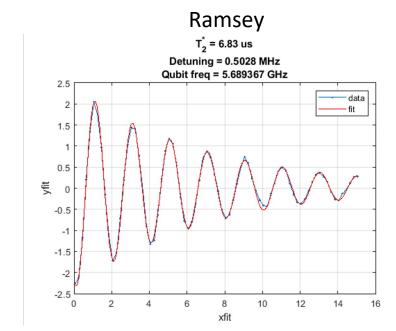


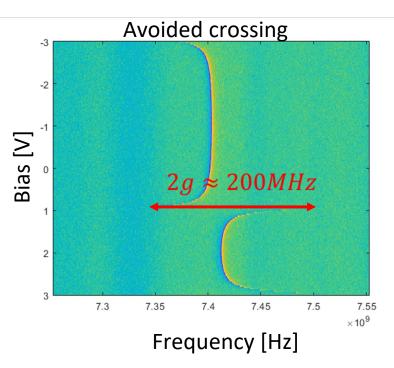
2d qubits

First version fabricated @ NIST Second Fab @ NIST currently under characterization Next fab @ FBK

Qubit spectroscopy



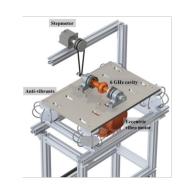


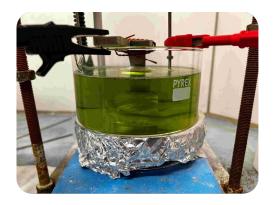


3D Cavity Fabrication









Mechanical machining

► Vibro-tumbling

► Electropolishing







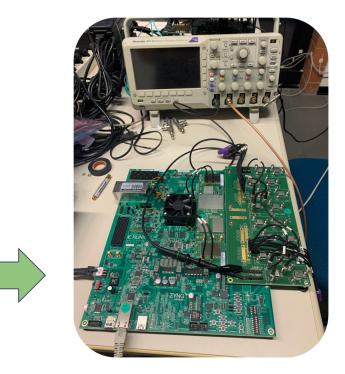


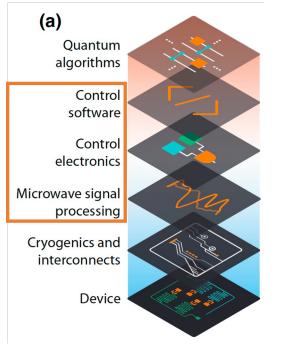
Qubit control with RFSoC hardware

QICK is a kit of firmware and software done to use the Xilinx RFSoC to control quantum systems.

Firmware exists for the ZCU111, ZCU216, and RFSoC4x2 evaluation boards.

Extending use of QICK also to the ZCU208 board.

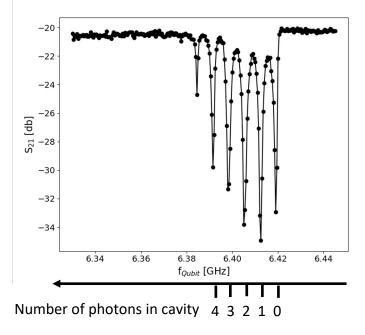




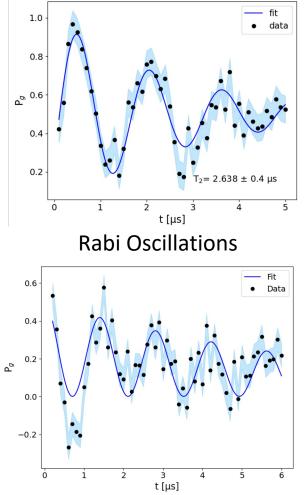


3D Resonator Coupled to a Superconducting Qubit

Photon Counting



Ramsey Spectroscopy



3D qubit manufactured at the Technology Innovation Institute of Abu Dhabi