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Modeling & Optimizing Plasma Accelerators

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It's a pleasure to be here!





Team DOLPHIN

Data-driven Optimization of Laser Physics and Interactions

PhD students

Sunny

Jannik

Chris

Jakob

Collaborators

in CALA

in Oxford

Stefan Karsch, Jörg Schreiber Peter Norreys, Robin Wang Faran, Nils, Jinpu

in Salamanca

Iñigo <u>Sola</u>, Benjamin Alonso





Modeling & Optimizing Plasma Accelerators

Modeling

- •Analytical Models
- Particle-In-Cell Models
- Machine-learning Models



Optimization

- Objective functions
- Multi-objective optimization

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• H. Ding, A. Döpp et al. Phys. Rev. E 101, 023209

The linear wakefield regime

Consider impulse response:

$$\left(\frac{\partial^2}{\partial\xi^2} + k_p^2\right)\phi = \delta(\xi)$$







Find the Green's function

$$G(\xi) = \frac{1}{2} \sin[k_p(\xi - \xi')]$$

Convolve the actual excitation with the Green's function:

$$\phi = \int_{-\infty}^{\xi} f(\xi') G(\xi - \xi') d\xi'$$

The non-linear wakefield regime







The non-linear wakefield regime



Longitudinal motion (mostly)







• H. Ding, A. Döpp et al. Phys. Rev. E 101, 023209



Modeling laser-plasma interaction The highly non-linear "blow-out" regime





No direct differential equation for the fields: **Need to model from** first principles!

Blow-out / Bubble / Ion Cavity

• H. Ding, A. Döpp et al. Phys. Rev. E 101, 023209



Modeling laser-plasma interaction The Vlasov equation

- We are interested in the evolution of a particle distribution in space, velocity and time $\rho(\vec{x}, \vec{v}, t)$
- In a collisionless system, where particles are neither created nor destroyed, the continuity equation is valid:

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \vec{v}_x \cdot \left(\nabla_x f\right) + \vec{v}_p \cdot \left(\frac{\vec{v}}{c} \nabla_p f\right) = 0$$
How does distribution change to movement of particles? How does distribution changes (forces)?
How does distribution changes to changes (forces)?
$$\frac{\partial f_e}{\partial t} + \vec{v}_e \cdot \nabla_x f_e - e\left(\vec{E} + \frac{\vec{v}_e}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \times \vec{B}\right) \cdot \frac{\partial f_i}{\partial t} + \vec{v}_i \cdot \nabla_x f_i + Z_i e\left(\vec{E} + \frac{\vec{v}_i}{c} \cdot \vec{B$$



nange due to momentum



The Maxwell equations

• At the same time we need to fulfill the Maxwell equations

$$\nabla \times \vec{B} = \frac{4\pi}{c}\vec{j} + \frac{1}{c}\frac{\partial E}{\partial t}$$
$$\nabla \cdot \vec{B} = 0$$

• With the charge density ρ and current j defined as

$$\rho = e \int (Z_i f_i - f_e) d^3 p$$







$$\nabla \cdot \overrightarrow{E} = 4\pi\rho$$

 $\vec{j} = e \left[(Z_i \vec{v}_i f_i - \vec{v}_e f_e) d^3 p \right]$

Modeling laser-plasma interaction The Particle-In-Cell Method





Modeling laser-plasma interaction The Particle-In-Cell Method



The PIC assumption is basically that particles initially close to another stay close to another



$f(\vec{x}, \vec{v}, t)$ 1 macro-particle represents 1000's of electrons / ions

























The Particle-In-Cell Method | Geometries





Cylindrical





Figures from FBPIC documentation

Modeling laser-plasma interaction The Particle-In-Cell Method | Reference frames





Laboratory frame (can be co-moving)



Modeling laser-plasma interaction The Particle-In-Cell Method | The Output

- Field data: Meshes of all fields, as well as the deposited density and currents
- Particle data: List of macro particles with their individual weights and the phase space values (position, momentum).
- We are talking about GBs of data. But often we are actually only interested in a single number that we can optimize.
- We do this via an objective function that reduces the entire distribution function to a single scalar number $O[f(\vec{x}, \vec{v}, t)] = z.$





Towards optimization



Choose new input \vec{x}



Towards optimization

PIC Input

 $\vec{\chi}$





Choose new input \vec{x}



Modeling laser-plasma accelerators

Which neural network architecture to choose?



• A. Döpp et al. Data-driven Science and Machine Learning Methods in Laser-Plasma Physics, High Power Laser Science and Engineering 11 55 (2023) | arXiv:2212.00026 (2022)



Modeling laser-plasma accelerators

The multilayer perceptron





1st Hidden

(30 neurons)

layer

2nd Hidden layer (30 neurons)



Output layer

(3 neurons) Prediction *y*

Charge Q_{pred}
 Median Energy \bar{E}_{pred} Energy spread ΔE_{pred}

Modeling laser-plasma accelerators The multilayer perceptron 1st Hidden layer (30 neurons) Input layer 20% dropout (15 values) Elaser λ_0 $\Delta\lambda$ Z_{foc} Zernike coefficients normalize ReLU activation



2nd Hidden layer (30 neurons)



Output layer

(3 neurons) Prediction *y*

Charge Q_{pred}
 Median Energy \bar{E}_{pred} Energy spread ΔE_{pred}

Loss function

 $\mathscr{C}_{1}(y_{pred}, y_{train}) = \sum \|y_{pred} - y_{train}\|$

$$\|Q_{pred} - Q_{train}\| + \|\bar{E}_{pred} - \bar{E}_{train}\| + \|\Delta E_{pred} - \Delta E_{train}\| + \|\Delta E_{pred} - \Delta E_{train}\|$$

Training (Backpropagation)



2nd Hidden layer (30 neurons)



Output layer

(3 neurons) Prediction *y*

Charge Q_{pred}
 Median Energy \bar{E}_{pred} Energy spread ΔE_{pred}

Loss function

 $\mathscr{C}_{1}(y_{pred}, y_{train}) = \sum \|y_{pred} - y_{train}\|$

$$= \|Q_{pred} - Q_{train}\| + \|\bar{E}_{pred} - \bar{E}_{train}\| + \|\bar{E}_{pred} - \bar{E}_{train}\| + \|\Delta E_{pred} - \Delta E_{train}\|$$

Training (Backpropagation)



Prior Distribution

- between points, capturing their correlations.

- y(x)0.5
 - 0.0
 - -0.5
- -1.0

 $\ell =$ $\lambda =$

y(x)

- -0.5

















• Irshad, F., Karsch, S., & Döpp, A. Multi-objective and multi-fidelity Bayesian optimization of laser-plasma acceleration. Phys. Rev. Research 5, 013063 (2023)





Modeling laser-plasma accelerators

Summarizing the different approaches





Choosing the right objective







Optimizing laser-plasma accelerators Choosing the right objective

- Say we want to optimize three electron beam parameters:
 - Charge Q (total charge, charge within FWHM, etc.)
 - Bandwidth (standard deviation $\sigma_{E'}$ median absolute deviation E_{MAD} , etc.)
 - Distance to a target energy $|E_{target} E|$ (using mean energy, median energy, peak energy, etc.)
- Choosing different metrics or weights for each objective changes the outcome in an a priori unknown way!
- How to solve this problem?





• Irshad, F., Karsch, S., & Döpp, A. Multi-objective and multi-fidelity Bayesian optimization of laser-plasma acceleration. Phys. Rev. Research 5, 013063 (2023)





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12/2	J.S.S.			
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1 Perlim	509		16	1 2 27 11540005641252 8 552415012210526
	508 MATILDE 229	4	13	1,3 37.11540996641253, -8.652415012219626
3 Way Point Restaur	ante 20	4,1	16	1 2 37 11108481832249 -8 672930593184326
4 Portofinos	622	4,2	18	1,4 37.1096076012194668.673146823866775
5 Bahia Beach Bar	602	4,2	28	2,3 37.10716980672824, -8.664919435513847
6 Tasca do Kiko	1105	4,6	22	1,8 37.106624656646474, -8.669308679690374
7 Munich	258	4,3	23	1,9 37.1105486992403, -8.677055276454404
8 Dom Vinho	66	4,5	25	2 37.110124919949605, -8.678738568043244
9 Marisqueira "O Pe	rceve" 531	4,6	28	2,2 37.1103292555888, -8.679428868607687
10 Avenida Restaurar	te 238	4,5	28	2,2 37.107751331422925, -8.675210473584627
11 Delhi Darbar India	n restaurant 1592	4,7	25	2 37.11005631516564, -8.676511807921088
12 Restaurante Italian	o Tradicional 6	4	31	2,4 37.10238156732768, -8.674098597903255
13 Restaurante Reis	1095	4,5	30	2,4 37.102100118183195, -8.673296645313904
14 Empanadas & Co	813	4,9	35	2,7 37.10118594915564, -8.676940858243556
15 Pomò - la pasta bio	bitaliana 546	4,5	32	2,5 37.10151921949556, -8.674465875996285
16 Goji Lounge Café	770	4,7	32	2,5 37.10112913445654, -8.673816704831587
17 Restaurante Chefe	Artur 200	4,6	29	2,3 37.10718924880119, -8.67721415085503
18 Don Sebastião	1711	4,3	30	2,4 37.10191899220491, -8.671770260655029
19 Império do Mar	919	3,9	33	2,6 37.10040313626079, -8.671166052702327
20 O Brito	31	4,6	40	3,1 37.097610926295324, -8.676978329408325
21 Casinha do Petisco	1494	4,7	32	2,5 37.10103289240326, -8.674086925149648
22 Restaurante Do Vi	lage 30	3,9	54	4,2 37.0893268911832, -8.672745936541522
23 Restaurante Onda	Norte 445	4,5	49	3,8 37.092670065508315, -8.673187088926209
24 Repolho Gastrobar	& Garrafeira 478	4,6	49	3,8 37.09245587585402, -8.673490131255601
25 Gato Pardo	760	4,4	50	3,8 37.092339508114776, -8.673569339208367
26 Alma Lusa	/10	4,6	49	3,7 37.09341069036754, -8.673941120173023
27 Asnoka Indian Tan	doori cuisine 396	4,6	49	3,7 37.093511400753044, -8.674017835514153
28 Palacio Da China	582	4,5	40	3,7 37.0940129780111, -8.074289890878971
29 Gwazi 20 Cantinho do Potiso	984	4,7	40	3,5 57.094033182190000, -8.072281103440300
31 Tasca lota Laros	1502	4,1		0 27 0963811/558107 -8 673687600726025
32 Maria Petisca	427	4,5		0 37 09722833886898 -8 675758265939718
33 Friends	427	4,7		0 37 09671488952284 -8 677378320069934
34 Pizzeria Bell'Itália	225	4,7		0 37 11373711219587 -8 66942132925955
35 Restaurant Emoce	an 38	5		0 37 11743303919973 -8 652298107857126
36 Gaivota Branca	1079	3.8		0 37 11256646601128 -8 657174428435493
37 Bar Quim	507	4.1		0 37 11832457767363 -8 645621928347584
	Sol Meia Praia Lagos 309	3.8		0 37 11832457767363 -8.64512538464262
				0 37.126865480443048.641923959940664

-8.64

-8.62

-8.60







Optimizing laser-plasma accelerators Multi-objective optimization | The conference dinner problem

Optimizing laser-plasma accelerators Multi-objective optimization | The conference dinner problem

PULSE

Optimizing laser-plasma accelerators Multi-objective optimization | PIC simulations

1. Irshad, F., Karsch, S., & Döpp, A. Leveraging trust for joint multi-objective and multi-fidelity optimization. Machine Learning: Science and Technology 5 (1), 015056 (2024) 2. Irshad, F., Karsch, S., & Döpp, A. Multi-objective and multi-fidelity Bayesian optimization of laser-plasma acceleration. Phys. Rev. Research 5, 013063 (2023)

Optimizing laser-plasma accelerators Multi-objective optimization | Experiments

Optimizing laser-plasma accelerators Multi-objective optimization | Experiments

Input space (8D)

- Jet focus & height,
- Blade focus & height,
- Dispersion (ϕ_2, ϕ_3, ϕ_4)
- Gas Pressure

Optimizing laser-plasma accelerators Multi-objective optimization | Experiments

- Inverse Optimization: Given a target (e.g. energy), we can search the model for most likely positions in input space.
- Tricky part: Inverse from 1D to 8D has no unique solution. Use Gaussian Mixture Network to predict tuning curves.
- Tuning also works in experiments!

Further reading

Review paper

- What to do with my data?
- What are established machine learning techniques?
- Which method is suitable for my application?
- Extensive review / tutorial paper (30+ pages) on data-driven science and machine learning methods in laser-plasma physics

• A. Döpp et al. Data-driven Science and Machine Learning Methods in Laser-Plasma Physics, High Power Laser Science and Engineering 11 55 (2023) | *arXiv:2212.00026* (2022)

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Further reading

Review paper

- What to do with my data?
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- Extensive review / tutorial paper (30+ pages) on data-driven science and machine learning methods in laser-plasma physics

• A. Döpp et al. Data-driven Science and Machine Learning Methods in Laser-Plasma Physics, High Power Laser Science and Engineering **11** 55 (2023) | *arXiv:2212.00026* (2022)

FIG. 2. Illustration of standard approaches to making predictive models in machine learning. The data was sampled from the function $y = x(1 + \sin x^2) + \epsilon$ with random Gaussian noise, ϵ , for which $\langle \epsilon^2 \rangle = 1$. The data has been fitted by a) nearest neighbour interpolation, b) cubic spline interpolation and d) Gaus

The sim tem, be it a tion, is to 1 that every as the near A straightfo tational rec with straig tion. Both ferentiable. tion of the tion IV). A training po quire the pi a certain d are continu tive. While in one-dime in multi-dir polation ap certainty or measureme as a part of

In some approaches For instance a polynomia mial are de i.e. minimi between th

FIG. 4. Sketch of a random forest, an architecture for regression or classification consisting of multiple decision trees, whose individual predictions are combined using into an ensemble prediction e.g. via majority voting or averaging.

FIG. 5. Example of gradient boosting with decision

in regression settings or entropy and information gain in a classification setting. At each decision point the data set is split and subsequently the metric is re-evaluated for the resulting groups, generating the next layer of decision nodes. This process is repeated until the leaves are reached. The more layers decision layers are used, called the depth of the tree, the more complex relationships can

FIG. 7. Real-world example of a multilayer perceptr consists of 15 input neurons, two hidden layers with 30 neuron The input is derived from parasitic laser diagnostics (laser p $\Delta \lambda$, longitudinal focus position z_{foc} and Zernike coefficients 20% of neurons drop out for regularization during training. evaluate the accuracy of the model, in this case using the m the loss function is then propagated back through the netwo median energy (\overline{E}) and (c) measured and predicted energy b-c adapted from Kirchen et al.29

model incorporating a trained neural network was used to provide an additional computation package to the Geant4 particle physics platform. Neural networks are also trained to assist hohlraum design for ICF experiments by predicting the time evolution of the radiation temperature, in the recent work by McClarren *et al.*¹¹². In the work by Simpson et al.¹¹³, a fully-connected neural network with three hidden layers is constructed to assist the analysis of a x-ray spectrometer, which measures the x-rays driven by MeV electrons produced from high-power laser-solid interaction.

7. Physics-informed machine learning models

The ultimate application of machine learning for modeling physics systems would arguably be to create an "artificial intelligence physicist", as coined by Wu and Tegmark¹¹⁴. One prominent idea at the backbone of how

train a deep neural network. An example of using decision tree as an initializer are Deep Jointly-Informed Neural Networks (DJINN) developed by Humbird et al.⁹⁵, which have been widely applied in the high power laser community, especially ertial confinement fusion datasets The algorithm first constructs a tree or a random forest with tree depth set as a tunable hyperparameter. It then maps the tree to a neural network, or maps the forest to an ensemble of networks. The structure of the network (number of neurons and hidden layer, initial weights, etc.) reflects the structure of the tree. The neural network is then trained using back-propagation. The use of decision trees for initialization largely reduces the computational cost while maintaining comparable performance to optimized neural network architectures. The DJINN algorithm has been applied to several classification and regression tasks

Author, Year	Laser type	Optimization Method(s)	Free Parameters	Optimizati
He et al., 2015 ¹⁹⁶	800 nm Ti:Sa, 15 mJ, 35 fs, 0.5 kHz	Genetic algorithm	deformable mir- ror (37 actuator voltages)	Electron file, energ & transver
Dann et al., 2019 ¹⁹⁷	800 nm Ti:Sa, 450 mJ, 40 fs, 5 Hz	Genetic & Nelder-Mead algorithms	deformable mirror or acousto-optic programmable dispersive filter	Electron bea
Shalloo et al., 2020^{198}	800 nm Ti:Sa, 0.245 J, 45 fs (bandwidth limit), 1 Hz	Bayesian optimization	Gas cell flow rate & length, laser dispersion $(\partial_{\omega}^2 \phi, \partial_{\omega}^3 \phi, \partial_{\omega}^4 \phi)$, focus position	Total electro Electron cha ceptance ang ray counts
Jalas et al., 2021 ¹⁹⁹	800 nm Ti:Sa, 2.6 J, 39 fs, 1 Hz	Bayesian optimization	Gas cell flow rates $(H_2 \text{ front and back}, N_2)$; focus position and laser energy	Spectral cha

TABLE I. Summary of a few representative papers on machine-learning-aided optimization in the context of laser-plasma acceleration and high-power laser experiments

distributions, in this case the electron energy distribution. While simple at the first glance, these objectives need to be properly defined and there are often different ways to do so^{201} . In the example above, energy and bandwidth are examples for the central tendency and the statistical dispersion of the energy distribution, respectively. These can be measured using different metrics such as weighted arithmetic or truncated mean, the median, mode, percentiles and so forth for the former; and full width at half maximum, median absolute deviation, standard deviation, maximum deviation, etc. for the latter. Each of these seemingly similar measures emphasises different features of the distribution they are calculated from, which can affect the outcome of optimization tasks. Sometimes one might also want to include higher order momenta as objectives, such as the skewness, or use integrals, e.g. the total beam charge.

2. Pareto optimization

multiple sometimes competing objectives g_i . As the objective function should only yield a single scalar value, one has to condense these objectives in a process known as scalarization. Scalarization can for instance take the form of a weighted product $q = \prod q_i^{\alpha_i}$ or sum $q = \sum \alpha_i q_i$ of the individual objectives q_i with the hyperparameters α_i describing its weight. Another common scalarization technique is ϵ -constraint scalarization, where one seeks to reformulate the problem of optimizing multiple objectives into a problem of single-objective optimization conditioned on constraints. In this method the goal is to optimize one of the q_i given some bounds on the other objectives. All of these techniques introduce some explicit bias in the optimization which may not necessarily repre-

FIG. 12. Pareto front. Illustration how a multi-objectiv function f(x) = y acts on a two-dimensional input space $x = (x_1, x_2)$ and transforms it to the objective space y = (y_1, y_2) on the right. The entirety of possible input positions is uniquely color-coded on the left and the resulting position in the objective space is shown in the same color on the right imal solutions form the Pareto front, indicate The Pareto-opt on the right, whereas the corresponding set of coordinates in the input space is called the Pareto set. Note that both Pareto front and Pareto set may be continuously defined locally, but can also contain discontinuities when local maxima get involved. Adapted from Irshad et al.²⁰²

sent the desired outcome. Because of this, the hyperparameters of the scalarization may have to be optimized themselves by running optimizations several times.

the entire vector of individual objectives $g = (g_1, \ldots, g_N)$ is optimized. To do so, instead of optimizing individual objectives, it is based on the concept of dominance. A

ion goal angular pro21

v distribution se emittance e compression am charge, total in energy range, m divergence

on beam energy. arge within ac gle, Betatron X-

arge density

A more general approach is Pareto optimization, where

Thank you for your attention!

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