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Introduction to Laser Wakefield Acceleration

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A European Research Infrastructure Consortium



Rationale

Electron accelerators are a cornerstone technology of modern society. In their variations, they are daily used as an instrument to enable scientific explorations, and as strategic tools at the disposal of the healthcare system for medical imaging and cancer treatment.

Electron beams are produced by radio-frequency machines, that accelerate up to 100 MV/m (typ. 10 MV/m). This becomes a limit for GeV accelerators, which are the key technology for last generation light source (synchrotron radiation and free electron lasers).

Laser Wakefield acceleration is the most compact technique to generate GeV electron beams. They also have other unique features.





Tajima and Dawson's paper

VOLUME 43, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JULY 1979

Laser Electron Accelerator

T. Tajima and J. M. Dawson Department of Physics, University of California, Los Angeles, California 90024 (Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density 10^{18} W/cm² shone on plasmas of densities 10^{18} cm⁻³ can yield gigaelectronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.



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Electron injection in the plasma wave

In LWFA both the "injector" and the "accelerating cavity" are re-created at every laser "shot". It is not hard to inject electrons in the plasma wave in a not-controlled way (self-injection). But to get high quality beams, the electron injection inside the accelerating plasma cavity must be controlled.





Catching a wave



Pre-acceleration required to catch the wave



They will oscillate and stay in the same position



Electron injection schemes

We need to help the electrons catch the wave. And we need to do it trying to keep the plasma wave formation process not affected. We have two ingredients: the laser and the plasma target.

Optical injection schemes

- Ionization injection
- Colliding pulse injection

Plasma injection schemes

- Downramp/shock injection
- Density bump



And their combination(s): i.e. shock-assisted ionization injection (Thaury, Sci. Rep. 2015)



Ionization injection







Colliding Pulse Injection





Downramp/Shock injection



Bulanov PRE 1998, Schmid PR STAB 2010, Grafenstein Sci. Rep. 2023



Density bump injection





Wang PRL 2016



Shock-assisted ionization injection





Thaury Sci. Rep. 2015

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LWFA Fundamental Limits: laser diffraction

In LWFA the accelerating cavity is generated in the laser focus. The electron beam energy gain is given by:

 $\Delta E_e[GeV] = V[GV/cm] \cdot L_{acc}[cm]$

Therefore, the strong accelerating gradient is not enough to get to GeV energies if we do not extend the region where the laser is focused as much as we can.

w_r(z_r) w_o w_o w(z) w_r(z_r) w_o k (z) intensity intensity intensity Radial Position Radial Position

Self-guiding

$$P_{c} = 17.5 \cdot \frac{n_{cr}}{n_{p}} GW = 17.5 \cdot \frac{\lambda_{p}^{2} [\mu m]}{\lambda_{L}^{2} [\mu m]} GW = 17.5 \cdot \frac{\lambda_{p}^{2} [\mu m]}{0.64} GW$$

Pre-formed plasma waveguide



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Ralph PRL 2009



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Capillary Discharge



Plasma waveguide

Bessel Beam



Leemans PRL 2014, Goncalves PRL 2019, Miao PRX 2022



Dephasing: the acceleration is over





LWFA Fundamental Limits: dephasing

The light inside a plasma travels slower than in vacuum:

$$v = c \cdot n = c \cdot \left(1 - \frac{n_e}{n_c}\right)^{1/2}$$

This implies a hard limit on the maximum speed of the accelerating cavity, and on the maximum energy that the electron beam can gain. From this fundamental argument, it is possible to derive the dephasing length for a laser wakefield accelerator:

$$L_d[cm] \sim \frac{3.9}{\lambda_L^2[\mu m] n_e [10^{18} cm^{-3}]}$$

Typically, single stage LWFA energy gain is limited by dephasing for laser energy of J-level and more. The dephasing length sets the length of the gas target (it has to be taken into account also at which longitudinal position injection is expected to happen).





Highest Energy LWFA



Goncalves PRL 2019, Aniculaesei MRE 2024



Demonstration of LWFA robustness





Highest Rep Rate LWFA - kHz





Guenot Nat. Phot. 2017, Salehi PRX 2021, Lazzarini PoP 2024 ²⁰



VERY GOOD!

YOU NOW KNOW ALL THE DETAILS OF LWFA!

IS TIME TO GO TO THE LAB!

Scaling laws and comparison with rf beams

$$E_e[GeV] \sim 1.7 \left(\frac{P[TW]}{100}\right)^{\frac{1}{3}} \left(\frac{1}{n_p[10^{18}cm^{-3}]}\right)^{\frac{1}{3}} \left(\frac{0.8}{\lambda_0[\mu m]}\right)^{\frac{4}{3}}$$

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Lu PR STAB 2006, Bulanov J. Plasma Phys 2016





The real size of a LWFA machine

Visit Visit Vi		Off-axis Parabola	a focal length	Laser pulse Gas Target leng what we refer	Gas target → gth (typicall to)
Parameter	10s mJ - TW	J – 100 TW	PW	10 PW	-
L _{laser}	1 m	3 m	10 m	30 m	
OAP length	0.1 m	1 m	10 m	100 m	
L _{acc}	0.0001 m	0.01 m	1 m	10 m	
Electron Energy	10s MeV	100s MeV	Multi-GeV	10s GeV	
Total	~ 1 m	~ 5 m	~ 20 m	~ 150 m	
Label	Table-top	University Lab	Research Lab	National Lab	24





The laser energy in the focal spot E_{FS} is estimated as follows:

$$E_{FS} = E_{fb} \epsilon_{sb} \epsilon_{wf}$$

where, E_{fb} is the laser energy in the full beam at the compressor output, ϵ_{sb} is the fraction of the energy that the laser beam can put into the focal spot, and ϵ_{wf} is the fraction of the laser energy that can go into the focal spot due to non-ideal wavefront. E.g. a square beam sets $\epsilon_{sb} = 0.85$, and ϵ_{wf} is typically in the 0.2 – 0.7 range.



Laser focusing



Full Power Characterization



How to measure laser focus?







Gas targets







Good density profile





Gas targets







fs probe beam



Imaging the accelerating cavity

Laser probe beam



Electron probe beam



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Savert PRL 2015, Wan Science Advences 2024, Lorenz 2024



Electron Beam Characterization





Screen	Absolute calibration $(10^9 \text{ photons/sr/pC})$	$\frac{N_{scint}/N_{CLS,20\ ms}/Q}{(pC^{-1})}$	$ \rho_{\text{sat}} $ (see Sec. III C) (pC/mm ²)
KODAK Biomax MS	14.8 ± 1.3	5.79 ± 0.26	21.8 ± 5.0
CAWO OG 16	12.4 ± 1.1	4.86 ± 0.21	32.9 ± 6.6
KODAK Biomax Transcreen HE	7.85 ± 0.67	3.02 ± 0.13	47 ± 10
KODAK Lanex Regular	6.95 ± 0.60	2.72 ± 0.12	66 ± 33
KONICA KR	6.58 ± 0.56	2.58 ± 0.11	>100
KODAK Biomax Transcreen LE	1.79 ± 0.15	0.700 ± 0.031	>100
KODAK Lanex Fine	1.75 ± 0.15	0.686 ± 0.030	>100
KONICA KF	1.54 ± 0.13	0.602 ± 0.027	>100



LWFA Applications

Compact and tunable accelerators

- Compact Free Electron Laser
- Compact and tunable light source for medical imaging and non-destructive testing
- Compact VHEE radiotherapy device
- Compact muon source

Enabling new science

- Laser Electron Colliders
- Ultra-fast radiobiology
- Dual electron beams



Laser-Driven Compact Free Electron Laser



400

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I. 13

н

н

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т

500

19





Compact X-ray sources





Ta Phuoc Nat. Phot. 2012, Khrennikov PRL 2015



Laser-driven VHEE radiotherapy device





Labate Sci. Rep. 2020, Svendsen Sci. Rep. 2021



Bethe-Heitler Pair Production: Create bremsstrahlung radiation by impacting electrons onto a high atomic number target and produce muon pairs



Compact Muon Source

Discovery of a hidden chamber in the Great Pyramid of Giza using Cosmic-ray Muons



Morishima Nature 2017

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Laser – Electron Colliders







10-12

Physical

Excitation

Ionization

Dissociation

Free -radical

reactions

10-18

Ultra-fast Radiobiology

20

16

[12 2] 8 8

0

20

16

12 18 Z [cm]

(e) E = 250 MeV

0

6

24 30

0

6

6

Time to reach 99% of the dose

12 18 Z [cm]

Time to reach 99% of the dose

12 18 Z [cm]

(b) E = 50 MeV

24

24

30

30

200

150 [sd] 100 amjr 50

200

150 [sd] 100 amil 50





Gauduel Eur. J. Phys 2010, Gauduel J. Phys. 2011, Horvath Phys. Med. Biol. 2023



Special Electron Beams





Wenz Nat. Phot. 2019, Pollock PRL 2015, Valenta PoP 2021



THANK YOU FOR YOUR ATTENTION!

