

# Thermal one-point blocks

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# Motivation

## Study Conformal Field Theories

- Four-point functions
- Multipoints
- Defects
- **Thermal CFTs:**

Universal properties (e.g. at high  $\Delta$ )

CFTs on other manifolds

# Outline

- Casimir for spinning representations on  $S^1 \times S^{d-1}$
- Solve it with expansion in  $T$  in  $d = 3 \rightarrow$  one-point thermal blocks
- Computations of OPE coefficients (holography)

# Different notion of temperature $T$

## Classical statistical theory

- Relevant coupling in effective action
- Description of second-order phase transitions at critical temperature  $T_c$
- Changing  $T$  means moving away from the CFT

## QFT at finite temperature

- Compactified dimension  $T \in S^1$
- Non-zero temperature physics of quantum critical points
- The only scale is  $T$  (or chemical potentials)

# Intro to thermal CFTs

- Since conformal invariance is broken,

$$\langle \phi \rangle_\beta \neq 0 \quad \text{for } \Delta \neq 0.$$

- For general thermal QFTs,

$$\langle \phi \rangle_\beta = \frac{1}{\mathcal{Z}} \text{tr}_{\mathcal{H}}(\phi e^{-\beta \mathcal{H}}) = \frac{1}{\mathcal{Z}} \sum_n \langle n | \phi e^{-\beta \mathcal{H}} | n \rangle.$$

	$S^1 \times \mathbb{R}^{d-1}$	$S^1 \times S^{d-1}$
$\langle \phi \rangle$	Fixed	$\sum_{\mathcal{O}} \lambda_{\phi\phi\mathcal{O}} \times \text{Blocks}$
$\langle \phi\phi \rangle$	$\sum_{\mathcal{O}} \lambda_{\phi\phi\mathcal{O}} \times \text{Blocks}$	...

# High-temperature limit

CFT correlators on  $S^1 \times S^{d-1} \xrightarrow{T \rightarrow \infty}$  correlators on  $S^1 \times \mathbb{R}^{d-1}$

Invariance under  $SO(d-1)$  implies: [Iliesiu, Kologlu, Mahajan, Perlmutter, Simmons-Duffin '18]

- Only STT representations
- By symmetry and dimensional analysis

$$\langle \phi^{\mu_1 \dots \mu_l}(x) \rangle_\beta \xrightarrow{T \rightarrow \infty} \frac{b_\phi}{\beta^\Delta} (e^{\mu_1} \dots e^{\mu_l} - \text{traces}).$$

## Definition of correlator

- Consider CFT on cylinder  $\mathbb{R} \times S^{d-1}$
- One-point function at finite  $T$  (on  $S^1 \times S^{d-1}$ ) as

$$G_1(x, \beta) = \langle \phi_{\text{cyl}}(x) \rangle_{\beta} = \text{tr}_{\mathcal{H}}(\phi_{\text{cyl}}(x) e^{-\beta D})$$

- We want non-zero **chemical potentials**,

$$G_1(x, \beta, \mu_i) = \langle \phi_{\text{cyl}}(x) \rangle_{\beta, \mu_i} = \text{tr}_{\mathcal{H}}(\phi_{\text{cyl}}(x) e^{-\beta D} e^{-\mu_1 H_1} \dots e^{-\mu_{d-2} H_{d-2}}),$$

with  $H_1, \dots, H_{d-2}$  Cartan generators of rotations.

We will denote  $q = e^{-\beta}$  and  $y_i = e^{-\mu_i}$ .

# Thermal blocks

Definition of blocks with projectors,

$$\lambda_{\mathcal{O}\mathcal{O}\phi} G_{\mathcal{O}}(x) = \text{tr}(P_{\mathcal{O}}\phi(x)q^D y_1^{H_1} \dots y_{d-2}^{H_{d-2}}).$$

Relation with **OPE coefficients** with shadow formalism,

$$\lambda_{\mathcal{O}\mathcal{O}\phi} G_{\mathcal{O}}(x) = q^{\Delta_{\phi}} y_1^{J_1} \dots y_{d-2}^{J_{d-2}} \int d^d x' \langle \tilde{\mathcal{O}}(x')\phi(x)\mathcal{O}(x'_{q,y}) \rangle.$$

Without chemical potentials  $\mu_i$ :

- Lost most of information about OPE coefficients,
- No clear way to obtain Casimir.



# Casimir with internal and external scalars

Fix  $d = 3 \rightarrow$  only STT operators and one chemical potential.

[Gobeil, Maloney, Ng, Wu '18]

$$C_2 = q^2 \partial_q^2 + y^2 \partial_y^2 + q \partial_q + y \partial_y + \frac{(1+y)y \partial_y}{y-1} + \frac{(1+q)q \partial_q}{q-1} + \frac{(q+y)(q \partial_q - y \partial_y)}{q-y} + \frac{(1+qy)(q \partial_q + y \partial_y)}{qy-1} + \mathcal{F}_1(q, y, p) \mathcal{F}_2(\partial_p^2, \partial_p),$$

with  $\mathcal{F}_1$  rational function and  $\mathcal{F}_2$  linear.

## Ansatz

$$\sum_{a,b,c} f_{a,b,c} q^a y^b p^c \quad \text{with } a \geq 0, b \in [-a, a], c \in [0, 2a].$$

# Spherical functions

- How to deal with general spinning representations?
- Spherical functions over  $(G, K)$  with  $K$  subgroup of  $G$
- Blocks  $\longleftrightarrow$  spherical functions  $\implies$  tools from Harmonic analysis

	$G$	$K$
Conformal blocks <small>[Schomerus, Sobko, Isachenkov '16] [Burić, Schomerus '22]</small>	$SO(d+1, 1)$	$SO(1, 1) \times SO(d)$
Partial waves <small>[Burić, FR, Vichi '23]</small>	$SO(d-1)$	$SO(d-2)$
Thermal blocks	$SO(d+1, 1) \times SO(d+1, 1)$	$SO(d+1, 1)$

# Spinning Casimir

- **Radial component map**  $\rightarrow$  reducing Casimir to four-variable  $(q, y, p, z)$  differential equation:

$$\begin{aligned} C_2 = & q^2 \partial_q^2 + y^2 \partial_y^2 + q \partial_q + y \partial_y + \frac{(1+y)y \partial_y}{y-1} + \frac{(1+q)q \partial_q}{q-1} + \\ & \frac{(q+y)(q \partial_q - y \partial_y)}{q-y} + \frac{(1+qy)(q \partial_q + y \partial_y)}{qy-1} + \\ & + \mathcal{F}_1(q, y, p, z, \sqrt{1+z^2}, \sqrt{1-p^2}) \mathcal{F}_2(\partial_p^2, \partial_p, \partial_z^2, \partial_z, \partial_p \partial_z), \end{aligned}$$

with  $\mathcal{F}_1$  rational function.

## Weight-shifting operator

$$\hat{q}_{\Delta, l} C_2(\Delta, l) = C_2(\Delta - 1, l + 1) \hat{q}_{\Delta, l}$$

## Spinning ansatz

$$\sum_{a,b,c} f_{a,b,c,d} q^a y^b p^c z^d (1 + \sqrt{1 - p^2})(1 + \sqrt{1 + z^2}), \quad \text{with } a \geq 0, \\ b \in [-a - l_{\mathcal{O}}, a + l_{\mathcal{O}}], c \in [0, 2(a + l_{\phi} + l_{\mathcal{O}})], d \in [0, l_{\phi}].$$

## Example

Order  $q$  for  $l_{\phi} = l_{\mathcal{O}} = 1$  (and  $\Delta_{\phi} = \Delta_{\mathcal{O}} = 3$ ),

$$\frac{1 - y^2}{y} + \frac{q(1 - y)}{16y^2} \left( 9p^2(y - 1)^2(y + 1) + 2(1 + y)(5 + y(14 + 5y)) \right. \\ \left. + 3p\sqrt{1 - p^2}(y - 1)(2(y^2 - 1)z + (2 + y)(1 + 2y)\sqrt{1 + z^2}) \right) + o(q)$$

# Potential applications involving $\langle T \rangle_{\beta, \mu}$

## General CFTs

From [Iliesiu, Kologlu, Mahajan, Perlmutter, Simmons-Duffin '18],

$$\langle T^{00} \rangle_{\beta} = \frac{1}{\text{vol}(S^{d-1})} \frac{\partial}{\partial \beta} \log \mathcal{Z}, \quad \mathcal{Z} = \sum_{\mathcal{O}} \chi_{\mathcal{O}}(e^{-\beta}).$$

- $T \rightarrow \infty$
- Knowledge of spectrum  $\rightarrow b_T$ .

Blocks for  $\langle T \rangle_{\beta, \mu}$  and spectrum  $\rightarrow$  OPE coefficients.

## Holography

- Consider a AdS Kerr black hole in the bulk

$$(m, J) \rightarrow (\Omega, T) \sim (\mu, T)$$

- Asymptotic expansion of the bulk metric [de Haro, Skenderis, Solodukhin '00]
- Compute  $\langle T \rangle_{\beta, \mu}$  [Cardoso, Dias, Hartnett, Lehner, Santos '13]
- One-point block expansion for the stress tensor
- OPE coefficients for the dual CFT.

# Conclusions

## Results

- Casimir for general spinning representations
- Solving perturbatively in  $q$
- Thermal blocks for one-point functions

## Future plans

- Computation of OPE coefficients both in holographic setting or not
- Higher-points: Crossing equations? Positivity?

Thanks!