Bootstrap meets Experimental Data

Ning Su Caltech & MIT

19/02/2024 @ 50+ ϵ years of conformal bootstrap

Today's talk

- ™ Review of numerical techniques for positivity bootstrap
- * Application in condensed matter physics : Helium superfluid phase transition

W List of unsolved problems in	* +	_
→ C 🖙 en.wikipedia.c	org/wiki/List_of_unsolved_problems_in_physics#Condensed_matter_physics	*
мистеат ризнез	While compressive mechanical stress is known to encourage whisker formation, the g mechanism has yet to be determined.	rowth
Fluid dynamics	 Superfluid transition in helium-4: Explain the discrepancy between the experimental^{[8} theoretical^[84][85][86] determinations of the heat capacity critical exponent a.^[87] 	^{3]} and
physics	Scharnhorst effect: Can light signals travel slightly faster than <i>c</i> between two closely s conducting plates purchase the Conjugate of the state of the	e whisker formation, the growth etween the experimental ^[83] and tical exponent α . ^[87] an <i>c</i> between two closely spaced
Quantum computing	conducting plates, exploring the casimir effectives	

[™] Application in particle physics : pion-pion scattering



Numerical techniques

1, Cutting surface algorithm [Chester, Landry, Liu, Poland, Simmons-Duffin, SN, Vichi 2019]

Scan over
$$\left\{ \frac{\lambda_{\phi\phi t}}{\lambda_{\phi\phi s}}, \frac{\lambda_{tts}}{\lambda_{\phi\phi s}}, \frac{\lambda_{sss}}{\lambda_{\phi\phi s}} \right\}$$
 with cost ~ dimension

2, Navigator function [Reehorst, Rychkov, Simmons-Duffin, Sirois, SN, van Rees 2021] SDP compution → sign of the objective indicating feasible / infeasible



Numerical techniques

3, Skydiving algorithm [Aike Liu, David Simmons-Duffin, NS, Balt van Rees, 2023] Treat two optimizations (SDP optimization, external parameter search) as a single optimization problem.



See [Rychkov, NS 2023] for a review of numerical conformal bootstrap

Today's talk

* Application in condensed matter physics : Helium superfluid phase transition

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Quantum computing	conducting plates, exploiting the Casimir effect? ^[88]	

[™] Application in particle physics : pion-pion scattering



Helium superfluid phase transition



Bootstrap Helium superfluid phase transition



[Chester, Landry, Liu, Poland, Simmons-Duffin, SN, Vichi 2019]

What's wrong in the experimental result?



Experimental error bar in global significance

The 95% confidence range : $\alpha \in (-0.031, -0.008)$ Bootstrap: -0.0152(3)



[Landry, Liu, Poland, Simmons-Duffin, SN, to appear]

Today's talk * Application in condensed matter physics : Helium superfluid phase transition W List of unsolved problems in \times + \rightarrow C \simeq en.wikipedia.org/wiki/List_of_unsolved_problems_in_physics#Condensed_matter_physics ☆ While compressive mechanical stress is known to encourage whisker formation, the growth тисски рнузка mechanism has yet to be determined. • Superfluid transition in helium-4: Explain the discrepancy between the experimental^[83] and Fluid dynamics theoretical^{[84][85][86]} determinations of the heat capacity critical exponent α .^[87] Condensed matter • Scharnhorst effect: Can light signals travel slightly faster than c between two closely spaced physics conducting plates, exploiting the Casimir effect?[88] Quantum computing Application in particle physics : pion-pion scattering 1.0 40 60 80 100 20 -0.2 0.5 -0.4 200 300 400 500 -0.6 -0.5

pion-pion scattering : the S matrix

Unitary :

$$1 = \left\langle p_1 p_2 \left| S S^{\dagger} \right| p_3 p_4 \right\rangle = \sum_X \left\langle p_1 p_2 \left| S \right| X \right\rangle \left\langle X \left| S^{\dagger} \right| p_3 p_4 \right\rangle$$

$$|S_{2\rightarrow 2}|^2 = 1 - \sum_{X \in \text{rest}} |S_{2\rightarrow X}|^2 \implies |S_{2\rightarrow 2}|^2 \le 1$$

 $1 + \begin{pmatrix} -\operatorname{Im}[S_{2 \to 2}] & \operatorname{Re}[S_{2 \to 2}] \\ \operatorname{Re}[S_{2 \to 2}] & \operatorname{Im}[S_{2 \to 2}] \end{pmatrix} \ge 0$

Consistency condition for S matrix

$$\pi + \pi \rightarrow \pi + \pi$$
 S matrix: $S \equiv S_{2 \rightarrow 2} (s, t, u) = 1 + i T$, $s + t + u = 4$ (unit: $m_{\pi} = 1$)

Partial waves : $T(s, t, u) = \sum_{\ell} T_{\ell}(s) P_{\ell}(x)$ with $x = \cos(\theta) = 1 + \frac{2t}{s-4}$

Unitarity : $2 \operatorname{Im}[T_{\ell}] \ge |T_{\ell}|^2$ or $|S_{\ell}(s)|^2 \le 1$ for all ℓ ($S_{\ell} = 1 + i T_{\ell}$)

Analyticity : $S_{\ell}(s)$ is analytic, except real $s \ge 4$ and its image under crossing

Exact O(3) isospin symmetry : vector \otimes vector \rightarrow singlet \oplus antisym \oplus sym (0 \oplus 1 \oplus 2 isospin) $\langle \pi_i \pi_j | \pi_k \pi_l \rangle = \sum_{r=0,1,2} P_{ijkl}^{(r)} S^{(r)}(s, t, u) = \sum_{r=0,1,2} P_{ijkl}^{(r)} \sum_{\ell} S_{\ell}^{(r)}(s)$ (Example : $P_{ijkl}^{(r=0)} = \delta_{ij} \delta_{kl}$)

Crossing : S(s, t, u) has crossing symmetry

Translate to SDP : analyticity

[Paulos, Penedones, Toledo, van Rees, Vieira 1708.06765; Guerrieri, Penedones, Vieira 1810.12849] Analyticity :

S(s) is analytic in s complex plane except $[4, \infty) \implies S(s) = \sum a_n \rho_s^n$, $\rho_s = \frac{\sqrt{m^2 - 4} - \sqrt{4 - s}}{\sqrt{m^2 - 4} + \sqrt{4 - s}}$



Translate to SDP : isospin symmetry

 $\langle \pi_i \pi_j | \pi_k \pi_l \rangle = S(s | t, u) \,\delta_{ij} \,\delta_{kl} + S(t | s, u) \,\delta_{ik} \,\delta_{jl} + S(u | s, t) \,\delta_{il} \,\delta_{jk}$

Under $i \leftrightarrow j, k \leftrightarrow l \implies t \leftrightarrow u \implies S(s \mid t, u)$ is symmetric under $t \leftrightarrow u$

Example: $S(s | t, u) = \sum a_{nm} (\rho_t^n \rho_u^m + \rho_u^n \rho_t^m) + \sum b_{nm} (\rho_t^n + \rho_u^n) \rho_s^m$ (co

constraint :
$$s + t + u = 4$$
)

 $S^{(r)}$ is linear combination of $S(\Box \mid \Box, \Box)$

Example : $S^{(r=1)}(s, t, u) = S(t | s, u) - S(u | s, t)$

Translate to SDP

$$S_{\ell}^{(r)}(s) = \sum a_{n\,m}(...) + \sum b_{n\,m}(...) \qquad (...): \text{ functions in } s$$

$$1 + \begin{pmatrix} -\ln[S_{\ell}^{(r)}(s)] & \operatorname{Re}[S_{\ell}^{(r)}(s)] \\ \operatorname{Re}[S_{\ell}^{(r)}(s)] & \ln[S_{\ell}^{(r)}(s)] \end{pmatrix} \ge 0 \quad \text{for } r = 0, 1, 2; \ \ell = 0, 1, \ ...; s \in [4, \infty)$$

SDP : search for
$$\alpha_i$$
 such that $\sum_i \alpha_i M_i^{(j)} > 0$ for $j = 1, ..., J$

Found α_i : we have a candidate $S_{\ell}(s)$ that satisfies (1) unitarity; (2) analyticity; (3) crossing; (4) isospin symmetry

Bootstrap pion scattering : old result

[Guerrieri, Penedones, Vieira 1810.12849]



"Adler zeros" $s_0, s_2 : T_0(s_0) = 0$, $T_2(s_2) = 0$ ($S_\ell(s) = 1 + i T_\ell(s)$)

Shaded : Found $lpha_i$. White : Can't found $lpha_i$

Bootstrap pion scattering

Bootstrap is the right way to analyze data in particle scattering experiment

How people found meson (resonance) mass in the past?

A model + experimental constraints + some theory constraints (unitarity/crossing) \downarrow An analytic continuation of $S_{\ell}(s)$ to *s* complex plane \downarrow Find pole on the complex plane : $M = m - i \frac{\Gamma}{2}$

Problem : different continuation \rightarrow different mass \rightarrow biggest source of error in PDG

Pion scattering : old approaches

A model + experimental constraints + some theory constraints (unitarity/crossing)

Problem : different continuation \rightarrow different mass \rightarrow biggest source of error in PDG

Complex mass for σ resonance :



Statistical error bar may have uncontrollable systematic errors

Bootstrap + experimental constraints

[Kelian Haring, Andrea Guerrieri, NS, ongoing work]

Bootstrap + exp input : systematic analysis of experimental constraints + full theory constraints

 $S = a_{nm} \left(\rho_t^n \rho_u^m + \rho_u^n \rho_t^m \right) + \dots$

Find $a_{n,m}$ such that $|S| \leq 1$, crossing symmetric, AND matching experimental data

Experimental data : $\text{Arg}\left[S_{\ell=0}^{(r=0)}\right]$, $\text{Arg}\left[S_{\ell=1}^{(r=1)}\right]$, $\text{Arg}\left[S_{\ell=0}^{(r=2)}\right]$



Bootstrap + experimental constraints

New ideas:

1, Navigator function

Scan over parameters $(s_0, s_2, ...)$ such that the result matches experimental data

 $\mathsf{SDP}(s_0, s_2, \ \dots) \ \rightarrow \ \text{fitness score on} (s_0, s_2, \ \dots) \ \xrightarrow{\text{minimize score}} \text{best fit } w.r.t. (s_0, s_2, \ \dots)$

Bootstrap + experimental constraints

New ideas:

2, An objective that encoded experimental data

objective = $\sum_{s} (\text{Im}[S^{\text{ans}}(s)], \text{Re}[S^{\text{ans}}(s)]) . (\text{Sin}[\phi(s)], \text{Cos}[\phi(s)])$

SDP tries to matches experimental data as much as possible, subject to theory constraints

Bootstrap pion : Preliminary results



Bootstrap pion : Preliminary results



Bootstrap pion : Preliminary results

*• Observation of higher spin Regge trajectory



* Observation of isospin 2 spin 0 meson (must be at least tetra-quark or more quarks)



Error bar for sigma

Statistical error bar may have uncontrollable systematic errors



Bootstrap value : $m_{\sigma} = 3.09 + 1.12 i$ (Red) . What's if σ is away from 3.09 + 1.12 i ?

Error bar for sigma



"robust error bar" : points outside the error bar can't match experimental data

Verifying experimental data



Verifying experimental data



Matches with José R. Peláez 's analysis

The original dream of bootstrap

← → C 🙄 en.wikipedia.org/wiki/Bootstrap model

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coin greater than 1 and without the then undiscovered phenomenon of confinement, it is naively inconsistent with the rved Regge behavior of hadrons.

Chew and followers believed that it would be possible to use crossing symmetry and Regge behavior to formulate a consistent S-matrix for infinitely many particle types. The Regge hypothesis would determine the spectrum, crossing and analyticity would determine the scattering amplitude (the forces), while unitarity would determine the self-consistent quantum corrections in a way analogous to including loops. The only fully successful implementation of the program required another assumption to organize the mathematics of unitarity (the narrow resonance approximation). This meant that all the hadrons were stable particles in the first approximation, so that scattering and decays could be thought of as a perturbation. This allowed a bootstrap model with infinitely many particle types to be constructed like a field theory — the lowest order scattering amplitude should show Regge behavior and unitarity would determine the loop corrections order by order. This is how Gabriele Veneziano and many others constructed string theory, which remains the only theory constructed from general consistency conditions and mild assumptions on the spectrum.

Many in the bootstrap community believed that field theory, which was plagued by problems of definition, was fundamentally inconsistent at high energies. Some believed that there is only one consistent theory which requires infinitely many particle species and whose form can be found by consistency alone. This is nowadays known not to be true, since there are many theories which are nonperturbatively consistent, each with their own S-matrix. Without the narrow-resonance approximation, the bootstrap program did not have a clear expansion parameter, and the consistency equations were often complicated and unwieldy, so that the method had limited success. It fell out of favor with the rise of quantum chromodynamics, which described mesons and baryons in terms of elementary particles called quarks and gluons.