

Lattice QCD for TMD Physics: Nonperturbative Collins Soper Kernel

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Pavia, Italy



transverse momentum dependent (TMD) functions are ubiquitous

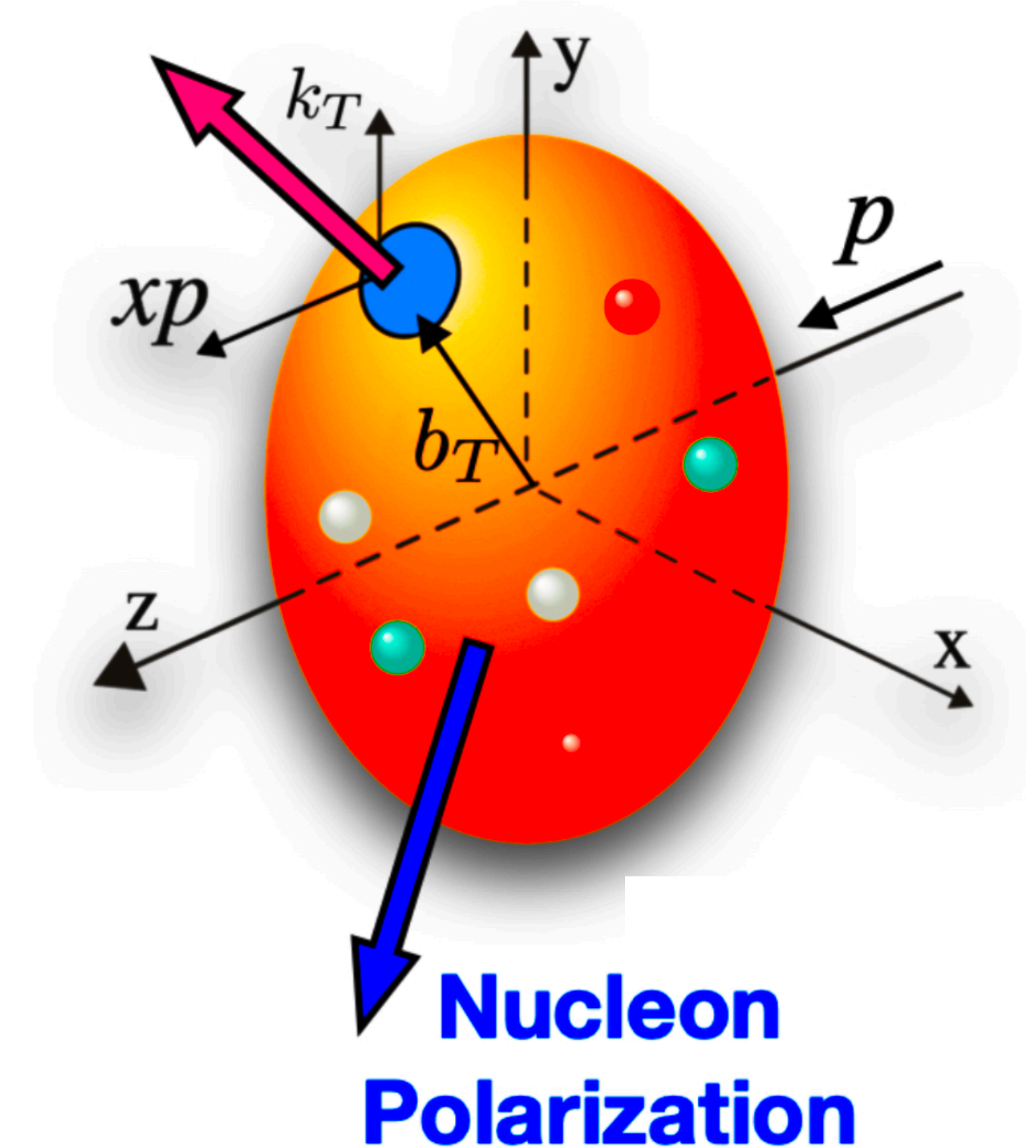
- ◆ TMD parton distribution function (TMDPDF)

$$\phi(x, k_{\perp}, \zeta, \mu)$$

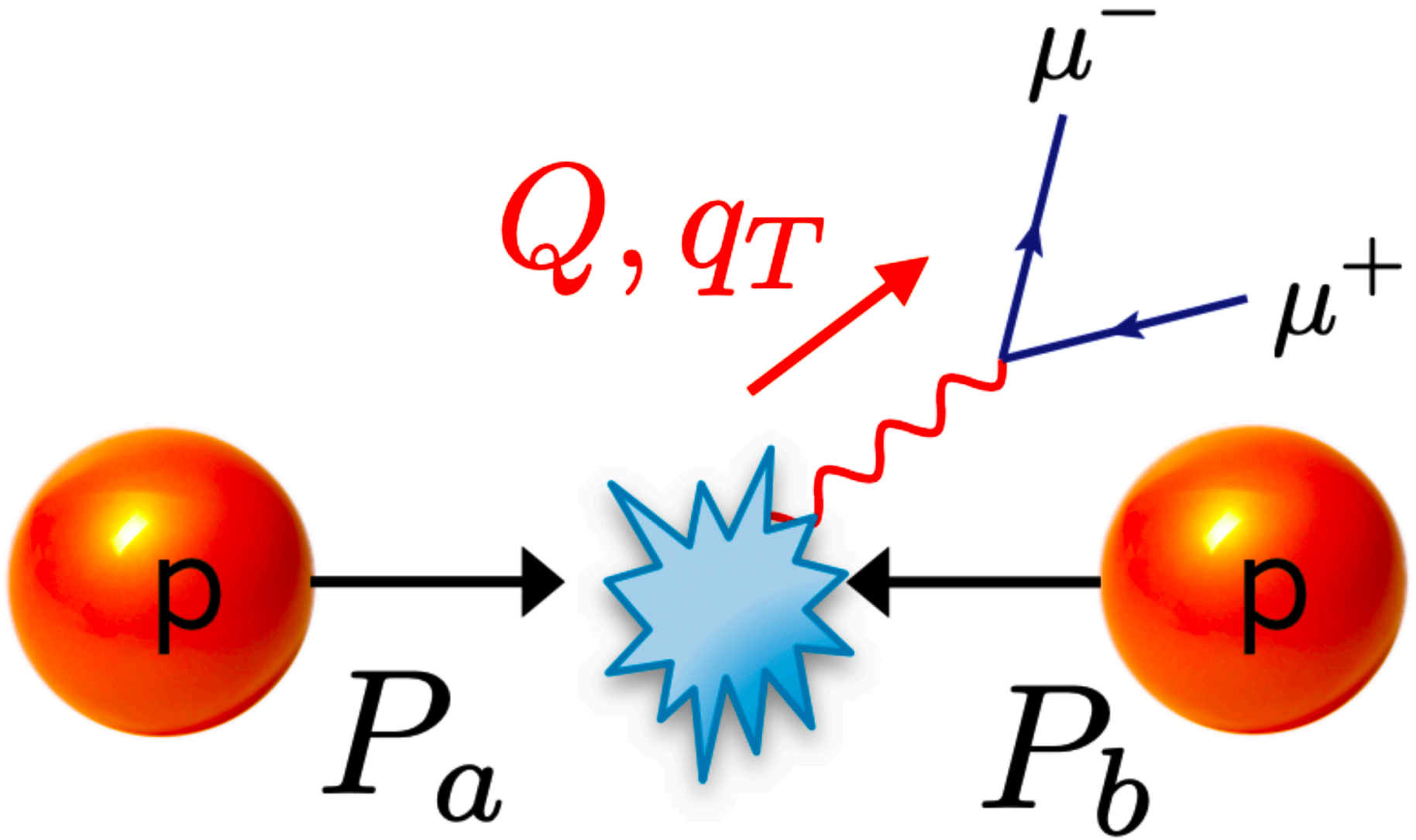
- ◆ TMD fragmentation function (TMDFF)

$$F(x, k_{\perp}, \zeta, \mu)$$

**Quark
Polarization**



transverse momentum dependent (TMD) functions are ubiquitous

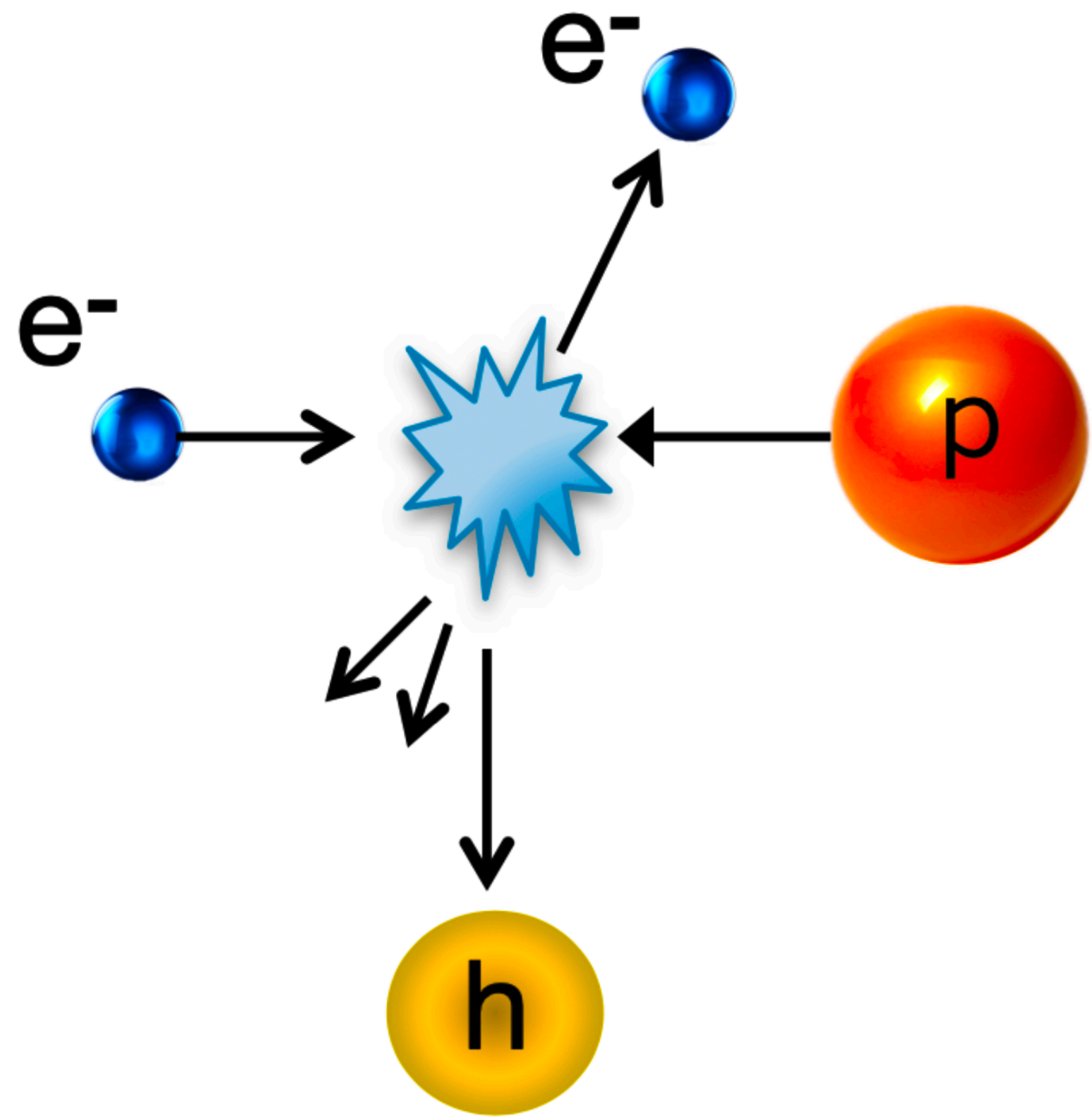


Drell-Yan processes

$$\sigma \sim \phi(x, k_{\perp}, \zeta, \mu) \otimes \phi(x, k_{\perp}, \zeta, \mu)$$

RHIC, LHC, ...

transverse momentum dependent (TMD) functions are ubiquitous



semi-inclusive deep inelastic scatterings

$$\sigma \sim F(x, k_{\perp}, \zeta, \mu) \otimes \phi(x, k_{\perp}, \zeta, \mu)$$

JLAB, EIC, ...

evolution of TMD functions across collision energies ...

Collins Soper (CS) kernel

$$\gamma^{\overline{\text{MS}}}(b_{\perp}, \mu) = \frac{\partial \phi(x, b_{\perp}, \zeta, \mu)}{\partial \ln \sqrt{\zeta}}$$

UV property of QCD— independent of hadronic state

JLAB

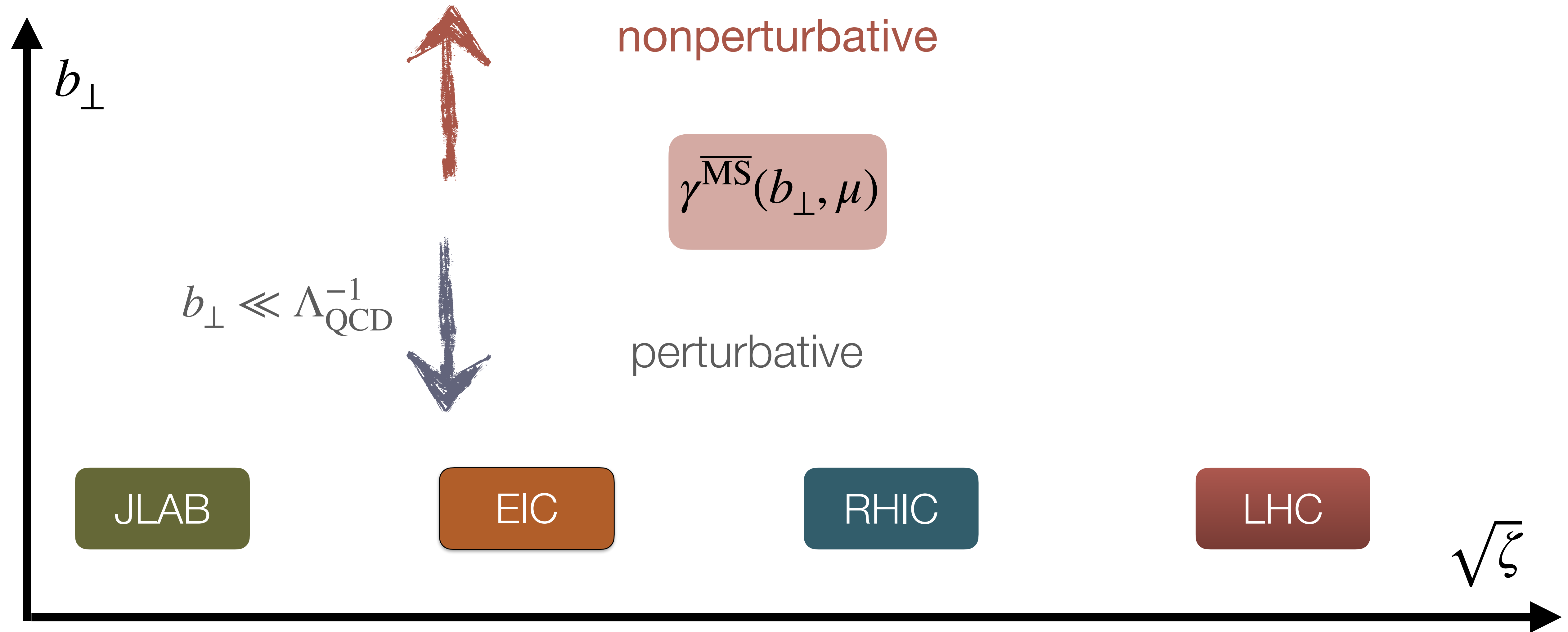
EIC

RHIC

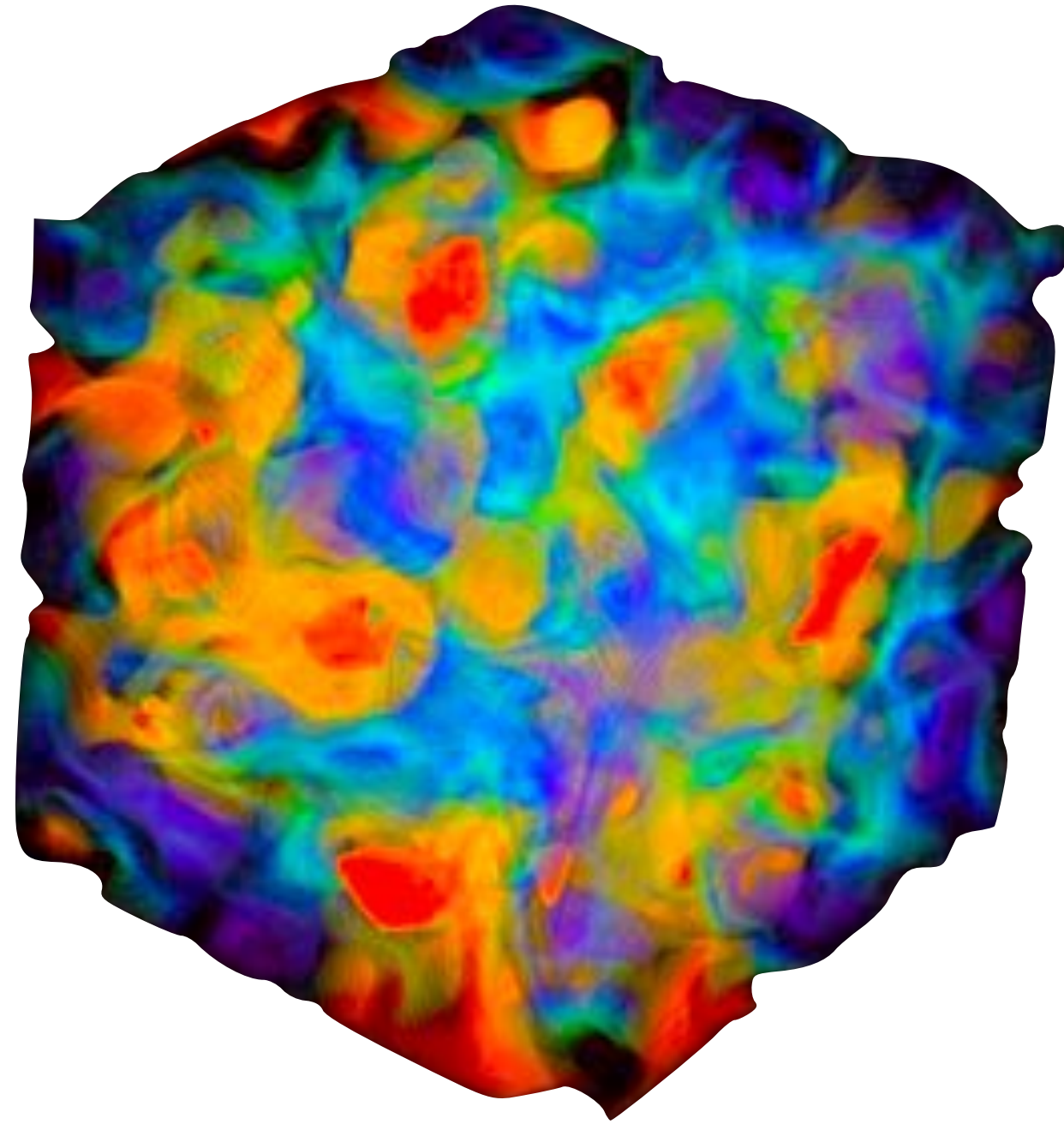
LHC

$\sqrt{\zeta}$

nonperturbative CS kernel



our dream ...



lattice QCD

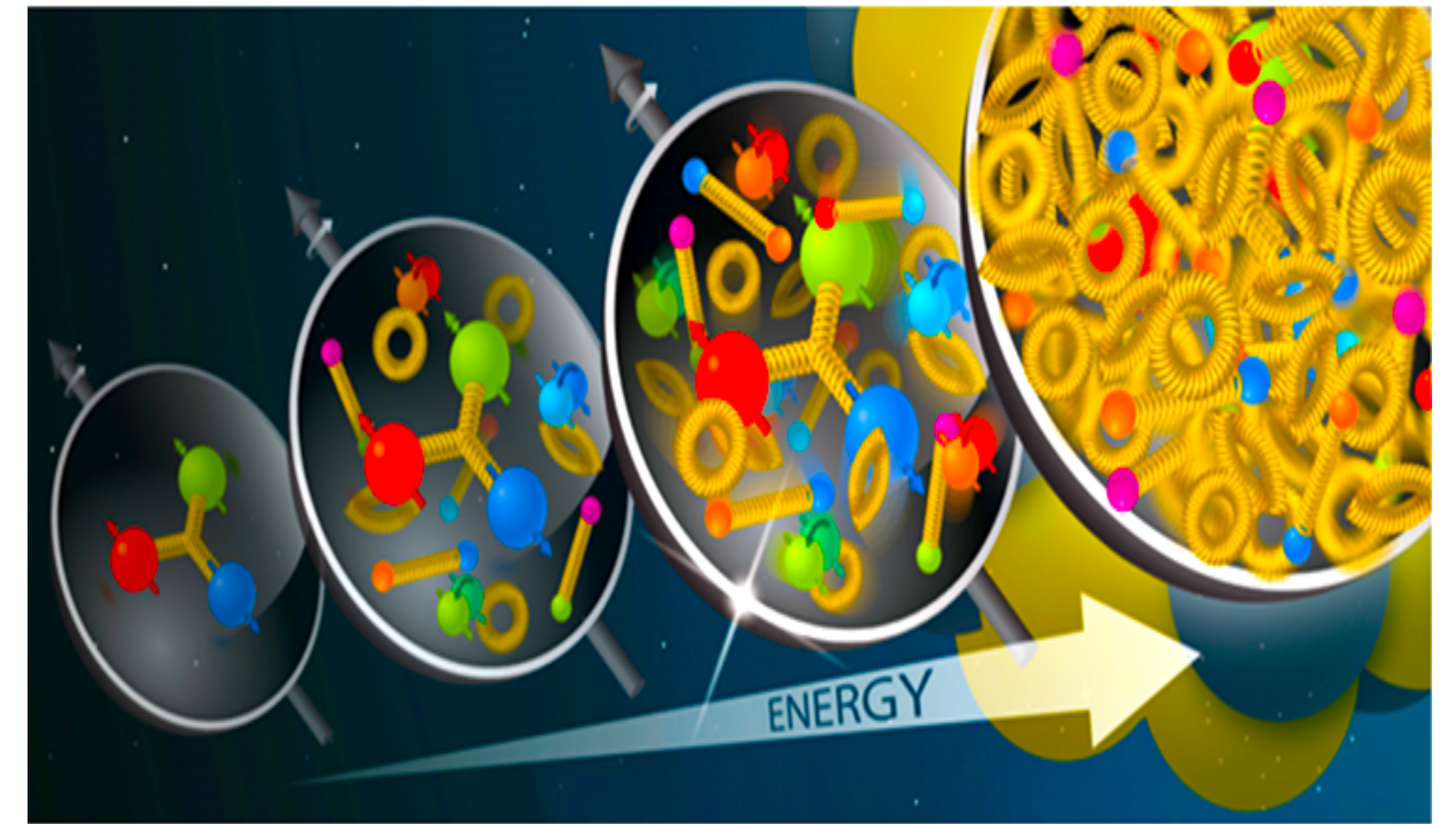
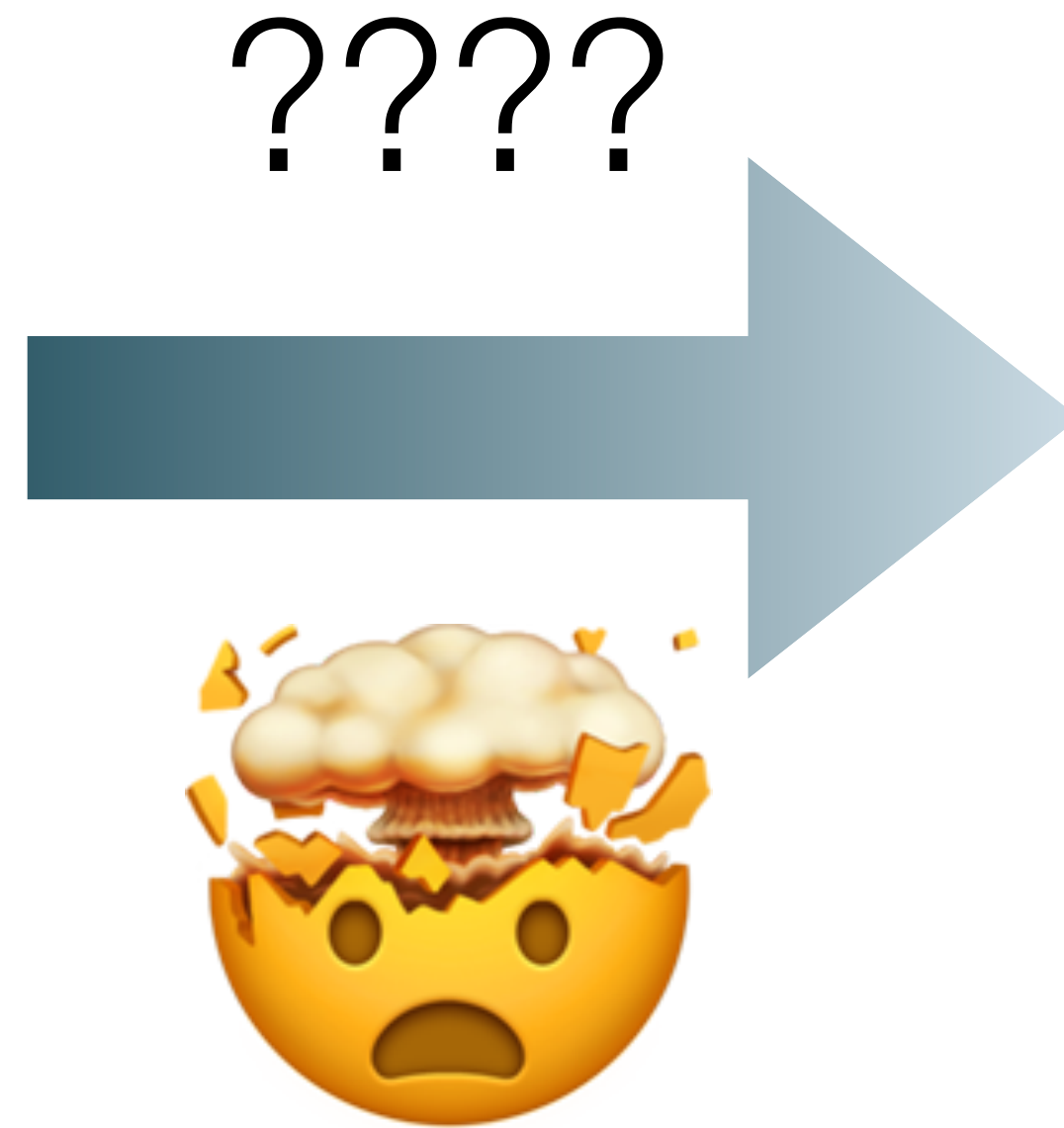
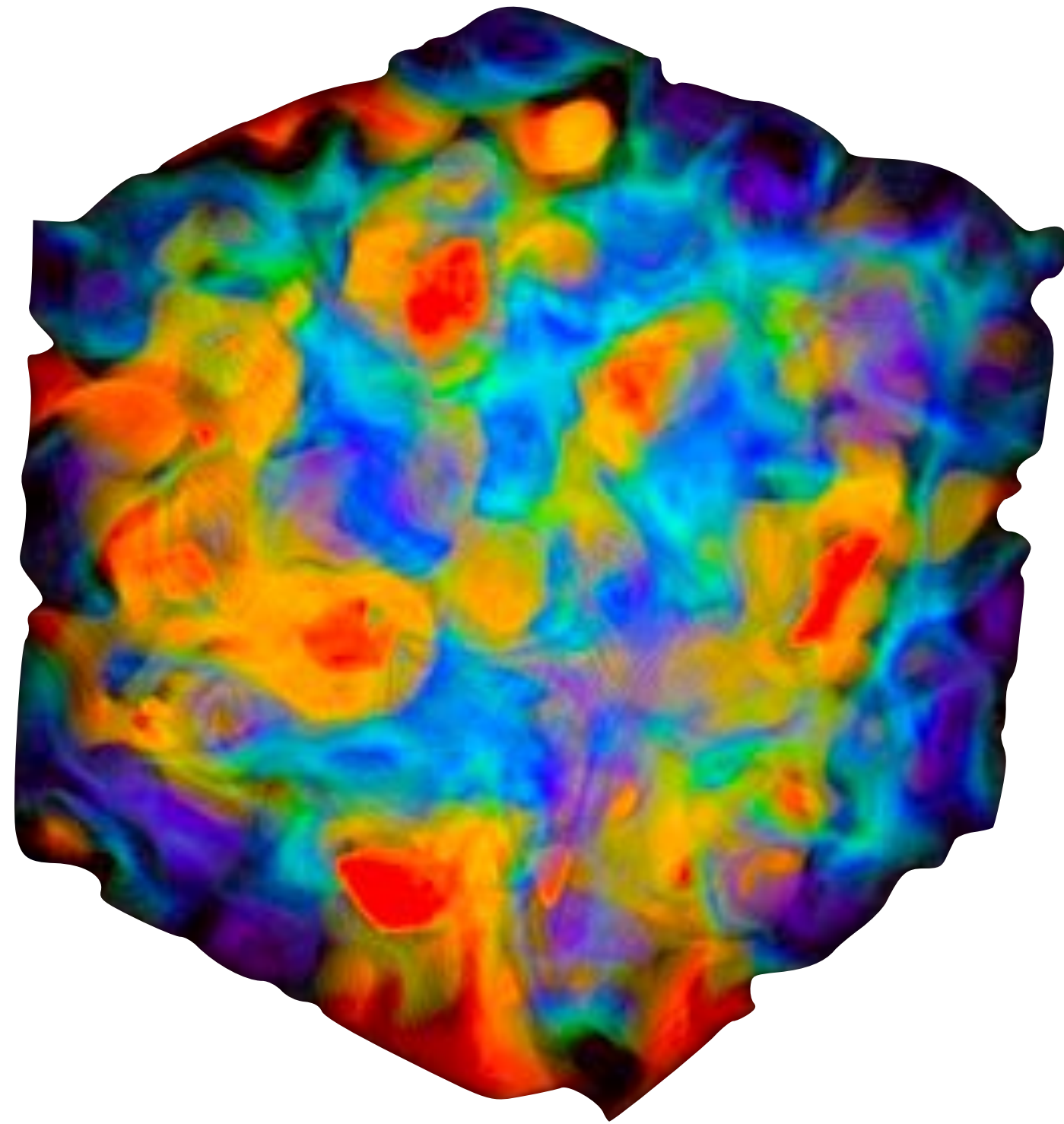


$$\gamma^{\overline{\text{MS}}}(b_{\perp}, \mu)$$

nonperturbative CS kernel

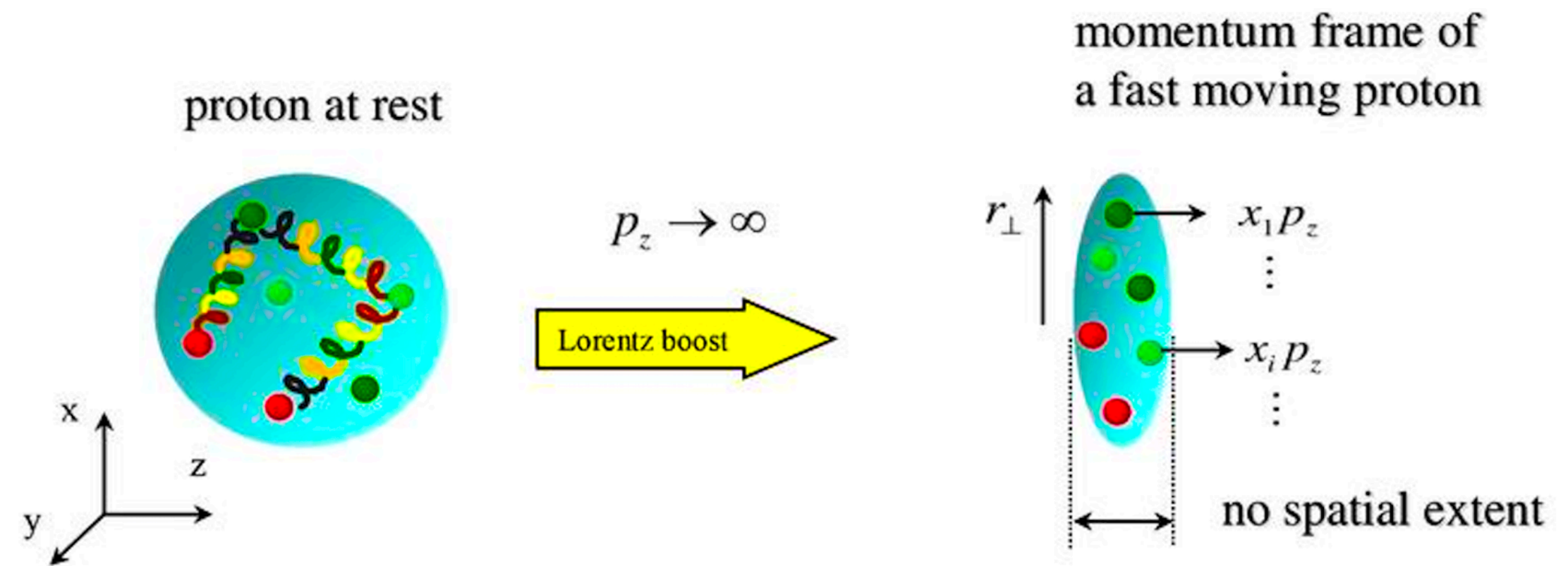
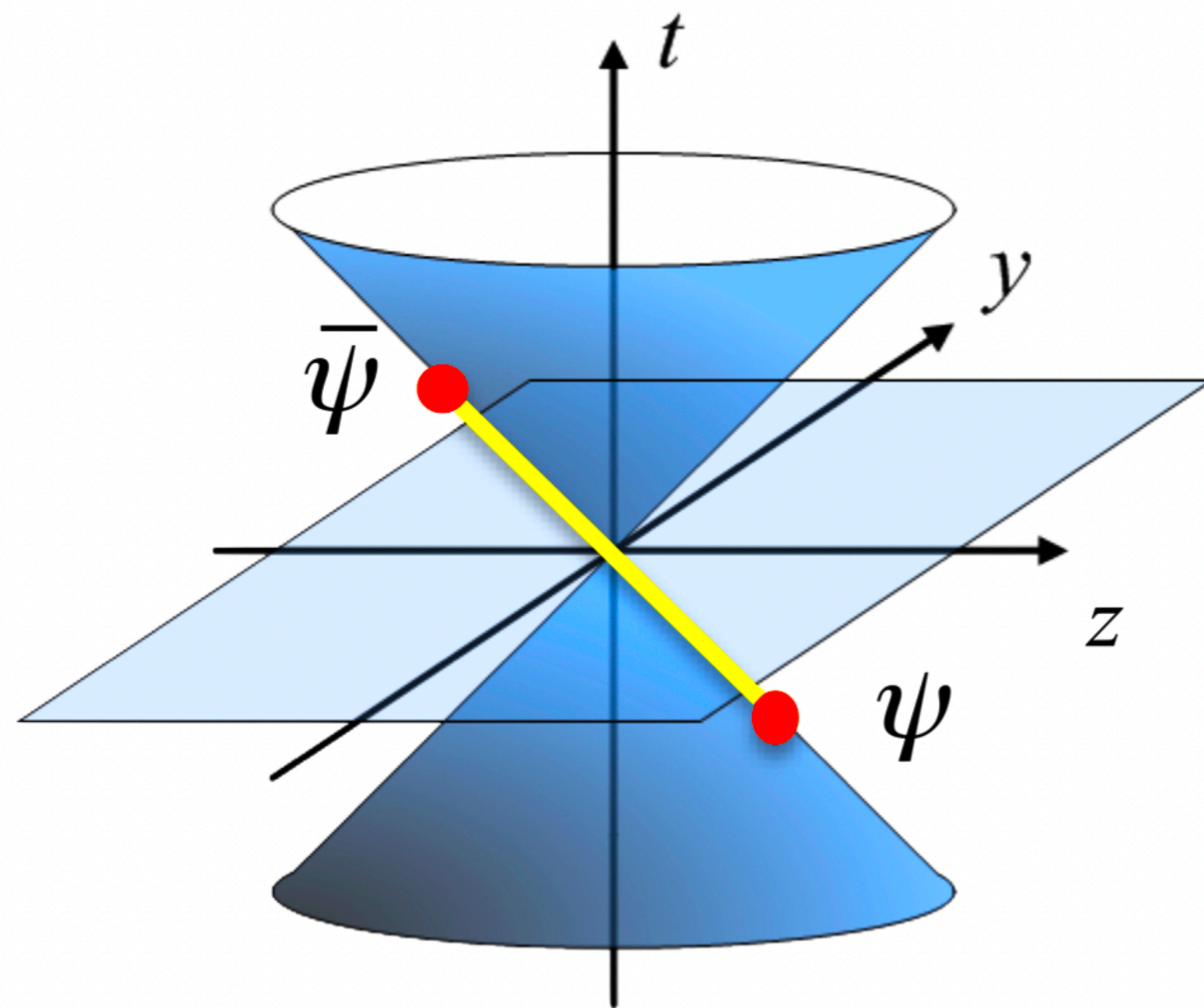
our challenge ...

how to 'see' a parton on the lattice ?



strongly coupled
quark, gluon fields

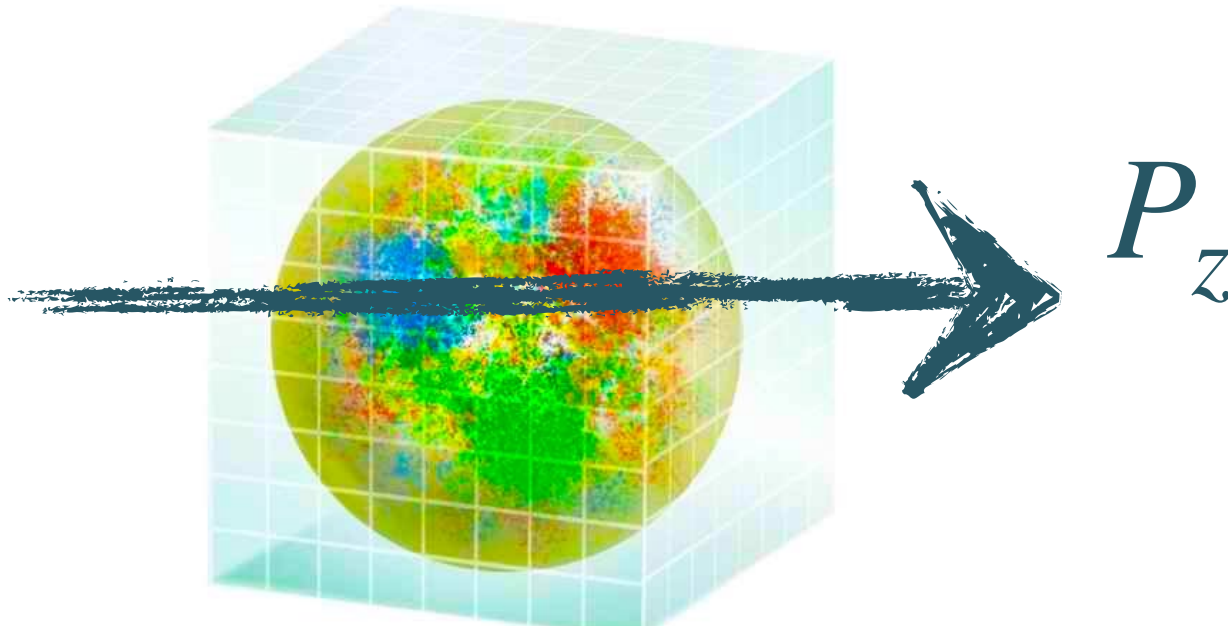
partonic picture



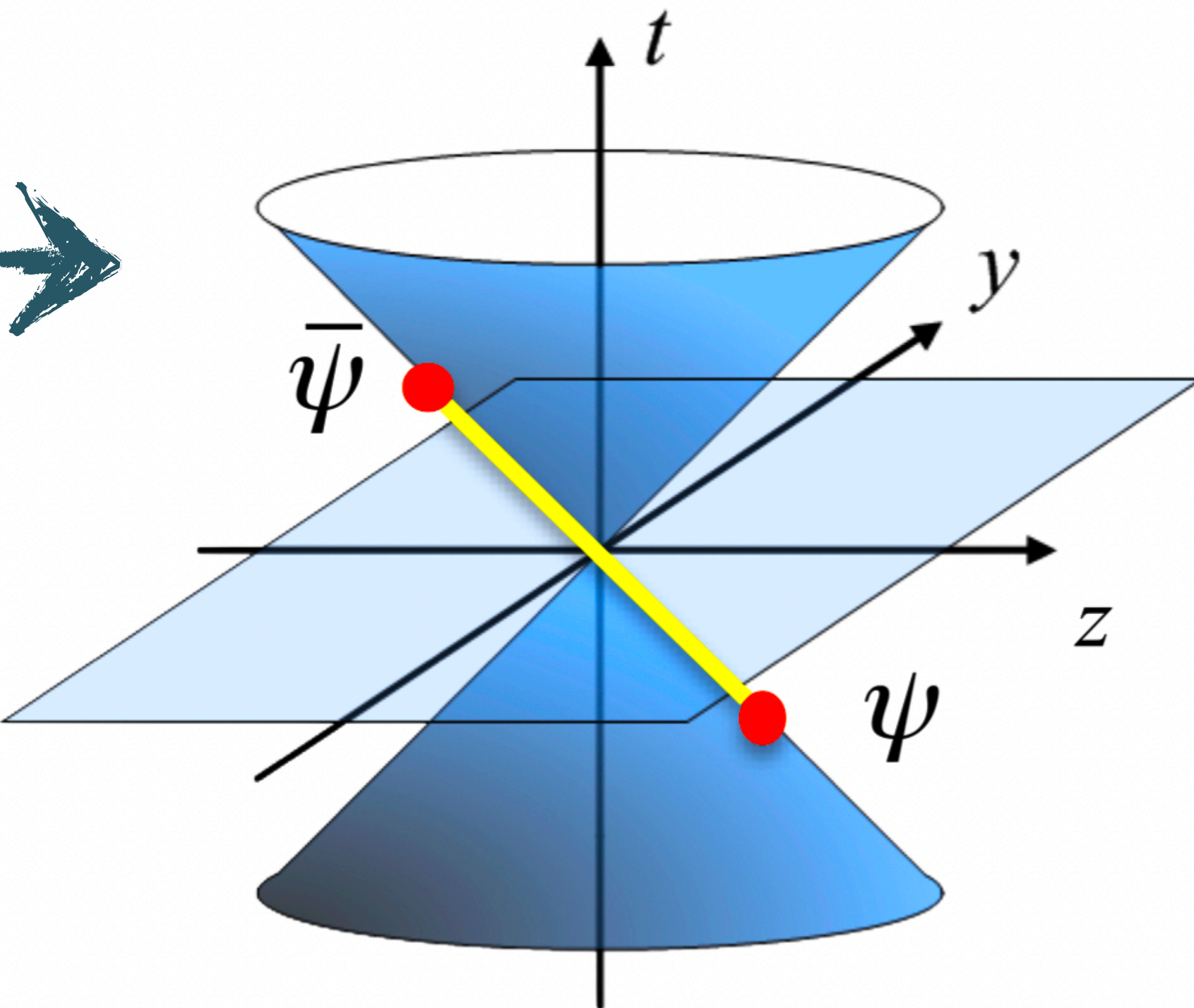
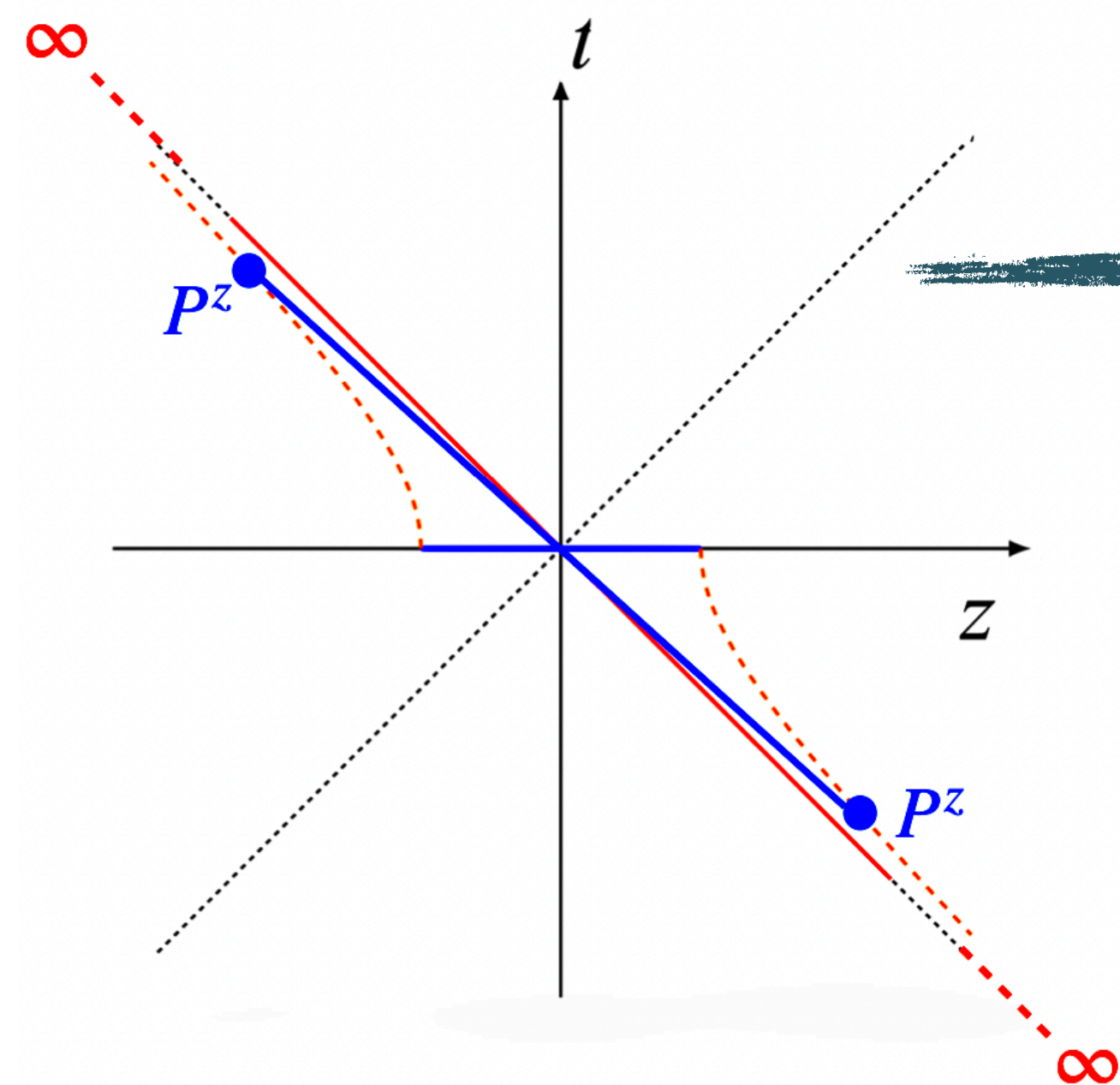
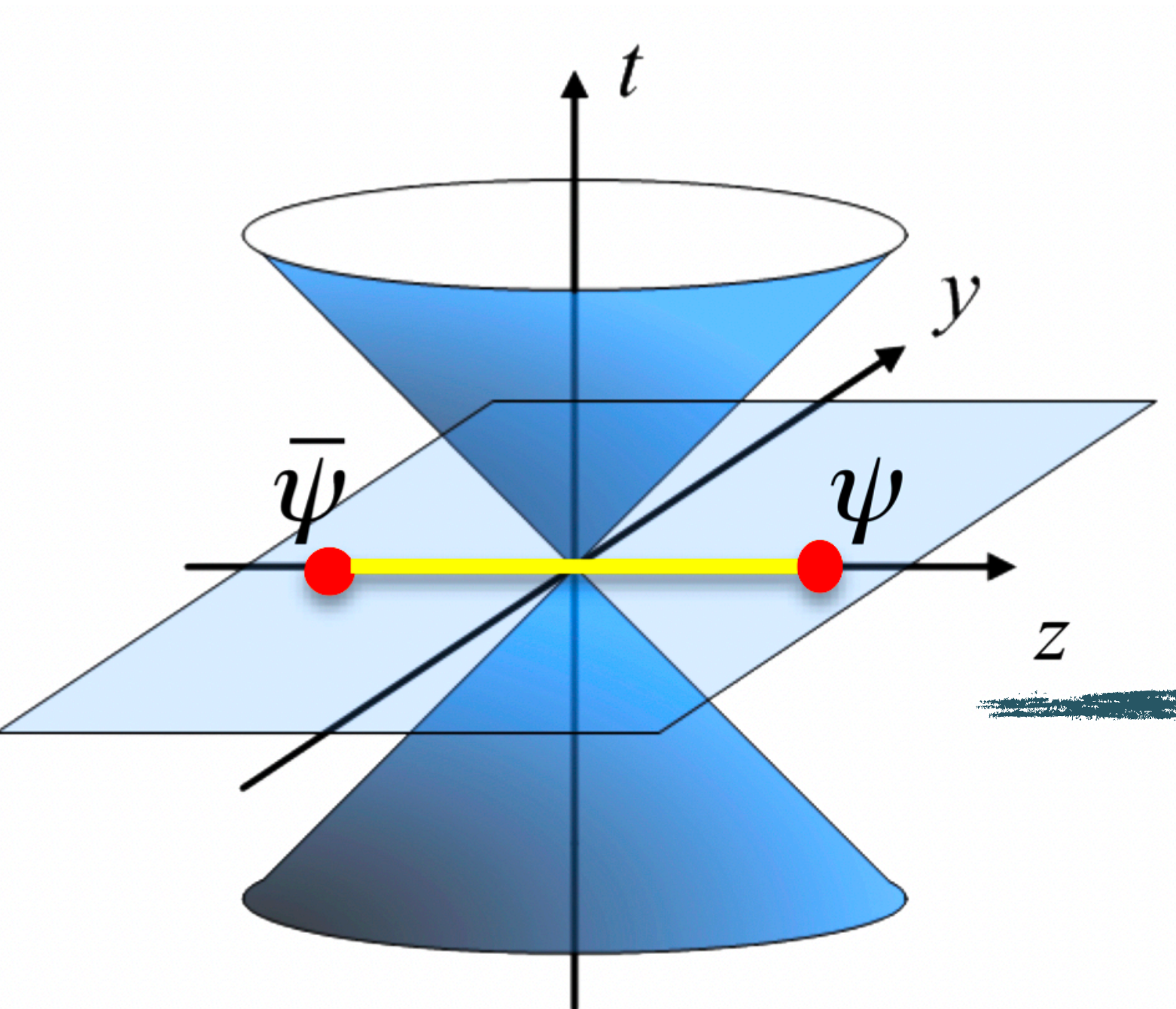
- ◆ QCD in infinite-momentum frame / probed at vanishingly short distances
- ◆ QCD simplified / effective description of on the lightcone
- ◆ $P_z \rightarrow \infty / z^2 \rightarrow 0$ first, regularize QFT later

partonic structure from lattice QCD

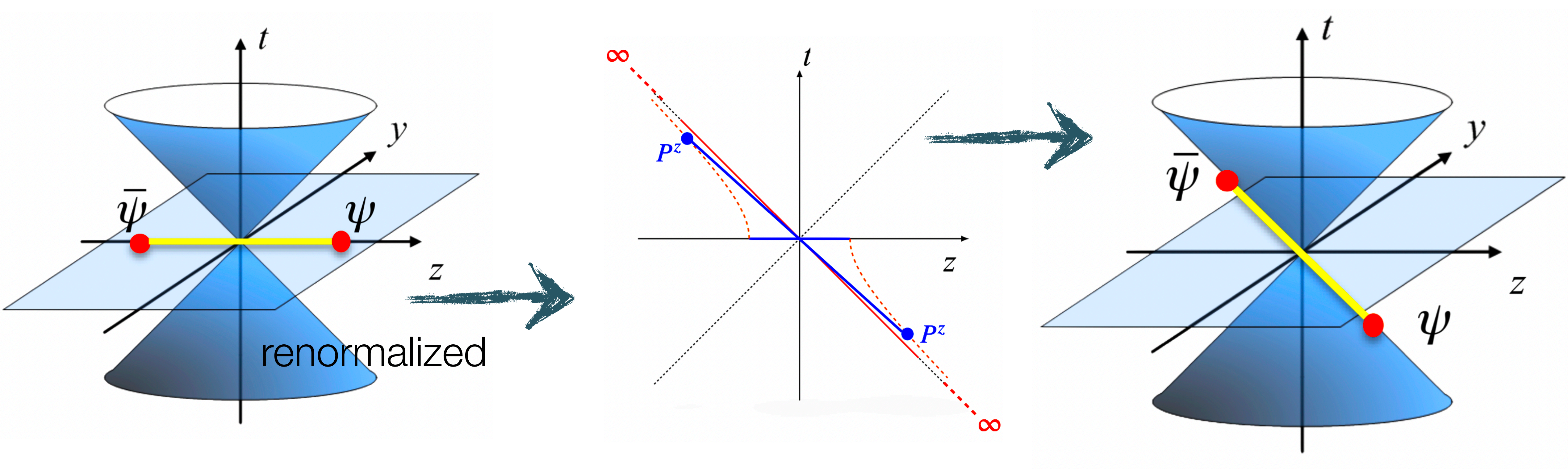
fast-moving hadron



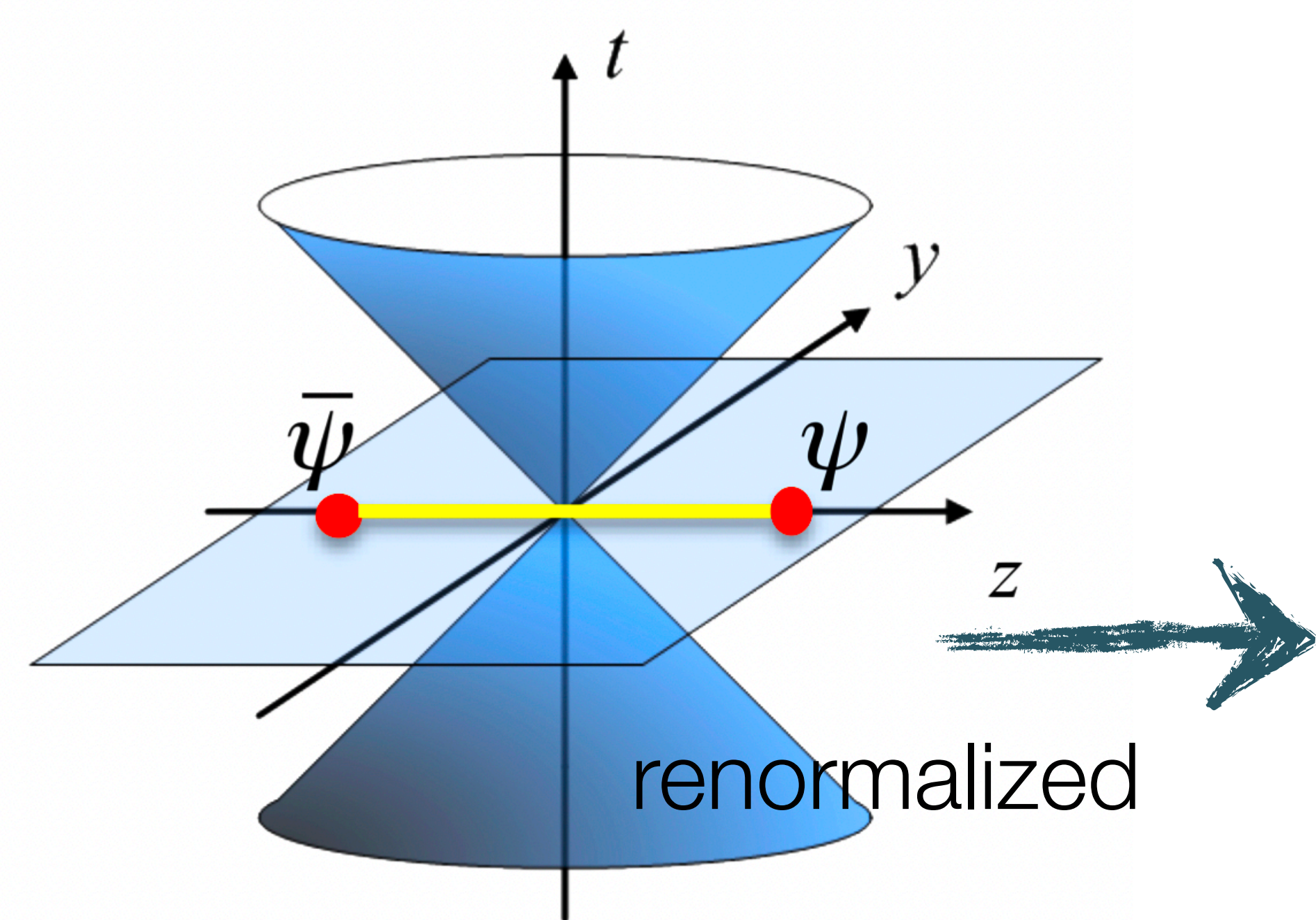
$$P_z \approx E$$



renormalize



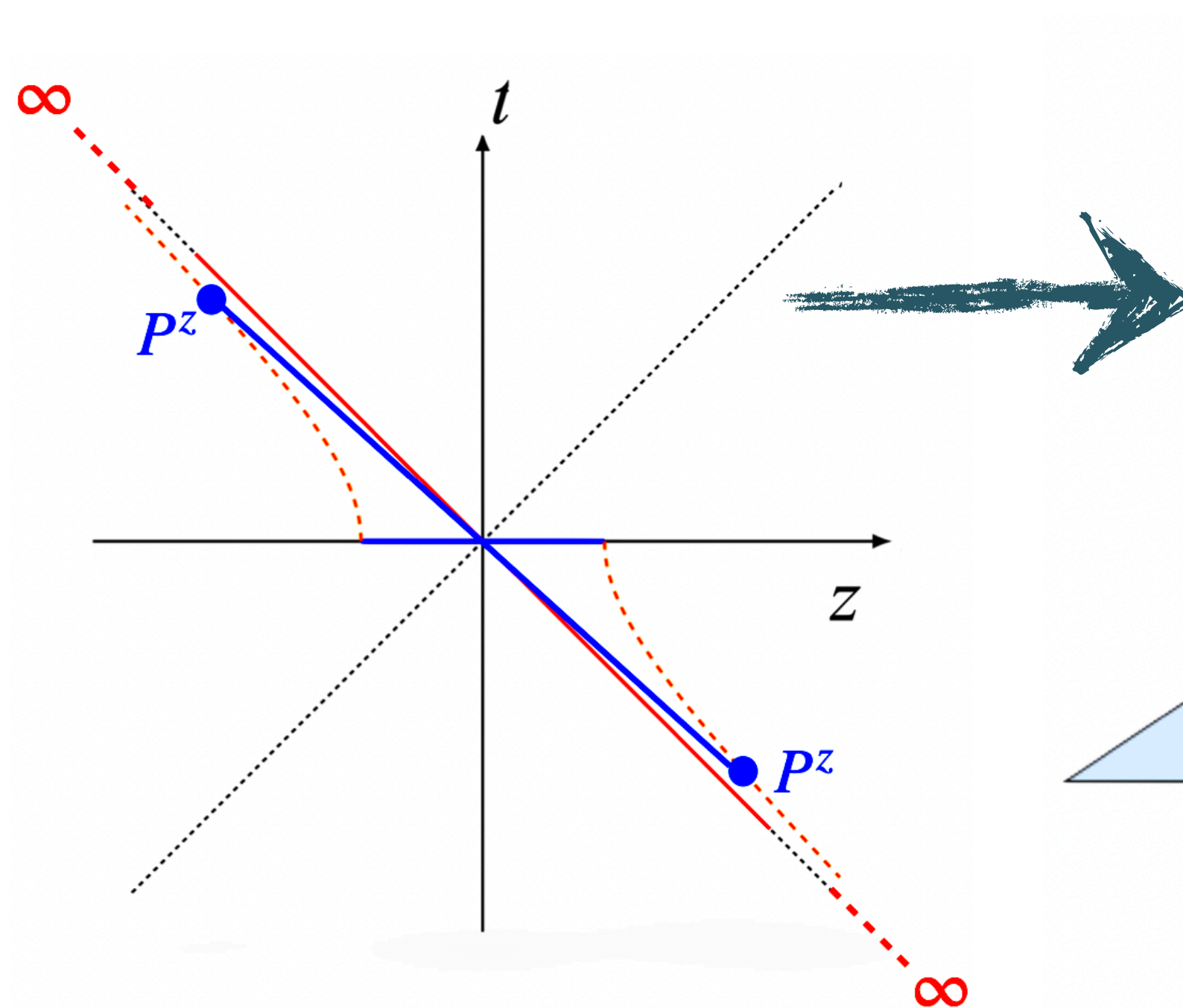
- ◆ first regularize QCD on a lattice, then $P_z \rightarrow \infty / z^2 \rightarrow 0$
- ◆ opposite order of limits; two limits don't commute
- ◆ difference is UV physics, can be taken care of through pQCD



parton physics

$$+\mathcal{O}\left[\frac{\Lambda_{\text{QCD}}^2}{x^2 P_z^2}, \frac{\Lambda_{\text{QCD}}}{(1-x)P_z}, \frac{M_H^2}{P_z^2}, \dots\right]$$

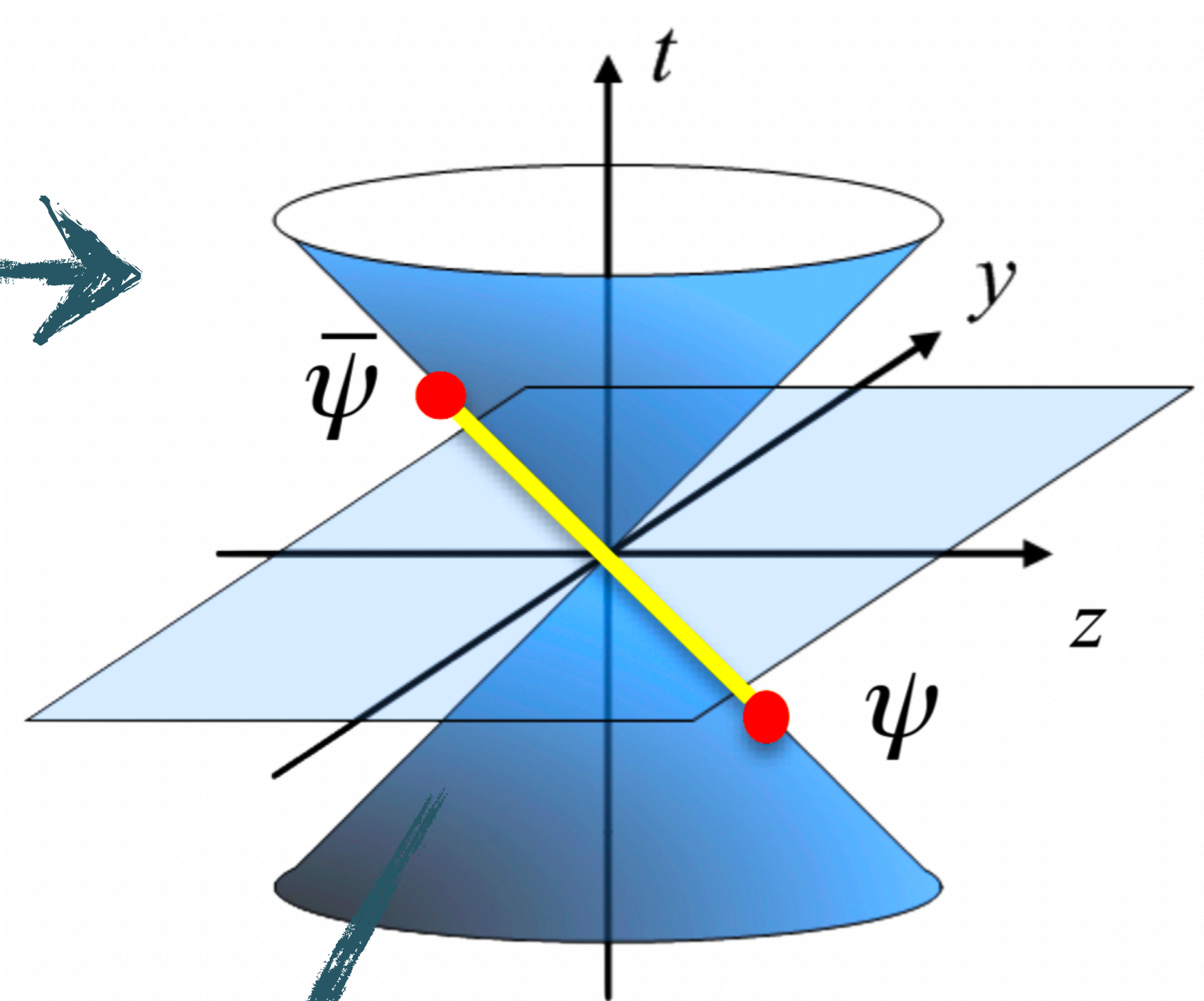
$$+\mathcal{O}\left[z^2 \Lambda_{\text{QCD}}^2, z^2 M_H^2, \dots\right]$$



pQCD

$$C(x, P_z, \mu) \otimes$$

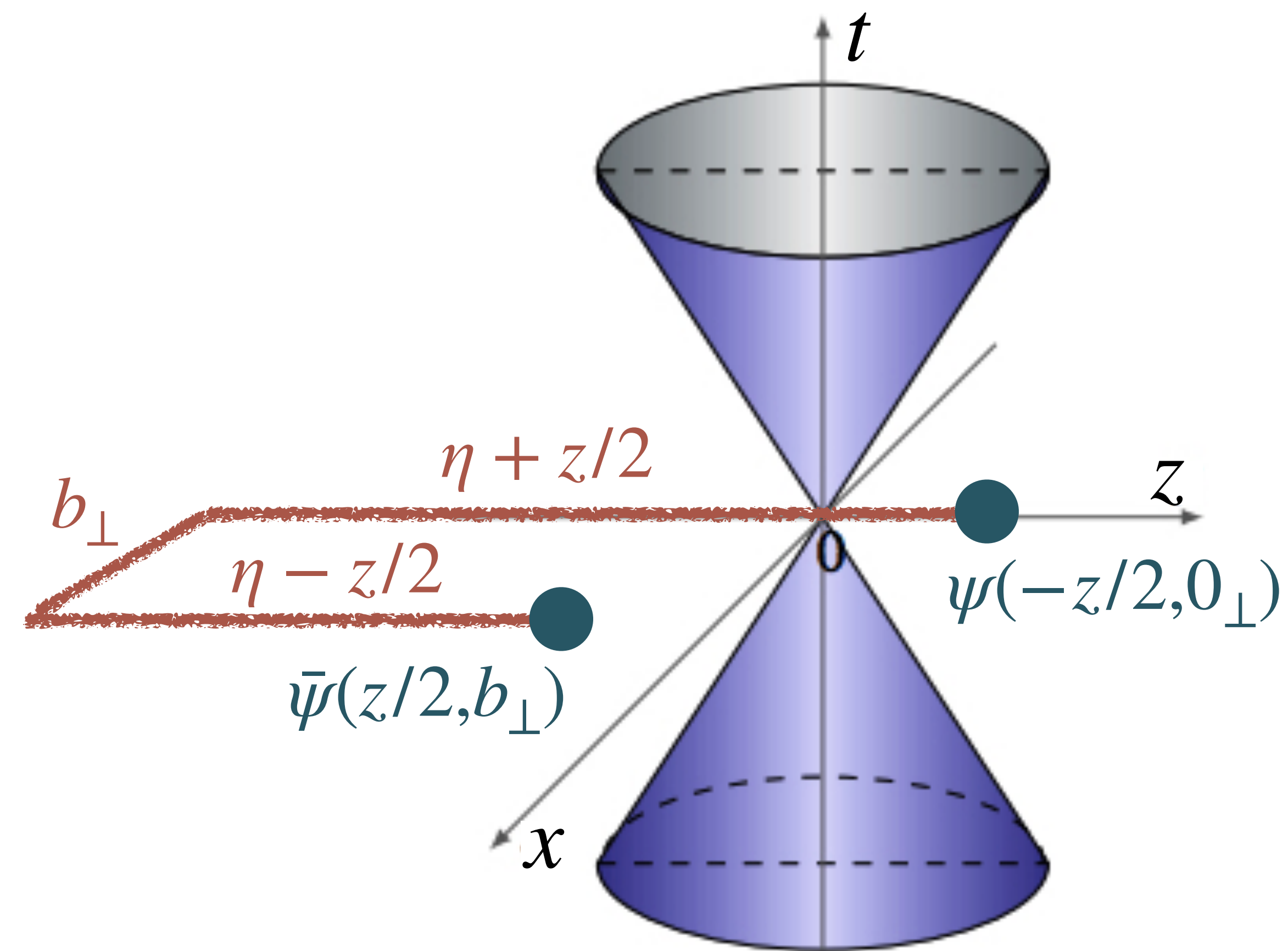
$$C(\alpha, z^2, \mu) \otimes$$



momentum space

position space

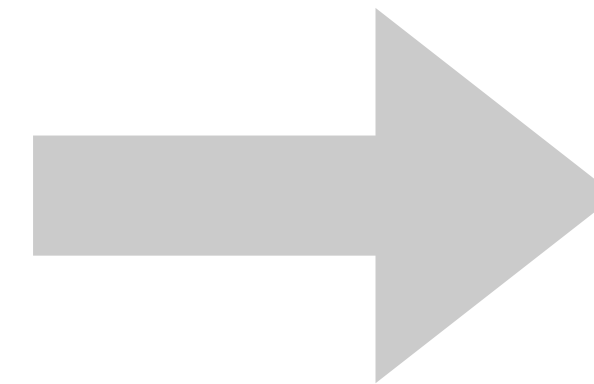
TMD distributions from lattice QCD



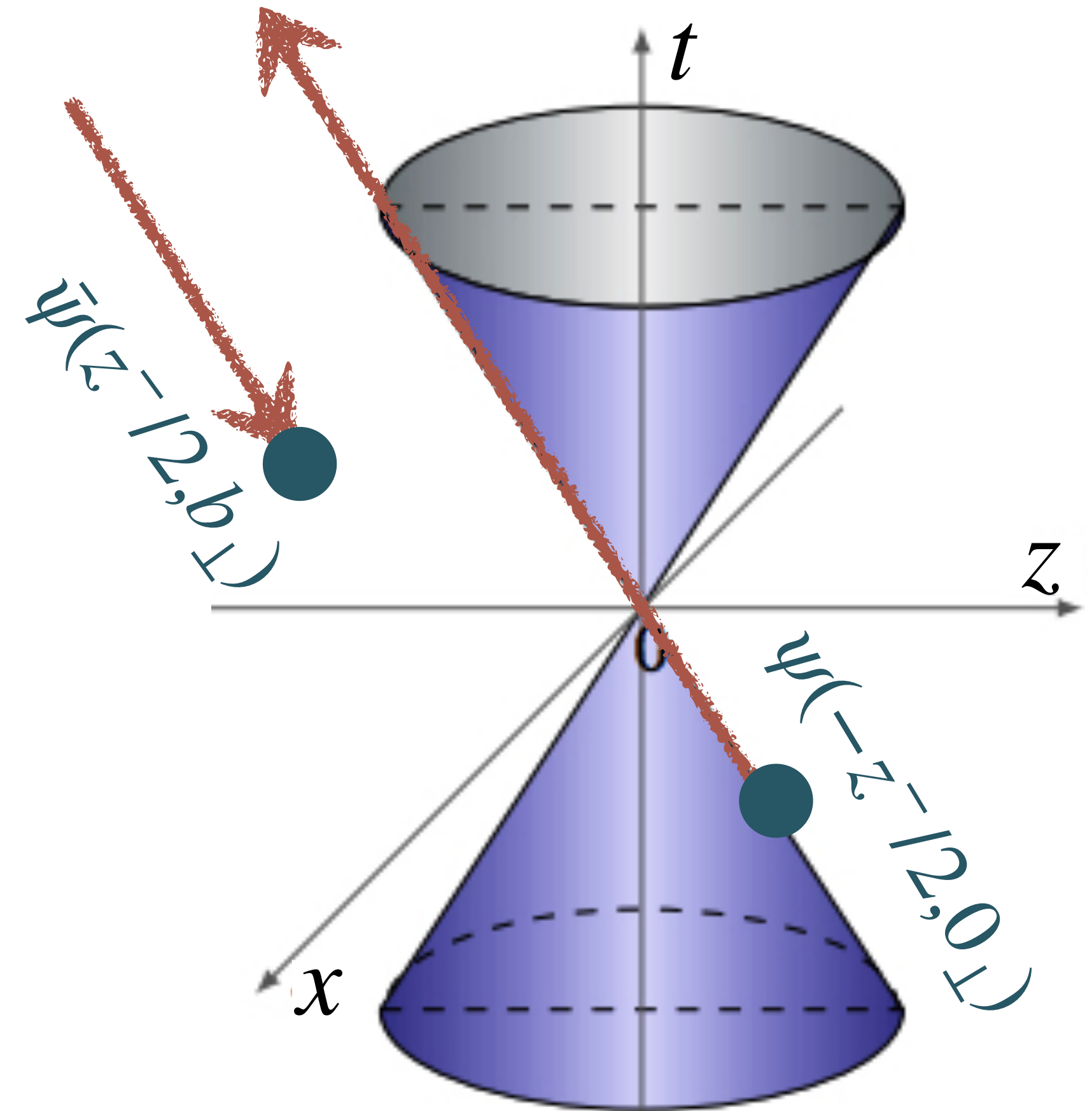
$$\tilde{\phi}(z, b_{\perp}, \eta, P_z)$$

quasi-TMD beam function

$$P_z \rightarrow \infty$$



$$\eta \rightarrow \infty$$



lightcone-TMD
beam function

◆ large η ; check for η -independence: $\tilde{\phi}(b_z, b_\perp, \eta, P_z) \longrightarrow \tilde{\phi}(b_z, b_\perp, P_z)$

◆ renormalize: $\tilde{\phi}(b_z, b_\perp, P_z) \longrightarrow \tilde{\phi}(b_z, b_\perp, P_z, \mu)$

◆ Fourier transform to momentum (x) space: $\tilde{\phi}(b_z, b_\perp, P_z, \mu) \longrightarrow \tilde{\phi}(x, b_\perp, P_z, \mu)$

◆ Perturbative matching:

Collins Soper kernel

$$\frac{\tilde{\phi}_\Gamma(x, b_\perp, P_z, \mu)}{\sqrt{S_r(b_\perp, \mu)}} = H(x, \bar{x}, P_z, \mu) \phi(x, b_\perp, \zeta, \mu) \exp \left[\frac{1}{4} \left(\ln \frac{(2xP_z)^2}{\zeta} + \ln \frac{(2\bar{x}P_z)^2}{\zeta} \right) \gamma^{\overline{\text{MS}}}(b_\perp, \mu) \right]$$

soft
function

perturbative
kernel

$$\bar{x} = 1 - x$$

lightcone
TMD
distribution

$$+ \mathcal{O} \left(\frac{\Lambda_{\text{QCD}}^2}{(xP_z)^2}, \frac{1}{(b_\perp(xP_z))^2}, \frac{\Lambda_{\text{QCD}}^2}{(\bar{x}P_z)^2}, \frac{1}{(b_\perp(\bar{x}P_z))^2} \right)$$

power corrections

CS kernel from lattice QCD

◆ ratios of quasi-TMD beam functions for 2 different boost momenta, P_1 & P_2

Collins Soper kernel

perturbative kernel

$$\gamma^{\overline{\text{MS}}}(b_{\perp}, \mu) = \frac{1}{\ln(P_2/P_1)} \ln \left[\frac{\tilde{\phi}(x, b_{\perp}, P_2, \mu)}{\tilde{\phi}(x, b_{\perp}, P_1, \mu)} \right] + \delta\gamma^{\overline{\text{MS}}}(x, \mu, P_1, P_2)$$

independent of x, P_1, P_2

+ power corrections

soft function cancels

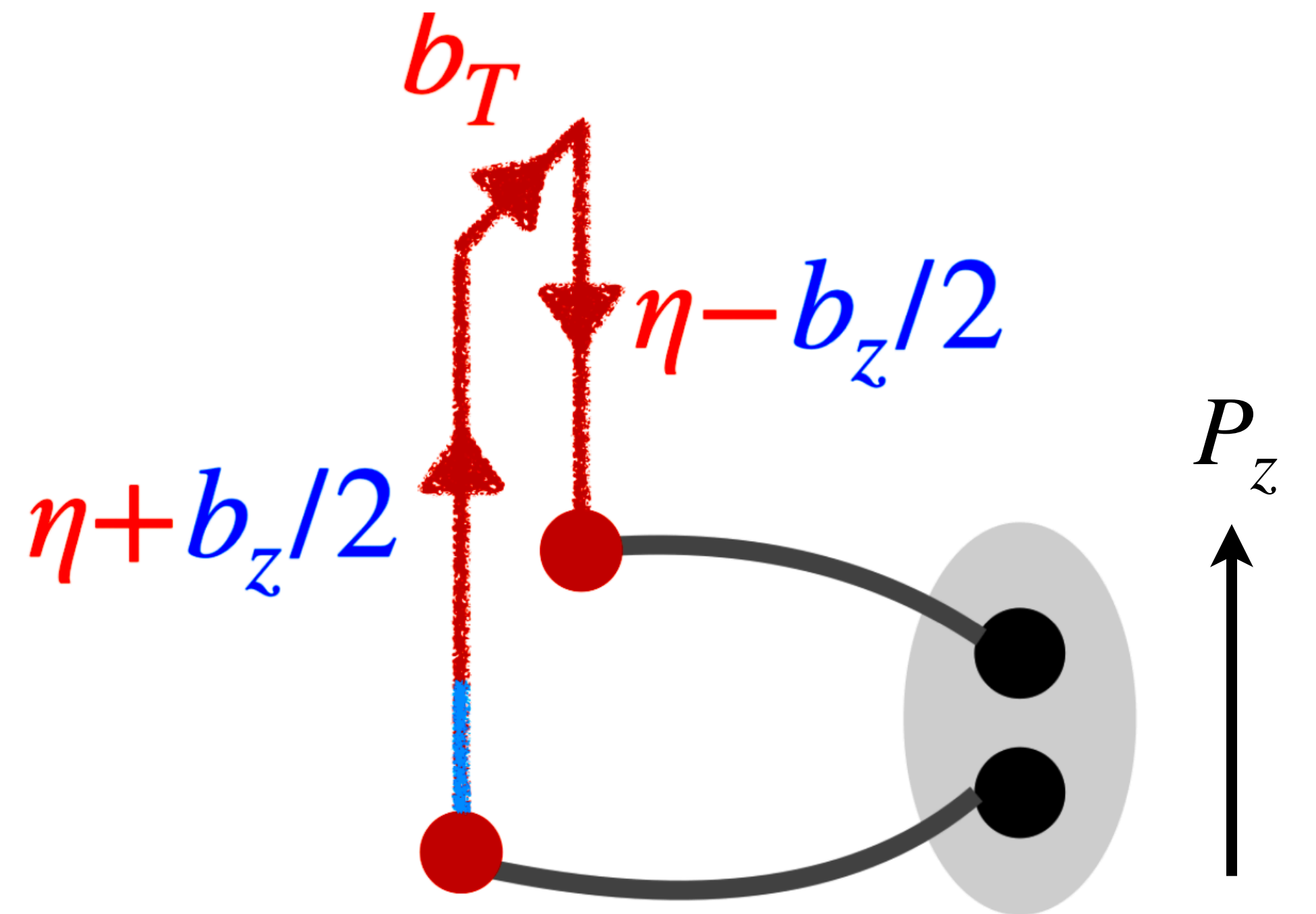
P_1 & P_2 both must be large to suppress power corrections,
such that CS kernel is indep. of those

lattice QCD calculations of CS kernel

- ◆ simplest choice for the quasi-TMD beam function $\tilde{\phi}(b_z, b_\perp, \eta, P_z)$

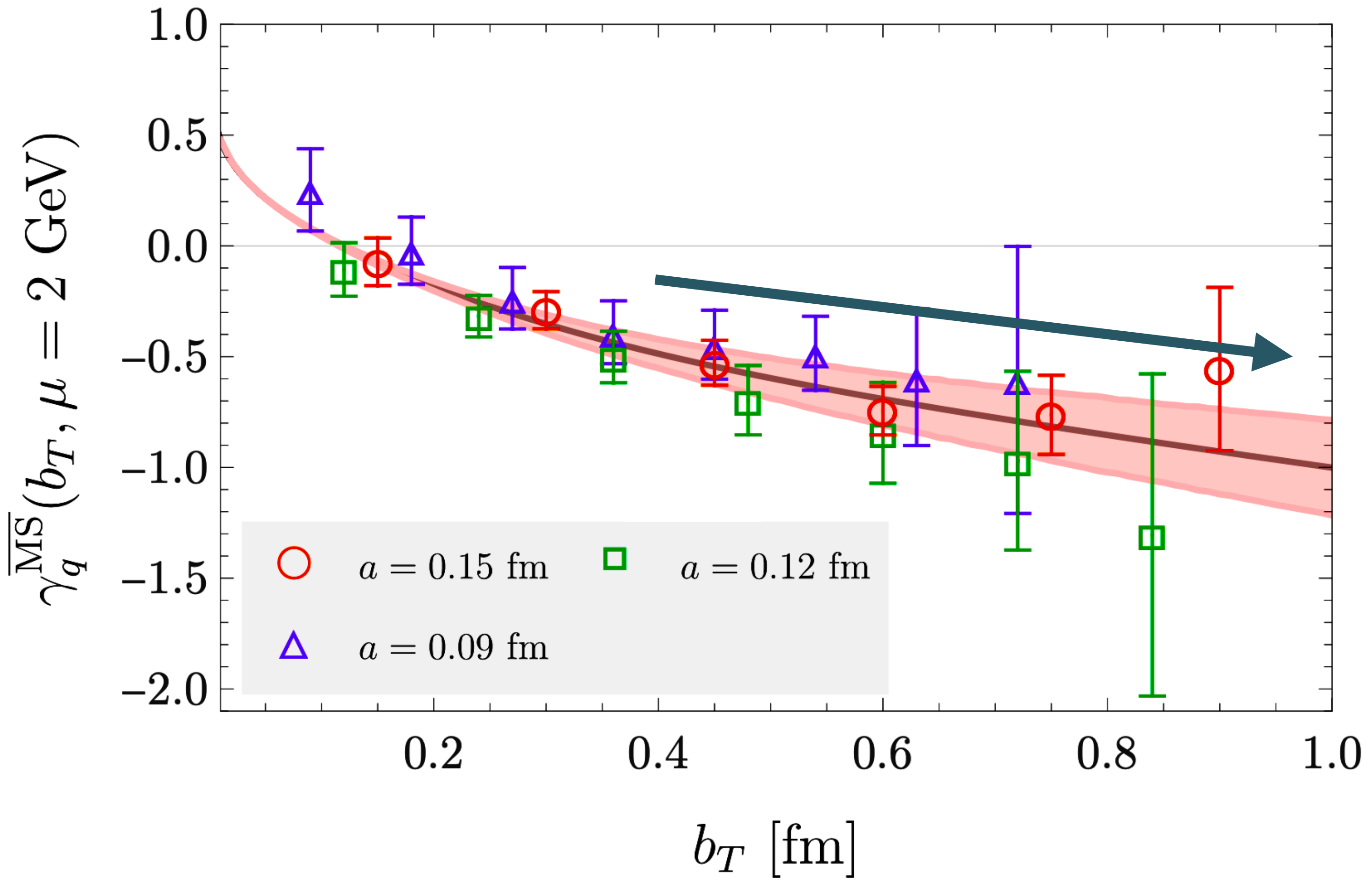
pion TMD wave function (TMDWF)

$$\langle \Omega | \bar{\psi}(\frac{b_z}{2}, b_\perp) \Gamma W_\square(\frac{\mathbf{b}}{2}, -\frac{\mathbf{b}}{2}, \eta) \psi(-\frac{b_z}{2}, 0) | \pi^+, P_z \rangle$$



difficulties in lattice QCD calculations

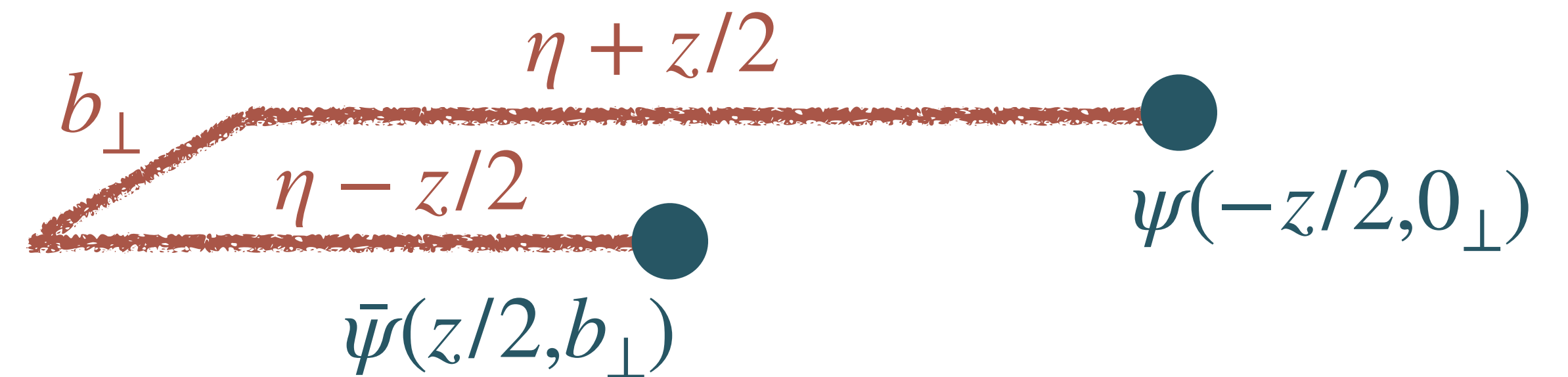
rapidly growing errors with increasing b_{\perp}



why difficult ?

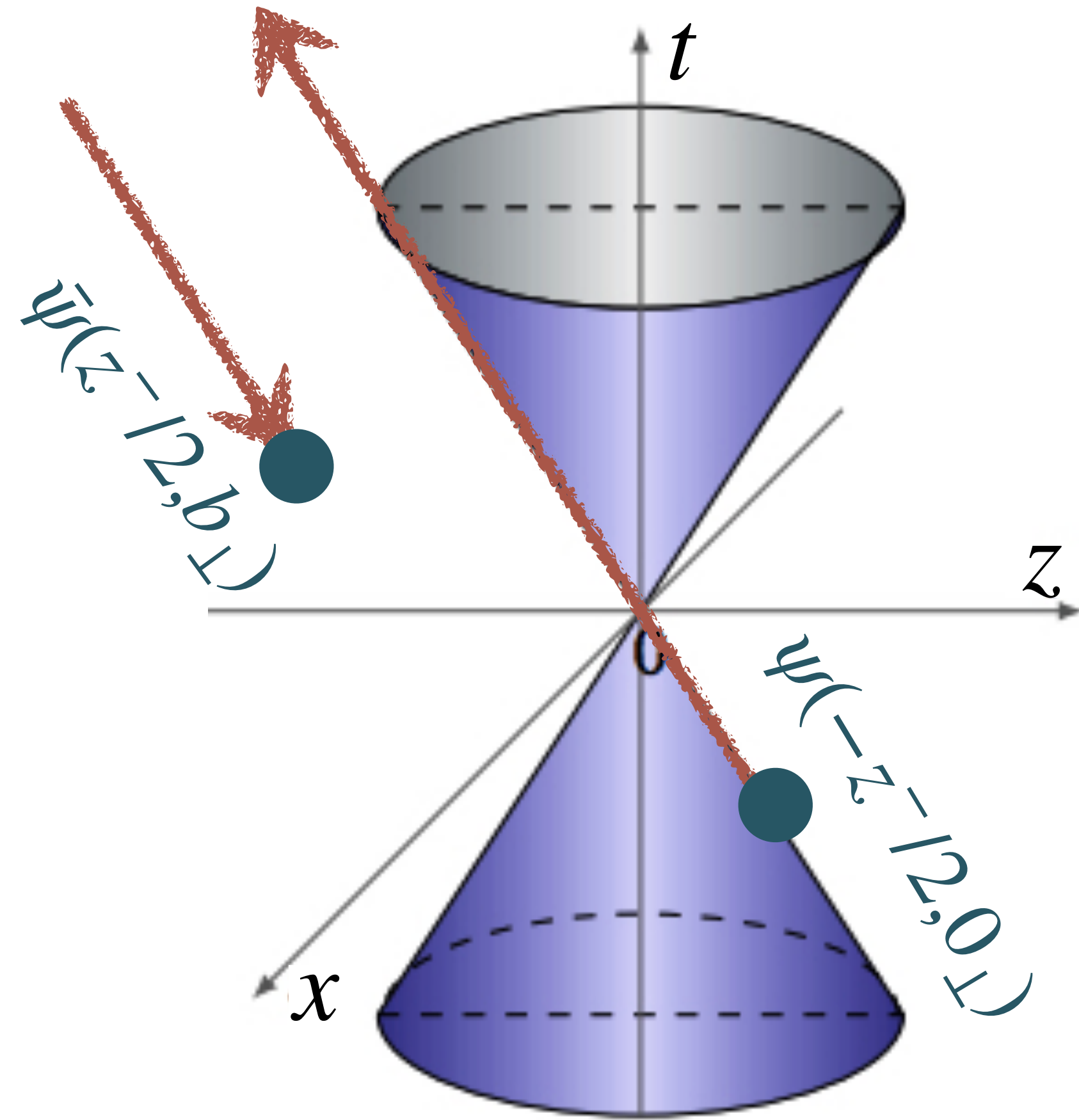
multiplicative renormalization factor of the Wilson line:

$$\sim e^{-\delta m(\eta + b_{\perp})}$$

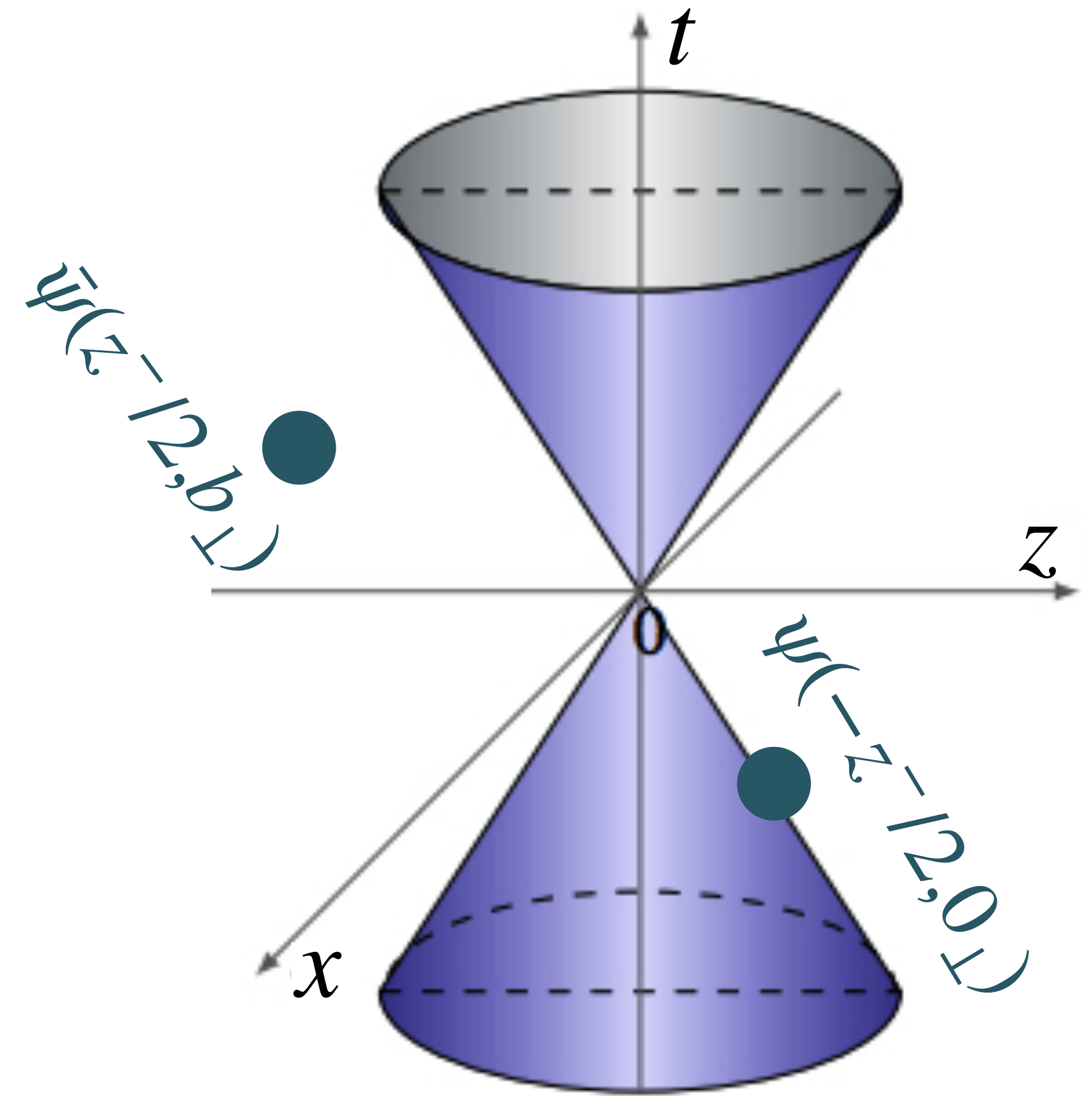


exponential decrease of signal for large η and increasing b_{\perp}

overcoming difficulties



≡



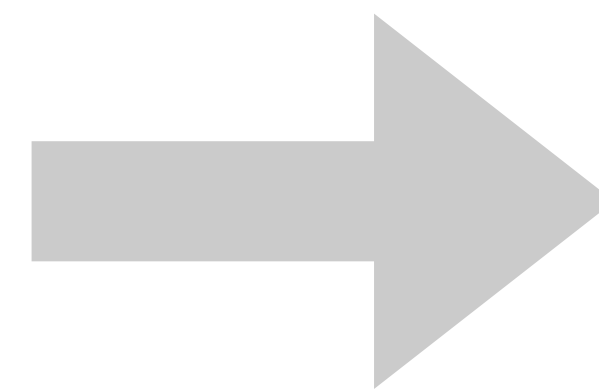
overcoming difficulties

but how can we access $A^+ = 0$ in lattice QCD calculations ?

find a gauge that becomes equivalent to $A^+ = 0$ in the limit $P_z \rightarrow \infty$

Coulomb gauge:

$$\vec{\nabla} \cdot \vec{A} = 0$$



$$P_z \rightarrow \infty$$

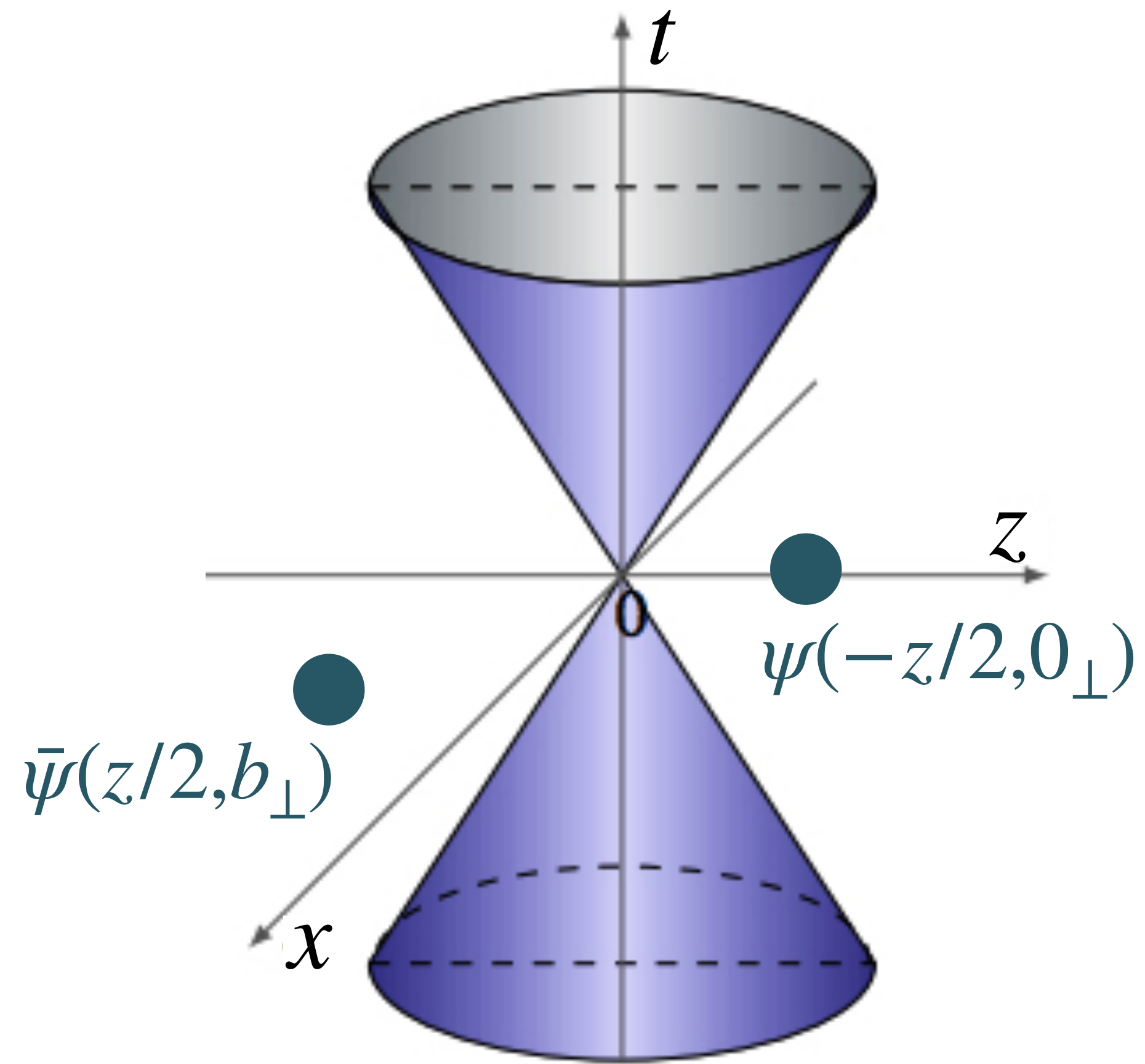
$$A^+ = 0$$

Ji et al., Phys. Rev. Lett. 111, 112002 (2013)

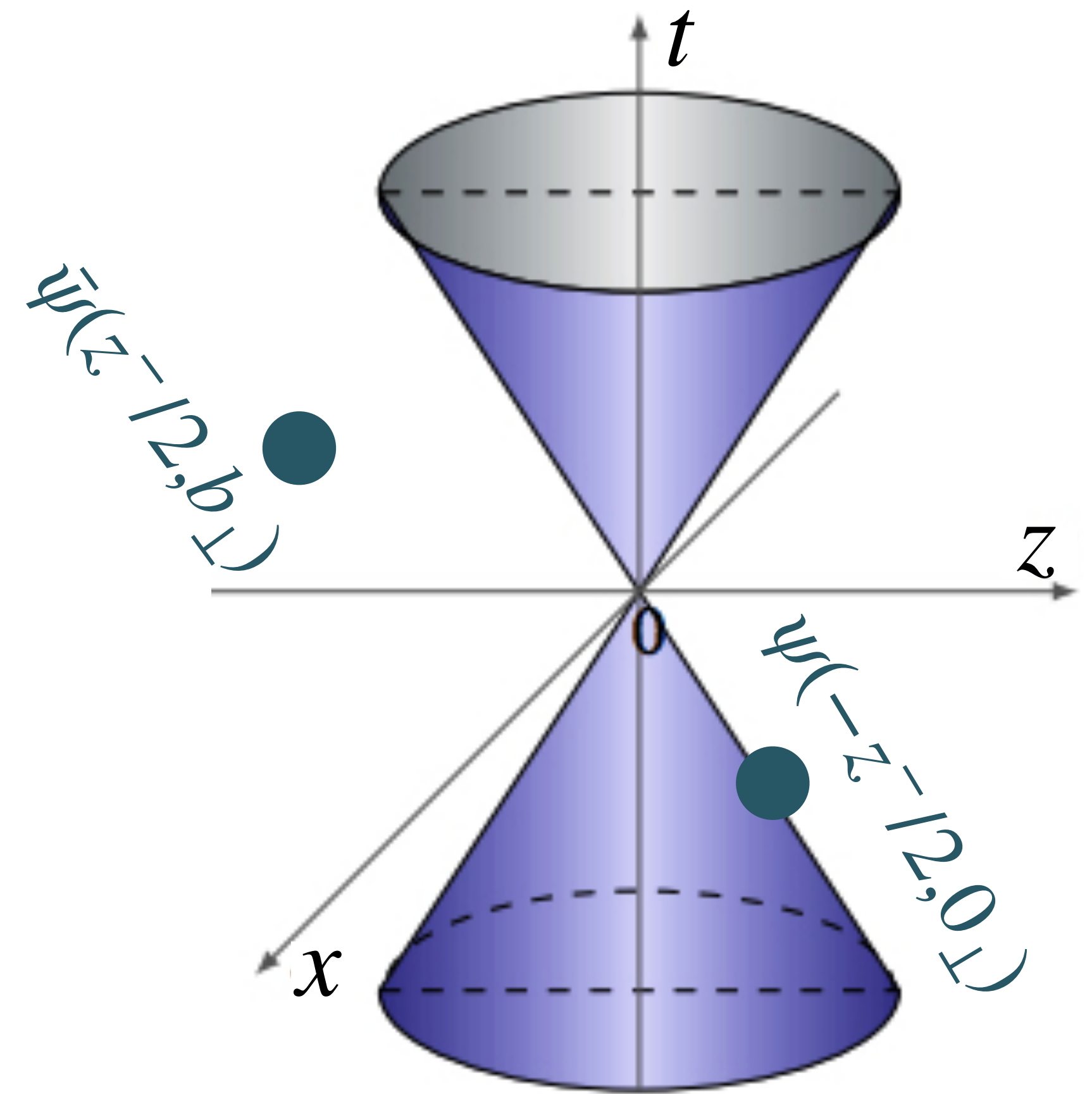
Hatta et al., Phys. Rev. D 89, no.8, 085030 (2014)

quasi-TMD beam function in Coulomb gauge (CG)

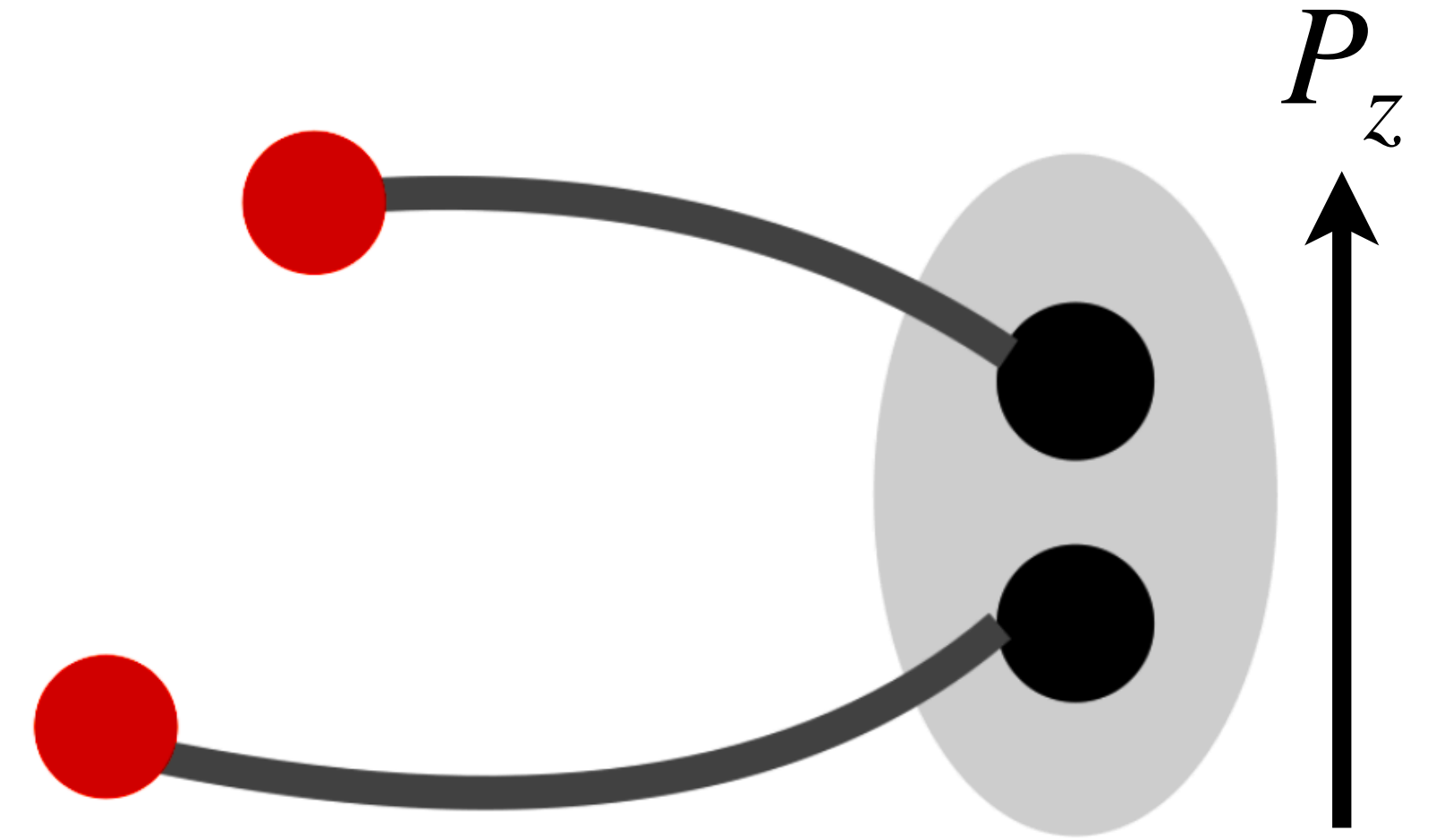
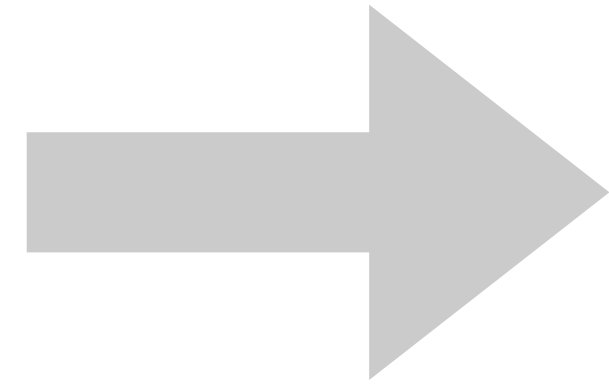
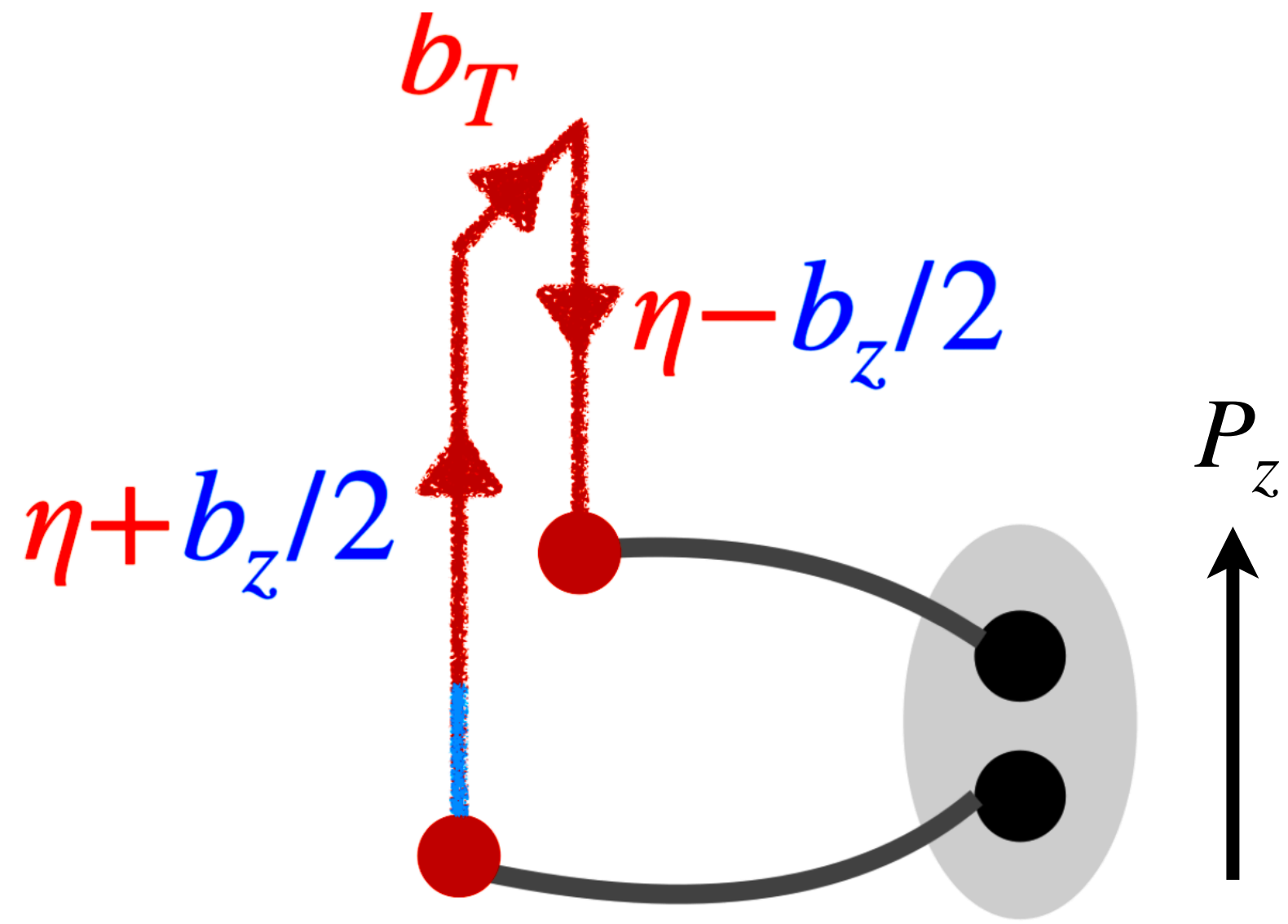
$$\bar{\psi}\left(\frac{\mathbf{b}}{2}\right)\Gamma\psi\left(-\frac{\mathbf{b}}{2}\right)\Big|_{\nabla\cdot\mathbf{A}=0}$$



$P_z \rightarrow \infty$



CG quasi-TMD beam function



$$\langle \Omega | \bar{\psi}(\frac{b_z}{2}, b_{\perp}) \Gamma W_{\square}(\frac{\mathbf{b}}{2}, -\frac{\mathbf{b}}{2}, \eta) \psi(-\frac{b_z}{2}, 0) | \pi^+, P_z \rangle$$

$$\langle \Omega | \bar{\psi}(\frac{b_z}{2}, b_{\perp}) \Gamma \psi(-\frac{b_z}{2}, 0) | \vec{\nabla} \cdot \vec{A} = 0 | \pi^+, P_z \rangle$$

+ re-computation of pQCD matching function $\delta\gamma^{\overline{\text{MS}}}(x, \mu, P_1, P_2)$

NLL accuracy

Zhao: 2311.01391

renormalized quasi-TMD beam functions

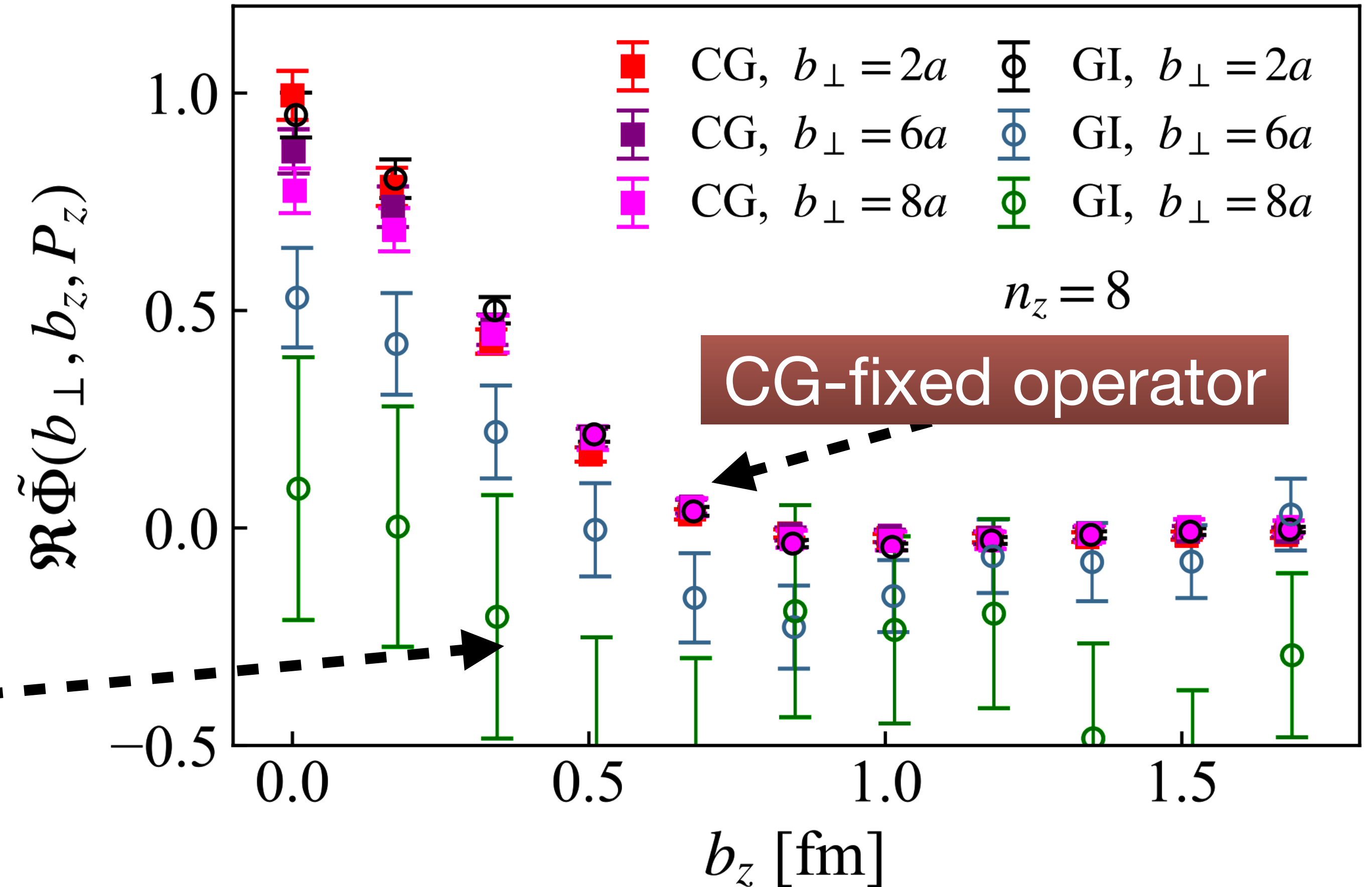
Bollweg et al.: Phys. Lett. B 852, 138617 (2024)

unitary chiral (Domain Wall)
fermions, physical pion mass

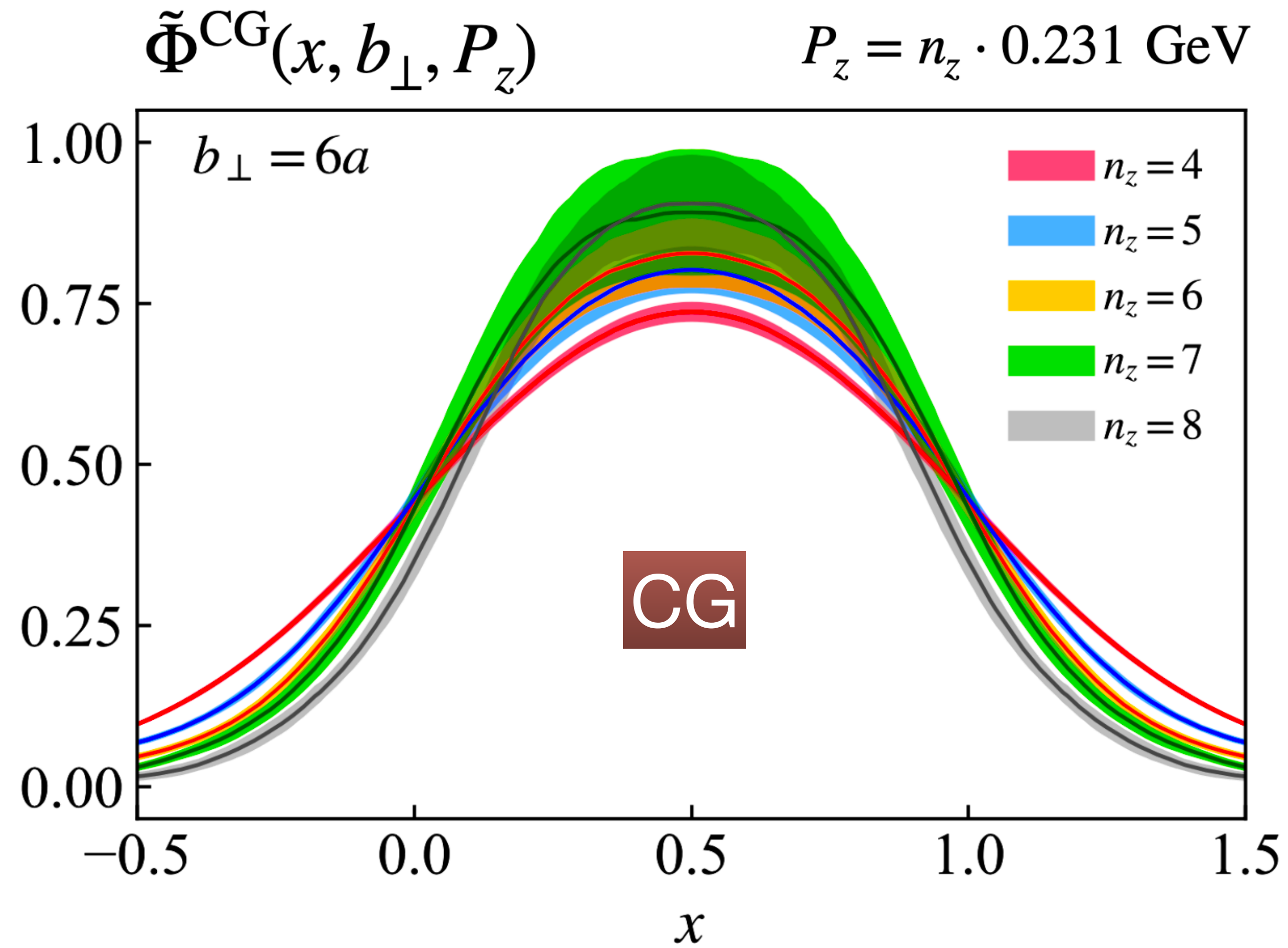
lattice spacing $a=0.085$ fm

gauge invariant (GI) operator

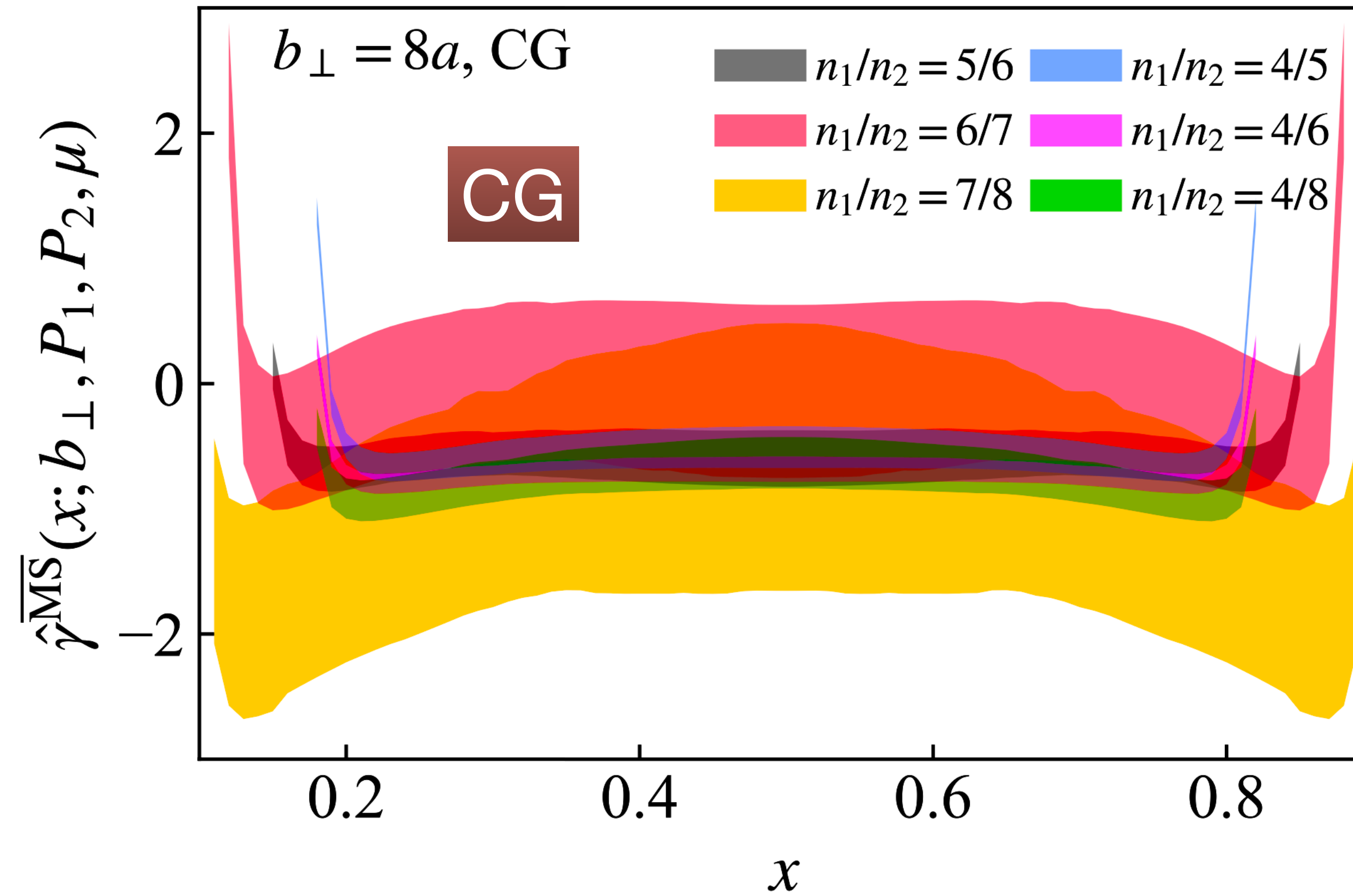
$$P_z = 0.231 n_z \text{ GeV}$$



quasi-TMD beam functions in momentum space



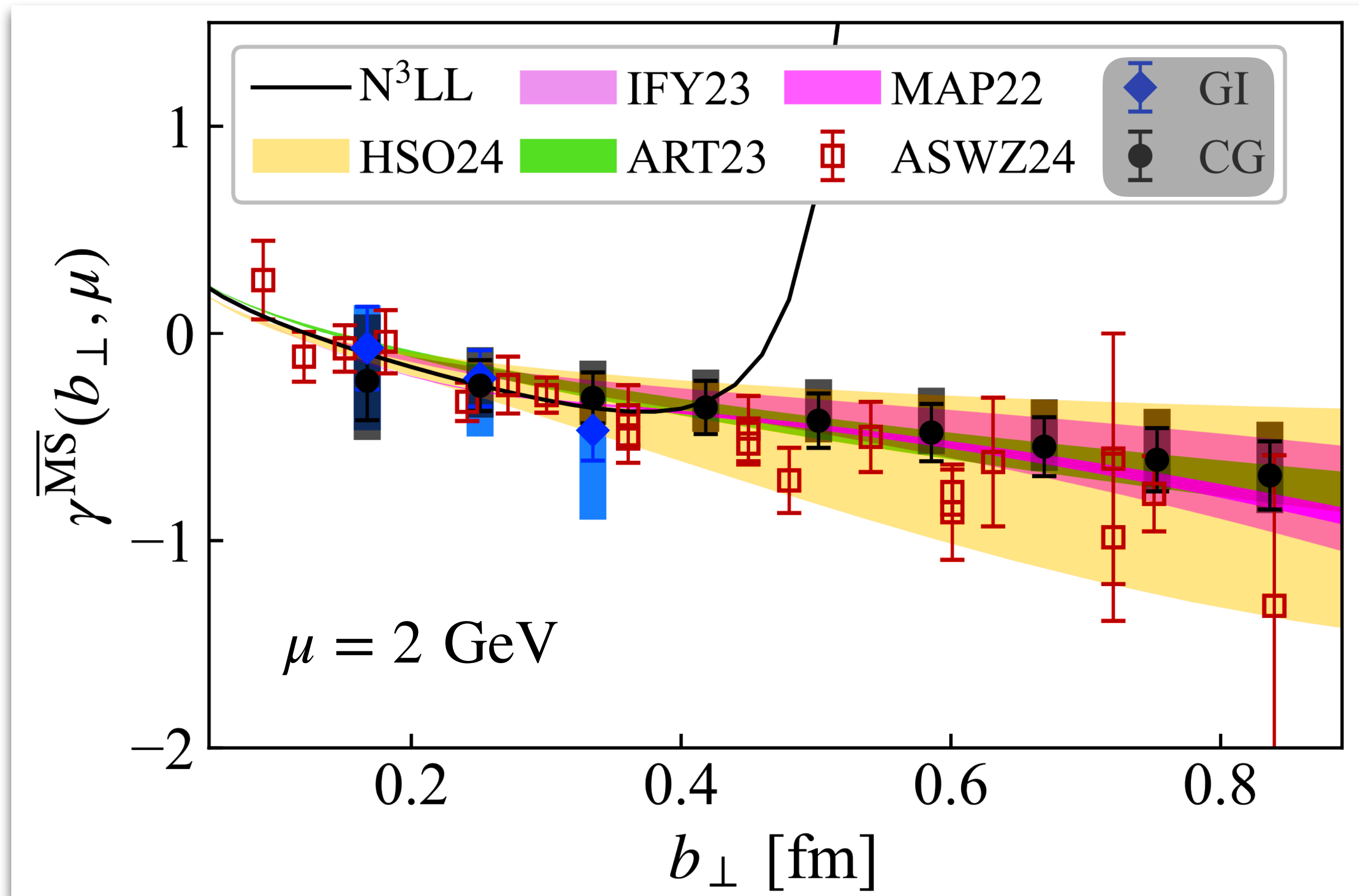
x and P independence of CS kernel



$$P_i = 0.231n_i \text{ GeV}, \mu = 2 \text{ GeV}$$

Summary: nonperturbative CS kernel from lattice QCD

Bollweg et al.: Phys. Lett. B 852, 138617 (2024)



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Bollweg et al.: Phys. Lett. B 852, 138617 (2024)

