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QCD Evolution 2024

Collins asymmetries in e^+e^- annihilation: a MC event generator study

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Istituto Nazionale di Fisica Nucleare

Studying hadronization

- ❑ Phenomenological QCD fits (SIDIS, e+e-, pp)
parametrize FFs and PDFs, extract from data e.g. Collins FF or IFF, and transversity

- ❑ Development of hadronization models

Studying hadronization

- Phenomenological QCD fits (SIDIS, e+e-, pp)
parametrize FFs and PDFs, extract from data e.g. Collins FF or IFF, and transversity
- Development of hadronization models
analytic calculation of FFs e.g. using spectator models Bacchetta et al., PLB 659 (2008) 234-243
multiproduction models
 - **String fragmentation** models
 - Artru, Menessier '74 no spin
 - Lund String Model '83 no spin
 - string+ 3P_0** spin [Artru '09; Artru, Belghobsi '13; Kerbizi, Artru, Belghobsi, Bradamante, Martin '18..]
 - Cluster model no spin [Marchesini, Webber '84]
 - quark jet models
 - Field, Feynman '77 no spin
 - NJL-jet model '17* spin [Matevosyan, Kotzinian, Thomas, '17]
- MC event generators needed to perform calculations with multiproduction models
general purpose generators (Pythia, Lepto, Herwig).. **no spin in hadronization**
for spin effects → develop a model of spin-dependent hadronization
then implement in MC

Modeling spin-dependent hadronization: the string+ 3P_0 model

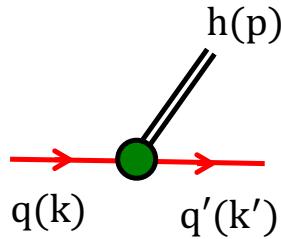
- Extension of the Lund string fragmentation model that includes the quark spin based on amplitudes, preserves positivity and entanglement different developments

Artru, DSPIN-09, arXiv:1001.1061	2009	toy model
AK, Artru, Belghobsi, Bradamante, Martin, PRD 97, 074010 (2018)	2018	PS mesons
AK, Artru, Belghobsi, Martin, PRD 100, 014003 (2019)	2019	PS mesons simplified
AK, Artru, Martin, PRD 104, 114038 (2021)	2021	PS mesons + VM

Modeling spin-dependent hadronization: the string+ 3P_0 model

- Extension of the Lund string fragmentation model to include the quark spin
- Basic quantity – quark (and antiquark) splitting amplitude

AK, Artru, Martin, PRD 104, 114038 (2021)



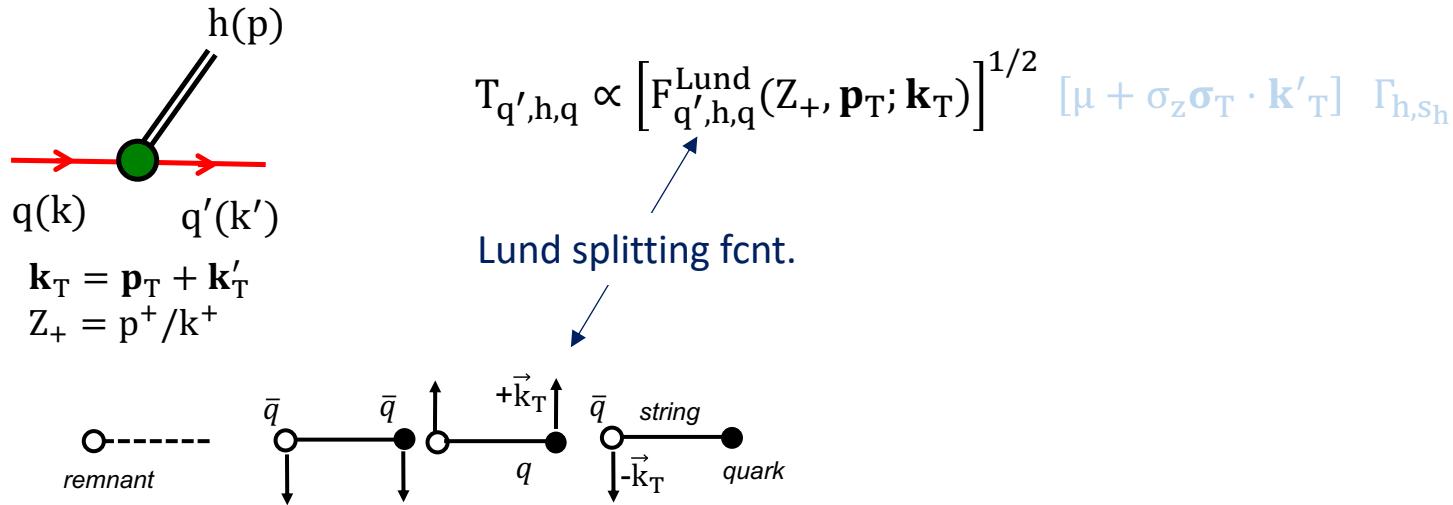
$$\mathbf{k}_T = \mathbf{p}_T + \mathbf{k}'_T$$
$$Z_+ = p^+/k^+$$

$$T_{q',h,q} \propto \left[F_{q',h,q}^{\text{Lund}}(Z_+, \mathbf{p}_T; \mathbf{k}_T) \right]^{1/2} [\mu + \sigma_z \boldsymbol{\sigma}_T \cdot \mathbf{k}'_T] \Gamma_{h,s_h}$$

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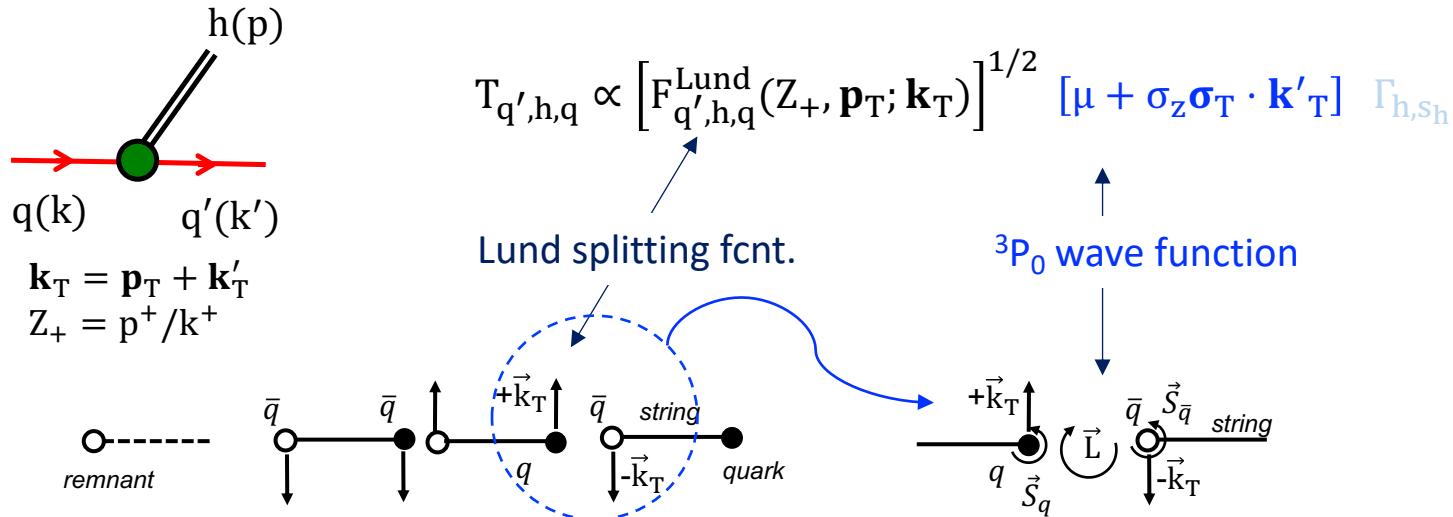
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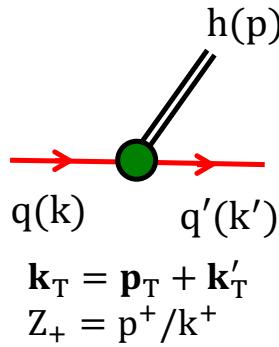
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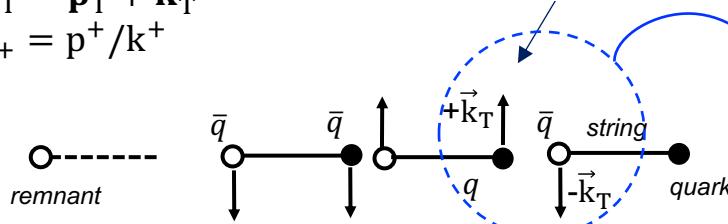
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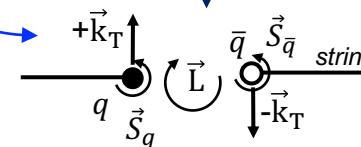


$$T_{q',h,q} \propto [F_{q',h,q}^{\text{Lund}}(Z_+, \mathbf{p}_T; \mathbf{k}_T)]^{1/2} [\mu + \sigma_z \boldsymbol{\sigma}_T \cdot \mathbf{k}'_T]$$

Lund splitting fcnt.



3P_0 wave function



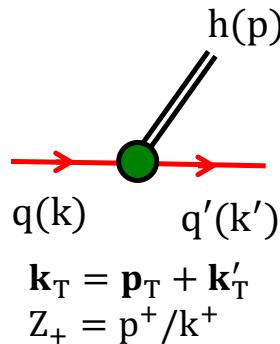
PS meson $\rightarrow \sigma_z$

$$\begin{aligned} & \Gamma_{h,s_h} \\ & \text{coupling} \\ & \text{VM with pol. } \vec{V} \\ & G_L V_L^* 1 + G_T \vec{\sigma}_T \cdot \vec{V}_T^* \sigma_z \end{aligned}$$

Modeling spin-dependent hadronization: the string+ 3P_0 model

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AK, Artru, Martin, PRD 104, 114038 (2021)



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↑
Lund splitting fcnt. ↑
 3P_0 wave function ↑
coupling

Free parameters:

as in Lund Model

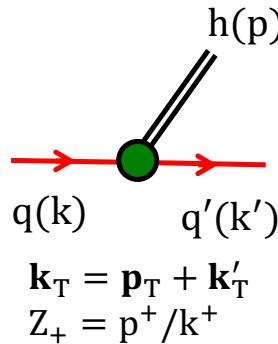
complex «mass»

$\text{Im}(\mu)$ responsible for
transverse spin effects,
e.g. Collins effect

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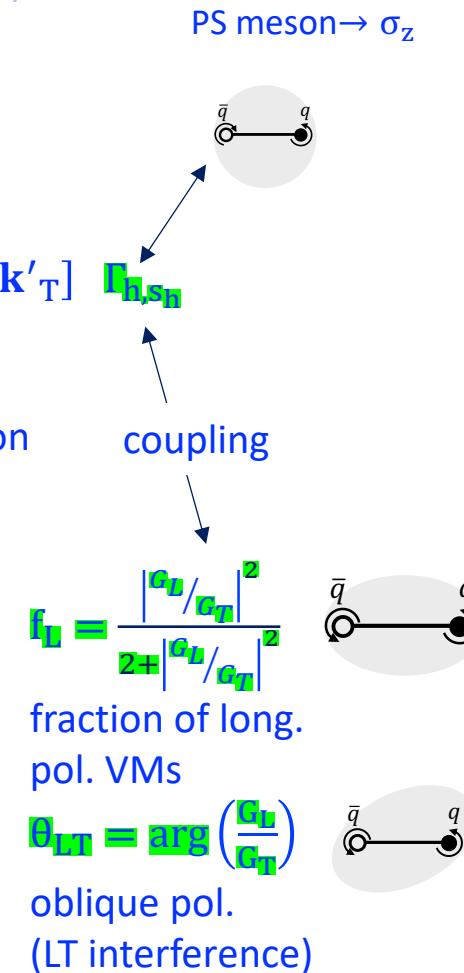
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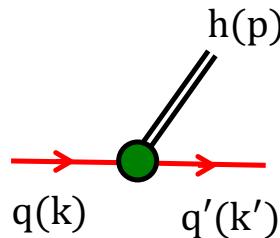
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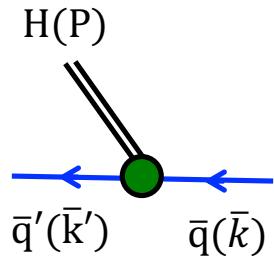
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For anti-quark splitting
 $\{q, h, q'\} \rightarrow \{\bar{q}, H, \bar{q}'\}, Z_+ \rightarrow Z_-, \{\mathbf{k}_T, \mathbf{p}_T, \mathbf{k}'_T\} \rightarrow \{\bar{\mathbf{k}}_T, \bar{\mathbf{P}}_T, \bar{\mathbf{k}}'_T\}$

$$\bar{\mathbf{k}}_T = \bar{\mathbf{P}}_T + \bar{\mathbf{k}}'_T$$

$$Z_- = P^-/\bar{k}^-$$

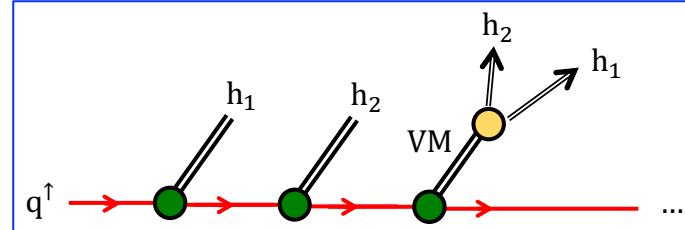
Applications of the string+ 3P_0 model

□ Standalone MC for polarized quark jets

initial studies

detailed study of the model predictions

(Collins asymmetry, dihadron asymmetry, G_1^\perp ..)



□ SIDIS

implemented in Pythia via **StringSpinner**

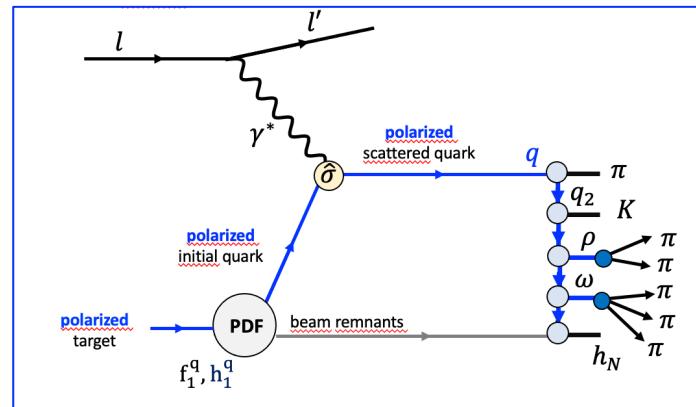
complete collisions

struck quark polarized, remnant unpolarized

parametrization of transversity PDF

reproduces TSA data (Collins, 2h)

[AK, L. Lönnblad, CPC **272** (2022) 108234; CPC **292** (2023) 108886]



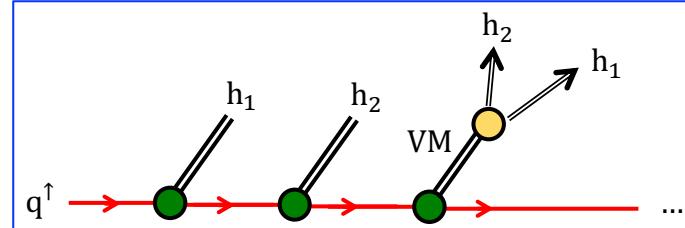
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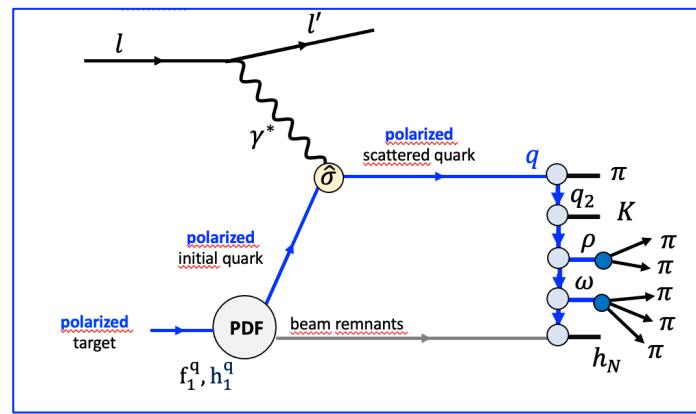
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□ e^+e^- annihilation to hadrons → next slides

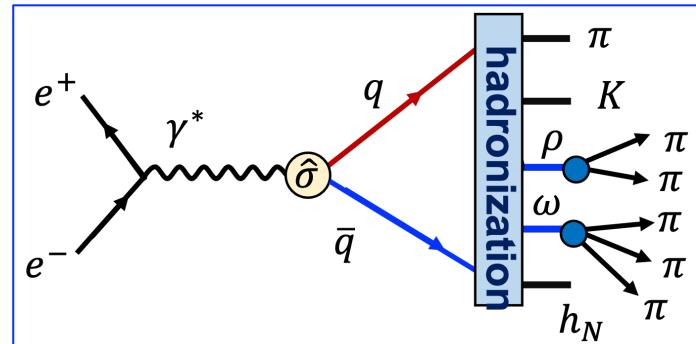
hadronize $q\bar{q}$ accounting for

i. correlated spin states of q and \bar{q}

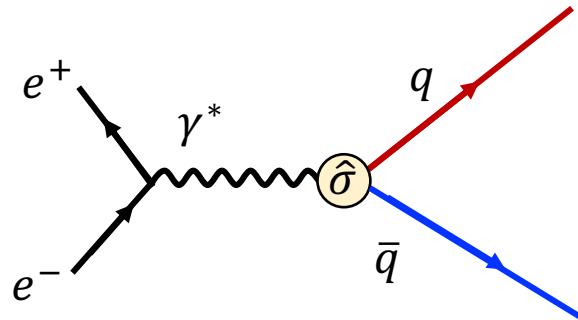
ii. quantum mechanical spin-correlations

in fragmentation chain

AK, X. Artru, PRD 109 (2024) 5, 054029



Recursive recipe for e^+e^-

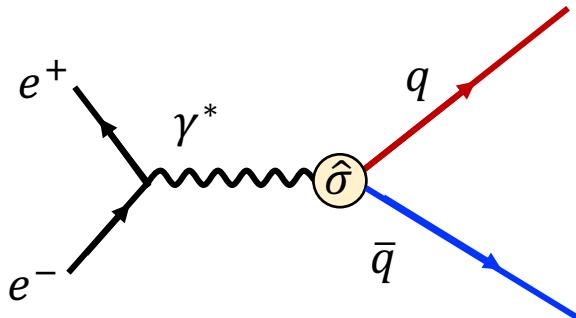


Steps:

1. Hard scattering
2. Joint spin density matrix
3. Hadron emission from q
4. Update density matrix
5. Hadron emission from \bar{q}
6. Exit condition

[AK, X. Artru, PRD 109 (2024) 5, 054029]

Recursive recipe for e^+e^-

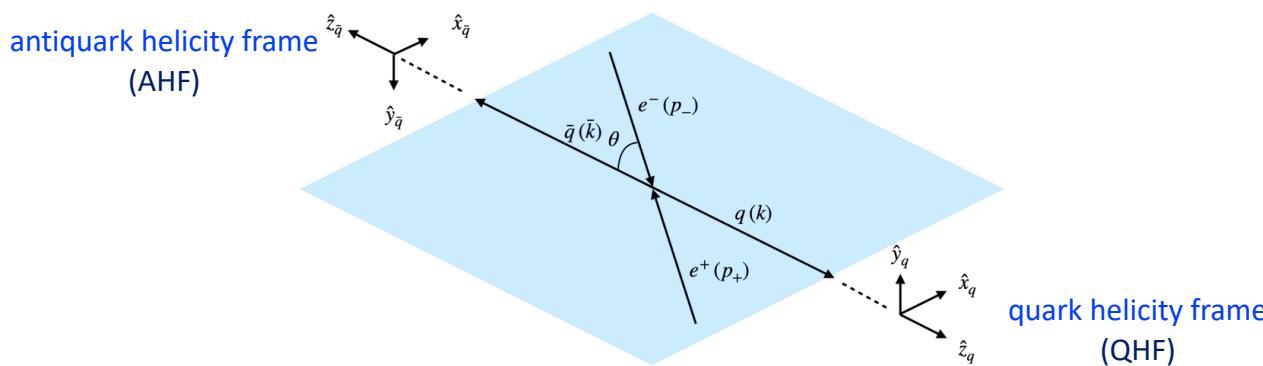


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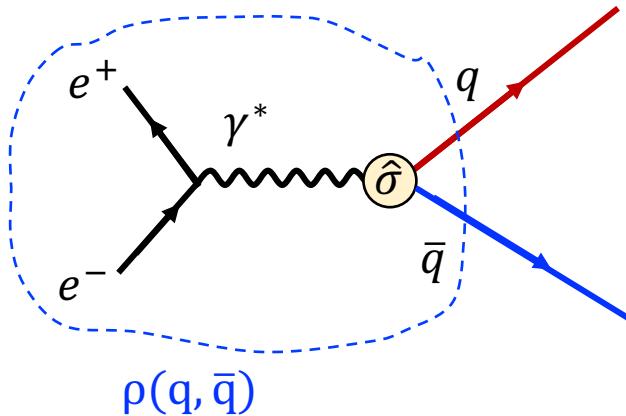
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[AK, X. Artru, PRD 109 (2024) 5, 054029]

Set up the scattering $e^+e^- \rightarrow q\bar{q}$ in the c.m.s
generate the quark flavors and kinematics using differential cross section



Recursive recipe for e^+e^-



- Steps:
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[AK, X. Artru, PRD 109 (2024) 5, 054029]

- Set up the joint spin density matrix of the $q\bar{q}$ pair

$$\rho(q, \bar{q}) = C_{\alpha\beta}^{q\bar{q}} \sigma_q^\alpha \otimes \sigma_{\bar{q}}^\beta$$

correlation coefficients Pauli matrices
 along QHF and AHF

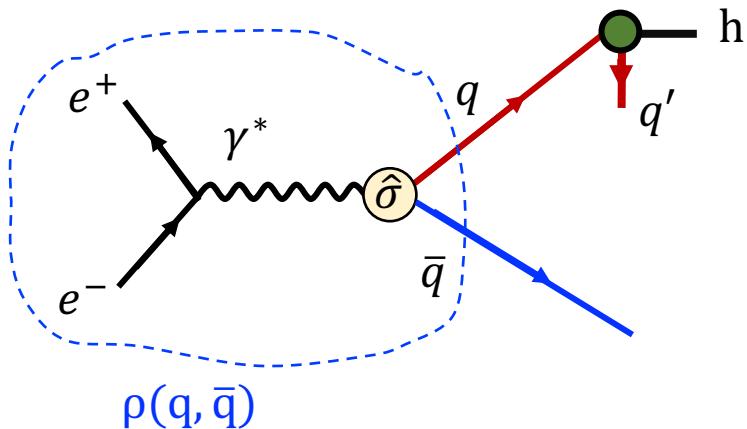
$$\alpha = 0, x_q, y_q, z_q$$

$$\beta = 0, x_{\bar{q}}, y_{\bar{q}}, z_{\bar{q}}$$

For γ^* exchange

$$\rho(q, \bar{q}) \propto 1_q \otimes 1_{\bar{q}} - \sigma_q^z \otimes \sigma_{\bar{q}}^z + \frac{\sin^2 \theta}{1 + \cos^2 \theta} [\sigma_q^x \otimes \sigma_{\bar{q}}^x + \sigma_q^y \otimes \sigma_{\bar{q}}^y]$$

Recursive recipe for e^+e^-



Steps:

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[AK, X. Artru, PRD 109 (2024) 5, 054029]

- Emit the first hadron using the splitting function
(emission probability density)

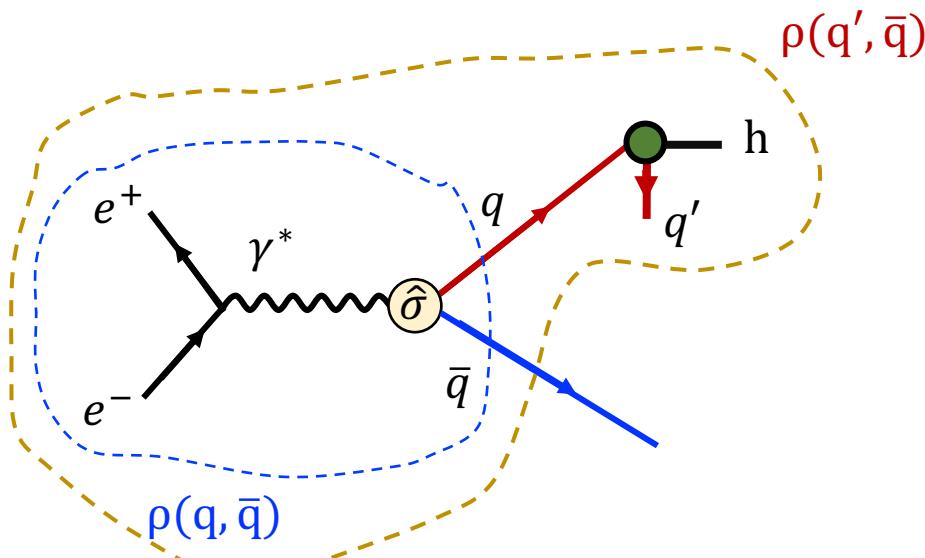
$$\frac{dP(q \rightarrow h + q'; q\bar{q})}{dZ_+ Z_+^{-1} d^2 p_T} = \text{Tr}_{q'\bar{q}} T_{q',h,q} \rho(q, \bar{q}) T_{q',h,q}^\dagger = F_{q',h,q}(Z_+, \mathbf{p}_T; \mathbf{k}_T, C^{q\bar{q}})$$

$$T_{q',h,q} \equiv T_{q',h,q} \otimes 1_{\bar{q}}$$

in the QHF

- VM emission → backup

Recursive recipe for e^+e^-



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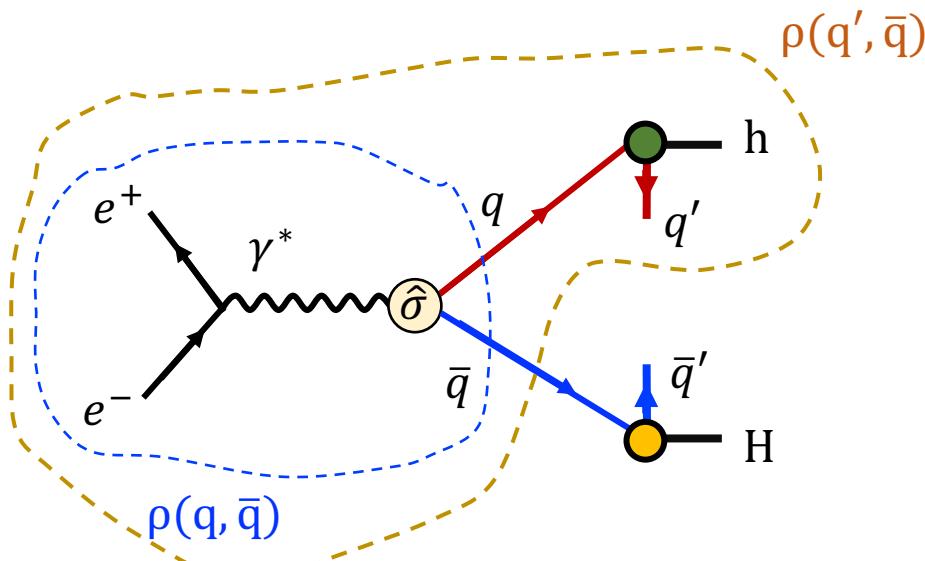
[AK, X. Artru, PRD 109 (2024) 5, 054029]

Evaluate the spin density matrix $\rho(q'\bar{q})$

$$\rho(q', \bar{q}) = T_{q', h, q} \rho(q, \bar{q}) T_{q', h, q}^\dagger$$

includes the information on the emission of h

Recursive recipe for e^+e^-



- Emit a hadron from the \bar{q} side using the splitting function

$$\frac{dP(\bar{q} \rightarrow H + \bar{q}'; q'\bar{q})}{dZ_- Z_-^{-1} d^2 P_T} = \text{Tr}_{q'\bar{q}'} T_{\bar{q}', H, \bar{q}} \rho(q', \bar{q}) T_{\bar{q}', H, \bar{q}}^\dagger = F_{\bar{q}', H, \bar{q}}(Z_-, P_T; \bar{k}_T, C^{q'\bar{q}})$$

conditional probability of emitting H , having emitted h
 → correlations between the transverse momenta

[Collins NPB, 304:794–804, 1988, Knowles NPB, 310:571–588, 1988]

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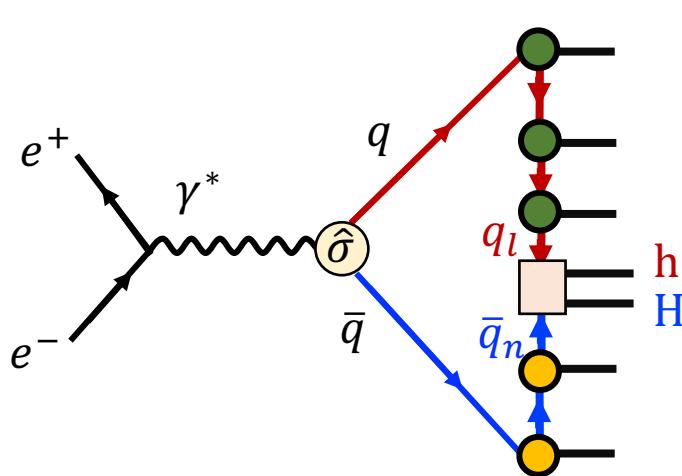
[AK, X. Artru, PRD 109 (2024) 5, 054029]

Depend on the azimuthal angle h

↓

Expressed in the AHF

Recursive recipe for e^+e^-



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[AK, X. Artru, PRD 109 (2024) 5, 054029]

- Iterate until the exit condition is called and the last quark pair is hadronized
more details in PRD 109 (2024) 5, 054029

Simulations of e^+e^- with spin effects

- ❑ Now possible in Pythia 8.3
 - StringSpinner package extended
- ❑ Photon exchange, only uds quarks produced, $\sqrt{s} = 10.6$ GeV
 - consistent with BELLE and BABAR data

Simulations of e^+e^- with spin effects

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StringSpinner package extended
- ❑ Photon exchange, only uds quarks produced, $\sqrt{s} = 10.6$ GeV
consistent with BELLE and BABAR data
- ❑ Free parameters
 - spin-less hadronization as in standard Pythia 8.3
 - spin-dependent hadronization
 - complex mass μ as in AK, Lonnblad, CPC 292 (2023) 108886
 - $f_L = 0.12$ $\sim T$ pol. VMs
 - $\theta_{LT} = -0.65$ interference between T and L pol. of VMs
- found to give a satisfactory agreement with e^+e^- data,
ok also for SIDIS
- ❑ Comparison with Collins asymmetries from BELLE and BABAR
AK, L. Lönnblad, A. Martin, in preparation

The A_{12} asymmetry

- Reference plane formed by the thrust axis $\hat{\mathbf{n}}$ and the beam e-back-to-back h_1 and h_2

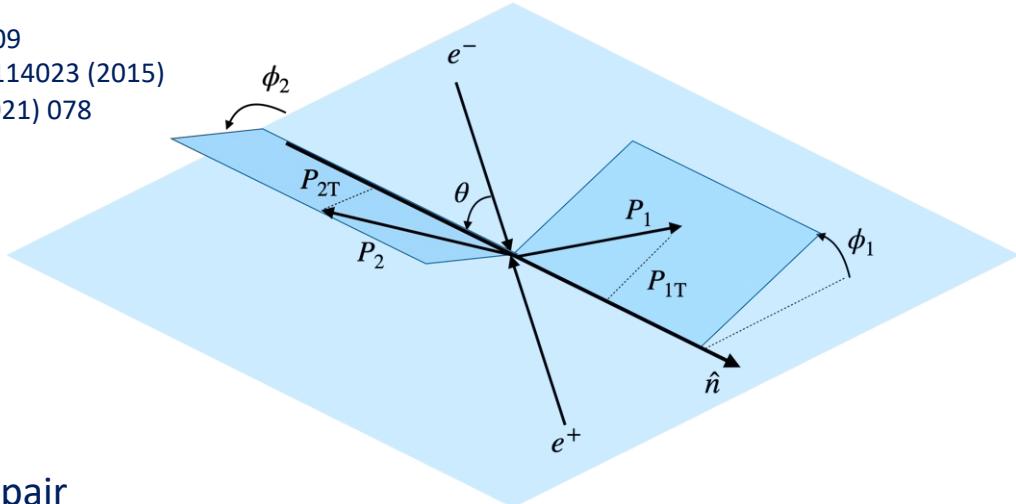
$$N_{h_1 h_2} \propto 1 + \frac{\langle \sin^2 \theta \rangle}{\langle 1 + \cos^2 \theta \rangle} A_{12} \cos(\phi_1 + \phi_2)$$

Boer, NPB, 806:23–67, 2009
 Anselmino et al., PRD 92, 114023 (2015)
 D'Alesio et al., JHEP 10 (2021) 078

Collins asymmetry

$$A_{12} = \frac{\sum_q e_q^2 H_{1q}^{\perp h_1} H_{1\bar{q}}^{\perp h_2}}{\sum_q e_q^2 D_{1q}^{h_1} D_{1\bar{q}}^{h_2}}$$

$$z_i = \frac{2E_{h_i}}{\sqrt{s}}$$



- Experimental asymmetry

$$R_{12}^{U,L,C} = N_{12}^{U,L,C} / \langle N_{12}^{U,L,C} \rangle$$

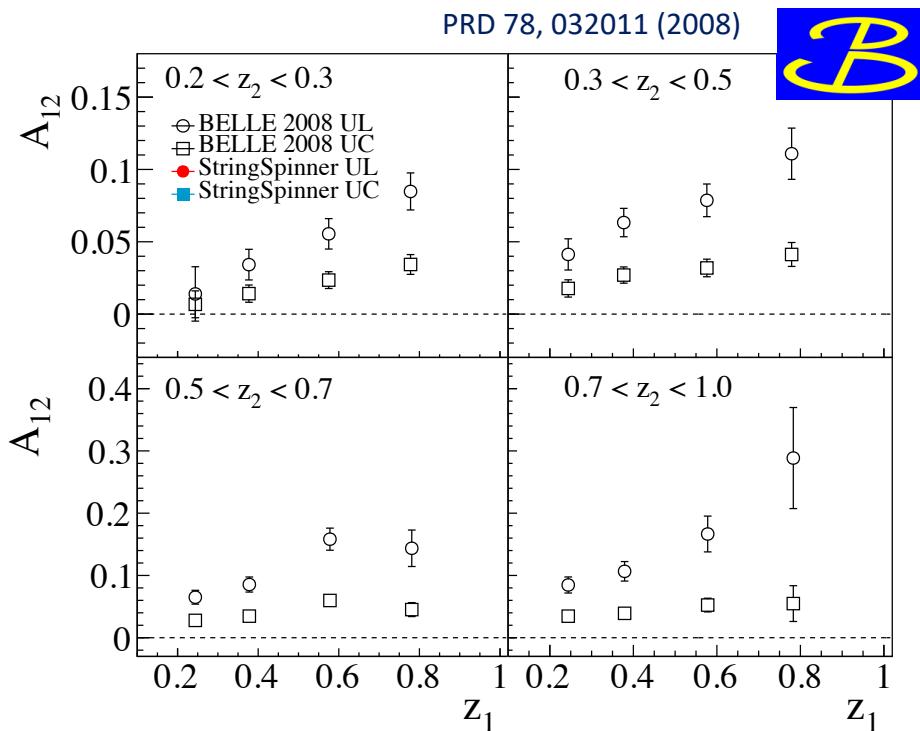
U = unlike sign pair
 L = like sign pair
 C = charged pair

$$R_{12}^{UL(UC)} = \frac{R_{12}^U}{R_{12}^{L(C)}} \approx 1 + \frac{\langle \sin^2 \theta \rangle}{\langle 1 + \cos^2 \theta \rangle} A_{12}^{UL(UC)} \cos(\phi_1 + \phi_2)$$

$$A_{12}^{UL(UC)} \simeq A_{12}^U - A_{12}^{L(C)}$$

- Measured by BELLE (2008, 2019) and BABAR (2014, 2015)

A_{12} asymmetry for charged $\pi \pi$ pairs

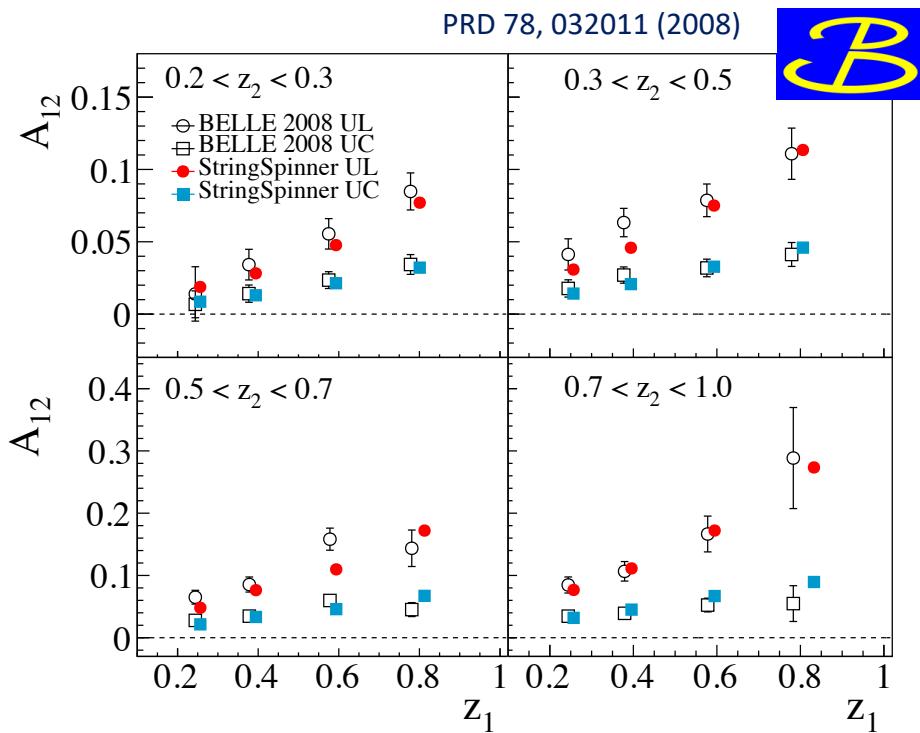


Belle asymmetries corrected for thrust smearing

Cuts:

$T > 0.8, z > 0.2, Q_T < 3.5 \text{ GeV}$

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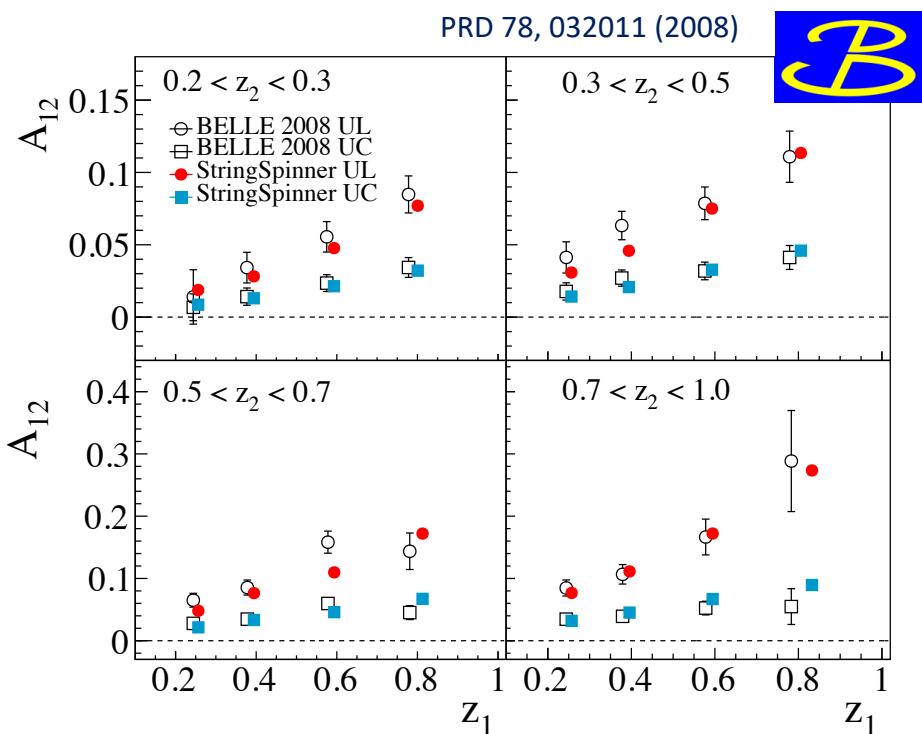
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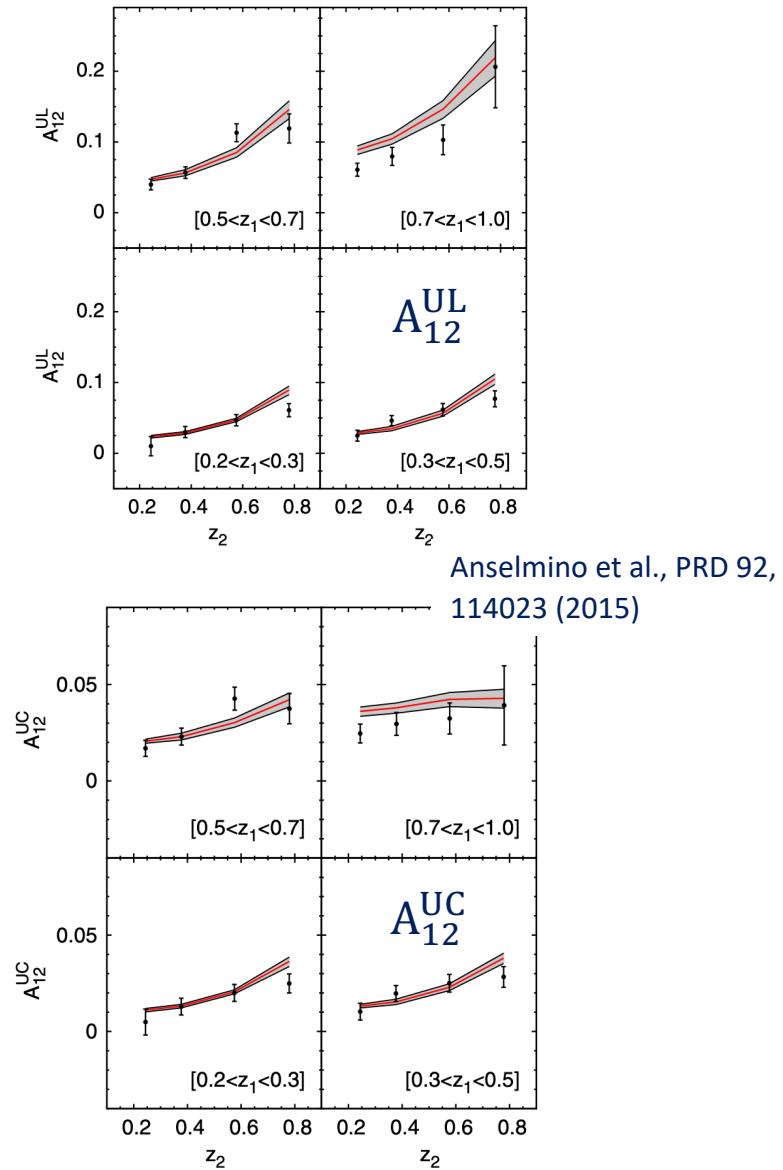
StringSpinner reproduces trend and size

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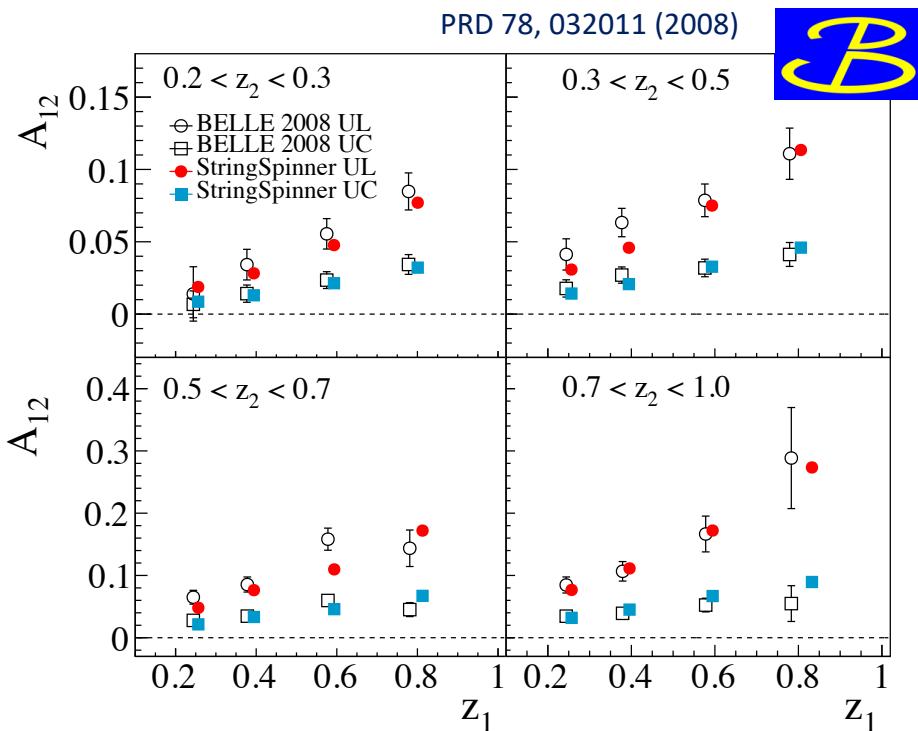


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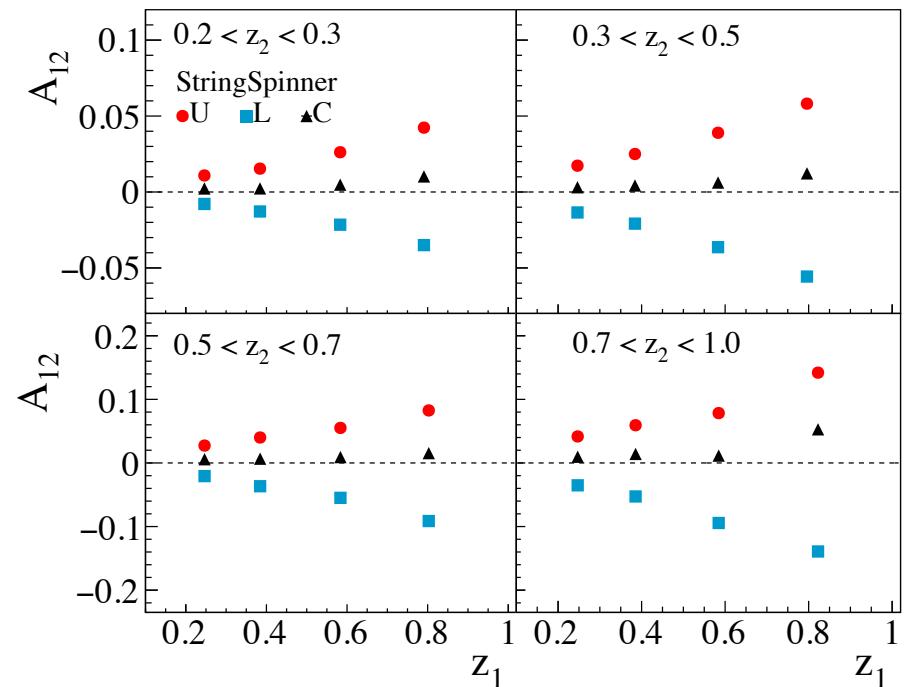
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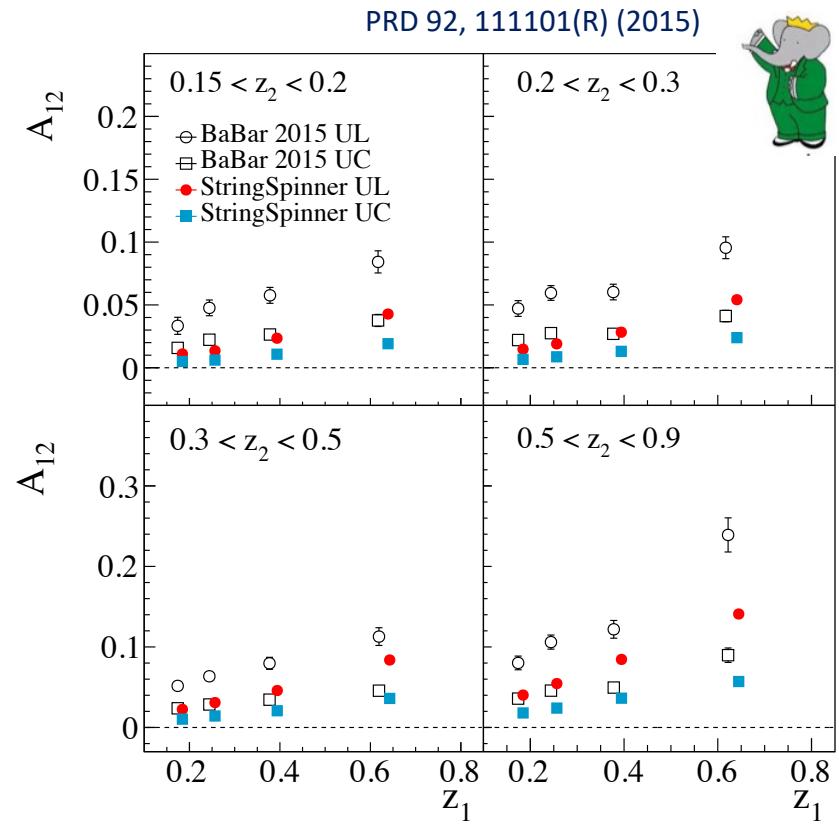
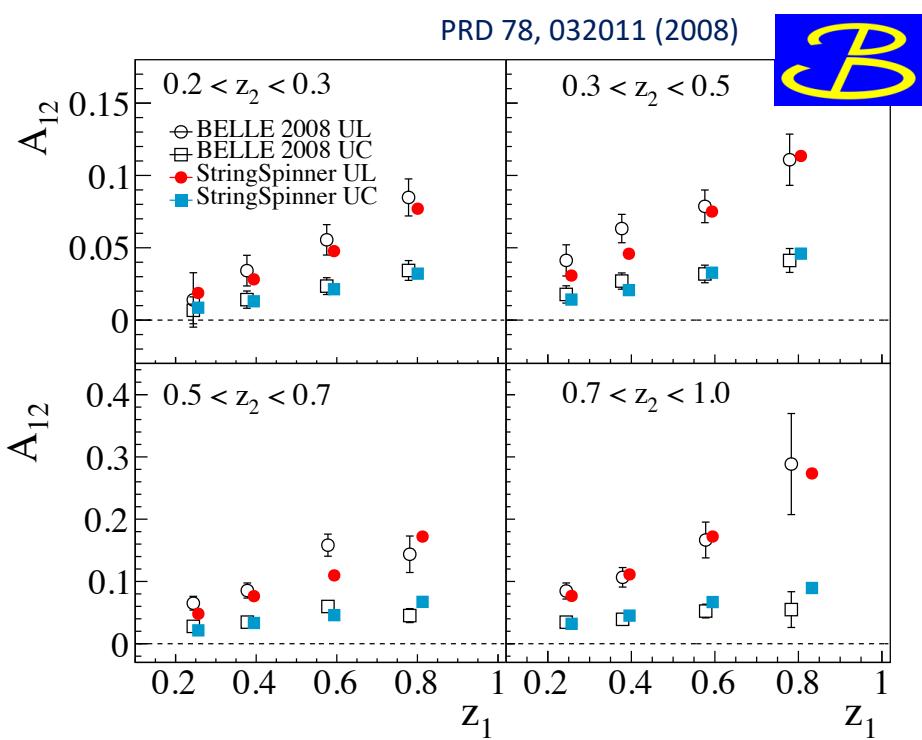


Opposite sign for A_{12}^U and A_{12}^L , as expected

$$A_{12}^C \simeq 0 \rightarrow \text{explains } A_{12}^{UC} < A_{12}^{UL}$$

$$A_{12}^{UL(UC)} \simeq A_{12}^U - A_{12}^{L(C)}$$

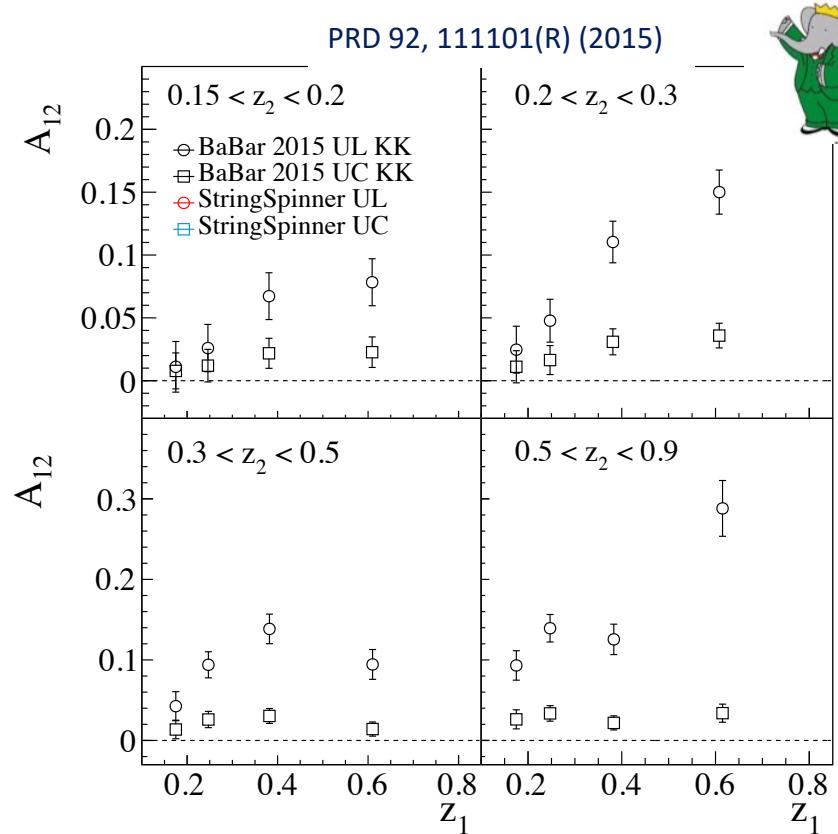
A_{12} asymmetry for charged $\pi \pi$ pairs



StringSpinner lower than BABAR
difference between BABAR and BELLE known

PRD 90, 052003 (2014)

A_{12} asymmetry for charged KK pairs

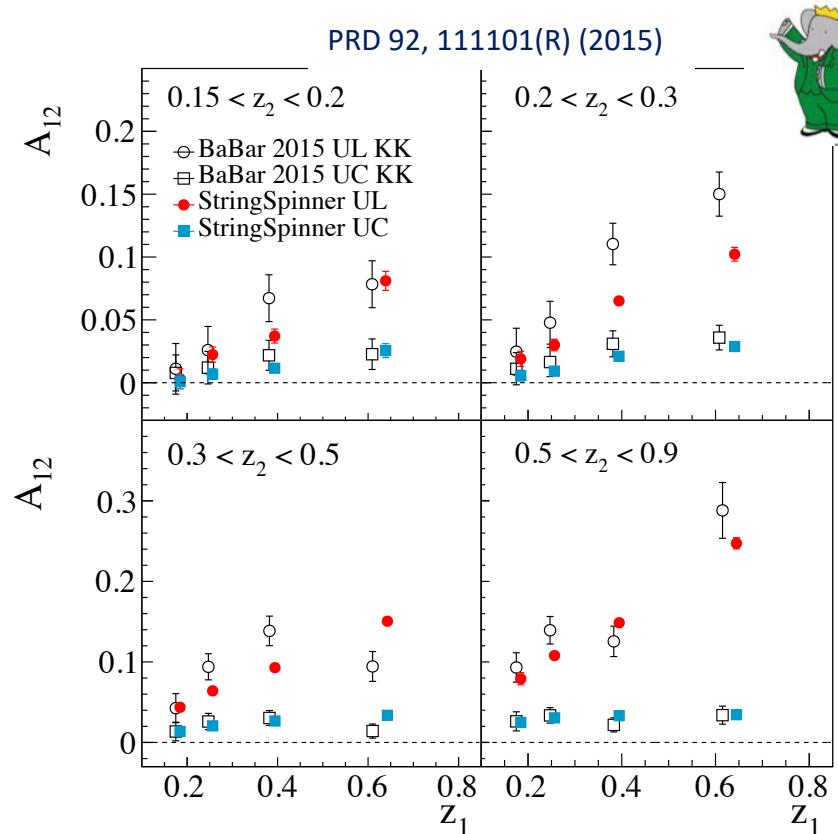


Corrected for thrust smearing

Cuts

$T > 0.8, z > 0.15, Q_T < 3.5 \text{ GeV}, \alpha_0 < \pi/4$

A_{12} asymmetry for charged KK pairs



Corrected for thrust smearing

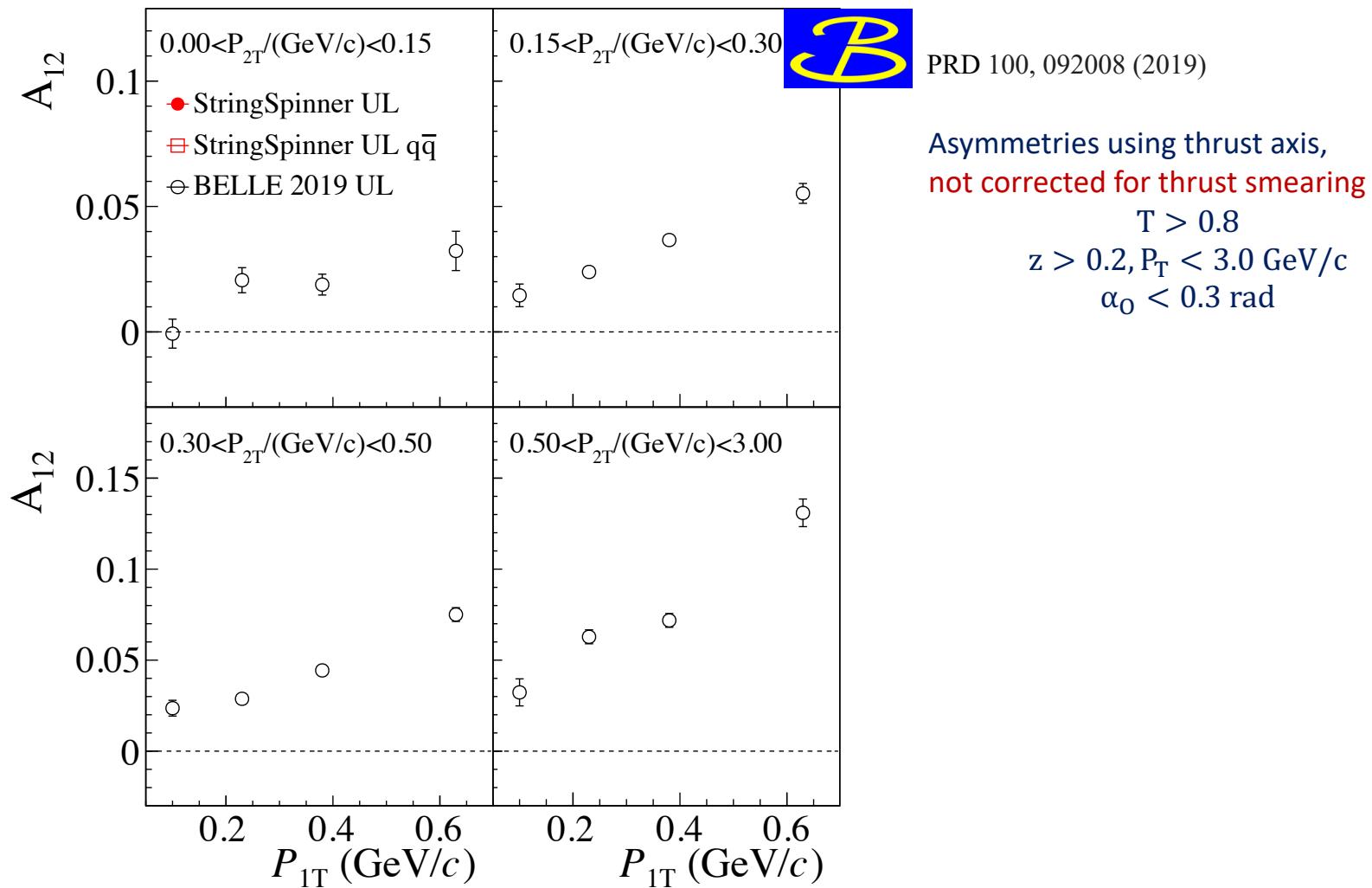
Cuts

$T > 0.8, z > 0.15, Q_T < 3.5 \text{ GeV}, \alpha_0 < \pi/4$

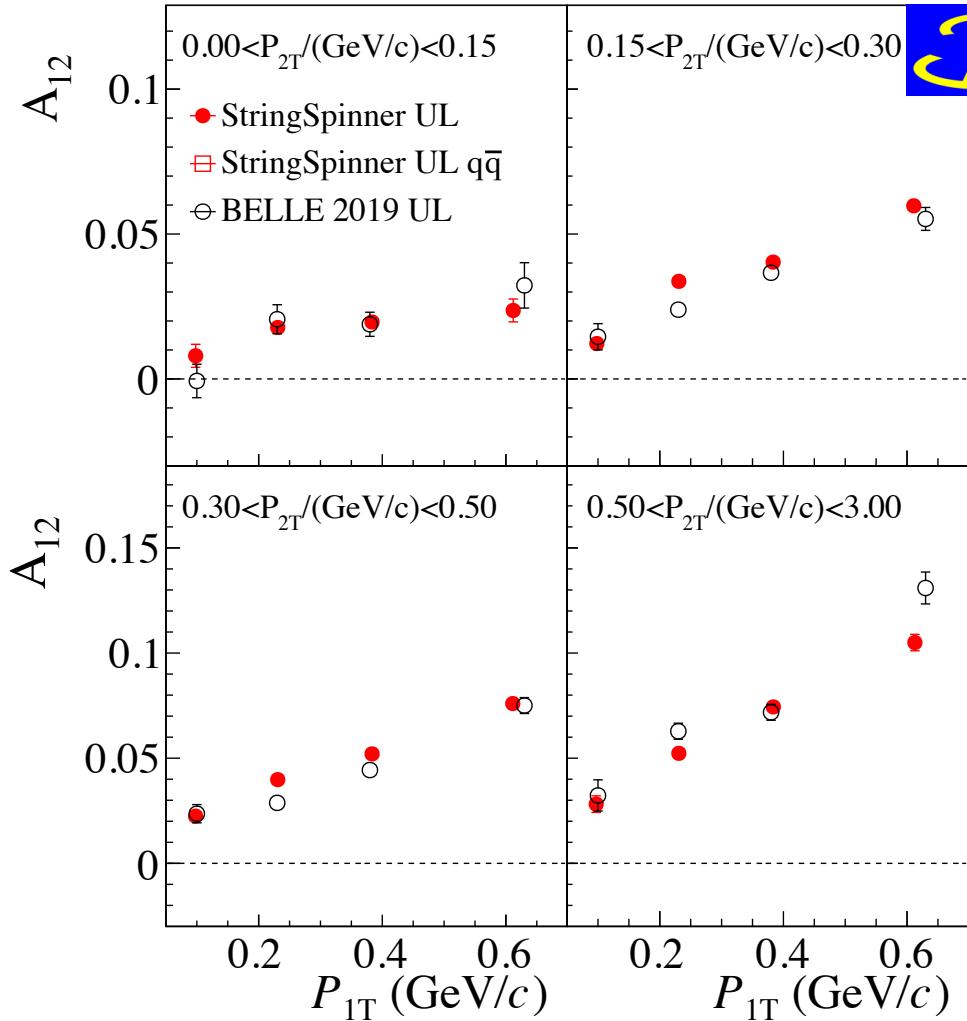
A_{12}^{UC} much smaller than A_{12}^{UL} at large z
reproduced by string+ 3P_0

A_{12}^{UL} asymmetry for charged $\pi \pi$ pairs

$P_{\text{T}1} \times P_{\text{T}2}$ - dependence w.r.t thrust



A_{12}^{UL} asymmetry for charged $\pi \pi$ pairs $P_{\text{T}1} \times P_{\text{T}2}$ - dependence w.r.t thrust



PRD 100, 092008 (2019)

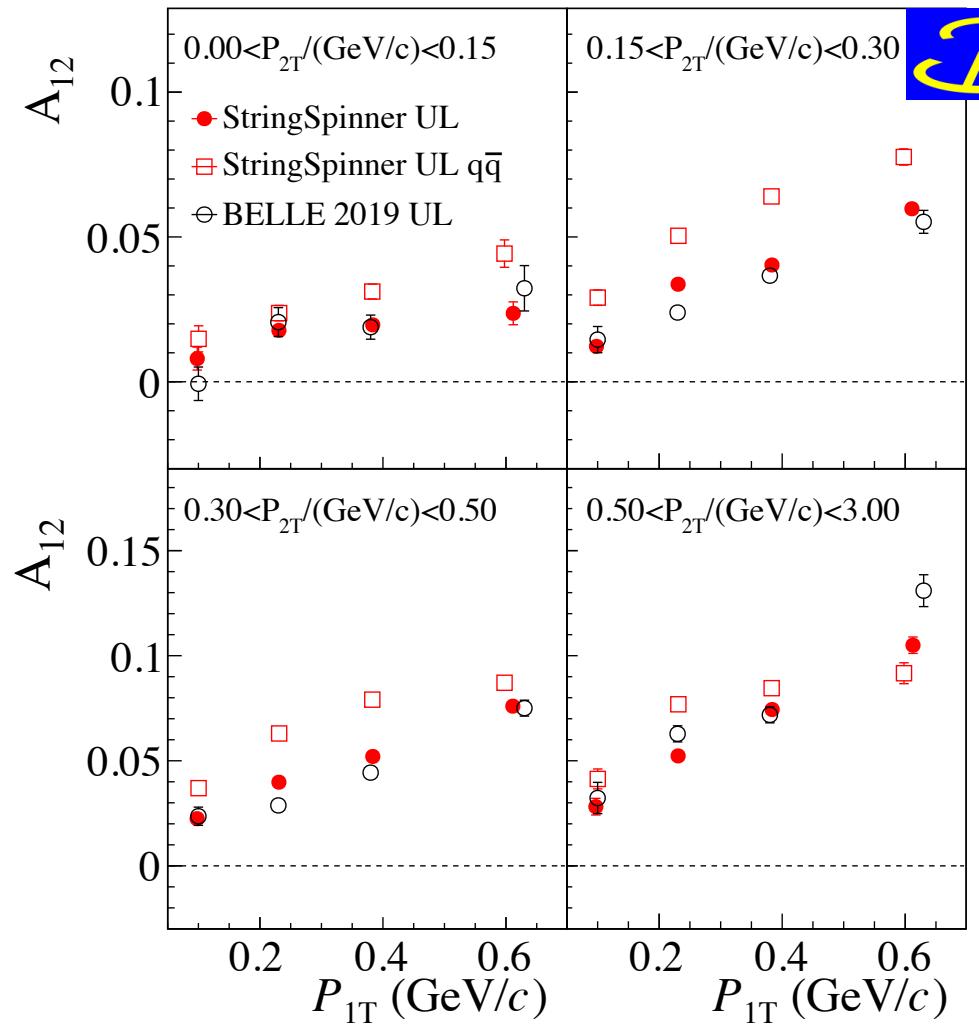
Asymmetries using thrust axis,
not corrected for thrust smearing

$T > 0.8$
 $z > 0.2, P_{\text{T}} < 3.0 \text{ GeV}/c$
 $\alpha_0 < 0.3 \text{ rad}$

StringSpinner reproduces the nearly linear trend observed by BELLE

A_{12}^{UL} asymmetry for charged $\pi \pi$ pairs

$P_{\text{T}1} \times P_{\text{T}2}$ - dependence w.r.t thrust



PRD 100, 092008 (2019)

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$T > 0.8$
 $z > 0.2, P_{\text{T}} < 3.0 \text{ GeV}/c$
 $\alpha_0 < 0.3 \text{ rad}$

StringSpinner reproduces the nearly linear trend observed by BELLE

Linear trend shows up as an effect of the misalignment between thrust and $q\bar{q}$ axis
strong effect by thrust

The A_0 asymmetry

- Reference plane formed by the momentum \mathbf{p}_2 of h_2 and the beam e^-

$$N_0 \propto 1 + \frac{\langle \sin^2 \theta_2 \rangle}{\langle 1 + \cos^2 \theta_2 \rangle} A_0 \cos(2\phi_0)$$

Boer, NPB, 806:23–67, 2009
 Anselmino et al., PRD 92, 114023 (2015)
 D'Alesio et al., JHEP 10 (2021) 078
 ..

Collins asymmetry

$$A_0 = \frac{\sum_q e_q^2 w H_{1q}^{\perp h_1} \otimes H_{1\bar{q}}^{\perp h_2}}{\sum_q e_q^2 D_{1q}^{h_1} \otimes D_{1\bar{q}}^{h_2}}$$

- Experimental asymmetry

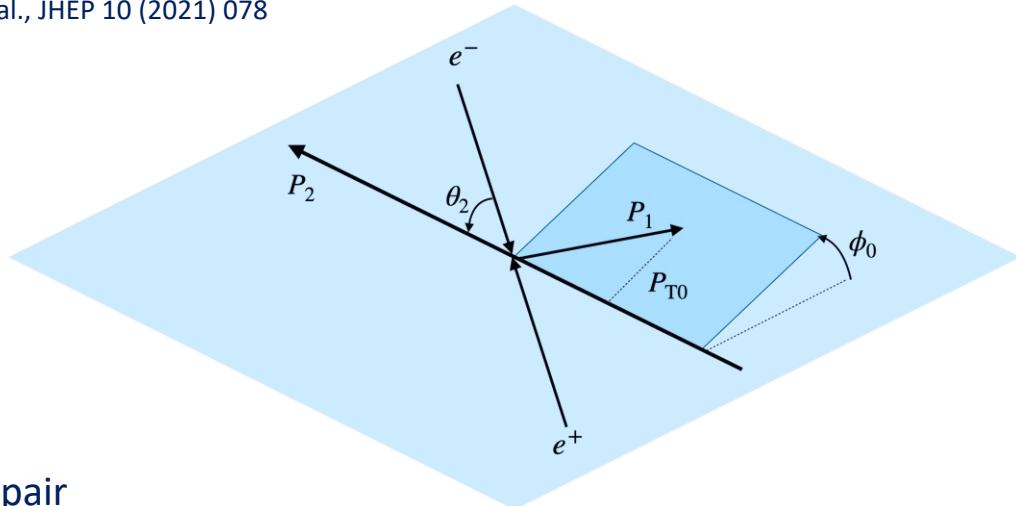
$$R_0^{U,L,C} = N_0^{U,L,C} / \langle N_0^{U,L,C} \rangle \quad U = \text{unlike sign pair}$$

$L = \text{like sign pair}$

$C = \text{charged pair}$

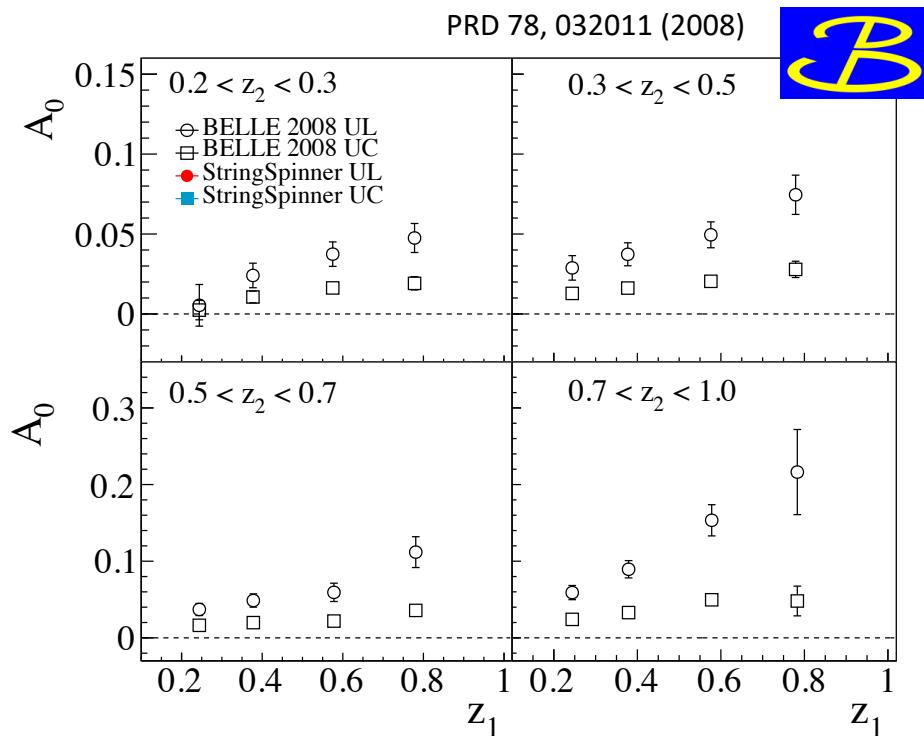
$$R_0^{UL(UC)} = \frac{R_0^U}{R_0^{L(C)}} \approx 1 + \frac{\langle \sin^2 \theta_2 \rangle}{\langle 1 + \cos^2 \theta_2 \rangle} A_0^{UL(UC)} \cos(2\phi_0)$$

$$A_0^{UL(UC)} \simeq A_0^U - A_0^{L(C)}$$



- Measured by BELLE (2008) and BABAR (2014, 2015)

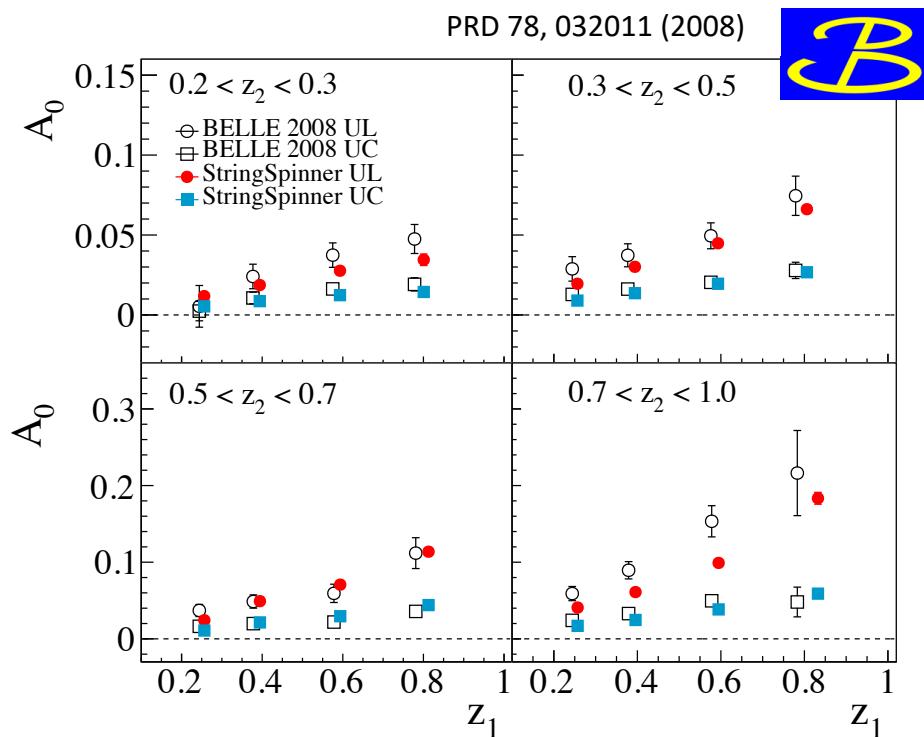
A_0 asymmetry for charged pions



Cuts:

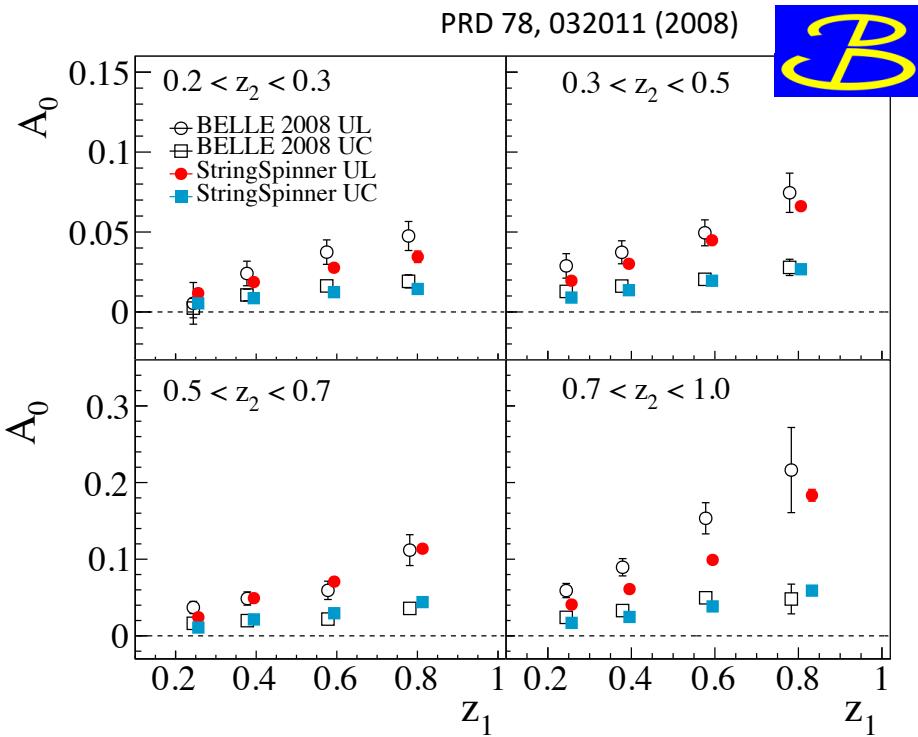
$T > 0.8, z > 0.2, Q_T < 3.5 \text{ GeV}$

A_0 asymmetry for charged pions

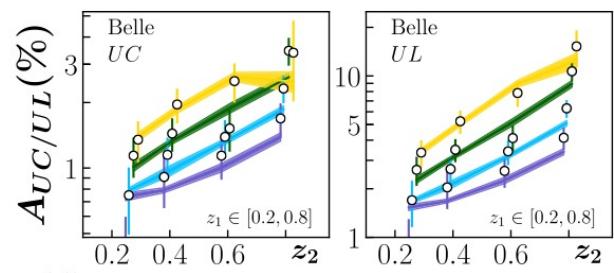
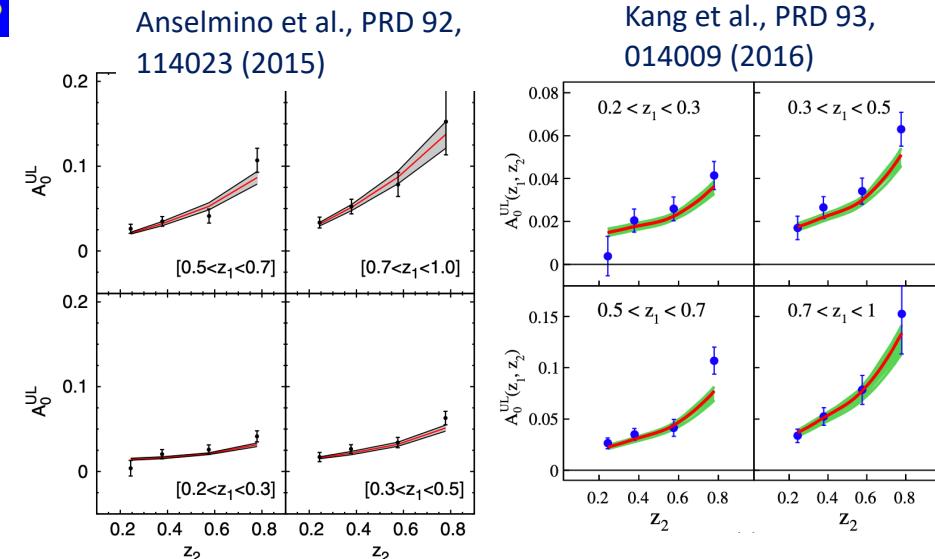


Trend reproduced by string+ 3P_0
somewhat lower values in the last z_2 bin

A_0 asymmetry for charged pions



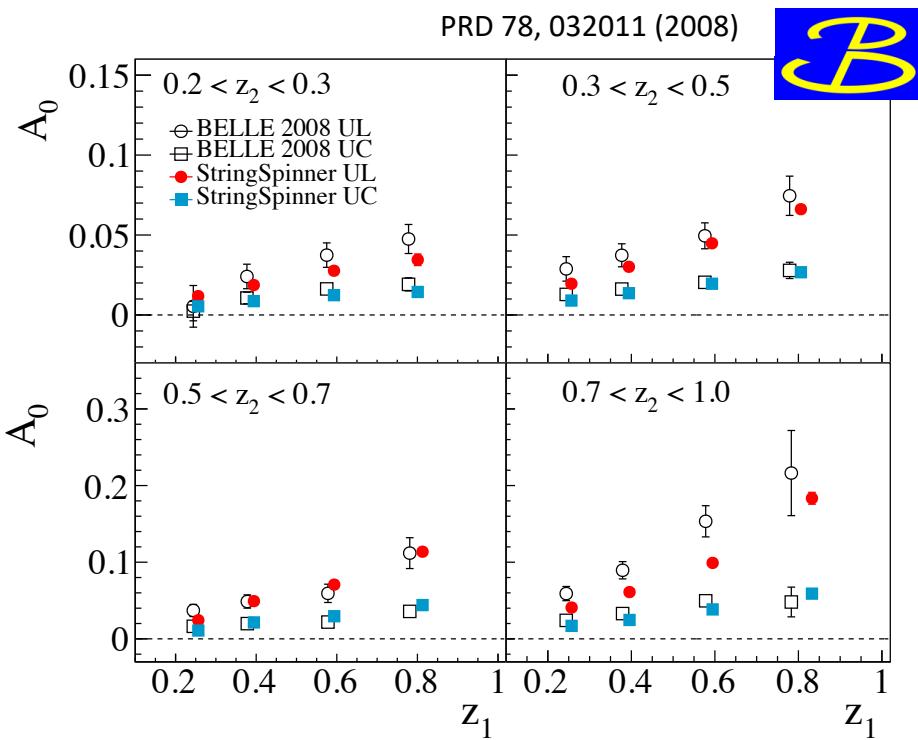
Trend reproduced by string+ 3P_0
 somewhat lower values in the last z_2 bin



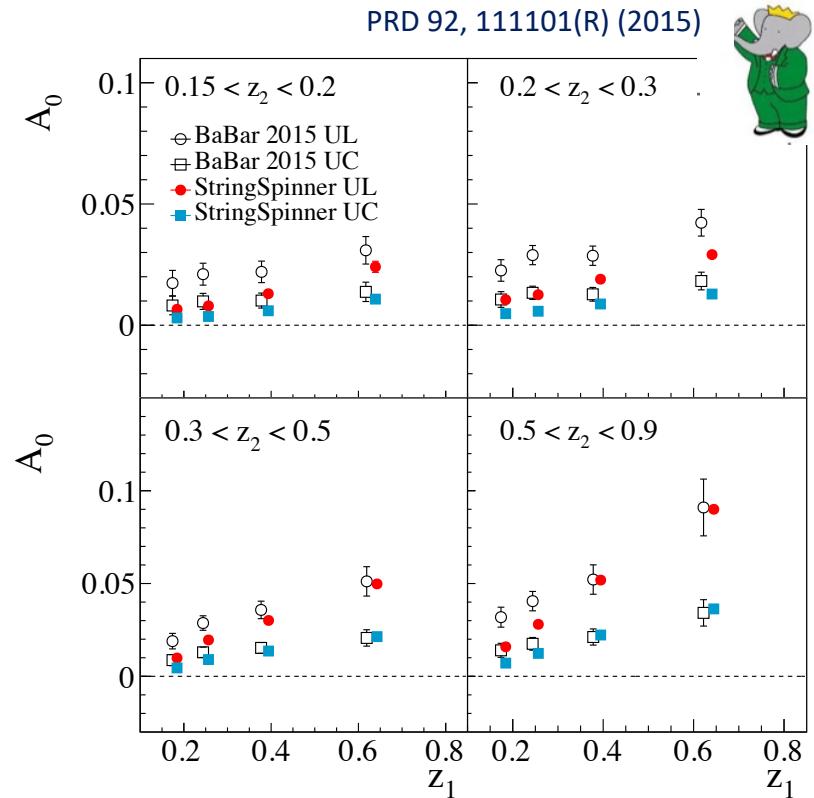
JAM, PRD 102, 054002 (2020)

A_0 asymmetry included in
 phenomenological fits

A_0 asymmetry for charged pions



Trend reproduced by string+ 3P_0
somewhat lower values in the last z_2 bin



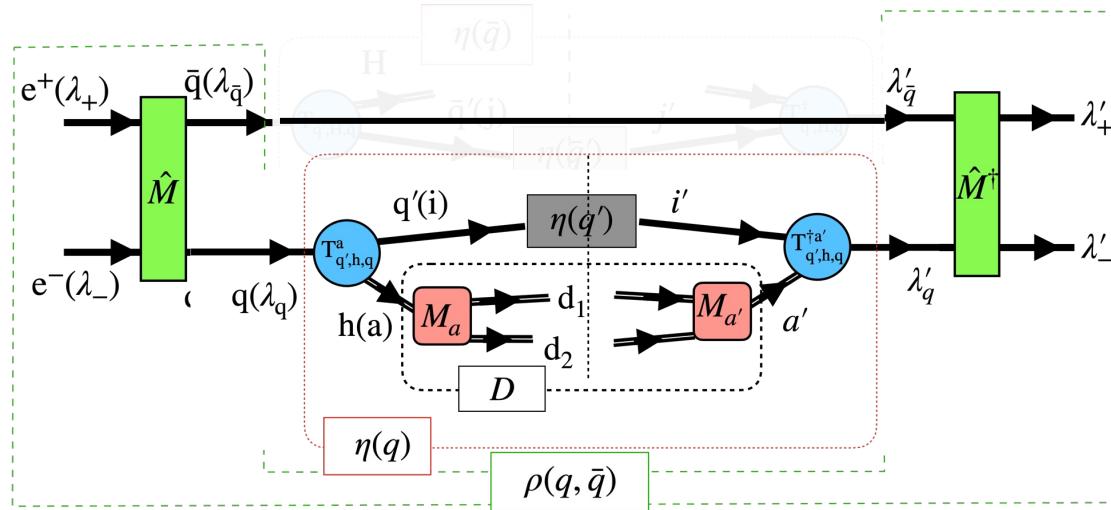
Conclusions

- ❑ The string+ 3P_0 model applied to the fragmentation of a string stretched between quarks with correlated spin states
 - these kind of strings are produced also in pp collisions..
- ❑ Implemented in Pythia 8.3 for e^+e^-
 - extension of StringSpinner to be published
- ❑ Encouraging results on Collins asymmetries in e^+e^-
 - study of Artru-Collins asymmetries ongoing
- ❑ More developments of the string+ 3P_0 model foreseen
 - baryon production
 - connection to parton shower
 - ...

Still a long way to an event generator fully implementing spin effects, but several steps done!

Backup

The recursive recipe for simulating e^+e^- annihilation: VM emission



For a vector meson $h=VM$

$$\rightarrow \eta(q) = T_{q',h=VM,q}^{a'\dagger} \eta(q') T_{q',h=VM,q}^a D_{a'a}, \quad \eta(q') = 1_{q'}, \text{ and } \eta(\bar{q}) = 1_{\bar{q}}$$

Steps:

i) Emission probability density (summing over decay information, i.e. $D_{a'a} = \delta_{a'a}$)

$$\frac{dP(q \rightarrow h = VM + q'; q\bar{q})}{dM^2 dZ_+ Z_+^{-1} d^2 p_T} = \text{Tr}_{q'\bar{q}} T_{q',h,q}^a \rho(q, \bar{q}) T_{q',h,q}^{a\dagger} = F_{q',h,q}(M^2, Z_+, p_T; k_T, C^{q\bar{q}})$$

ii) Calculate the spin density matrix of $h=VM$, and decay the meson

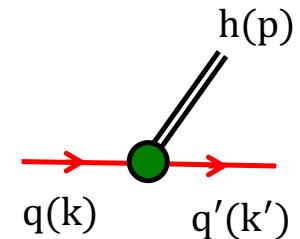
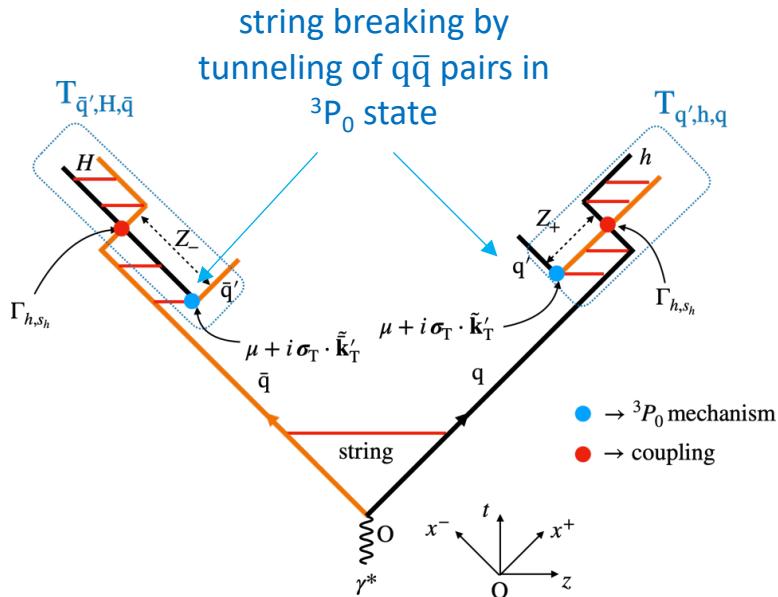
$$\rho_{aa'}(h) = \text{Tr}_{q'\bar{q}} T_{q',h,q}^a \rho(q, \bar{q}) T_{q',h,q}^{a\dagger}$$

iii) Decay the meson $p \rightarrow p_1 p_2 ..$

$$dN(p_1, p_2, \dots) / d\Omega \propto M_{\text{dec.}}^a(p \rightarrow p_1 p_2, \dots) \rho_{aa'}(h) M_{\text{dec.}}^{a\dagger a'}(p \rightarrow p_1 p_2, \dots)$$

iv) Build the decay matrix $D_{a'a}(p_1, p_2, \dots) = M_{\text{dec.}}^{a\dagger a'}(p \rightarrow p_1 p_2, \dots) M_{\text{dec.}}^a(p \rightarrow p_1 p_2, \dots)$

The hadronization model: string+ 3P_0



quark splitting $q \rightarrow h + q'$

Relevant variables:
 $\mathbf{k}_T = \mathbf{p}_T + \mathbf{k}'_T$
 $Z_+ = p^+/k^+$
 $\varepsilon_h^2 = M^2 + p_T^2$

Transverse vectors
defined w.r.t. string axis

Quark splitting amplitude in the string+ 3P_0 model

$$T_{q'hq} \propto C_{q'hq} D_h(M^2) \left(\frac{1 - Z_+}{\varepsilon_h^2} \right)^{\frac{a}{2}} \underbrace{\exp \left[-\frac{b_L \varepsilon_h^2}{2Z_+} \right]}_{\text{longitudinal momentum}} N_a^{-\frac{1}{2}}(\varepsilon_h^2) e^{-\frac{b_T k'^2_T}{2}} \underbrace{\varepsilon_h^2}_{\text{transverse momentum (w.r.t string axis)}}$$

flavor mass

Free param. Lund

Free param. string+ 3P_0

$[u + \sigma_z \sigma_T \cdot \mathbf{k}'_T]$
 3P_0 mechanism
 $[\mu \text{ complex mass parameter}]$

Γ_{hsh}
Coupling
e.g.
 $\Gamma_{hPS} = \sigma_z$

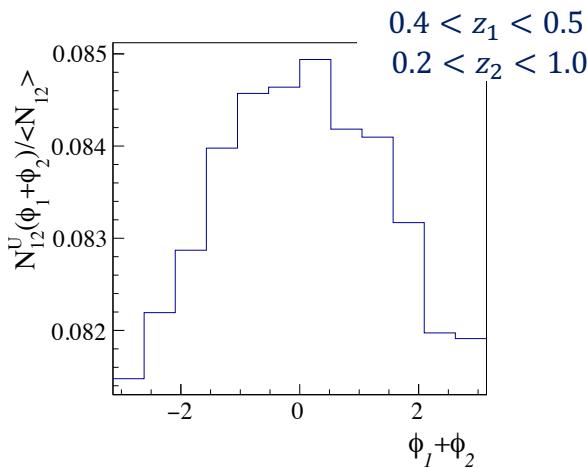
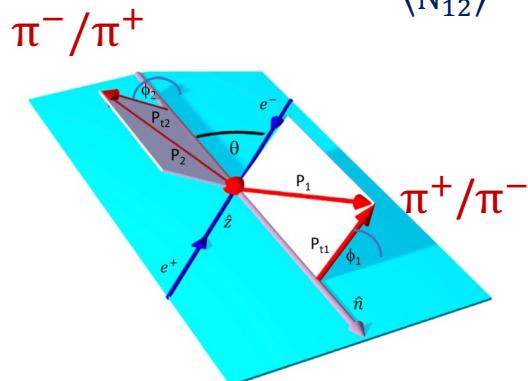
AK, Artru, Martin, PRD 104, 114038 (2021)

Steps for the extraction of Collins asymmetries

Example of $e^+e^- \rightarrow \pi^+\pi^-X$

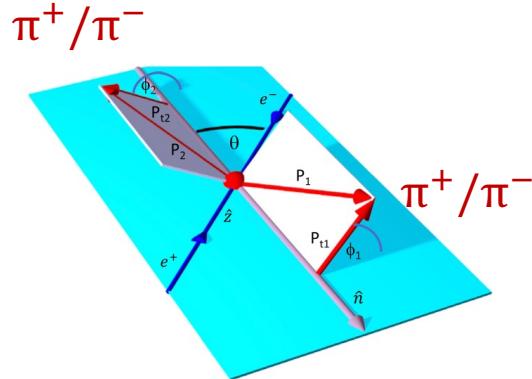
i) Evaluate normalized yields for
 $\pi^\pm - \pi^\mp$ "Unlike pairs"

$$R_{12}^U = \frac{N_{12}^U(\phi_1 + \phi_2)}{\langle N_{12} \rangle}$$



ii) Evaluate normalized yields for
 $\pi^\pm - \pi^\pm$ "Like pairs"

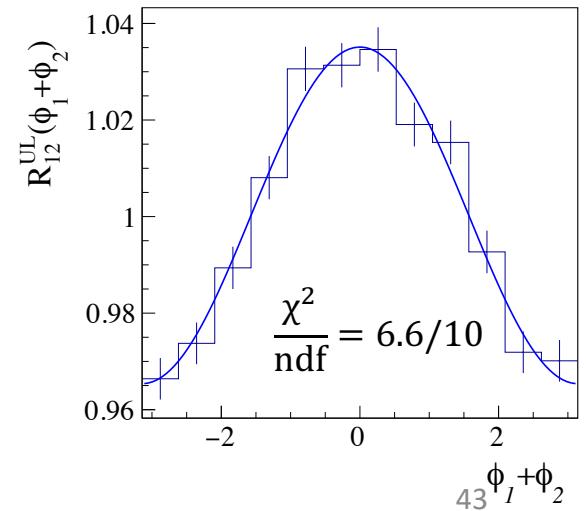
$$R_{12}^L = \frac{N_{12}^L(\phi_1 + \phi_2)}{\langle N_{12} \rangle}$$



ii) Evaluate the ratio $\frac{R_{12}^U}{R_{12}^L}$
and fit the asymmetry

$$R_{12}^{UL} = \frac{R_{12}^U}{R_{12}^L} \approx 1 + A_{12}^{UL} \cos(\phi_1 + \phi_2)$$

Fit function
 $f(\phi_1 + \phi_2) = p_0 + p_1 \cos(\phi_1 + \phi_2)$



A_{12}^{UL} asymmetry for back-to-back $\pi^\pm - \pi^\mp$

$z_1 \times z_2$ - dependence

