State of the art of observables for Generalized TMDs

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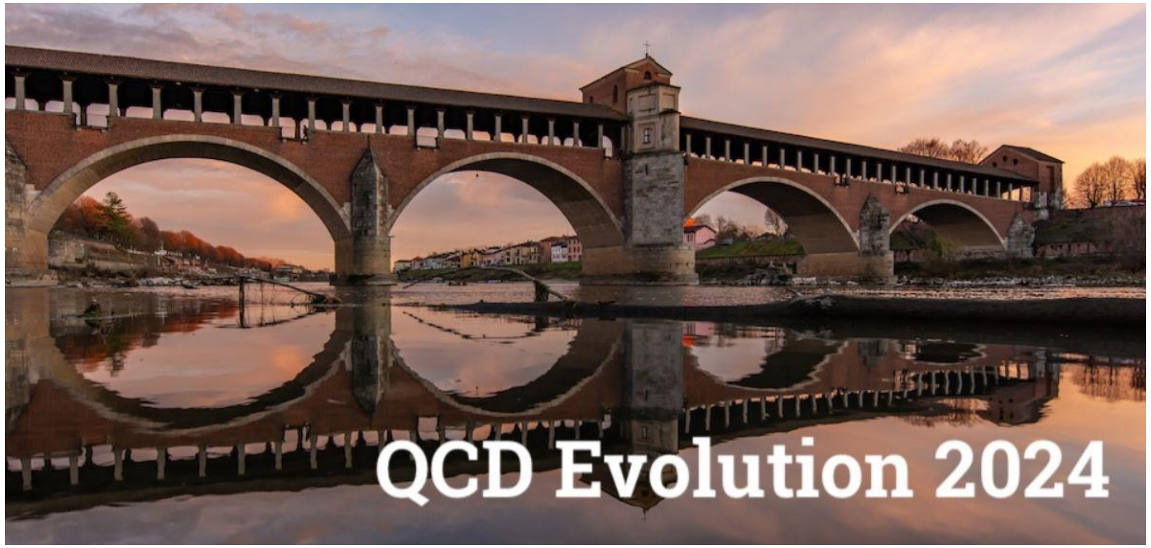
May 30, 2024



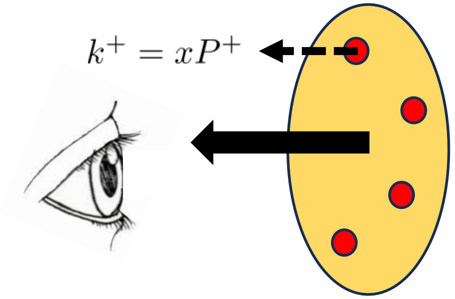
In Collaboration with:

Duxin Zheng, Jian Zhou (arXiv: 2312.01309)

Renaud Boussarie, Yoshitaka Hatta (arXiv: 2201.08709, 2404.04208, 2404.04209)

Wigner function - The “mother function”

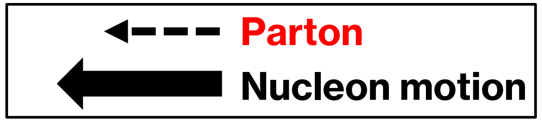
Parton
Nucleon motion

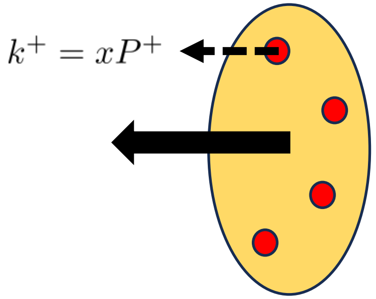


Parton Distribution Functions

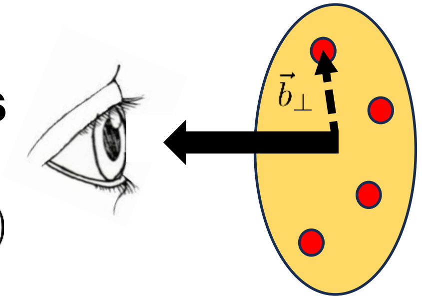
PDFs (x)

Wigner function - The "mother function"





Form Factors



PDFs (x)

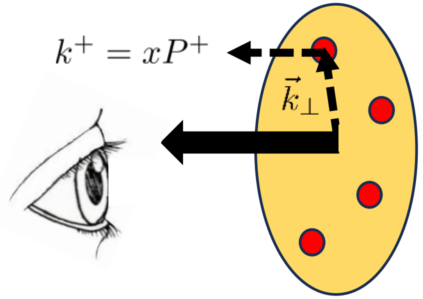
FFs (Δ)

Wigner function - The “mother function”

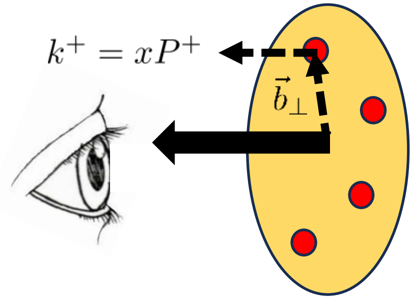
Parton

Nucleon motion

Transverse Momentum-dependent Distributions



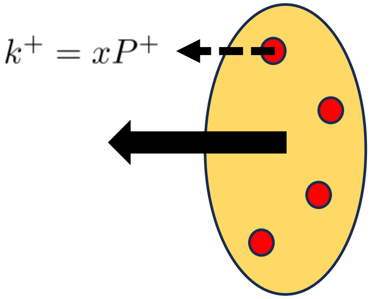
Generalized Parton Distributions

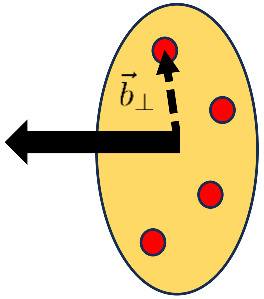


TMDs $(x, \vec{k}_\perp)$ $\xrightarrow{\int d^2\vec{k}_\perp}$ **PDFs** (x)

GPDs (x, Δ) $\xrightarrow{\Delta = 0}$ **PDFs** (x)

GPDs (x, Δ) $\xrightarrow{\int dx}$ **FFs** (Δ)





Wigner function - The “mother function”

Generalized Transverse Momentum-dependent Distributions

(Meissner, Metz, Schlegel, 2009)

$$
\begin{tikzcd}
 & \textbf{GTMDs}\ (x, \vec{k}_\perp, \Delta) \arrow[dl, "\Delta = 0"'] \arrow[dr, "\int d^2\vec{k}_\perp"] & & \\
\textbf{TMDs}\ (x, \vec{k}_\perp) \arrow[dr, "\int d^2\vec{k}_\perp"'] & & \textbf{GPDs}\ (x, \Delta) \arrow[dl, "\Delta = 0"'] \arrow[dr, "\int dx"] & \\
 & \textbf{PDFs}\ (x) & & \textbf{FFs}\ (\Delta)
\end{tikzcd}
$$

Wigner function - The “mother function”

Wigner functions $(x, \vec{k}_\perp, \vec{b}_\perp)$ (Belitsky, Ji, Yuan, 2003)

2-D Fourier Transform

$(\vec{\Delta}_\perp)$

(Meissner, Metz, Schlegel, 2009) **GTMDs** $(x, \vec{k}_\perp, \Delta)$

$\Delta = 0$

$\int d^2\vec{k}_\perp$

TMDs $(x, \vec{k}_\perp)$

GPDs (x, Δ)

$\int d^2\vec{k}_\perp$

$\Delta = 0$

$\int dx$

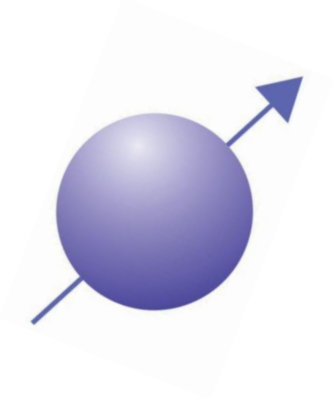
PDFs (x)

FFs (Δ)

Spin of proton

Jaffe-Manohar spin decomposition

An incomplete story:

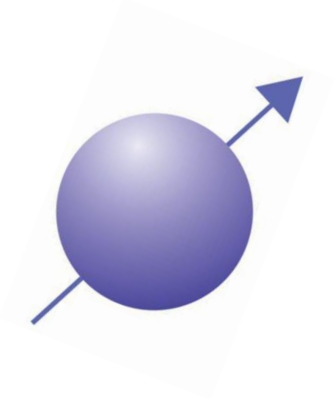


$$\frac{1}{2}$$

Spin of proton

Jaffe-Manohar spin decomposition

An incomplete story:



$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G$$

Best known — **How well do we know?**

Quark helicity $\sim 30\%$ — **Gluon helicity** $\sim 40\%$

Spin of proton

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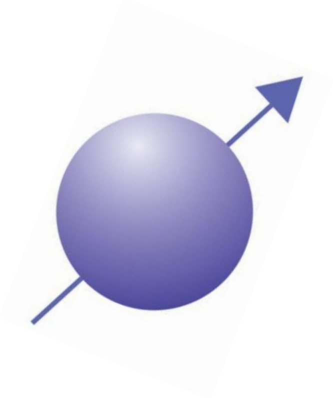

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L^q + L^g$$

Best known | **How well do we know?** | **???????**

Quark helicity $\sim 30\%$ | **Gluon helicity** $\sim 40\%$ | **OAM of quarks & gluons**

Wigner functions & Orbital Angular Momentum

Wigner functions in Quantum Mechanics

(Wigner, 1932)

- Calculate from wave functions:

$$W(x,k) = \int \frac{dx'}{2\pi} e^{-ikx'} \psi(x+\frac{x'}{2})\, \psi^*(x-\frac{x'}{2})$$

- Expectation value of observables:

$$\langle \mathcal{O} \rangle = \int dx \int dk \mathcal{O}(x,k) W(x,k)$$

Wigner functions & Orbital Angular Momentum

Wigner functions in Quantum Mechanics **(Wigner, 1932)**	**Wigner functions in parton physics** **(Belitsky, Ji, Yuan, 2003)**
• Calculate from wave functions: $W(x,k) = \int \frac{dx'}{2\pi} e^{-ikx'} \psi(x+\frac{x'}{2})\,\psi^*(x-\frac{x'}{2})$ • Expectation value of observables: $\langle \mathcal{O} \rangle = \int dx \int dk \mathcal{O}(x,k) W(x,k)$	• Calculate from fourier transform of GTMD correlator: $W^{[\Gamma]}(x,\vec{k}_\perp,\vec{b}_\perp)$ • Application: **O**rbital **A**ngular **M**omentum (**OAM**) $L_z^{q,g} = \int dx \int d^2k_\perp d^2b_\perp (\vec{b}_\perp \times \vec{k}_\perp)_z \, W^{q,g} \; (x,\vec{b}_\perp,\vec{k}_\perp)$ (Lorcé, Pasquini, 2011 / Hatta, 2011)

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Wigner functions in parton physics

(Belitsky, Ji, Yuan, 2003)

- Calculate from fourier transform of GTMD correlator:

$$W^{[\Gamma]}(x, \vec{k}_\perp, \vec{b}_\perp)$$

- Application: Relation between GTMD $F_{1,4}^{q,g}$ & OAM

$$\boxed{L_z^{q,g} = -\int dx \int d^2\vec{k}_\perp \frac{\vec{k}_\perp^2}{M^2} F_{1,4}^{q,g}(x, k_\perp, \xi = 0, \Delta_\perp = 0)}$$

(Lorcé, Pasquini, 2011 / Hatta, 2011)

Wigner functions & Orbital Angular Momentum

Wigner functions in Quantum Mechanics (Wigner, 1932)	Wigner functions in parton physics (Belitsky, Ji, Yuan, 2003)

- Calculate from wav

$$W(x,k) = \int \frac{dx'}{2\pi}$$

- Expectation value of observables:

$$\langle \mathcal{O} \rangle = \int dx \int dk \mathcal{O}(x,k) W(x,k)$$

Big question:

Experimental observable?

- of GTMD correlator:

- Application: Relation between GTMD $F_{1,4}^{q,g}$ & OAM

$$\boxed{L_z^{q,g} = -\int dx \int d^2\vec{k}_\perp \frac{\vec{k}_\perp^2}{M^2} F_{1,4}^{q,g}(x, k_\perp, \xi = 0, \Delta_\perp = 0)}$$

(Lorcé, Pasquini, 2011 / Hatta, 2011)

Developments

arXiv: 1612.02438 (2016) Feng Yuan’s talk

Hunting the Gluon Orbital Angular Momentum at the Electron-Ion Collider

Xiangdong Ji,[1,2] Feng Yuan,[3] and Yong Zhao[1,3]

Developments

arXiv: 1612.02438 (2016)

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Exclusive double quarkonium production and generalized TM

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Probing the Weizsäcker-Williams gluon Wigner distribution in *pp* collisions

Renaud Boussarie,[1] Yoshitaka Hatta,[2] Bo-Wen Xiao,[3,4] and Feng Yuan[5]

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Probing the gluon Sivers function with an unpolarized target: GTMD distributions and the Odderons

Renaud Boussarie,[1] Yoshitaka Hatta,[1] Lech Szymanowski,[2] and Samuel Wallon[3,4]

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Signature of the gluon orbital angular momentum

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Angular correlations in exclusive dijet photoproduction in ultra-peripheral PbPb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

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arXiv: 1802.10550 (2018)

Exclusive double quarkonium production and generalized TM

Exclusive double Drell-Yan:

Until now, this has been the sole known process sensitive to quark GTMDs

tion in pp collisions

d Feng Yuan[5]

ction of di-pions

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Probing quark OAM through double Drell-Yan

Main findings

arXiv: 1702.04387 (2017)

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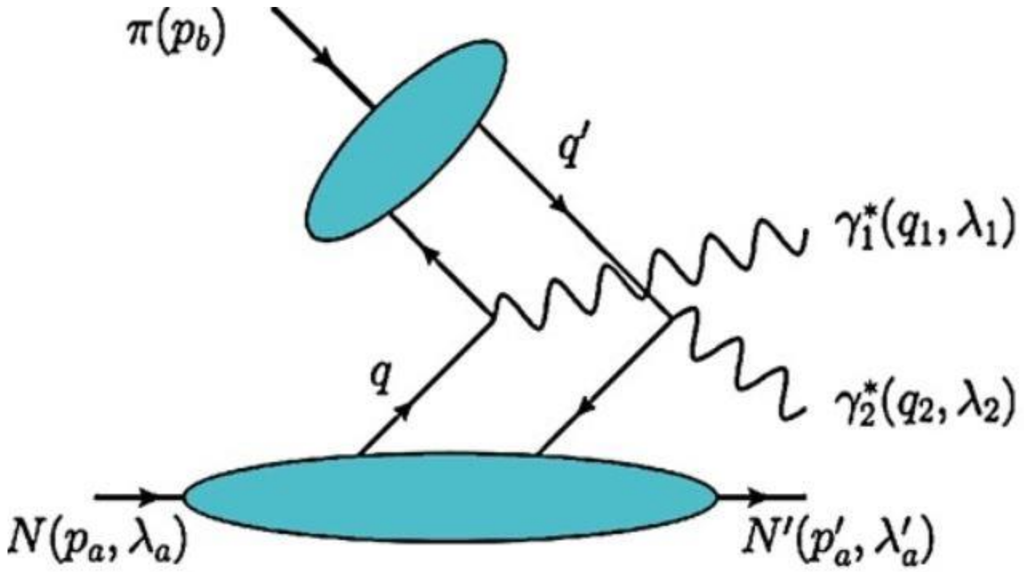

Probing quark OAM through double Drell-Yan

Main findings

Example of an observable sensitive to OAM & spin-orbit correlation :

$$\frac{1}{2}\left(\tau_{XY} - \tau_{YX}\right) = \frac{4}{M_a^2}\,(\varepsilon_\perp^{ij}\Delta q_\perp^i \Delta_{a\perp}^j)\,\mathrm{Re.}\left\{C^{(-)}\left[F_{1,1}\,\phi_\pi\right]C^{(+)}\left[\vec{\beta}_\perp \cdot \vec{k}_{a\perp}\,\boldsymbol{F}^*_{\mathbf{1,4}}\,\phi^*_\pi\right] \right\}$$

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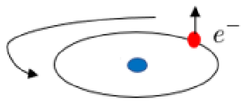
Spin-orbit entanglement in the Color Glass Condensate

Shohini Bhattacharya,[1,*] Renaud Boussarie,[2,†] and Yoshitaka Hatta[3,4,‡]

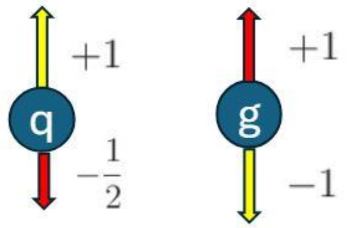
2404.04208

Recall Spin-Orbit coupling in H atom!



$$\boxed{\boldsymbol{G^{q/g}_{1,1}} \to L^{q/g}\cdot S^{q/g}}$$

Perfect spin-orbit anti-correlation



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Main findings

Challenges:

- Low count rate (Amplitude $\sim \alpha_{em}^2$)

Probing quark OAM through double Drell-Yan

Main findings

Challenges:

- Low count rate (Amplitude $\sim \alpha_{em}^2$)
- Sensitivity to GTMDs only in the ERBL region $-\xi < x < \xi$

OAM density: $$L^{q/g}(x, \boldsymbol{\xi}) = -\int d^2\vec{k}_\perp \frac{\vec{k}_\perp^2}{M^2} F_{1,4}^{q,g}(x, k_\perp, \boldsymbol{\xi}, \Delta_\perp = 0)$$

OAM: $$L^{q/g} = \int dx\, L^{q/g}(x, \boldsymbol{\xi} = \mathbf{0})$$

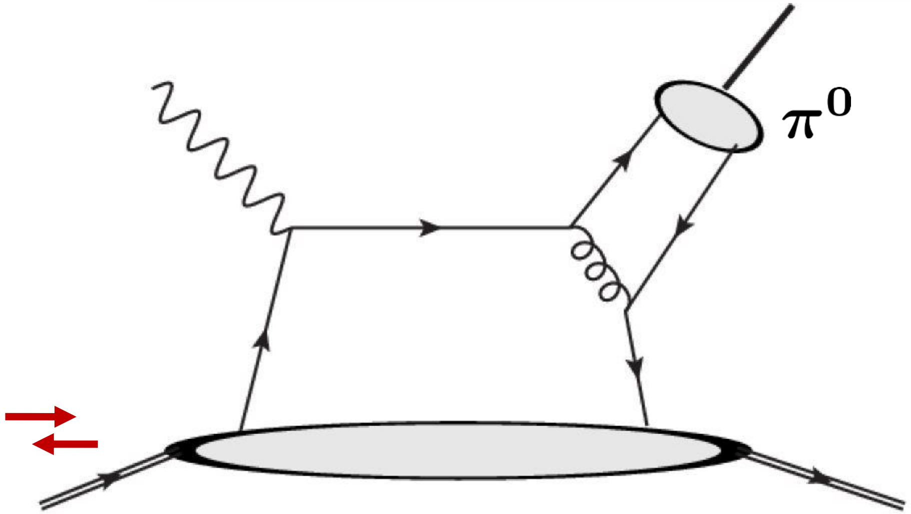
The challenge lies in extrapolating the distribution to the forward limit, where the OAM equation is applicable

Our work

arXiv: 2312.01309 (2023)

Probing quark orbital angular momentum at EIC and EicC

Shohini Bhattacharya,[1] Duxin Zheng,[2] and Jian Zhou[3]

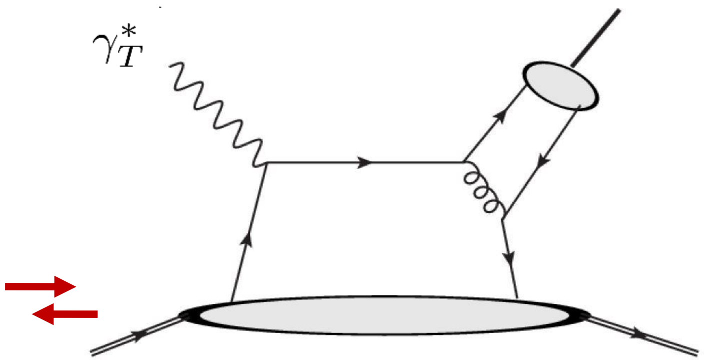


Main Observable:

Longitudinal single-target spin asymmetry

Probing quark OAM through π^0 production in ep collisions

Scattering amplitude



4 leading-order Feynman diagrams

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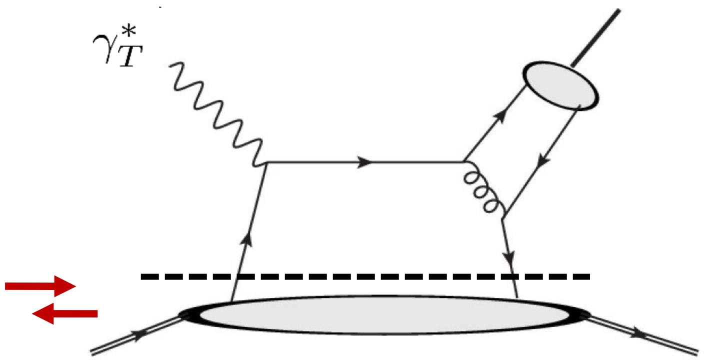

Scattering amplitude:

$$\boxed{A \propto \int dx \int d^2k_\perp \, H(x,\xi,z,k_\perp,\Delta_\perp)\, f^q(x,\xi,k_\perp,\Delta_\perp) \int dz \phi_\pi(z)}$$

Hard part

Soft part from proton

Pion Distribution Amplitude

Probing quark OAM through π^0 production in ep collisions

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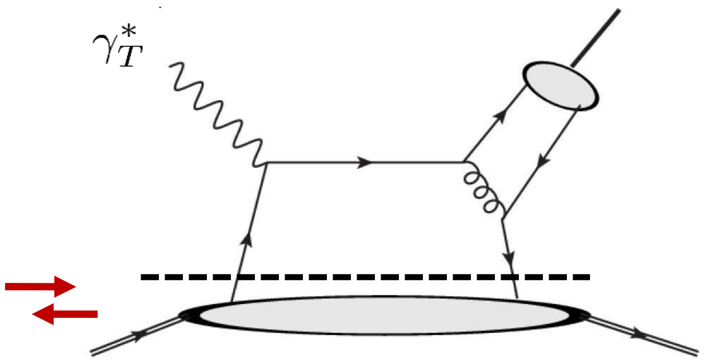

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Collinear twist-expansion of hard part:

$$H(k_\perp,\Delta_\perp) = H(k_\perp=0,\Delta_\perp=0) + \left.\frac{\partial H(k_\perp,\Delta_\perp=0)}{\partial k_\perp^\mu}\right|_{k_\perp=0} k_\perp^\mu + \left.\frac{\partial H(k_\perp=0,\Delta_\perp)}{\partial \Delta_\perp^\mu}\right|_{\Delta_\perp=0} \Delta_\perp^\mu + \ldots$$

Probing quark OAM through π^0 production in ep collisions

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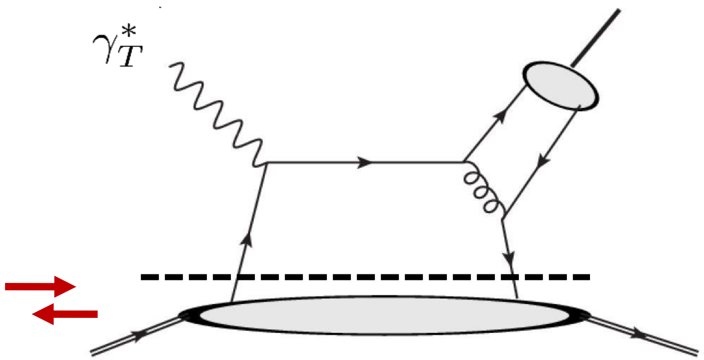

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Twist 2 term vanishes

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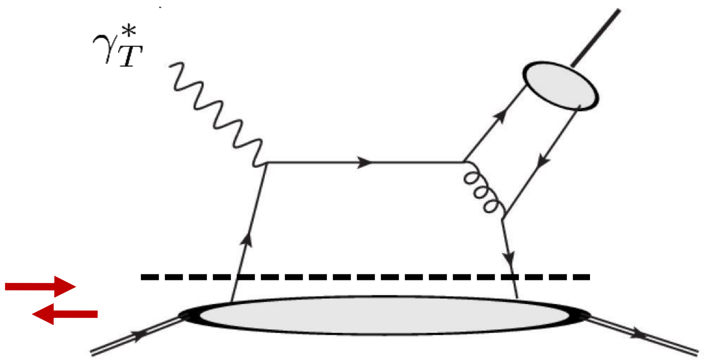

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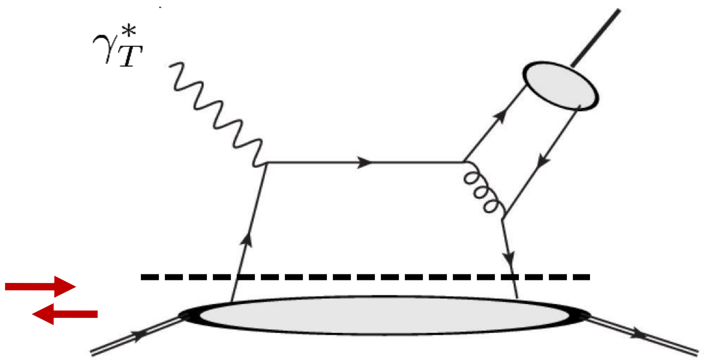
Use special-propagator technique to ensure electromagnetic gauge invariance

(J. W. Qiu, 1990)

Probing quark OAM through π^0 production in ep collisions

Scattering amplitude



Scattering amplitude:

$$\boxed{A \propto \int dx \int \boldsymbol{d^2 k_\perp}\, H(x, \xi, z, k_\perp, \Delta_\perp)\, \boldsymbol{f^q}(x, \xi, k_\perp, \Delta_\perp) \int dz \phi_\pi(z)}$$

Collinear twist-expansion of hard part:

$$H(k_\perp, \Delta_\perp) = H(k_\perp = 0, \Delta_\perp = 0) + \frac{\partial H(k_\perp, \Delta_\perp = 0)}{\partial k_\perp^\mu}\Big|_{k_\perp = 0} \boldsymbol{k_\perp^\mu} + \frac{\partial H(k_\perp = 0, \Delta_\perp)}{\partial \Delta_\perp^\mu}\Big|_{\Delta_\perp = 0} \Delta_\perp^\mu + \ldots$$

$$\boxed{A \propto \int \boldsymbol{d^2 k_\perp k_\perp^2}\ \mathbf{GTMD}}$$

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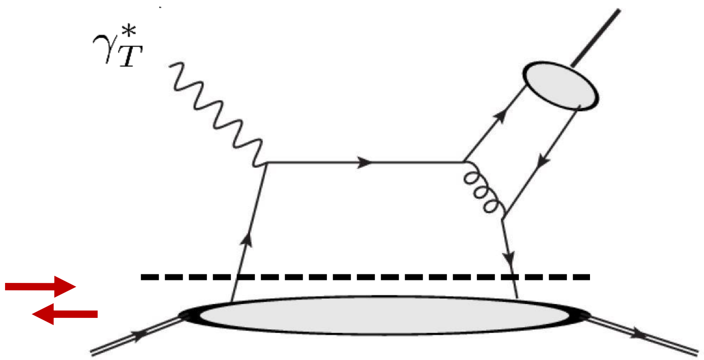

Scattering amplitude:

$$\boxed{A \propto \int dx \int \boldsymbol{d^2 k_\perp}\, H(x,\xi,z,k_\perp,\Delta_\perp)\, \boldsymbol{f^q}(x,\xi,k_\perp,\Delta_\perp) \int dz \phi_\pi(z)}$$

Collinear twist-expansion of hard part:

$$H(k_\perp,\Delta_\perp) = H(k_\perp=0,\Delta_\perp=0) + \frac{\partial H(k_\perp,\Delta_\perp=0)}{\partial k_\perp^\mu}\bigg|_{k_\perp=0} k_\perp^\mu + \frac{\partial H(k_\perp=0,\Delta_\perp)}{\partial \Delta_\perp^\mu}\bigg|_{\Delta_\perp=0} \boldsymbol{\Delta_\perp^\mu} + \ldots$$

$$\boxed{A \propto \mathbf{GPD}}$$

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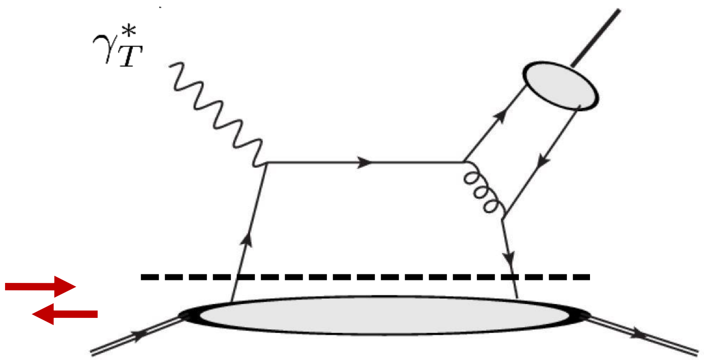
Scattering amplitude



Scattering amplitude:

$$A \propto \int dx \int d^2k_\perp \, H(x,\xi,z,k_\perp,\Delta_\perp) \, f^q(x,\xi,k_\perp,\Delta_\perp) \int dz \phi_\pi(z)$$

Collinear twist-expansion of hard part:

Consequently, the scattering amplitudes are a convolution of moments of GTMDs and GPDs and are of twist-3 nature

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Angular correlations

Scattering amplitudes depend on different angular correlations:

$$\mathcal{M}_1 = \frac{g_s^2 e f_\pi}{2\sqrt{2}} \frac{(N_c^2-1)2\xi}{N_c^2\sqrt{1-\xi^2}} \delta_{\lambda\lambda'} \frac{\boldsymbol{\epsilon}_\perp \times \boldsymbol{\Delta}_\perp}{Q^2} \{\mathcal{F}_{1,1} + \mathcal{G}_{1,1}\}$$

$$\mathcal{M}_2 = \frac{g_s^2 e f_\pi}{2\sqrt{2}} \frac{(N_c^2-1)2\xi}{N_c^2\sqrt{1-\xi^2}} \delta_{\lambda,-\lambda'} \frac{M\boldsymbol{\epsilon}_\perp \cdot \boldsymbol{S}_\perp}{Q^2} \{\mathcal{F}_{1,2} + \mathcal{G}_{1,2}\} \qquad S_\perp^\mu = (0^+, 0^-, -i, \lambda)$$

$$\mathcal{M}_4 = \frac{i g_s^2 e f_\pi}{2\sqrt{2}} \frac{(N_c^2-1)2\xi}{N_c^2\sqrt{1-\xi^2}} \lambda\delta_{\lambda\lambda'} \frac{\boldsymbol{\epsilon}_\perp \cdot \boldsymbol{\Delta}_\perp}{Q^2} \{\mathcal{F}_{1,4} + \mathcal{G}_{1,4}\}$$

Probing quark OAM through π^0 production

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$$\mathcal{M}_4 = \frac{i g_s^2 e f_\pi}{2\sqrt{2}} \frac{(N_c^2-1)2\xi}{N_c^2\sqrt{1-\xi^2}} \lambda\delta_{\lambda\lambda'} \frac{\epsilon_\perp \cdot \Delta_\perp}{Q^2} \boxed{\{\mathcal{F}_{1,4} + \mathcal{G}_{1,4}\}}$$

Compton Form Factors:

$$\mathcal{F}_{1,1} = \int_{-1}^{1} dx \frac{x^2 \int d^2k_\perp F_{1,1}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \tag{8}$$

$$\mathcal{G}_{1,1} = \int_{-1}^{1} dx \int_0^1 dz \frac{\phi_\pi(z)(x^2+2x^2z+\xi^2)}{z^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int d^2k_\perp \frac{k_\perp^2}{M^2} G_{1,1}^{u+d}(x,\xi,\Delta_\perp,k_\perp), \tag{9}$$

$$\mathcal{F}_{1,2} = \int_{-1}^{1} dx x \frac{\xi(1-\xi^2)\int d^2k_\perp k_\perp^2 F_{1,2}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \tag{10}$$

$$\mathcal{G}_{1,2} = \int_{-1}^{1} dx \int_0^1 dz \frac{\phi_\pi(z)(x^2+2x^2z+\xi^2)(1-\xi^2)}{z^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int d^2k_\perp \frac{k_\perp^2}{M^2} G_{1,2}^{u+d}(x,\xi,\Delta_\perp,k_\perp), \tag{11}$$

$$\mathcal{F}_{1,4} = \int_{-1}^{1} dx \frac{x\xi \int d^2k_\perp k_\perp^2 F_{1,4}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \tag{12}$$

$$\mathcal{G}_{1,4} = \int_{-1}^{1} dx \int_0^1 dz \frac{x(4\xi^2 z+\xi^2-2x^2z+x^2)}{z^2\xi(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2}\phi_\pi(z) \times \int d^2k_\perp G_{1,4}^{u+d}(x,\xi,\Delta_\perp,k_\perp). \tag{13}$$

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$$\mathcal{M}_4 = \frac{i g_s^2 e f_\pi}{2\sqrt{2}} \frac{(N_c^2-1)2\xi}{N_c^2\sqrt{1-\xi^2}} \lambda\delta_{\lambda\lambda'} \frac{\epsilon_\perp \cdot \Delta_\perp}{Q^2} \boxed{\{\mathcal{F}_{1,4} + \mathcal{G}_{1,4}\}}$$

Sensitivity to quark OAM

$$\mathcal{F}_{1,1} = \int_{-1}^{1} dx \frac{x^2 \int d^2k_\perp F_{1,1}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \quad (8)$$

$$\mathcal{G}_{1,1} = \int_{-1}^{1} dx \int_0^1 dz \frac{\phi_\pi(z)(x^2+2x^2z+\xi^2)}{z^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int d^2k_\perp \frac{k_\perp^2}{M^2} G_{1,1}^{u+d}(x,\xi,\Delta_\perp,k_\perp), \quad (9)$$

$$\mathcal{F}_{1,2} = \int_{-1}^{1} dx\, x \frac{\xi(1-\xi^2)\int d^2k_\perp k_\perp^2 F_{1,2}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \quad (10)$$

$$\mathcal{G}_{1,2} = \int_{-1}^{1} dx \int_0^1 dz \frac{\phi_\pi(z)(x^2+2x^2z+\xi^2)(1-\xi^2)}{z^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int d^2k_\perp \frac{k_\perp^2}{M^2} G_{1,2}^{u+d}(x,\xi,\Delta_\perp,k_\perp), \quad (11)$$

$$\mathcal{F}_{1,4} = \int_{-1}^{1} dx \frac{x\xi \int d^2k_\perp k_\perp^2 F_{1,4}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}, \quad (12)$$

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Cross section

$$\frac{d\sigma}{dtdQ^2dx_Bd\phi} = \frac{(N_c^2-1)^2\alpha_{em}^2\alpha_s^2 f_\pi^2 \xi^3 \Delta_\perp^2}{2N_c^4(1-\xi^2)Q^{10}(1+\xi)} \left[1+(1-y)^2\right]$$

$$\times \left\{ \left[|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2 + 2\frac{M^2}{\Delta_\perp^2}|\mathcal{F}_{1,2}+\mathcal{G}_{1,2}|^2 \right] + \cos(2\phi)a\left[-|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2\right] \right.$$

$$\left. +\lambda \sin(2\phi)\, 2a\, \mathrm{Re}\left[(i\mathcal{F}_{1,4}+i\mathcal{G}_{1,4})\left(\mathcal{F}_{1,1}^* + \mathcal{G}_{1,1}^*\right)\right] \right\}$$

$$a = \frac{2(1-y)}{1+(1-y)^2}$$

Probing quark OAM through π^0 production in ep collisions

Cross section

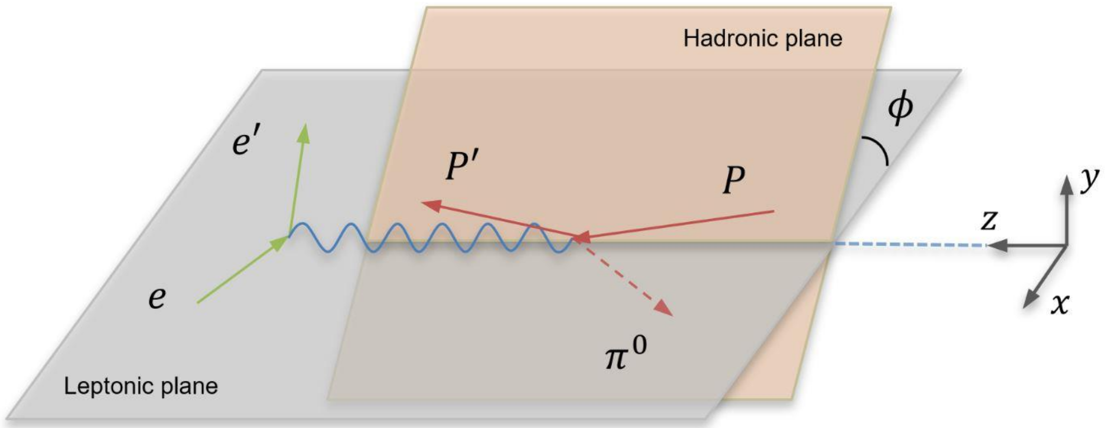
$$\frac{d\sigma}{dtdQ^2dx_Bd\phi} = \frac{(N_c^2-1)^2\alpha_{em}^2\alpha_s^2 f_\pi^2\xi^3\Delta_\perp^2}{2N_c^4(1-\xi^2)Q^{10}(1+\xi)}\left[1+(1-y)^2\right]$$

$$\times\left\{\left[|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2+|\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2+2\frac{M^2}{\Delta_\perp^2}|\mathcal{F}_{1,2}+\mathcal{G}_{1,2}|^2\right]+\cos(2\phi)a\left[-|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2+|\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2\right]\right.$$

$$\boxed{\left.+\lambda\,\mathbf{sin(2\phi)}\,2a\,\mathrm{Re}\left[\left(i\boldsymbol{\mathcal{F}_{1,4}}+i\mathcal{G}_{1,4}\right)\left(\mathcal{F}_{1,1}^*+\mathcal{G}_{1,1}^*\right)\right]\right\}}$$

Distinguished experimental signature of quark OAM

$$\phi = \phi_{l_\perp} - \phi_{\Delta_\perp}$$



Probing quark OAM through π^0 production in ep collisions

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$$\boxed{\times\left\{\left[|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2+|\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2+2\frac{M^2}{\Delta_\perp^2}|\boldsymbol{\mathcal{F}_{1,2}}+\mathcal{G}_{1,2}|^2\right]+\cos(2\phi)a\left[-|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2+|\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2\right]\right.}$$

$$\left.+\lambda\sin(2\phi)\,2a\,\mathrm{Re}\left[(i\mathcal{F}_{1,4}+i\mathcal{G}_{1,4})\left(\mathcal{F}_{1,1}^*+\mathcal{G}_{1,1}^*\right)\right]\right\}$$

Surprise!

- Probe quark Sivers function through an unpolarized target

$$\boxed{\mathrm{Im}\left[\boldsymbol{F_{1,2}}\right]\Big|_{\Delta=0} = -\boldsymbol{f_{1T}^\perp}}$$

(Similar to the gluon GTMD $F_{1,2}$, as discussed in Boussarie, Hatta, Szymanowski, Wallon, 2019)

Probing quark OAM through π^0 production in ep collisions

Cross section

$$\frac{d\sigma}{dt dQ^2 dx_B d\phi} = \frac{(N_c^2-1)^2 \alpha_{em}^2 \alpha_s^2 f_\pi^2 \xi^3 \Delta_\perp^2}{2N_c^4(1-\xi^2)Q^{10}(1+\xi)} \left[1+(1-y)^2\right]$$

$$\boxed{\times \left\{ \left[|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2 + 2\frac{M^2}{\Delta_\perp^2}|\boldsymbol{\mathcal{F}_{1,2}}+\boldsymbol{\mathcal{G}_{1,2}}|^2 \right] + \cos(2\phi) a \left[-|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2 \right]}$$

$$+\lambda \sin(2\phi)\, 2a\, \mathrm{Re}\left[(i\mathcal{F}_{1,4}+i\mathcal{G}_{1,4})\left(\mathcal{F}_{1,1}^*+\mathcal{G}_{1,1}^*\right) \right] \Big\}$$

Surprise!

- Probe quark Sivers function through an unpolarized target $\quad \boxed{\mathrm{Im}\left[\boldsymbol{F_{1,2}}\right]\big|_{\Delta=0} = -\boldsymbol{f_{1T}^{\perp}}}$
- Probe quark worm gear function through an unpolarized target $\quad \boxed{\mathrm{Re}\left[\boldsymbol{G_{1,2}}\right]\big|_{\Delta=0} = \boldsymbol{g_{1T}}}$

Probing quark OAM through π^0 production in ep collisions

Cross section

$$\frac{d\sigma}{dt\,dQ^2\,dx_B\,d\phi} = \frac{(N_c^2-1)^2\alpha_{em}^2\alpha_s^2 f_\pi^2 \xi^3 \boldsymbol{\Delta_\perp^2}}{2N_c^4(1-\xi^2)Q^{10}(1+\xi)}\left[1+(1-y)^2\right]$$

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$$\left.+\lambda\sin(2\phi)\,2a\,\mathrm{Re}\left[\left(i\mathcal{F}_{1,4}+i\mathcal{G}_{1,4}\right)\left(\mathcal{F}_{1,1}^*+\mathcal{G}_{1,1}^*\right)\right]\right\}$$

Helicity flip terms persist even when $\Delta_\perp \to 0$

Probing quark OAM through π^0 production in ep collisions

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$$\times \left\{ \left[|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2 + 2\frac{M^2}{\Delta_\perp^2} |\mathcal{F}_{1,2}+\mathcal{G}_{1,2}|^2 \right] + \cos(2\phi) a \left[-|\mathcal{F}_{1,1}+\mathcal{G}_{1,1}|^2 + |\mathcal{F}_{1,4}+\mathcal{G}_{1,4}|^2 \right] \right.$$

$$\left. + \lambda \sin(2\phi)\, 2a\, \mathrm{Re}\left[(i\mathcal{F}_{1,4} + i\mathcal{G}_{1,4}) \left(\mathcal{F}_{1,1}^* + \mathcal{G}_{1,1}^* \right) \right] \right\}$$

Since both unpolarized and polarized cross sections contribute at twist-3, the magnitudes of the asymmetries are not power-suppressed

Probing quark OAM through π^0 production in ep collisions

Model input for numerical estimations

Probing quark OAM through π^0 production in ep collisions

Model input for numerical estimations

Ingredients for non-perturbative functions:

- Model $(H^q,\, \tilde{H}^q)$ according to the Double distribution approach (see Radyushkin, 9805342)

Example:

$$\boxed{H^q(x,\xi,t) = \int_{-1}^{1} d\beta \int_{-1+|\beta|}^{1-|\beta|} d\alpha\, \delta(\beta+\xi\alpha-x) \times \frac{3}{4}|\beta|^{-1.3\,t}\, \frac{\left[(1-|\beta|)^2-\alpha^2\right]}{(1-|\beta|)^3}\, q(|\beta|)}$$

JAM 2022 PDFs

Probing quark OAM through π^0 production in ep collisions

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The t-dependence is determined based on a fit to CLAS data

Probing quark OAM through π^0 production in ep collisions

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- Model $(H^q,\ \tilde{H}^q)$ according to the Double distribution approach (see Radyushkin, 9805342)

- Model for OAM:

  1. “OAM density”: (Hatta, Yoshida, 1207.5332)

$$L^q_{can}(\boldsymbol{x}) = x\int_x^1 \frac{dx'}{x'}q(x') - x\int_x^1 \frac{dx'}{x'^2}\Delta q(x') + \text{ genuine twist-three}$$

Probing quark OAM through π^0 production in ep collisions

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- Model $(H^q, \tilde{H}^q)$ according to the Double distribution approach (see Radyushkin, 9805342)

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Probing quark OAM through π^0 production in ep collisions

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2. Use the Double distribution approach to construct $xL^q(x,\boldsymbol{\xi})e^{\boldsymbol{t}/\Lambda}$ from $xL^q(x)$

Model input for numerical estimations

Ingredients for non-perturbative functions:

- Pion distribution amplitude:

Asymptotic form $\boxed{\phi_\pi(z) = 6z(1-z)}$

Probing quark OAM through π^0 production in ep collisions

Model input for numerical estimations

End-point singularity & discontinuity:

$$\mathcal{F}_{1,4} = \int_{-1}^{1} dx \frac{x\xi \int d^2k_\perp k_\perp^2 F_{1,4}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}$$

Model-dependent method:

$$\boxed{\int_{\langle p_\perp^2\rangle/Q^2}^{1-\langle p_\perp^2\rangle/Q^2} dz}$$

$\langle p_\perp^2\rangle = 0.04\,\mathrm{GeV}^2$ determined based on a fit to CLAS data

S. V. Goloskokov and P. Kroll, 2005

Probing quark OAM through π^0 production in ep collisions

Model input for numerical estimations

End-point singularity & discontinuity:

$$\mathcal{F}_{1,4} = \int_{-1}^{1} dx \frac{x\xi \int d^2k_\perp k_\perp^2 F_{1,4}^{u+d}(x,\xi,\Delta_\perp,k_\perp)}{M^2 (x+\xi-i\epsilon)^2 (x-\xi+i\epsilon)^2} \times \int_0^1 dz \frac{\phi_\pi(z)(1+z^2-z)}{z^2(1-z)^2}$$

Model-dependent method:

$$\int_{\langle p_\perp^2\rangle/Q^2}^{1-\langle p_\perp^2\rangle/Q^2} dz$$

S. V. Goloskokov and P. Kroll, 2005

$$\boxed{\frac{1}{(x-\xi+i\epsilon)^2} \to \frac{1}{(x-\xi-\langle p_\perp^2\rangle/Q^2+i\epsilon)^2}}$$

I. V. Anikin, O. V. Teryaev, 2003

Probing quark OAM through π^0 production in ep collisions

Numerical results

Kinematics:

	$Q^2(\mathrm{GeV}^2)$	$\sqrt{s}_{ep}(\mathrm{GeV})$
EIC	10	100
EicC	3	16

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	$Q^2(\mathrm{GeV}^2)$	$\sqrt{s_{ep}}(\mathrm{GeV})$
EIC	10	100
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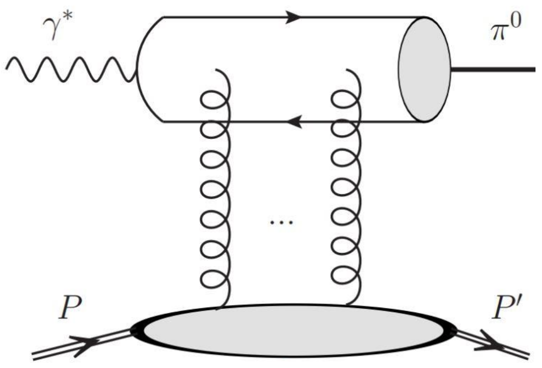
- We focus on large skewness (ξ) region to suppress gluon contribution



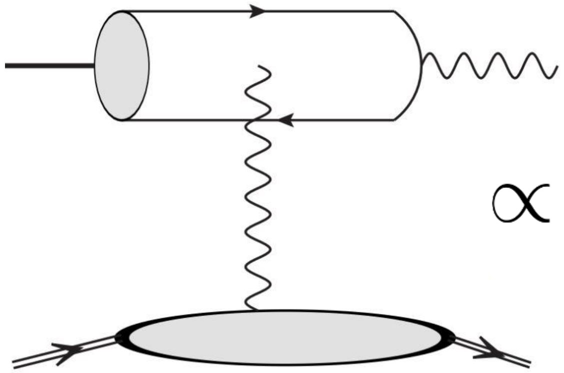
Probing quark OAM through π^0 production in ep collisions

Numerical results

Kinematics:

	$Q^2(\mathrm{GeV}^2)$	$\sqrt{s_{ep}}(\mathrm{GeV})$
EIC	10	100
EicC	3	16

- We focus on large skewness (ξ) region to suppress gluon contribution
- We focus on large momentum transfer (t) region to suppress contribution from Primakoff process



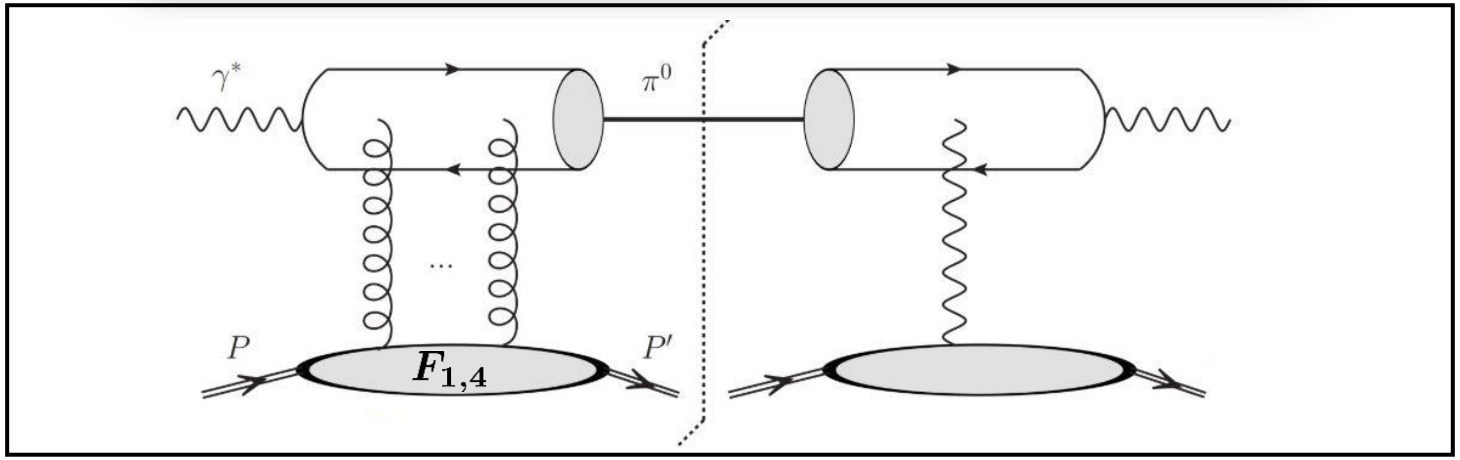
$$\propto \frac{1}{t}$$

Probing quark OAM through π^0 production in ep collisions

Remark:

Accessing the gluon GTMD $F_{1,4}$ in exclusive π^0 production in ep collisions

Shohini Bhattacharya,[1] Duxin Zheng,[2] and Jian Zhou[3]



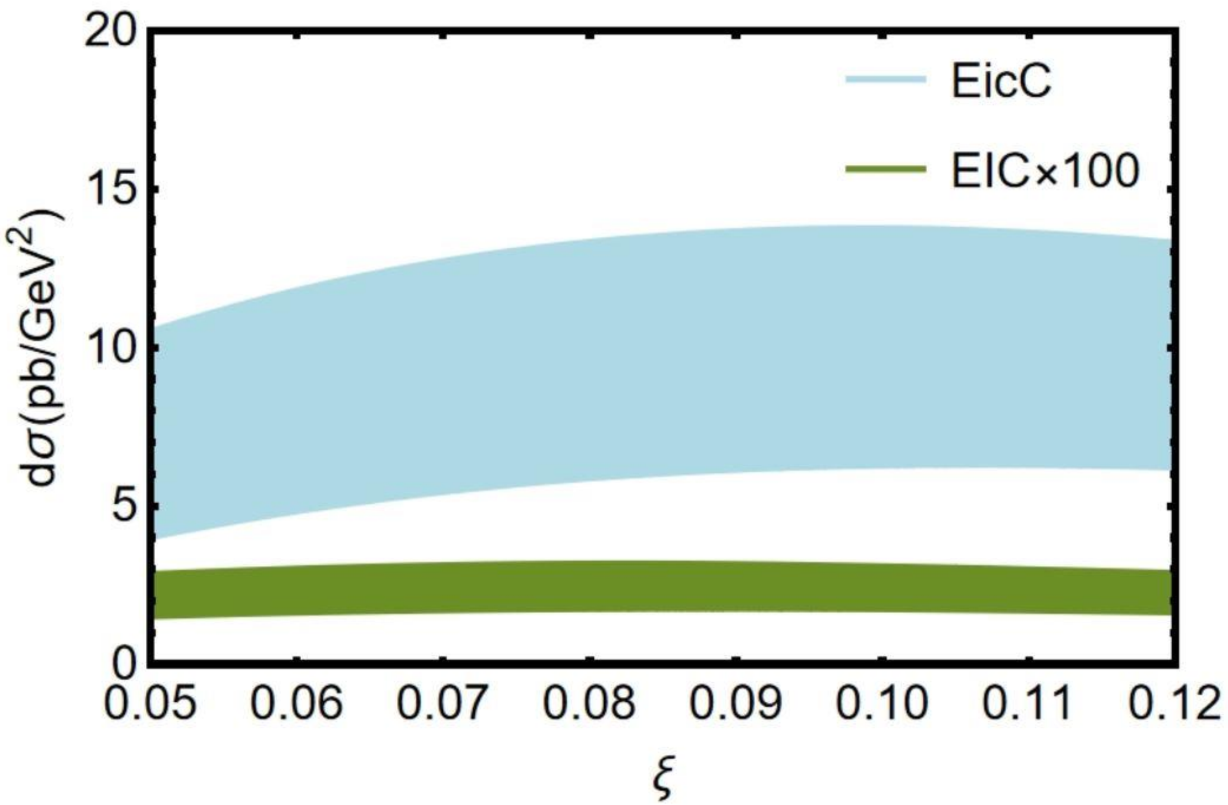
$$\frac{d\Delta\sigma}{dt\,dQ^2\,dx_B\,d\phi} = -\sin(2\phi)\frac{\alpha_{em}^3\alpha_s f_\pi^2(1-y)\xi x_B\mathcal{F}(t)}{3Q^8N_c}\left[\int_0^1 dz\,\frac{\phi_\pi(z)}{z(1-z)}\right]^2 \mathrm{Im}\left[\int_{-1}^1 dx\,\frac{F_{1,4}^{(1)}(x,\xi,\Delta_\perp)/M^2}{(x+\xi-i\epsilon)^2(x-\xi+i\epsilon)^2}\right]$$

The same azimuthal asymmetry, precisely mirroring what we observe in this study, emerges from the interference between the Primakoff process and the contribution from the gluon GTMD

Probing quark OAM through π^0 production in ep collisions

Numerical results

Unpolarized cross section



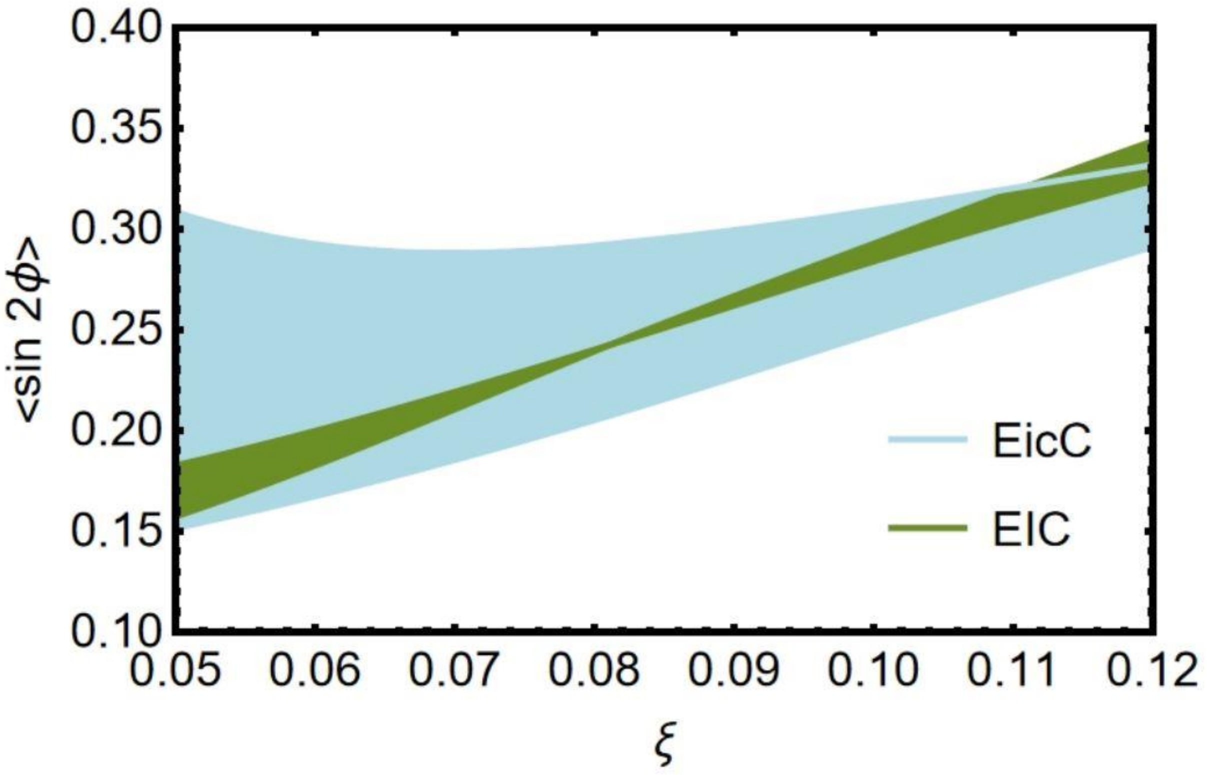
Findings:

- The unpolarized cross section exhibits a notable magnitude at EicC energy
- Relatively small at EIC energy

Probing quark OAM through π^0 production in ep collisions

Numerical results

Asymmetry



$$\langle \sin(2\phi) \rangle = \frac{\int \frac{d\Delta\sigma}{d\mathcal{P}.\mathcal{S}.} \sin(2\phi)\, d\mathcal{P}.\mathcal{S}.}{\int \frac{d\sigma}{d\mathcal{P}.\mathcal{S}.} d\mathcal{P}.\mathcal{S}.}$$

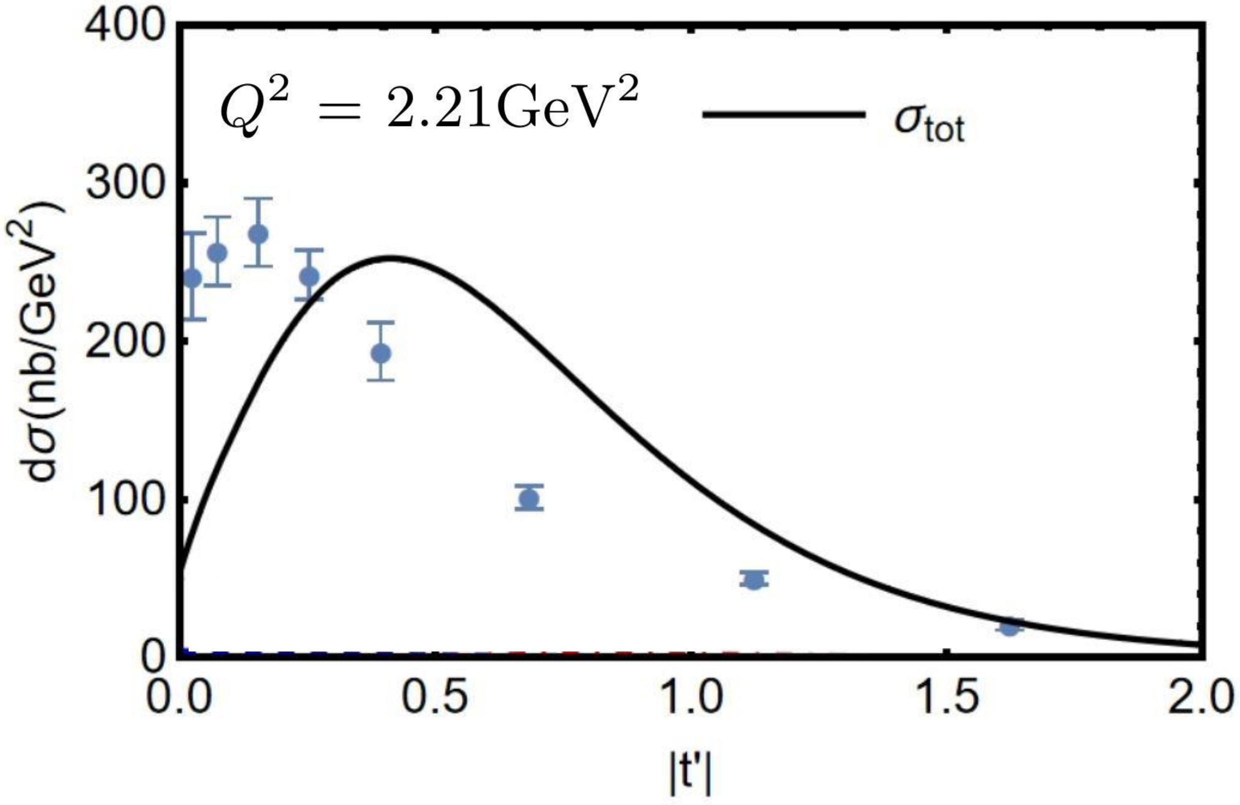
Findings:

The asymmetries are substantial for both EIC & EicC kinematics

Probing quark OAM through π^0 production in ep collisions

Numerical results

Comparison with CLAS data



Unpolarized cross section:

$$\boxed{\frac{d\sigma_T}{dt} + a\frac{d\sigma_L}{dt}}$$

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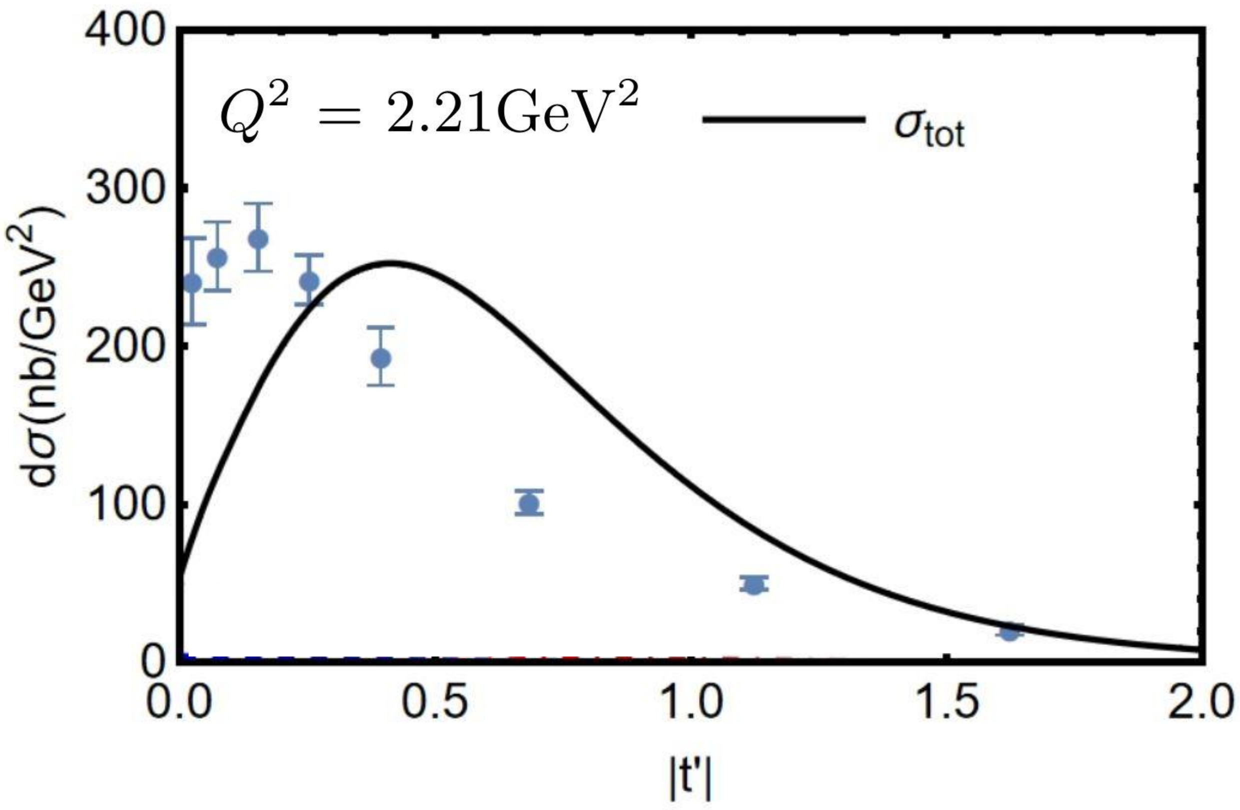


Findings:

- Our theoretical model is in reasonable agreement with experimental data

Developments

arXiv: 1612.02438 (2016)

Hunting the Gluon Orbital Angular Momentum at the Electron-Ion Collider

Xiangdong Ji,[1,2] Feng Yuan,[3] and Yong Zhao[1,3]

arXiv: 1702.04387 (2017)

Generalized TMDs and the exclusive double Drell-Yan process

Shohini Bhattacharya,[1] Andreas Metz,[1] and Jian Zhou[2]

arXiv: 1802.10550 (2018)

Exclusive double quarkonium production and generalized TM

Shohini Bhattacharya,[1] Andreas Metz,[1] Vikash Kumar Ojha,[2] Jeng-Yuan Tsai,[1]

arXiv: 1807.08697 (2018)

Probing the Weizsäcker-Williams gluon Wigner distribution in pp collisions

Renaud Boussarie,[1] Yoshitaka Hatta,[2] Bo-Wen Xiao,[3,4] and Feng Yuan[5]

arXiv: 1912.08182 (2019)

Probing the gluon Sivers function with an unpola GTMD distributions and the Odderon

Renaud Boussarie,[1] Yoshitaka Hatta,[1] Lech Szymanowski,[2] and S

arXiv: 2106.13466 (2021)

Probing the gluon tomography in photoproduction of di-pions

Yoshikazu Hagiwara, Cheng Zhang, Jian Zhou, and Ya-jin Zhou

arXiv: 2201.08709 (2022/2024)

Signature of the gluon orbital angular momentum

Shohini Bhattacharya,[1,*] Renaud Boussarie,[2,†] and Yoshitaka Hatta[1,3,‡]

rXiv: 2205.00045 (2022)

ngular correlations in exclusive dijet photoproduction in ultra-peripheral PbPb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

Selected works on gluon GTMDs

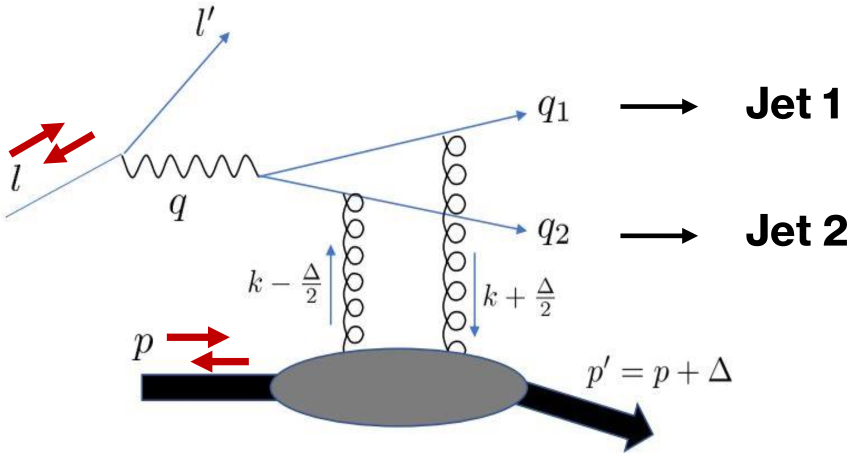
Probing gluon OAM through exclusive di-jet production

PHYSICAL REVIEW LETTERS **128,** 182002 (2022)

DOE Highlight

Signature of the Gluon Orbital Angular Momentum

Shohini Bhattacharya,[1,*] Renaud Boussarie,[2,†] and Yoshitaka Hatta[1,3,‡]



Selected works on gluon GTMDs

Probing gluon OAM through exclusive di-jet production

PHYSICAL REVIEW LETTERS 128, 182002 (2022)

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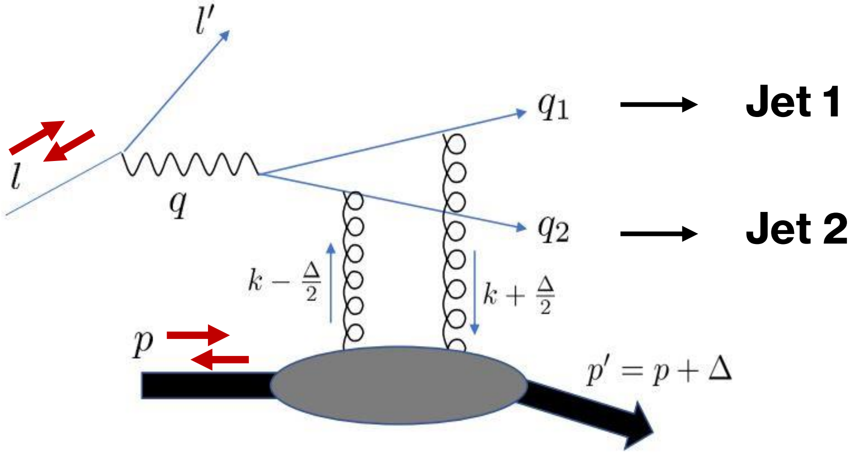

Main result (double spin asymmetry):

Signature of gluon OAM is cosine angular modulation

$$d\sigma^{\text{asym}} \sim -\,\Re\mathrm{e}\left[\left\{\mathcal{H}_g^{(1)*}(\xi) + \frac{4q_\perp^2}{q_\perp^2+\mu^2}\mathcal{H}_g^{(2)*}(\xi)\right\}\mathcal{L}_g(\xi)\right]\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$$

$$+\,\Re\mathrm{e}\left[\mathcal{H}_g^{(1)*}(\xi)\,\tilde{\mathcal{H}}_g^{(2)}(\xi)\right]\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$$

Selected works on gluon GTMDs

Probing gluon OAM through exclusive di-jet production

PHYSICAL REVIEW LETTERS 128, 182002 (2022)

DOE Highlight

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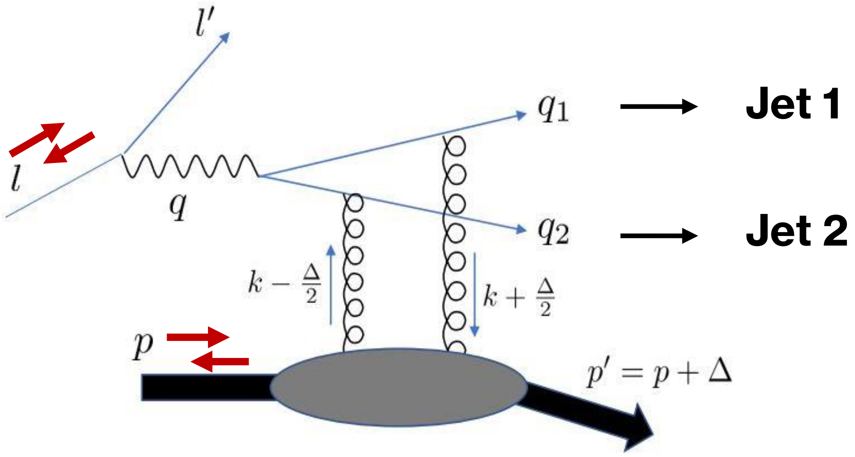
Shohini Bhattacharya[1,*] Renaud Boussarie,[2,†] and Yoshitaka Hatta[1,3,‡]



Gluon helicity contributes to the same angular modulation as that of OAM

$$d\sigma^{\text{asym}} \sim -\Re e\left[\left\{\mathcal{H}_g^{(1)*}(\xi) + \frac{4q_\perp^2}{q_\perp^2+\mu^2}\mathcal{H}_g^{(2)*}(\xi)\right\}\mathcal{L}_g(\xi)\right]\cos(\phi_{l_\perp}-\phi_{\Delta_\perp})$$

$$+\Re e\left[\mathcal{H}_g^{(1)*}(\xi)\tilde{\mathcal{H}}_g^{(2)}(\xi)\right]\cos(\phi_{l_\perp}-\phi_{\Delta_\perp})$$

Helicity GPD
(intrinsic spin)

Selected works on gluon GTMDs

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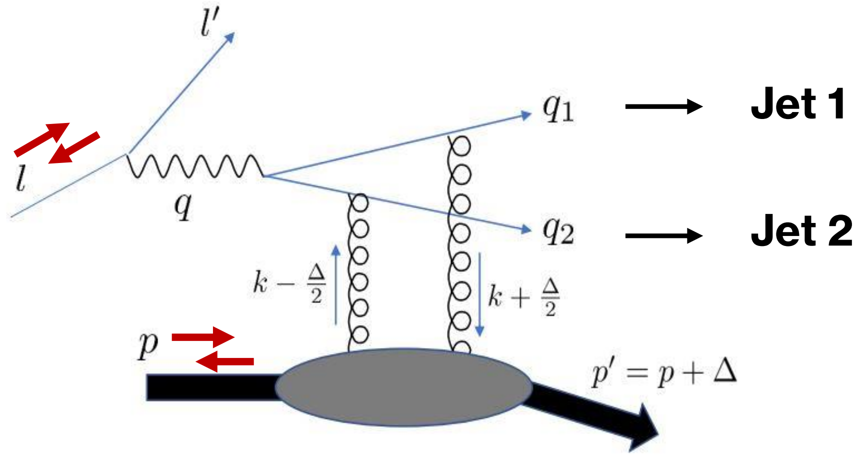


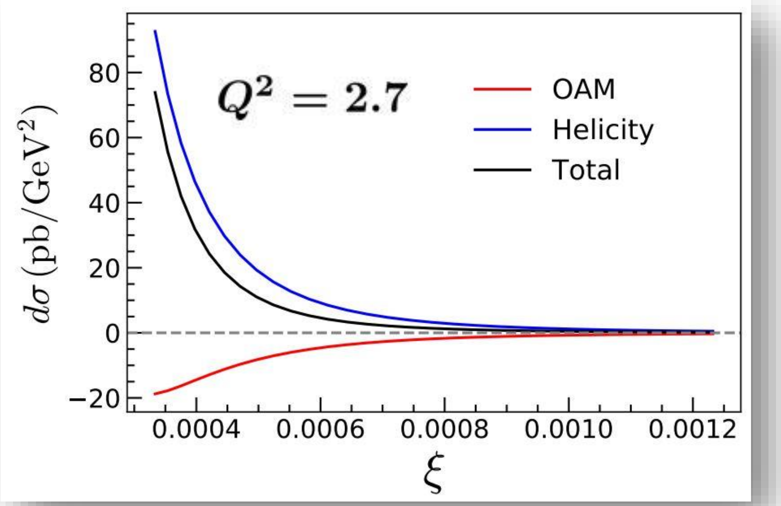
Our observable is a simultaneous probe of gluon OAM & it's helicity

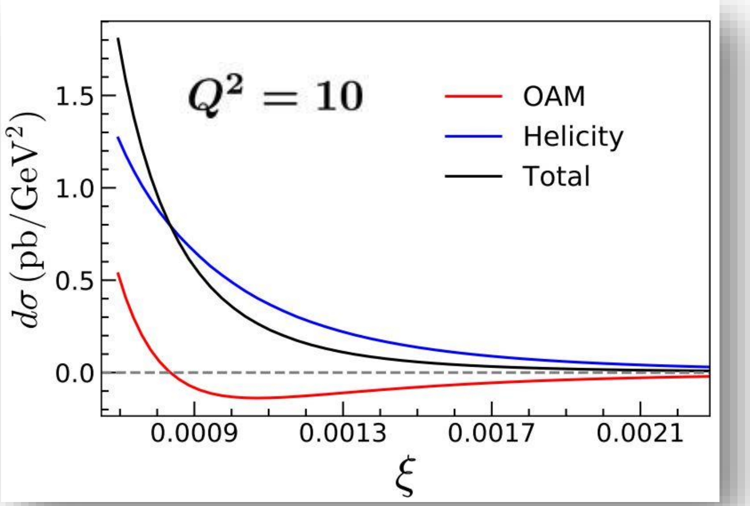
$$d\sigma^{\mathbf{asym}} \sim -\,\mathfrak{Re}\left[\left\{\mathcal{H}_g^{(1)*}(\xi) + \frac{4q_\perp^2}{q_\perp^2+\mu^2}\mathcal{H}_g^{(2)*}(\xi)\right\}\mathcal{L}_g(\xi)\right]\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$$

$$+\,\mathfrak{Re}\left[\mathcal{H}_g^{(1)*}(\xi)\,\tilde{\mathcal{H}}_g^{(2)}(\xi)\right]\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$$

Selected works on gluon GTMDs

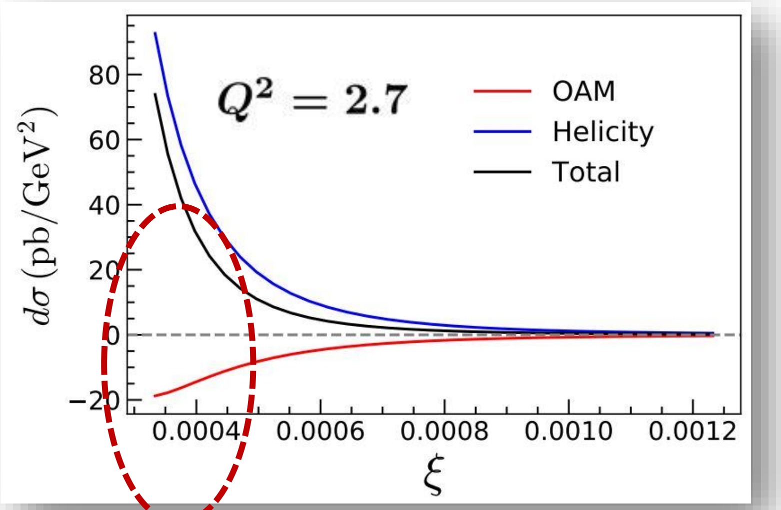
Interplay between OAM and helicity at small x

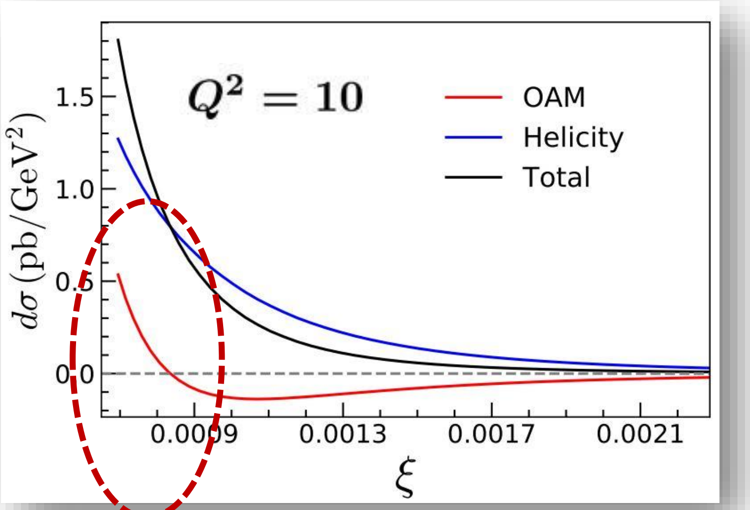




Selected works on gluon GTMDs

Interplay between OAM and helicity at small x





Schematic structure of our observable:

$$d\sigma^{\text{asym}} \sim \mathcal{H}_g^{(1)*}(\xi)\left(\tilde{\mathcal{H}}_g^{(2)}(\xi) + \frac{q_\perp^2 - Q^2/4}{q_\perp^2 + Q^2/4}\,\mathcal{L}_g(\xi)\right)$$

$$\tilde{\mathcal{H}}_g^{(2)}(\xi) \longrightarrow \Delta G(x) \qquad \mathcal{L}_g(\xi) \longrightarrow L_g(x)$$

Yuri Kovchegov's talk

Cancellation expected between helicity & OAM at small x

$$\boxed{\Delta G(x) \approx -\, L_g(x)}$$

Boussarie, Hatta, Yuan (2019)
Kovchegov, Manley (2023, 2024)

Selected works on gluon GTMDs

Contribution from spin-orbit correlation at small x?

Yet another contribution to the process:

$$\boxed{d\sigma^{\mathbf{asym}} \sim \frac{4q_\perp^2}{q_\perp^2 + \mu^2} C_g^{(2)}(\xi) \tilde{\mathcal{H}}_g^{(2)}(\xi)}$$

Spin-orbit correlation:

$$C^g(x) = \int d^2\vec{k}_\perp \frac{\vec{k}_\perp^2}{M^2} \boldsymbol{G^g_{1,1}(x, \vec{k}^2_\perp)}$$

Selected works on gluon GTMDs

Contribution from spin-orbit correlation at small x?

Yet another contribution to the process:

$$d\sigma^{\text{asym}} \sim \frac{4q_\perp^2}{q_\perp^2 + \mu^2} C_g^{(2)}(\xi)\tilde{\mathcal{H}}_g^{(2)}(\xi)$$

First insight into the small-x behavior of spin-orbit correlation:

$$C^g(x) \approx -2x \int_x^1 \frac{dx'}{x'^2} G(x') + \ldots \propto -G(x)$$

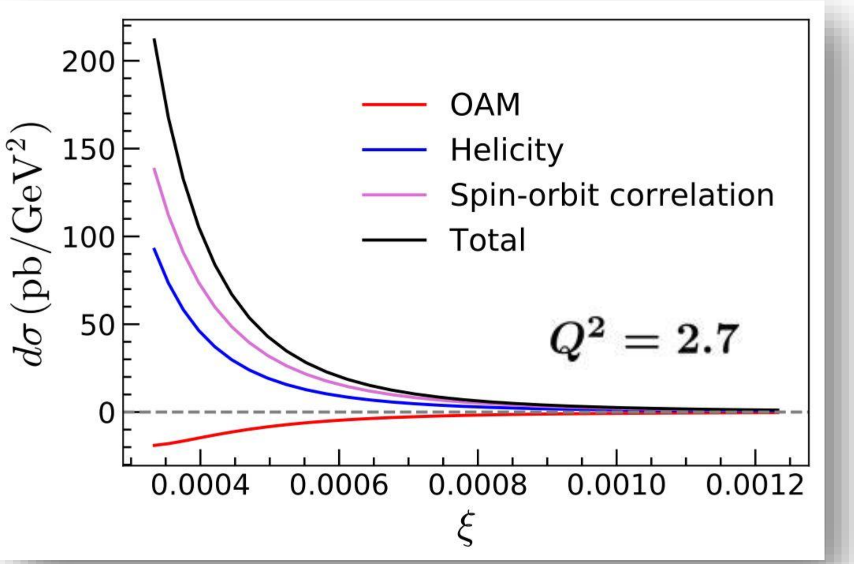
(SB, Boussarie, Hatta, 2404.04208 , 2404.04209)

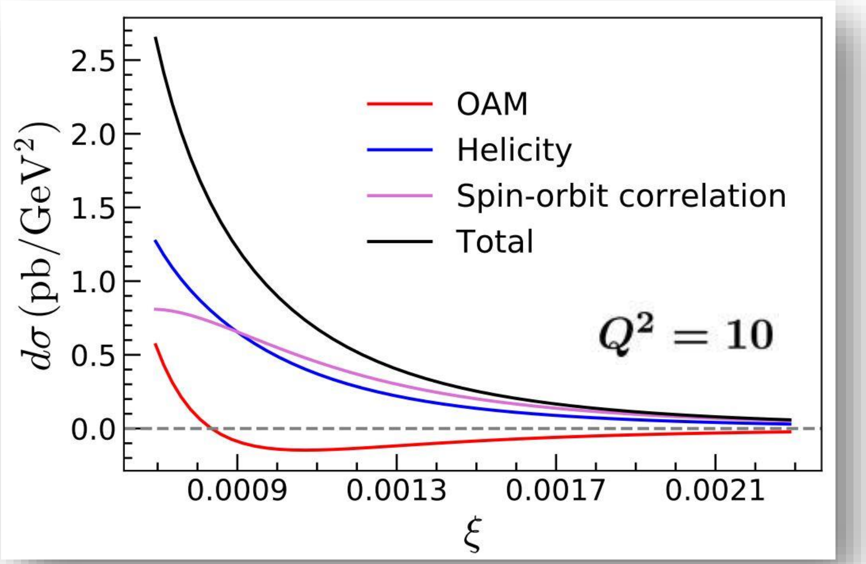
For a complete twist structure of spin-orbit correlation, see Hatta, Schoenleber, 2404.18872

Selected works on gluon GTMDs

Probing gluon OAM & spin-orbit correlation at small x

Updated numerical results (SB, Boussarie, Hatta, 2404.04209)**:**



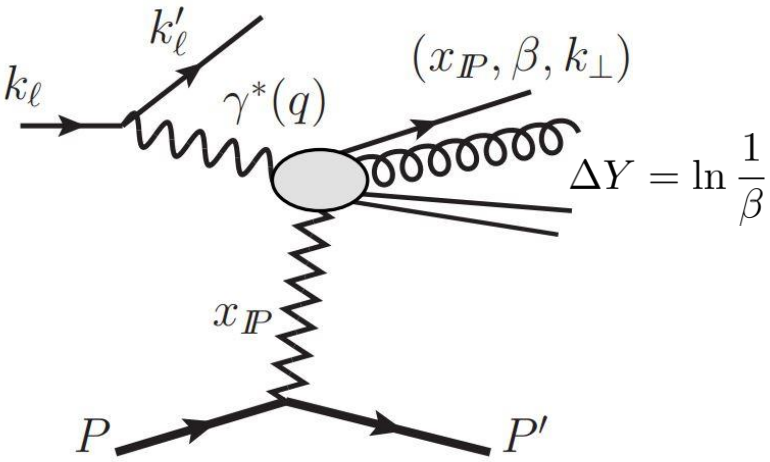


Spin-orbit correlation is more accurately constrained than **OAM** because the latter necessitates the precise determination of both unpolarized and polarized gluon distributions

Selected works on gluon GTMDs

Probing gluon OAM through Semi Inclusive Diffractive Deep Inelastic Scattering

Feng Yuan's talk



Hatta, Xiao, Yuan (2022)

- Measure invariant mass of diffractively produced system instead of reconstructing jets

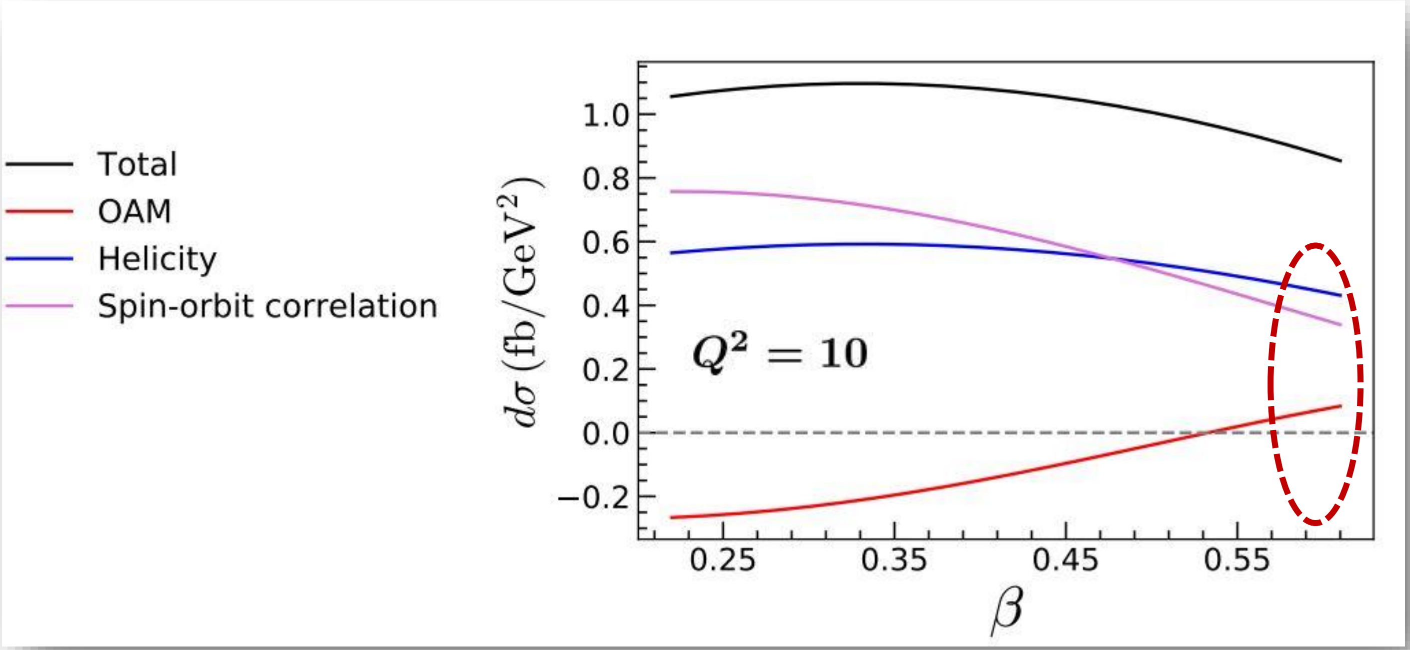
$$M_X^2 = \frac{q_\perp^2}{z\bar{z}} = \frac{1-\beta}{\beta} Q^2$$

- Tag hadron species out of the diffractively produced system

Selected works on gluon GTMDs

Probing gluon OAM through **S**emi **I**nclusive **D**iffractive **D**eep **I**nelastic **S**cattering

Numerical results (SB, Boussarie, Hatta, 2404.04209)**:**



Challenging, yet there is no requirement to reconstruct jets & we still maintain sensitivity to OAM

Summary

- Generalized TMDs/Wigner functions are the holy grail of spin physics

Summary

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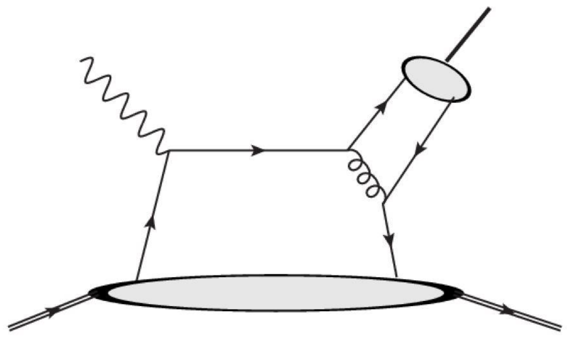
- Probe **quark OAM** via exclusive π^0 production in ep collisions
- Circumvent challenges associated with double Drell-Yan process



Summary

- Generalized TMDs/Wigner functions are the holy grail of spin physics
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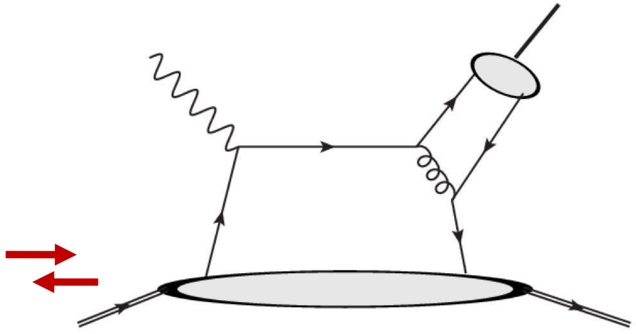
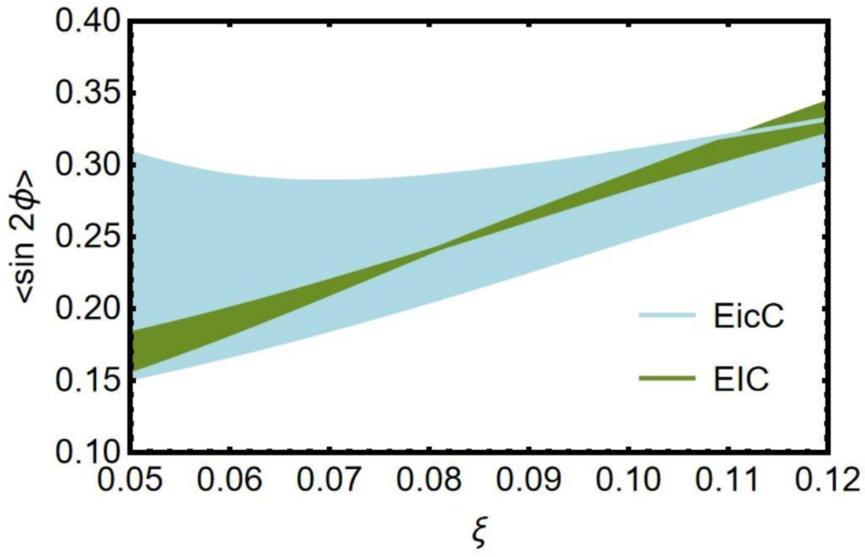
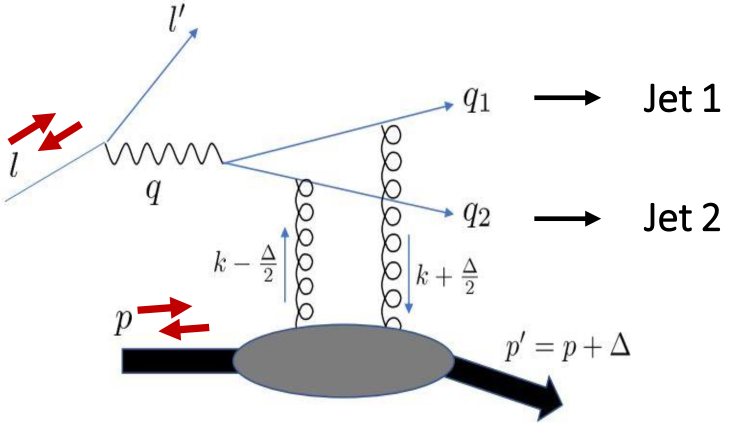


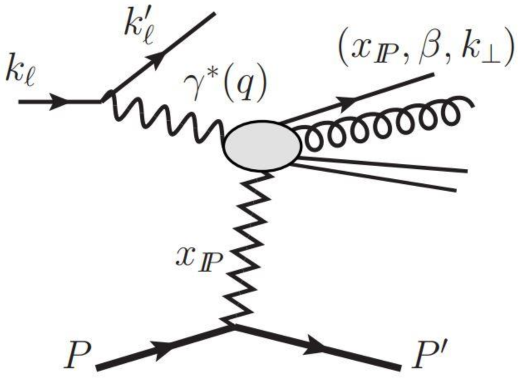


- Longitudinal single-target spin asymmetry is not power suppressed
- Asymmetry is substantial & thus exclusive π^0 production in ep collisions maybe a promising route to constrain quark OAM

Summary

- Probe **gluon OAM** via exclusive di-jet production/ SIDDIS in ep collisions





Summary

- Probe **gluon OAM** via exclusive di-jet production/ SIDDIS in ep collisions