

numerical implementation for twist-3 PDF evolution equations

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Introduction



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Fundamentals

The operators

 $\bar{q}_i(z_1n)gF_a^{\mu+}(z_2n)q_j(z_3n) \qquad F_a^{\mu+}(z_1n)gF_b^{\nu+}(z_2n)F_c^{\rho+}(z_3n)$



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Symmetries and independent distributions

From Parity, Time-Reversal and Charge Conjugation:

2 independent quark-gluon-quark distributions and 2 independent 3-gluon distributions

 $\mathfrak{S}^{\pm}(x_1, x_2, x_3) = \pm \mathfrak{S}^{\pm}(-x_1, -x_2, -x_3)$

 $\mathfrak{F}^{\pm}(x_1, x_2, x_3) = \mp \mathfrak{F}^{\pm}(-x_1, -x_2, -x_3) \quad \mathfrak{F}^{\pm}(x_1, x_2, x_3) = \mp \mathfrak{F}^{\pm}(x_1, x_3, x_2)$



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Support

From momentum conservation in intermediate state

 $x_i \in [-1, 1]$ i = 1, 2, 3

 \Rightarrow cube in 3D space

From total momentum conservation

 $x_1 + x_2 + x_3 = 0$

 \Rightarrow slice of the cube

hexagon





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Partonic interpretation



Evolution equations

Flavor separation, singlet
$$= \mathfrak{S}_S^{\pm} = u + d + s + \cdots$$

non-singlet

$$\begin{split} \mathfrak{S}_{\mathrm{NS}_{1}}^{\pm} &= \mathfrak{S}_{u}^{\pm} - \mathfrak{S}_{d}^{\pm} \\ \mathfrak{S}_{\mathrm{NS}_{2}}^{\pm} &= \mathfrak{S}_{u}^{\pm} + \mathfrak{S}_{d}^{\pm} - 2\mathfrak{S}_{s}^{\pm} \end{split}$$

. . .

$$\begin{split} \mu^{2} \frac{\partial}{\partial \mu^{2}} \mathfrak{S}_{\mathrm{NS}_{i}}^{\pm} &= -\frac{\alpha_{s}(\mu)}{4\pi} \mathbb{H}_{\mathrm{NS}} \mathfrak{S}_{\mathrm{NS}_{i}}^{\pm} \\ \mu^{2} \frac{\partial}{\partial \mu^{2}} \begin{pmatrix} \mathfrak{S}_{S}^{\pm} \\ \mathfrak{F}^{\pm} \end{pmatrix} &= -\frac{\alpha_{s}(\mu)}{4\pi} \begin{pmatrix} \mathbb{H}_{qq}^{\pm} & \mathbb{H}_{qg}^{\pm} \\ \mathbb{H}_{gq}^{\pm} & \mathbb{H}_{gg}^{\pm} \end{pmatrix} \begin{pmatrix} \mathfrak{S}_{S}^{\pm} \\ \mathfrak{F}^{\pm} \end{pmatrix} \\ \mu^{2} \frac{\partial}{\partial \mu^{2}} [E, H] &= -\frac{\alpha_{s}(\mu)}{4\pi} \mathbb{H}_{\mathrm{co}} [E, H] \end{split}$$

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Kernels action



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Grids



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The goal
$$\mathfrak{S}_{\rm NS}({\pmb x}) = \sum_{i,j} \mathscr{G}_{ij}({\pmb x}) \mathfrak{S}_{\rm NS}^{ij} \qquad \mathfrak{S}_{\rm NS}^{ij} = \mathfrak{S}_{\rm NS}({\pmb x}(i,j))$$
so that

$$\mu^2 \frac{d}{d\mu^2} \mathfrak{S}^{ij}_{\mathrm{NS}} = -a_s(\mu) \sum_{i'j'} \mathbb{H}_{\mathrm{NS}}{}^{ij}_{i'j'} \mathfrak{S}^{i'j'}_{\mathrm{NS}} \qquad \mathbb{H}_{\mathrm{NS}}{}^{ij}_{i'j'} = \mathbb{H}_{\mathrm{NS}} \mathcal{G}_{i'j'}(\pmb{x}(i,j))$$

Two independent steps:

1) fix the discretization $\boldsymbol{x}(i,j)$ 2) fix the weights $\mathcal{G}_{ij}(\boldsymbol{x})$



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Nice properties to have:

1) Denser grid towards x = 0

2) Compatible with hexagonal symmetry

First guess: equispaced grid

$$\boldsymbol{x}(i,j) = \left(\frac{i}{N}, \frac{j}{N}, \frac{-i-j}{N}\right) \qquad i, j, i+j \in [-N, N]$$

Denser grid towards x = 0

We can do better!



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Step 1: change coordinate system $\label{eq:coordinate} \boldsymbol{x} \to (\phi,r)$ $r = \|\boldsymbol{x}\|_\infty = \max(|x_1|,|x_2|,|x_3|) \in [0,1]$



Now we can discretize angle/radius:

 $(i,j) \rightarrow (\phi_i, r_j)$ $i \in [0,6n)$ $j \in [0,M]$

$$\begin{split} \phi_i &= \frac{i \mod n}{n} + \left[\frac{i}{n}\right] &\implies \quad i = n\phi_i \\ r_j &= \left[\cosh\left(\frac{j-M}{Mc}\right)\right]^{-\alpha} &\implies \quad j = M\left[1 + c \, \operatorname{acosh}\left(\frac{1}{\sqrt[\alpha]{r_j}}\right)\right] \end{split}$$

Generalizing to continuous variables in $(\mathfrak{f},\mathfrak{r})\in[0,6n)\times[0,M]$

$$\begin{split} &\mathfrak{f}(\boldsymbol{x}) = n\phi(\boldsymbol{x}) \\ &\mathfrak{r}(\boldsymbol{x}) = M\left[1 + c \, \operatorname{acosh}\left(\frac{1}{\sqrt[\alpha]{r(\boldsymbol{x})}}\right)\right] \end{split}$$

 $r_0 = x_{\min}$

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Weights



Weights



One kernel

$$\boldsymbol{x}^{ij} = \boldsymbol{x}(\phi_i, r_j)$$

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One kernel

$$\begin{aligned} \widehat{\mathcal{H}}_{12}\mathcal{G}_{i'j'}^{+}(\boldsymbol{x}^{ij}) &= \delta_{i \neq i' \text{ or } j \neq j'} \int_{v_{\text{min}}^{i'j'}}^{v_{\text{max}}^{ij}} dv \frac{x_{1}^{ij} \Theta(x_{1}^{ij}, -v)}{v(x_{1}^{ij} - v)} \mathcal{G}_{i'j'}^{+}(x_{1}^{ij} - v, x_{2}^{ij} + v, x_{3}^{ij}) \\ &+ \delta_{ii'} \delta_{jj'} \mathcal{G}_{ij}^{+}(\boldsymbol{x}^{ij}) \left[\theta(x_{1}^{ij}) \log \left(1 - \frac{x_{1}^{ij}}{v_{\text{min}}^{ij}} \right) + \theta(-x_{1}^{ij}) \log \left(1 - \frac{x_{1}^{ij}}{v_{\text{max}}^{ij}} \right) \right] \\ &+ \delta_{ii'} \delta_{jj'} \int_{v_{\text{min}}^{ij}}^{v_{\text{max}}^{ij}} dv \frac{x_{1}^{ij} \Theta(x_{1}^{ij}, -v)}{v(v - x_{1}^{ij})} \left[\mathcal{G}_{ij}^{+}(\boldsymbol{x}^{ij}) - \mathcal{G}_{ij}^{+}(x_{1}^{ij} - v, x_{2}^{ij} + v, x_{3}^{ij}) \right] \end{aligned}$$

Λ ...

1.1.

Off-diagonal Integral of plus-prescription outside the support Plus-prescription, diagonal part



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Few numbers

- With our $\mathcal{G}_{ij}(\boldsymbol{x})$ only few % of entries $\mathbb{H}_{_{NS},ij}^{i'j'}$ are non-zero \Rightarrow custom sparse-algebra implementation
- (a) with 'just' 120×25 points, kernels sizes of $O(10 \sim 20 \text{MB})$
- average interpolation accuracy ~ 10^{-3} Worse towards the edge of the physical region: for $r \leq 0.9$ accuracy improves to ~ 10^{-4}

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Evolution results and comments



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3-gluon



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Conclusions and outlooks

- The two codes are available on github at VladimirovAlexey/Snowflake 🎯 and at QCDatHT/honeycomb 🜑
- 2 Both should be easily integrable in existing programs
- Evolution effects are not overwhelmingly strong, but also not negligible 3



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