

Angular structure of Drell-Yan reaction in TMD factorization

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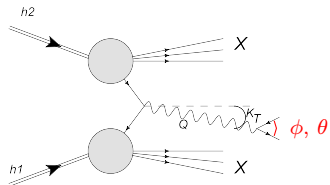
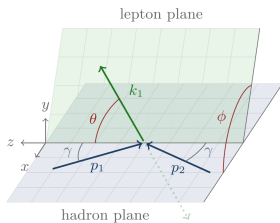


figure from [Gehrmann–De Ridder et al.
2301.11827]



Q^2 momentum of hard probe
 q_T transverse momentum
 y rapidity of hard probe
 θ, ϕ angles of lepton pair (in CS frame)



Angular decomposition of EW-boson production

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

8 angular structure functions $A_0, \dots, 7$ + 1 integrated (normalization)

alternative notations:

$$\frac{d\sigma}{d^4q d\Omega} = \frac{\alpha^2}{2\pi N_c Q^2 s^2} (W_T (1 + \cos^2\theta) + W_L (1 - \cos^2\theta) + W_\Delta \sin 2\theta \cos\phi + W_{\Delta\Delta} \sin^2\theta \cos 2\phi) \\ \frac{dN}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda + 3} \left(1 + \lambda \cos^2\theta + \mu \sin 2\theta \cos\phi + \frac{\nu}{2} \sin^2\theta \cos 2\phi \right. \\ \left. + \tau \sin\theta \cos\phi + \eta \cos\theta + \xi \sin^2\theta \sin 2\phi + \zeta \sin 2\theta \sin\phi + \chi \sin\theta \sin\phi \right)$$



Angular decomposition of EW-boson production

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

8 angular structure functions $A_{0,\dots,7}$ + 1 integrated (normalization)

Measurements (proton Drell-Yan):

CMS	$A_{0,\dots,7}$	large q_T -bins, 2 bins in y
LHCb	$A_{0,\dots,3}$	large q_T -bins, $y > 2$
ATLAS	$A_{0,\dots,7}$	2.5GeV bins, 3 bins in y
E866	$A_{0,\dots,3}$	1D measurement



Angular decomposition of EW-boson production

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

8 angular structure functions $A_0, \dots, 7$ + 1 integrated (normalization)

Practically, they are not studied in TMD approach

double-BM effects

[Barone, Melis, Prokudin, 1009.3423]

[Lu, Schmidt, 1107.4693]

[Liu, Ma, 1201.2472]

[Wang, Mao, Lu, 1805.03017]



Angular decomposition of EW-boson production

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

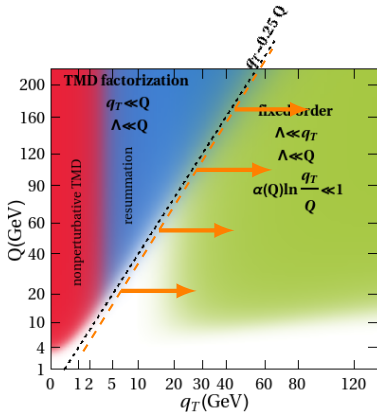
8 angular structure functions $A_0, \dots, 7$ + 1 integrated (normalization)

Leading power TMD factorization
gives only 3 non-zero coefficients $A_{2,4,5}$.

The majority of coefficients are power-suppressed



Different types of power corrections



1. q_T/Q -corrections
Y-term

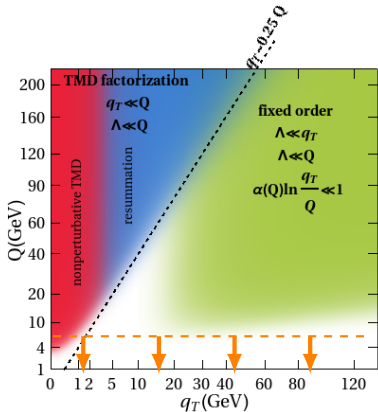
2. Λ/Q & M/Q -corrections
higher-twist
target-mass

3. k_T/Q -corrections
kinematic

[AV,2307.13054]



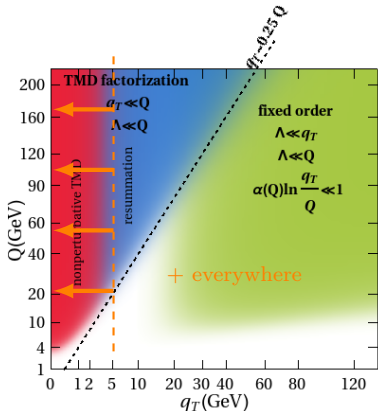
Different types of power corrections



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 2. Λ/Q & M/Q -corrections
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 3. k_T/Q -corrections
kinematic
- [AV,2307.13054]



Different types of power corrections



1. q_T/Q -corrections
Y-term
 2. Λ/Q & M/Q -corrections
higher-twist
target-mass
 3. k_T/Q -corrections
kinematic
- [AV,2307.13054]

Kinematic power corrections (KPC) are largest corrections in the “natural” domain of TMD factorization $q_T \ll Q$



$$W^{\mu\nu} = \frac{1}{N_c} \int \frac{d^2b}{(2\pi)^2} e^{-i(q_T b)} \left\{ \right.$$

$$\Phi_2 \times \Phi_2$$

$$+ \frac{1}{Q} \left(D \Phi_2 \times \Phi_2 + \Phi_2 \times \Phi_3 \right)$$

$$+ \frac{1}{Q^2} \left(D^2 \Phi_2 \times \Phi_2 + D \Phi_2 \times \Phi_3 + \Phi_3 \times \Phi_3 + \Phi_2 \times \Phi_4 + \frac{\Phi_2 \times \Phi_2}{b^2} \right)$$

$$+ \frac{1}{Q^3} \left(D^3 \Phi_2 \times \Phi_2 + \dots + \Phi_3 \times \Phi_4 + \dots + \frac{D \Phi_2 \times \Phi_2}{b^2} \right)$$

$$+ \frac{1}{Q^4} \left(D^4 \Phi_2 \times \Phi_2 + \dots + \Phi_2 \times \Phi_5 + \dots + \frac{D \Phi_2 \times \Phi_2}{b^2} + \frac{\Phi_2 \times \Phi_2}{b^4} \right)$$

$$+ \dots \left. \right\}^{\mu\nu},$$

$$D \sim \frac{\partial}{\partial b^\mu}$$

Φ_N is TMD of twist-N



$$D \sim \frac{\partial}{\partial b^\mu}$$

Φ_N is TMD of twist-N

$$W^{\mu\nu} = \frac{1}{N_c} \int \frac{d^2b}{(2\pi)^2} e^{-i(q_T b)} \left\{ \right.$$

$$\Phi_2 \times \Phi_2$$

LP ✓

$$+ \frac{1}{Q} \left(D \Phi_2 \times \Phi_2 + \Phi_2 \times \Phi_3 \right)$$

NLP ✓

$$+ \frac{1}{Q^2} \left(D^2 \Phi_2 \times \Phi_2 + D \Phi_2 \times \Phi_3 + \Phi_3 \times \Phi_3 + \Phi_2 \times \Phi_4 + \frac{\Phi_2 \times \Phi_2}{b^2} \right)$$

NNLP

Balitsky, tomorrow

$$+ \frac{1}{Q^3} \left(D^3 \Phi_2 \times \Phi_2 + \dots + \Phi_3 \times \Phi_4 + \dots + \frac{D \Phi_2 \times \Phi_2}{b^2} \right)$$

$$+ \frac{1}{Q^4} \left(D^4 \Phi_2 \times \Phi_2 + \dots + \Phi_2 \times \Phi_5 + \dots + \frac{D \Phi_2 \times \Phi_2}{b^2} + \frac{\Phi_2 \times \Phi_2}{b^4} \right)$$

$$+ \dots \left. \right\}^{\mu\nu},$$



$$W^{\mu\nu} = \frac{1}{N_c} \int \frac{d^2b}{(2\pi)^2} e^{-i(q_T b)} \left\{ \right.$$

- ▶ Same TMDs as at LP
- ▶ Same perturbative part as LP
- ▶ Non-vanishing $q_T = 0$, and larger than $\frac{\Lambda}{Q}$
- ▶ Restore gauge and frame-invariance

$$\begin{aligned}
 & \Phi_2 \times \Phi_2 \\
 & + \frac{1}{Q} \left(D \Phi_2 \times \Phi_2 + \Phi_2 \times \Phi_3 \right) \\
 & + \frac{1}{Q^2} \left(D^2 \Phi_2 \times \Phi_2 + D \Phi_2 \times \Phi_3 + \Phi_3 \times \Phi_3 + \Phi_2 \times \Phi_4 + \frac{\Phi_2 \times \Phi_2}{b^2} \right) \\
 & + \frac{1}{Q^3} \left(D^3 \Phi_2 \times \Phi_2 + \dots + \Phi_3 \times \Phi_4 + \dots + \frac{D\Phi_2 \times \Phi_2}{b^2} \right) \\
 & + \frac{1}{Q^4} \left(D^4 \Phi_2 \times \Phi_2 + \dots + \Phi_2 \times \Phi_5 + \dots + \frac{D\Phi_2 \times \Phi_2}{b^2} + \frac{\Phi_2 \times \Phi_2}{b^4} \right) \\
 & + \dots \left. \right\}^{\mu\nu},
 \end{aligned}$$

resummed KPC
this talk



DY-hadron tensor with resummed KPC [AV,2307.13054]

$$\begin{aligned}
 W_{\text{KPC}}^{\mu\nu} = & -\frac{4p_1^+ p_2^-}{N_c} C_0 \left(\frac{Q^2}{\mu^2} \right) \int d\xi_1 d\xi_2 \int d^4 k_1 d^4 k_2 \delta^4(q - k_1 - k_2) \\
 & \delta(k_1^+ - \xi_1 p_1^+) \delta(k_2^- - \xi_2 p_2^-) \delta(k_1^2) \delta(k_2^2) \left\{ \right. \\
 & ((k_1 k_2) g^{\mu\nu} - k_1^\mu k_2^\nu - k_2^\mu k_1^\nu) \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^-]} + \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^- \gamma^5]} \right) \\
 & \left. + i \epsilon^{\mu\nu\alpha\beta} k_1^\alpha k_2^\beta \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^- \gamma^5]} - \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^-]} \right) + t_{\alpha\beta}^{\mu\nu} \Phi_{\bar{n}11}^{[i\sigma^{\alpha+} \gamma^5]} \Phi_{n11}^{[i\sigma^{\beta-} \gamma^5]} \right\} \\
 & \Phi^{[\Gamma]} \text{ is ordinary TMD distribution } \sim \langle p | \bar{q} \dots \Gamma \dots q | r \rangle
 \end{aligned}$$



DY-hadron tensor with resummed KPC [AV,2307.13054]

$$W_{\text{KPC}}^{\mu\nu} = -\frac{4p_1^+ p_2^-}{N_c} C_0 \left(\frac{Q^2}{\mu^2} \right) \int d\xi_1 d\xi_2 \int d^4 k_1 d^4 k_2 \delta^4(q - k_1 - k_2)$$

$$\delta(k_1^+ - \xi_1 p_1^+) \delta(k_2^- - \xi_2 p_2^-) \delta(k_1^2) \delta(k_2^2) \left\{$$

$$\begin{aligned}
 & ((k_1 k_2) g^{\mu\nu} - k_1^\mu k_2^\nu - k_2^\mu k_1^\nu) \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^-]} + \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^- \gamma^5]} \right) \\
 & + i \epsilon^{\mu\nu\alpha\beta} k_1^\alpha k_2^\beta \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^- \gamma^5]} - \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^-]} \right) + t_{\alpha\beta}^{\mu\nu} \Phi_{\bar{n}11}^{[i\sigma^{\alpha+} \gamma^5]} \Phi_{n11}^{[i\sigma^{\beta-} \gamma^5]} \left. \right\}
 \end{aligned}$$


$$\int d^2 k_{1T} d^2 k_{2T} \delta^{(2)}(q_T - k_{1T} - k_{2T})$$

at LP



DY-hadron tensor with resummed KPC [AV,2307.13054]

$$\begin{aligned}
 W_{\text{KPC}}^{\mu\nu} = & -\frac{4p_1^+ p_2^-}{N_c} C_0 \left(\frac{Q^2}{\mu^2} \right) \int d\xi_1 d\xi_2 \int d^4 k_1 d^4 k_2 \delta^4(q - k_1 - k_2) \\
 & \delta(k_1^+ - \xi_1 p_1^+) \delta(k_2^- - \xi_2 p_2^-) \delta(k_1^2) \delta(k_2^2) \left\{ \right. \\
 & \left. \left((k_1 k_2) g^{\mu\nu} - k_1^\mu k_2^\nu - k_2^\mu k_1^\nu \right) \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^-]} + \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^- \gamma^5]} \right) \right. \\
 & \left. + i \epsilon^{\mu\nu\alpha\beta} k_1^\alpha k_2^\beta \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^- \gamma^5]} - \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^-]} \right) + t_{\alpha\beta}^{\mu\nu} \Phi_{\bar{n}11}^{[i\sigma^{\alpha+} \gamma^5]} \Phi_{n11}^{[i\sigma^{\beta-} \gamma^5]} \right\}
 \end{aligned}$$



$$g_T^{\mu\nu}$$

at LP



DY-hadron tensor with resummed KPC [AV,2307.13054]

$$W_{\text{KPC}}^{\mu\nu} = -\frac{4p_1^+ p_2^-}{N_c} C_0 \left(\frac{Q^2}{\mu^2} \right) \int d\xi_1 d\xi_2 \int d^4 k_1 d^4 k_2 \delta^4(q - k_1 - k_2) \left\{ \begin{aligned} &\delta(k_1^+ - \xi_1 p_1^+) \delta(k_2^- - \xi_2 p_2^-) \delta(k_1^2) \delta(k_2^2) \left\{ \right. \\ &((k_1 k_2) g^{\mu\nu} - k_1^\mu k_2^\nu - k_2^\mu k_1^\nu) \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^-]} + \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^- \gamma^5]} \right) \\ &\left. + i\epsilon^{\mu\nu\alpha\beta} k_1^\alpha k_2^\beta \left(\Phi_{\bar{n}11}^{[\gamma^+]} \Phi_{n11}^{[\gamma^- \gamma^5]} - \Phi_{\bar{n}11}^{[\gamma^+ \gamma^5]} \Phi_{n11}^{[\gamma^-]} \right) + t_{\alpha\beta}^{\mu\nu} \Phi_{\bar{n}11}^{[i\sigma^{\alpha+} \gamma^5]} \Phi_{n11}^{[i\sigma^{\beta-} \gamma^5]} \right\} \end{aligned} \right.$$

Same as at LP!

- ▶ Argument Q^2 [instead of $2q^+ q^-$ at LP]
- ▶ μ -dependence canceled by $\Phi(x, k_T; \mu, Q^2)$
- ▶ $\zeta \bar{\zeta} = Q^4$ [instead of $(2q^+ q^-)^2$ at LP]
- ▶ The summation of KPC is done at $\zeta/\bar{\zeta} = 1$ (general form see [AV,2307.13054])



Main difference with LP is the convolution integral

$$\text{KPC} \quad \frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2) \\ \rho_{pp}(\text{polarizations}) F(\xi_1(x_1, \mathbf{k}_{1,2}), \mathbf{k}_1^2, Q, Q^2) F(\xi_2(x_2, \mathbf{k}_{1,2}), \mathbf{k}_2^2, Q, Q^2)$$

$$\text{LP} \quad \frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2) \\ \rho_{pp}^{\text{LP}}(\text{polarizations}) F(x_1, \mathbf{k}_1^2, Q, Q^2) F(x_2, \mathbf{k}_2^2, Q, Q^2)$$



Main difference with LP is the convolution integral

$$\text{KPC} \quad \frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2)$$

$$\underline{\rho_{pp}(\text{polarizations})} F(\xi_1(x_1, \mathbf{k}_{1,2}), \mathbf{k}_1^2, Q, Q^2) F(\xi_2(x_2, \mathbf{k}_{1,2}), \mathbf{k}_2^2, Q, Q^2)$$

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$$\underline{\rho_{pp}^{\text{LP}}(\text{polarizations})} F(x_1, \mathbf{k}_1^2, Q, Q^2) F(x_2, \mathbf{k}_2^2, Q, Q^2)$$

$$\rho_{pp} = \rho_{pp}^{\text{LP}} + \underbrace{\frac{\mathbf{k}_T}{Q} + \frac{\mathbf{k}_T^2}{Q^2} + \dots}_{\text{restoration of g.inv.}}$$



Main difference with LP is the convolution integral

$$\text{KPC} \quad \frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2) \\ \rho_{pp}(\text{polarizations}) F(\xi_1(x_1, \mathbf{k}_{1,2}), \mathbf{k}_1^2, Q, Q^2) F(\xi_2(x_2, \mathbf{k}_{1,2}), \mathbf{k}_2^2, Q, Q^2)$$

$$\text{LP} \quad \frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2) \\ \rho_{pp}^{\text{LP}}(\text{polarizations}) F(\underline{x}_1, \mathbf{k}_1^2, Q, Q^2) F(\underline{x}_2, \mathbf{k}_2^2, Q, Q^2)$$

$$\xi_1 = \frac{x_1}{2} \left(1 + \frac{\mathbf{k}_1^2 - \mathbf{k}_2^2 + \sqrt{\lambda(\mathbf{k}_1^2, \mathbf{k}_2^2, \tau^2)}}{\tau^2} \right), \quad \xi_2 = \frac{x_2}{2} \left(1 + \frac{\mathbf{k}_2^2 - \mathbf{k}_1^2 + \sqrt{\lambda(\mathbf{k}_1^2, \mathbf{k}_2^2, \tau^2)}}{\tau^2} \right)$$



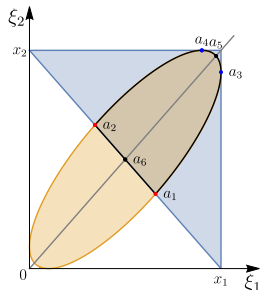
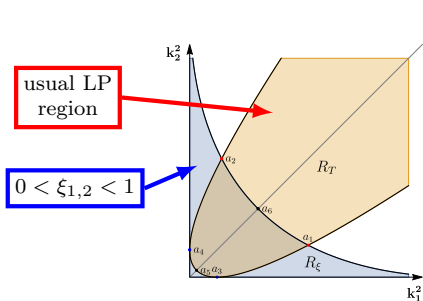
Main difference with LP is the convolution integral

KPC
$$\frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2)$$

$$\rho_{pp}(\text{polarizations}) F(\xi_1(x_1, \mathbf{k}_{1,2}), \mathbf{k}_1^2, Q, Q^2) F(\xi_2(x_2, \mathbf{k}_{1,2}), \mathbf{k}_2^2, Q, Q^2)$$

LP
$$\frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2)$$

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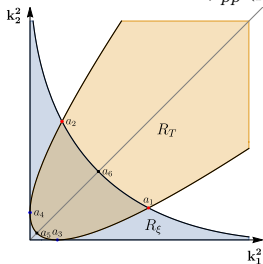
Main difference with LP is the convolution integral

KPC
$$\frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2)$$

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LP
$$\frac{d\sigma}{dq_T} = |C_V(Q)|^2 \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(\mathbf{q}_T - \mathbf{k}_1 - \mathbf{k}_2)$$

$$\rho_{pp}^{\text{LP}}(\text{polarizations}) F(x_1, \mathbf{k}_1^2, Q, Q^2) F(x_2, \mathbf{k}_2^2, Q, Q^2)$$



effectively **+1 integration**

artemide v3.0 α
 main DY set (~ 600 pt)
 +bin integrations +accept.corr.

LP

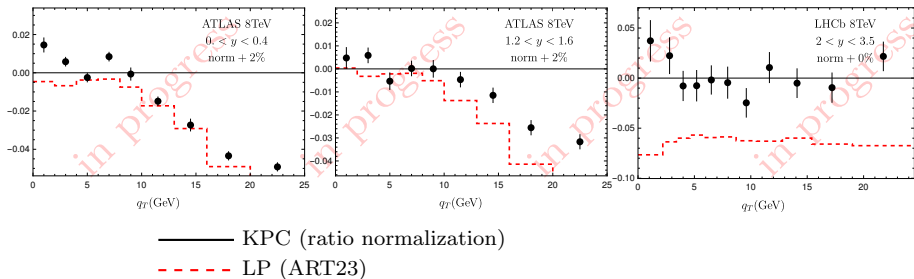
accurate ~ 1 min
 approx. ~ 5 sec.

KPC

accurate ~ 2 hours
 approx. ~ 4 min



Unpolarized DY = normalization for angular coefficients
 Computation with ART23 input [V.Moos,et al,2305.07473]

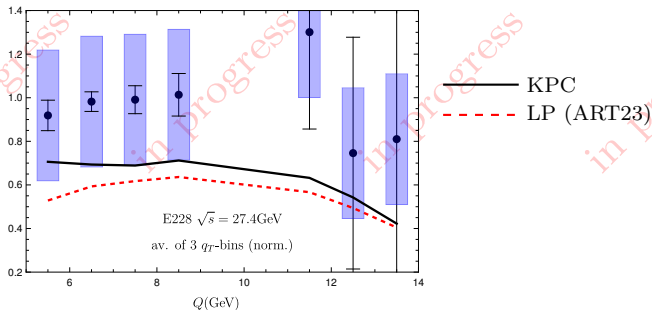


Main effects are at

- ▶ Large q_T -tail
- ▶ Normalization at forward rapidity & low energy



Unpolarized DY = normalization for angular coefficients
 Computation with ART23 input [V.Moos,et al,2305.07473]

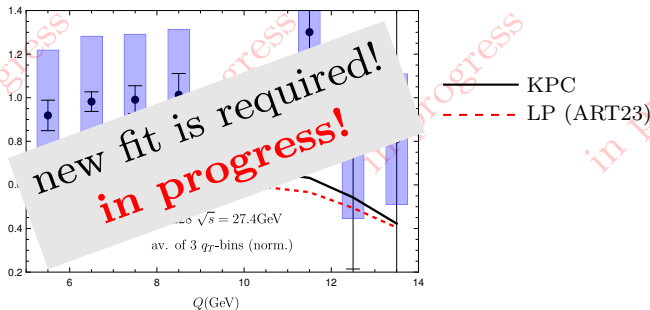


Main effects are at

- ▶ Large q_T -tail
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Unpolarized DY = normalization for angular coefficients
Computation with ART23 input [V.Moos,et al,2305.07473]



Main effects are at

- ▶ Large q_T -tail
- ▶ Normalization at forward rapidity & low energy



$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

I use ART23 TMDPDF as unpolarized input.
Some tension with data is due to it.



A_4

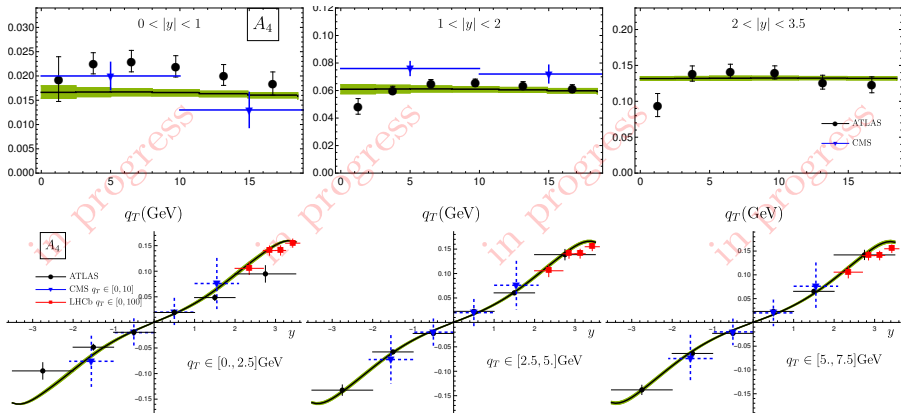
$$\frac{d\sigma}{dp_T^Z dy^Z dm^2 d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^2} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

$$A_4 \sim (av)_l(av)_q \left[2 - \frac{\mathbf{k}_1^2 + \mathbf{k}_2^2}{Q^2} \right] [f_{1q}f_{1\bar{q}} - f_{1\bar{q}}f_{1q}]$$

- ▶ Leading power (!!!)
- ▶ Sensitive to a different EW-coupling
- ▶ Only unpolarized contribution
- ▶ Sensitive to a **difference** of quark/anti-quark distributions (!!!)

why we did not include it into standard TMD fits?



A_4 

A_2

$$\frac{d\sigma}{dp_1^Z dy^Z dm^2 d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_1^Z dy^Z dm^2} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0 (1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

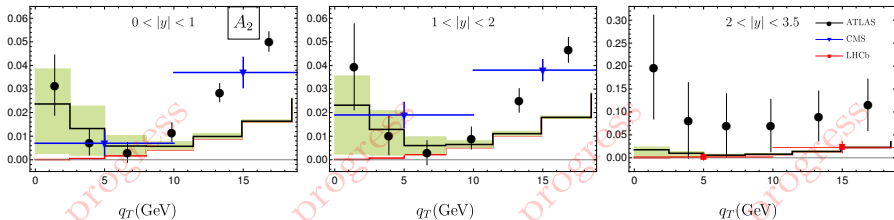
$$A_2 \sim (v^2 + a^2)_l \left[2(\hat{q}\mathbf{k}_1)(\hat{q}\mathbf{k}_2) - (\mathbf{k}_1\mathbf{k}_2) + \dots \right] \left[(v^2 - a^2)_q \frac{h_1^\perp h_1^\perp}{M^2} + (v^2 + a^2)_q \frac{f_1 f_1}{Q^2} \right]$$

- ▶ Double BM at leading power
- ▶ Mixing with $f_1 f_1$

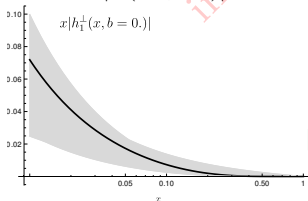


Double Boer-Mulders part at LHC

A_2



$$q_T(\text{GeV}) \quad x|E(-x, 0, x)|$$



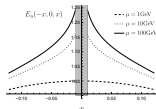
$$h_{1\perp}^{\perp} = N e^{-0.1b^2} x^a (1-x)^5$$

$$h_{1q}^{\perp} = -h_{1\bar{q}}^{\perp}$$

$$N \simeq (3 \pm 1) \cdot 10^{-3}, \quad a = -1.5 \pm 0.2$$

[Kovchegov, Santiago, 2209.03538] $\Rightarrow a \sim +1$

but see Rodini's talk on Wednesday

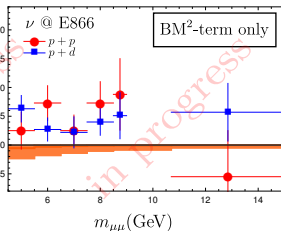
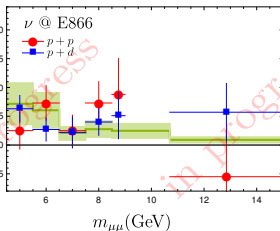
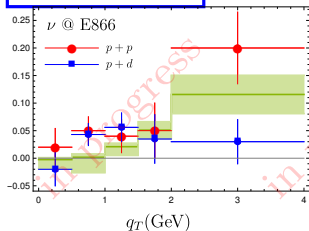


Double Boer-Mulders part at low-energy

$$\nu = \frac{2A_2}{2 + A_0}$$

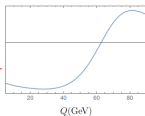
← complicated composition of terms...

Warning!
E866 is 1D measurement
i.e. partially outside of fac.region

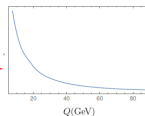


$$A_2 \sim \left[\underbrace{(v^2 - a^2)_q}_{< 0 \text{ for } Z} \frac{h_1^\perp h_1^\perp}{M^2} + (v^2 + a^2)_q \underbrace{\frac{f_1 f_1}{Q^2}}_{\text{grows at } Q \rightarrow 0} \right]$$

< 0 for Z
> 0 for γ



grows at
 $Q \rightarrow 0$



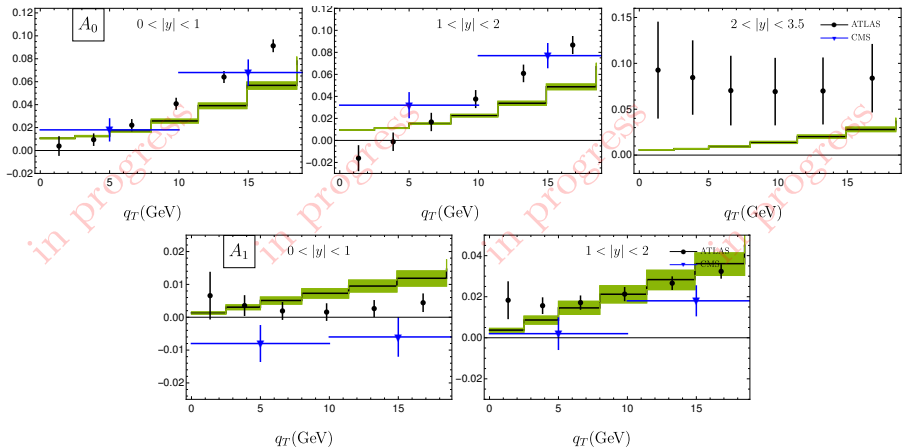
$A_{0,1}$

$$\frac{d\sigma}{dp_T^2 dy^2 dm^2 d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^2 dy^2 dm^2} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0 (1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

$$A_0 \sim (v^2 + a^2)_l \left[\frac{\mathbf{k}_1^2 + \mathbf{k}_2^2}{Q^2} \right] \left[(v^2 - a^2)_q \frac{h_1^\perp h_1^\perp}{M^2} + (v^2 + a^2)_q \frac{f_1 f_1}{Q^2} \right] \\ A_1 \sim (v^2 + a^2)_l \left[\frac{(\hat{q}\mathbf{k}_1) - (\hat{q}\mathbf{k}_2)}{Q} \right] \left[(v^2 - a^2)_q \frac{h_1^\perp h_1^\perp}{M^2} + (v^2 + a^2)_q \frac{f_1 f_1}{Q^2} \right]$$

- ▶ A_0 and A_1 are similar.
- ▶ $A_0 \sim Q^{-2}$, symmetric in y
- ▶ $A_1 \sim Q^{-1}$, anti-symmetric in y
- ▶ Non-zero at $q_T = 0$ (although small)



$A_{0,1}$


Double Boer-Mulders part is negligible, $A_0^{hh} \sim 10^{-4}$
 $A_1^{hh} = 0$ since $h_q = h_{\bar{q}}$

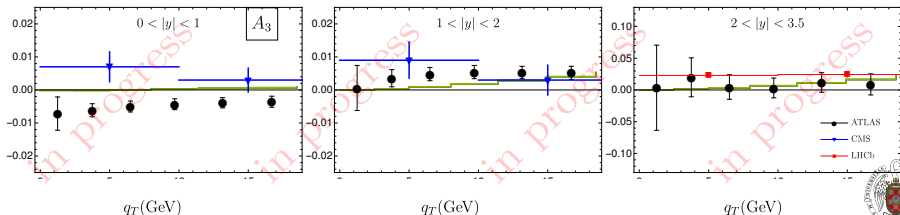


A₃

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

$$A_3 \sim (av)_l(av)_q \left[\frac{(\hat{q}\mathbf{k}_1) - (\hat{q}\mathbf{k}_2)}{Q} \right] [f_{1q}f_{1\bar{q}} - f_{1\bar{q}}f_{1q}]$$

► Similar to A_4 but with Q^{-1} suppression



A_{5,6,7}

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos\theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0(1 - 3\cos^2\theta) + A_1 \sin 2\theta \cos\phi \right. \\ \left. + \frac{1}{2} A_2 \sin^2\theta \cos 2\phi + A_3 \sin\theta \cos\phi + A_4 \cos\theta \right. \\ \left. + A_5 \sin^2\theta \sin 2\phi + A_6 \sin 2\theta \sin\phi + A_7 \sin\theta \sin\phi \right\}.$$

$$A_5 \sim (v^2 + a^2)_l(va - av)_q \left[\frac{2(\hat{q}\mathbf{k}_1)(\hat{q}\mathbf{k}_2) - (\mathbf{k}_1\mathbf{k}_2)}{M^2} \right] \left[h_{1q}^\perp h_{1\bar{q}}^\perp - h_{1\bar{q}}^\perp h_{1q}^\perp \right] \\ A_6 \sim (v^2 + a^2)_l(va - av)_q \left[\frac{(\mathbf{k}_1^2 + \mathbf{k}_2^2)((\hat{q}\mathbf{k}_1) + (\hat{q}\mathbf{k}_2))}{QM^2} \right] \left[h_{1q}^\perp h_{1\bar{q}}^\perp - h_{1\bar{q}}^\perp h_{1q}^\perp \right] \\ A_7 = 0$$

- ▶ A_{5,6,7} are T-odd
- ▶ A₅ is leading power [A₅ for A₂ alike A₄ for σ],
- ▶ A₆ ∼ Q⁻¹
- ▶ A_{5,6} are non-zero only due to Z/γ interference, thus **extra Γ_Z/Q-factor**
- ▶ ATLAS measured A₅ ≲ 0.2 – 1%, A₆ ≲ 0. – 0.3%, A₇ ≲ 0. – 0.1% (+uncertainties)



Conclusion

Angular coefficients provide extra information for TMD studies.
Each coefficient gives access to a unique portion of information.

Main observations

- ▶ TMD factorization with included KPC work very well but an update of main fits is required.
- ▶ A_4 coefficient should be added to the main f_1 -DY-data pool.
- ▶ A_2 measured at LHC indicates non-zero Boer-Mulders function (with $h_q h_{\bar{q}} < 0$)
- ▶ Power suppressed effects (such as terms $\sim f_1 f_1$ in A_2) became numerically significant/dominant at $Q \sim 10\text{GeV}$.

Work still in progress!



Backup



Remainder: EM gauge invariance

$$q^\mu W_{\mu\nu} = 0, \quad \Leftrightarrow \quad \left(q^+ \bar{n}^\mu + q^- n^\mu - i \frac{\partial}{\partial b_\mu} \right) \widetilde{W}_{\mu\nu} = 0$$

Remainder: frame (reparametrization) invariance

$$\text{I: } \begin{cases} n^\mu \rightarrow n'^\mu = n^\mu + \frac{\Delta^\mu}{q^-} - \frac{\Delta^2}{2(q^-)^2} \bar{n}^\mu, \\ \bar{n}^\mu \rightarrow \bar{n}'^\mu = \bar{n}^\mu, \end{cases} \quad \text{II: } \begin{cases} n^\mu \rightarrow n'^\mu = n^\mu, \\ \bar{n}^\mu \rightarrow \bar{n}'^\mu = \bar{n}^\mu + \frac{\bar{\Delta}^\mu}{q^+} - \frac{\bar{\Delta}^2}{2(q^+)^2} n^\mu, \end{cases}$$

For physical variables in implies (here for I)

$$k_{1T}^\mu \rightarrow k'_{1T}{}^\mu = k_{1T}^\mu, \quad k_{2T}^\mu \rightarrow k'_{2T}{}^\mu = k_{2T}^\mu - \Delta^\mu - ((k_2 \Delta) - \Delta^2) \frac{\bar{n}^\mu}{q^-}.$$

$$x_1 \rightarrow x'_1 = x_1 \left(1 + \frac{2(k_2 \Delta) - \Delta^2}{2q^+ q^-} \right), \quad x_2 \rightarrow x'_2 = x_2,$$

These invariances are restored by KPCs!