



# Glimmers from the Axiverse

David J. E. Marsh  
w/ Naomi Gendler, Jakob Moritz, Liam McAllister, 2309.13145

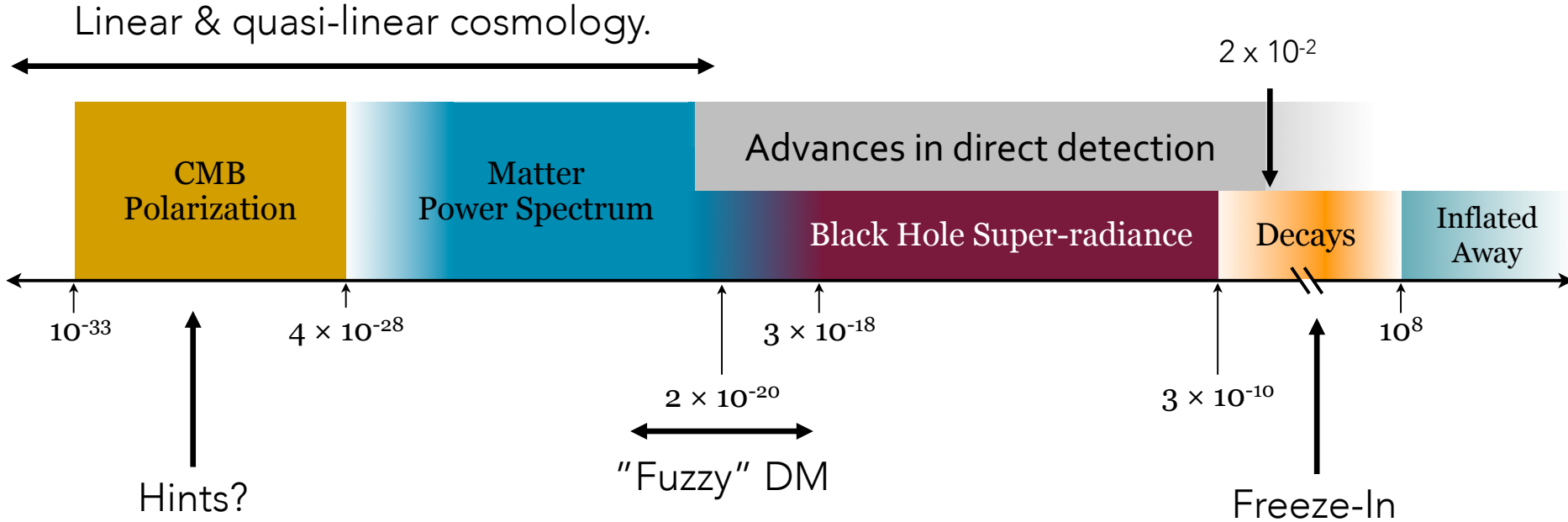


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# Axiverse in 2023



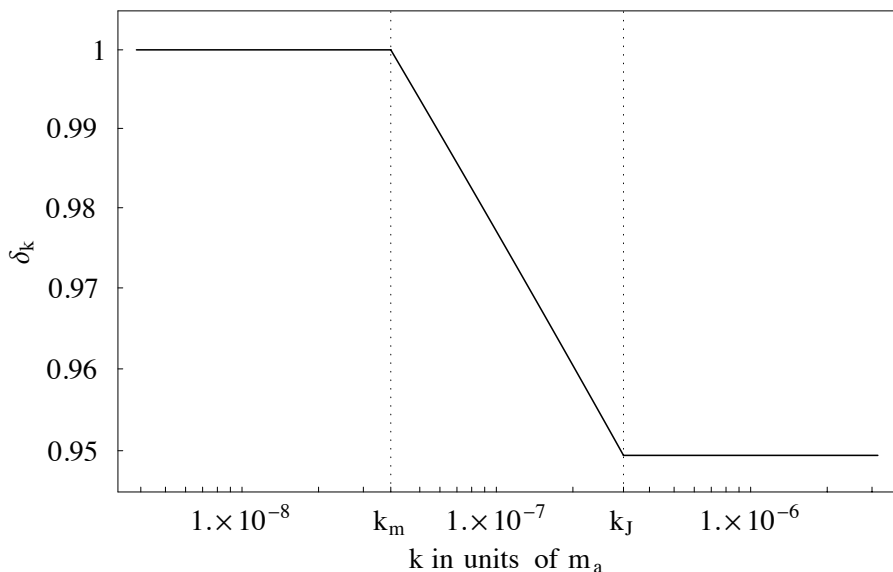
All of the suggested probes have been explored and become precise, along with some (unexpected?) new developments in cosmology and direct detection.

# LINEAR++ COSMOLOGY

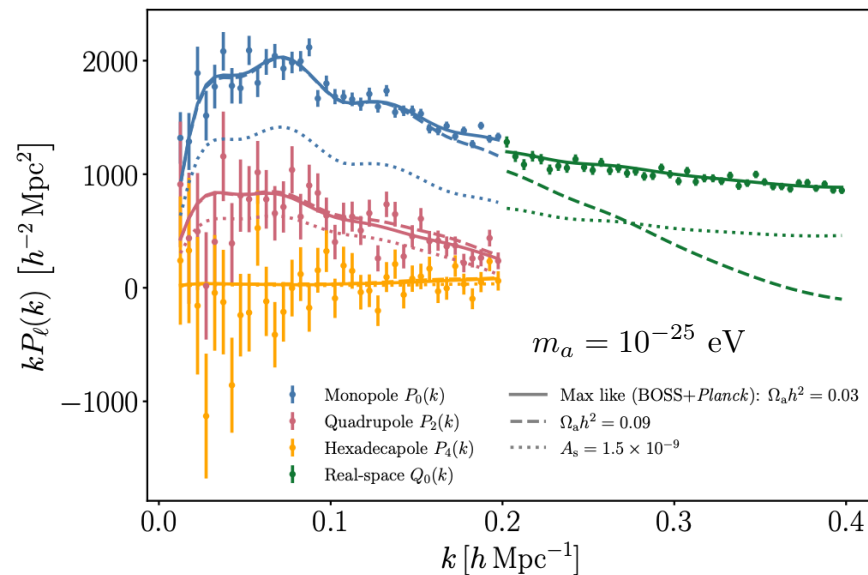


Matter  
Power Spectrum

Recent advances in **quasi-linear and non-linear modelling** (halo models, EFTofLSS, emulators) allow precision limits from smaller scales  $\rightarrow$  **probe larger axion masses**.



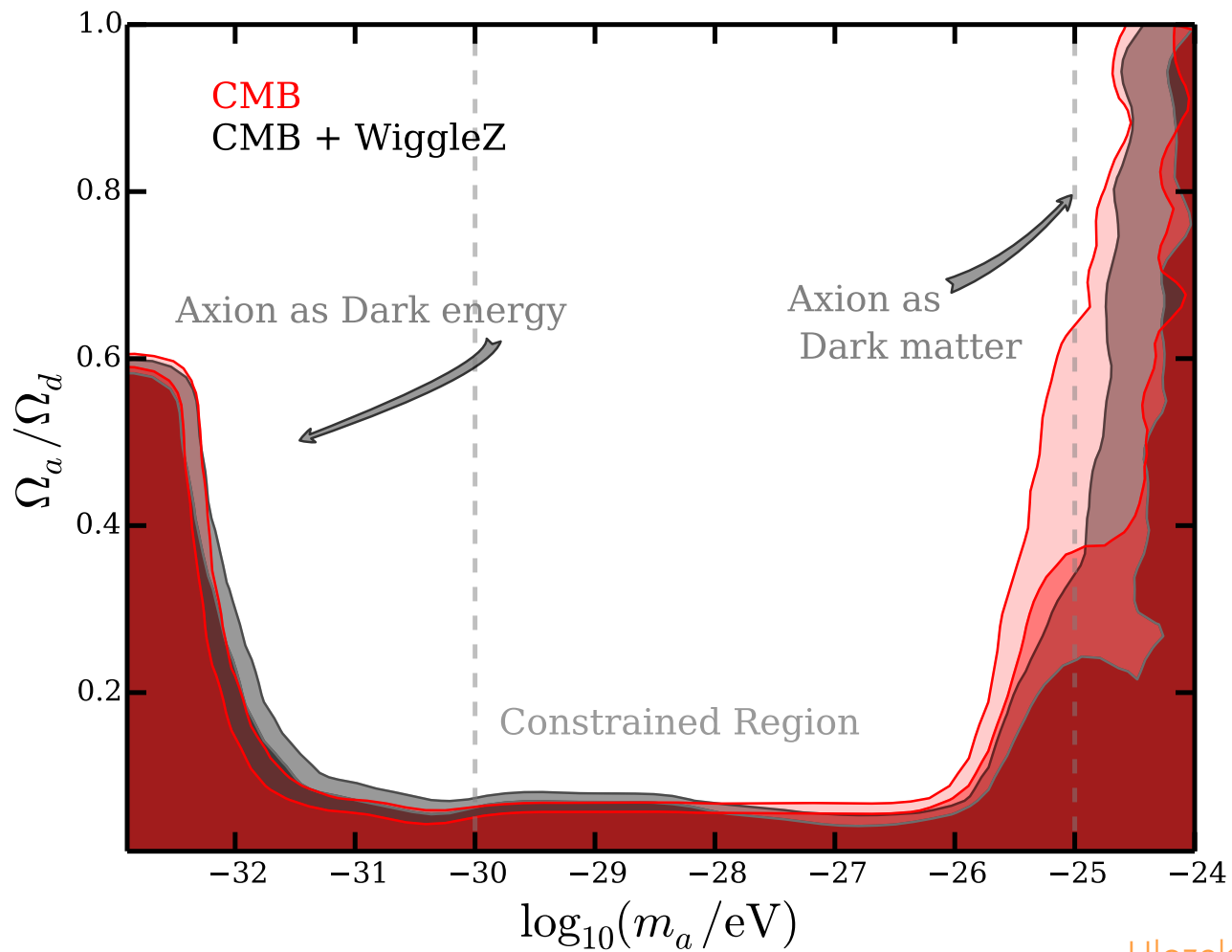
Axiverse 2009: "step in  $P(k)$ ".



Axiverse 2023: EFTofLSS in the BOSS DR12  $P(k)$  multipoles.

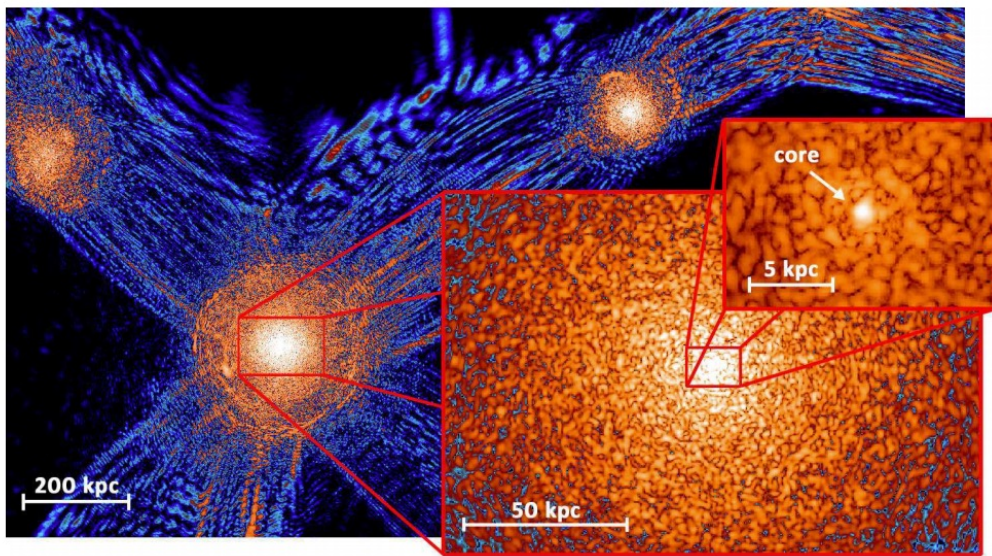
Lague DJEM et al (2021); Rogers, DJEM et al (2023).  
Codes: axionCAMB, multineust, cosmosis, PeakPatch, CLASS-PT



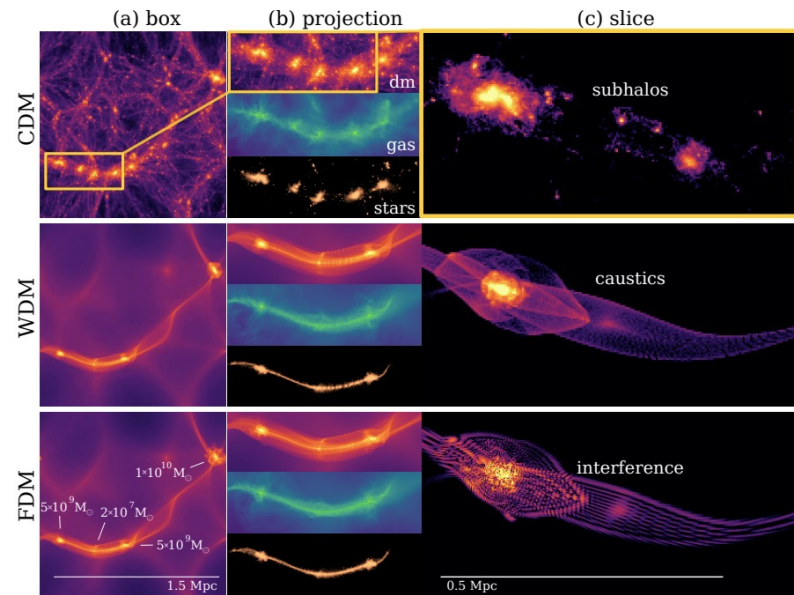


# Simulating Light Axions

Key advance since 2014: the cosmic web with wave effects at  $m \sim 10^{-22}$  eV.



Schive et al (2014)



Mocz et al (2019)

Deeper understanding of dynamics (condensation, relaxation) + new soliton pheno.

# FREEZE-IN AXIONS



Decays

# Cosmological constraints on decaying axion-like particles: a global analysis

Csaba Balázs,<sup>a</sup> Sanjay Bloor,<sup>b</sup> Tomás E. Gonzalo,<sup>c,d</sup>  
Will Handley,<sup>e,f</sup> Sebastian Hoof,<sup>d,g</sup> Felix Kahlhoefer,<sup>c,d</sup>  
Marie Lecroq,<sup>a,h</sup> David J.E. Marsh,<sup>i</sup> Janina J. Renk,<sup>b,j,k</sup>  
Pat Scott<sup>b,k</sup> and Patrick Stöcker<sup>c,l</sup>



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PHYSICAL REVIEW LETTERS **129**, 241101 (2022)

## Irreducible Axion Background

Kevin Langhoff,<sup>1,2</sup> Nadav Joseph Outmezguine<sup>1,2</sup> and Nicholas L. Rodd<sup>3</sup>



SUPERRADIANCE!



Black Hole Super-radiance



# Black Hole Superradiance

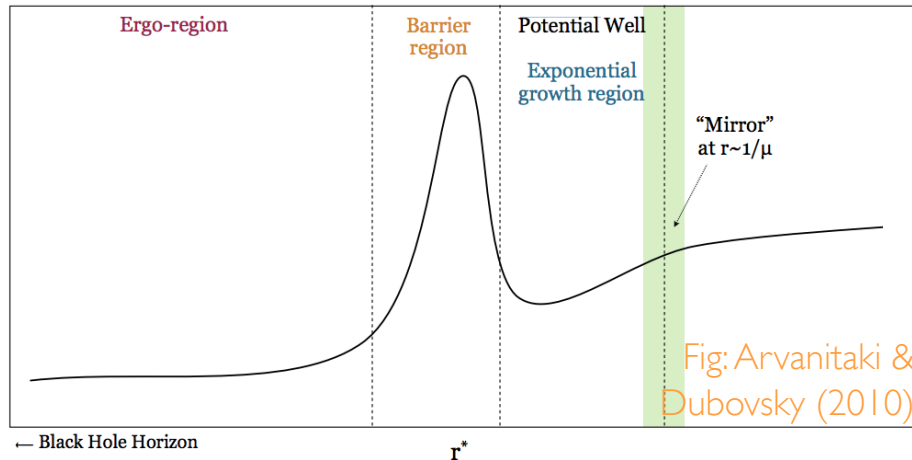
Review: Brito et al (2015)

Solve for instabilities of KG equation on

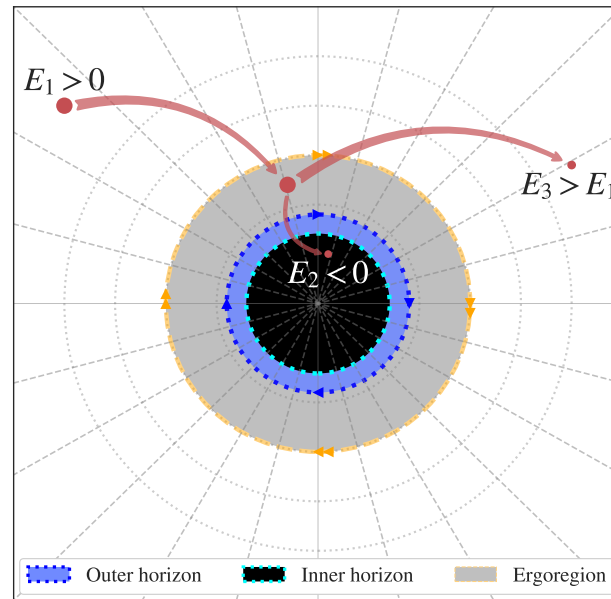
Kerr:  $\square\phi - \partial_\phi V(\phi) = 0$

Non-relativistic limit in "tortoise coords", find instability ( $\omega^2 < 0$ ):

$$\frac{d^2\psi_{lm}}{dr^{*2}} = [\omega^2 - V(r, \omega)] \psi_{lm}.$$



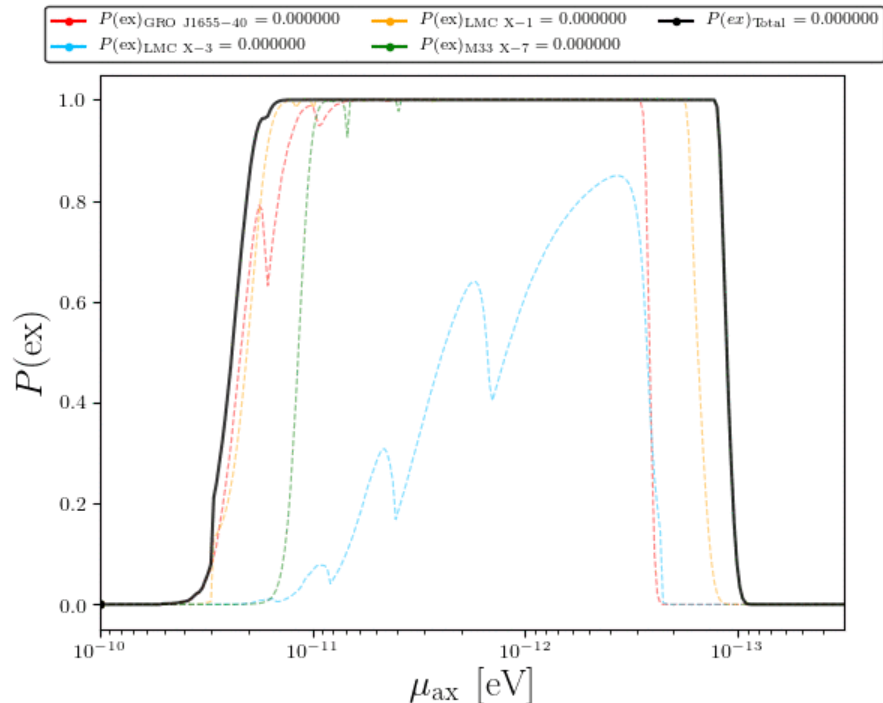
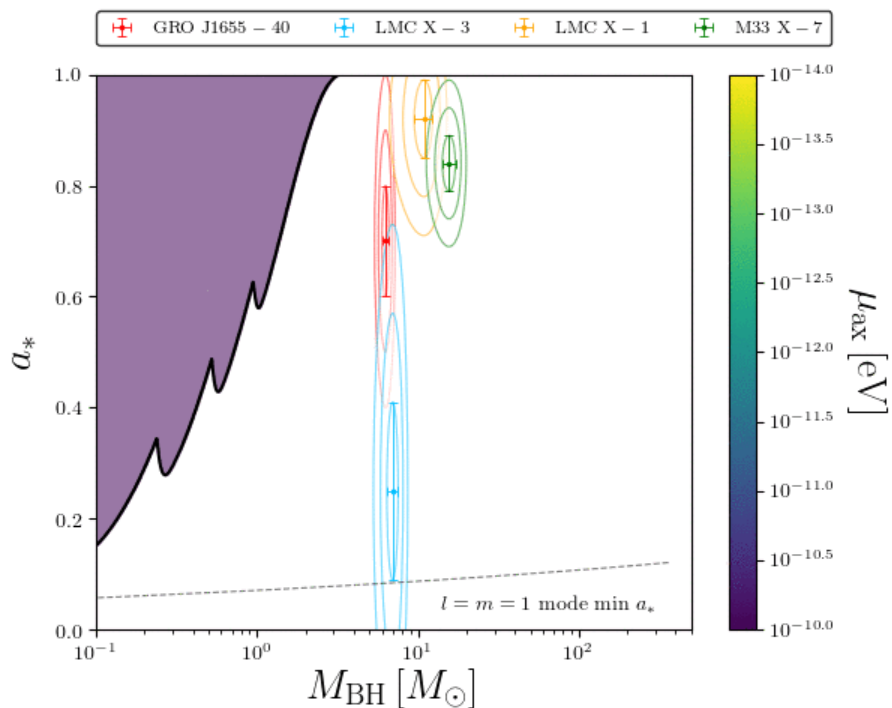
Physical picture: "Penrose process/  
black hole bomb"



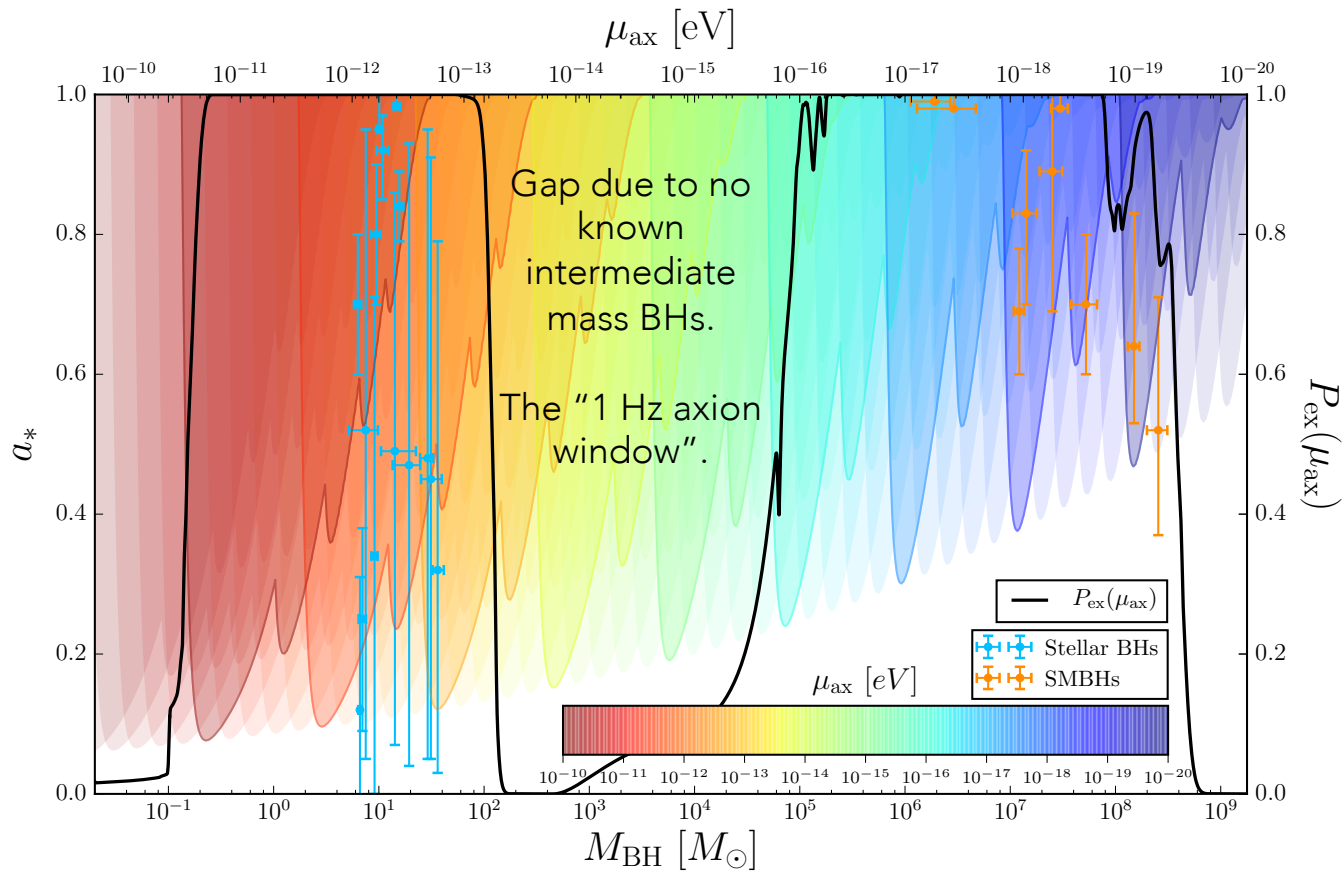
Resonant bosons extract spin from  
astrophysical BHs, if  $\Gamma_{SR} > \Gamma_{others}$

Exclude axion masses where known BHs exist in the superradiant forbidden region.  
 This sample: X-ray stellar BHs. Gaussian composite likelihood.

GIF by Matthew J. Stott



“Exclusion probability” is marginal likelihood. Statistically robust constraints.



Stott & DJEM (2018)

(NB: difference to Baryakhtar+ mass limits due to overly conservative stats model. See backup slides.)

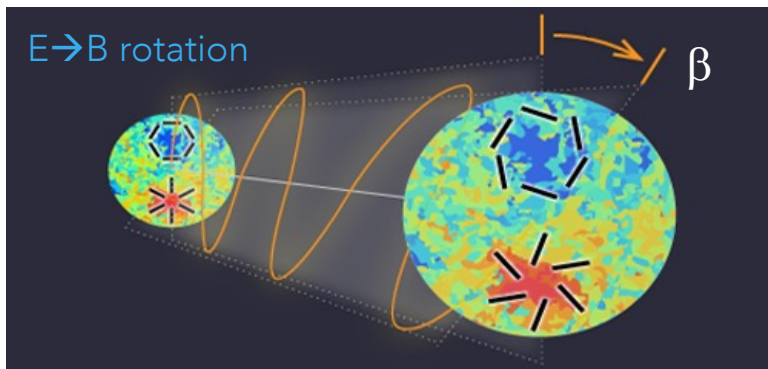
New analysis forthcoming w/Seb Hoof+

# COSMIC BIREFRINGENCE



CMB  
Polarization

# Birefringence



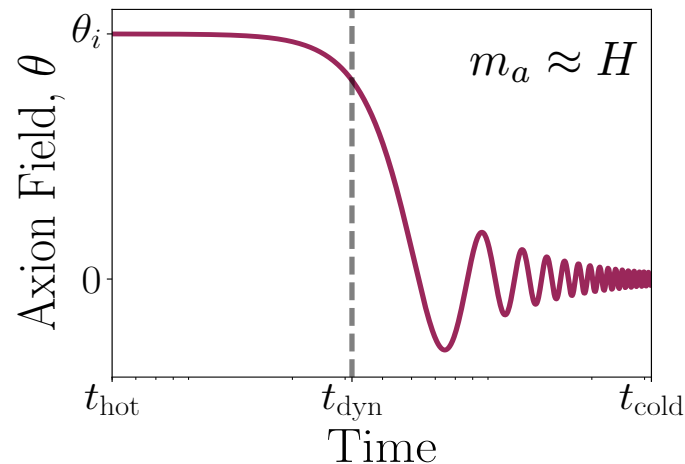
E, B are CMB polarization states (Stokes)

$$\beta = 5.2 \pm 1.9 \times 10^{-3} \quad \begin{array}{l} \text{Minami \& Komatsu (2020)} \\ \text{Planck collab. (2022)} \end{array}$$

Calibrate absolute polarization angle with galactic measurement. Mask dependence consistent with cosmic signal (?).

Isotropic birefringence can be caused by an **ultralight axion** via:

$$\mathcal{L} = g\phi F_{\mu\nu} \tilde{F}^{\mu\nu} \Rightarrow \beta = \int_{\eta_{\text{CMB}}}^{\eta_0} g \frac{d\phi}{d\eta} d\eta$$



**Fixes the axion mass:**

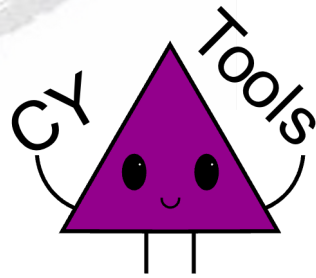
$$10^{-33} \text{ eV} \lesssim m \lesssim 10^{-28} \text{ eV}$$

$H_0$   $H_{\text{CMB}}$

# STRING THEORY PROGRESS

Specifically, type IIB on CY3's.

Demirtas, Rios-Tascon, McAllister





10D SUGRA has **p-form fluxes**.

Consider IIB 4-form,  $C_4$ :

$$S = -\frac{1}{2} \int F_5 \wedge \star F_5, \quad F_5 = dC_4$$

Decompose field into harmonic forms:

$$C_4 = \frac{1}{2\pi} \sum_i a_i(x) \omega_{4,i}(y)$$

Basis of harmonic forms given by closed 4-cycles (divisors) in X:

$$a_i(x) = \int_{D_i} C_4$$

# basis elements given by h<sup>1,1</sup> Hodge number = topological

**Compactify → massless fields in 4D:**

$$S = -\frac{1}{8} \int da_i \mathcal{K}_{ij} \wedge \star da_j,$$

$$\mathcal{K}_{ij} = \frac{\partial^2 \mathcal{K}}{\partial \sigma_i \partial \sigma_j}, \quad \mathcal{K} \propto \ln \mathcal{V}_X$$

$\mathcal{V}$  from "triple intersections" = topological

$$\tau_i = \sigma_i + i a_i \quad \text{SUSY} \rightarrow \tau = \text{Kähler modulus.}$$

**Eigenvalues of K give kinetic term** → "decay constant". Parametrically:

$$\text{Eig}(K) \sim \frac{M_{pl}^2}{(\text{Vol} D_i)^2}$$

Axion **potential generated by ED3 instantons** wrapping D:

$$V = \sum_j \Lambda_j^4 (1 - \cos Q_i^j a^i)$$

$Q$  = instanton charge = topological

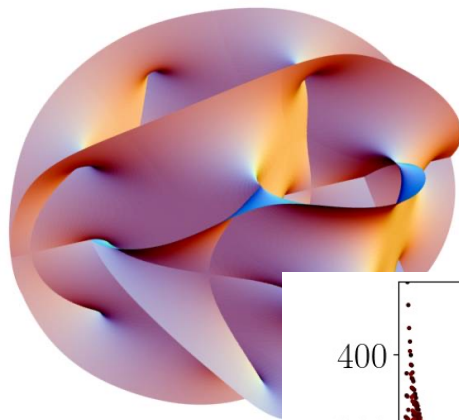
$$\Lambda_j \sim M_{pl}^3 m_{\text{SUSY}} \exp[-\text{Vol} D_i]$$

**→ massive "closed string" axions from gravity sector unavoidable.**

This discussion then suggests the following scenario for the distribution of  $f_a$  and  $m$  for different axions. The values of  $f_a$  are inversely proportional to the area of the corresponding cycle, so they do not change much from one axion to another. Given that the compactification is such that  $S \gtrsim 200$  for string contributions to the QCD axion, and no special fine-tuning is allowed, *all axion decay constants in this scenario are likely to be close to the GUT scale*  $M_{\text{GUT}} \simeq 2 \times 10^{16}$  GeV. *On the other hand, axion masses are exponentially sensitive to the area of the cycles, so that we expect their values to be homogeneously distributed on a log scale.* Given that, as argued

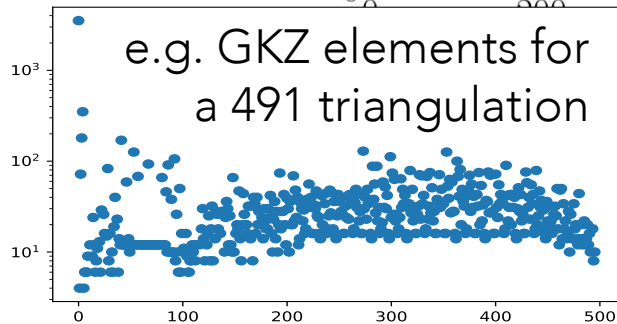
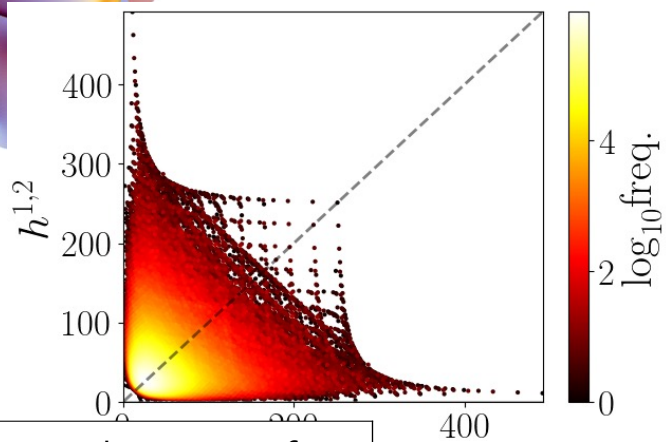
# KS Axiverse

- CY3s constructed as hypersurfaces in "ambient toric varieties". E.g. Fermat quintic in  $\mathbb{CP}^4$ .
- KS database gives all 4d reflexive polytopes  $\sim 4 \times 10^8$ .
- Triangulation of these gives ambient toric varieties. Unique polynomial  $\rightarrow$  CY with  $h_{11}$  Kähler moduli.
- Automated fun with CY-Tools!
- Axions:  $Q$  unique for polytope.  $K_{ij}$  fixed by CY. Saxion,  $\sigma$ , must be fixed in "stretched Kähler cone" where all curve and divisor volumes  $> 1$ .



$h_{11}=27$ : most polytopes  
 $h_{11}=491$ : most CYs  $< 10^{428}$

Year	$h^{1,1}$	CPU time
2014	25	a few hours
2017	491	2s
2019	491	20ms



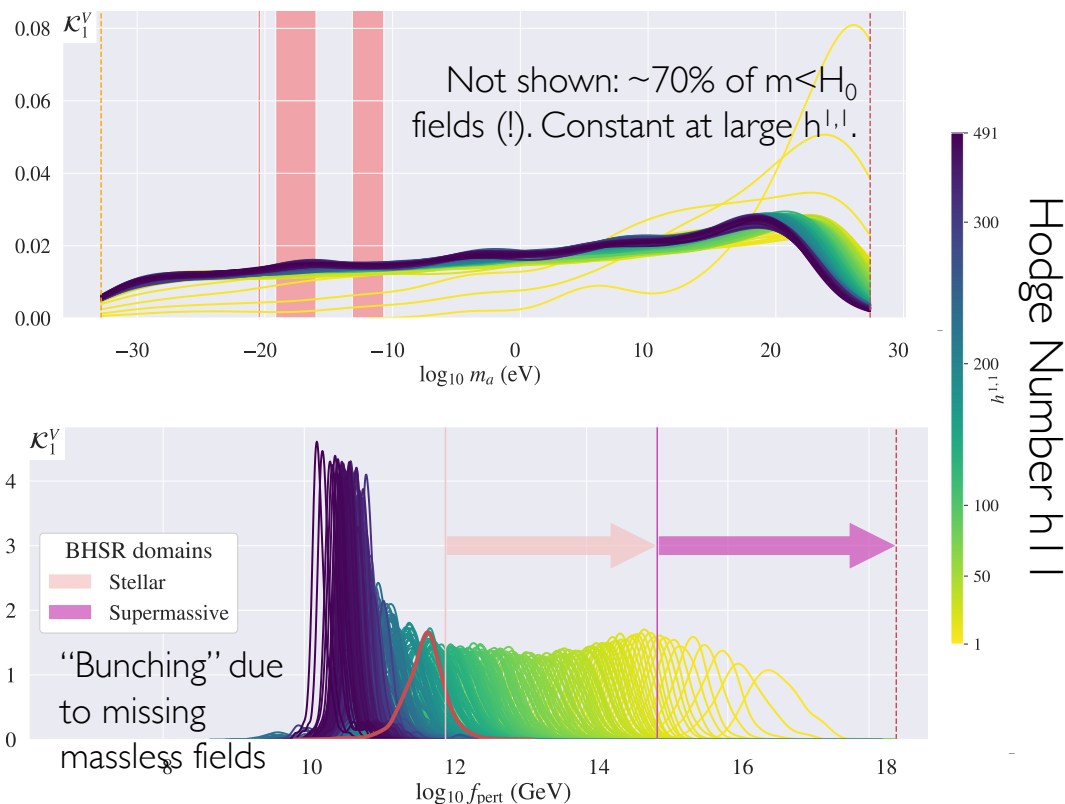
# Axion Spectra from KS

Mehta, DJEM et al (2021)

Find vacua of  $V$  in fundamental domain. Expand to quartic order  $\rightarrow$  masses + quartics ("fpert").

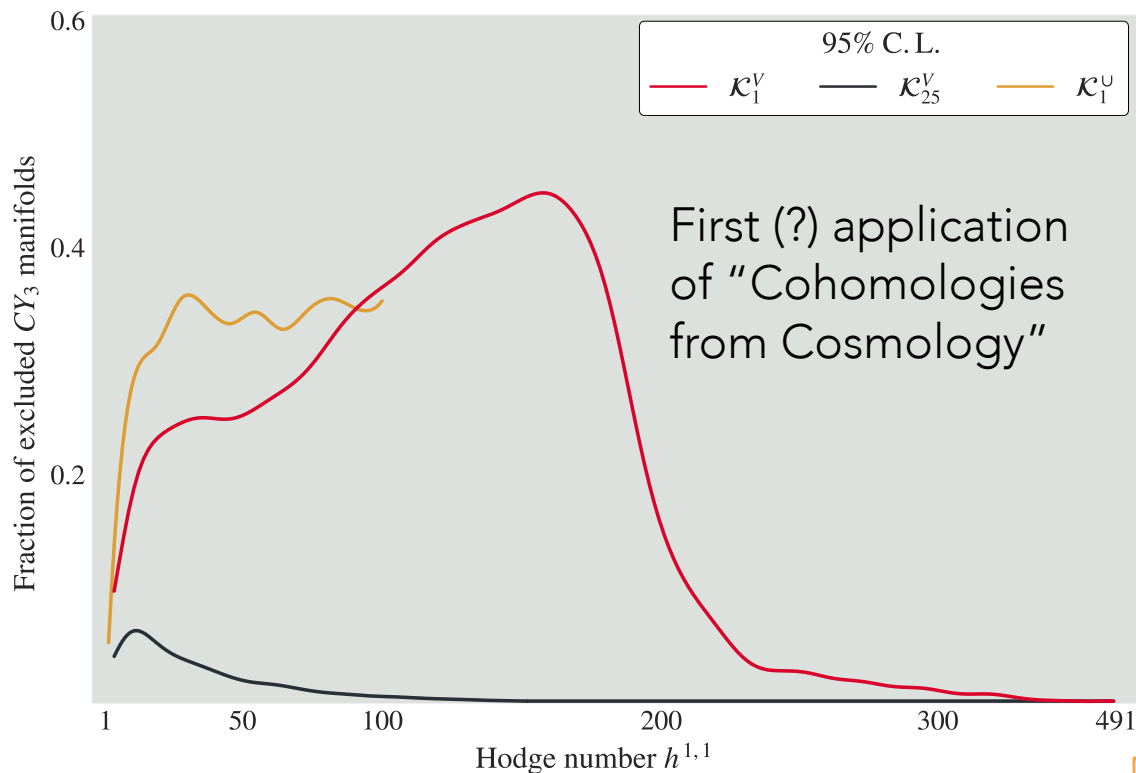
Trends: Kähler cones become very narrow at large  $h_{11} \rightarrow$  cycles in the CY have large volumes  $\rightarrow$  (ultra)light axions and smaller decay constants.

Mass spectrum "blue tilted".  
Decay constants log-normal, becoming smaller at large  $h_{11}$ .



# Constraints on IIB CY Vacua

Ensemble of  $O(10^5)$  CYs. All up to  $h_{1,1}=5$ . 100 per  $h_{1,1}$  up to 176. Few per  $h_{1,1}$  to 491.



(NB: this limit does NOT include LIGO BHs, which require different mass and spin model)

Above  $h_{1,1} \sim O(\text{few})$  limit driven by stellar BHs with well measured spin.

Trend easily understood from falling K eigs at large volume  $\rightarrow$  Bosenova shut-off for stellar BH limits

# GLIMMERS FROM THE AXIVERSE

Gendler, Marsh, McAllister, Moritz arXiv:2309.13145





# Axion-Photon Couplings

## Recipe:

1. Choose a CY. We take  $2 \times 10^5$  over  $h_{11}=50, 100, 200, 491$ .
2. Choose a divisor to host QCD.
3. Dilate divisor to  $\text{Vol}=40 \rightarrow$  QCD gauge coupling in UV.
4. Pick a divisor for QED. Same  $\rightarrow$  GUT. Intersecting  $\rightarrow$  non-GUT.
5. QED divisor  $\rightarrow$  linear combination of axions coupling to EM.
6. Diagonalise matrices and compute couplings of eigenstates.

## Consequences:

“Kinetic isolation”: K matrices are sparse  $\rightarrow$  suppressed kinetic mixing of all axions into EM linear comb.

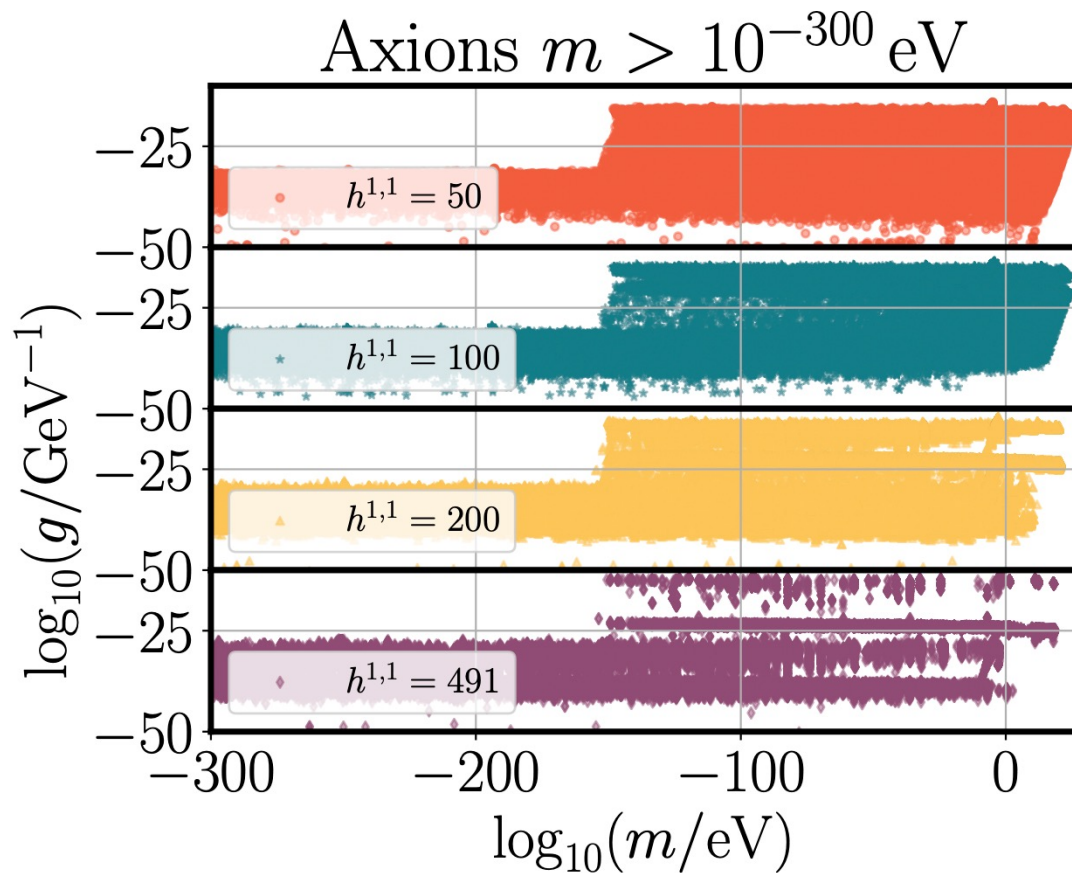
“Light threshold”: mass scale generated by ED3 branes on divisor  $\rightarrow$  hierarchically suppressed coupling of axions lighter than this.

c.f. similar effect in GUTs found by Agrawal et al.

Stringy effects! Need CY to get K matrix. No “U(1) instantons” in field theory.

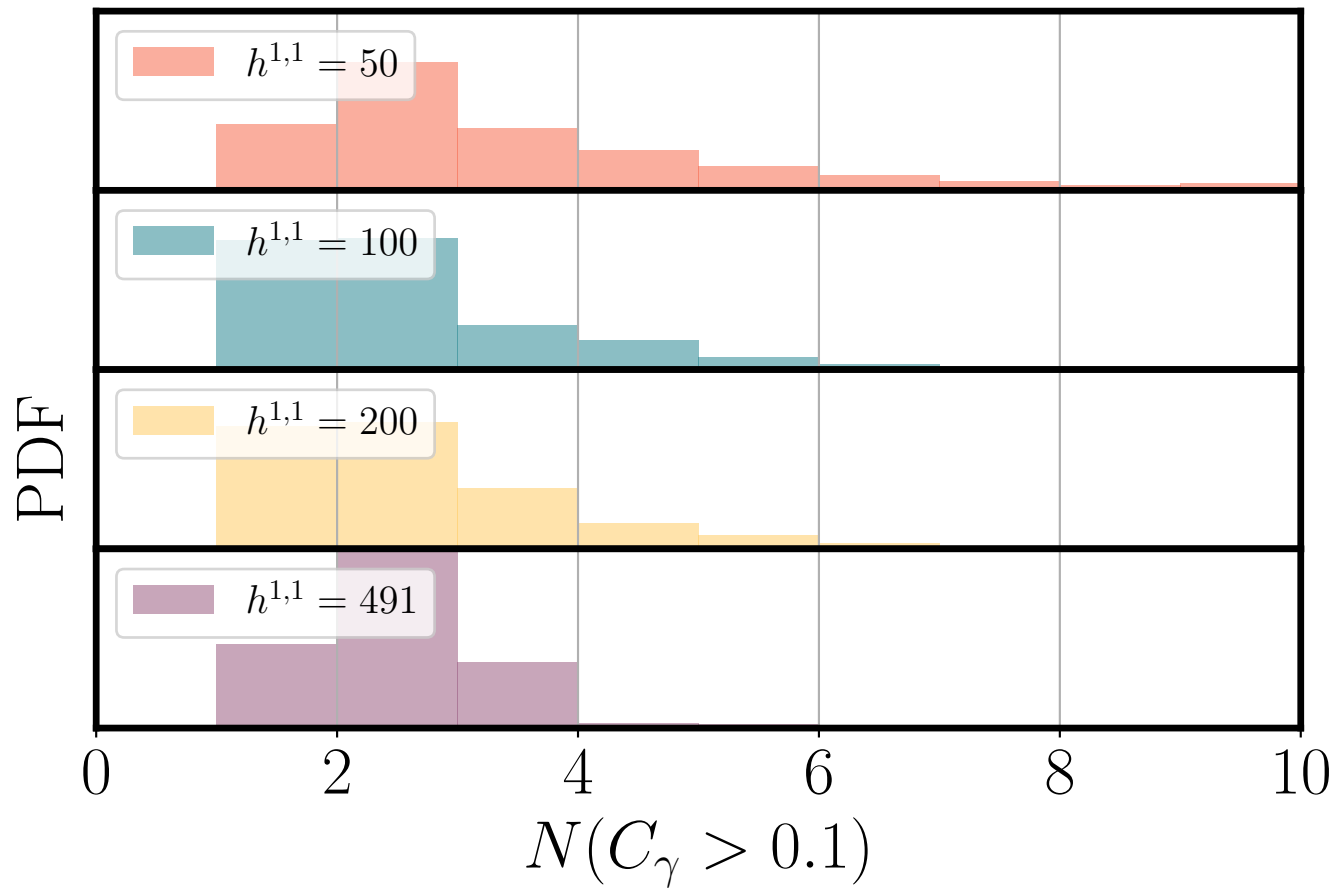
Restricting QED gauge coupling running  
decouples  $m < 10^{-150}$  eV  $\rightarrow$  not all axions couple.

$$\alpha(\Lambda_{\text{UV}}) > 1/128$$

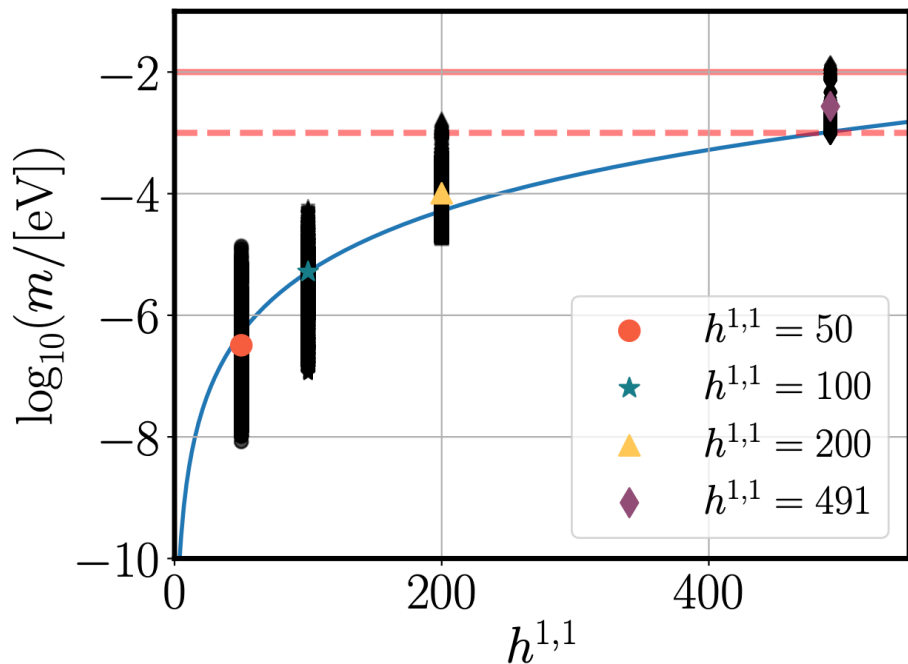


Number of axions with  $O(1)$  coupling to photons is small.

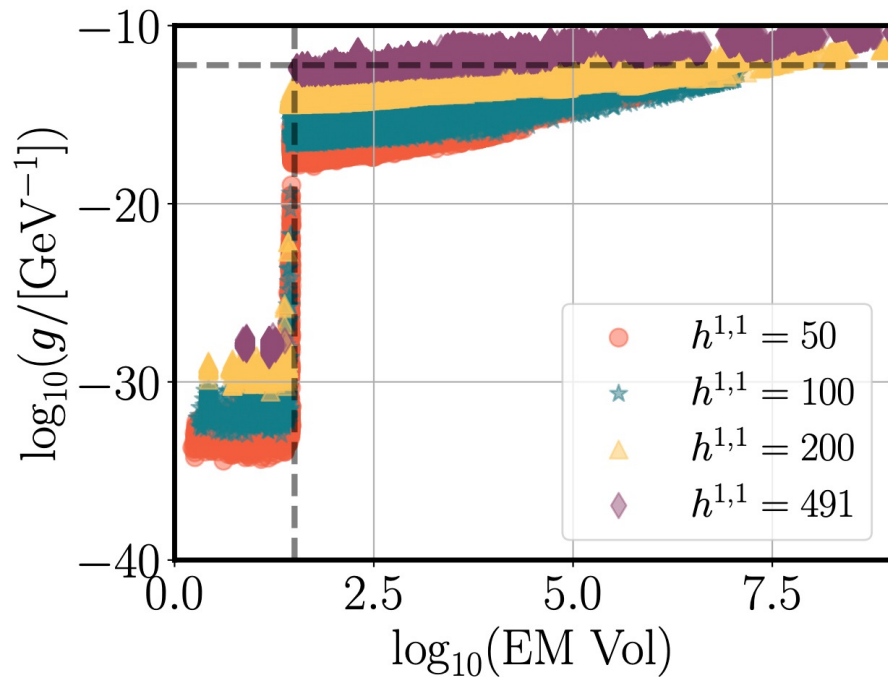
$$g = C_\gamma \frac{\alpha}{2\pi f_a}$$



QCD axion mass strongly correlated with the Hodge number  $\rightarrow$  discover QCD axion and measure topology of CY (!).

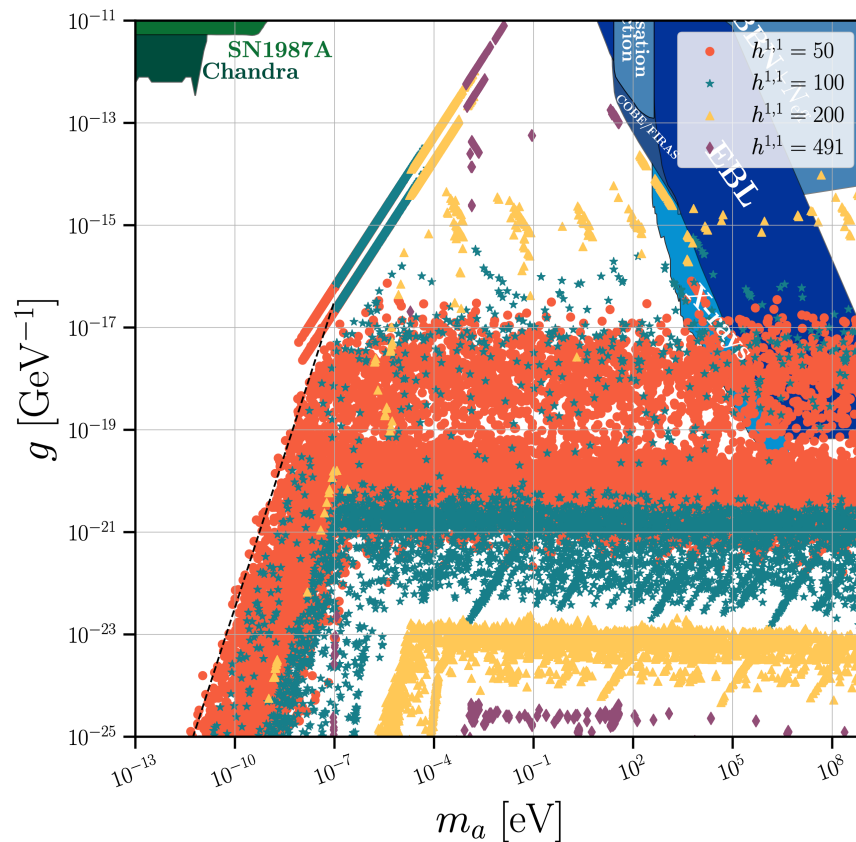


Light threshold at  $\text{Vol} \sim 30 \rightarrow$  X-ray spectra probe non-GUTs. May exclude some models with large  $h^{1,1}$



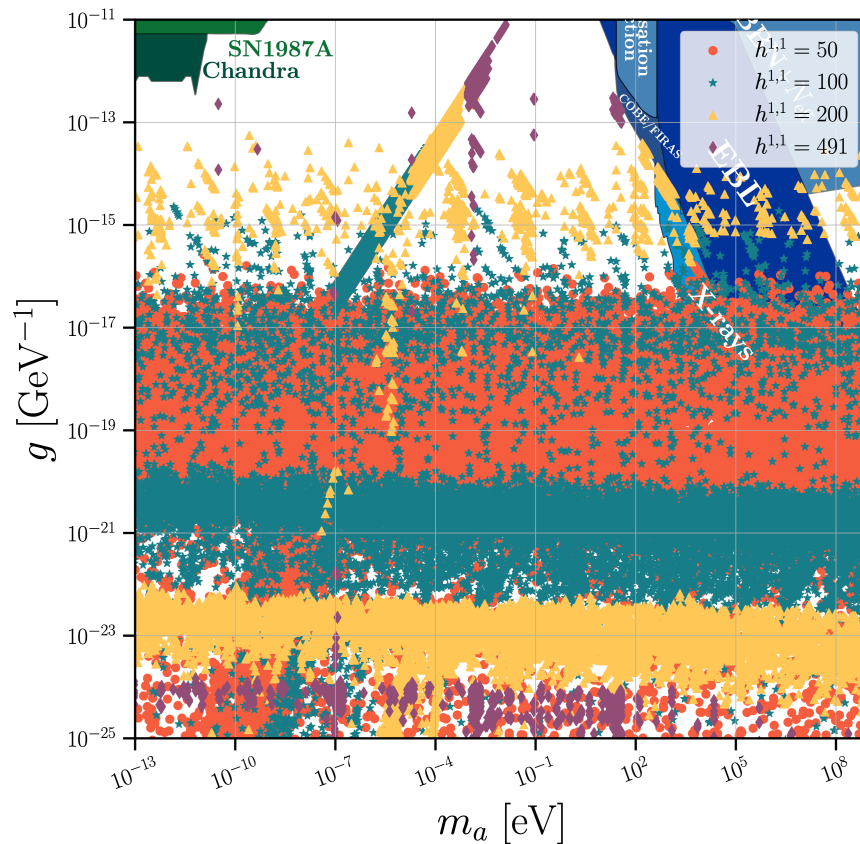
# GUTs vs non-GUTs

- GUT: fix cycle volume for group containing EM and QCD at  $1/\alpha_{\text{GUT}}$
- The QCD induced mass is bigger than the ED3 one.
  - QCD axion provides the light threshold.
  - No axions above the QCD line (as Agrawal et al in field theory).



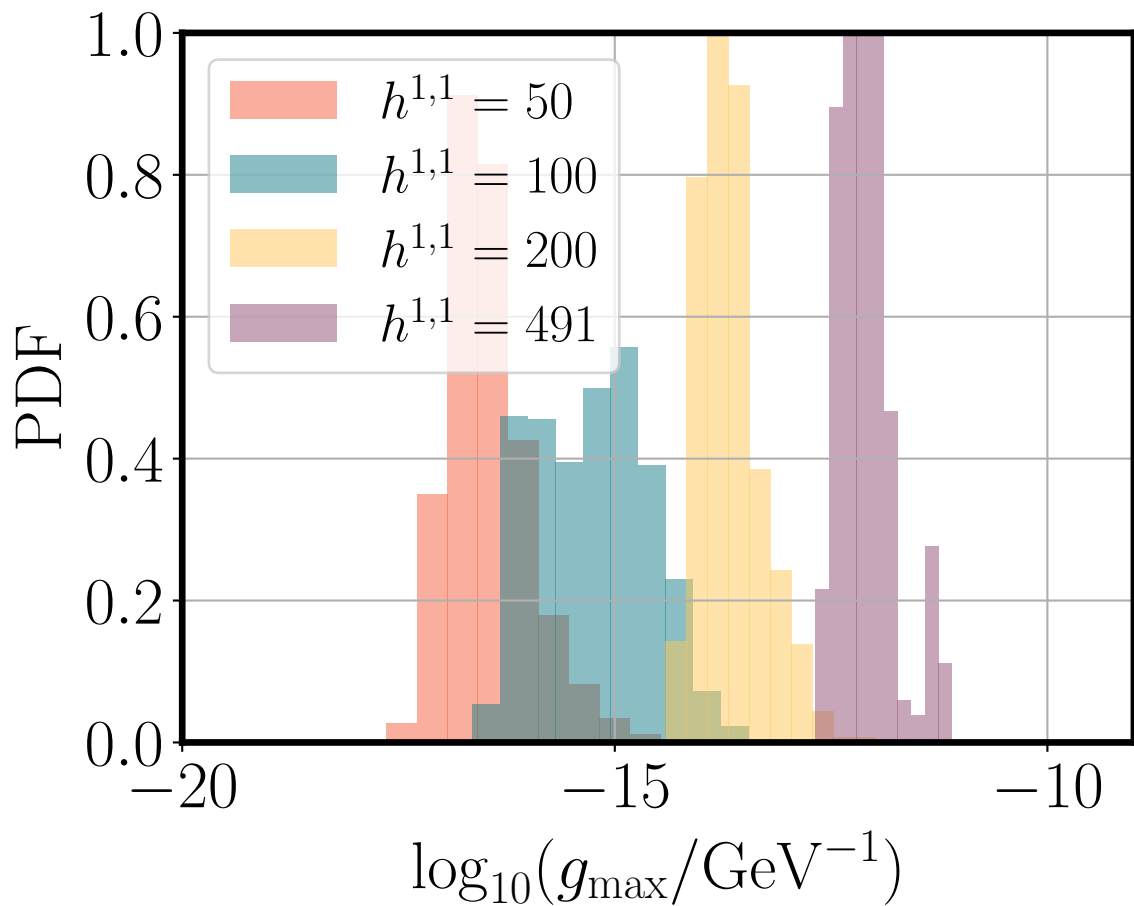
# GUTs vs non-GUTs

Non-GUTs: (almost) anything goes.

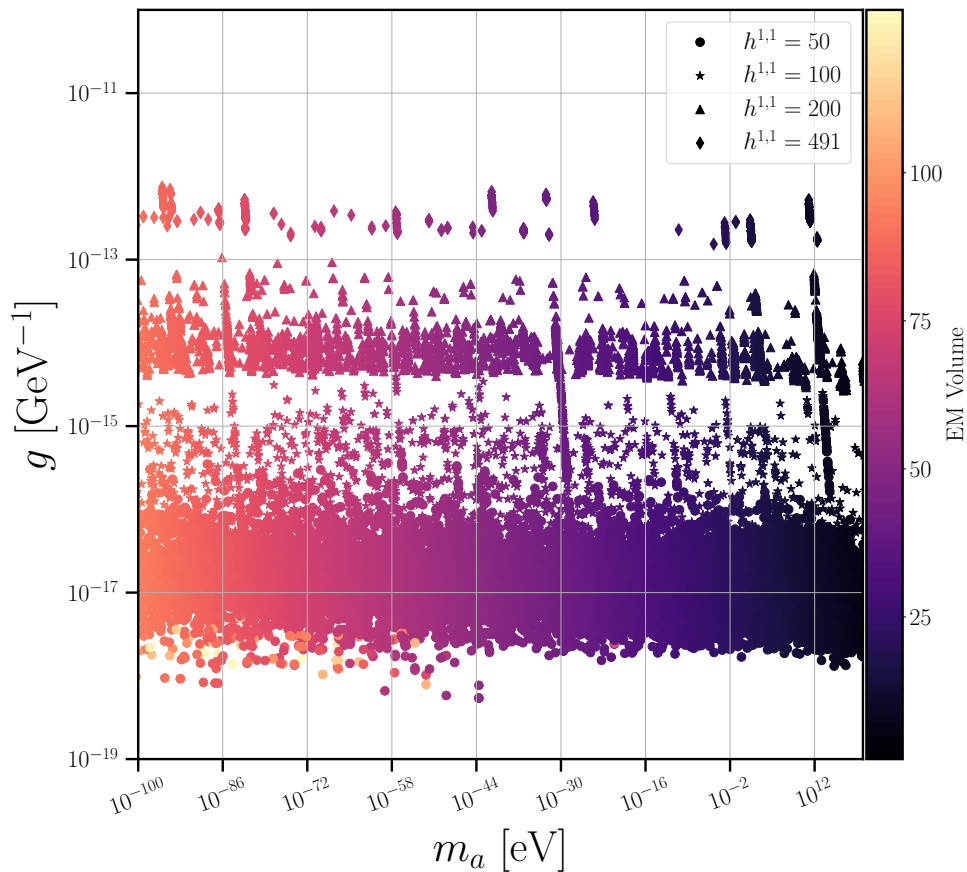




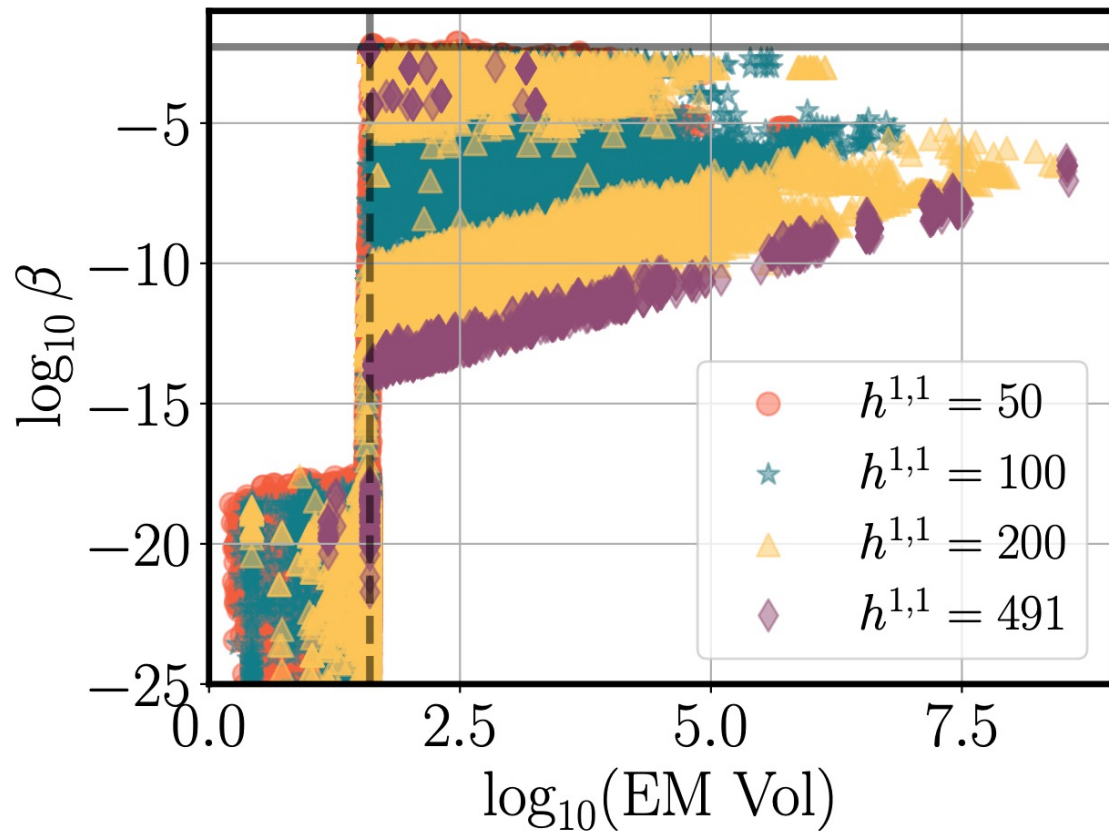
The maximum value of  $g$  increases with  $h^{1,1}$  due to increasing cycle volumes.



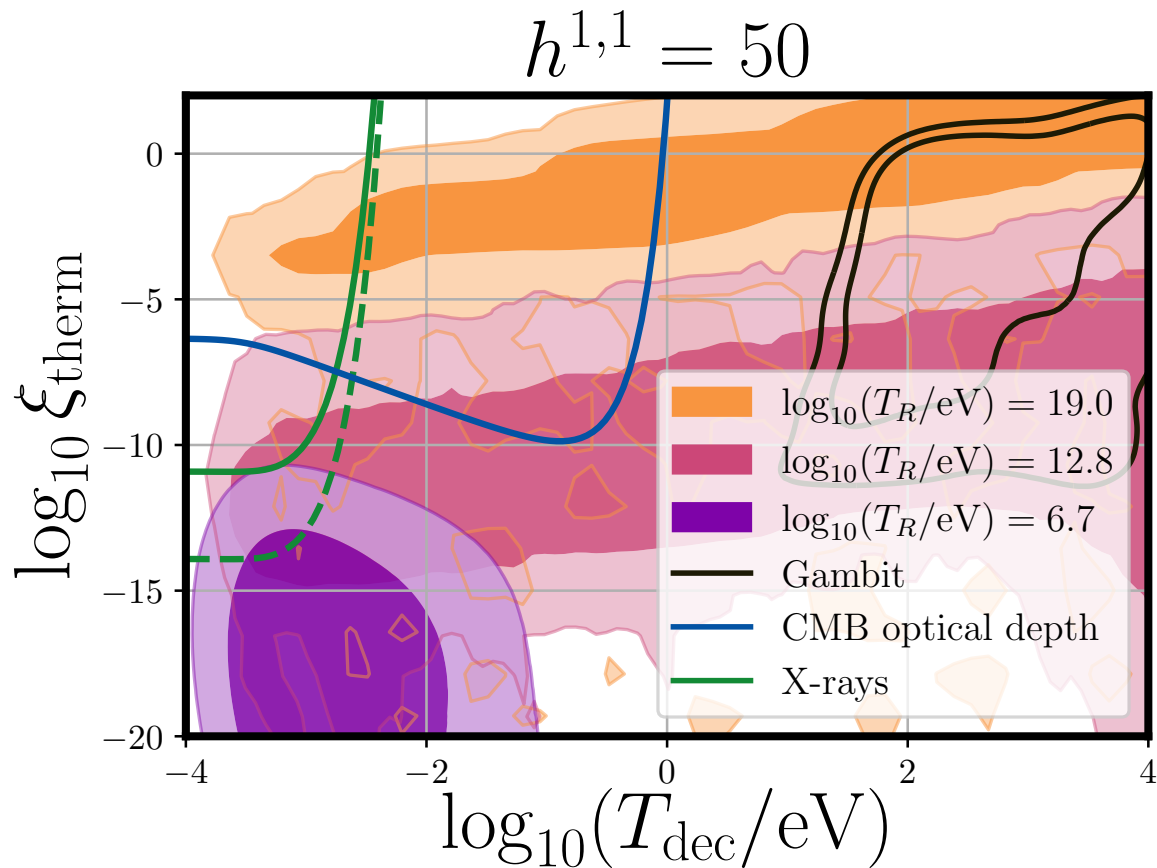
Which axions couple  $O(1)$  to EM depends on the cycle volume  $\rightarrow$  UV info.

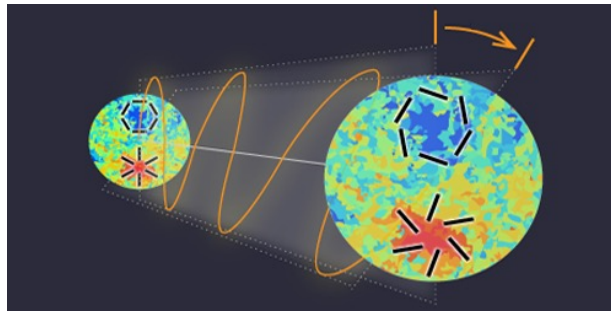


Birefringence possible but non-generic. Only get the right value if  $\text{Vol} \sim 40$  ( $H_{\text{CMB}}$  light threshold). Interesting coincidence?



keV – GeV axions Primakoff freeze-in then decay. Disfavours reheating  $T > 10^{10}$  GeV  
→ upper limit on SUSY?





- Birefringence has a hint from Planck: have we seen evidence of the axiverse already?
- Cosmological probes have matured in precision, and in the next decades will test GUT scale axions.
- Superradiance has been used to test the axiverse up to  $h^{1,1} \sim 200$  in explicit constructions on CYs.
- Advances in constructing the visible sector in type IIB offer promise to discover the axiverse.

