

Hagedorn temperature in string and gauge theories

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March 25, 2024

Based on:

- F. Bigazzi, **T. Canneti**, A. L. Cotrone, [arXiv:2210.09893 [hep-th]]
- F. Bigazzi, **T. Canneti**, W. Mück, [arXiv:2306.00588 [hep-th]]
- F. Bigazzi, **T. Canneti**, A. L. Cotrone, [arXiv:2306.17126 [hep-th]]
- F. Bigazzi, **T. Canneti**, F. Castellani, A. L. Cotrone, W. Mück, [arXiv:24xx.xxxxx [hep-th]]

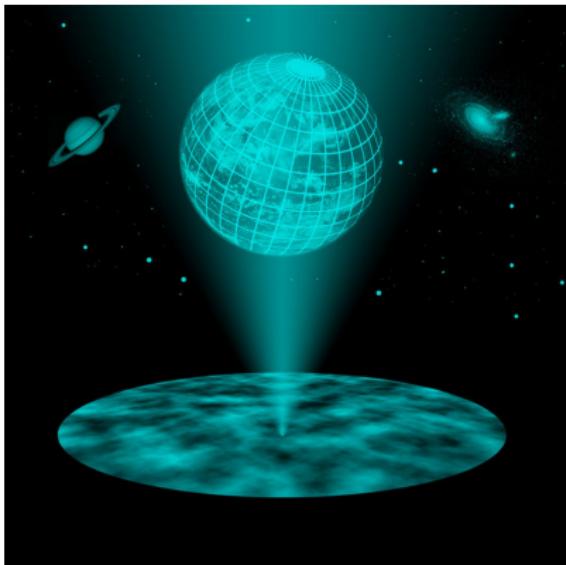
Holographic Correspondence

[Maldacena '97, Witten '98, Gubser et al. '98]

theory of gravity
in $d + 1$ -dimension
(string theory)

exact
 \equiv

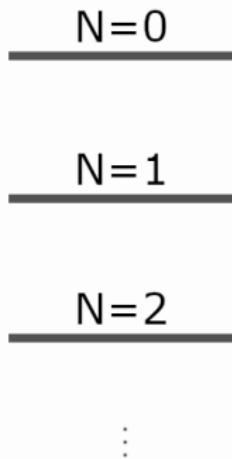
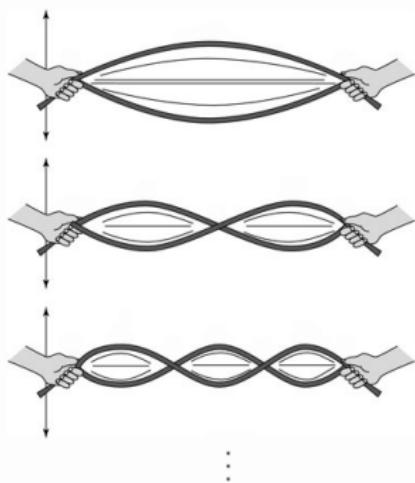
gauge theory
in d -dimension



- Keywords: holographic, weak/strong, conjecture, **duality**
- String theory and gauge theory share the very same properties

String theory

- Building block of nature: one-dimensional **string** (length scale l_s)
- Emergence of elementary particles/fundamental interactions
- String spectrum:

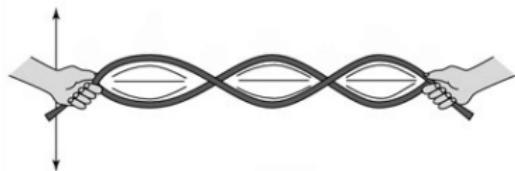


$$l_s^2 M^2 \approx N$$

(mass-shell condition)

↓
states belonging to the
same level N have the
same mass M

Asymptotic density of states $d(N)$ [Huang et al. '70, Sundborg '85]



$$\hat{H} \approx \omega_0 \hat{N}, \quad \hat{N} = \sum_{k=1}^{+\infty} k \hat{a}_k^\dagger \hat{a}_k$$

$$|\psi\rangle = (\hat{a}_1^\dagger)^{n_1} (\hat{a}_2^\dagger)^{n_2} \dots (\hat{a}_k^\dagger)^{n_k} |0\rangle : \quad \hat{N}|\psi\rangle = N|\psi\rangle, \quad N = \sum_{k=1}^{+\infty} k n_k$$

N	$ \psi\rangle$	$d(N)$
1	\hat{a}_1^\dagger	1
2	$\hat{a}_2^\dagger \quad (\hat{a}_1^\dagger)^2$	2
3	$\hat{a}_3^\dagger \quad \hat{a}_2^\dagger \hat{a}_1^\dagger \quad (\hat{a}_1^\dagger)^3$	3
4	$\hat{a}_4^\dagger \quad \hat{a}_3^\dagger \hat{a}_1^\dagger \quad (\hat{a}_2^\dagger)^2 \quad \hat{a}_2^\dagger (\hat{a}_1^\dagger)^2 \quad (\hat{a}_1^\dagger)^4$	5
\vdots	\vdots	\vdots

$$d(N) \stackrel{N \gg 1}{\approx} e^{\#\sqrt{N}}$$

[Hardy et al. '18]

The Hagedorn temperature [Hagedorn '68]

Partition function:

$$Z \sim \int dE \rho(E) e^{-\beta E}$$

- $\beta = 1/T$ (inverse temperature)
- Mass-shell condition $\Rightarrow E^2 \sim N$ (mass $M \sim$ energy E)
- $\rho(E)dE = d(N)dN \Rightarrow \rho(E) \sim e^{\beta_H E}$

Definition

The Hagedorn temperature $T_H = 1/\beta_H$ is defined as temperature above which Z does not exist (**divergent!**)

Methods

- String in flat space: Z is well-known (end of the story)
- String in curved space: we cannot compute Z in general

The **lightest state** of the spectrum of a string winding once the thermal (compact) direction is **tachyonic above T_H**

$$M^2(|0\rangle; \beta \leq \beta_H) \leq 0$$

[Sathiapalan '87, Kogan '87, O'Brien et al. '87]

Idea

[Atick&Witten '88]: flat space computation relying just on the mass-shell condition
⇒ we extend this method **to curved space** for a large class of models

Example [Urbach '22, Ekhammar et al. '23, Bigazzi et al. '23]

String sitting at the center of $d + 1$ -dim Anti de Sitter (AdS_{d+1}) space

$$ds^2 = R_{AdS}^2 (\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{d-1})$$

Deduce T_H by fixing $M^2 = 0$:

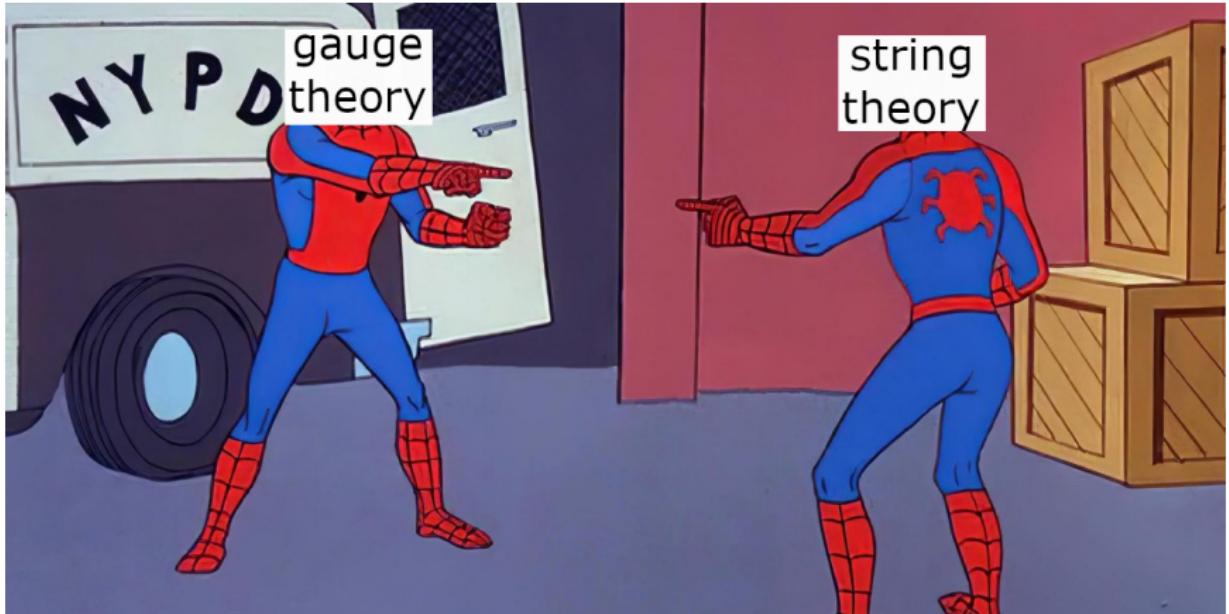
$$T_H = \sqrt{\frac{g}{2\pi}} + \frac{d}{8\pi} + \frac{d^2 + d - 8d \log 2}{32\sqrt{2}\pi^{3/2}\sqrt{g}} + \frac{4d^3 + 7d^2 - 2d}{1024\pi^2 g} + \mathcal{O}(g^{-3/2}),$$

where $g = 1/4\pi l_s^2$.

Here, $R_{AdS} = 1$: to get physical result $T_H \mapsto T_H R_{AdS}$ and $l_s \mapsto l_s/R_{AdS}$

Holography

$$T_H^{\text{gauge theory}} \equiv T_H^{\text{string theory}}$$



Comparison with numerical results

$d = 4$: (Super) Yang-Mills theory on $S^1 \times S^3$

$$\begin{aligned} T_H &= \sqrt{\frac{g}{2\pi}} + \frac{1}{2\pi} + \frac{5 - 8 \log 2}{8\pi\sqrt{2\pi}\sqrt{g}} + \frac{45}{128\pi^2 g} + \mathcal{O}(g^{-3/2}) \\ &\approx 0.39894\dots \sqrt{g} + 0.15916\dots - \frac{0.00865\dots}{\sqrt{g}} + \frac{0.0356\dots}{g} + \mathcal{O}(g^{-3/2}) \end{aligned}$$

vs

$$\begin{aligned} T_H^{[\text{Ekhammar et al.'23}]} &\approx (0.39894 \pm 0.00001) \sqrt{g} + (0.15916 \pm 0.00001) \\ &\quad - \frac{(0.00865 \pm 0.00001)}{\sqrt{g}} + \frac{(0.0356 \pm 0.0001)}{g} + \dots \end{aligned}$$

That's an outstanding and non-trivial **test of holography!**

Conclusions

- T_H for (confining) gauge theories up to next-to-next-to-leading order (NNLO) or NNNLO in the string length
- Agreement with expectation in the literature and numerical analysis
- Some ambiguities fixed in an analytical way from the string

Thanks a lot for your attention!