# Holographic thermal entropy from geodesic bit threads

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Based on: S. C, F. Gentile, D. Seminara and E. Tonni [arXiv:2403.03930]

# Entanglement Entropy

Divide the total system in two subsystems A and B

 $S_A = -tr_A \rho_A \log \rho_A$ 

 $\rho_A = tr_B |\Psi\rangle\langle\Psi| = \text{reduced density matrix of } A \rightarrow \text{trace out } B \text{ from the total} \\ \text{density matrix } \rho_{\text{tot}} = |\Psi\rangle\langle\Psi|$ 



- If the total system is in a pure state, *S*<sub>A</sub> is a good measure of the entanglement
- If the total system is in a mixed state, *S*<sub>A</sub> is counting not only the entanglement, but also the initial entropy of the system

# Holographic Entanglement Entropy

Ryu and Takayanagi (2006)

$$S_A = \frac{\text{Area of } m(A)}{4G_N}$$

m(A) = d – dimensional static minimal surface in AdS<sub>d+2</sub> ("RT surface") anchored on A



# **Bit Threads**

Freedman and Headrick (2017):

- How one should think about the minimal surface m(A)?
- Can we rewrite the HEE formula without involving m(A)?

They introduced a vector field V in the bulk:

 $abla_{\mu}V^{\mu} = 0$ ,  $|V| \le C$ , Flow lines = Bit Threads

- Flux of V through a surface  $\gamma$ :  $\phi(\gamma) = \int_{\gamma} \sqrt{h} n_{\mu} V^{\mu}$
- They demonstrated:

$$\max_{V} \phi(A) = C \min_{m \sim A} \operatorname{Area}(m)$$

There exists a vector field Vwhose flux through A equals the area of m(A)



## Reformulation of the Ryu-Takayanagi formula

If we set  $C = \frac{1}{4G_N}$ 

$$S(A) = \max_{V} \phi(A)$$

 $\rightarrow$  Each flow line (bit thread) leaving the region A carries one bit of information about the microstate of A

 $\rightarrow$  The minimal surface acts as a bottleneck limitating the number of threads emanating from A

 $\rightarrow$  The minimal surface is unique, but there are infinite classes of bit threads that maximizes the flux

## Sphere in S-AdS $_{d+2}$ black hole

• Constant time slice of Schwarzschild AdS<sub>*d*+2</sub> black hole:

$$ds_{|\mathsf{t}=\mathsf{const}}^2 = \frac{1}{z^2} \left( dr^2 + \frac{dz^2}{f(z)} + r^2 d\Omega_{d-1} \right) \quad \text{with } f(z) = 1 - \left( \frac{z}{z_h} \right)^{d+1}$$

• Minimal surface  $\rightarrow$  solution of a second order differential equation

$$\frac{z''(r)}{f(z)} + \frac{(d-1) \ z'(r) \left(\frac{z'(r)^2}{f(z)} + 1\right)}{rf(z)} + \frac{d\left(\frac{z'(r)^2}{f(z)} + 1\right)}{z(r)} - \frac{z'(r)^2 \ \frac{\partial f(z)}{\partial z(r)}}{2f(z)^2} = 0$$

No analytic solution known  $\rightarrow$  we find a <u>numerical solution  $z_m(r_m)$ </u>

• Flow lines of the vector field V = geodesics

$$z'(r) = \pm \frac{1}{z} \sqrt{f(z) (C^2 - z^2)}$$

The integration constant *C* is determined by imposing that geodesics intersect orthogonally the RT surface at a point  $z_m(r_m)$ 

The two branches (ascending and descending) of each geodesic can be expressed as

$$r_{\gtrless}(z) = r_m(z_m) + \int_{z_m}^{\bar{z}_*} \frac{v}{\sqrt{f(v)}(C^2 - v^2)} dv \pm \int_{z}^{\bar{z}_*} \frac{v}{\sqrt{f(v)}(C^2 - v^2)} dv$$

#### Three types of geodesics:

- 1) Green: two endpoints on the boundary (maximum height  $\tilde{z}_* = C$ )
- 3) Magenta: one endpoint on the boundary and its reaches the horizon only for (limiting case of the other two)



Magnitude of the vector field V:

$$\left| \boldsymbol{V}_{\geq} \right| = \left( \frac{z}{z_m} \right)^d \frac{\sqrt{C^2 - z_m^2}}{\sqrt{C^2 - z^2}} \left( \frac{r_m}{r_{\geq}} \right)^{d-1} \left. \frac{\left( \partial_{z_m} r_{<} \right) \right|_{z=z_m}}{\partial_{z_m} r_{\geq}}$$

The entanglement entropy  $S_A$  is computed as the flux of bit threads through the region A:

$$S_{A} = \phi(A) = \frac{L_{\mathsf{AdS}}^{d} \Omega_{d-1}}{4G_{N}} \int_{0}^{b-\epsilon} (n_{z} V^{z} \frac{1}{z^{d}})_{|_{z=0}} r^{d-1} dr = \frac{L_{\mathsf{AdS}}^{d} \Omega_{d-1}}{4G_{N}} \int_{0}^{b-\epsilon} C_{A}(r) r^{d-1} dr$$





Does the flux through the green interval  $A_{\beta}$  have a meaning?  $\rightarrow$  **Yes!** 

If we rename as  $\tilde{S}_{A,th}$  the flux of bit threads through the green interval

$$\tilde{S}_{A,th} = \frac{L^{d}_{AdS}\Omega_{d-1}}{4G_N} \int_0^{b_\beta} (n_z V^z \frac{1}{z})_{|_{z=0}} r^{d-1} dr$$

We find, at least numerically, that the geodesic bit threads identify a subinterval  $A_{\beta}$  that contains the information regarding the thermal entropy

$$S_{A,th} = \frac{L_{AdS}^d}{4G_N} V_d \left(\frac{b}{z_h}\right)^d \text{ where } V_d = \frac{\pi^{d/2}}{\Gamma(d/2+1)}$$





# Thank you for the attention!