The arctic curves of the four-vertex model Journal of Physics A (2023)

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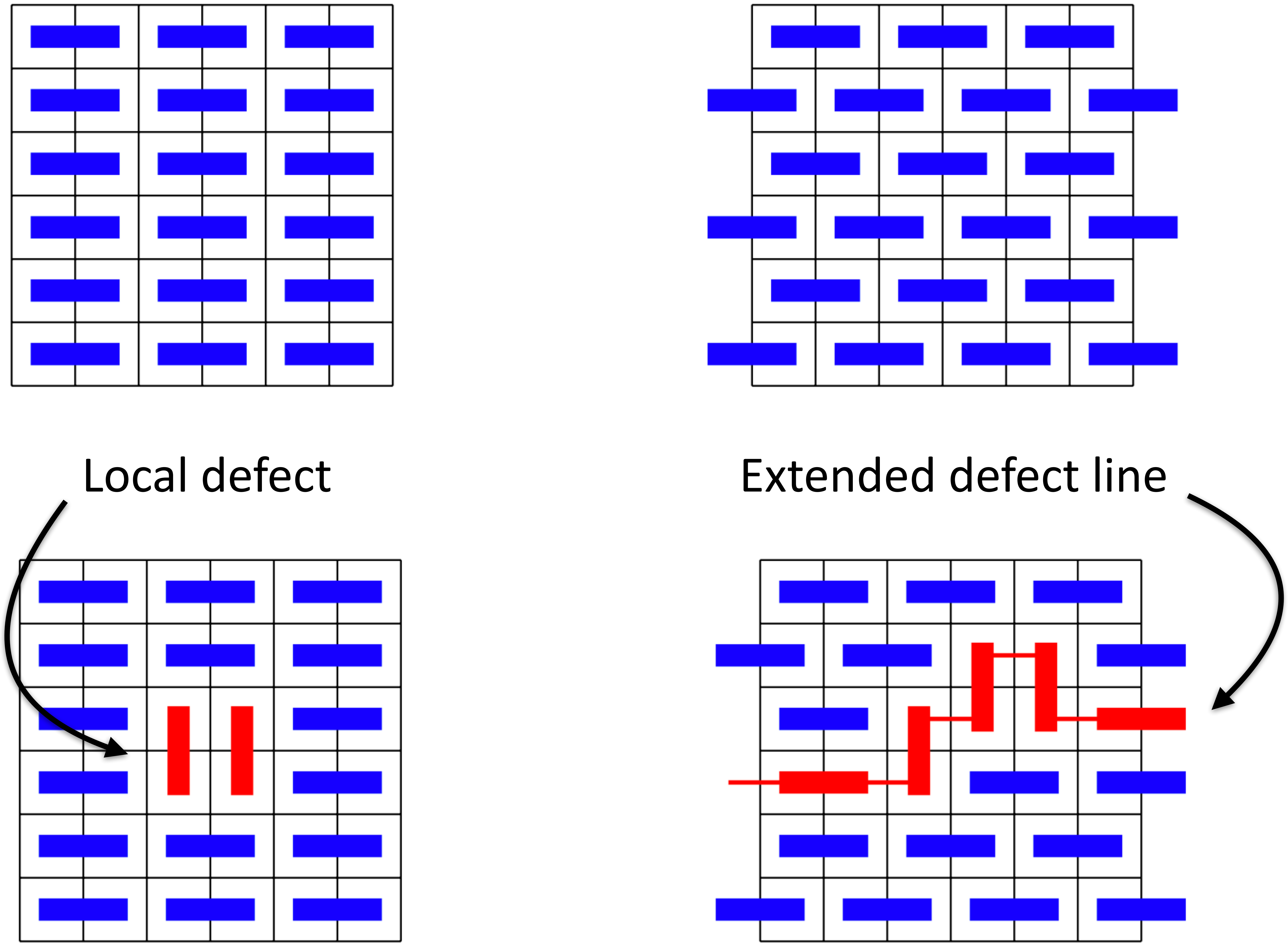
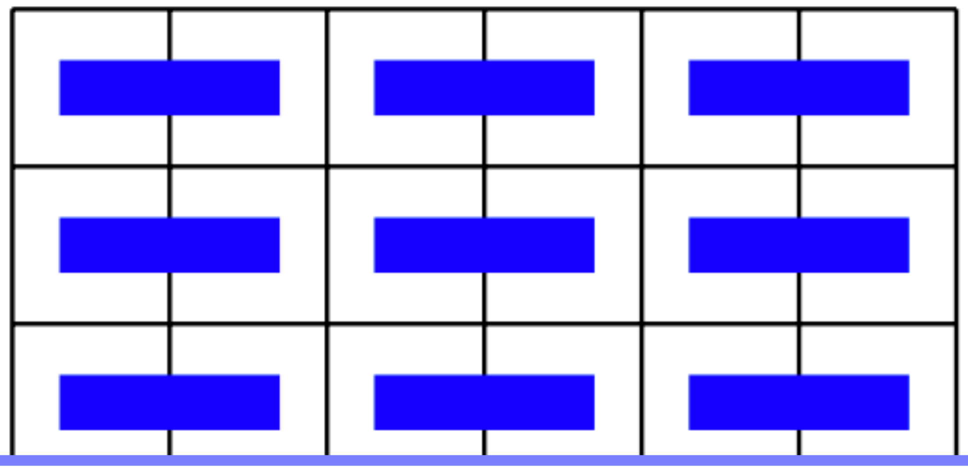
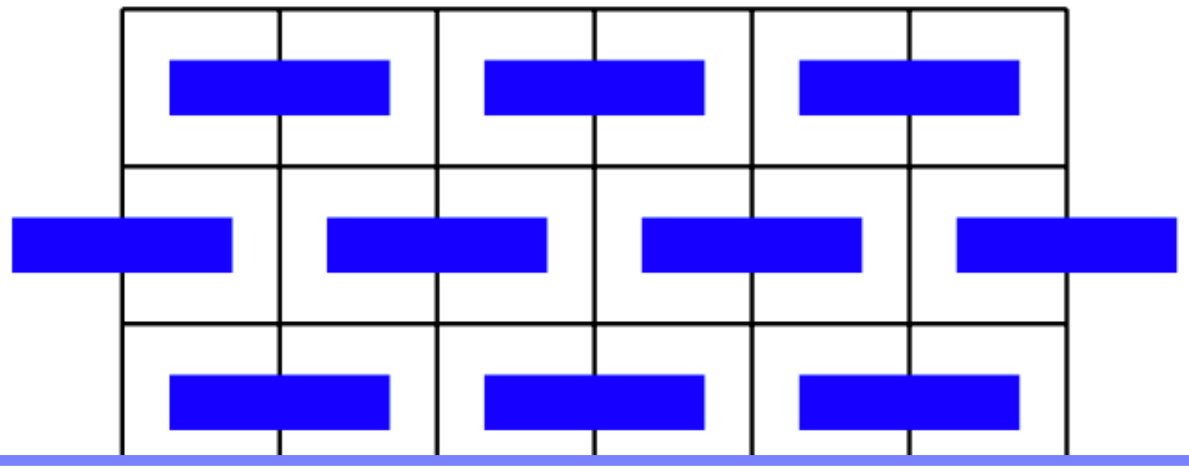
  - Based on "*The arctic curves of the four-vertex model*"
    I. Burenev, F. Colomo, AM, A. Pronko (Journal of Physics A)

Tiling of a square



Tiling of a square



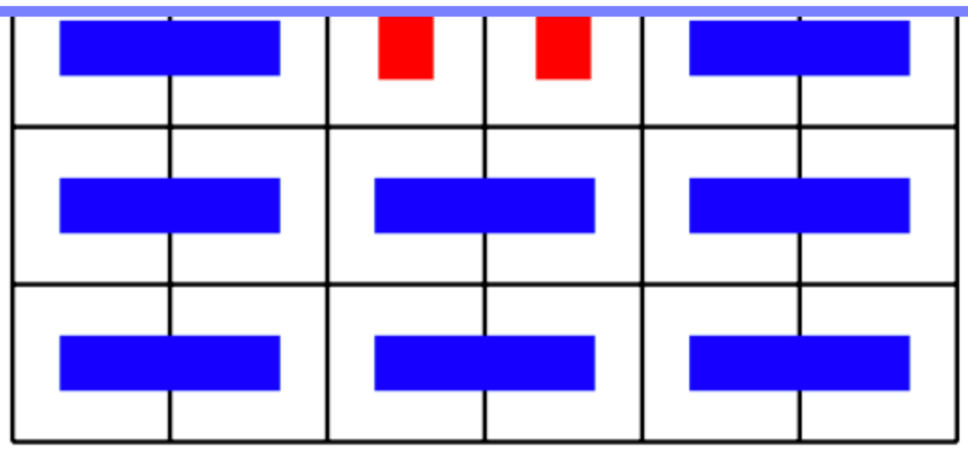


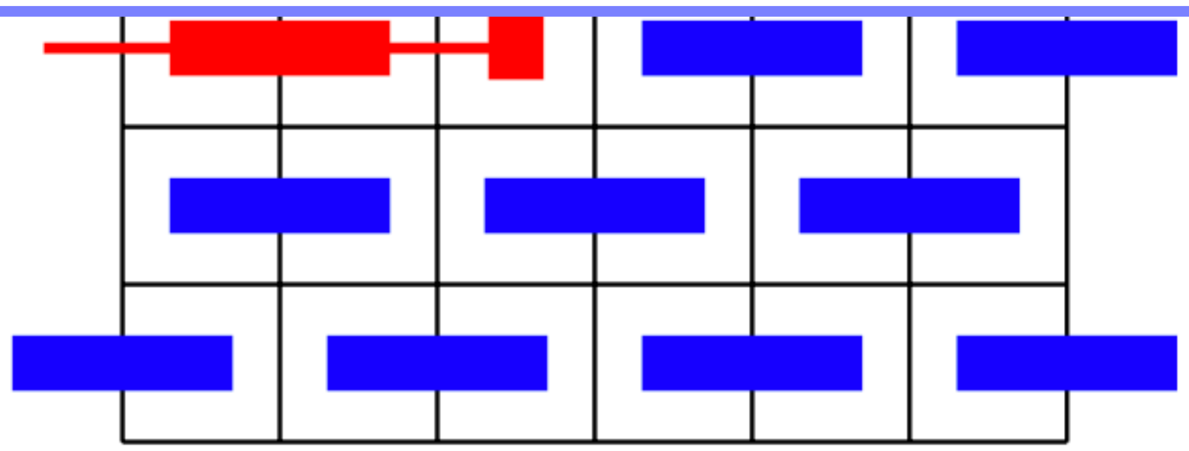
Geometric constraints can induce long range effects

Suitable boundary conditions may modify thermodynamics

Order parameters (free energy, etc…) may acquire spatial dependence

Possibility of **spatial phase separation** in the **scaling limit**

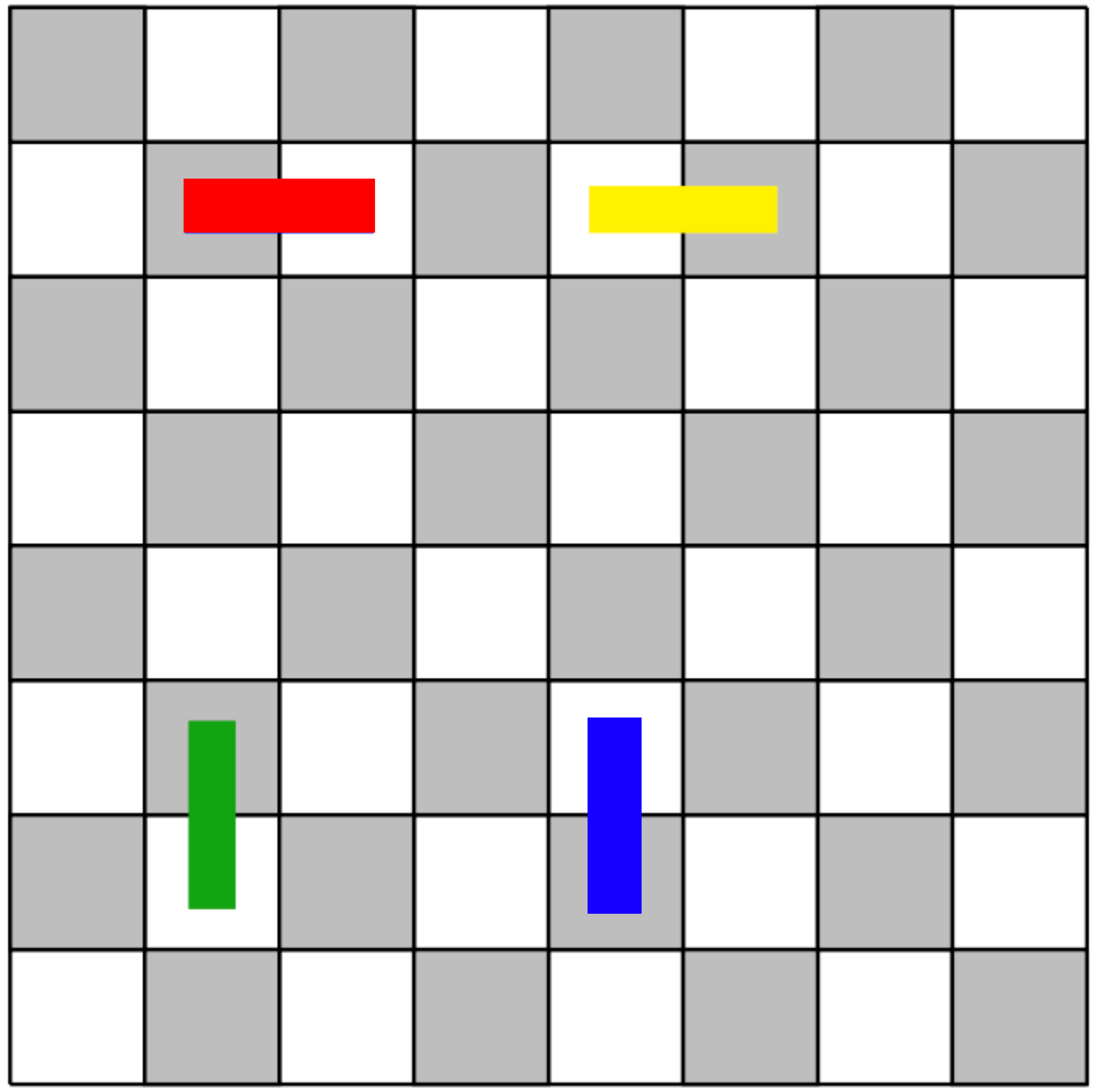




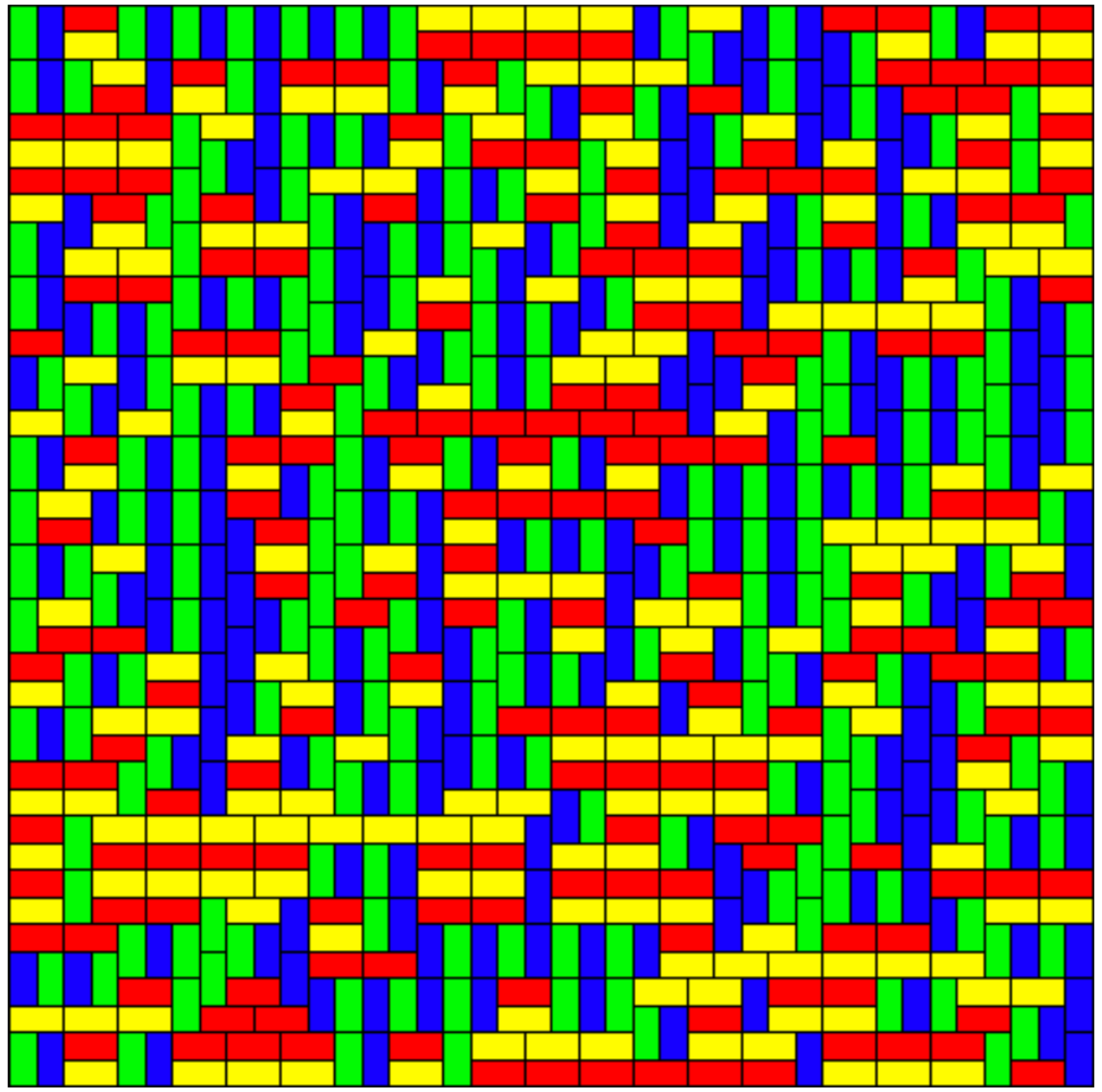
Tiling of a bipartite square

Square $N \times N$

Color coding of the tiles



Random tiling

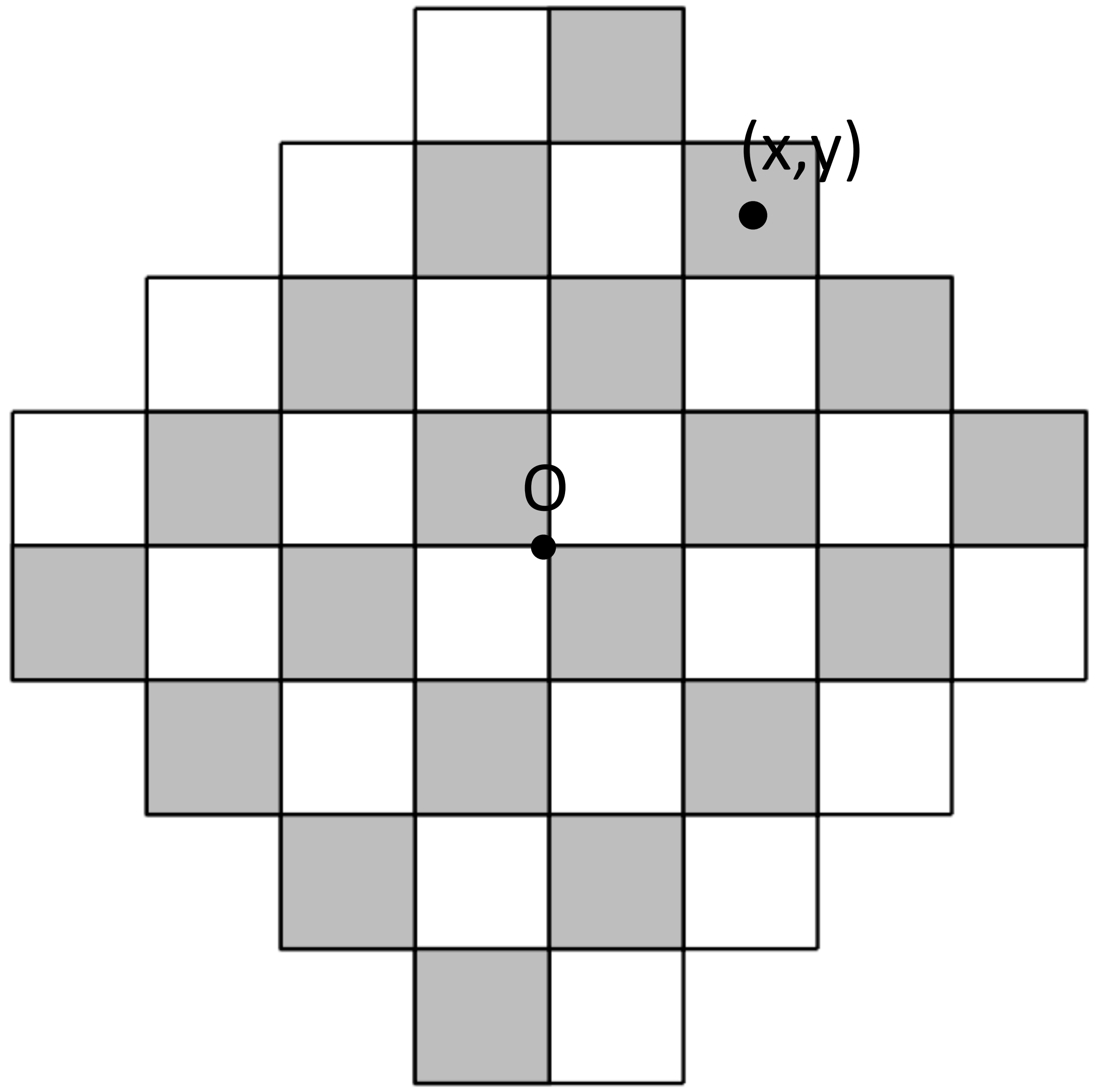


The arctic circle Theorem

Aztec diamond

of order *N*

$$|x| + |y| \leq N$$

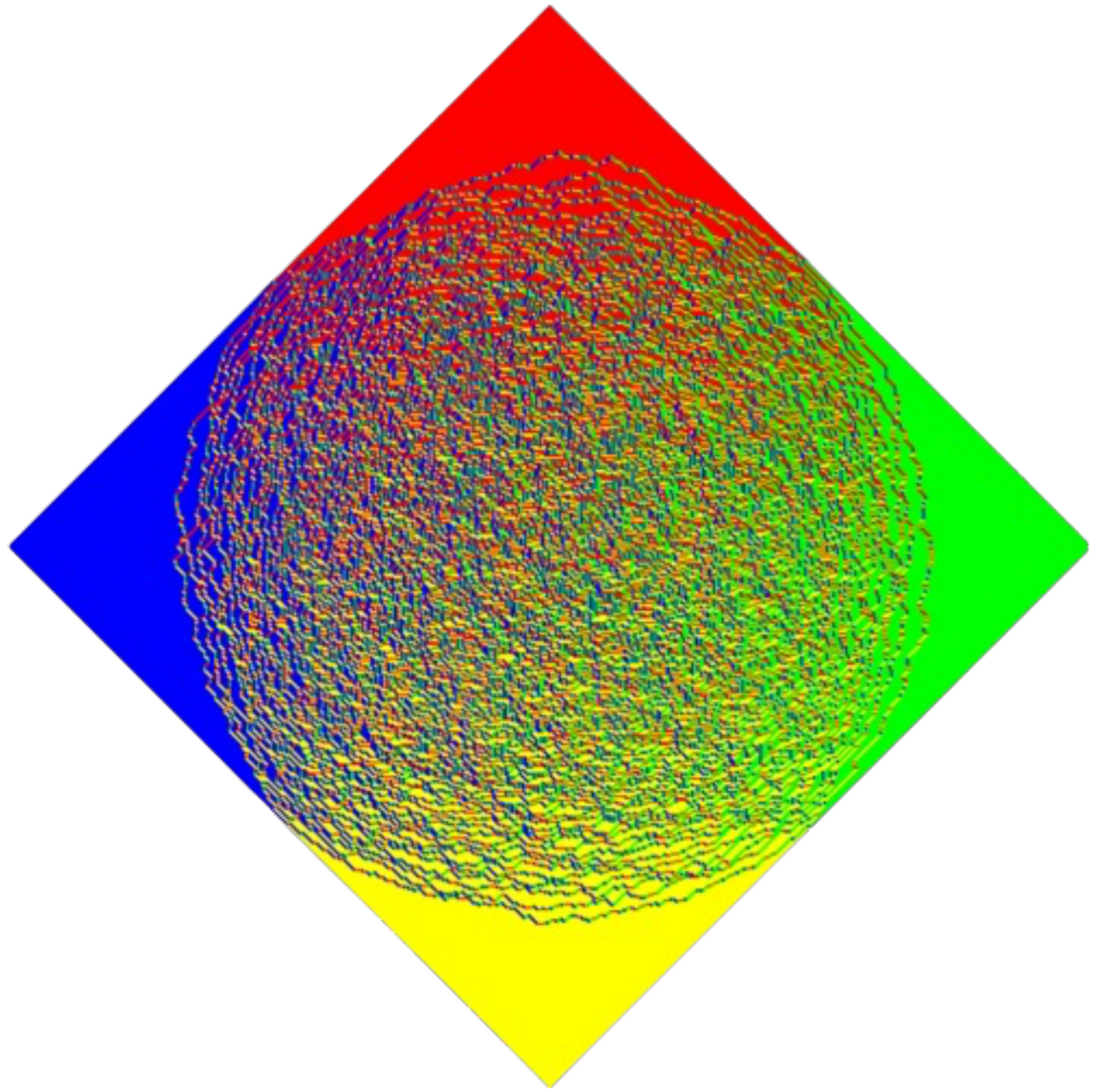


N = 4

Arctic Circle Theorem

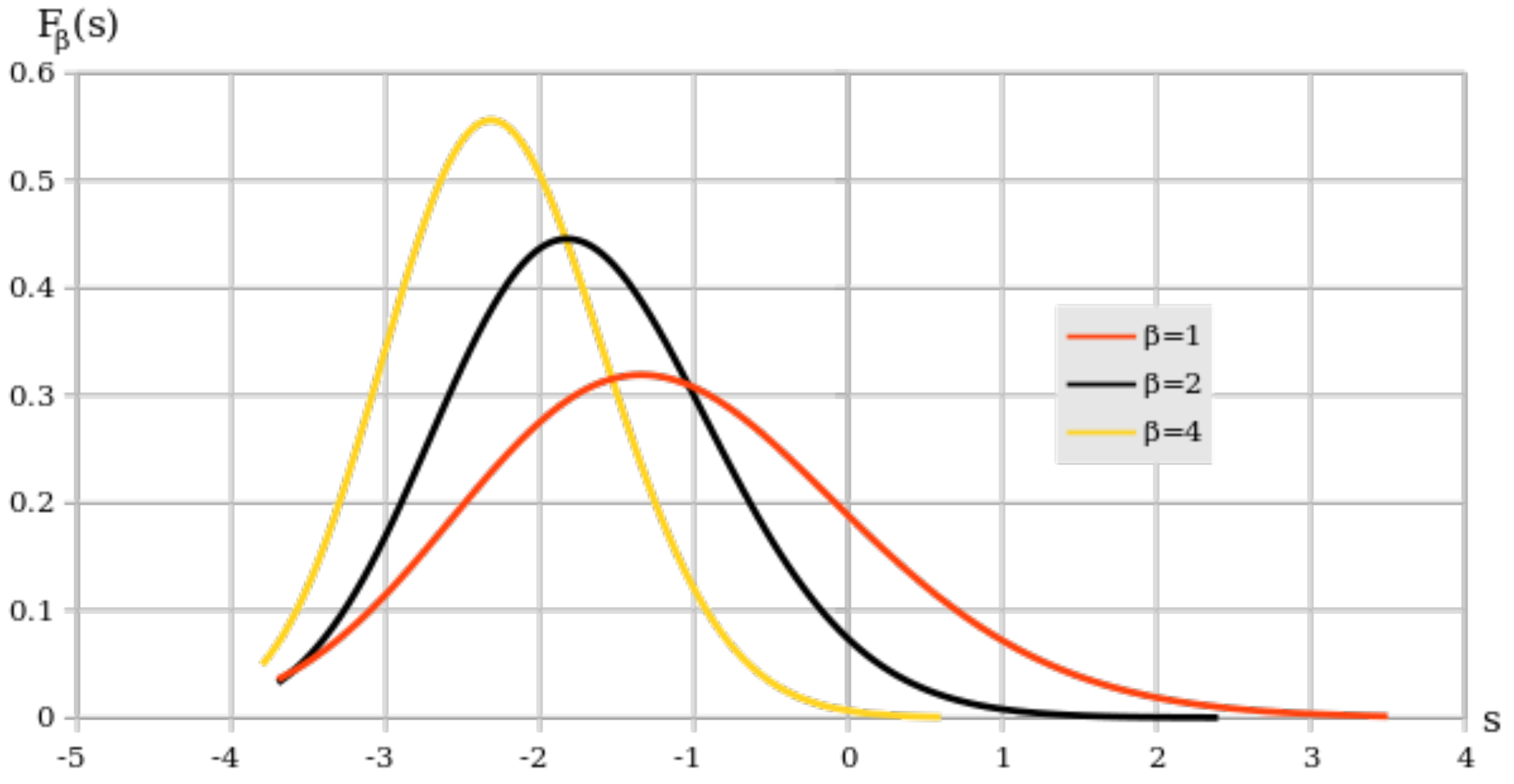
[Jockush - Shor - Propp, 1995]

“In the scaling limit the arctic curve is a circle”

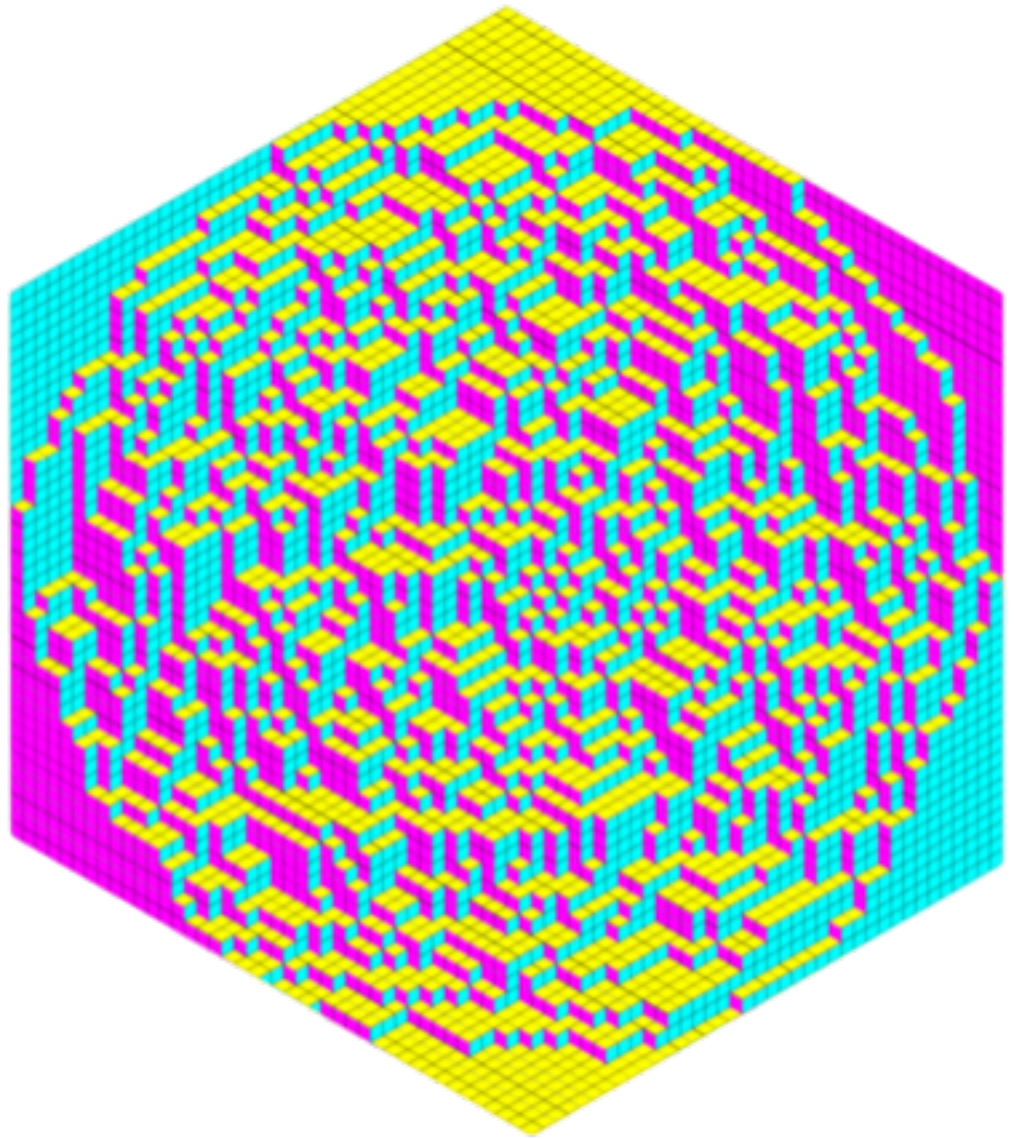
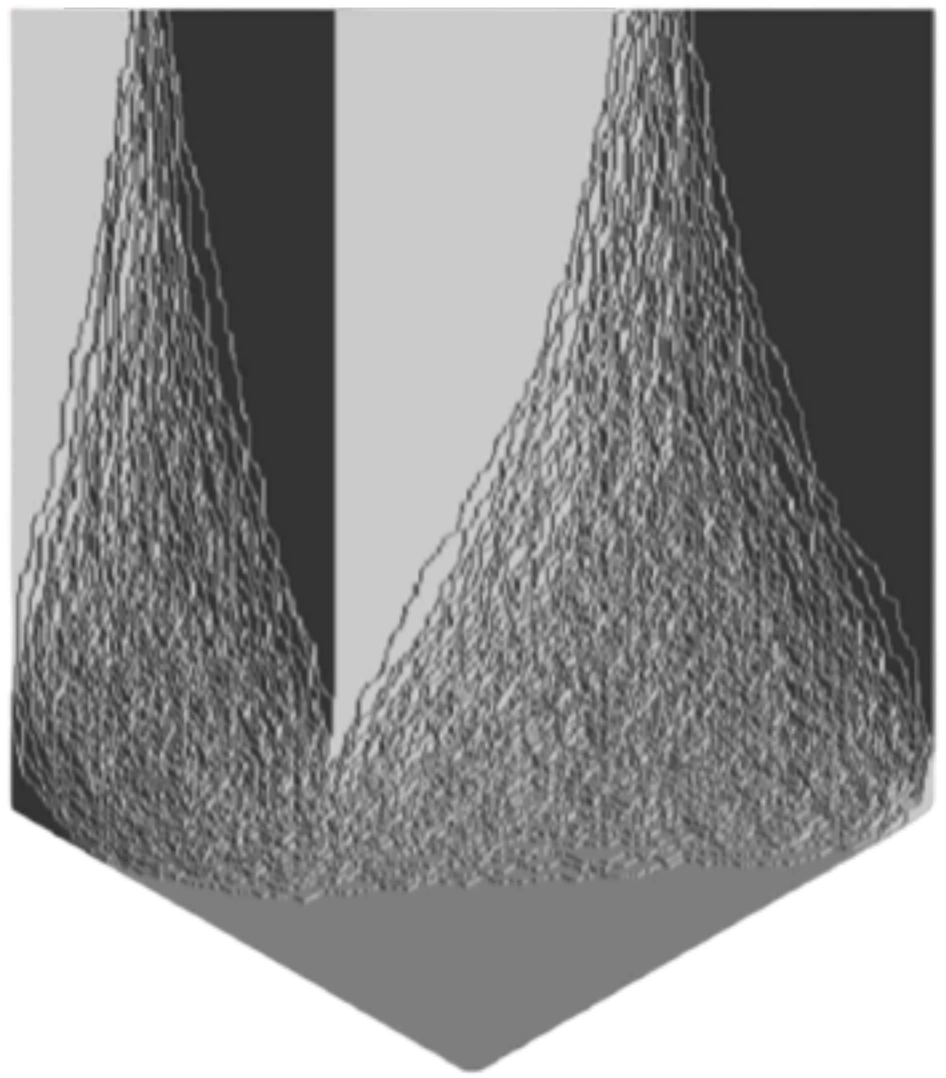
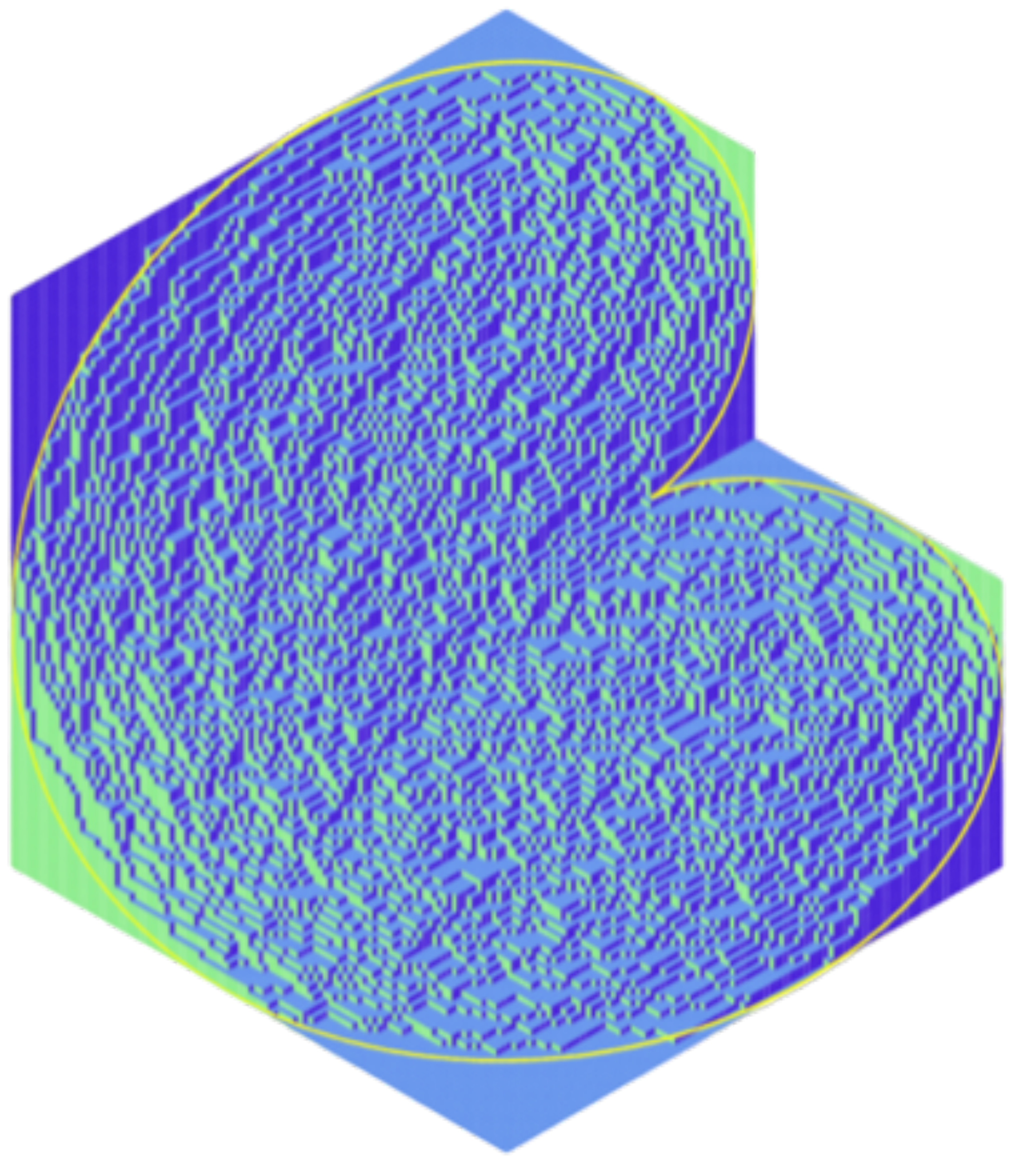


Fluctuations around the arctic circle

- fluctuations $\propto N^{1/3}$ [Johansson, 2000]
- fluctuations along the diagonal are governed by Tracy-Widom distribution [Johansson, 2002]



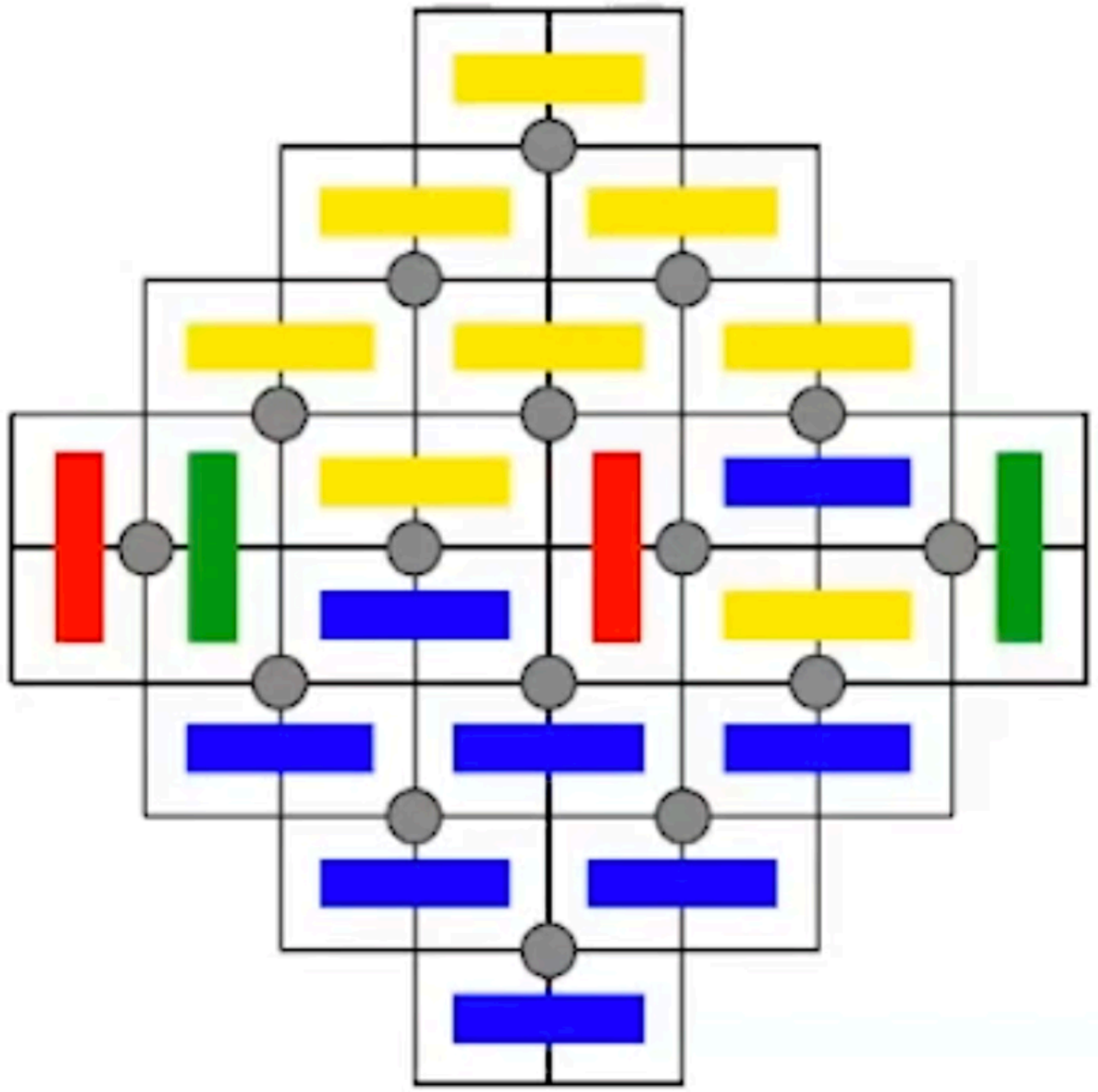
Other tiling models

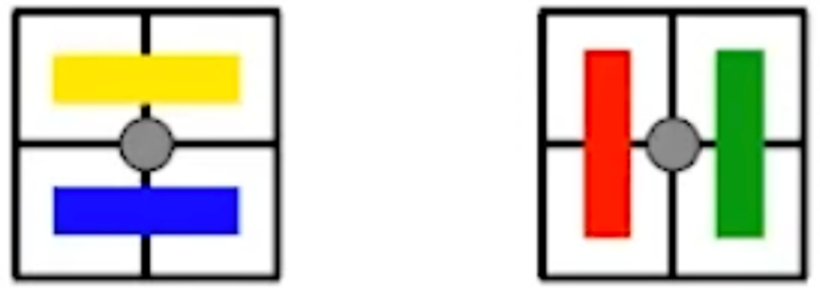
Kenyon, Okounkov and Sheffield formulated a general theory for all dimer models on bipartite graphs with generic boundary conditions and domains. However, it appears that all these models may be viewed as discrete free fermionic models.

[Kenyon - Okounkov - Sheffield, 2006]

Adding the interactions



Assign a Boltzmann weight e^{δ} only to

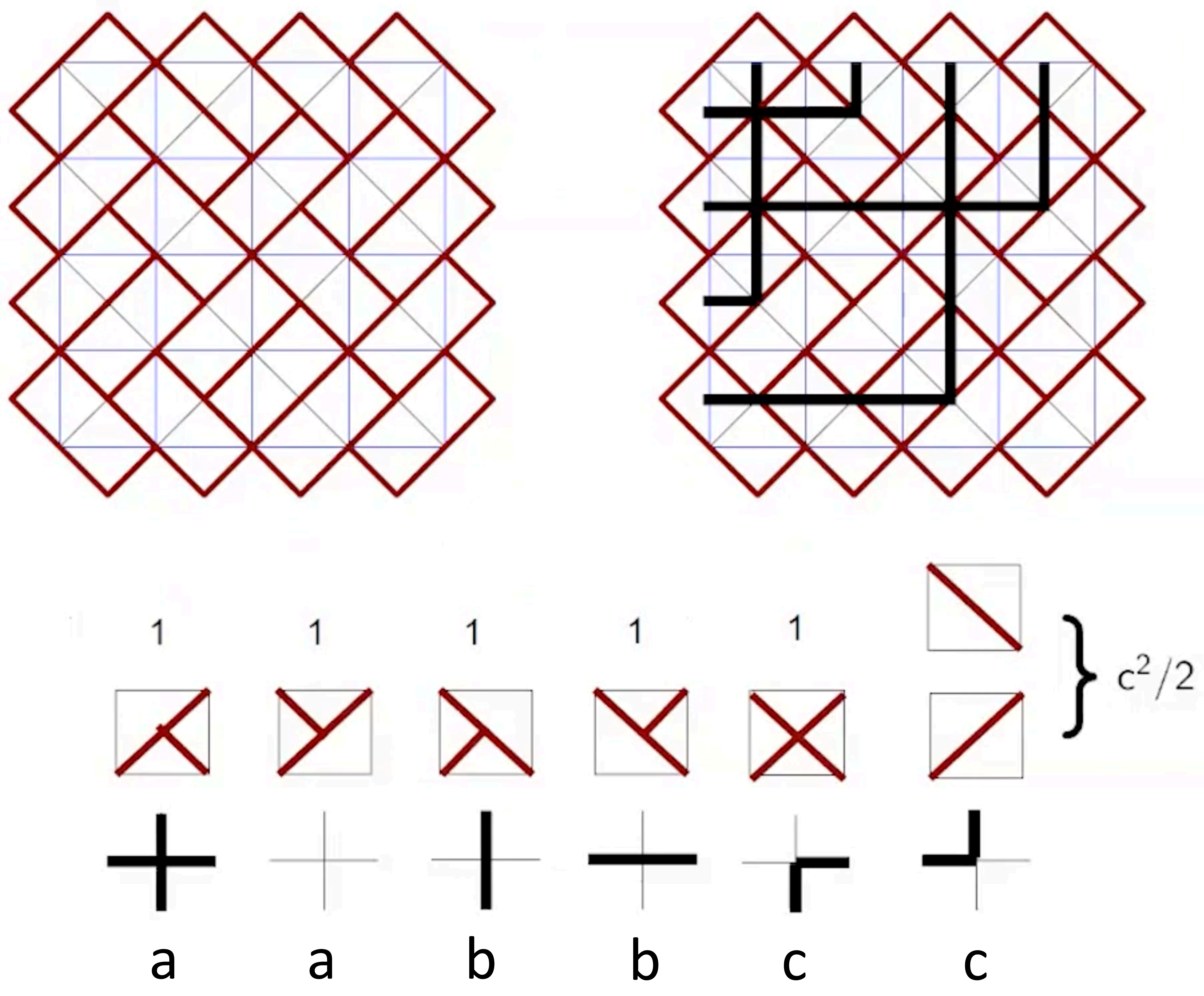


and you get an exactly solvable model:

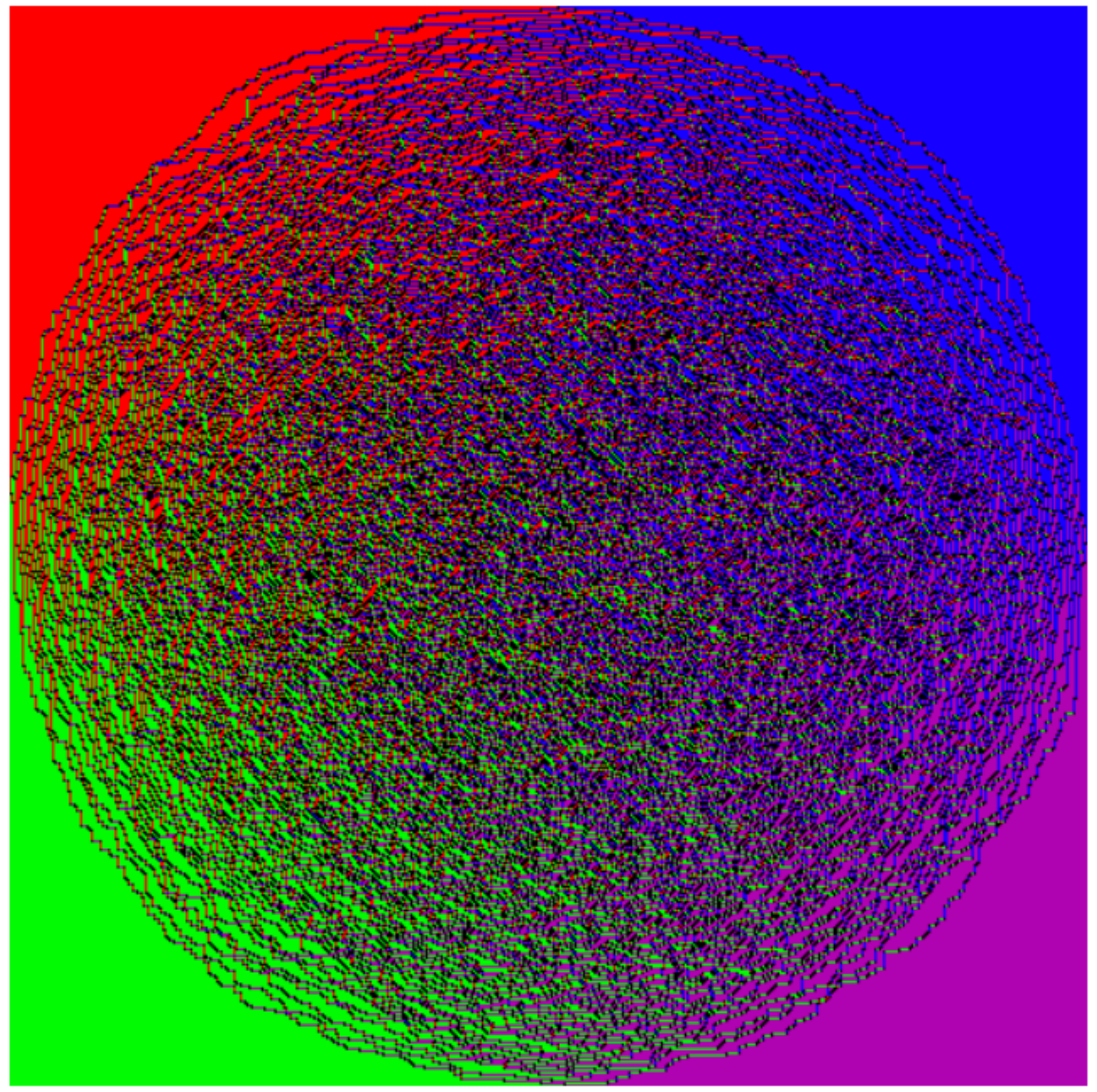
the six-vertex model, with $\Delta = 1 - e^{\delta} (<1)$

[Kuperberg, 1996]

Map from the Aztec diamond to the six-vertex model



$\Delta \neq 0$: Interface fluctuations

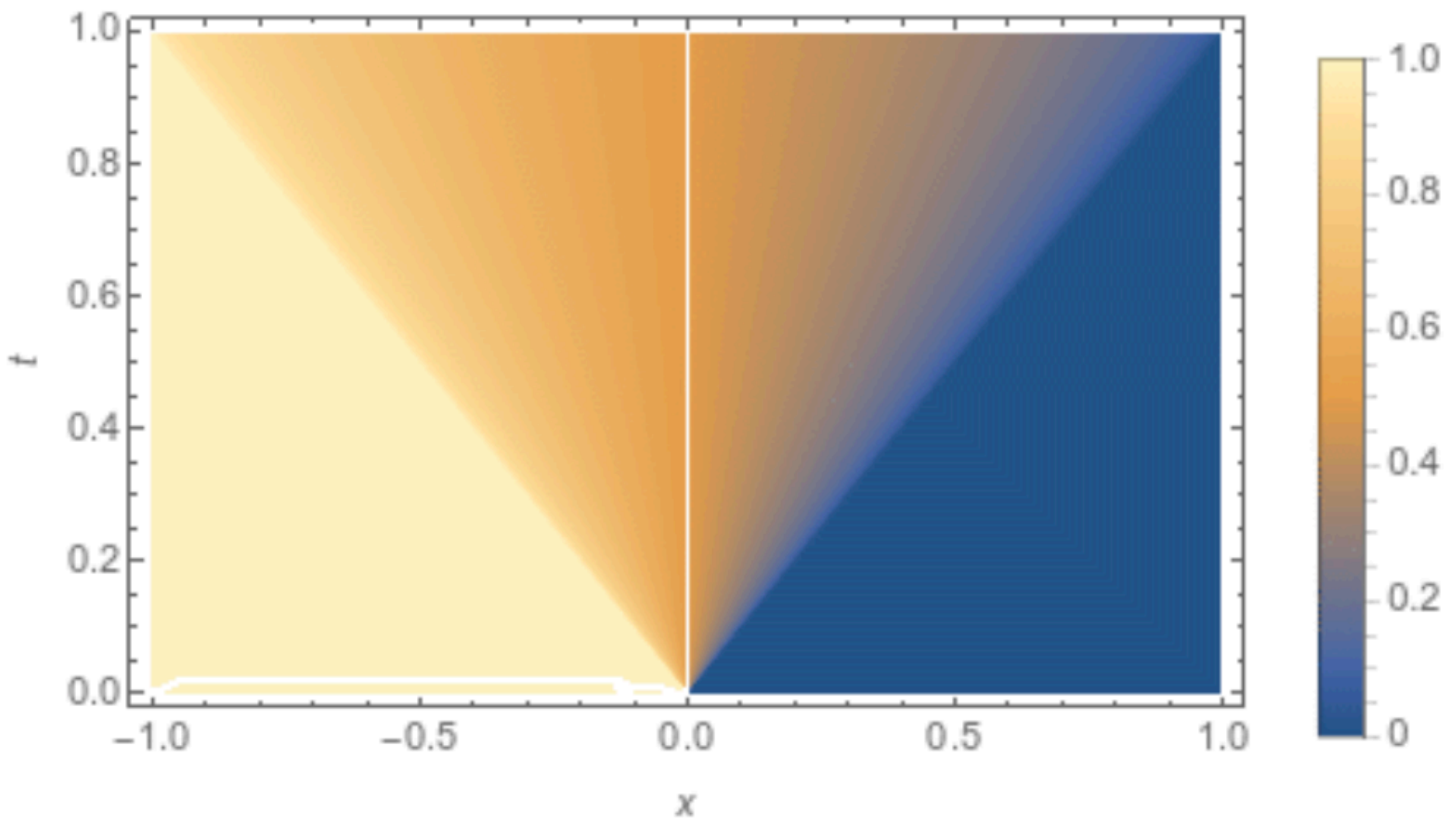


Recent numerics:

T-W scaling is observed with excellent accuracy

[Prauhofer - Spohn, 2023]

[Korepin - Lyberg - Viti, 2023]

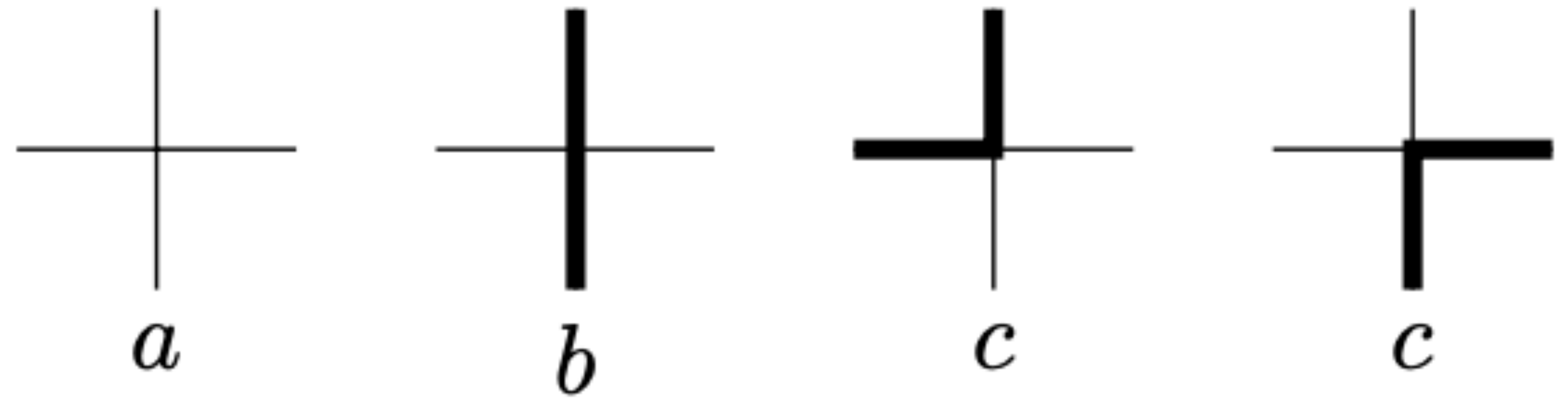


Analytic prediction:

T-W scaling is destroyed

[Collura - De Luca - Viti, 2018]

The four-vertex model [I. Burenev, F. Colomo, AM, A. Pronko, Journal of Physics A]



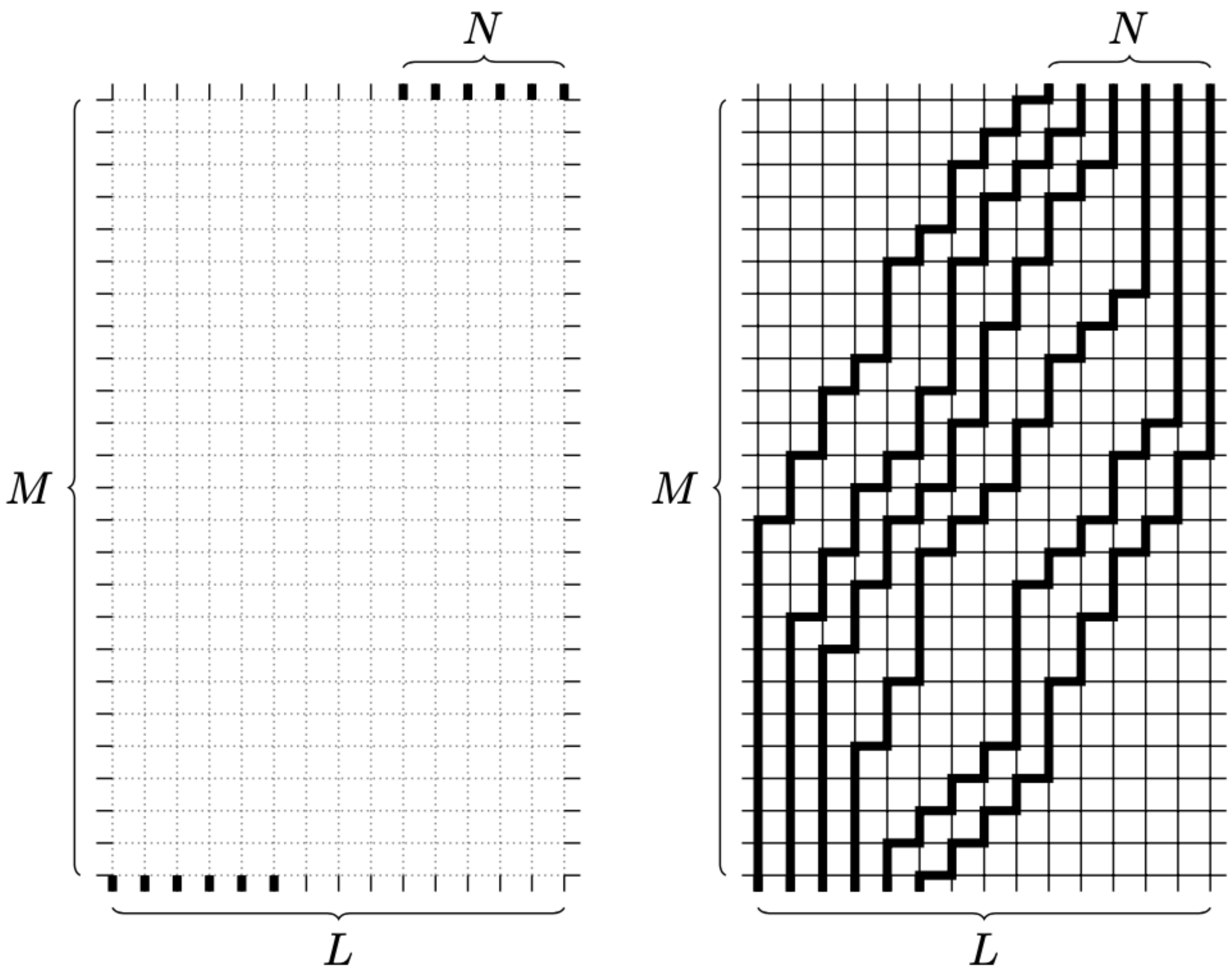
The number of vertices of each type does not depend on the configuration:

$\#a = (L - N)(M - N)$

$\#b = N(M - L + N)$

$\#c = 2N(L - N)$

So we can set $a = b = c = 1$ without loss of generality.



The arctic curves of the four-vertex model

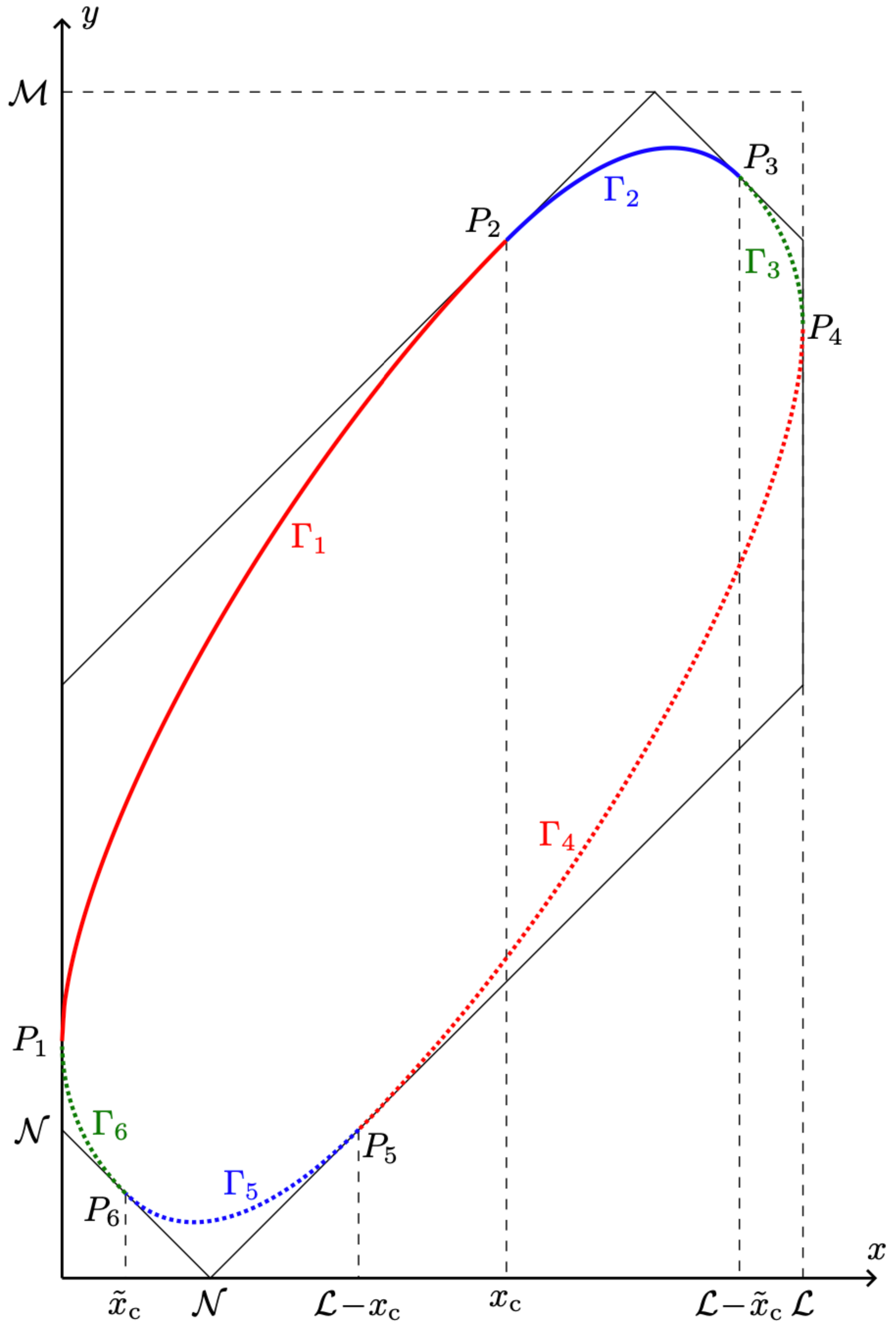
The scaling limit is achieved with:

$$L=[\mathcal{L}\ell],\quad M=[\mathcal{M}\ell],\quad N=[\mathcal{N}\ell],\quad \ell\to\infty.$$

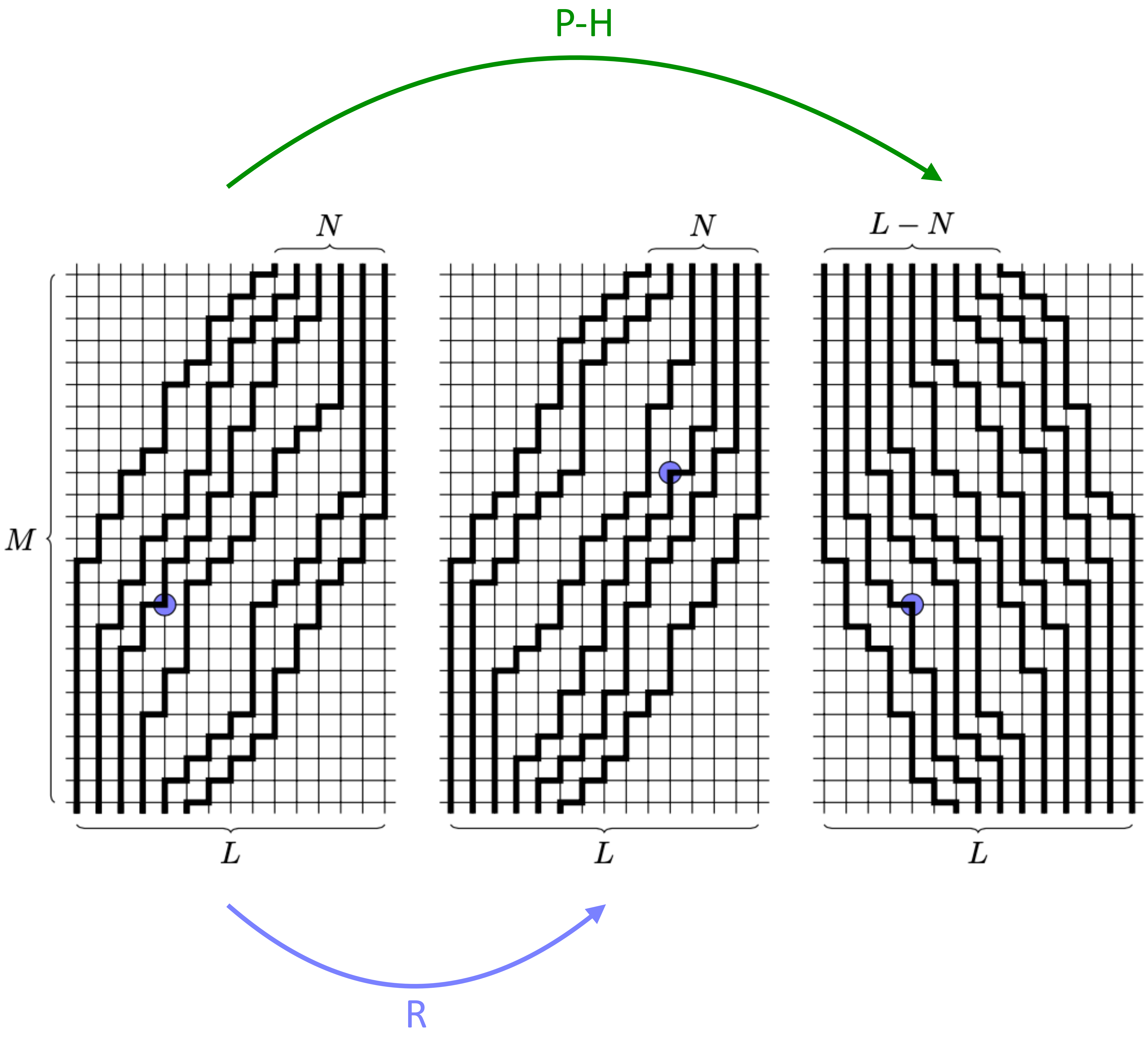
The arctic curve is made of six consecutive arcs, joined end by end at six contact points P_i.

$$\Gamma_1:\ y_1=\frac{\mathcal{M}\mathcal{N}(\mathcal{L}-2x)+(\mathcal{M}+\mathcal{N})\mathcal{L}x}{\mathcal{L}^2}+2\frac{\sqrt{\mathcal{M}\mathcal{N}(\mathcal{L}-\mathcal{N})(\mathcal{M}-\mathcal{L})(\mathcal{L}-x)x}}{\mathcal{L}^2}$$

$$\Gamma_2:\ y_2=(\mathcal{L}-\mathcal{M}-\mathcal{N}-x)+2y_1$$



Particle-hole and reflection symmetries



P-H

Swapping the state (thick ↔ thin) of each vertical edge, and reflecting with respect a vertical axis, we obtain a four-vertex model with L-N lines.

$$y_3(\mathcal{L},\mathcal{M},\mathcal{N};x) = y_1(\mathcal{L},\mathcal{M},\mathcal{L}-\mathcal{N};\mathcal{L}-x)$$

R

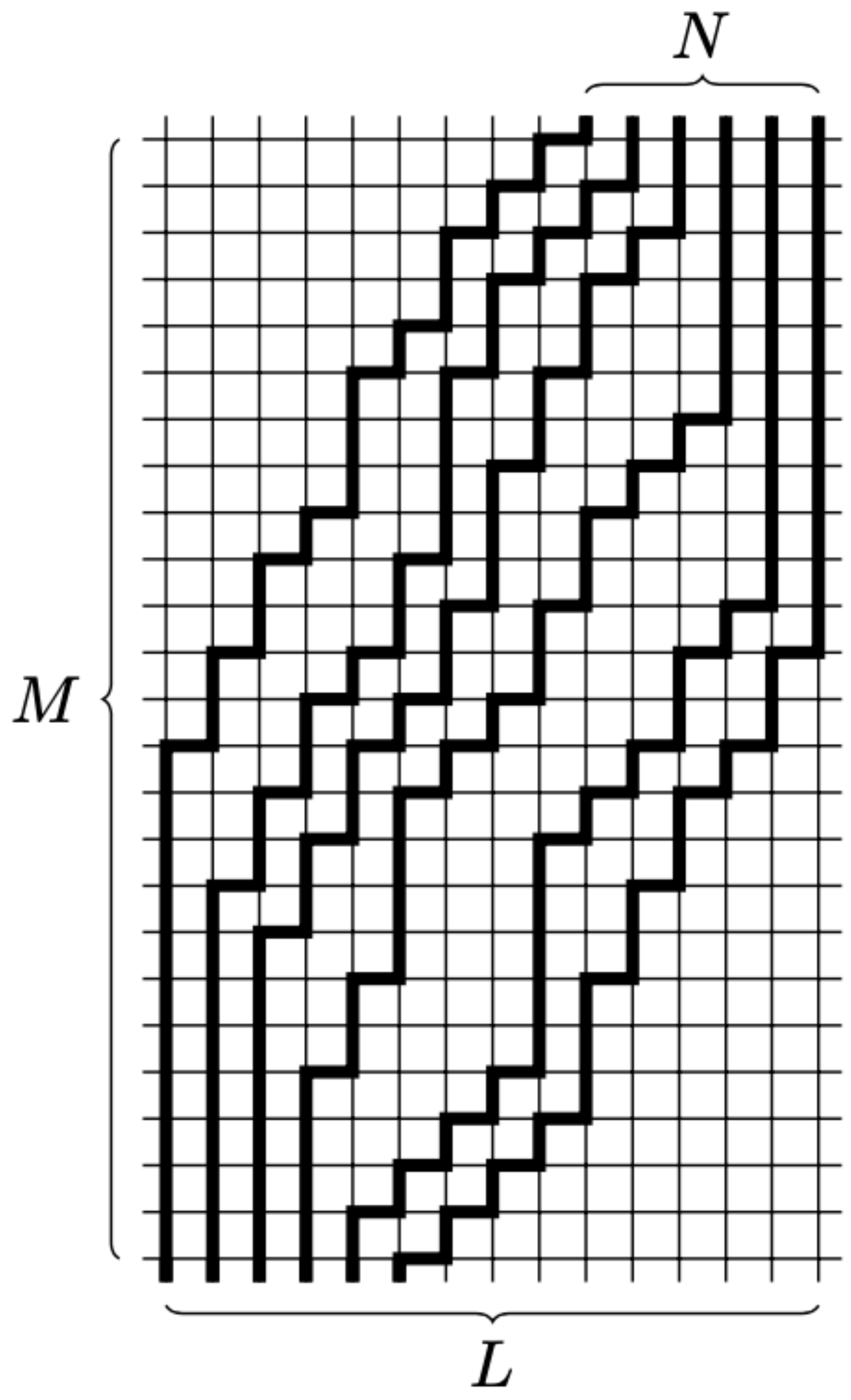
Another symmetry is the simultaneous reflection under the vertical and horizontal axes.

$y_4,\ y_5,\ y_6$ from $y_1,\ y_2,\ y_3$ through:

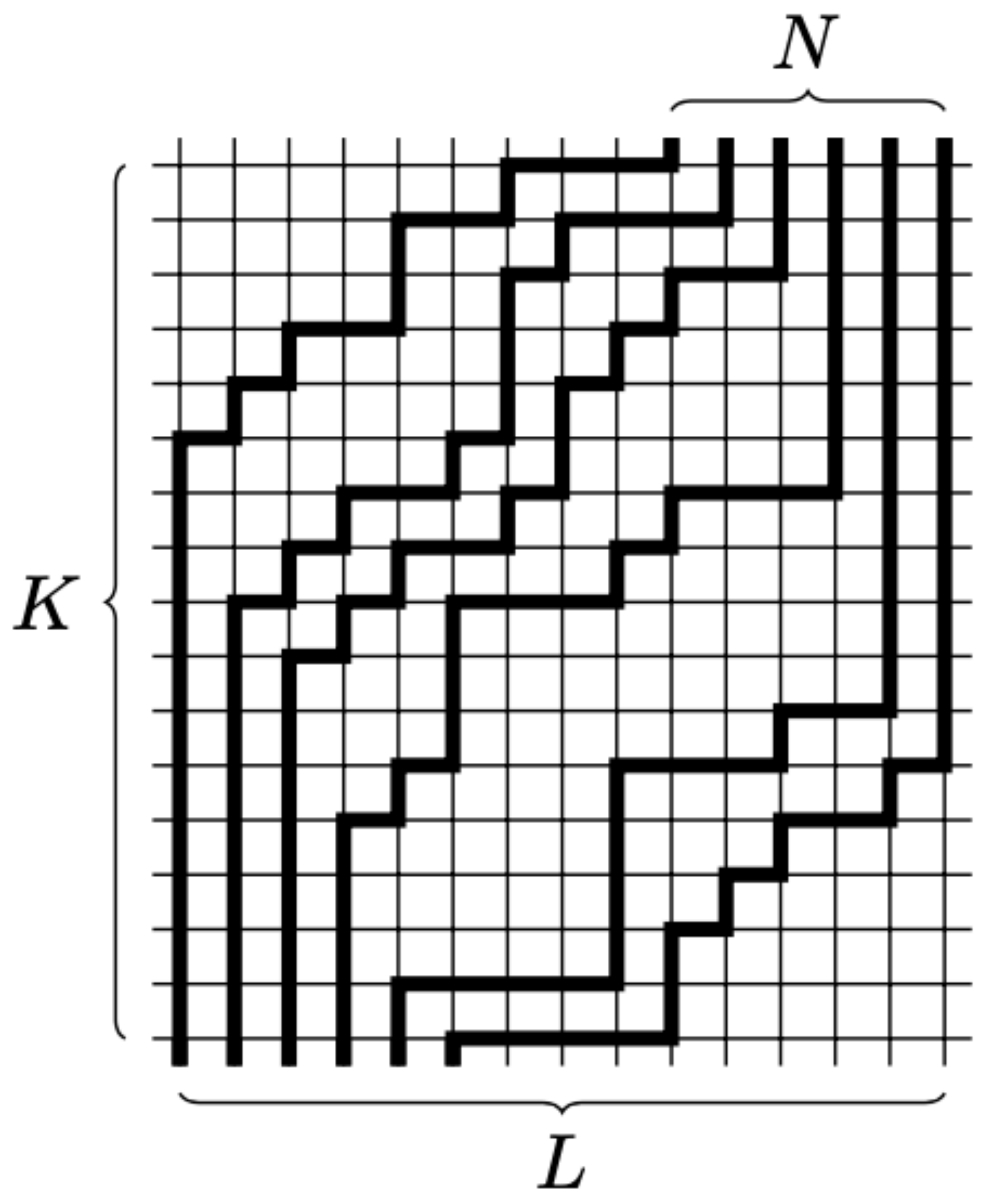
$$x \to \mathcal{L}-x,\ y \to \mathcal{M}-y$$

Methods: bijection to non-intersecting lattice paths

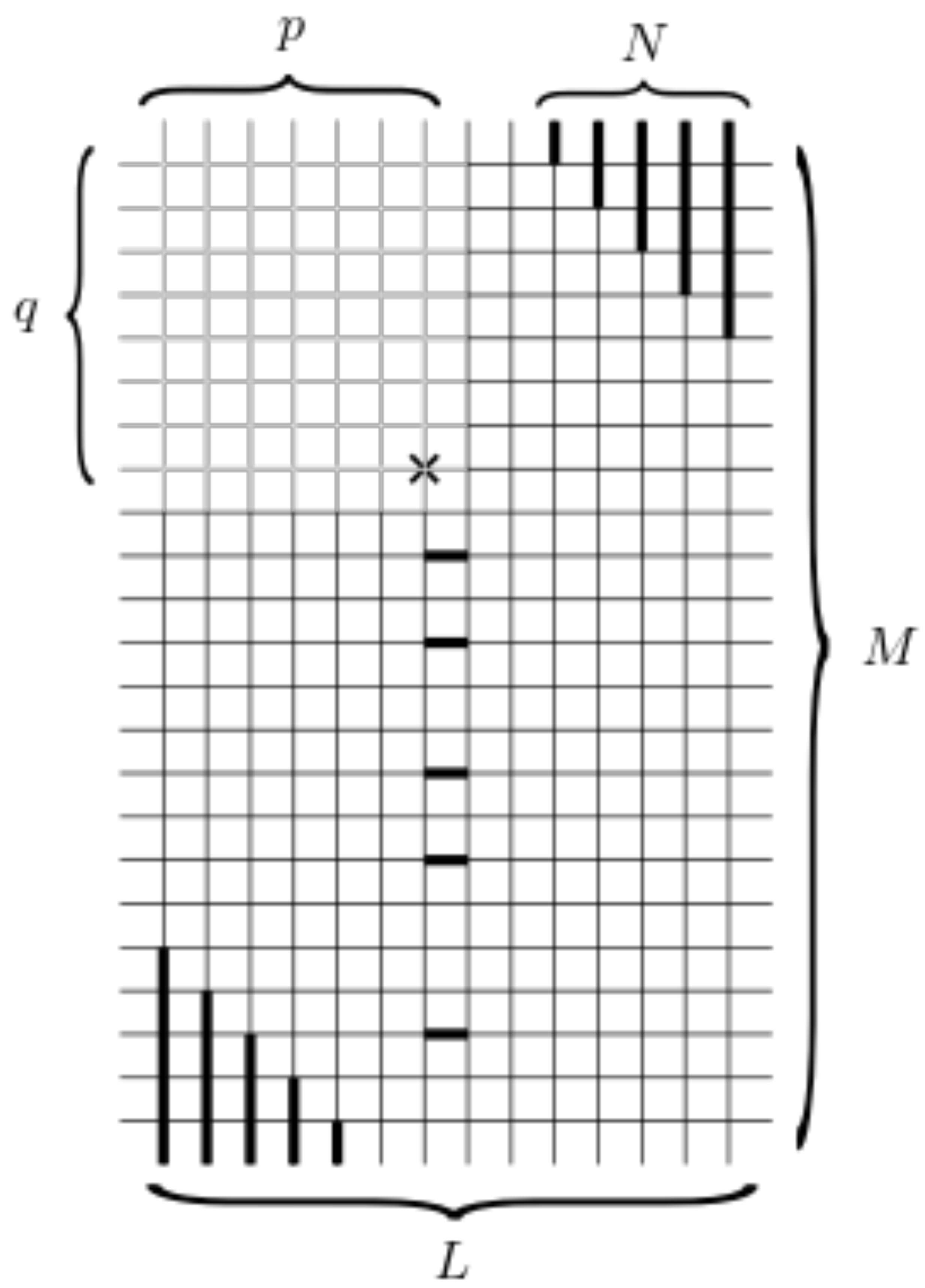
If we shift the *i*-th horizontal edges (*i*=1,…,L-N) of each path by *i*-1 lattice spacing southward, then we have a non-intersecting lattice paths model

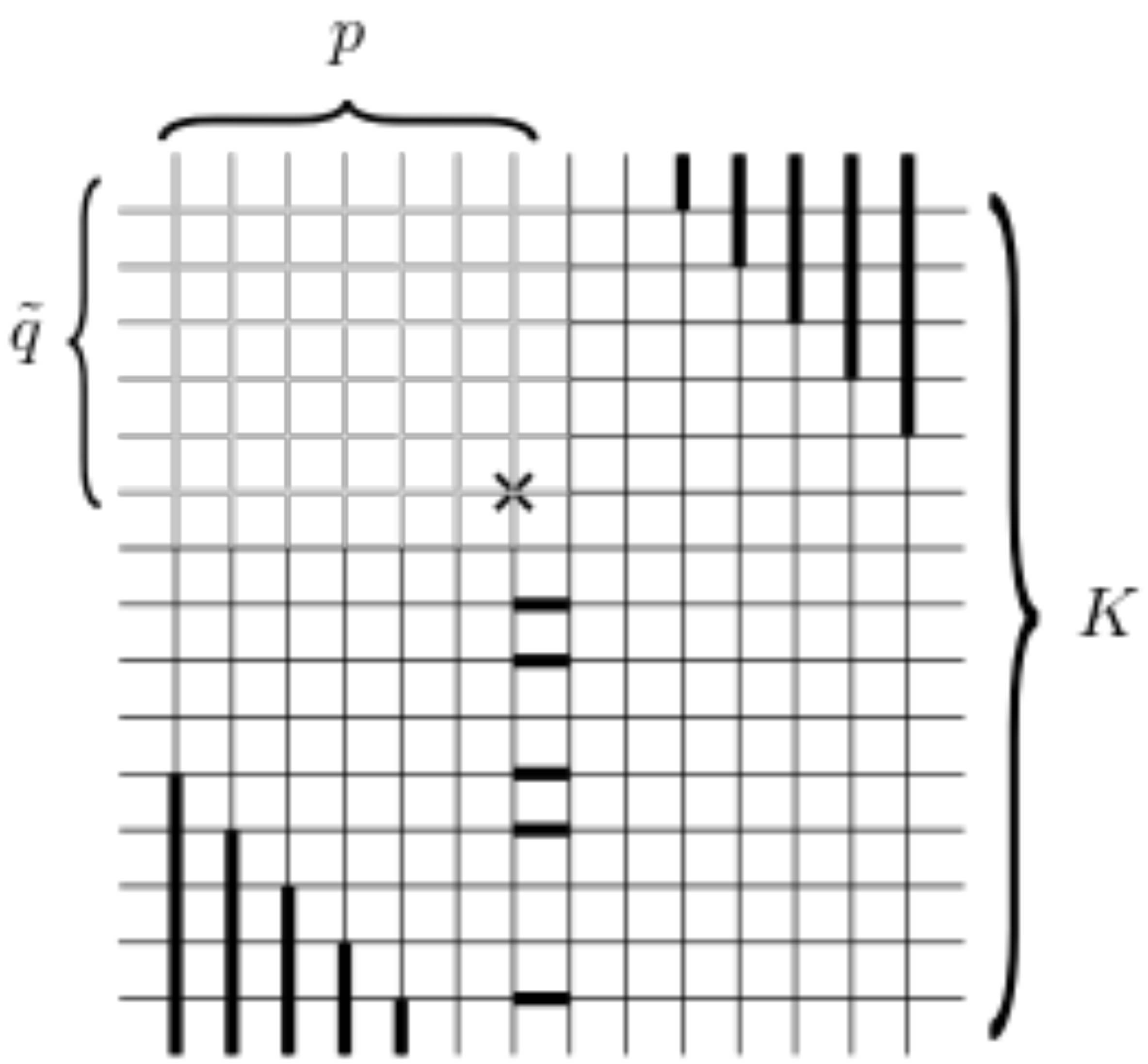


$$K := M - L + N + 1$$



The emptiness formation probability (EFP)





The EFP is the probability to have at least q vertices of type *a* starting from the point (p,M)

After the bijection, it is the probability that no paths pass through the point $(p, K - \tilde{q} + 1)$, with $\tilde{q} = p + q - L + N$

$$H(d,\alpha,\beta,s,n) := \frac{1}{H_0(\alpha,\beta,s,n)} \sum_{0 \le x_1,\ldots,x_s \le d} \prod_{1 \le i < j \le s} (x_j - x_i)^2 \prod_{i=1}^{s} \binom{\alpha + x_i}{x_i} \binom{\beta + n - x_i}{n - x_i}$$

Hahn measure

$$d = M - N + \min(p,N) - p - q, \qquad \alpha = |N - p|, \qquad \beta = L - N - p,$$
$$s = \min(p,N), \qquad n = M - L + \min(p,N).$$

EFP in a log-gas representation

Let $\mathbf{x} = \{x_1, \ldots, x_s\}$ with $0 \leq x_1 < \cdots < x_s \leq n$ denote the position of s particles on the interval $[0, n]$ and consider the probability measure on $[0, n]^s$:

$$P_{n,s}^{(\alpha,\beta)}[\mathbf{x}] = \frac{1}{Z(\alpha, \beta, s, n)} \prod_{1 \leq i < j \leq s} (x_i - x_j)^2 \prod_{i=1}^{s} w_n^{(\alpha,\beta)}(x_j)$$

$$w_n^{(\alpha,\beta)}(x) = \binom{\alpha + x}{x} \binom{\beta + n - x}{n - x}$$

$$H(d, \alpha, \beta, s, n) := \sum_{0 \leq \mathbf{x} \leq d} P_{n,s}^{(\alpha,\beta)}[\mathbf{x}].$$

The **EFP is the probability that**, in a discrete log-gas of *s* particles, associated to the Hahn measure, **no particle has coordinate larger than *d*.**

The arc of the arctic curve

The **scaling limit** is achieved with: $d=[d_0\ell],\ \alpha=[\alpha_0\ell],\ \beta=[\beta_0\ell],\ s=[s_0\ell],\ n=[n_0\ell],\ \ell\to\infty$

$$\textbf{EFP} = \begin{cases} 1 & \text{Frozen region} \\ 0 & \text{Otherwise} \end{cases}$$

To find the arctic arc, it is sufficient the **support [L,R]** of the density.

We have the **arctic curve** when:

$$R(\alpha_0,\beta_0,s_0,n_0)=d_0 \longrightarrow y_1=\frac{\mathcal{M}\mathcal{N}(\mathcal{L}-2x)+(\mathcal{M}+\mathcal{N})\mathcal{L}x}{\mathcal{L}^2}+2\frac{\sqrt{\mathcal{M}\mathcal{N}(\mathcal{L}-\mathcal{N})(\mathcal{M}-\mathcal{L})(\mathcal{L}-x)x}}{\mathcal{L}^2}$$

(setting $x=p_0,\ y=\mathcal{M}-q_0$)

Conclusions and Outlooks

- Main Results
  - Calculation of the arctic curve for the four-vertex model
  - Observation that the fluctuations are governed by T-W

- Further investigations
  - Generalization to the five-vertex model

Thank you for your time!