

GRAVITY

Stars orbiting the Galactic Black Hole

Stefan Gillessen
on behalf of

The GRAVITY
collaboration



The mass measurement of Sgr A*,
the massive black hole candidate in the Galactic Center,
was honored with the Nobel prize



Outline

- Part I The mass of Sgr A*
 - Astronomical black holes
 - Measuring mass
 - Infrared observations
 - Adaptive Optics
 - Stellar Orbits
- Part II Errors, Fitting and all that
- Part III Testing the black hole paradigm
 - The black hole nature of Sgr A*
 - Interferometry
 - Testing General Relativity in the Galactic Center
 - SgrA* flares
 - (A funny gas cloud)

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$$R_S = \frac{2GM}{c^2}$$

How to find a black hole

1. Find an object that has an event horizon
2. Show that a certain amount of mass is confined to a volume smaller than R_S
3. Show that an object is unstable against its self-gravity








Note: Density is ill-defined for black holes

$$\rho \sim \frac{M}{\frac{4\pi}{3} R_S^3} \sim \frac{1}{M^2}$$

Black holes are astronomical objects

- 1kg: $R_S = 10^{-27}$ m;
(proton: 10^{-15} m) $\rho = 10^{77}$ g/cm³
- Earth: $R_S = 1$ cm; $\rho = 10^{27}$ g/cm³
- Sun: $R_S = 3$ km; $\rho = 10^{16}$ g/cm³
- Star cluster:
(10^5 stars) $R_S = 3 \times 10^8$ m;
(distance to the moon) $\rho = 10^6$ g/cm³
- Galaxy
(10^{11} stars) $R_S = 3 \times 10^{14}$ m;
(2000 AU) $\rho = 10^{-6}$ g/cm³
(thinner than air)

Types of black holes

- Particle physics scale
 - LHC, cosmic radiation
 - Should be bright & unstable due to their Hawking radiation
- Planetary mass – primordial black holes
 - primordial
- Stellar mass black holes ($1 - 10 M_{\text{sun}}$)
 - End products of massive stars
- Intermediate mass black holes ($10^3 - 10^5 M_{\text{sun}}$)
 - Primordial, or first generation of stars
- Supermassive black holes ($10^6 - 10^9 M_{\text{sun}}$)
 - in the center of galaxies
 - Quasars, Active Galactic Nuclei

A short history of black holes

1783 John Michell

Basic idea

1915 Albert Einstein:

General relativity

1916 Karl Schwarzschild

Black hole solution

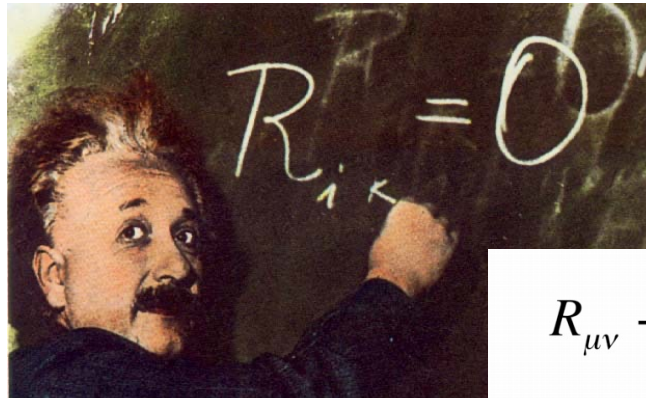
1964 John Wheeler

“Black hole”

1969 Quasars speculated

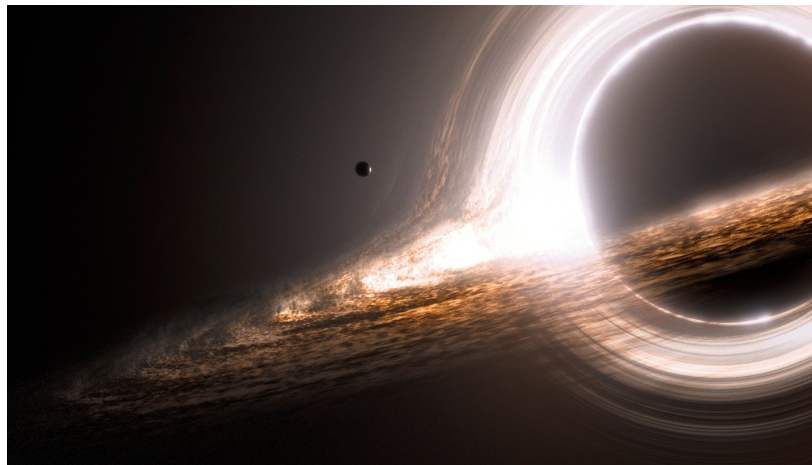
1971 Cyg X-1's mass

2002 Galactic Center



$$R_{\mu\nu} - \frac{1}{2}R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$ds^2 = c^2 \left(1 - \frac{2GM}{c^2 r}\right) dt^2 - \left(1 - \frac{2GM}{c^2 r}\right)^{-1} dr^2 - r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$



PHILOSOPHICAL
TRANSACTIONS
OF THE
ROYAL SOCIETY
OF
LONDON.

VOL. LXXIV. For the Year 1784.

PART I. 2



LONDON,

SOLD BY LOCKYER DAVIS, AND PETER ELMSLY,
PRINTERS TO THE ROYAL SOCIETY.

MDCCLXXXIV.

VII. *On the Means of discovering the Distance, Magnitude, &c. of the Fixed Stars, in consequence of the Diminution of the Velocity of their Light, in case such a Diminution should be found to take place in any of them, and such other Data should be procured from Observations, as would be farther necessary for that Purpose. By the Rev. John Michell, B. D. F. R. S. In a Letter to Henry Cavendish, Esq. F. R. S. and A. S.*

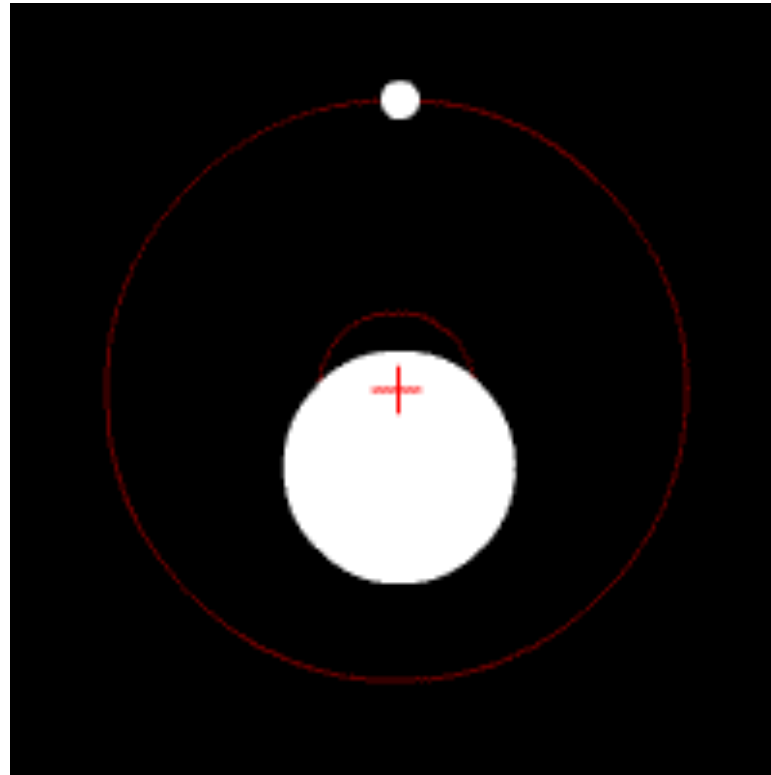
P. 35.

16. Hence, according to article 10, if the semi-diameter of a sphere of the same density with the sun were to exceed that of the sun in the proportion of 500 to 1, a body falling from an infinite height towards it, would have acquired at its surface a greater velocity than that of light, and consequently, supposing light to be attracted by the same force in proportion to its vis inertiae, with other bodies, all light emitted from such a body would be made to return towards it, by its own proper gravity.

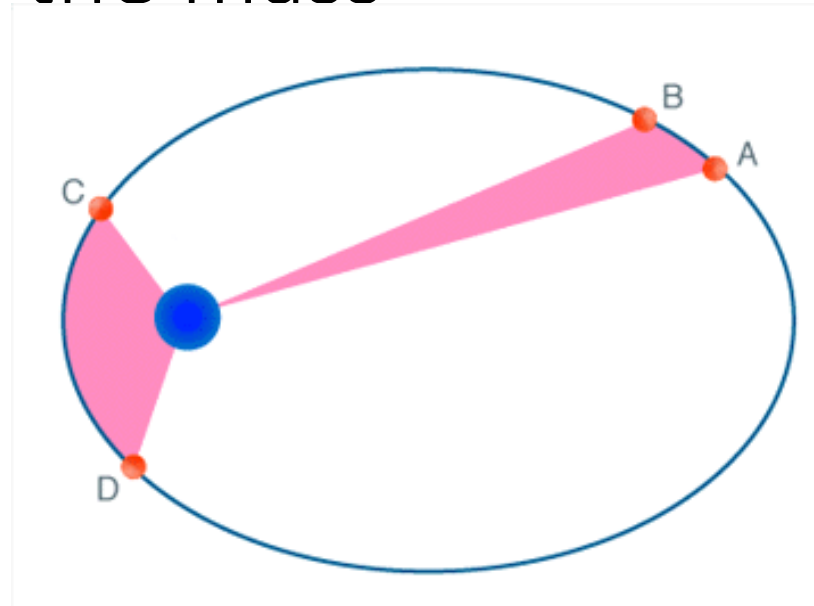
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Unseen mass can shake a visible object

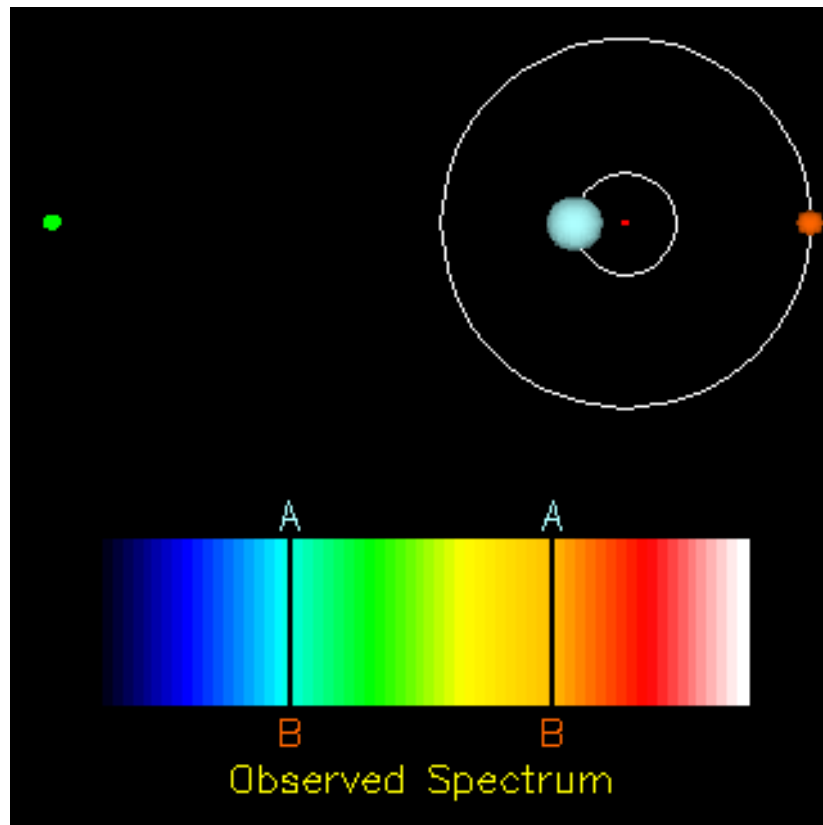


Kepler ellipses yield the mass

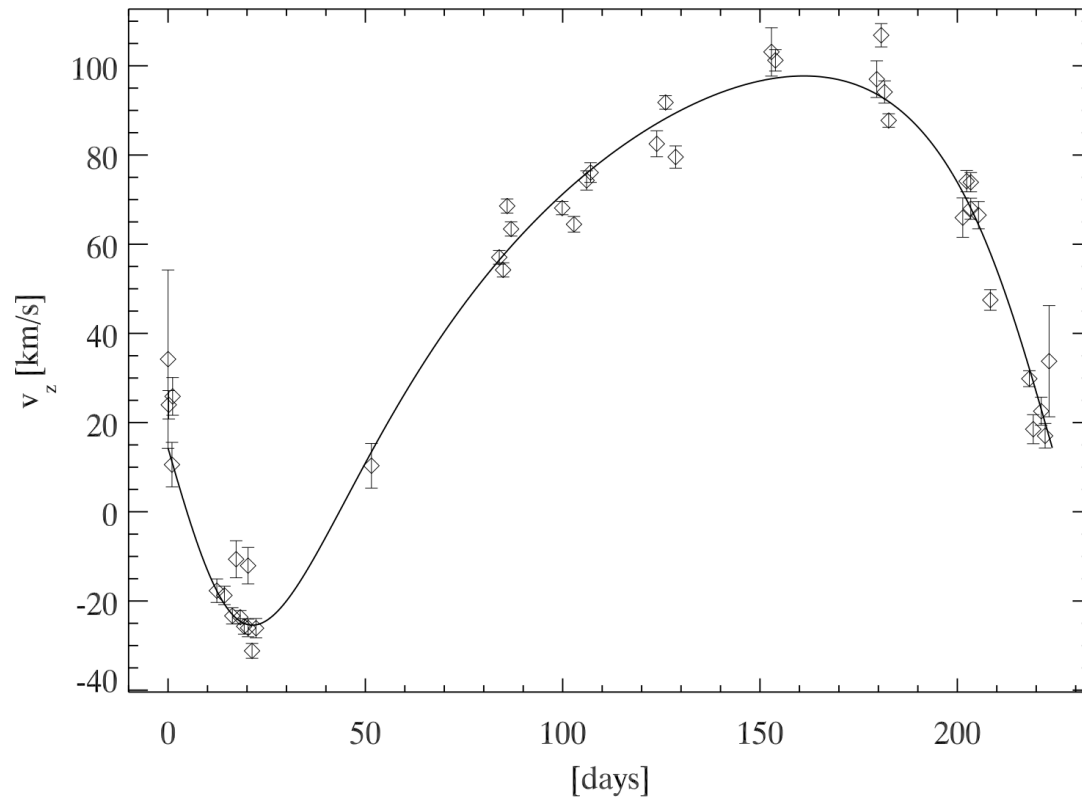


$$M = \frac{4\pi^2}{G} \frac{a^3}{T^2}$$

Most successful route:
Radial velocity measurements using the Doppler effect

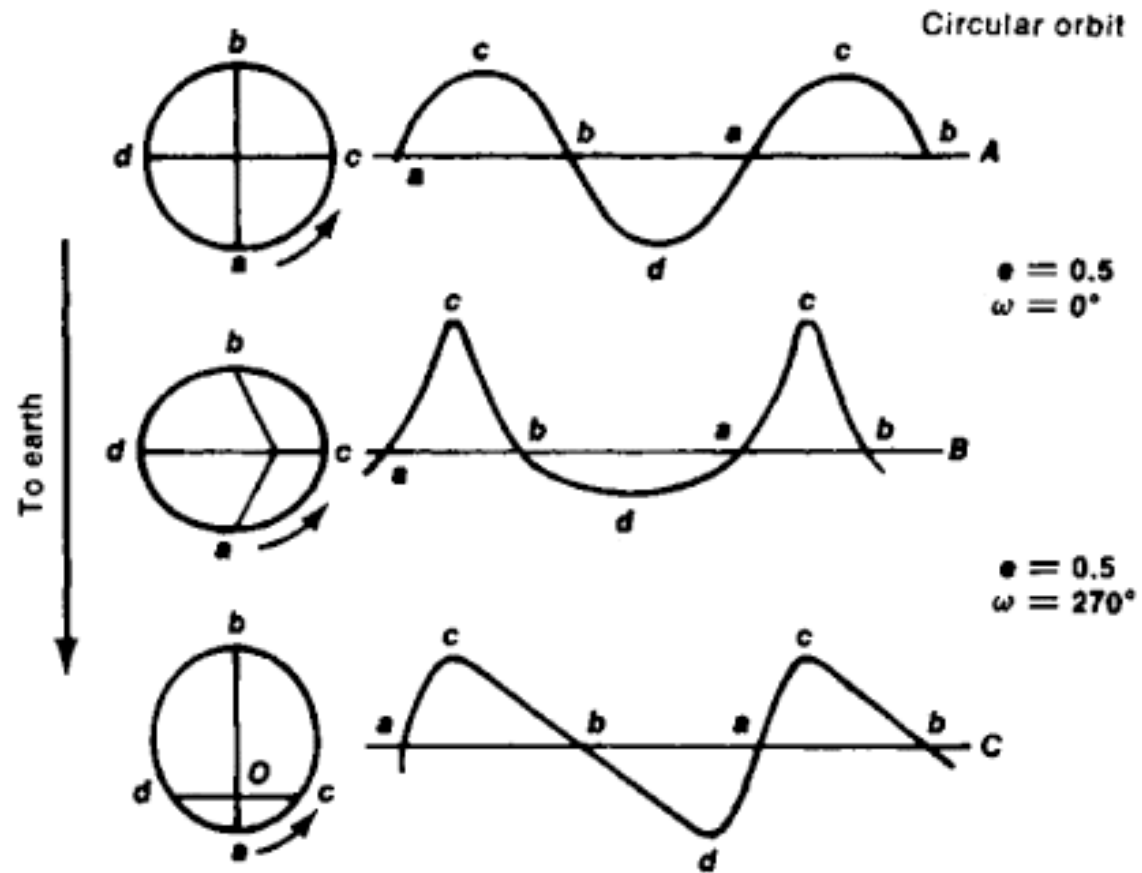


A typical example



Pfuhl et al. 2014

Shape and orientation of the orbit matter



Determine the orbital elements

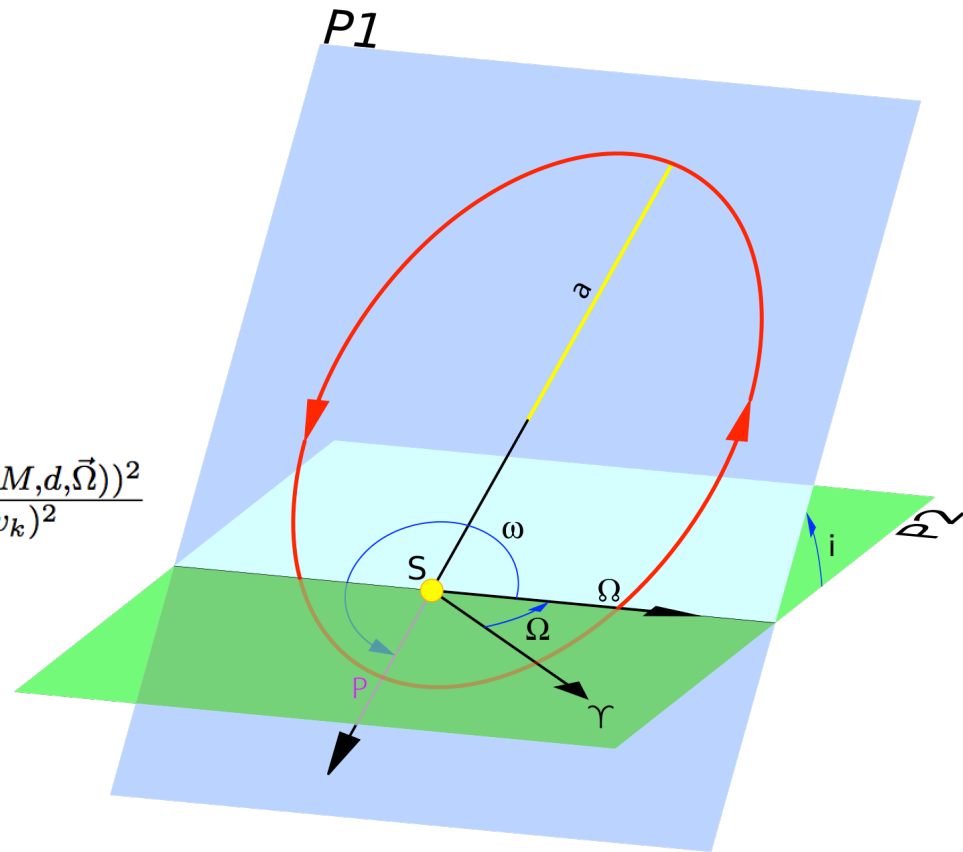
$$\{v_k \pm \Delta v_k\}$$

$$\vec{\Omega} = \{a, e, i, \Omega, \omega, t_0\}$$

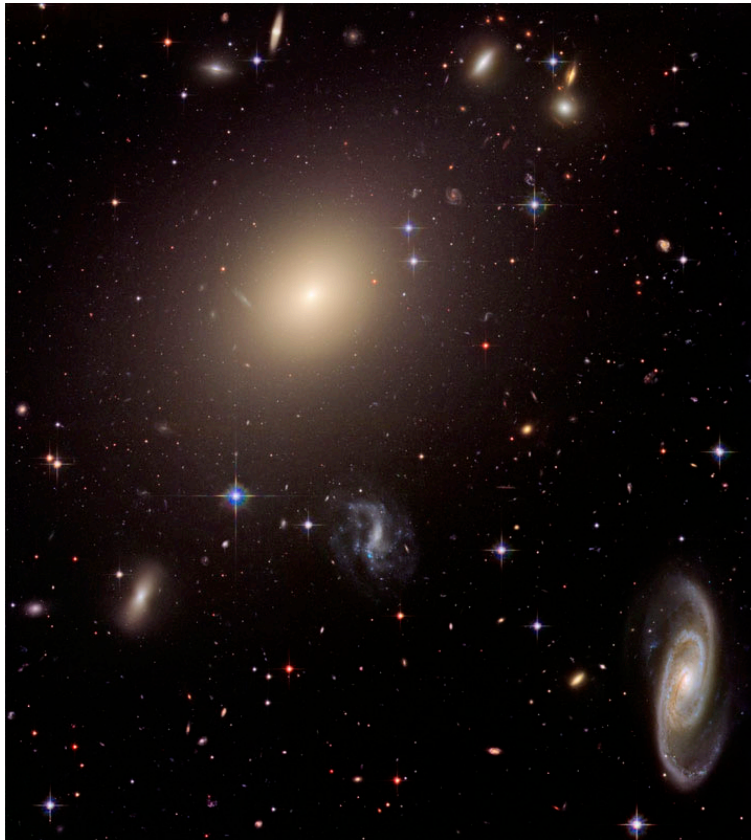
$$v(t) = v(t; M, d, \vec{\Omega})$$

$$\chi^2(M, d, \vec{\Omega}) = \sum_{k=1}^{\max} \frac{(v_k - v(t; M, d, \vec{\Omega}))^2}{(\Delta v_k)^2}$$

$$\min_{(M, d, \vec{\Omega})} \chi^2(M, d, \vec{\Omega})$$



More difficult: Unresolved stellar system



You can measure:

- Projected light profile
→ proxy for $n(R)$
- Radial velocity profile
→ $v(R)$, $\sigma(R)$

Solution: Jeans modeling

Distribution function: $f = f(\vec{x}, \vec{v}, t)$

“How many stars are at any given time
at a certain position with a certain velocity”

Number density of particles: $n(\vec{x}, t) = \int f(\vec{x}, \vec{v}, t) d^3v$

Mean velocity: $\langle v \rangle = \int v f d^3v$

Solution: Jeans modeling

Full differential:

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \frac{\partial f}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial f}{\partial v} \frac{\partial v}{\partial t}$$

Newton's law:

$$\frac{\partial v}{\partial t} = a = -\frac{\partial \Phi}{\partial x}$$

yields:

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + v \frac{\partial f}{\partial x} - \frac{\partial \Phi}{\partial x} \frac{\partial f}{\partial v}$$

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla f - \nabla \Phi \cdot \frac{\partial f}{\partial \vec{v}}$$

Solution: Jeans modeling

Collisionless, incompressible system:

$$\text{“CBE”} \quad 0 = \frac{\partial f}{\partial t} + \vec{v} \cdot \nabla f - \nabla \Phi \cdot \frac{\partial f}{\partial \vec{v}}$$

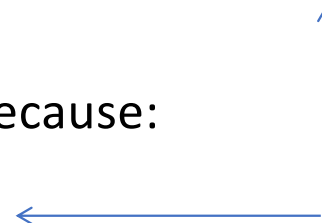
$$\int \text{CBE} d^3v \text{ yields continuity equation:} \quad 0 = \frac{\partial n}{\partial t} + \nabla \cdot (n \langle \vec{v} \rangle)$$

$\int v \text{ CBE} d^3v$ yields Jeans equation

$$-\nabla \Phi = \frac{\partial \langle \vec{v} \rangle}{\partial t} + \langle \vec{v}_i \rangle \frac{\partial \langle \vec{v} \rangle}{\partial x_i} + \frac{1}{n} \nabla (n \sigma^2)$$

Velocity dispersion appears here, because:

$$\langle v^2 \rangle = \langle v \rangle^2 + \sigma^2$$



Solution: Jeans modeling

Simplest case: steady state, spherical symmetric, non-rotating system:

The potential is Newtonian $\Phi(r) = \frac{G M(r)}{r}$

Only radial part of Jeans equations survives, and many terms are 0:

$$-\frac{d\Phi}{dr} = \frac{1}{n} \frac{d(n\sigma^2)}{dr}$$

Integrating:

$$n(r)\sigma^2(r) = \int_r^\infty n(r') \frac{d\Phi(r')}{dr'} dr'$$

Solution: Jeans modeling

Slight complication: We measure **projected** values only

$$r^2 = R^2 + z^2 \quad \Sigma(R) = \int_{-\infty}^{\infty} n(R, z) dz$$

change of variables $z \rightarrow r$

$$\Sigma(R) = \int_R^{\infty} \frac{r}{\sqrt{r^2 - R^2}} n(r) dr$$

Similar for velocity dispersion

$$\Sigma(R) \sigma_P^2(R) = \int_R^{\infty} \frac{r}{\sqrt{r^2 - R^2}} n(r) \sigma^2(r) dr$$

Some maths ...

$$\sigma_P^2(R) = G \frac{\int_R^{\infty} (r^2 - R^2)^{1/2} r^{-2} M(r) n(r) dr}{\int_R^{\infty} (r^2 - R^2)^{-1/2} r n(r) dr}$$

Now we can “Jeans model”

1. Define your model

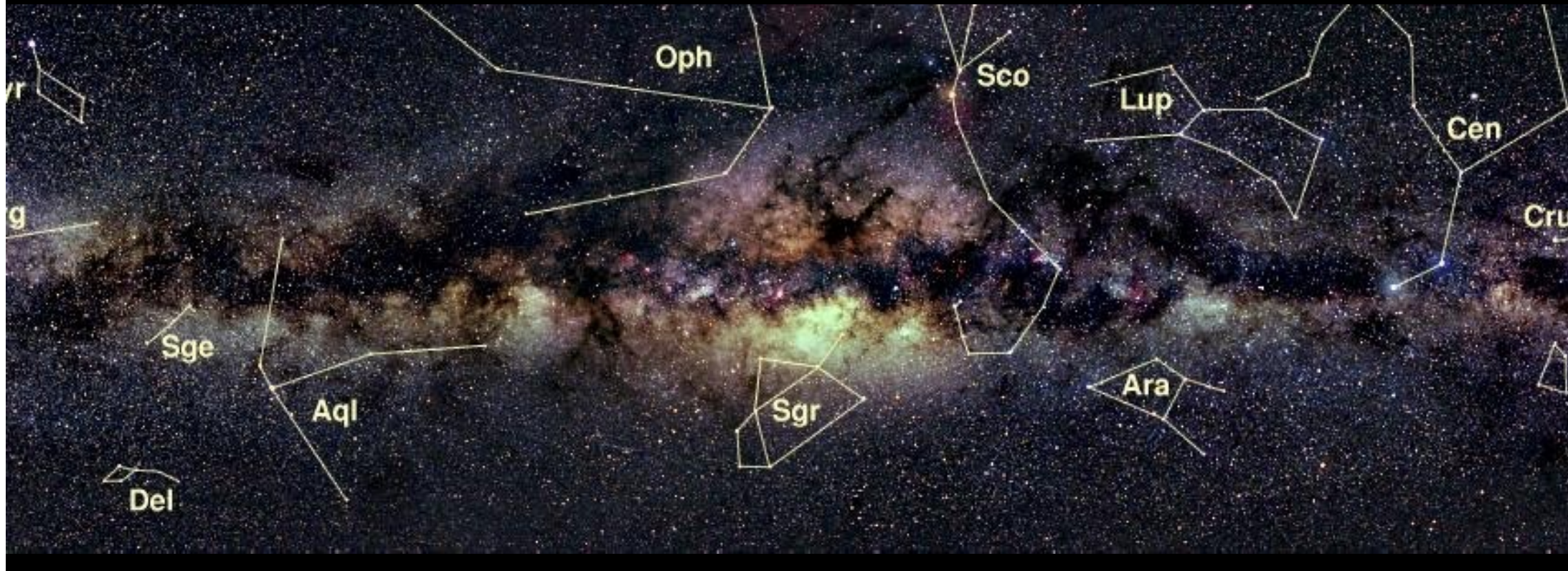
$$M(r) = M_{\text{central}} + 4\pi \int_0^r r^2 \rho(r) dr$$
$$\rho(r) = \rho_0 r^\gamma$$
$$n(r) \propto r^\Gamma$$

1. Determine $n(r)$ from $\Sigma(R)$,
i.e. fit Γ to reproduce measured density profile
2. Remaining free 3 parameters: $M_{\text{central}}, \rho_0, \gamma$.
3. For each set of these 3, one can evaluate $\sigma_p^2(R)$
4. Minimize χ^2 to measured values of $\sigma_p(R_k)$
5. You have determined $M(r)$!

Outline

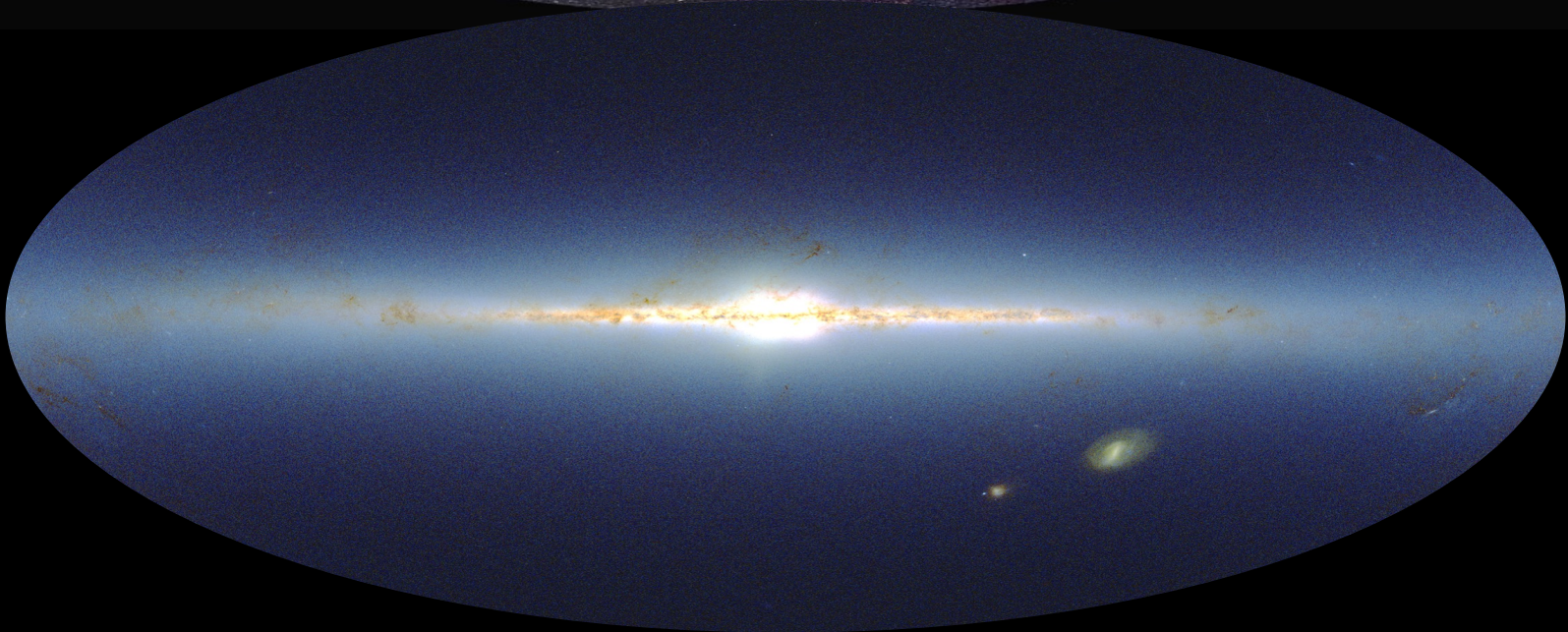
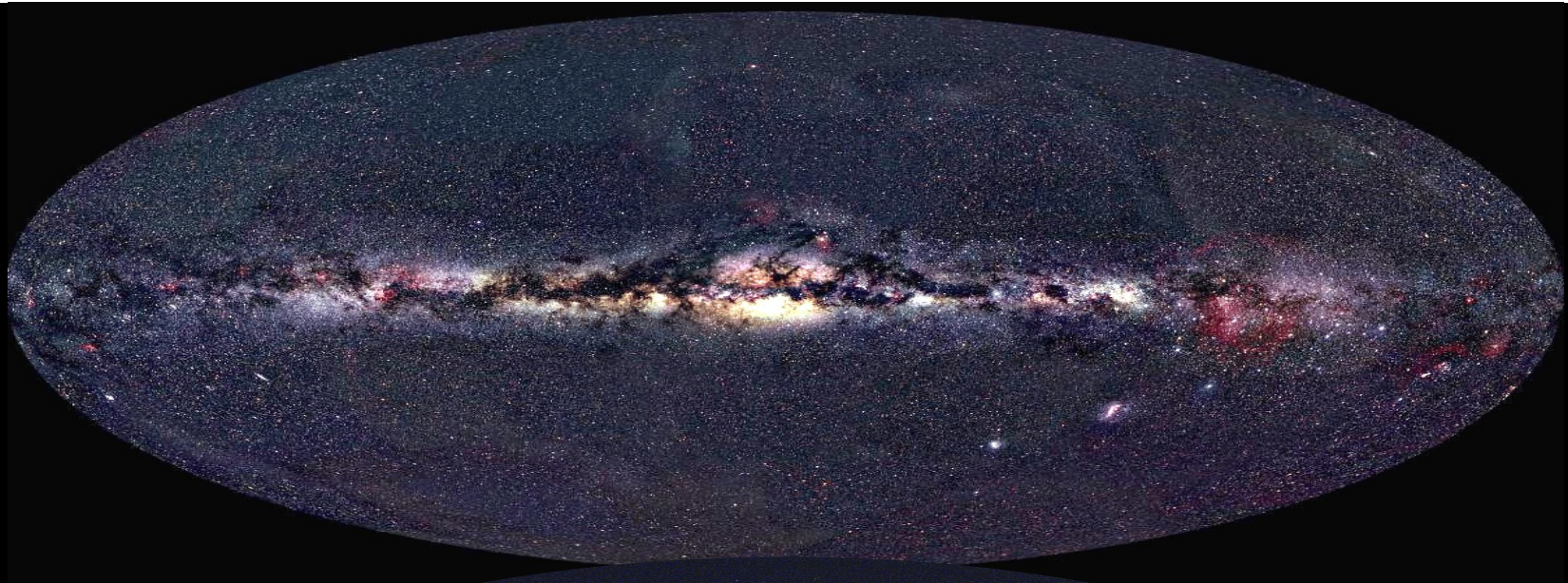
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The Galactic Center is highly obscured

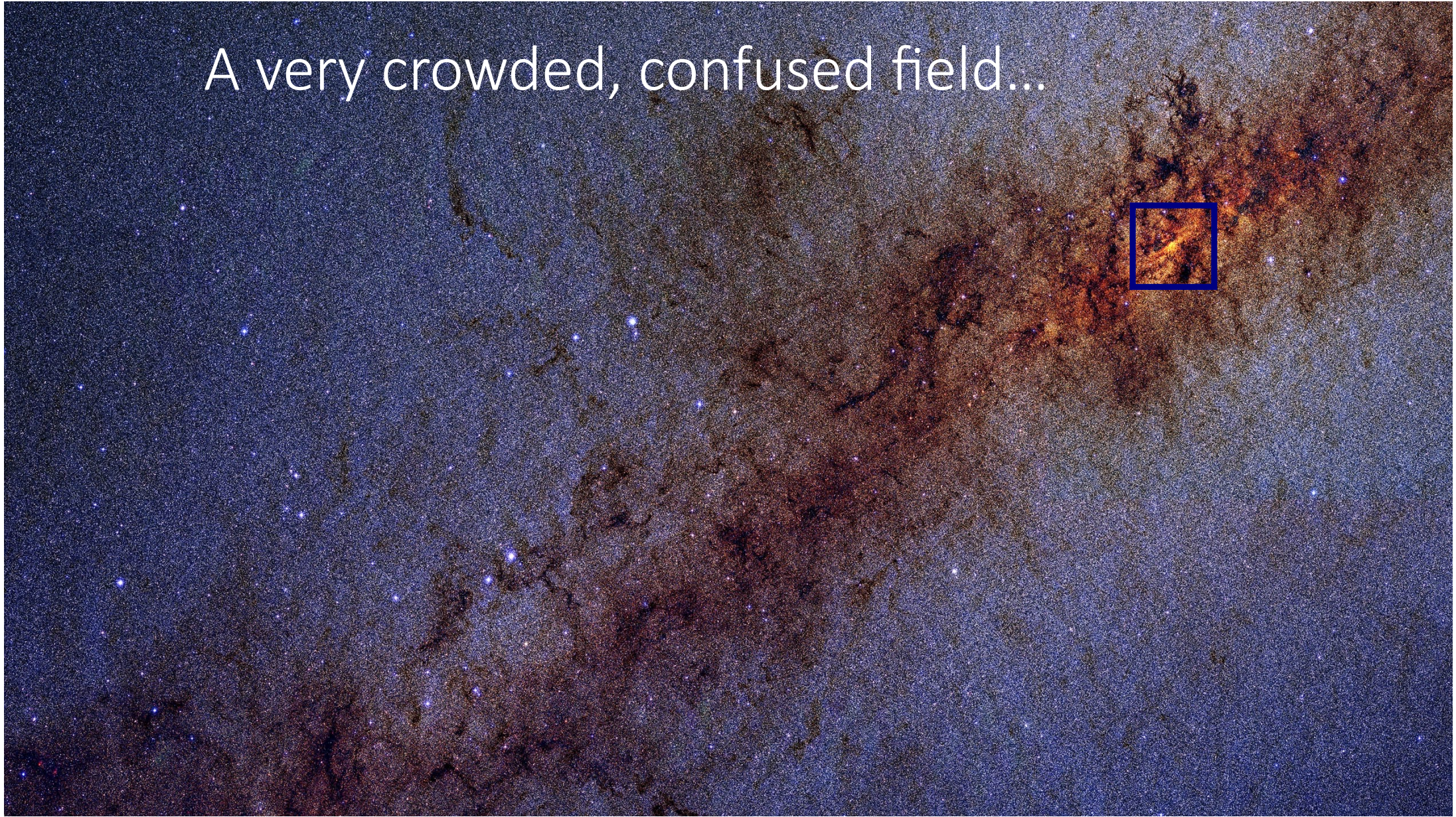


Using infrared light allows one to look through the dust

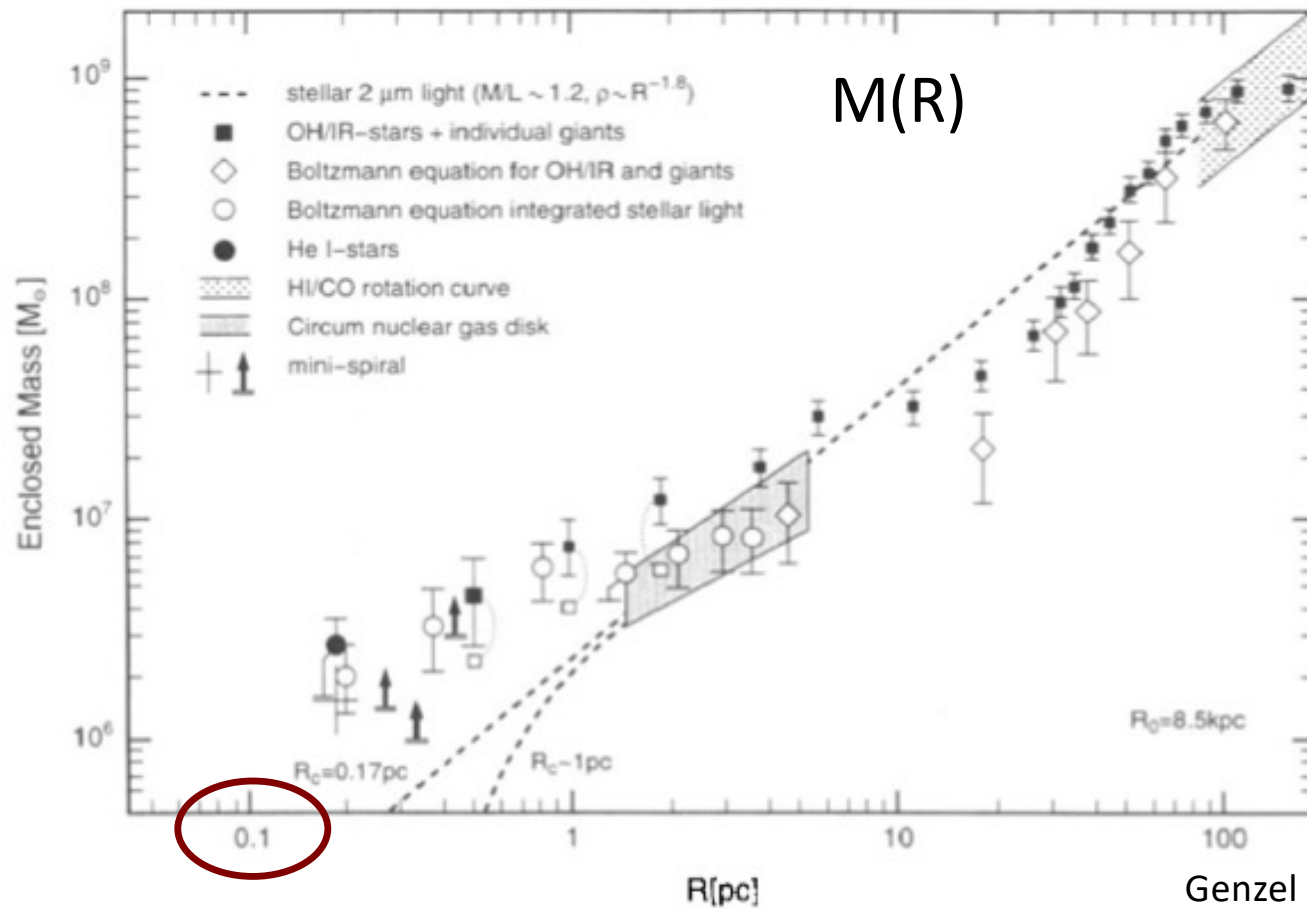




A very crowded, confused field...

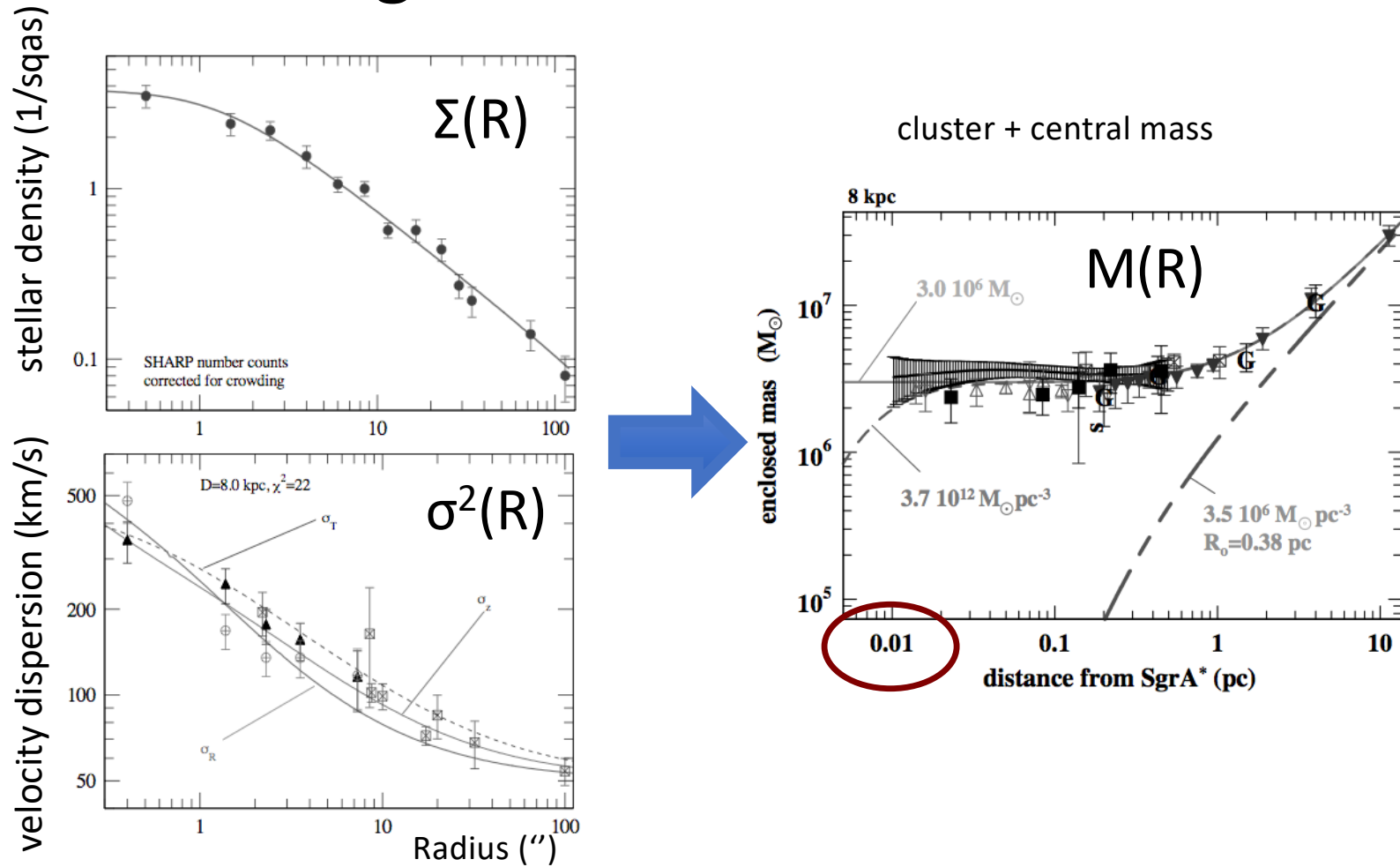


Jeans modeling in the Galactic Center – 1990's



Genzel et al. 1994

Jeans modeling in the Galactic Center – 2000

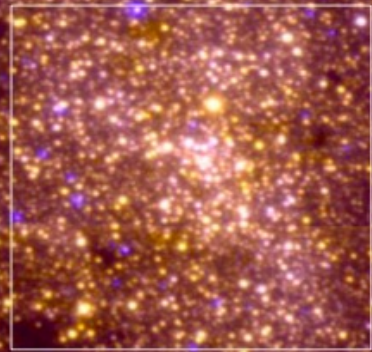


Genzel et al. 2000

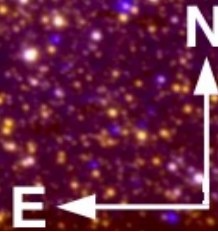
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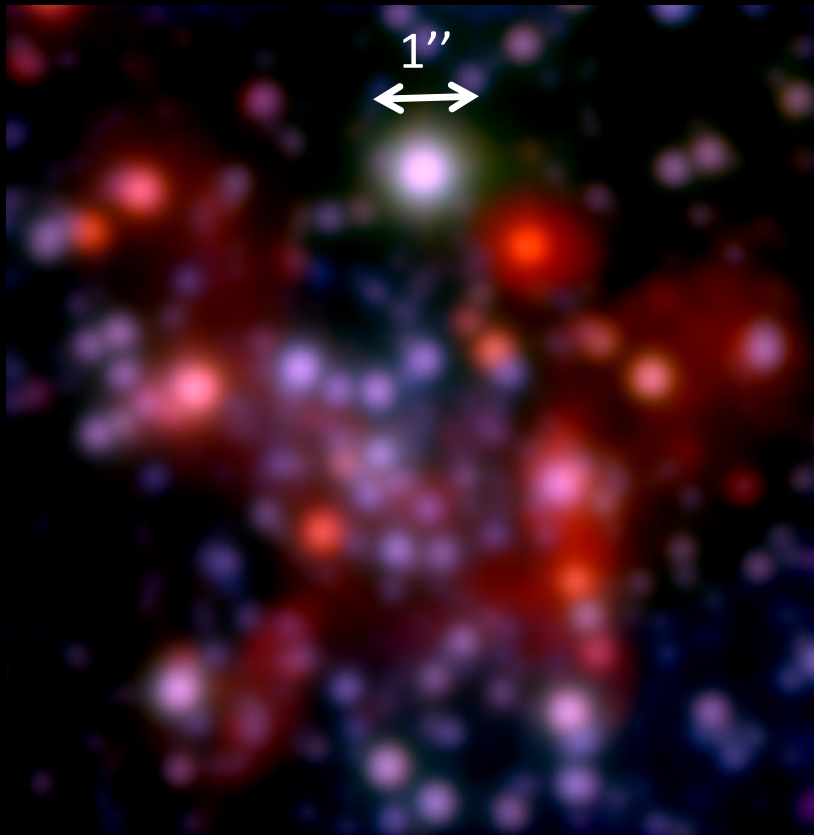
The Galactic Center contains an extremely dense star cluster



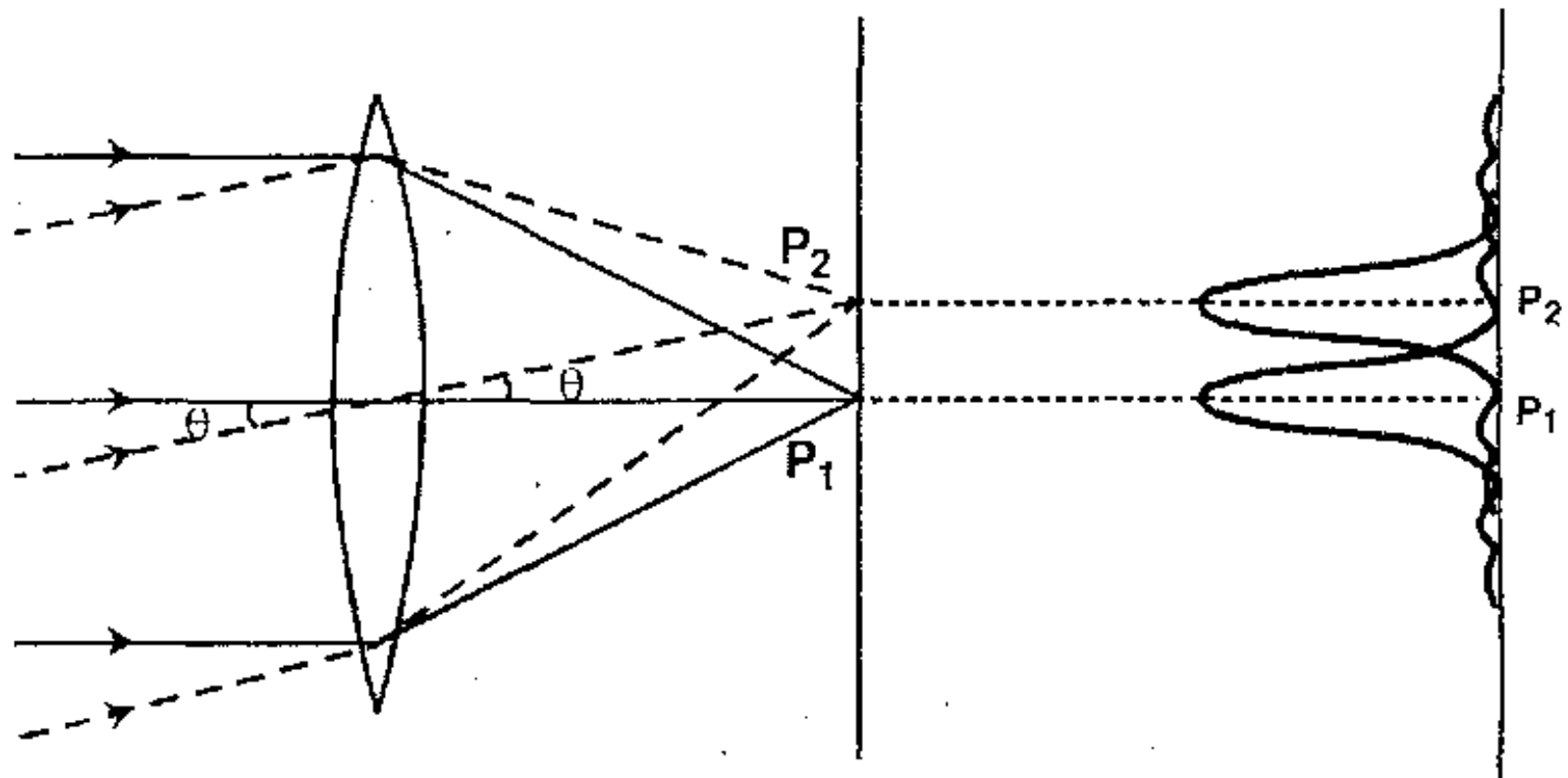
30'' = 4 lightyears



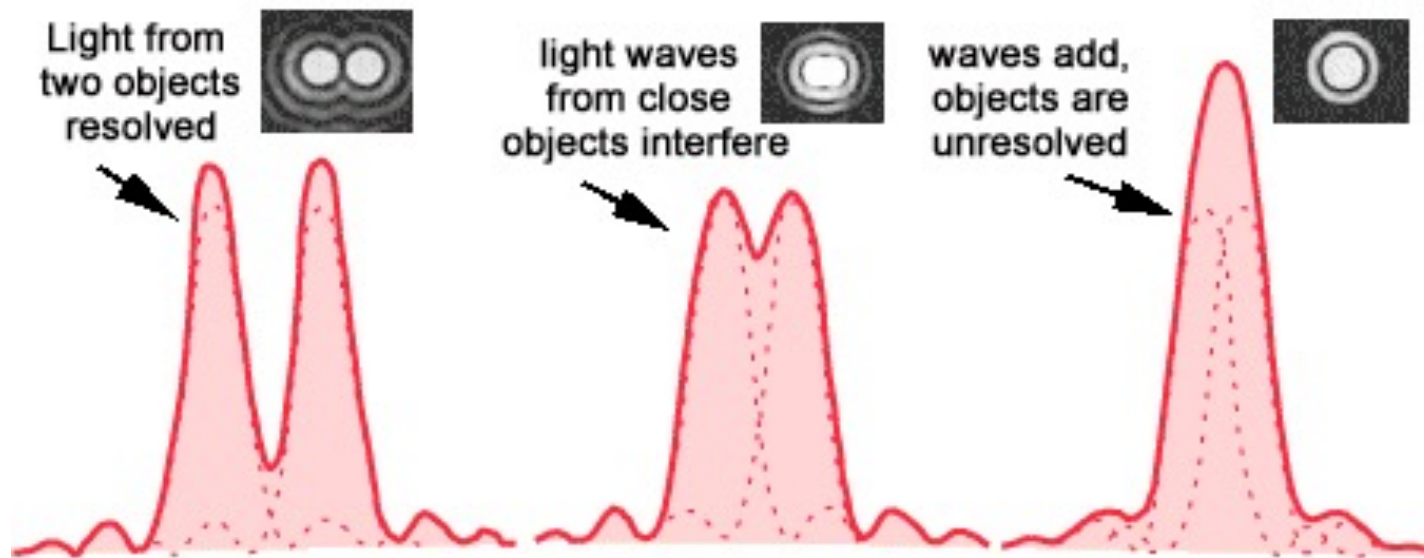
The central 20'' : Not very sharp



Why are the point-like stars not images as points?



Objects are resolved when they are further apart than $1.22 \times \lambda / D$



Rayleigh criterion: $\theta > 1.22 \times \lambda / D$

$\lambda = 2\mu\text{m}$; $D = 8\text{m}$ should be: $\theta = 0.05''$

but we observe $\theta = 1''$?!

The "Seeing" smears out the images



That is severely limiting

Light wavefronts are aberrated

Wavefront error
for telescope diameter D :

$$\sigma^2 = 1.030 \left(\frac{D}{r_0} \right)^{5/3}$$

complete decorrelation: $\sigma^2 = 1 \text{ rad}$ $\rightarrow D \approx r_0$

Turbulence limits the resolution of a telescope to an effective size of r_0 .

$R \sim \lambda/r_0$ instead of $R \sim \lambda/D$.

$$\lambda = 2.2 \mu\text{m}, \theta = 1''$$



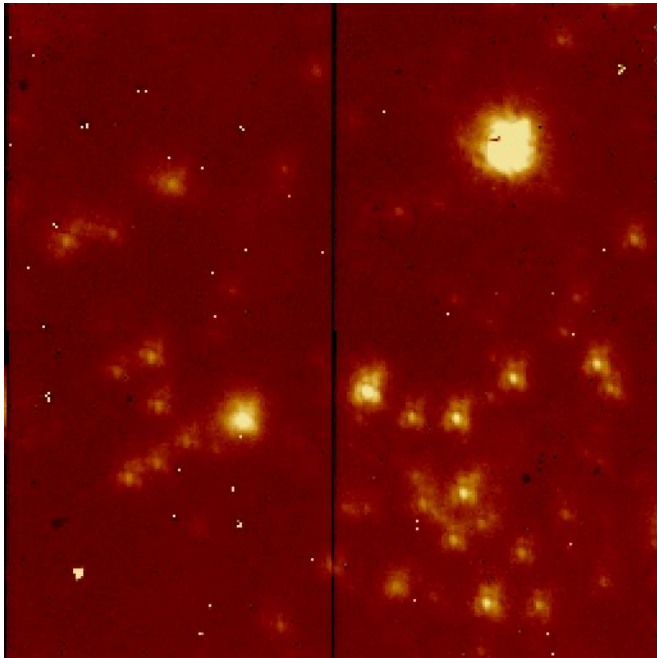
$$r_0 = 40 \text{ cm}$$

La Silla, Chile

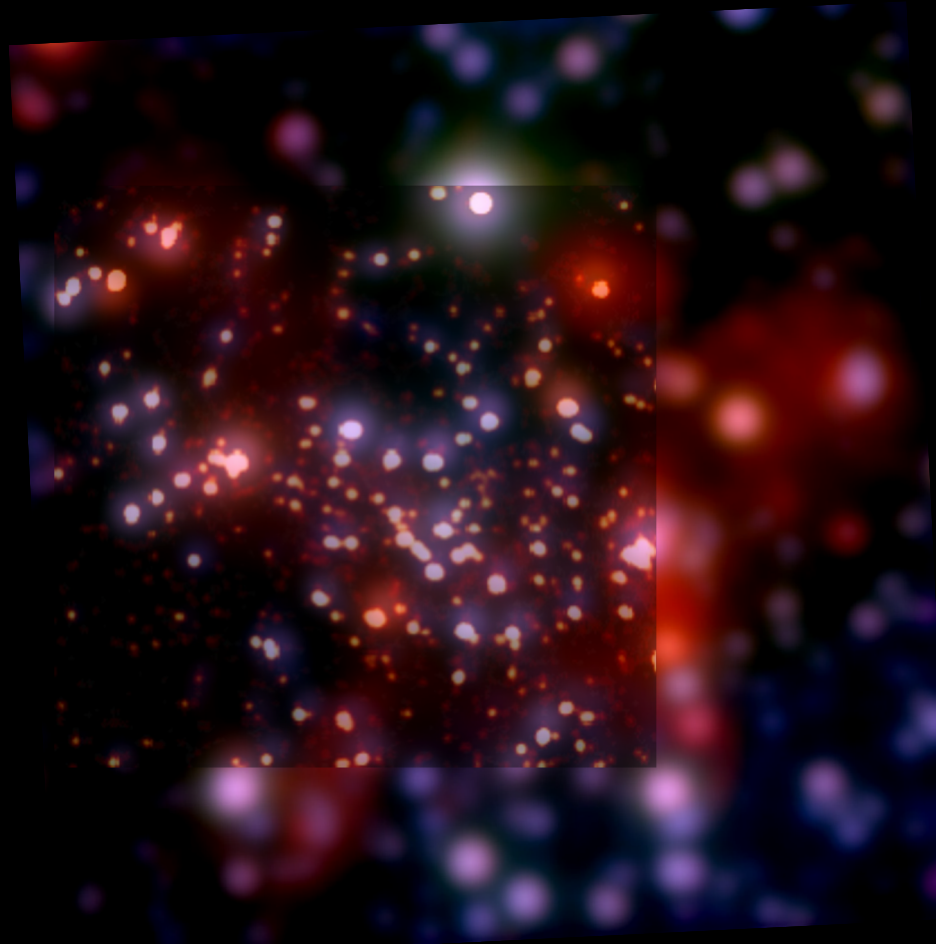


© Serge Brunier

1992 - 2001: Speckle-Imaging



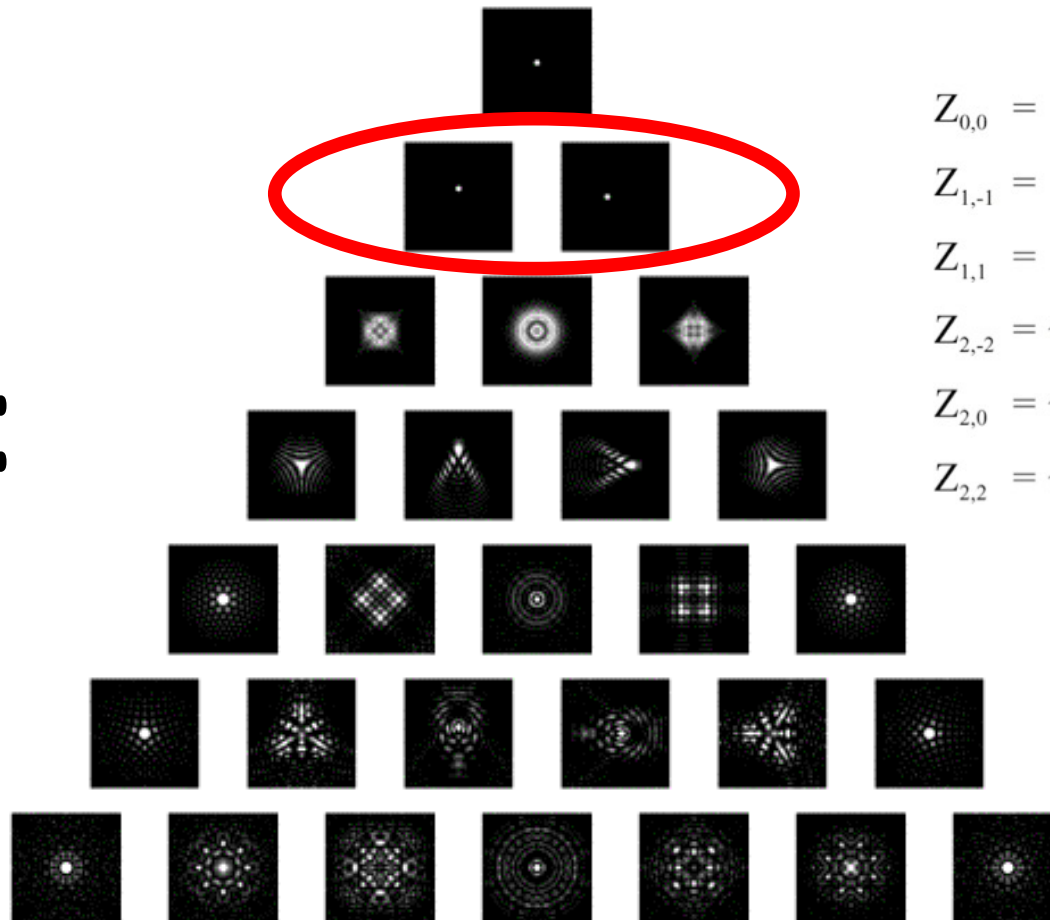
From seeing-limited to to diffraction-limited



“Simple shift and add” only corrects two terms of wavefront aberrations



=



$$Z_{0,0} = 1$$

piston

$$Z_{1,-1} = 2 r \sin\theta$$

$$Z_{1,1} = 2 r \cos\theta$$

tip/tilt

$$Z_{2,-2} = \sqrt{6} r^2 \sin 2\theta$$

astigmatism

$$Z_{2,0} = \sqrt{3} (2r^2 - 1)$$

focus

$$Z_{2,2} = \sqrt{6} r^2 \cos 2\theta$$

astigmatism

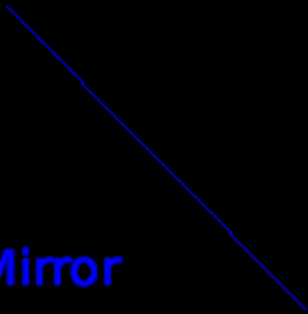
“Zernike” polynomials



Atmosphere



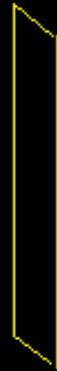
Mirror



Lens



Detector

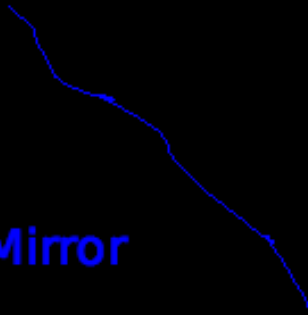




Atmosphere



Mirror



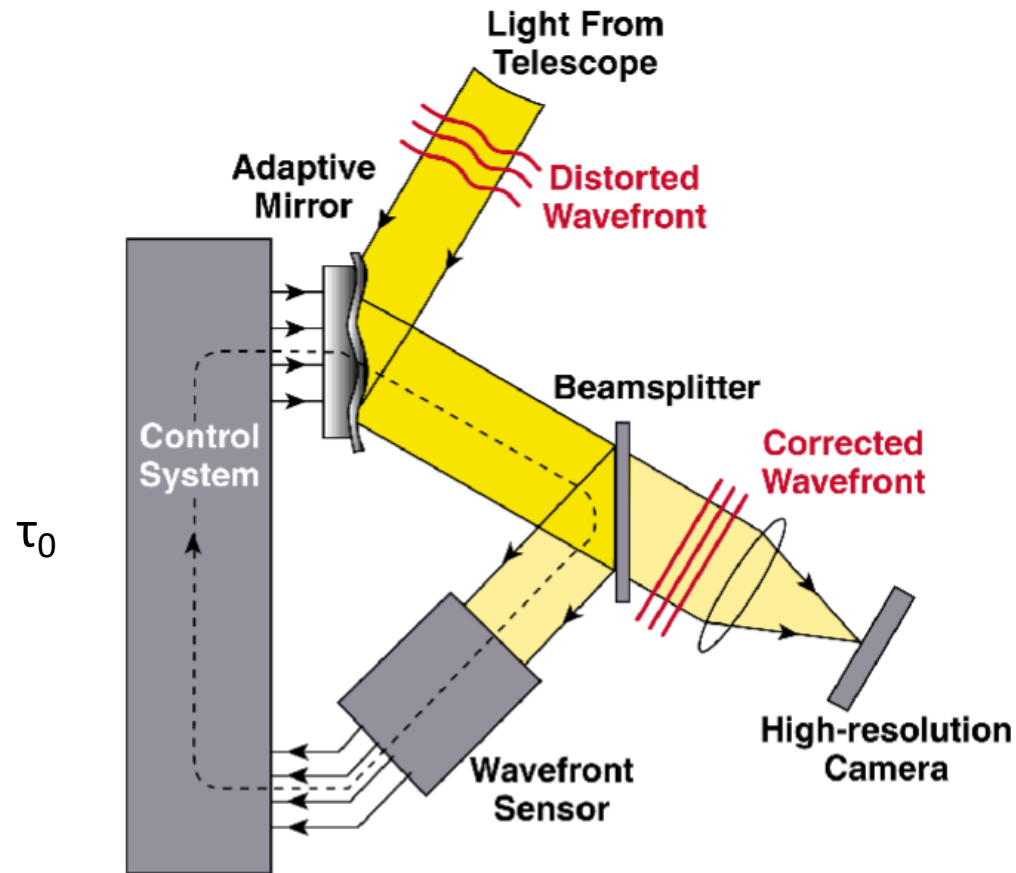
Lens



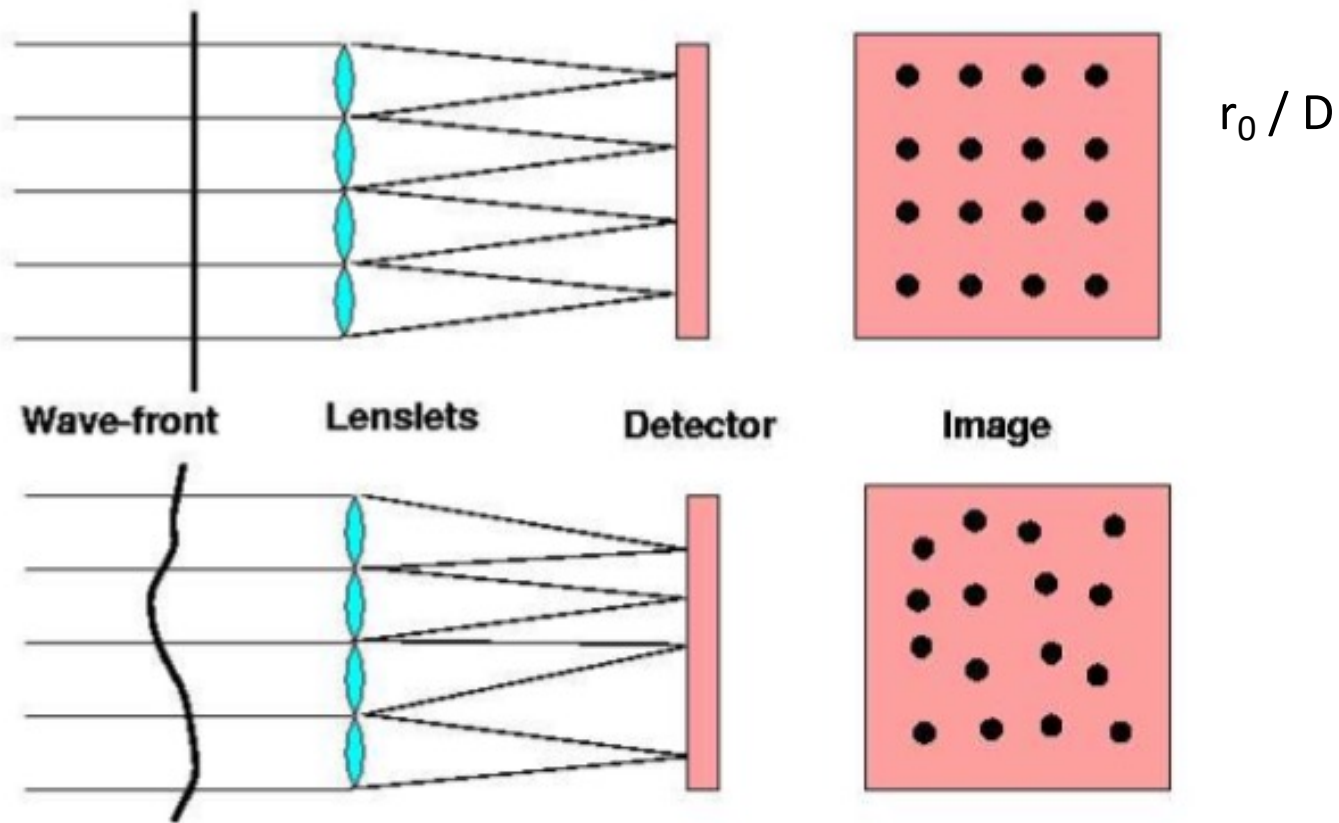
Detector



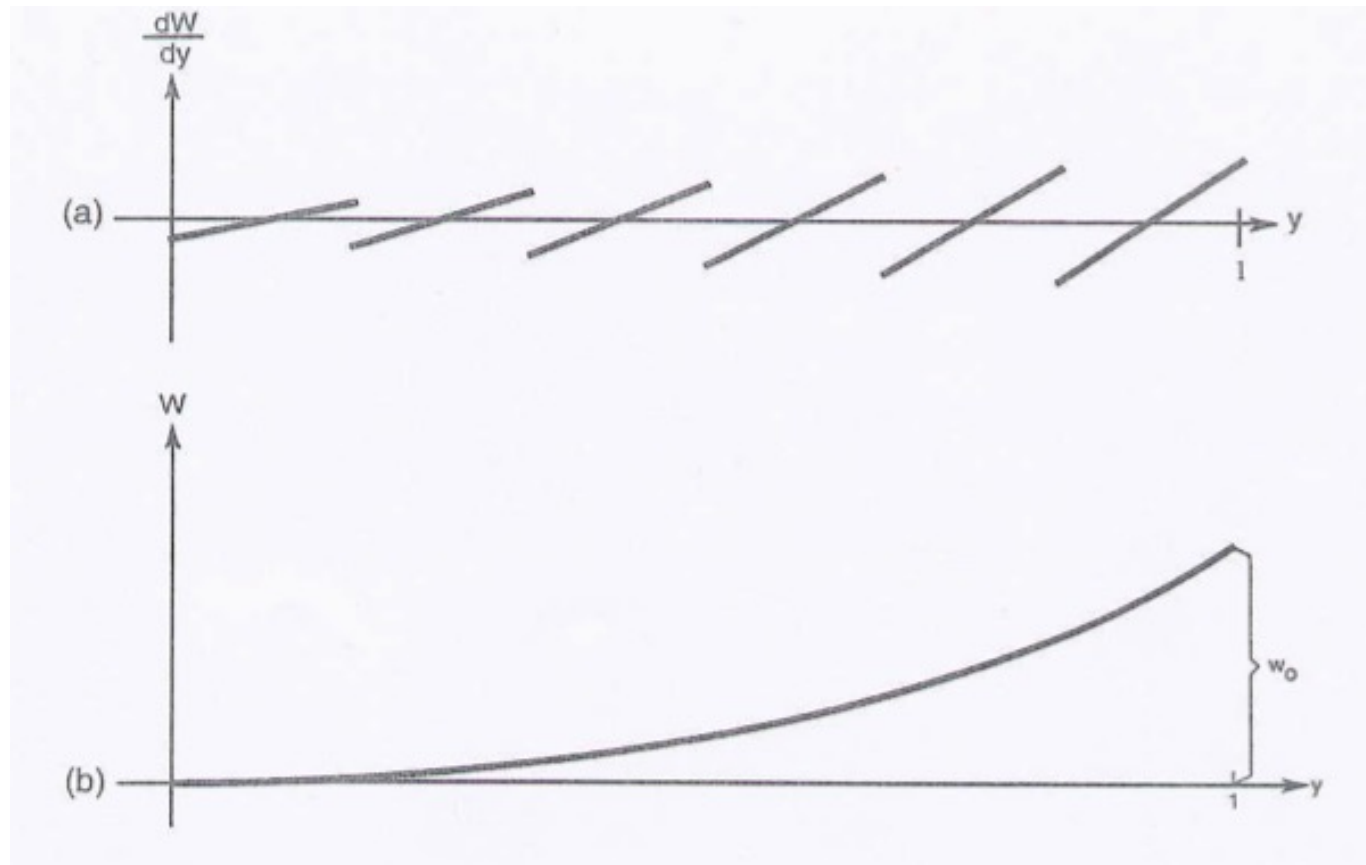
We need... An Adaptive Optics system



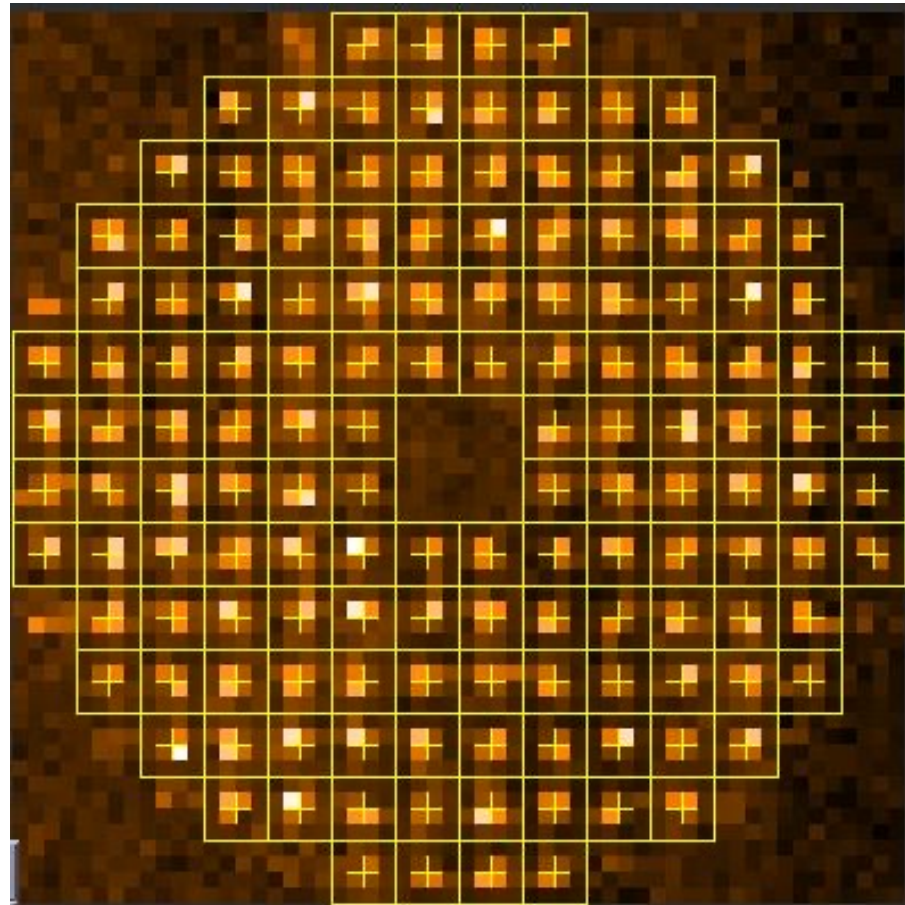
Shack-Hartmann-WFS: samples **slope** of wave front



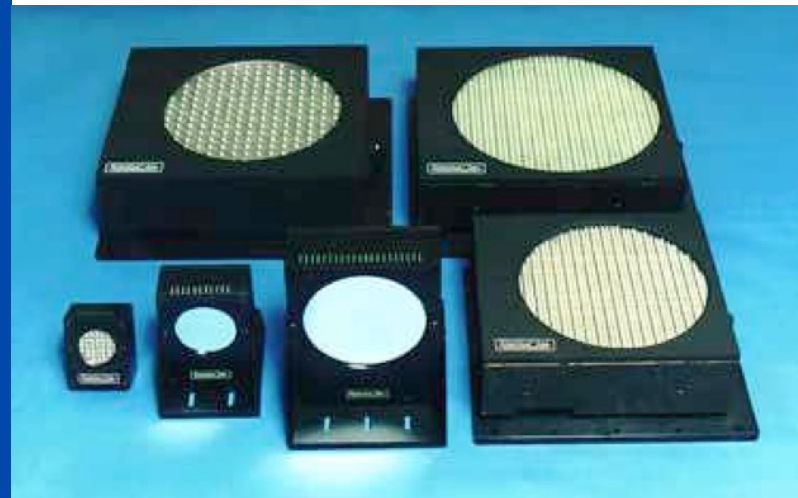
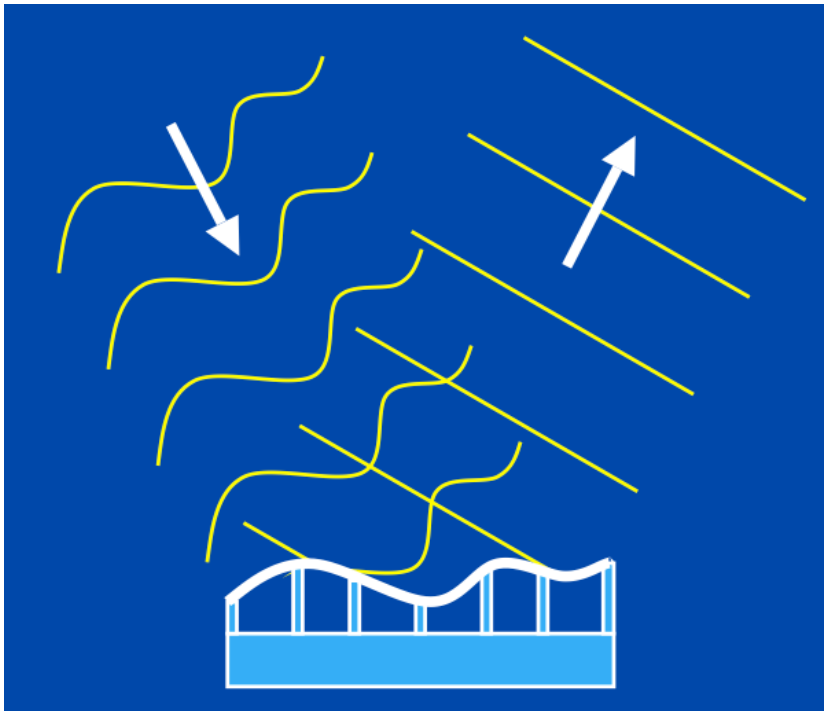
From slopes to the wavefront



A real world detector: NACO-WFS



Deformable Mirrors



Paranal, Chile



NACO & SINFONI: For almost 20 years the Adaptive optics instruments at the VLT

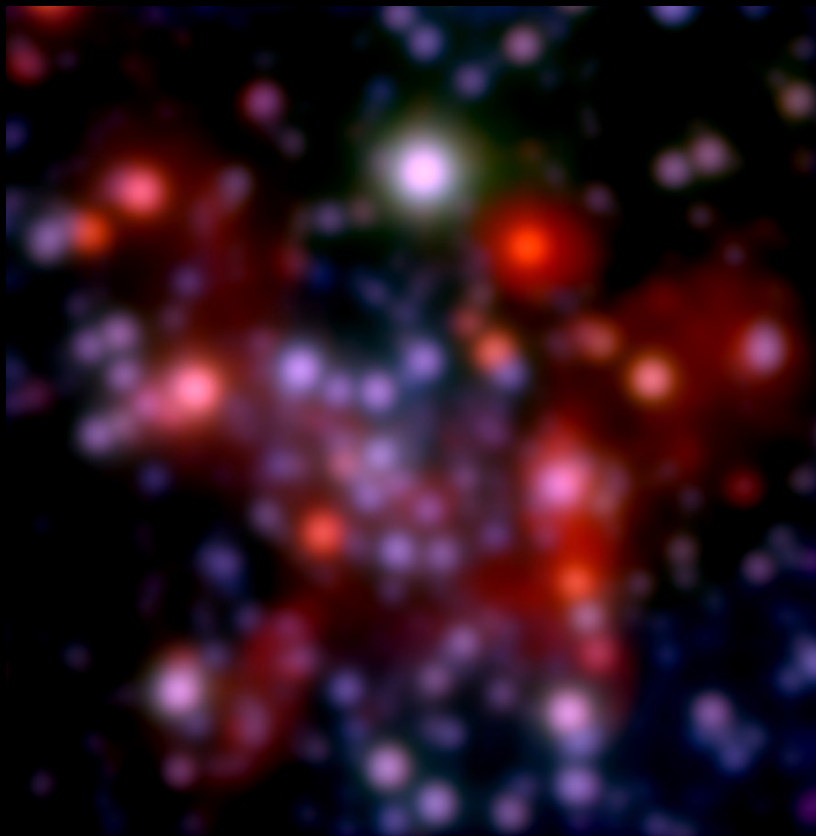


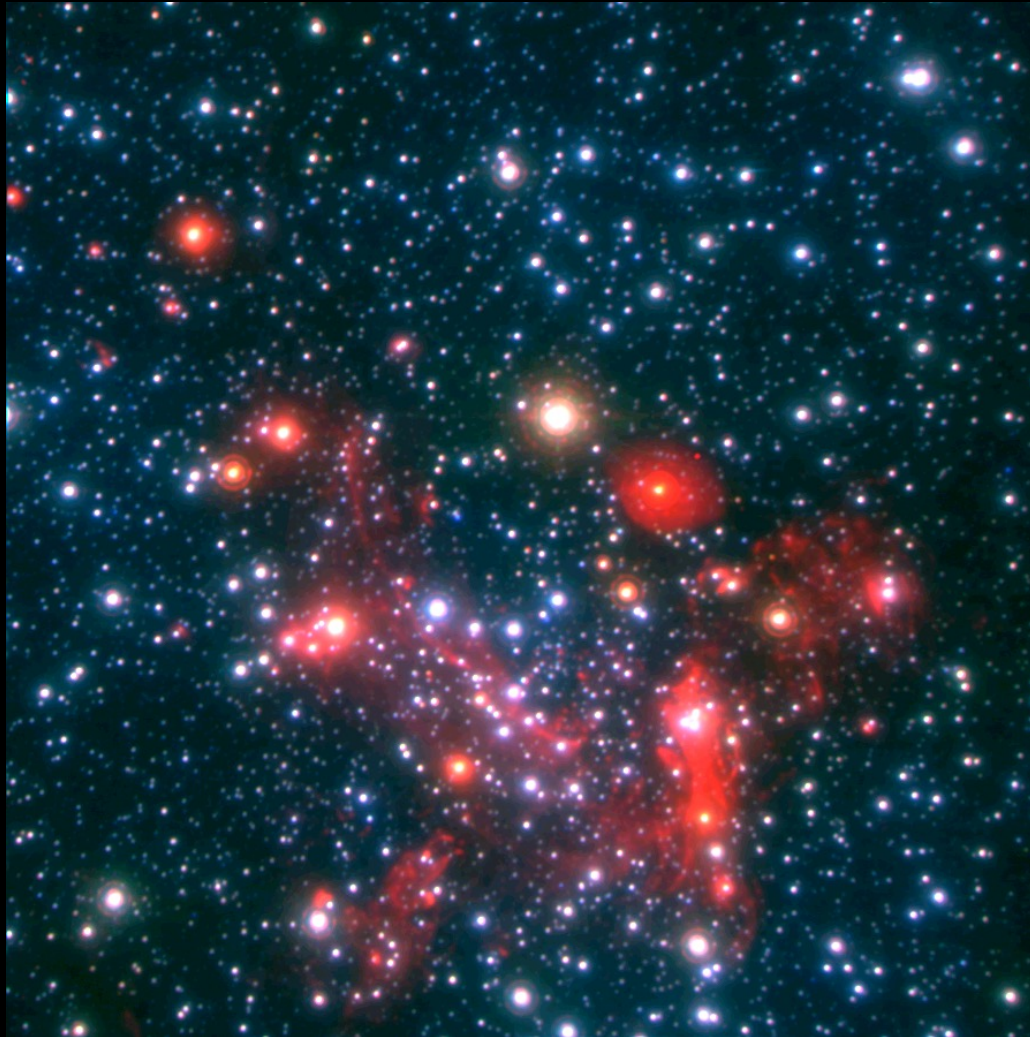
NACO (2002 – 2019):
Astrometry with $300 \mu\text{s}$



SINFONI (2003 – 2019):
Spectroscopy with 7 km/s

Really a big step forward: AO



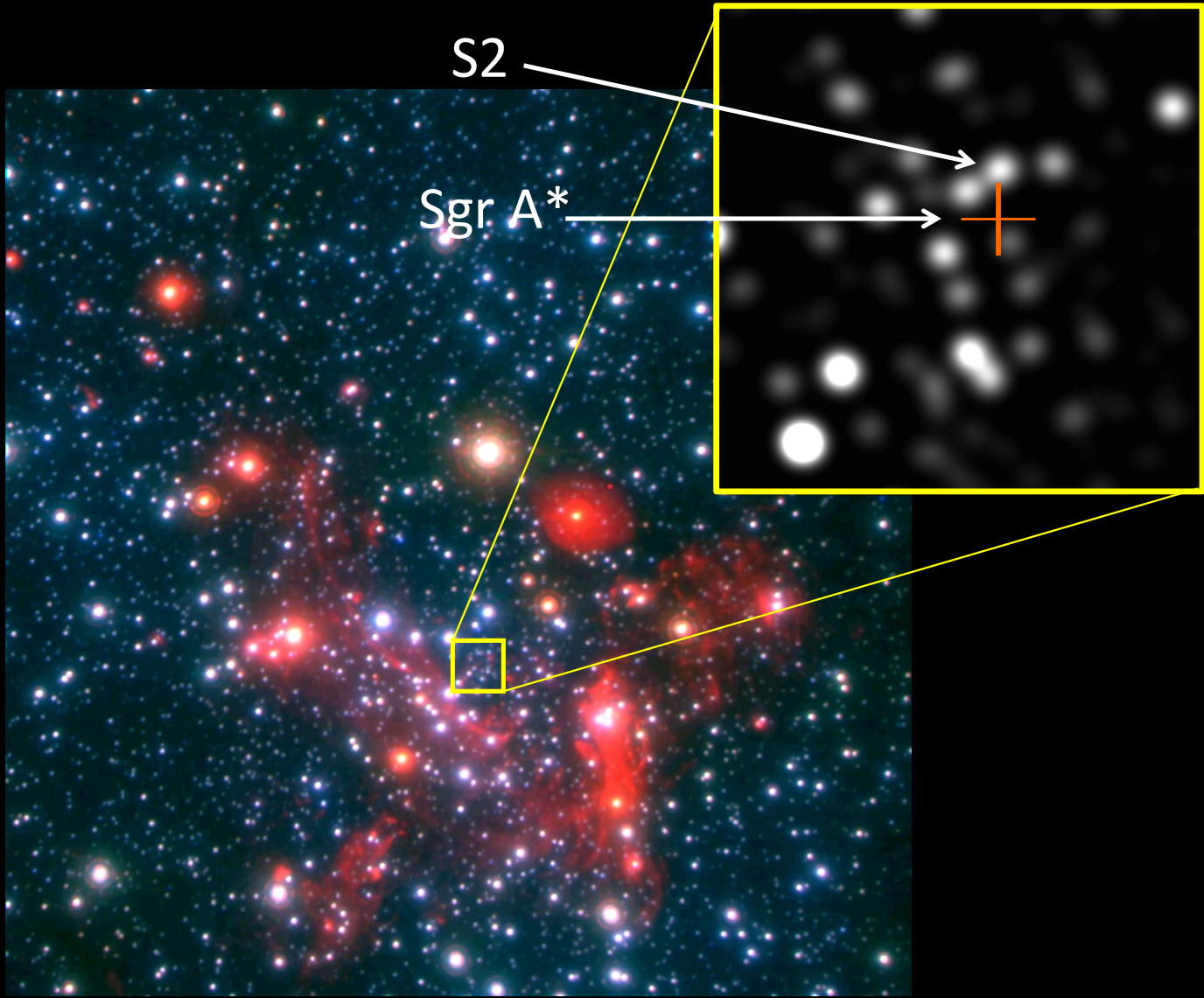


Strehl ratio
40%

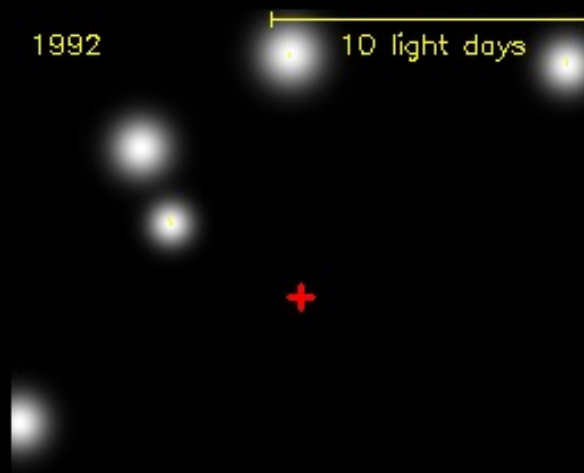
NACO,
HKL color composite

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The stars move on Keplerian ellipses



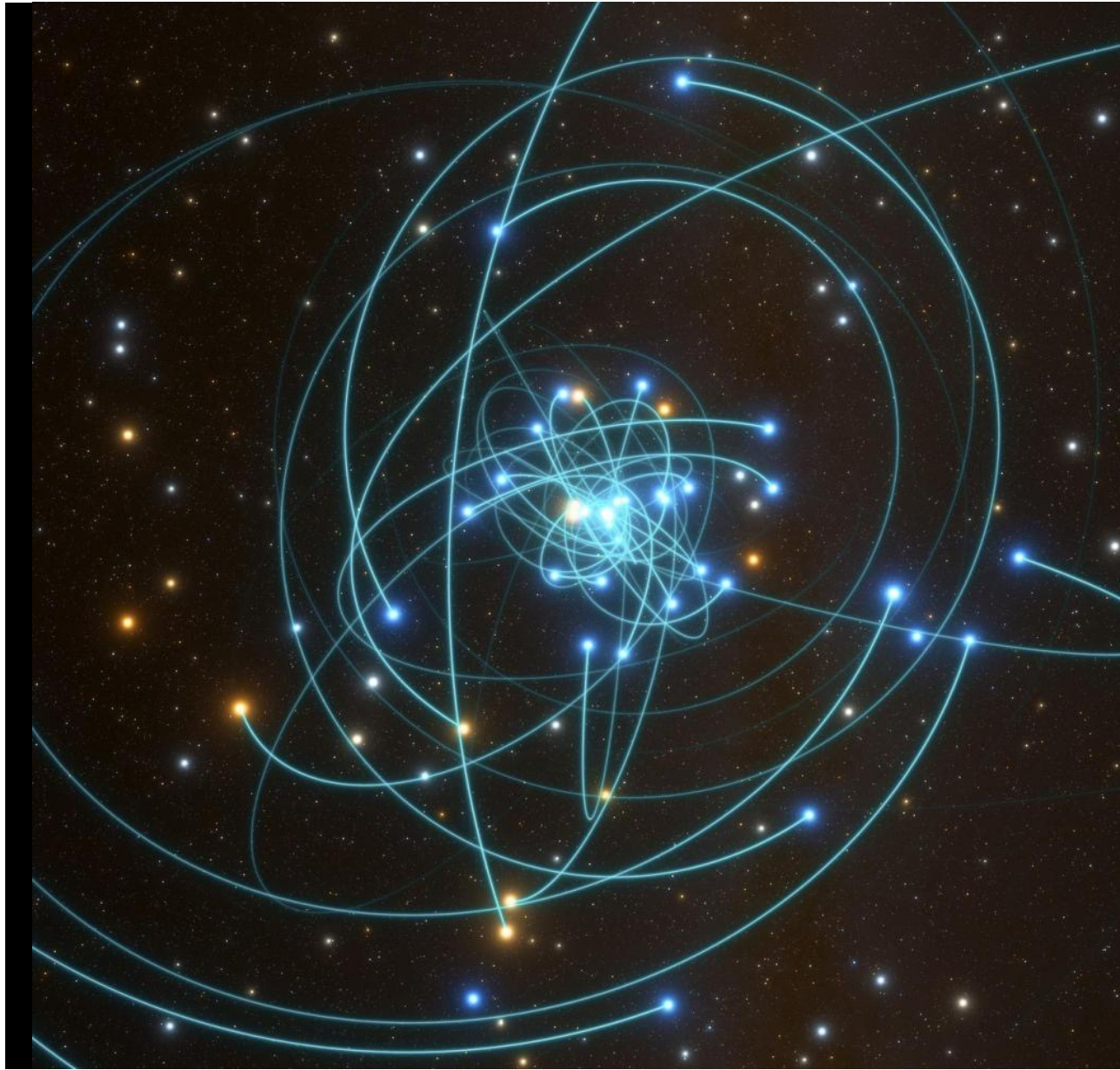
Resulting mass:

$$M = \frac{4\pi^2 a^3}{G P^2}$$

$$M = \frac{4\pi^2}{6.67 \times 10^{-11} \frac{\text{Nm}^2}{\text{kg}^2}} \frac{(125'' \times 8.3 \text{ kpc})^3}{(16 \text{ yr})^2} =$$

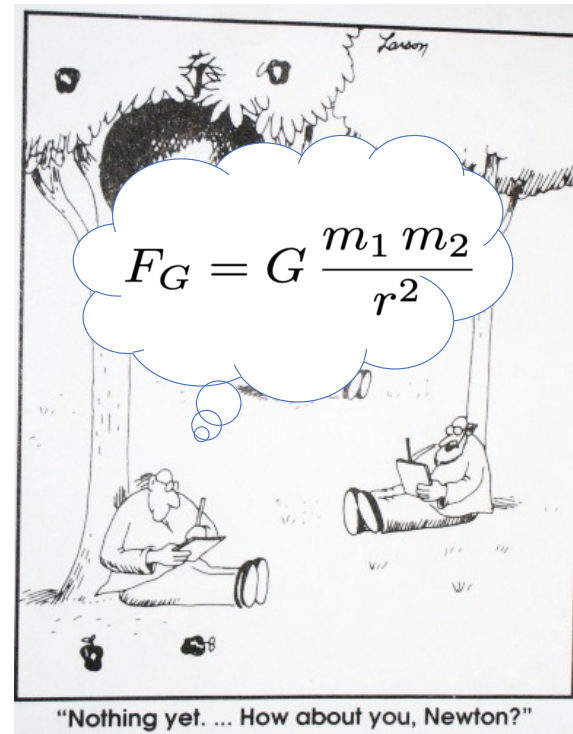
$$4.3 \times 10^6 M_{\odot}$$

WOW!

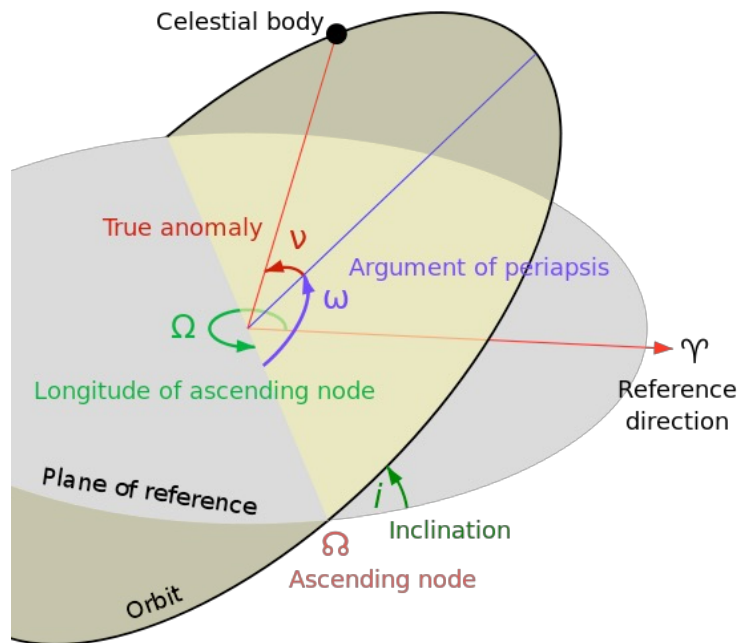


Currently known:
around 50 orbits

Newton would have understood this
but he had no chance to see it



How to measure stellar orbits



- 6 orbital elements:
(a , e , i , ω , Ω , t)
- Need 6 dynamical quantities
- more if information on potential needs to be inferred
- Imaging data:
 x , y , v_x , v_y
- Spectroscopy:
 v_z
- At least one more number needed:
an acceleration
 - a_{2D} from imaging
 - $a_z = dv_z / dt$ from spectroscopy

Determine the orbital elements and black hole parameters

$$\{x_j \pm \Delta x_j, y_j \pm \Delta y_j, v_k \pm \Delta v_k\}$$

$$\vec{\Omega} = \{a, e, i, \Omega, \omega, t_0\}$$

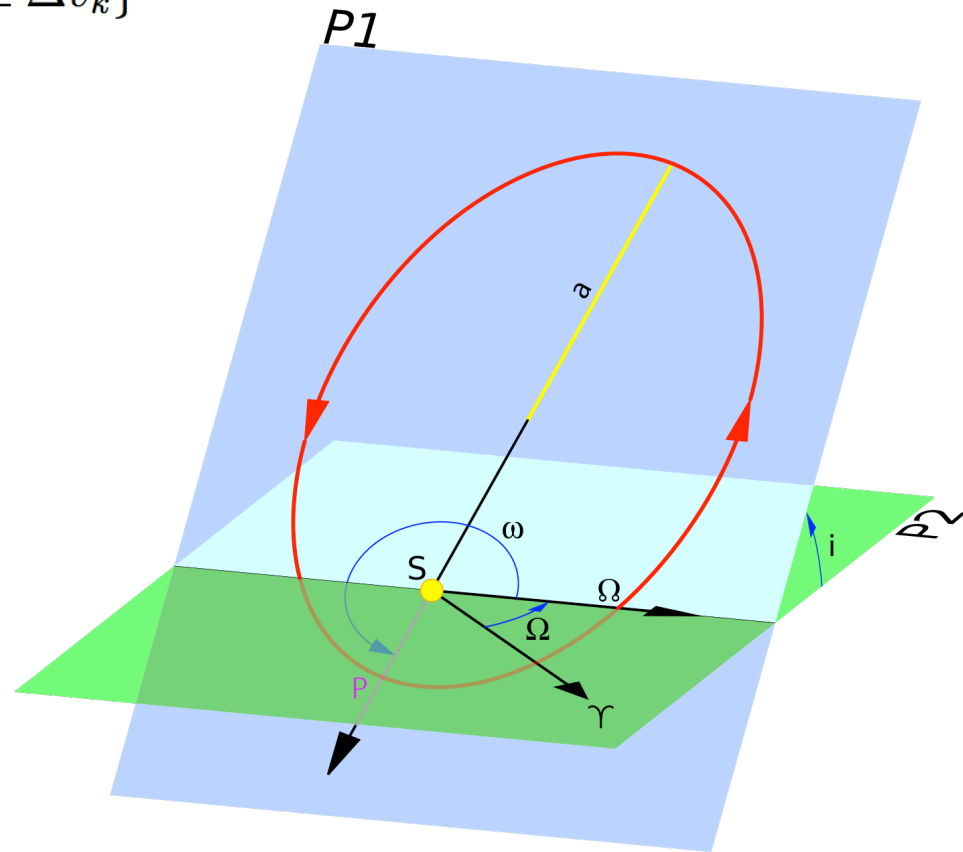
$$x(t) = x(t; M, d, \vec{\Omega})$$

$$y(t) = y(t; M, d, \vec{\Omega})$$

$$v(t) = v(t; M, d, \vec{\Omega})$$

$$\begin{aligned} \chi^2(M, d, \vec{\Omega}) = & \\ & \sum_{j=1}^{\max} \frac{(x_j - x(t; M, d, \vec{\Omega}))^2}{(\Delta x_j)^2} + \\ & \sum_{j=1}^{\max} \frac{(y_j - y(t; M, d, \vec{\Omega}))^2}{(\Delta y_{kj})^2} + \\ & \sum_{k=1}^{\max} \frac{(v_k - v(t; M, d, \vec{\Omega}))^2}{(\Delta v_k)^2} \end{aligned}$$

$$\min |_{(M, d, \vec{\Omega})} \chi^2(M, d, \vec{\Omega})$$



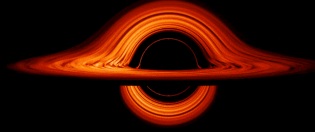
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- Part II Errors, Fitting and all that
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 - SgrA* flares
 - (A funny gas cloud)

Outline

- Part I The mass of Sgr A*
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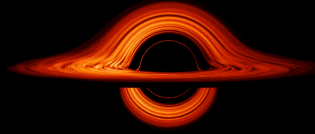
Why we are sure



Sgr A* is a black hole

	mass	position	faintness	size	
Eckart & Genzel 1996, Ghez+ 1998, 2003, 2005, Schödel+ 2002, Eisenhauer+ 2005, Gillessen+ 2009, 2017, Böhle+ 2016, GRAVITY Coll 2018, 2022	IR stellar orbits $4.3 \times 10^6 M_{\odot}$ $\pm 0.3\%$	IR + VLA	radio + IR	IR + X-ray timescales $\approx 0.1 \text{ AU}$	Baganoff+ 2001, Genzel+ 2003, Dodds-Eden+ 2009
Reid & Brunthaler 2004, 2020	VLBA : radio Sgr A* $> 1 \times 10^6 M_{\odot}$	mass and radio Sgr A* coincide	Sgr A* cannot have a surface	VLBI /EHT 0.4 AU	Doeleman+ 2008, Akiyama+ 2022,
		Reid+ 2007 Plewa+ 2015	Broderick & Narayan 2006 Broderick+ 2009	GRAVITY flares 0.5 AU	GRAVITY Coll. 2018b

Why we are sure



Sgr A* is a black hole

mass

**IR
stellar orbits
 $4.3 \times 10^6 M_{\odot}$
 $\pm 0.3\%$**

**VLBA :
radio Sgr A*
 $> 1 \times 10^6 M_{\odot}$**

Eckart & Genzel
1996, Ghez+ 1998,
2003, 2005,
Schödel+ 2002,
Eisenhauer+ 2005,
Gillessen+ 2009,
2017, Böhle+
2016, GRAVITY
Coll 2018, 2022

Reid & Brunthaler
2004, 2020

position

**IR + VLA

mass and
radio Sgr A*
coincide**

Reid+ 2007
Plewa+ 2015

faintness

**radio + IR

Sgr A*
cannot have
a surface**

Broderick &
Narayan 2006
Broderick+ 2009

size

**IR + X-ray
timescales
 ≈ 0.1 AU**

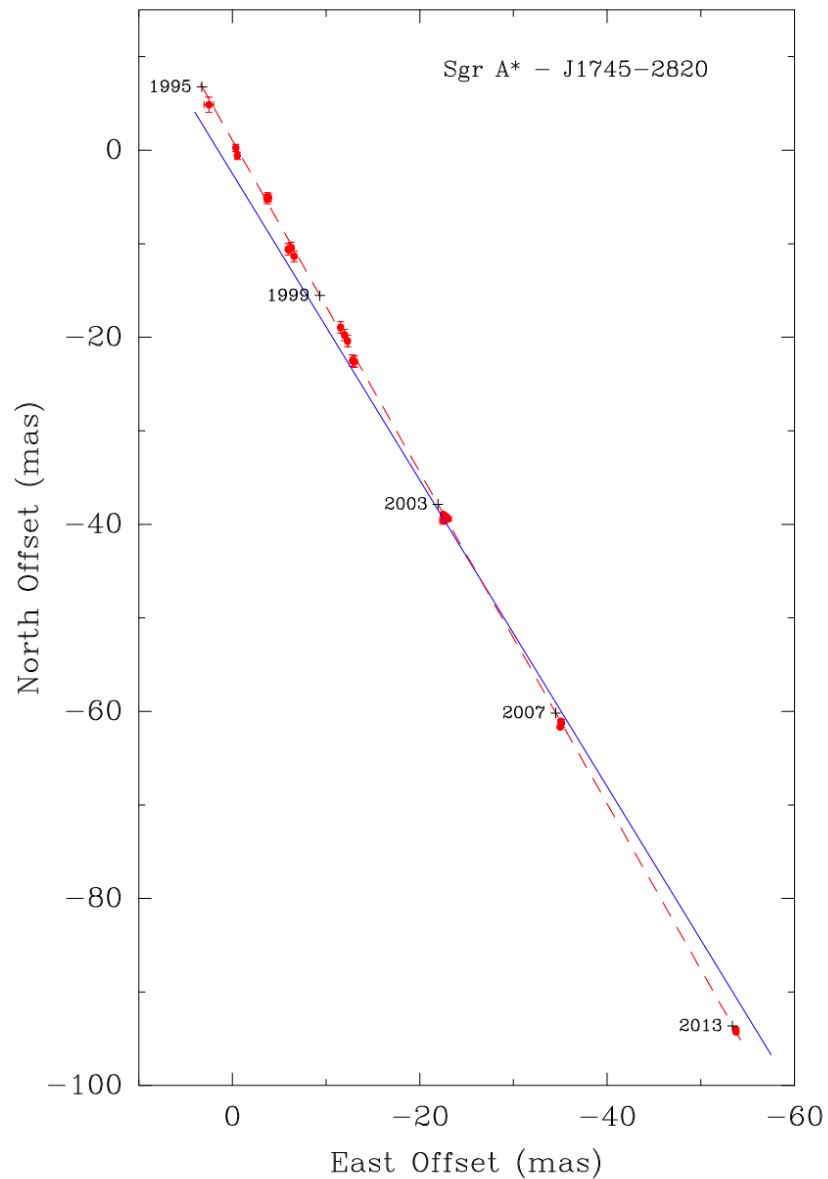
**VLBI /EHT
0.4 AU**

**GRAVITY
flares
0.5 AU**

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Genzel+ 2003,
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2009

Doeleman+
2008, Akiyama+
2022,

GRAVITY Coll.
2018b

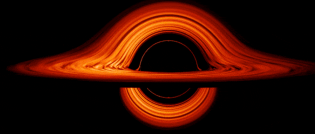


Sgr A* must be very heavy

- perfectly linear motion
 - reflex motion of Sun (~200 km/s)
- intrinsic motion
 - gal. l : -7.2 ± 8.5 km/s
 - gal. b: -0.4 ± 0.9 km/s
- Sgr A* is much heavier than surrounding stars
 - mass $> 10^6 M_{\odot}$

Reid 2007, 2020

Why we are sure



Sgr A* is a black hole

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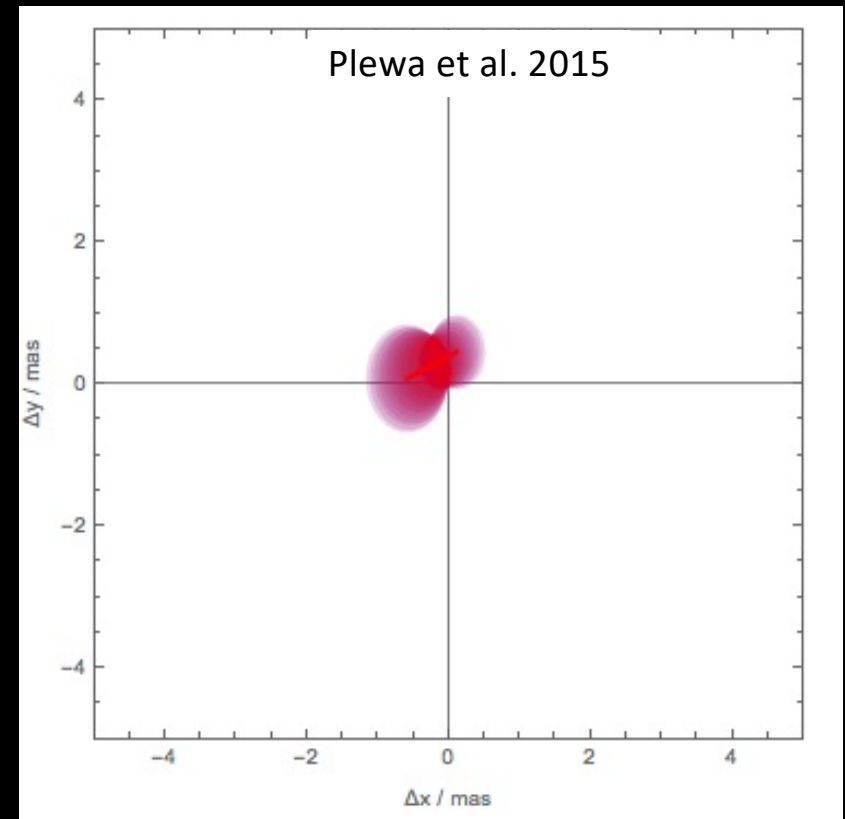
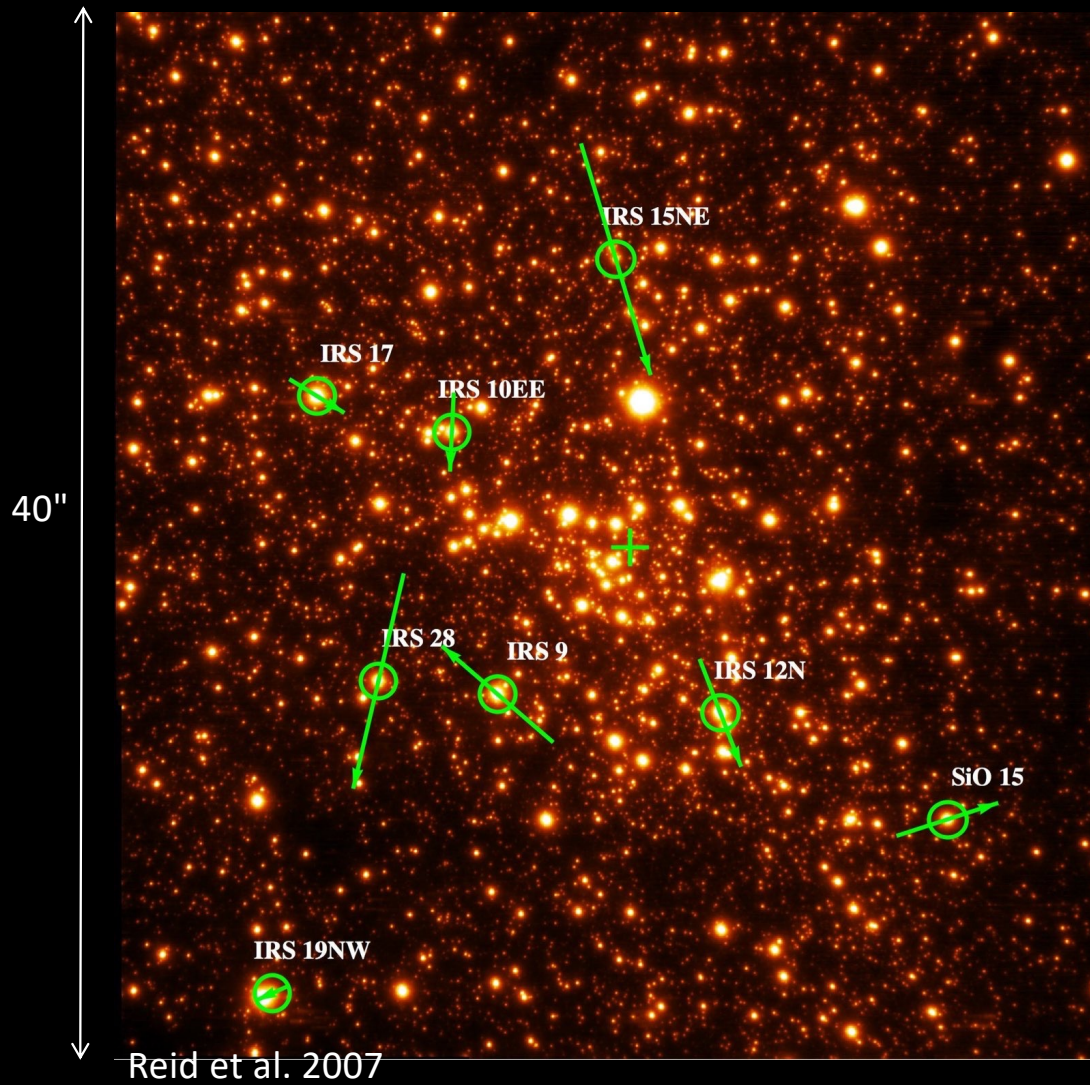
Baganoff+ 2001,
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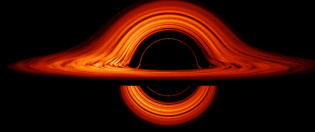
GRAVITY Coll.
2018b

The positions of mass & radio-Sgr A* agree to within 1 mas

- Sgr A*
 - radio source
- SiO maser stars
 - IR sources
 - radio sources



Why we are sure



Sgr A* is a black hole

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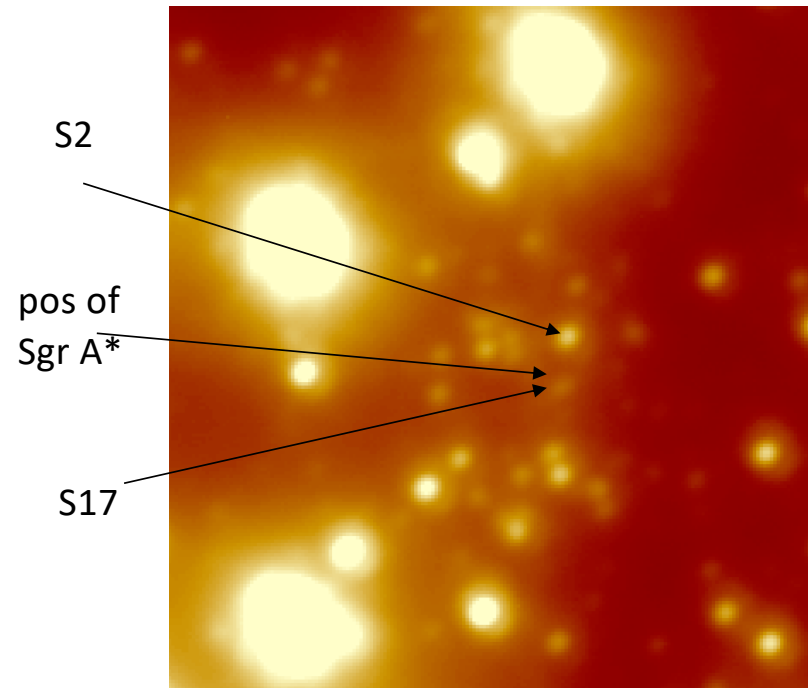
Sgr A* should be bright - but is not

Limit: Eddington luminosity
radiation pressure = gravitation

$$L_{Edd} = \frac{4\pi G m_p c}{\sigma_T} M$$

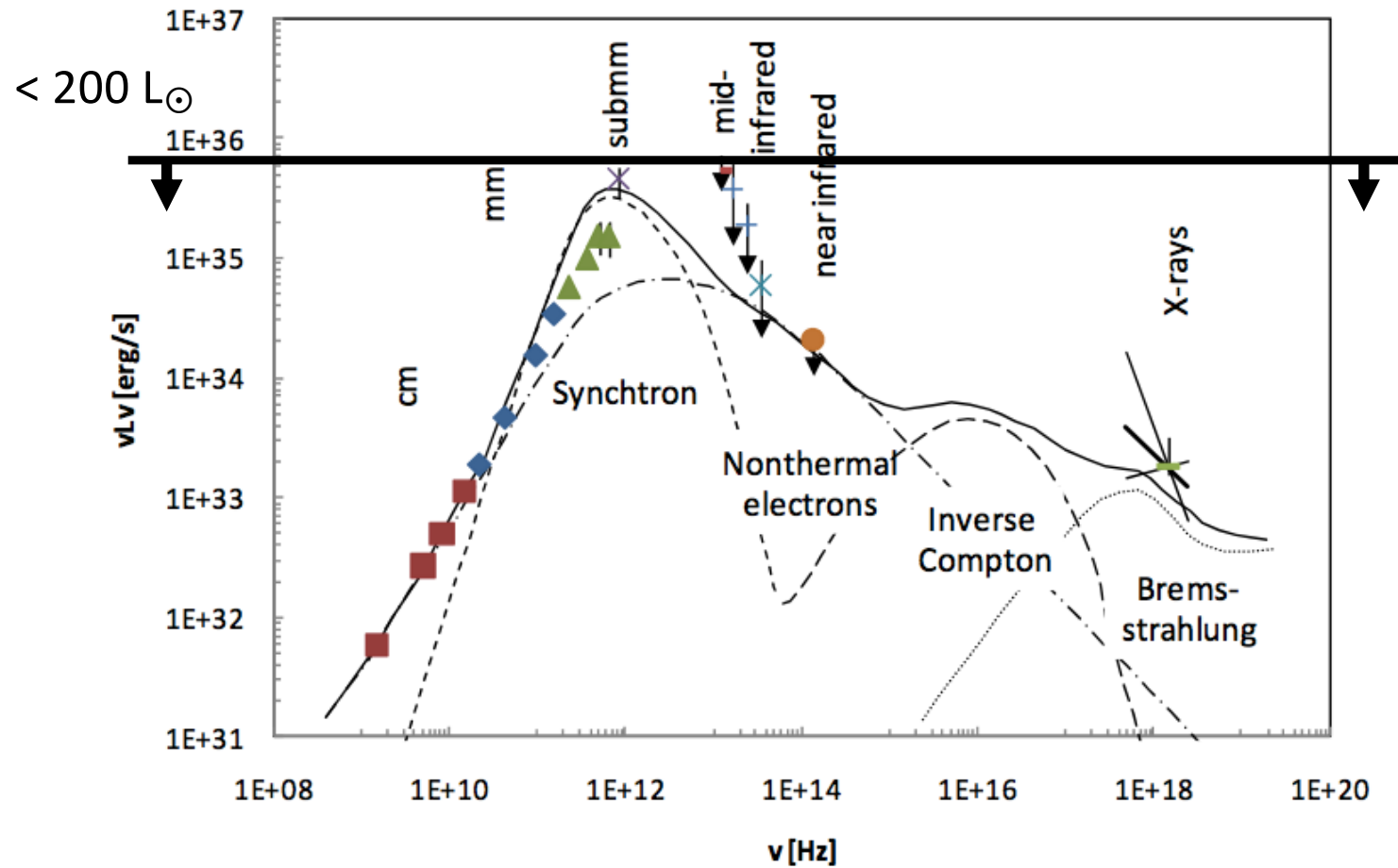
$$L = \eta \times 5 \times 10^{44} \text{ erg/s}$$

$$= \eta \times 10^{11} L_{\odot}$$



Sgr A* is dim at all wavelengths:
 $\eta \sim 10^{-8}$

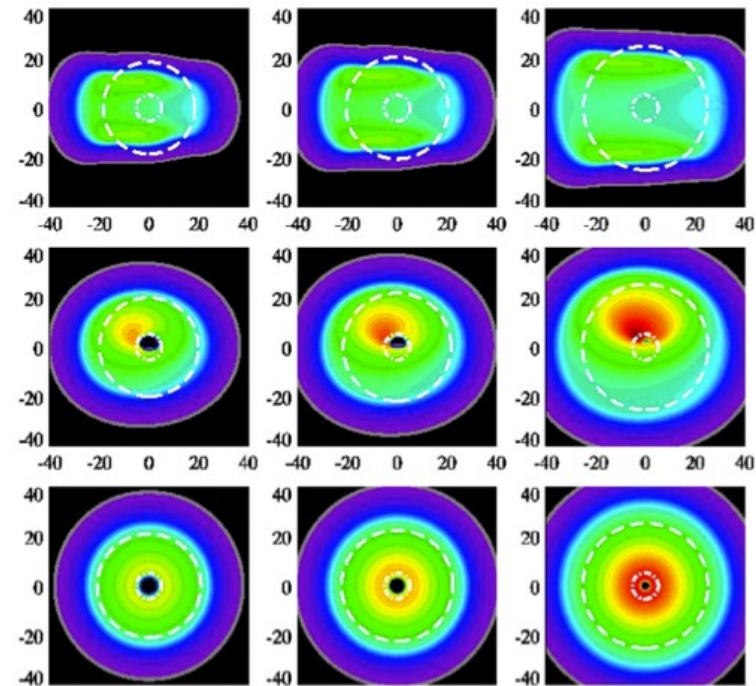
Genzel et al. 2010



Radiatively Inefficient Accretion Flow

low L/L_{Edd} is a combination of:

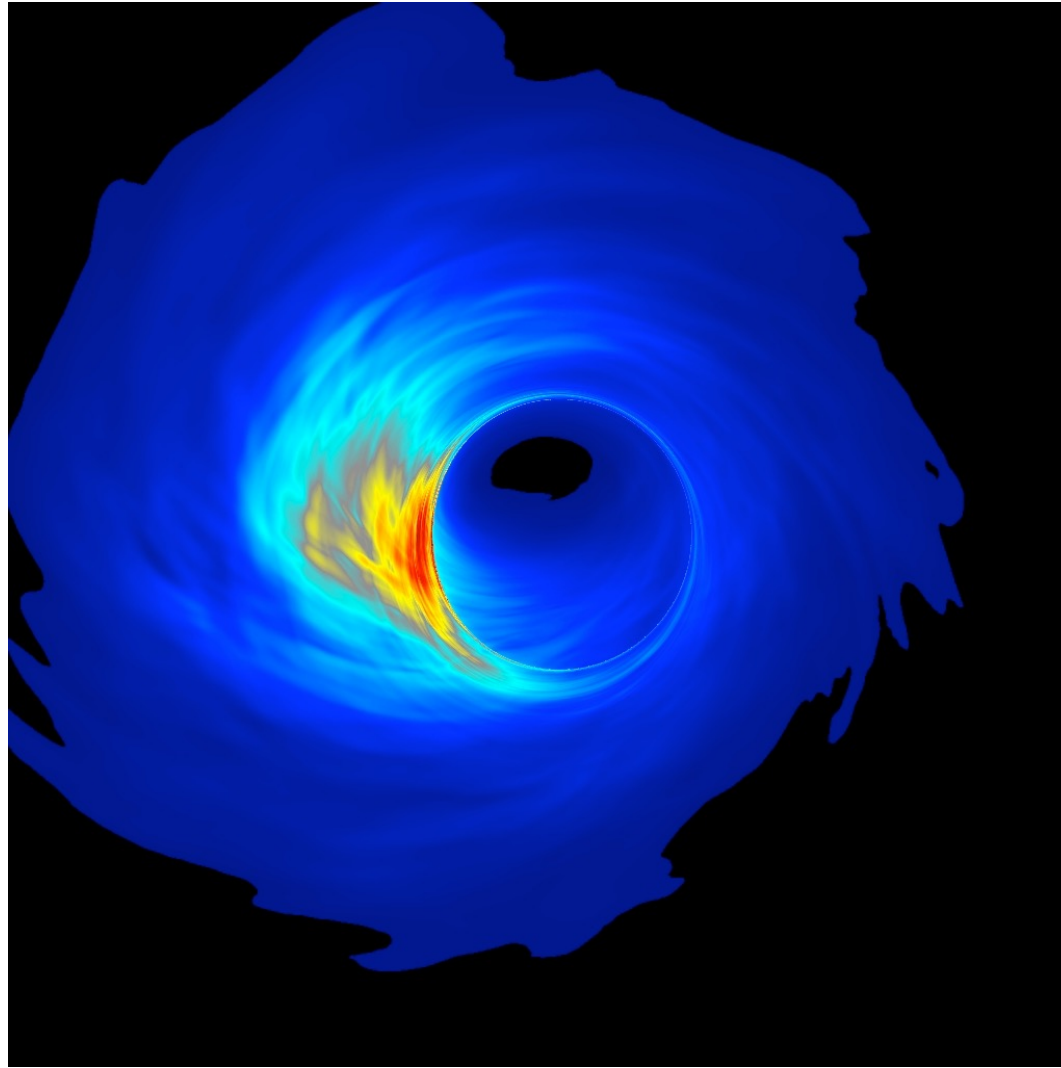
- low accretion rate at Bondi radius
- low efficiency angular momentum transport
- low efficiency energy transfer protons to electrons
- most of the gas arriving at a few R_s ejected back out



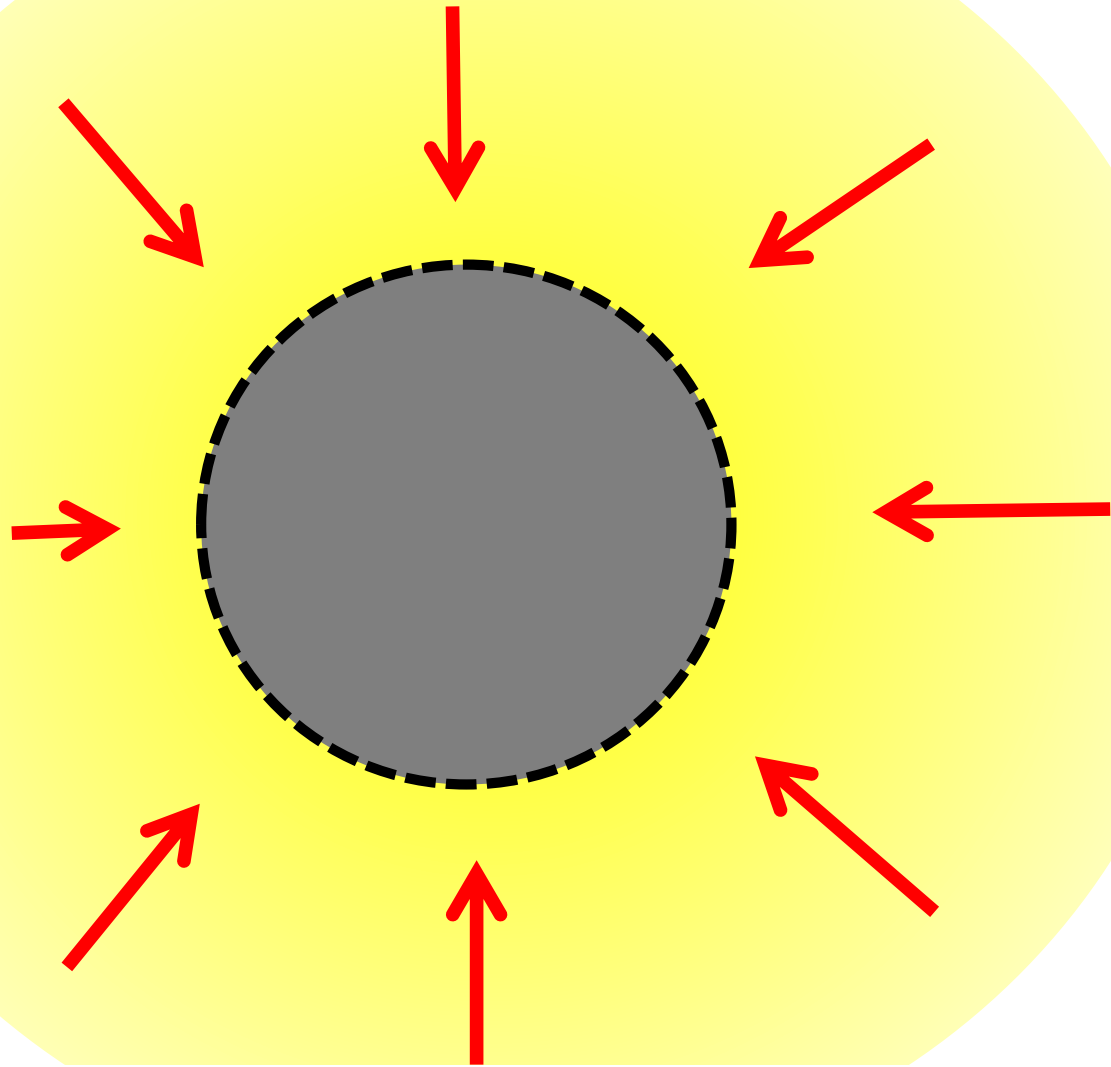
Yuan et al. 2009

Begelman, Bower, Blandford, Cuadra, De Villers, Falcke, Hawley, Krolik, Liu, Melia, Markoff, Marrone, Narayan, Quataert, Rees, Revnitsev, Stone, Wang, Yuan 1995-2013

MHD simulation with GR ray-tracing



RIAF predicts accretion rate



accretion flow

radio emission:
RIAF

accretion rate:
 $\sim 5 \times 10^{-8} M_{\odot}/\text{yr}$

99% of that
reaches
“surface”

Assume, material is crashing onto surface

$$P = \sigma A T^4$$

↑ ↑ ↑ ↑

$P = \dot{M}c^2$ constant area temperature

↑ ↑

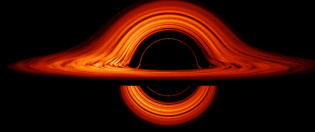
Accretion Rate measured size < 4 Rs 20000 K

Sgr A* has an event horizon

- Hot surface: 20000 K
 - infrared light
- Size: 1/10 Earth orbit
 - like a normal star
 - extremely luminous
- But normal stars can be seen
- Sgr A* cannot have a surface



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 $\approx 0.1 \text{ AU}$

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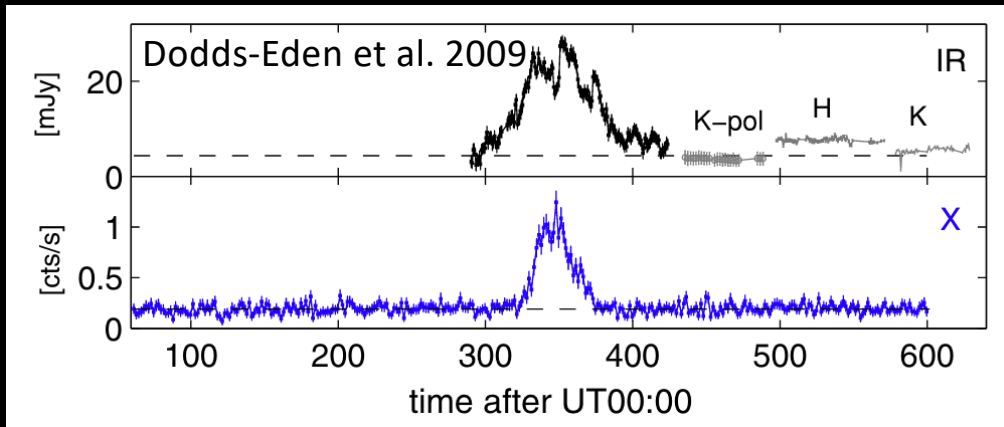
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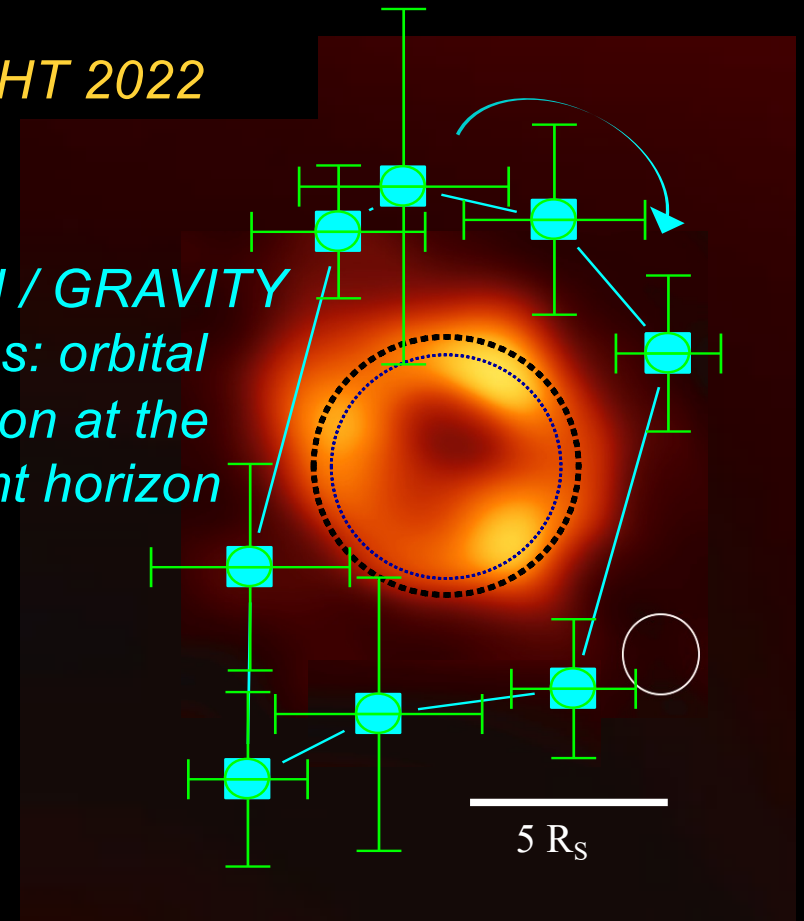
GRAVITY Coll.
2018b

Sgr A* is very compact



EHT 2022

*VLT / GRAVITY
flares: orbital
motion at the
event horizon*



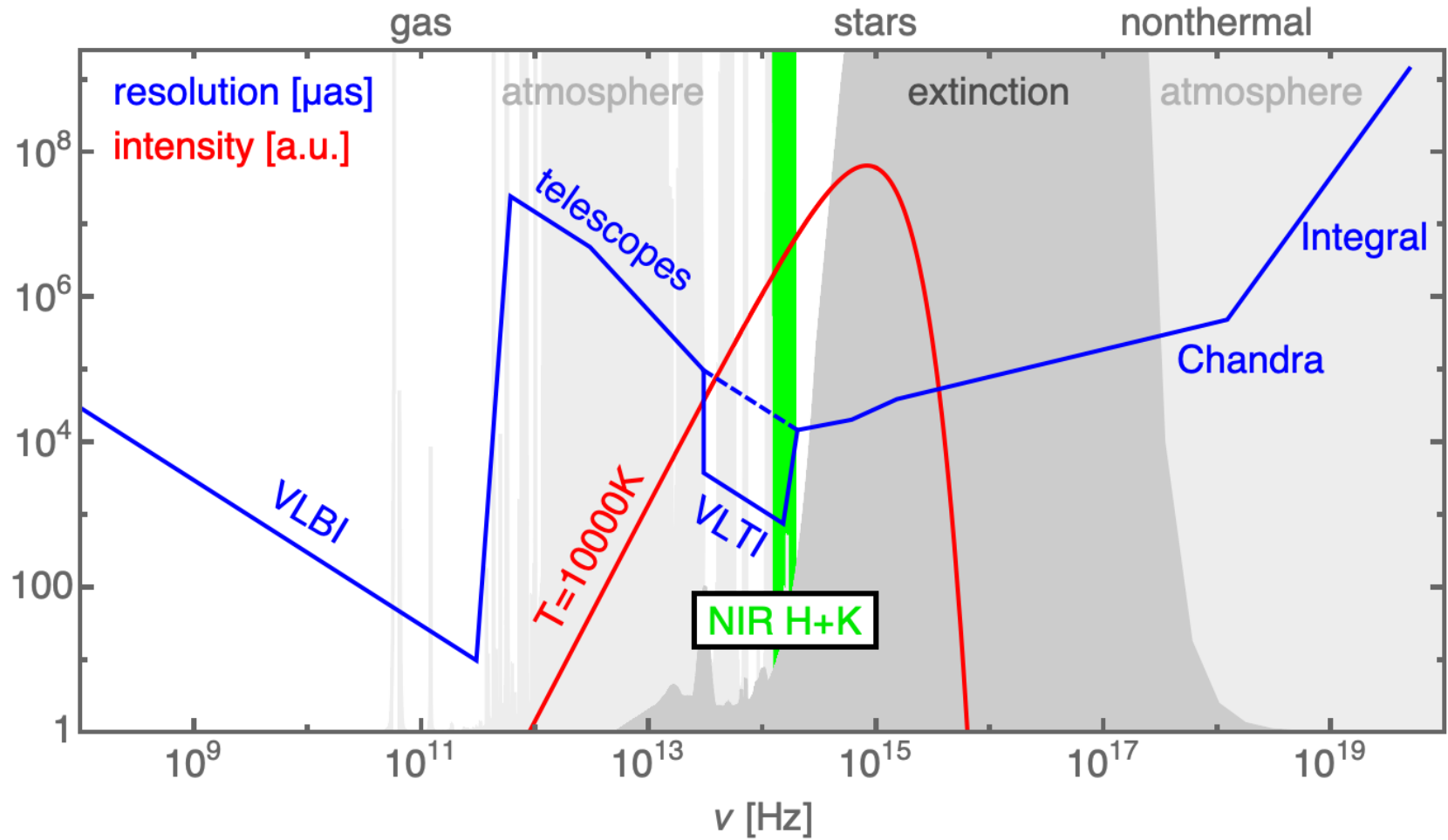
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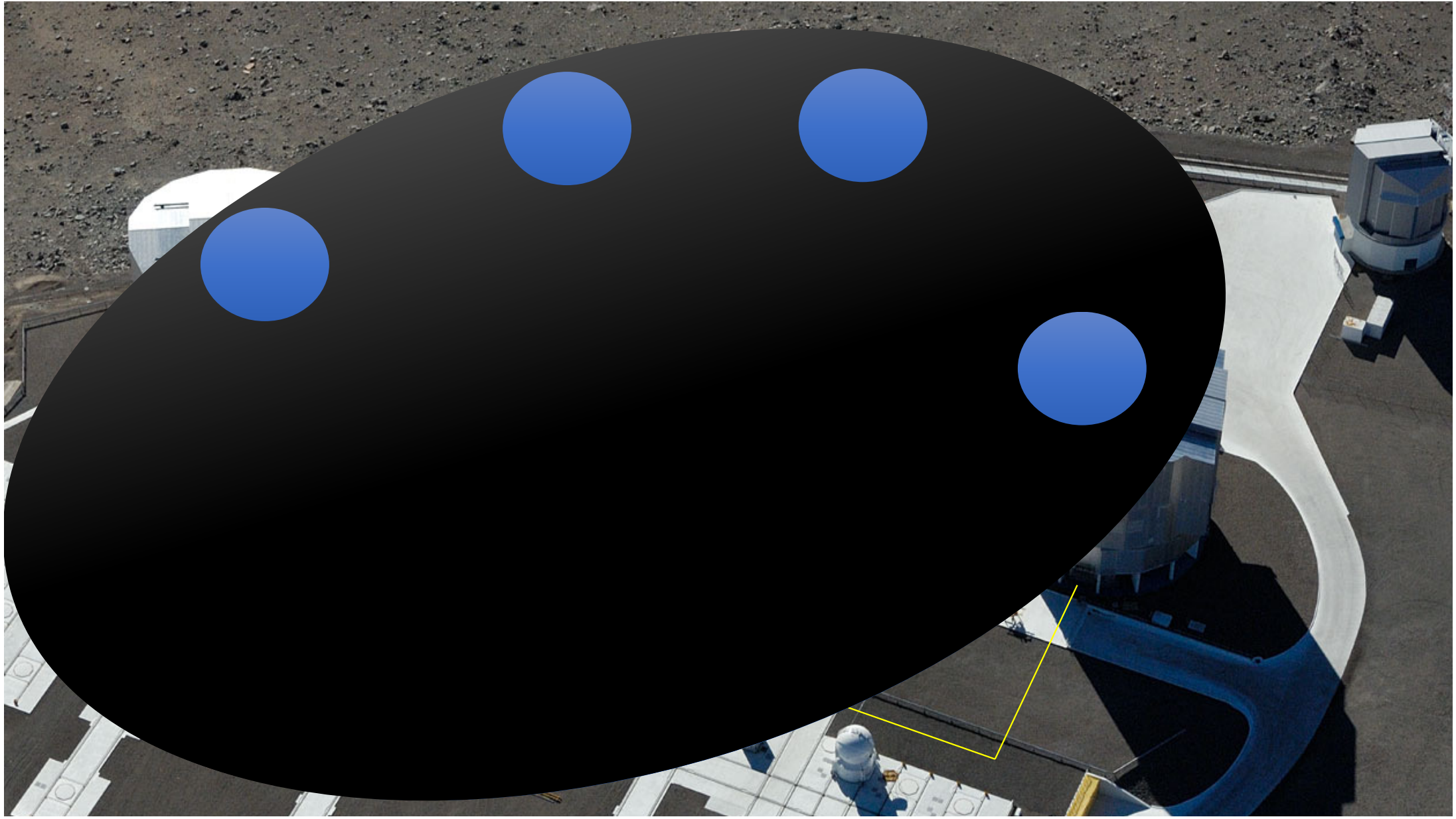
How to zoom in further ?



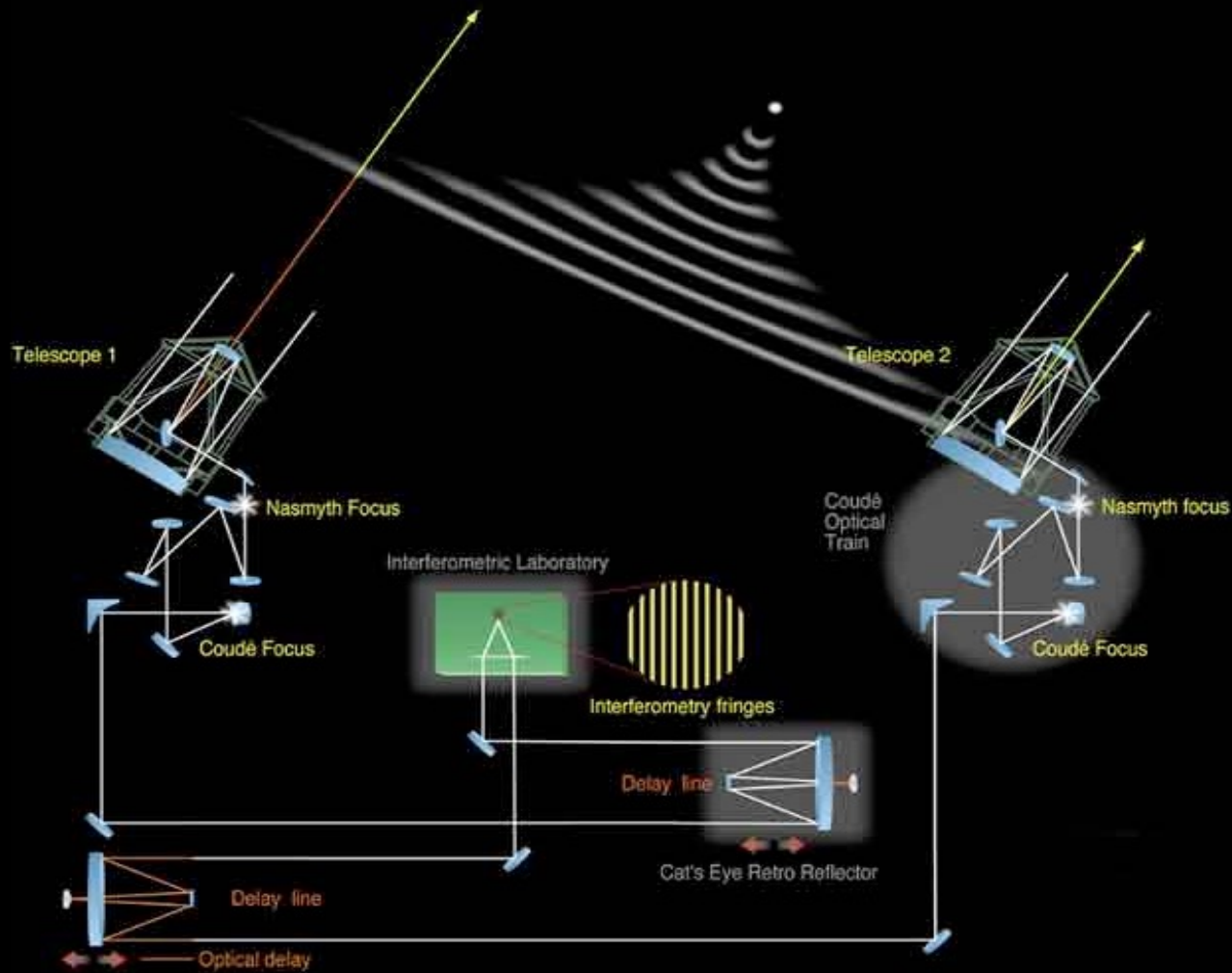
The sweet spot in wavelength: The near-infrared







But: Optical interferometry is hard



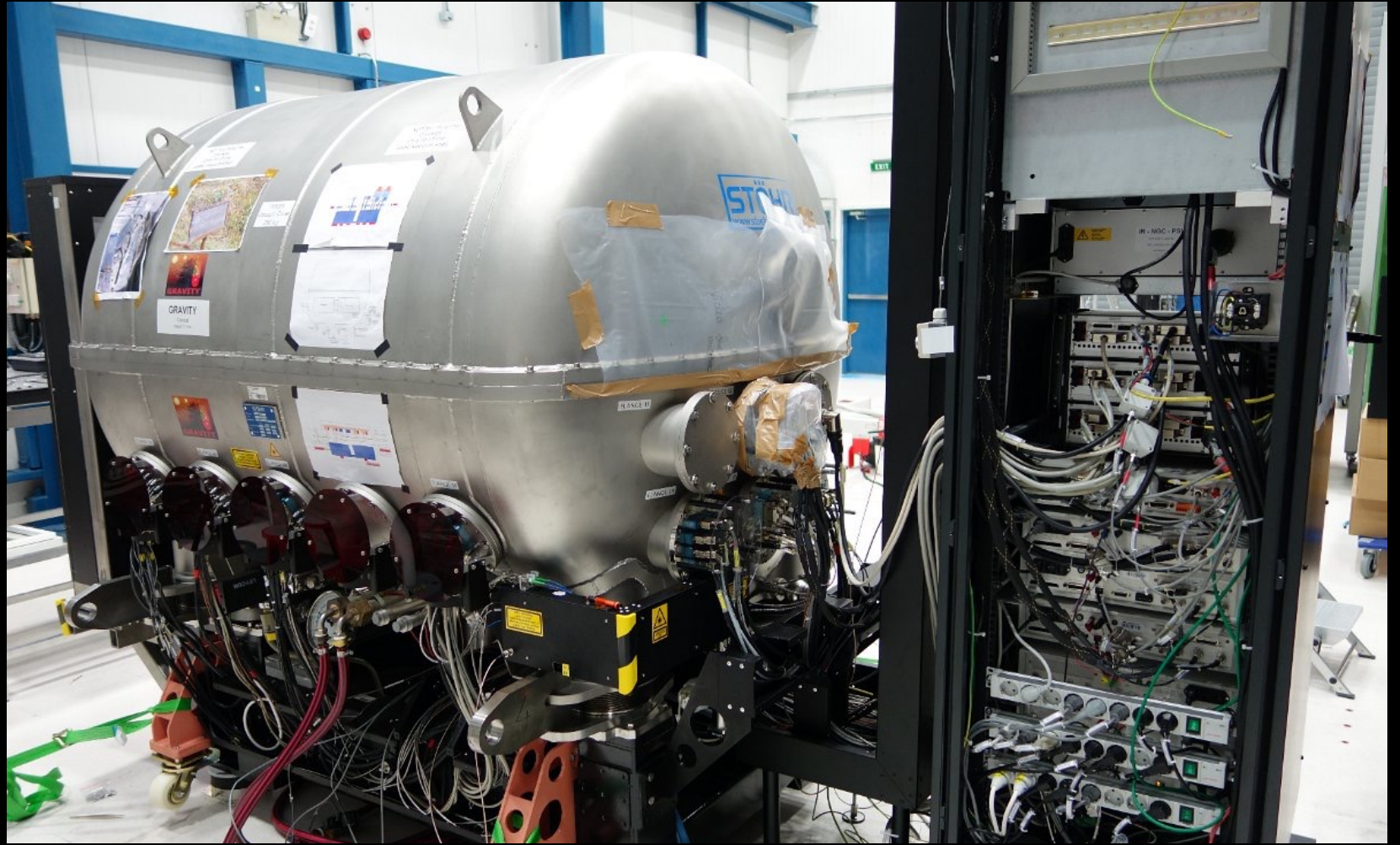
Online
cophasing to
nano-meters



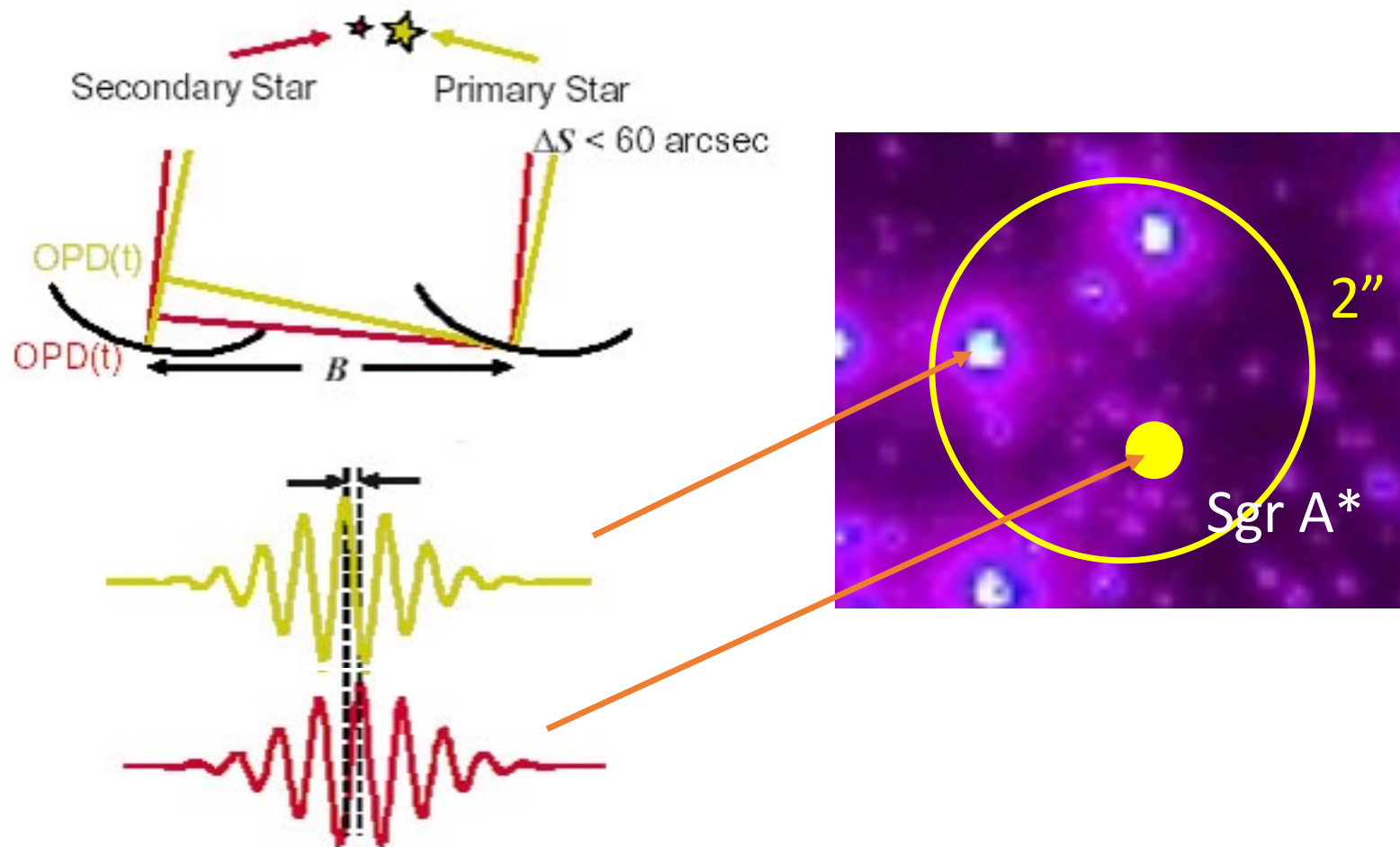
Dual feed, 4-telescope, adaptive optics-assisted, fringe-tracking beam combiner instrument



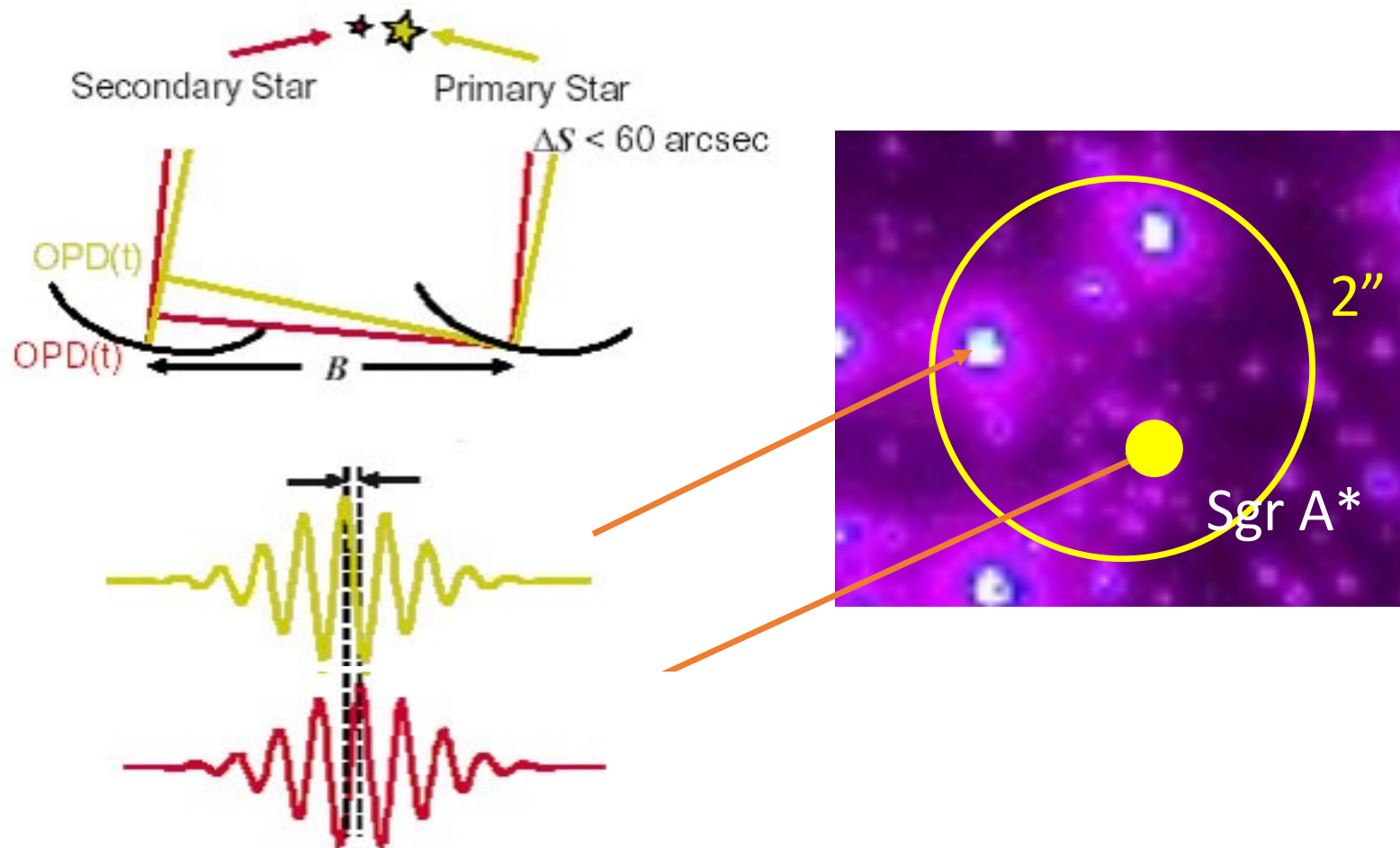
Partenariat Haute résolution Angulaire Sol-Espace

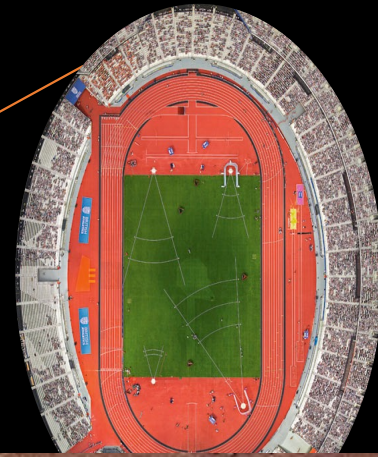


Phase referenced imaging



Dual beam astrometry



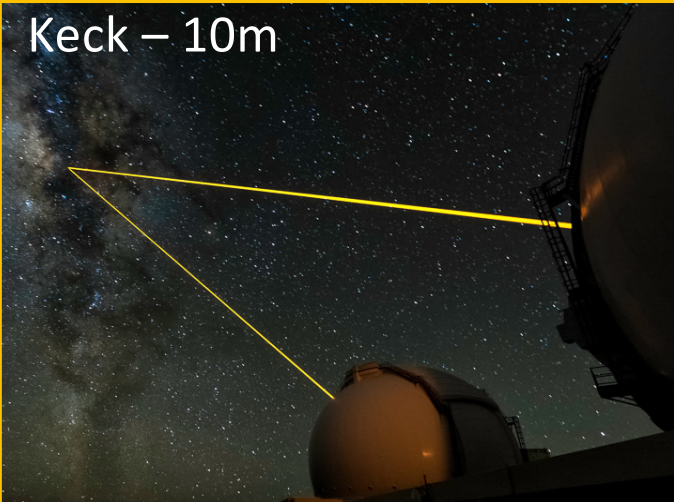


S2
200 mas

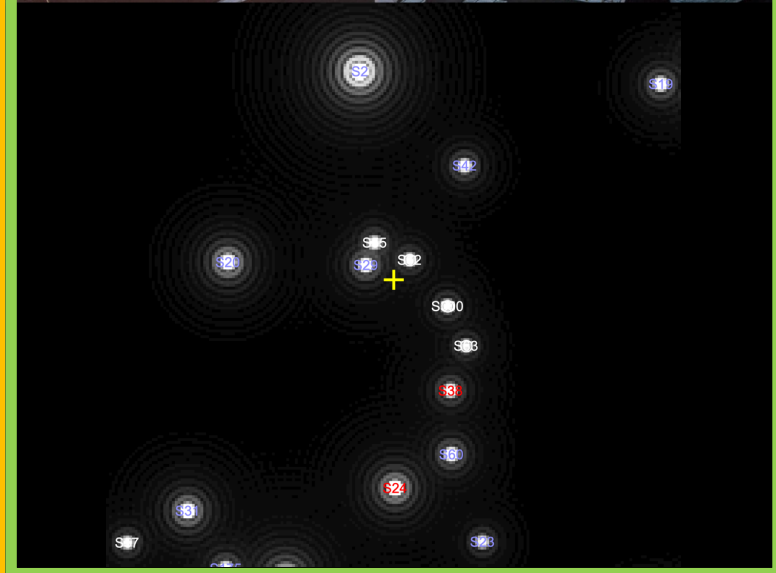
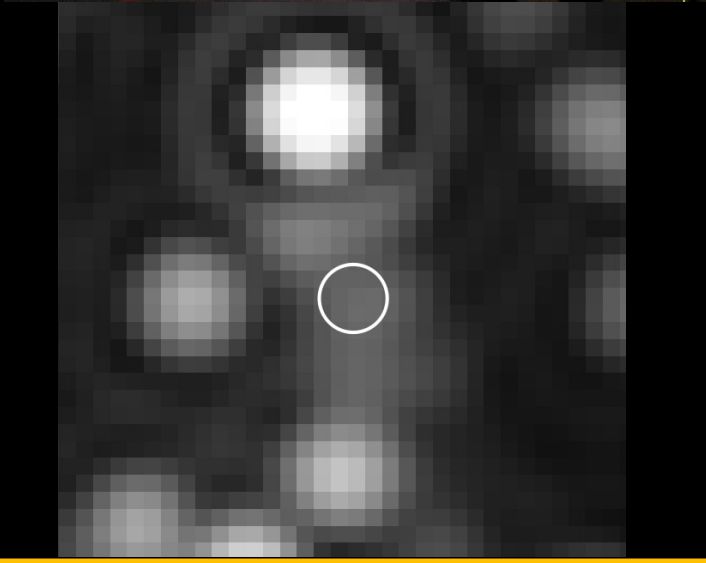
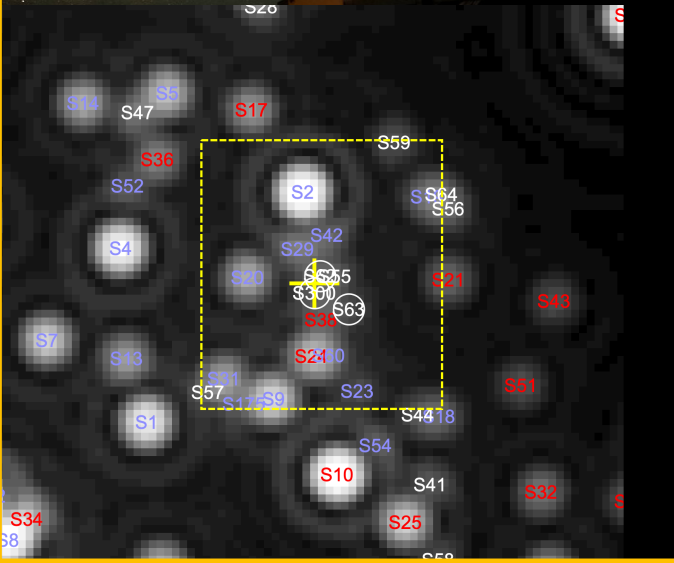
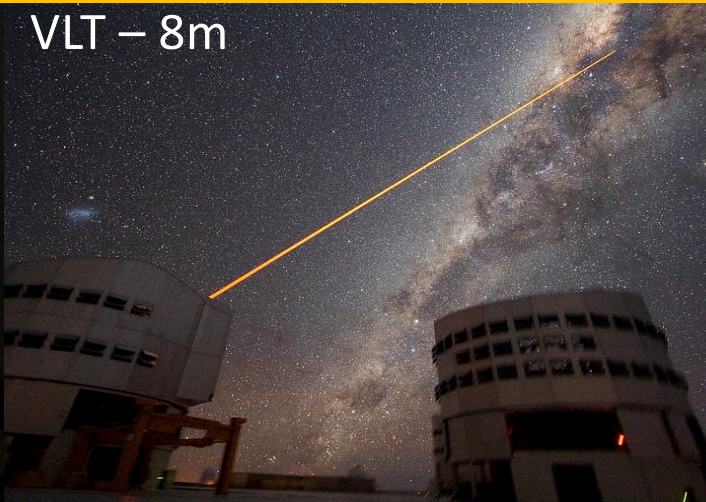


Laser guide star adaptive optic is 15x less sharp

Keck – 10m



VLT – 8m

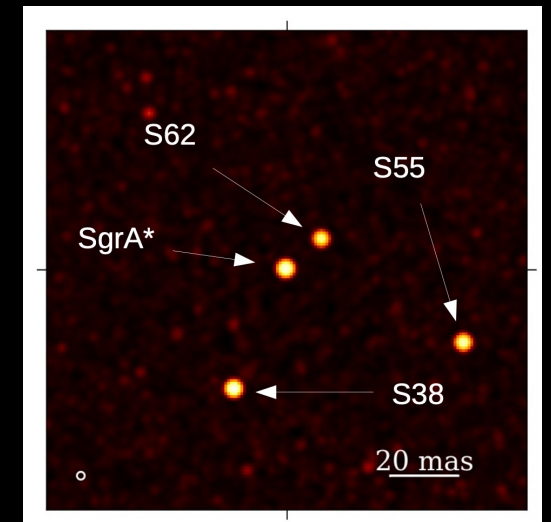
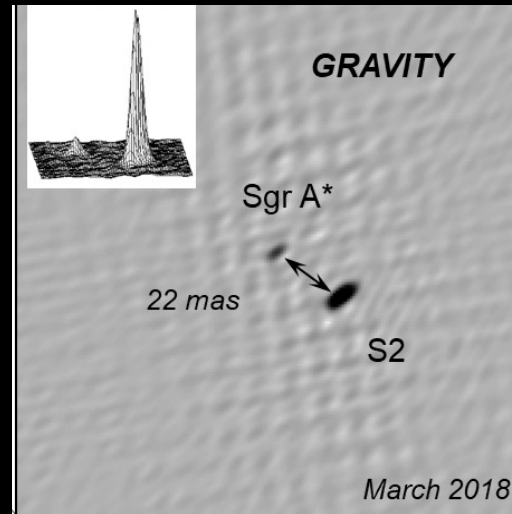
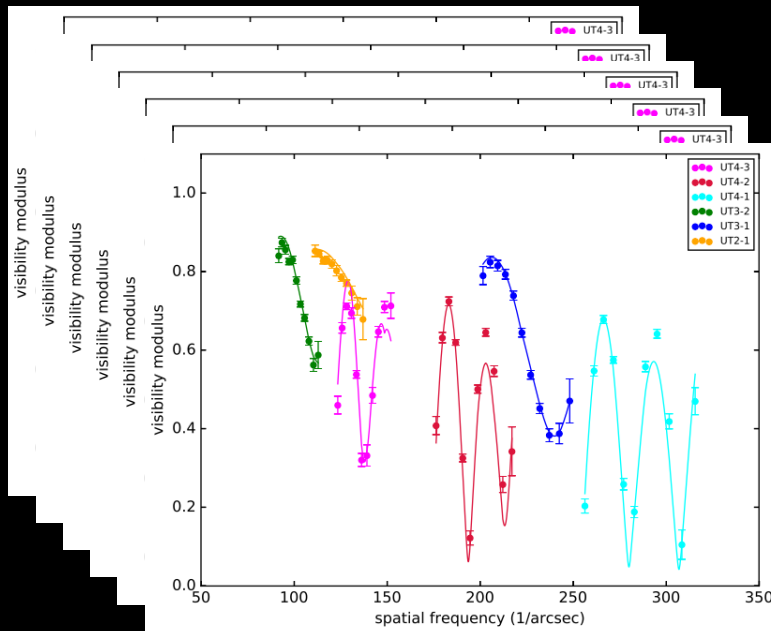


GRAVITY interferometry of Sgr A* and S2

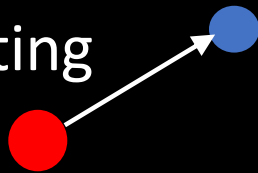
Fringe contrast and phase

Images

GRAVITY coll. 2017, 2018a



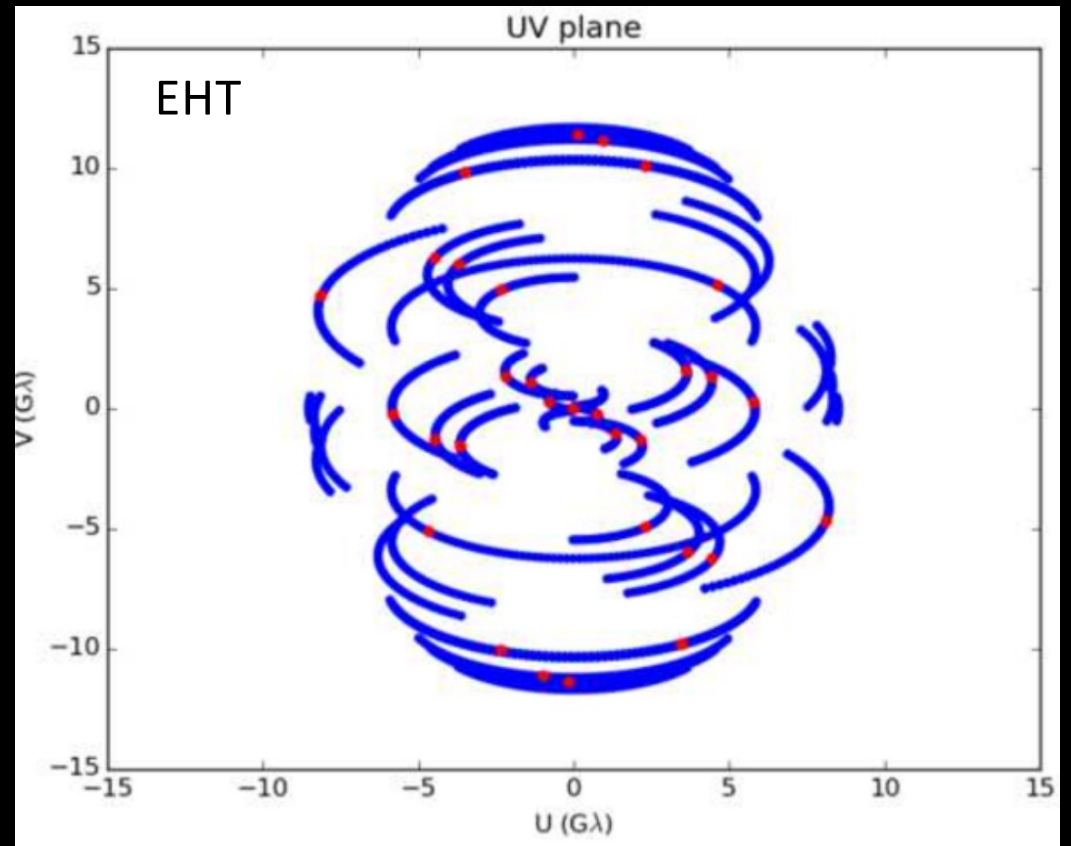
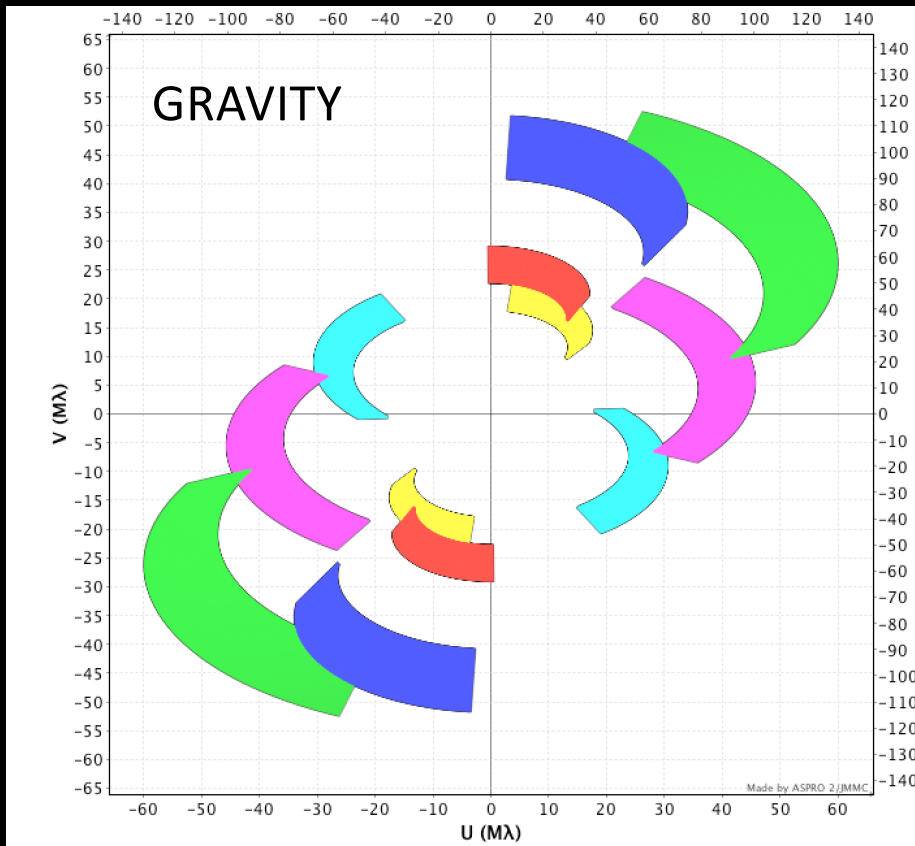
Model fitting

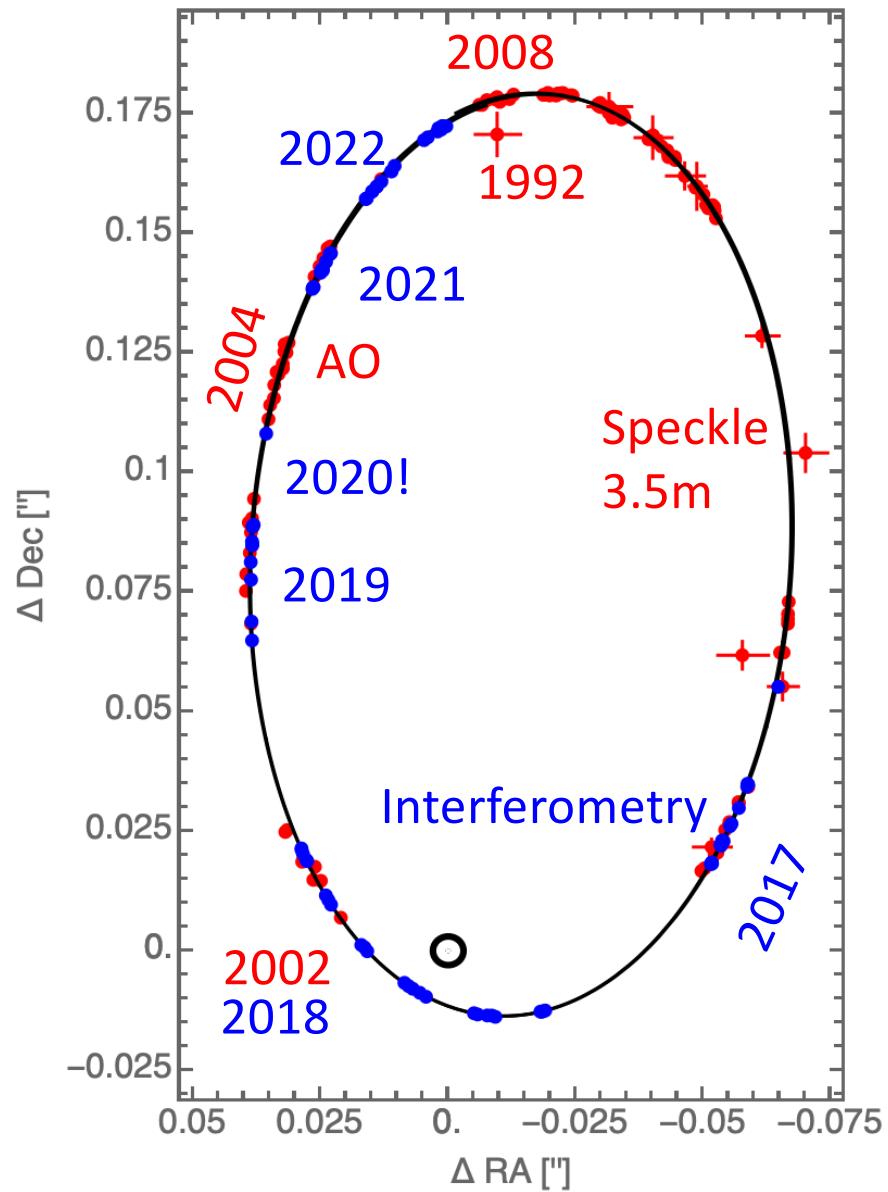


Source & instrument model:

- Separation vector
- Brightness-ratio
- source colors
- Bandwidth smearing
- Injection per telescope

UV coverage good even with 6 baselines
- due to the 20% width of the K-band



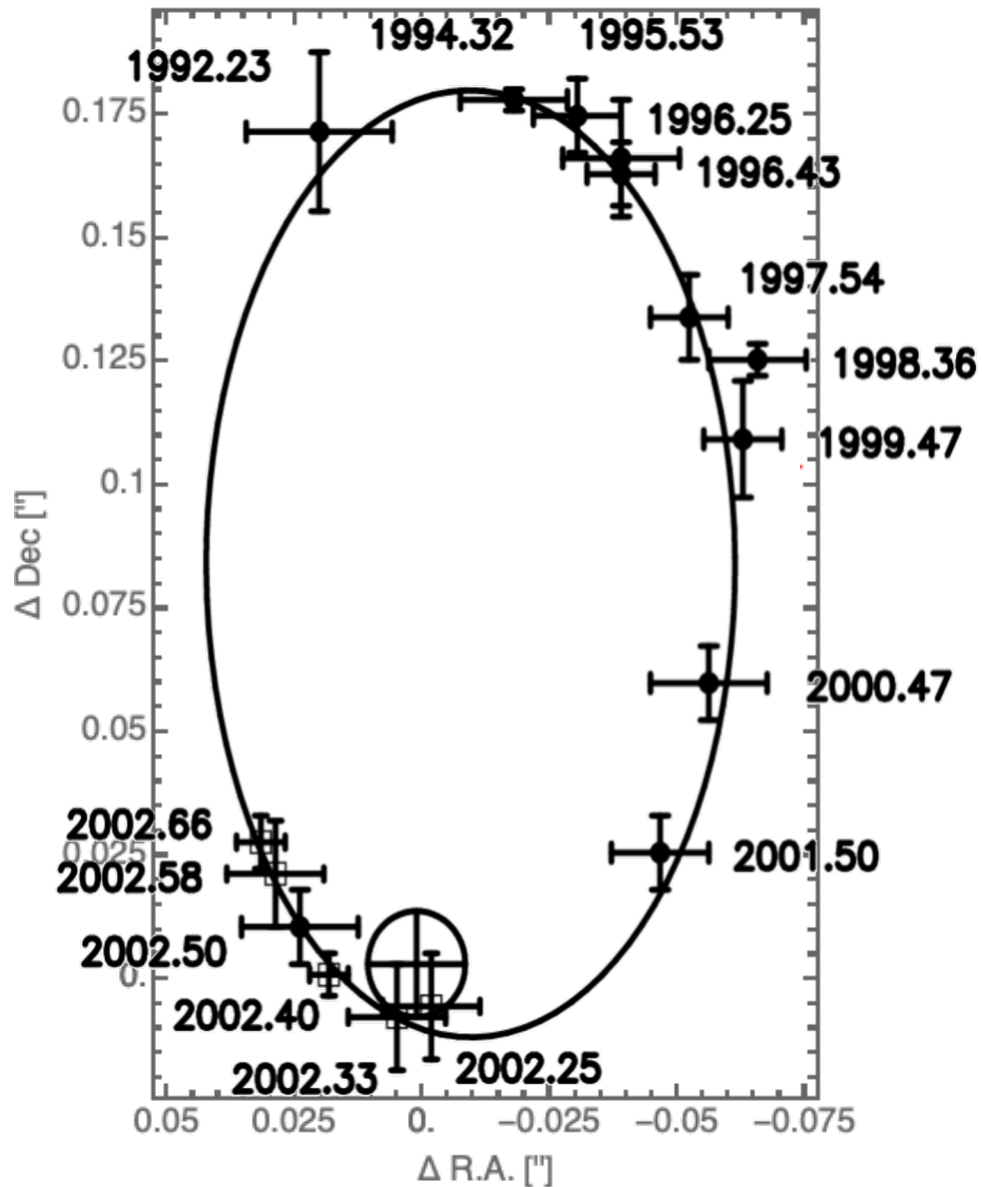


S2: the showcase star

Adaptive Optics &
Interferometry data

(Gillessen et al. 2017, Boehle et al. 2016, GRAVITY collaboration 2018a, 2019a, 2020, 2021, 2022)

- period: 16 years
- semi major axis: 125 mas
- eccentricity 0.88
- angular momentum and energy have errors of 0.2%
- pericenter: 19 May 2018



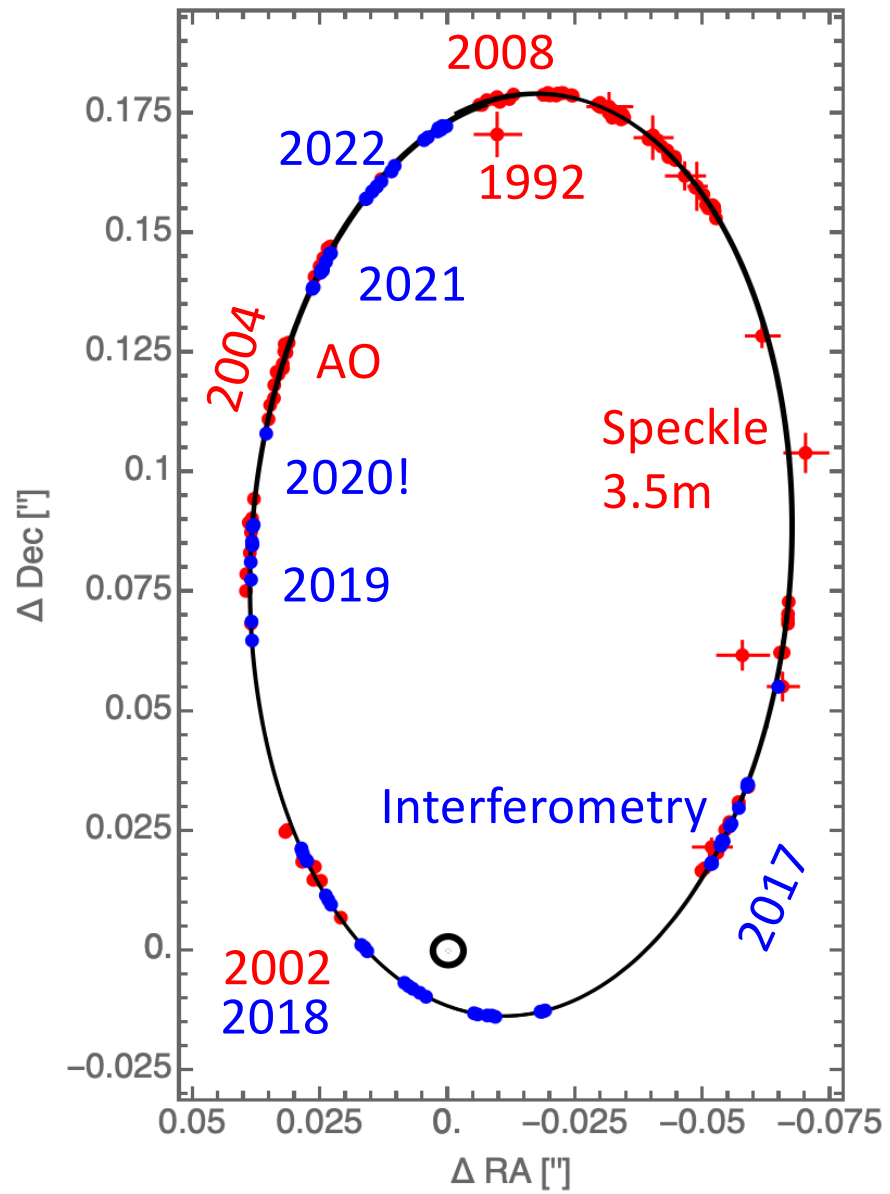
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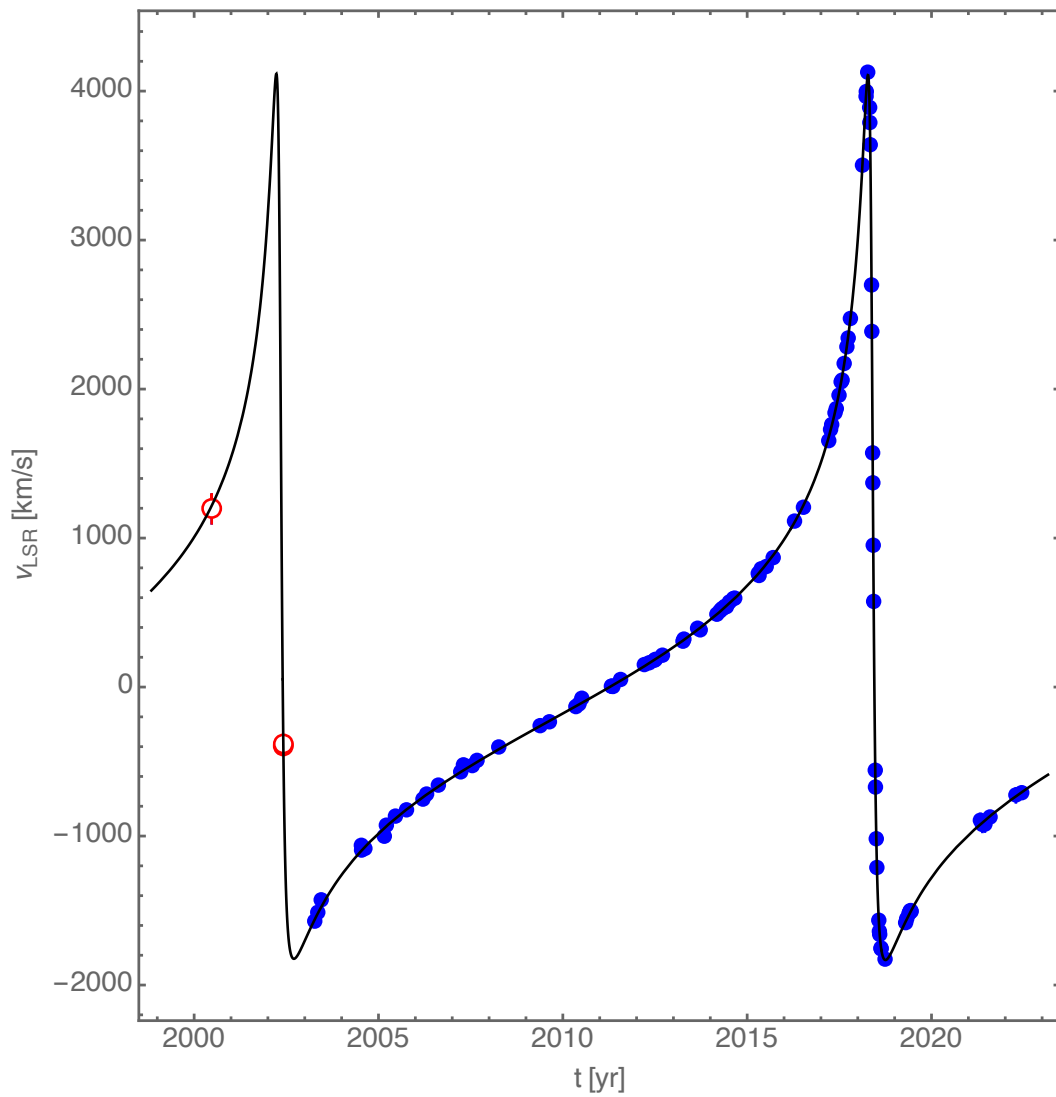
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The S2 orbit is very fruitful

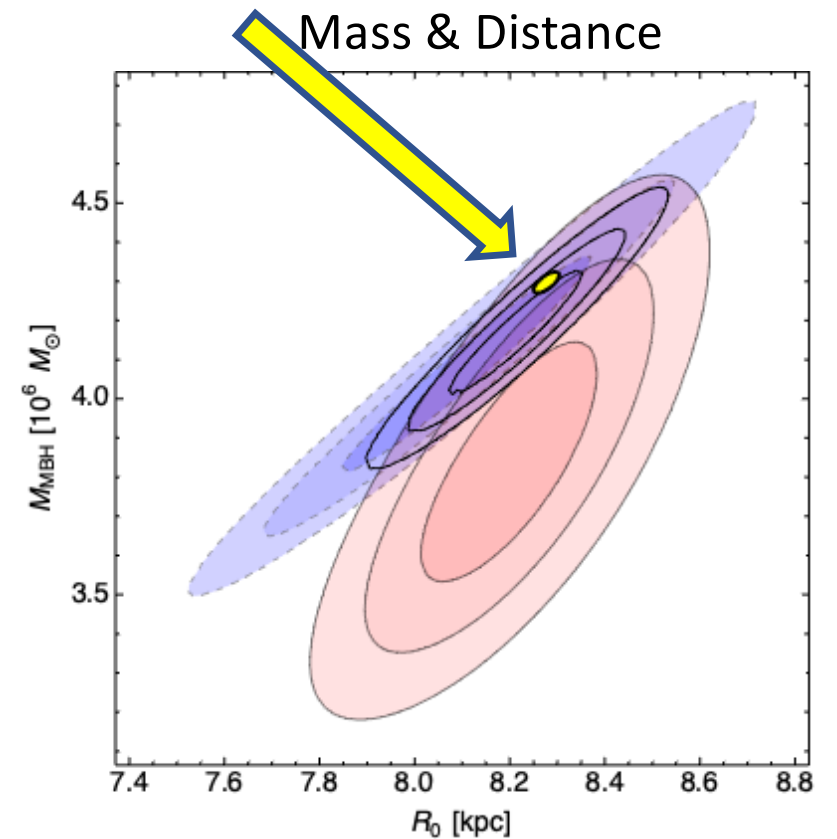
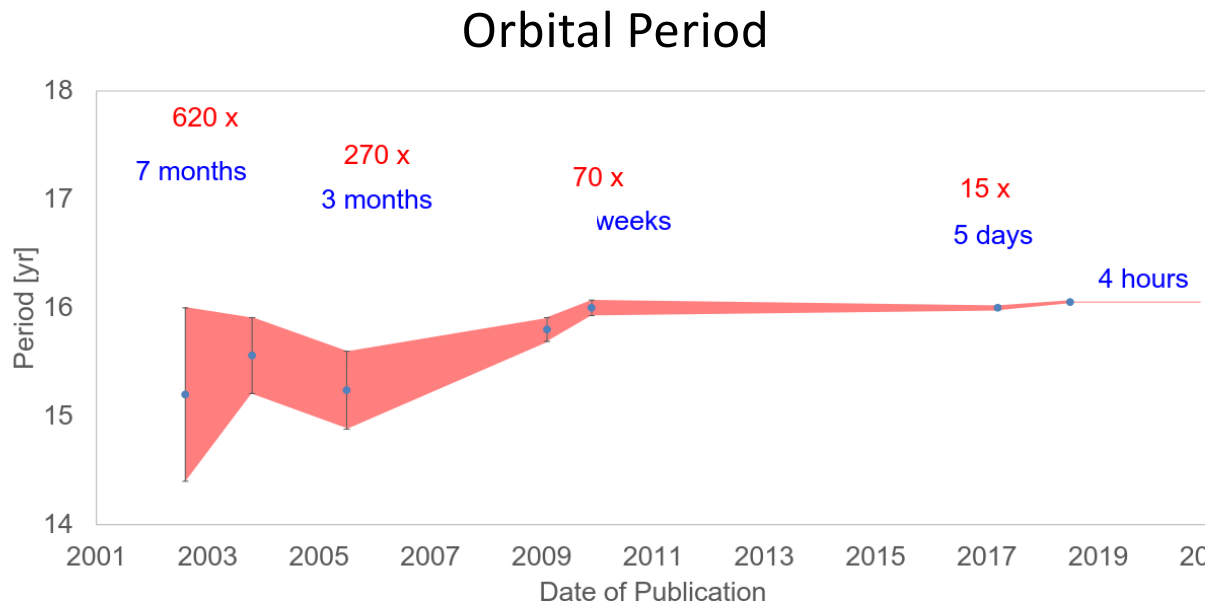




The S2 orbit is very fruitful

- geometric distance
- relativistic redshift
- test of equivalence principle
- relativistic precession
- limits on extended mass

Improvements by orders of magnitudes



Images: Positions, proper motions, angular velocity in mas/yr

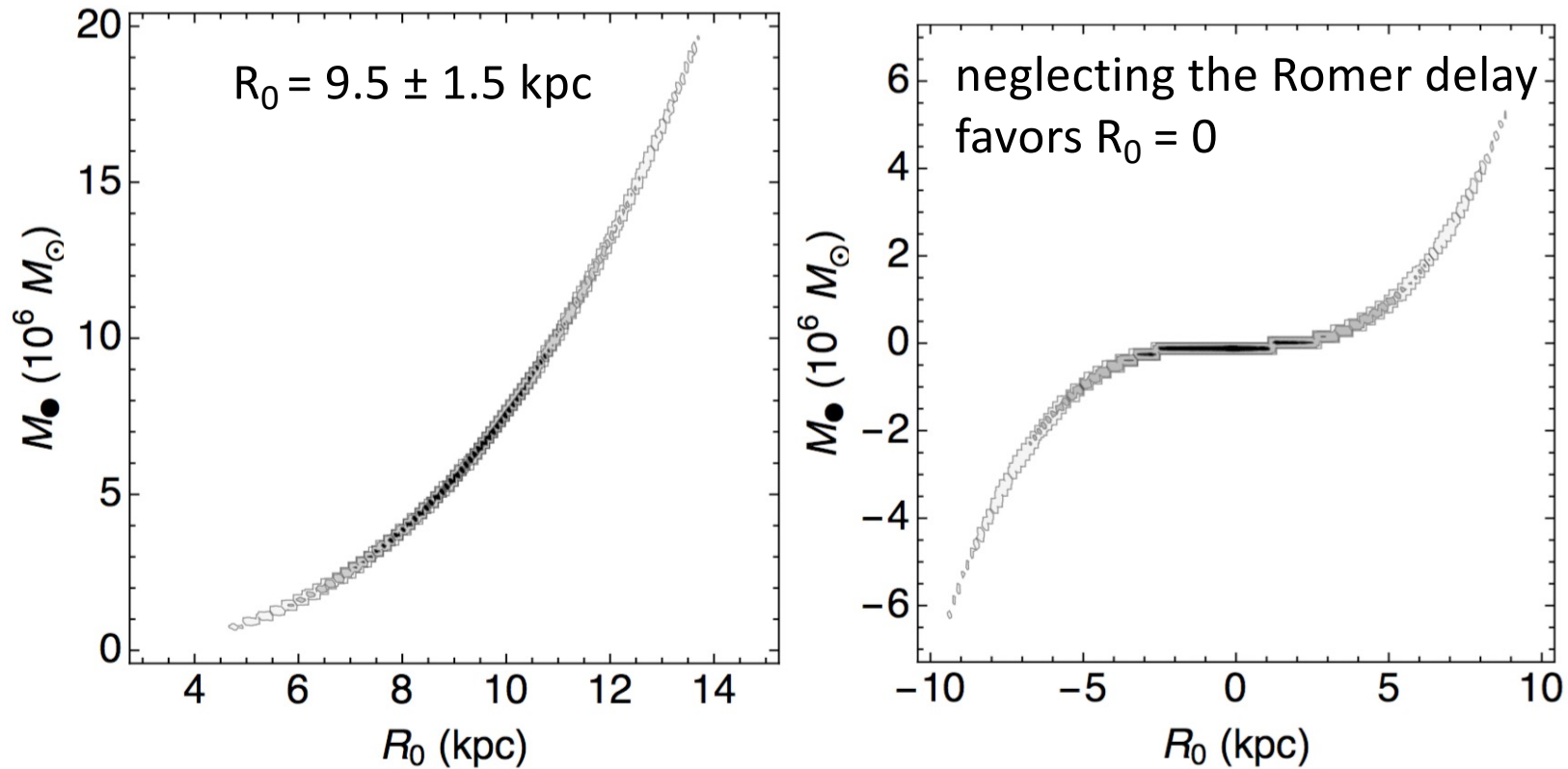
Spectra: Radial velocity in km/s

Conversion of angles to absolute length: **The distance**

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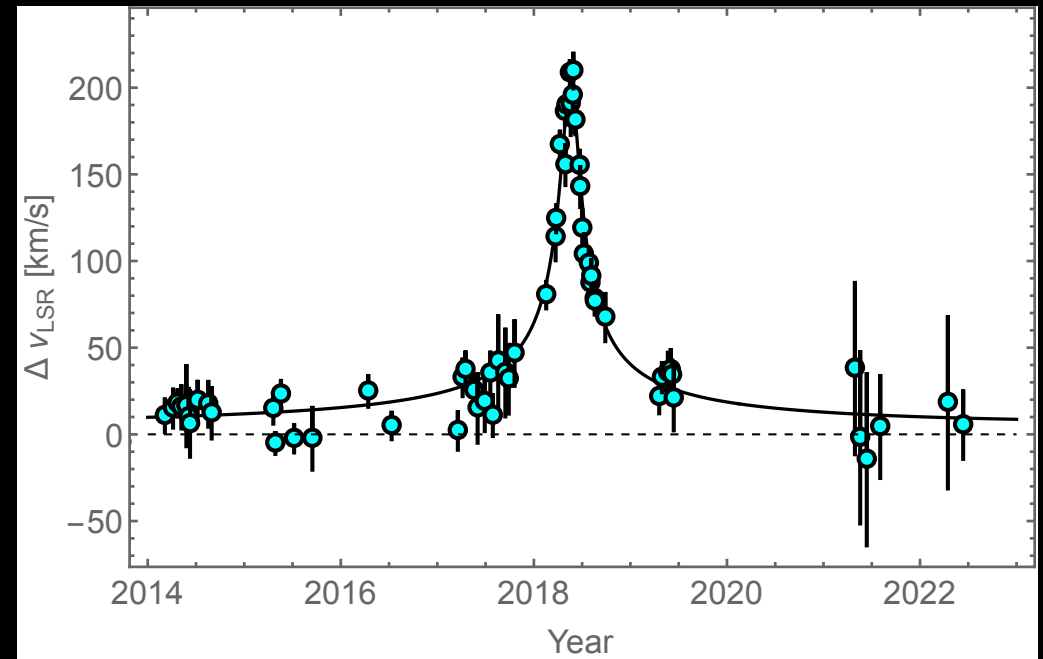
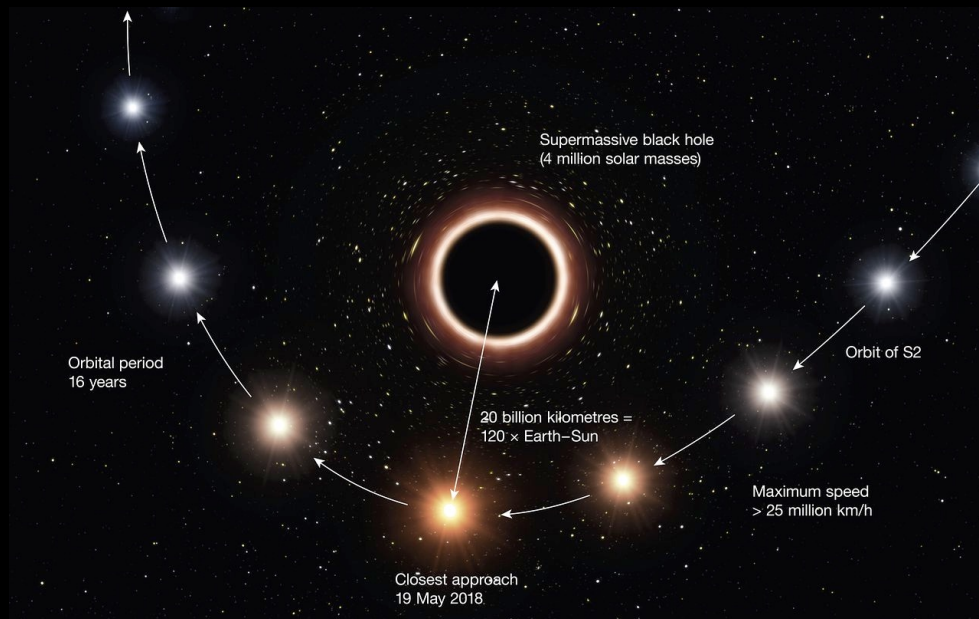
We can measure R_0 **without** radial velocities
via the Romer delay in the astrometry



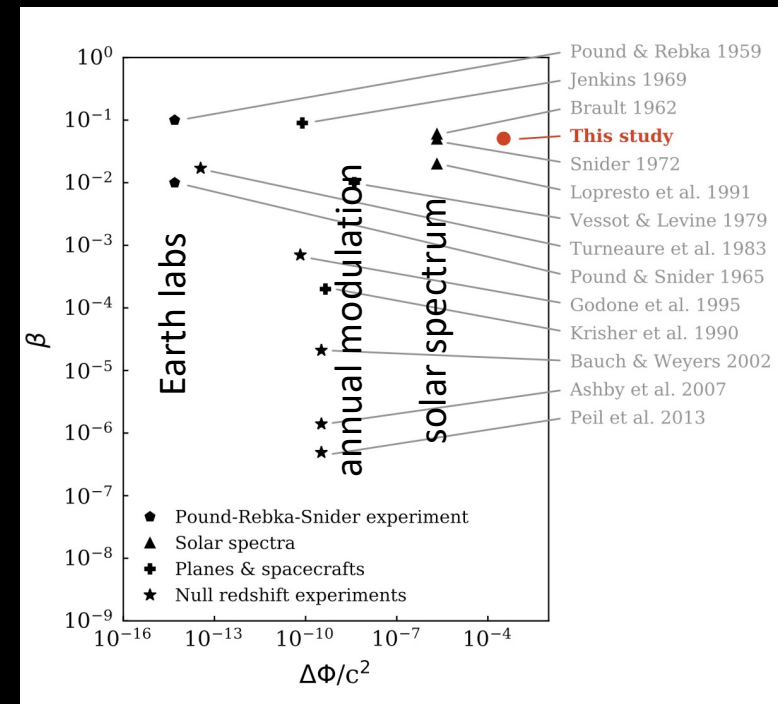
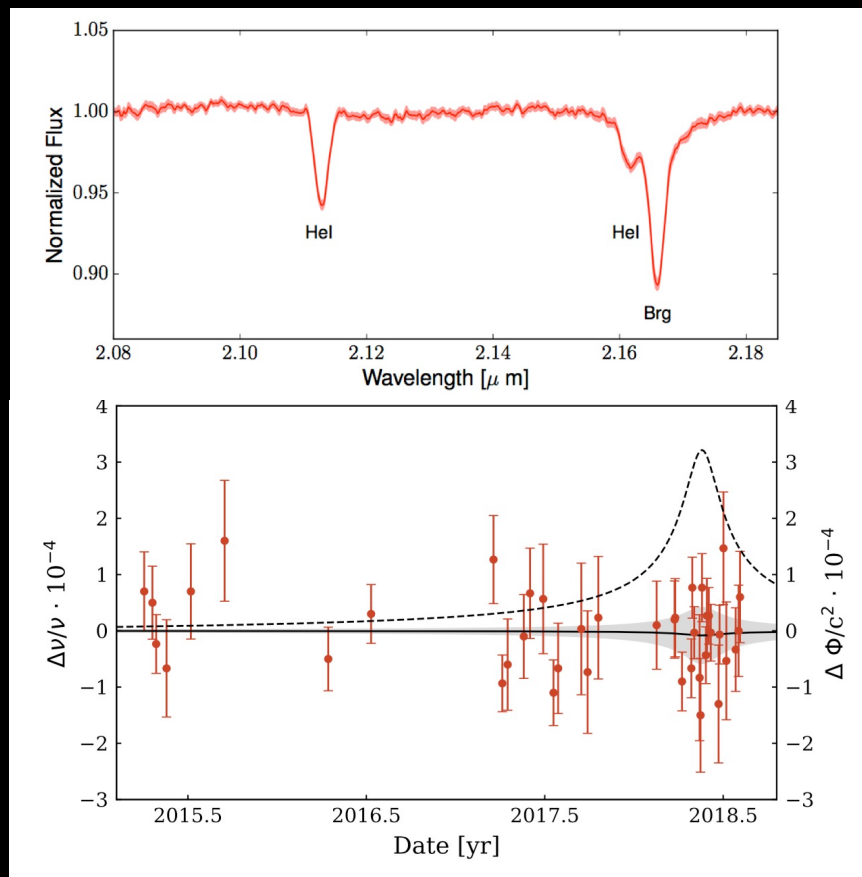
At closest approach of 120 AU, S2 reached 8000 km/s

- **transverse Doppler effect**
- **gravitational redshift from Sgr A***

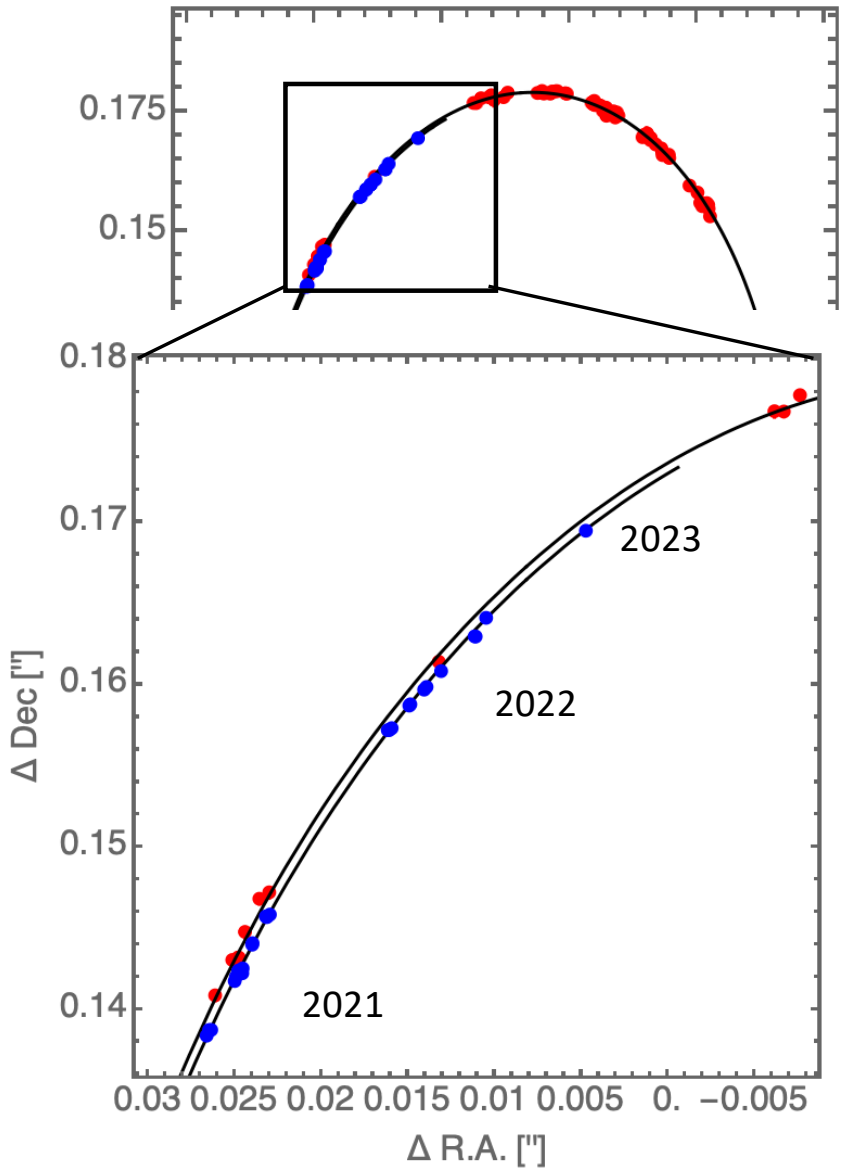
GRAVITY Coll. 2018a



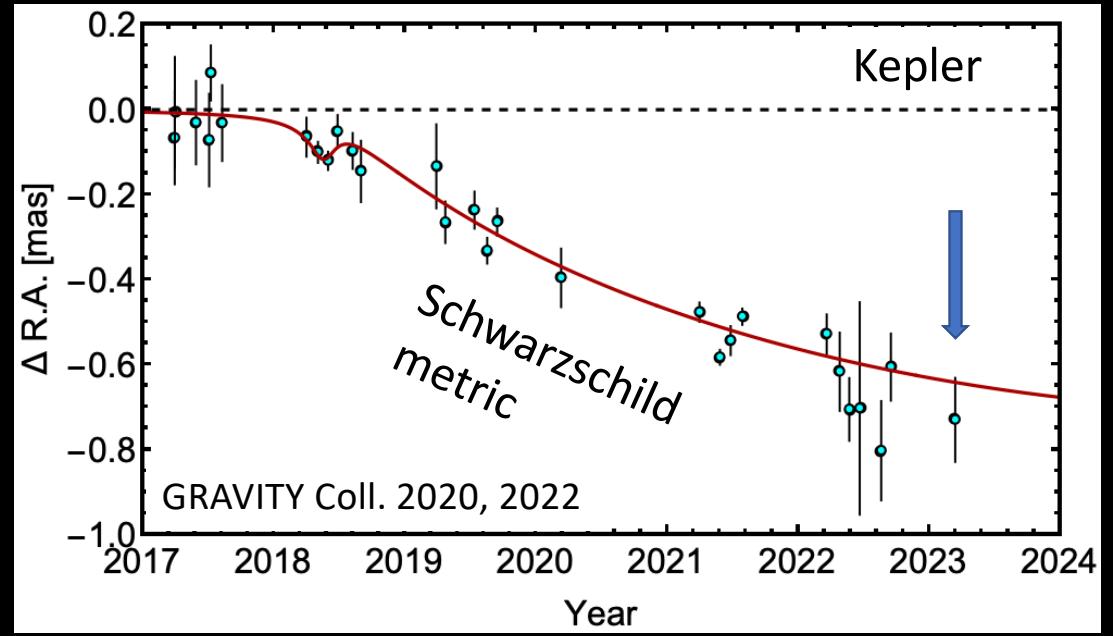
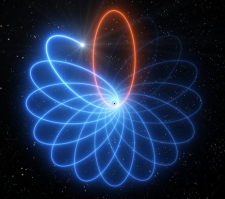
Comparing the redshift for two atomic lines tests the equivalence principle



GRAVITY Coll. 2019

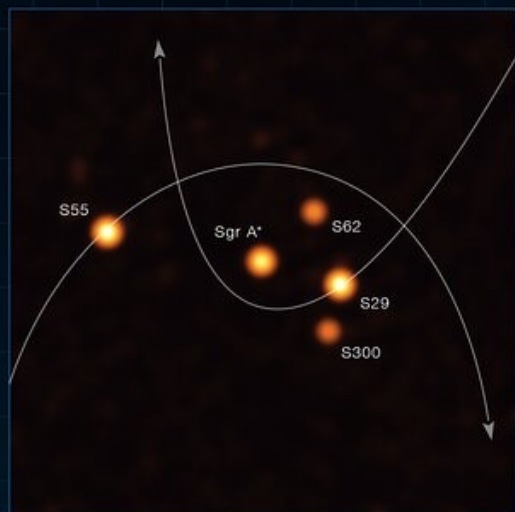


The Schwarzschild precession is detected at $>10\sigma$



New since 2021: GRAVITY astrometry of more stars (S29, S38, S42, S55, ...)

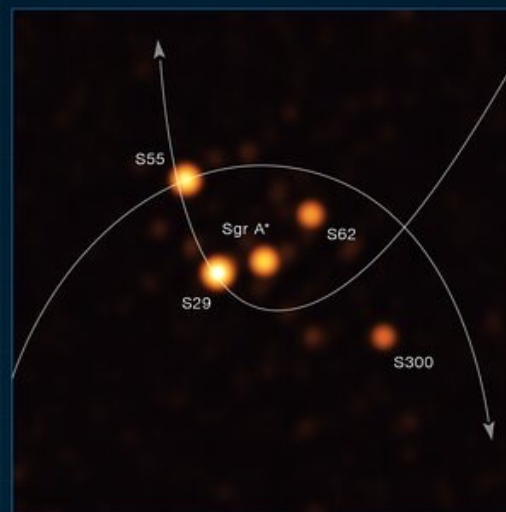
March



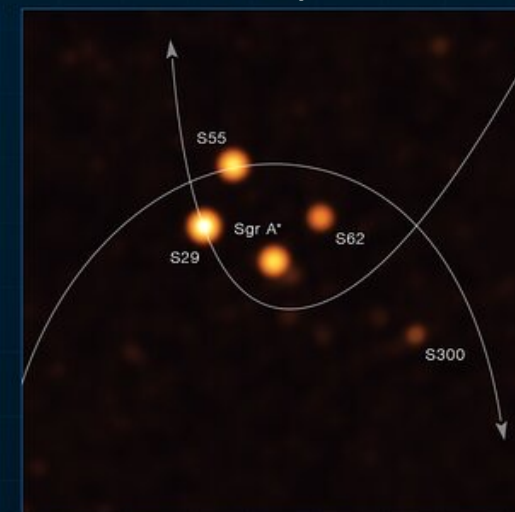
May

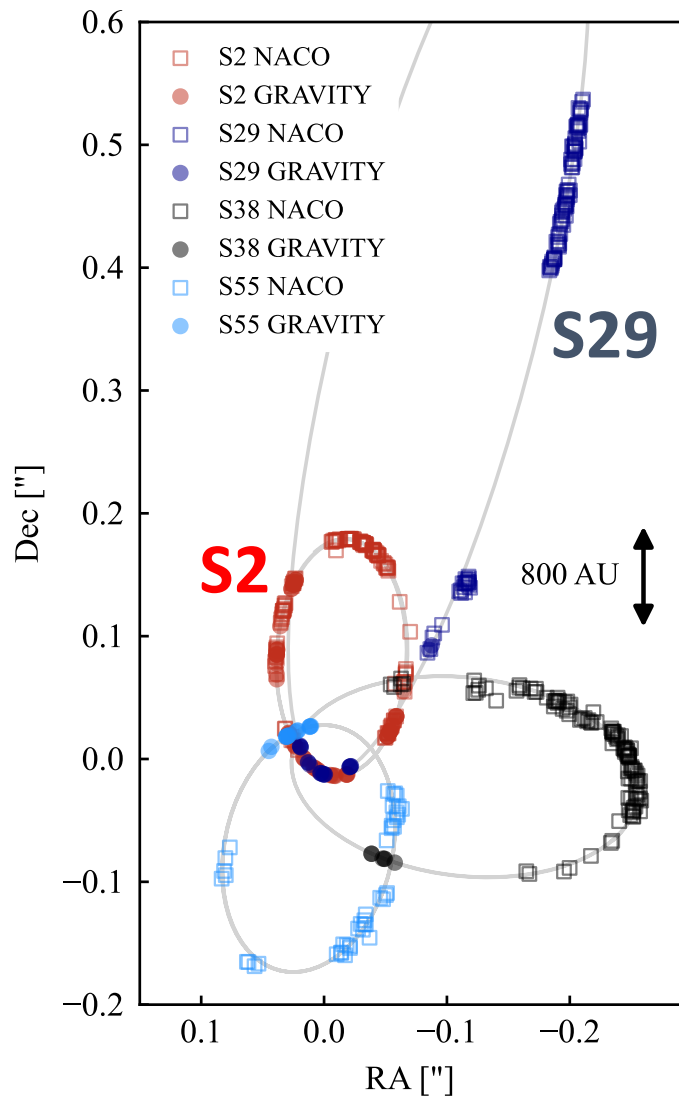


June

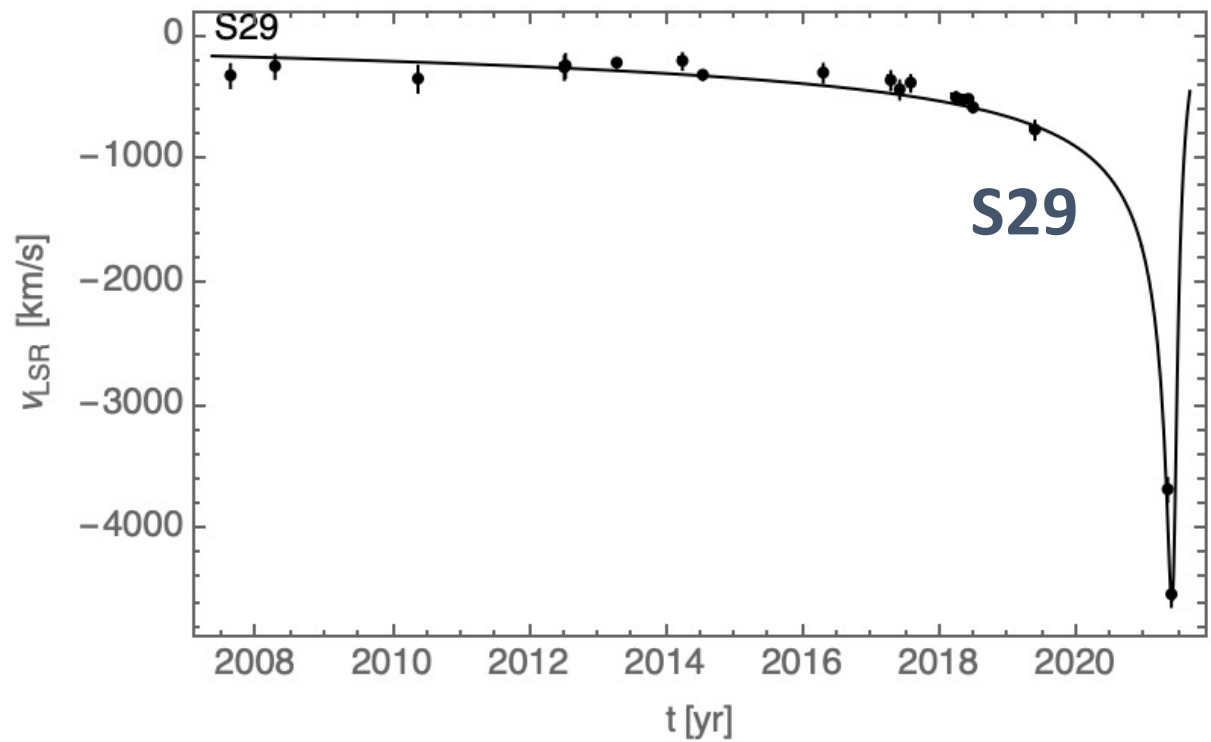


July

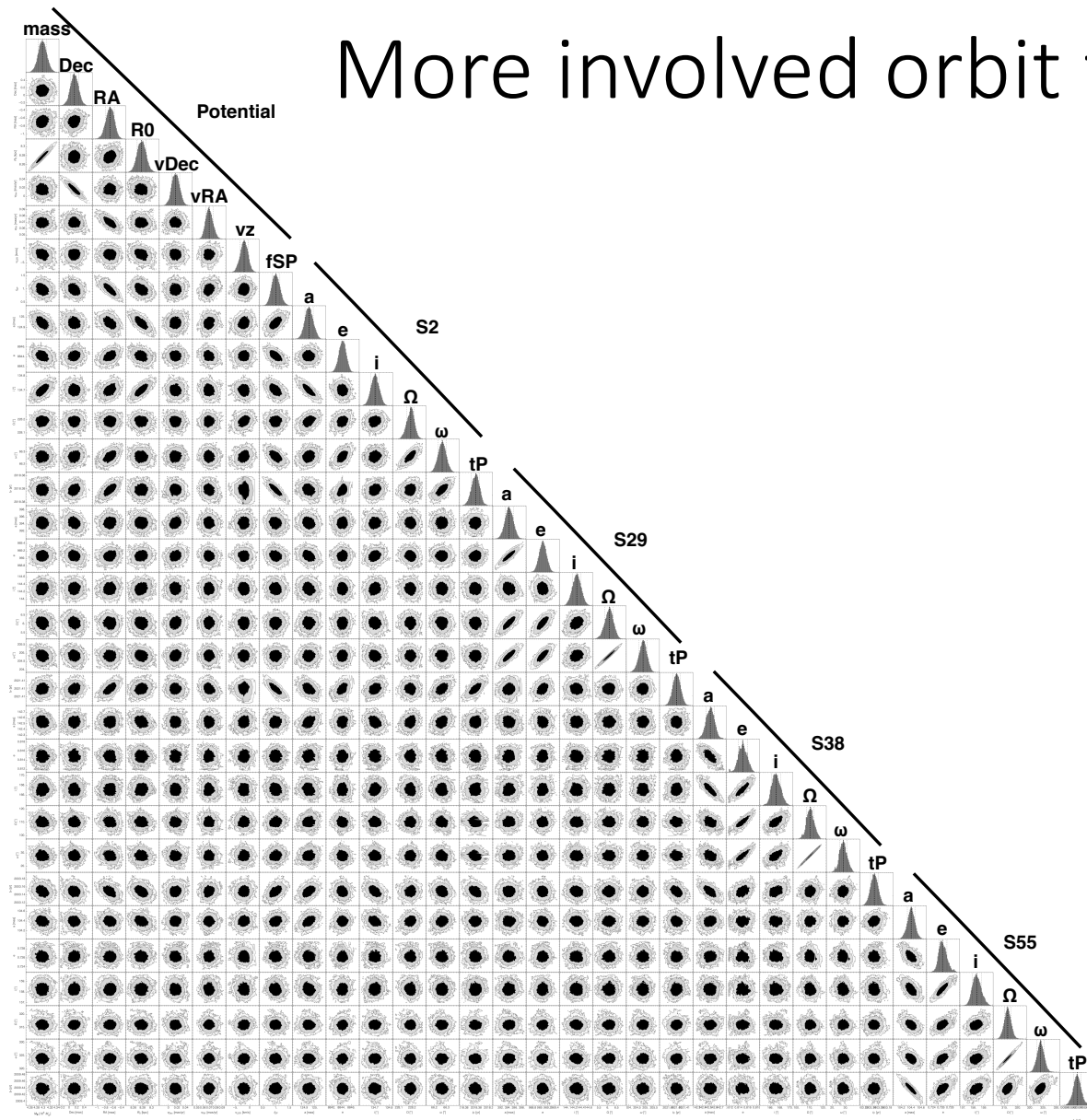




We had previous coverage for these stars

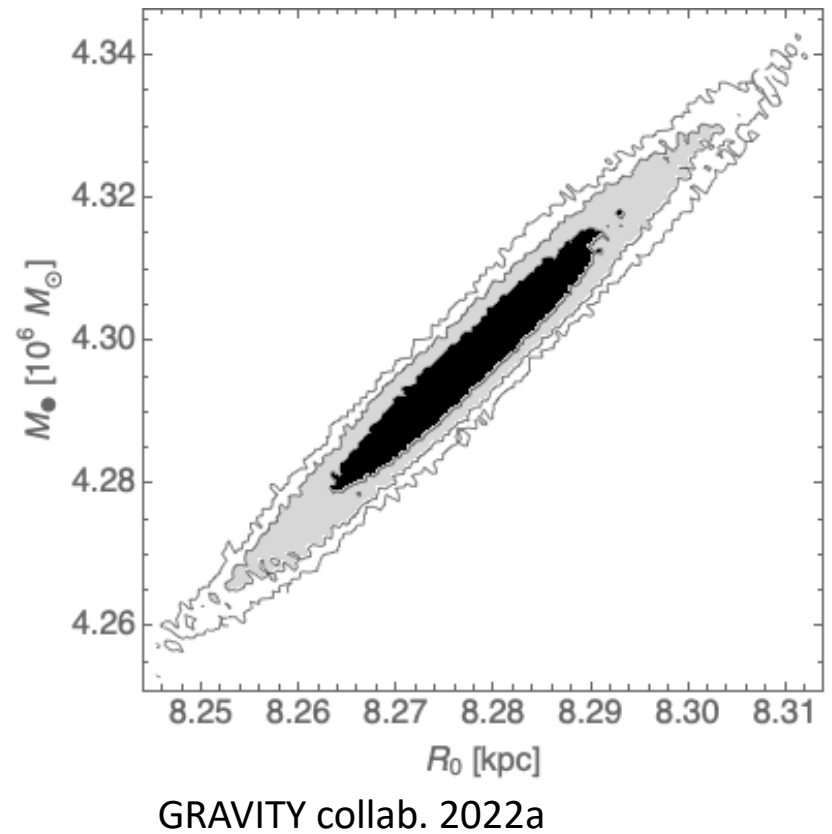


More involved orbit fit...



$$R_0 = 8.277 \pm 0.009_{\text{stat}} \pm 0.035_{\text{sys}} \text{ kpc}$$

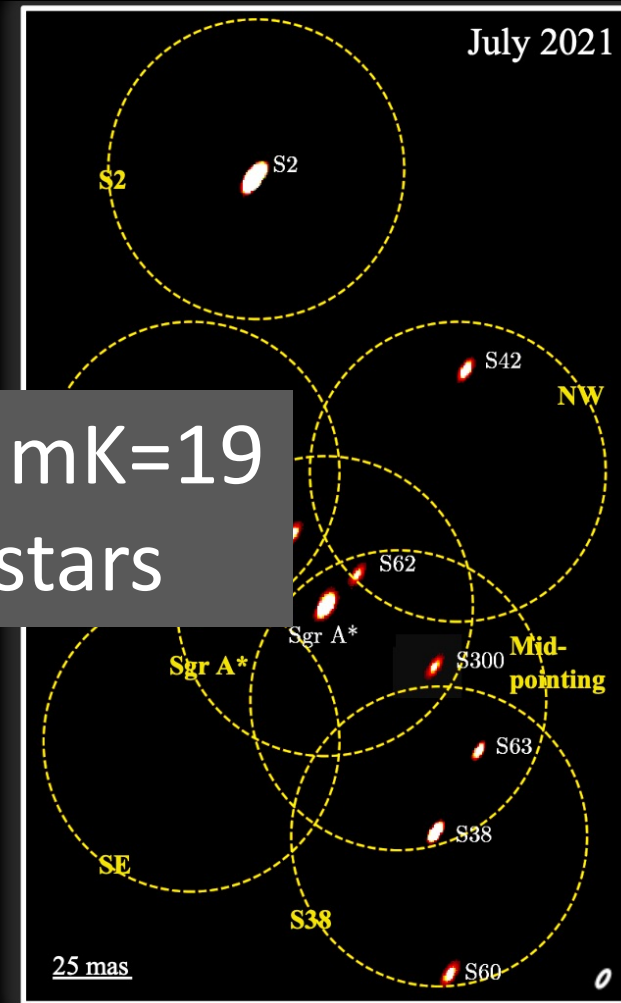
$$M = 4.297 \pm 0.012_{\text{stat}} \pm 0.040_{\text{sys}} \times 10^6 M_{\odot}$$



Imaging at milli-arcsec resolution

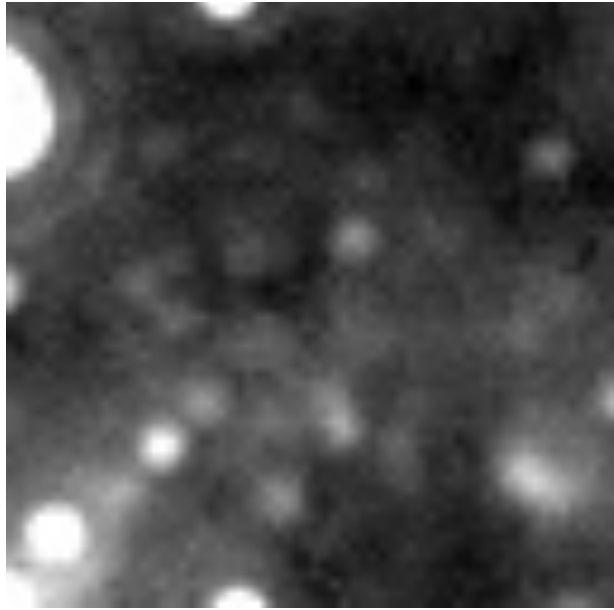


- Tracking stars down to $mK=19$
- Discovering new, faint stars

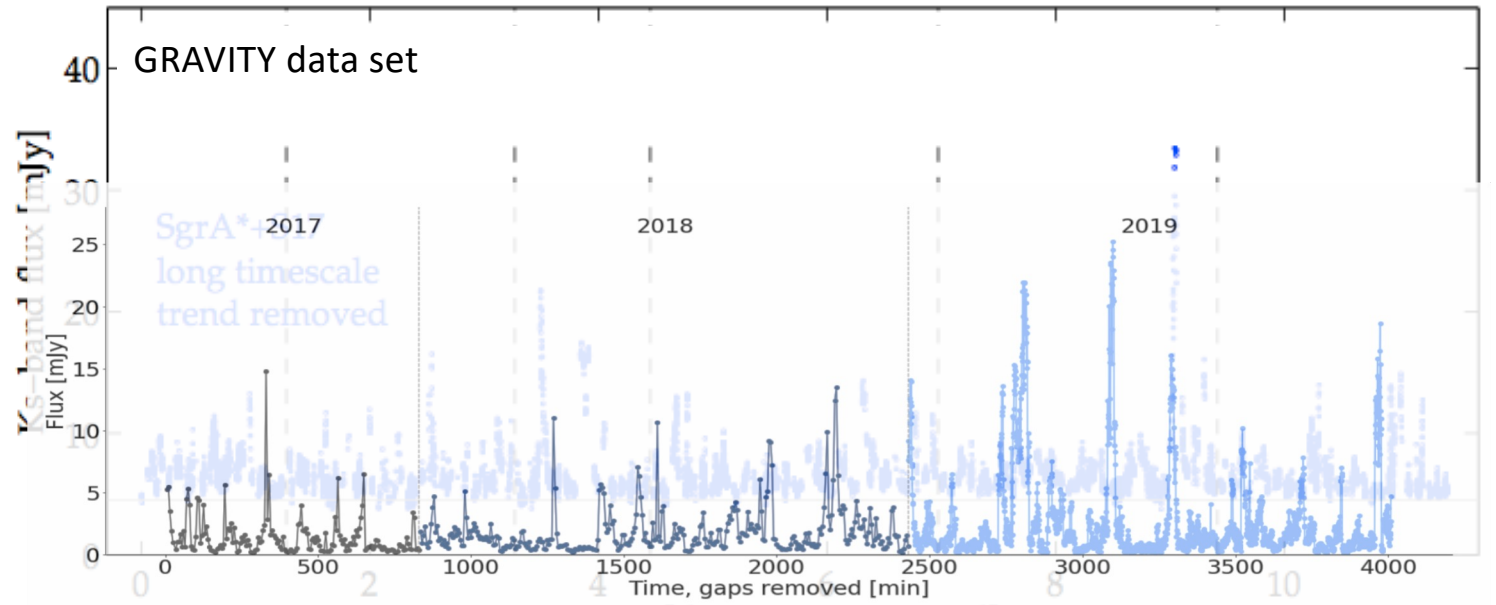


Outline

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 - The black hole nature of Sgr A*
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 - Testing General Relativity in the Galactic Center
 - SgrA* flares
 - (A funny gas cloud)

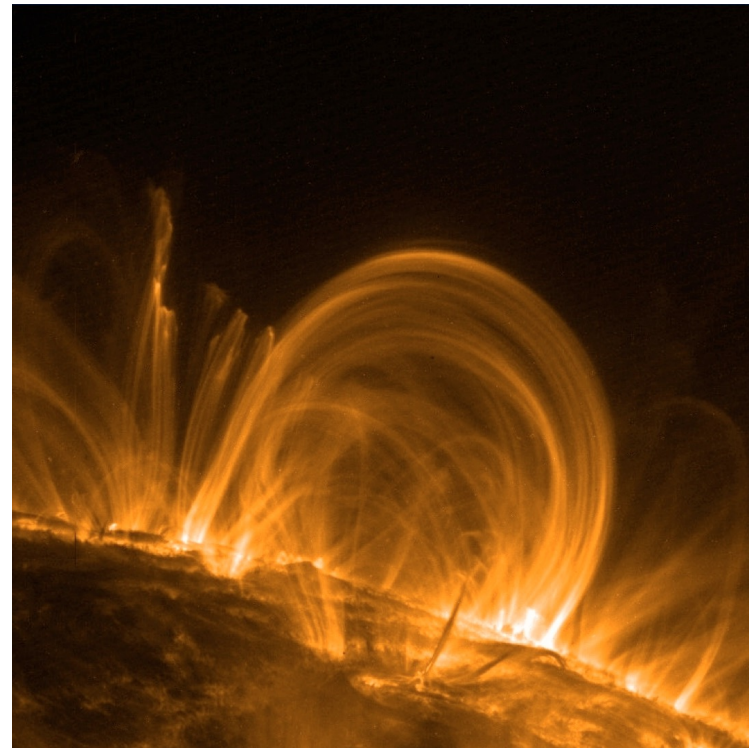
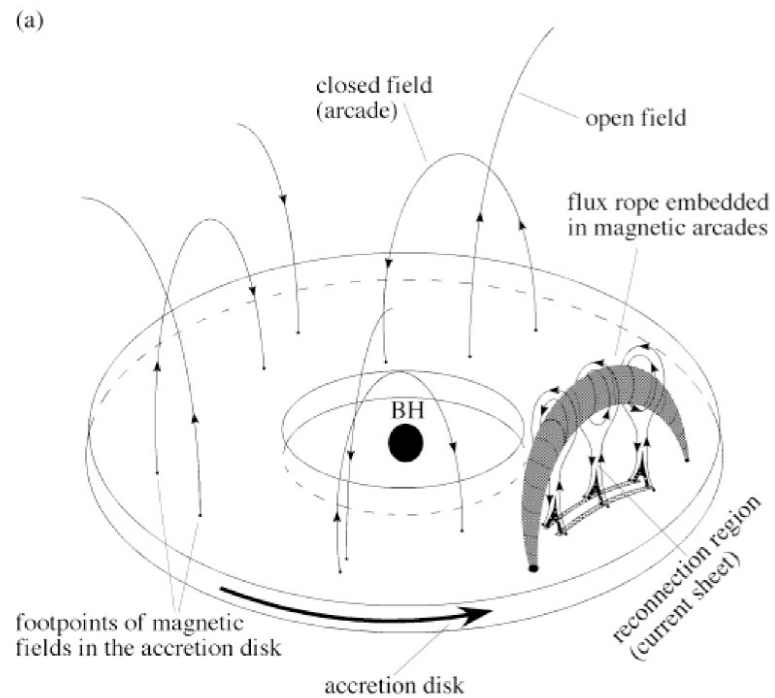


Sgr A* is flaring in the infrared



Genzel et al. 2003, Dodds-Eden et al. 2011, Witzel et al. 2012, 2018, GRAVITY coll. 2020b

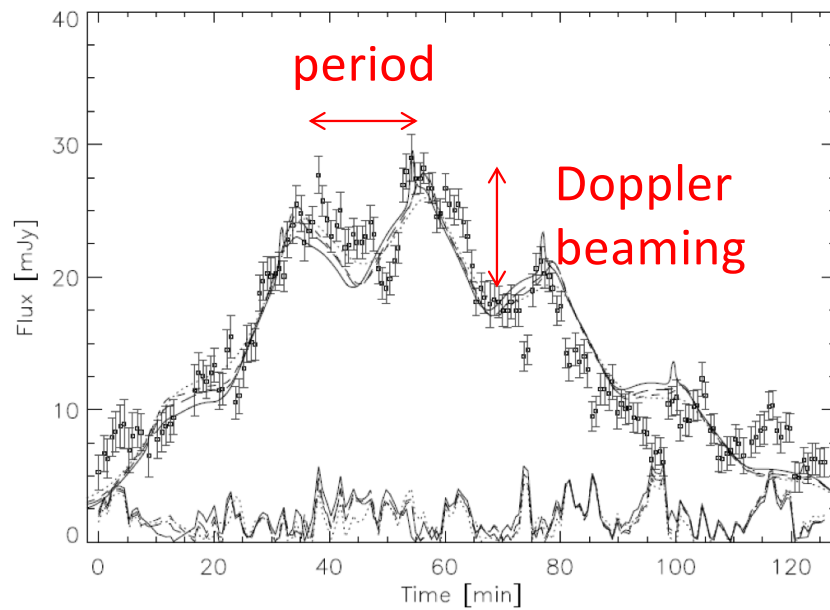
Speculation: An analog to solar flares



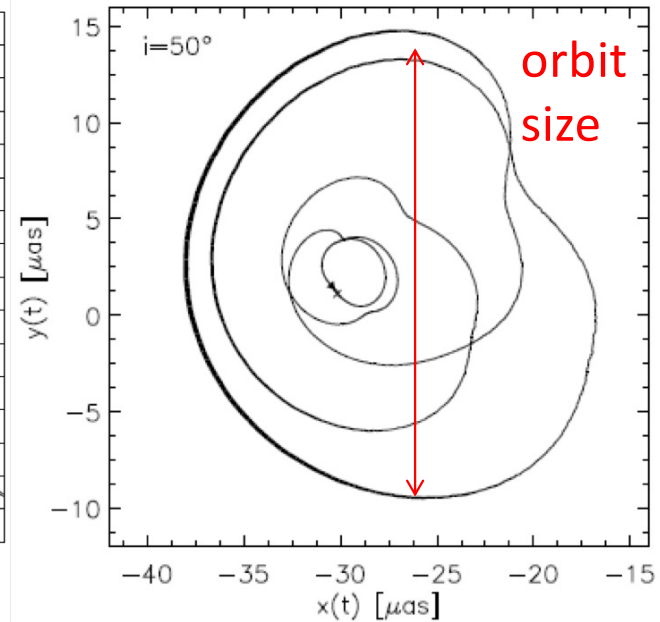
Yuan et al. 2009

What can one learn from flare orbits?

Light curve



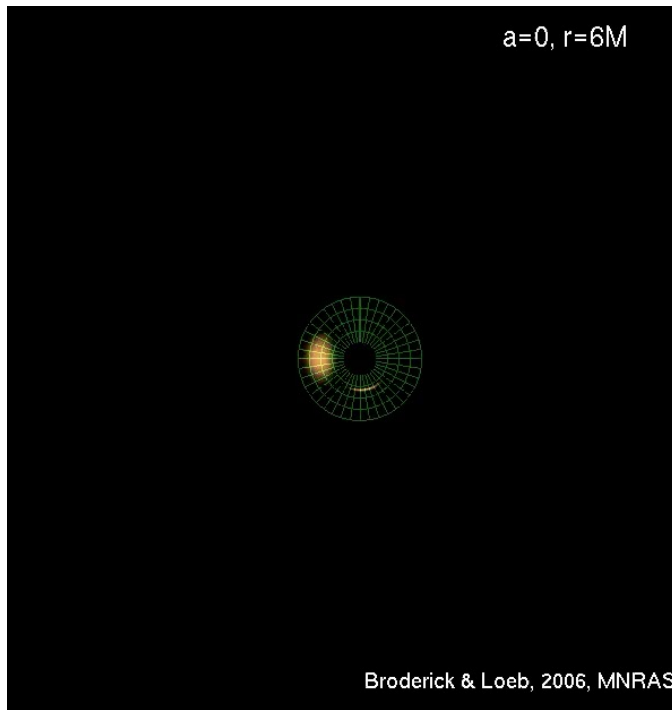
Centroid track



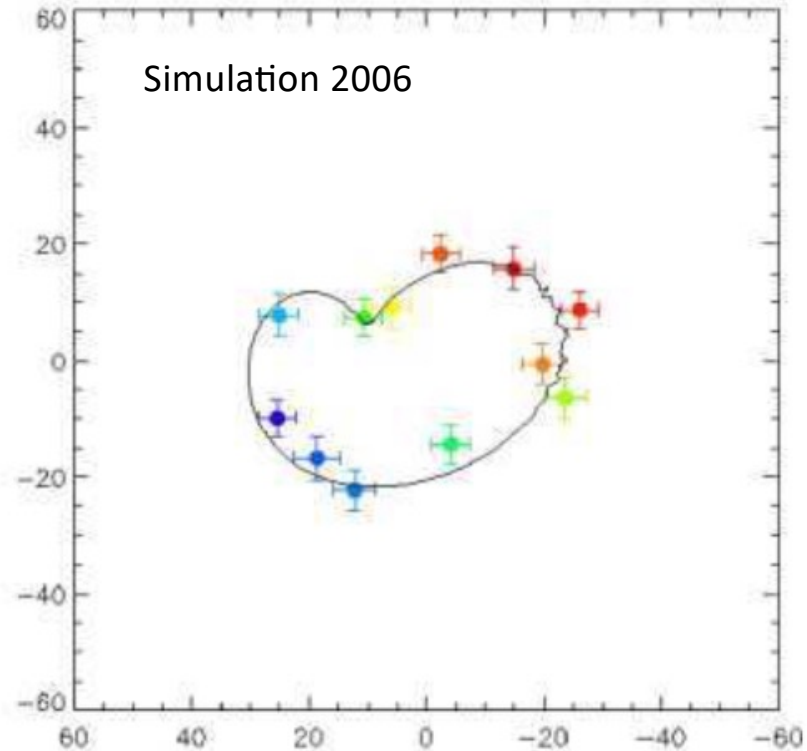
Hamas et al. 2009

→ Radius, Inclination, Spin

The flares should orbit at a few $10\mu\text{as}$ radius -
in reach for GRAVITY astrometry

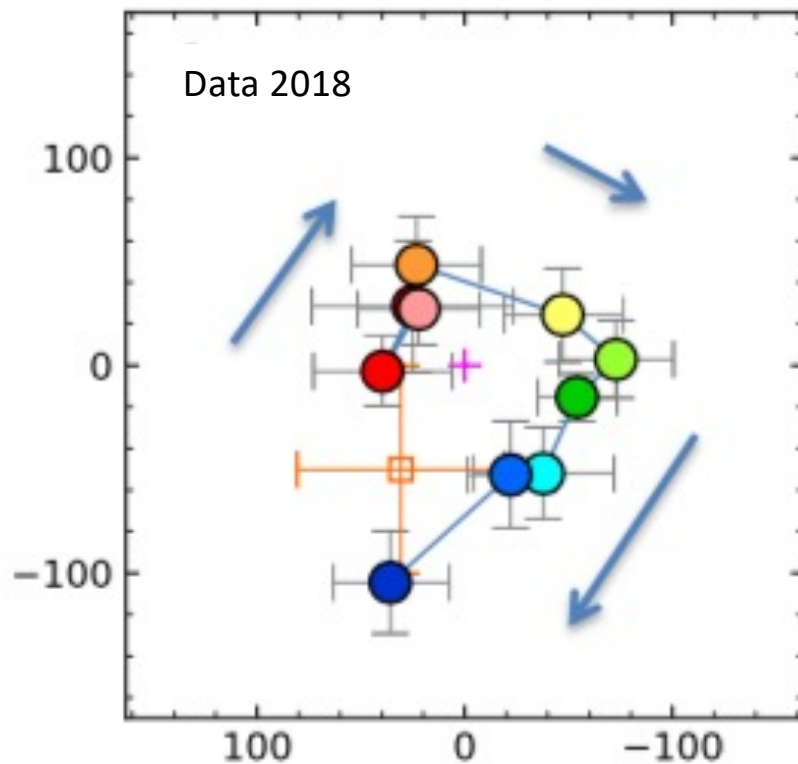


Broderick & Loeb 2005, 2006

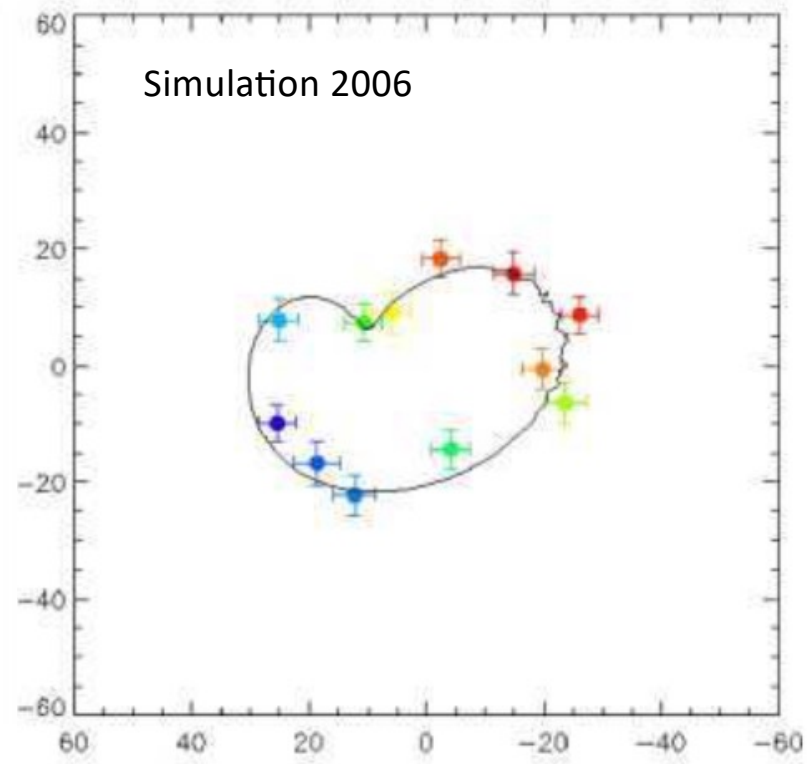


GRAVITY science proposal 2006

GRAVITY observations 2018: Orbiting hot spots (4x)



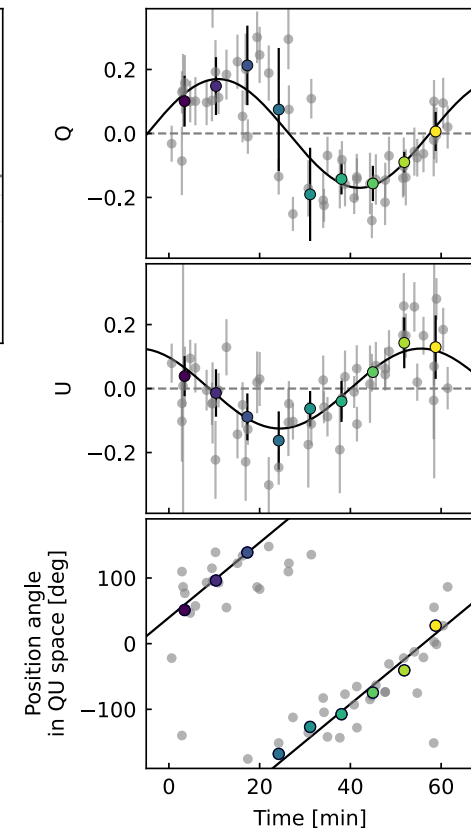
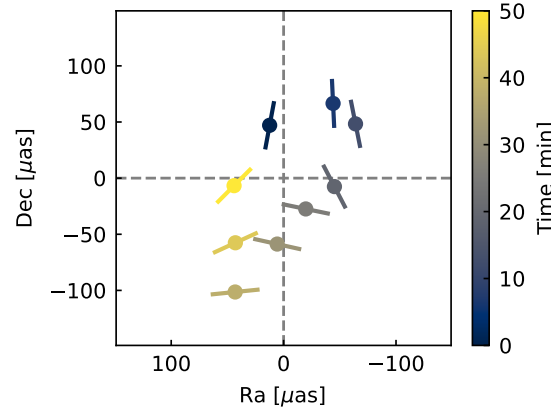
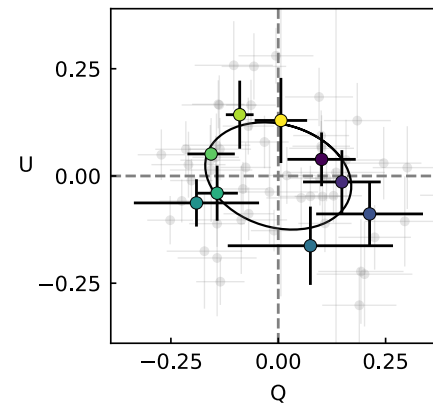
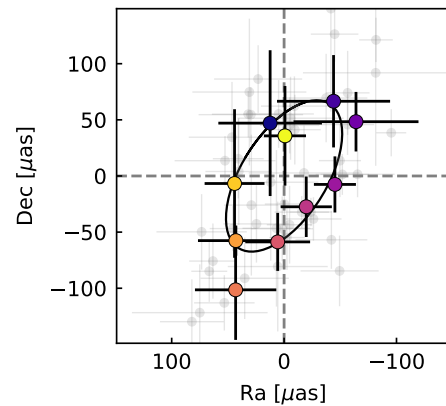
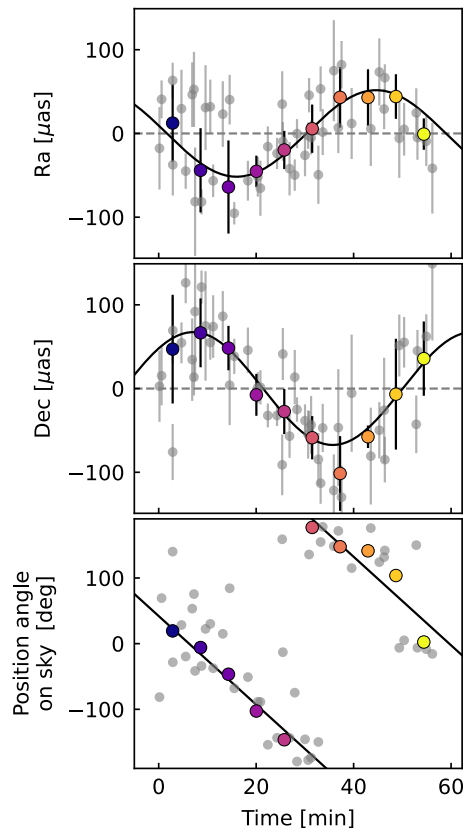
GRAVITY collaboration 2018b



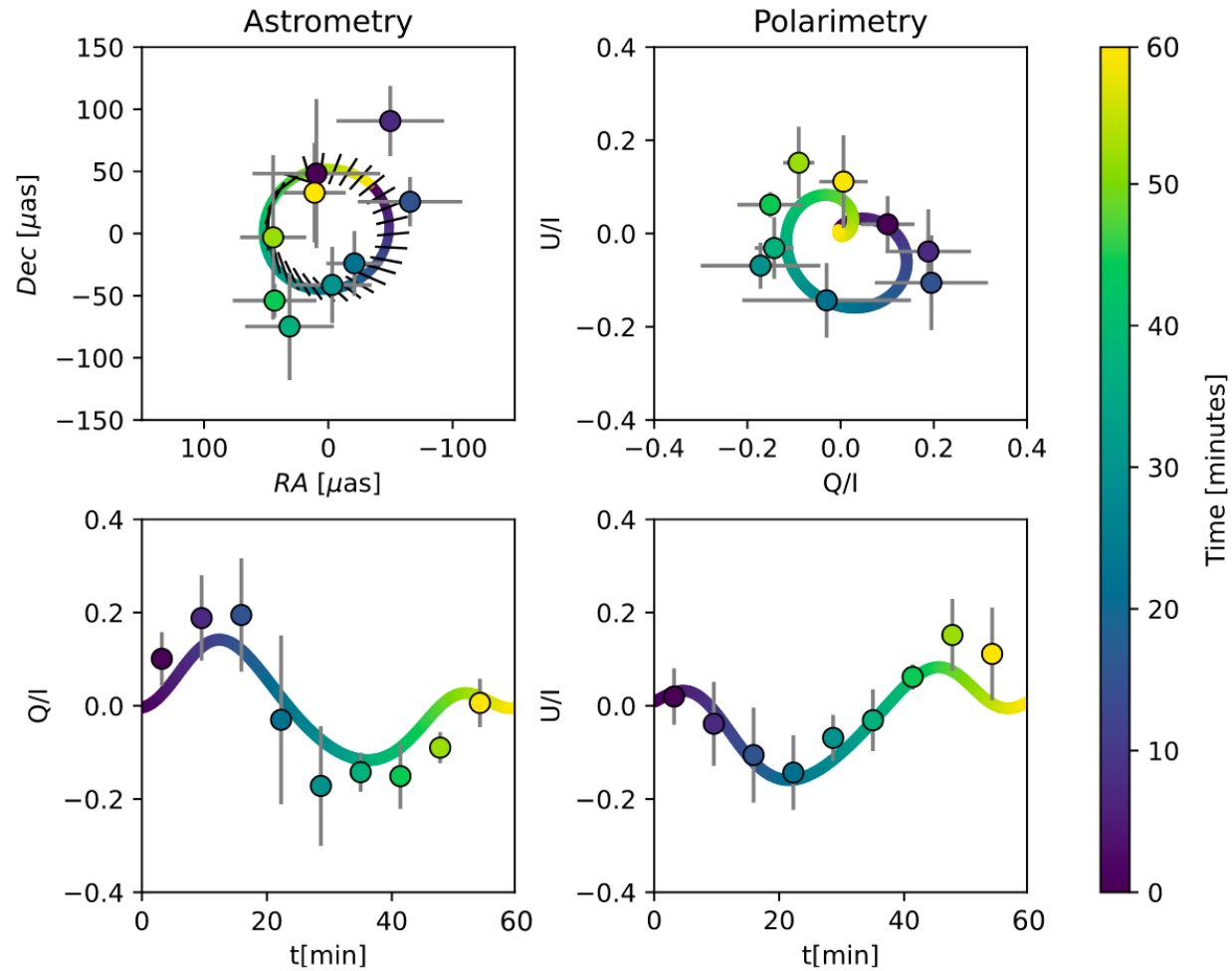
GRAVITY science proposal 2006

Status 2023:

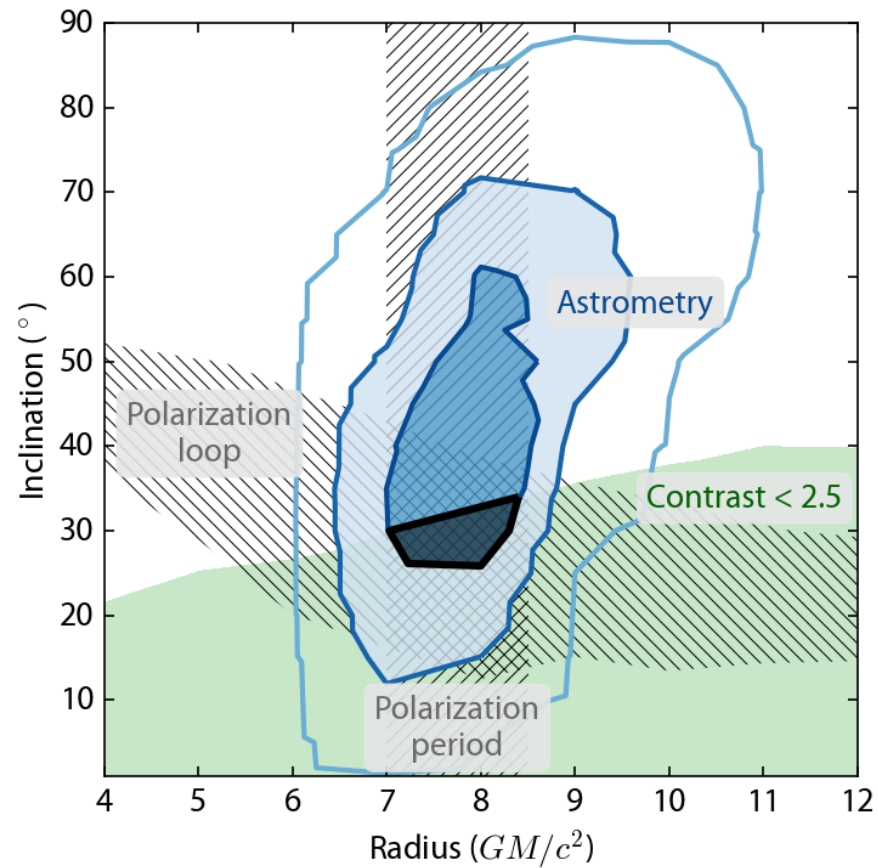
4 flares with astrometric loop
6 flares with polarimetric loop



An orbiting, polarized hot spot can describe the data

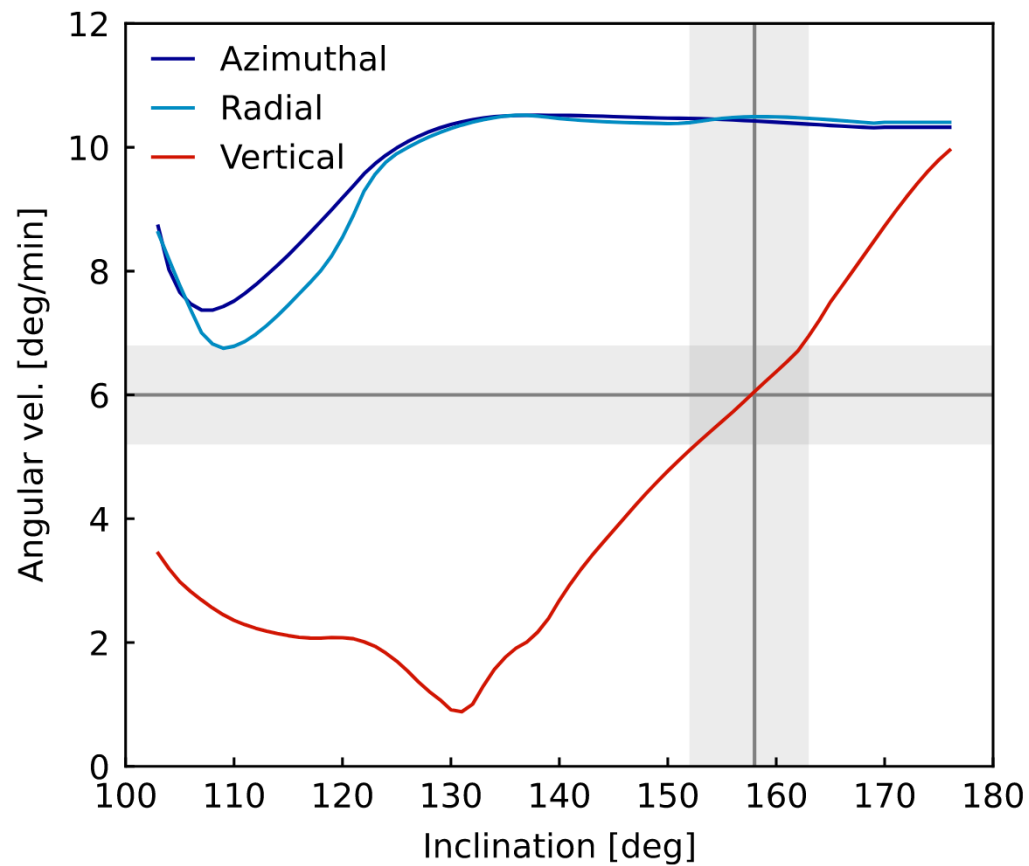


A face-on view and a vertical magnetic field for SgrA* is favored

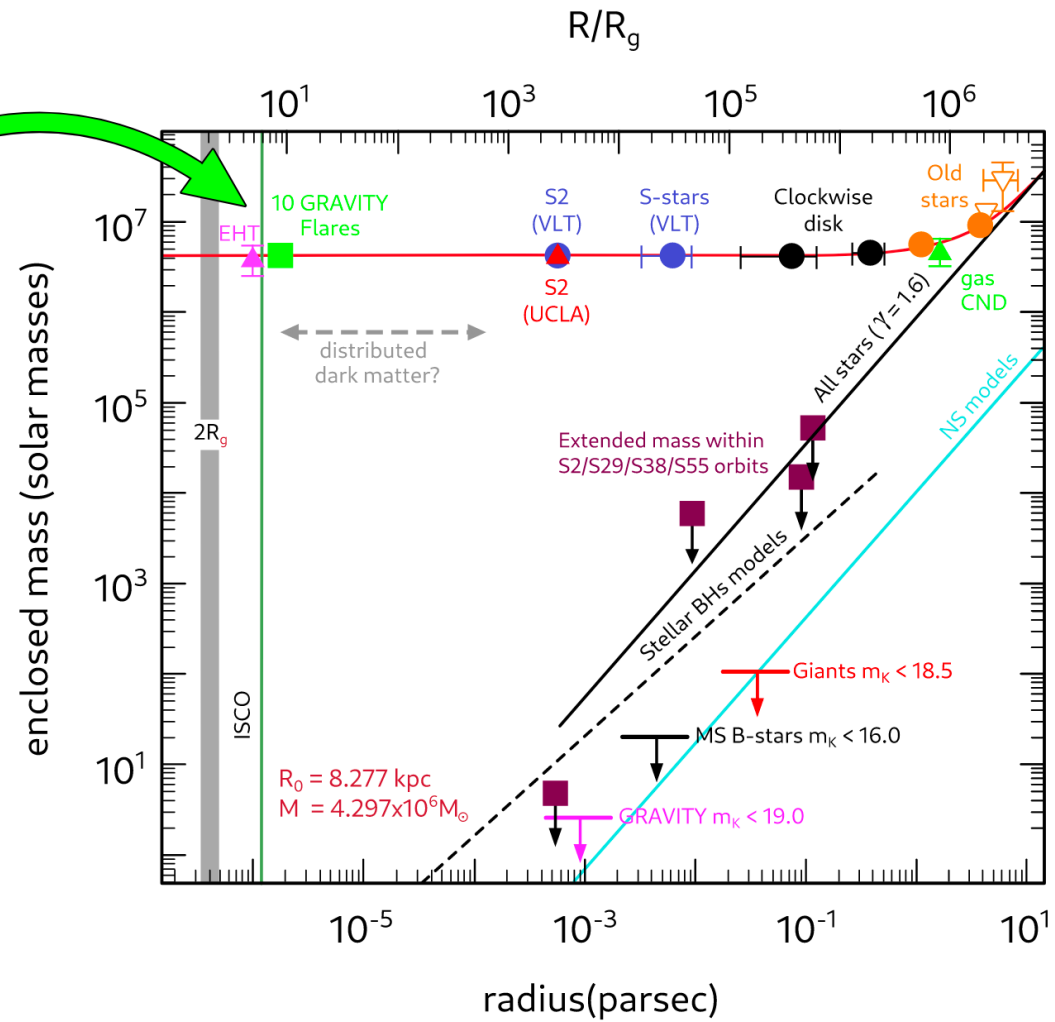
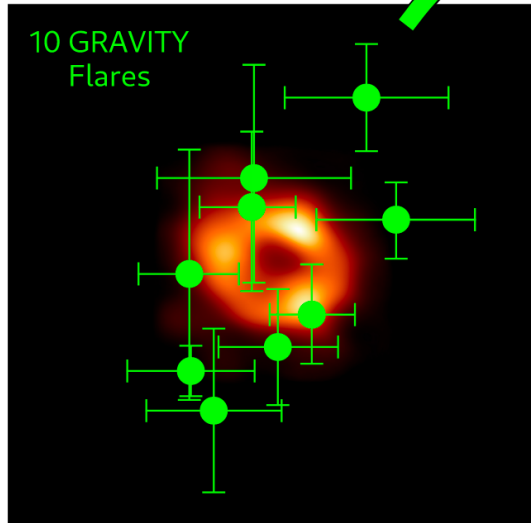


GRAVITY collaboration 2018b

Key for the inclination: Angular velocity of polarization vector rotation

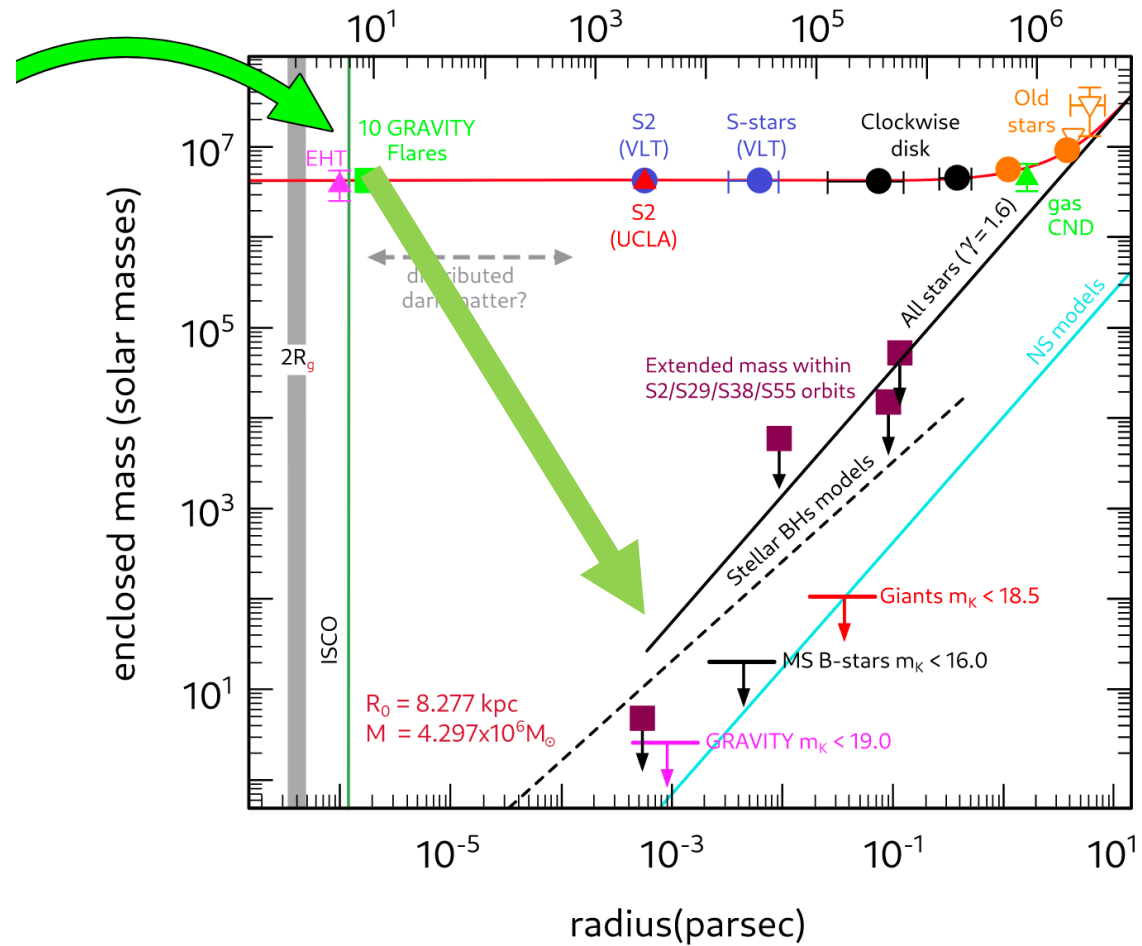


Period known, radius known \rightarrow mass known

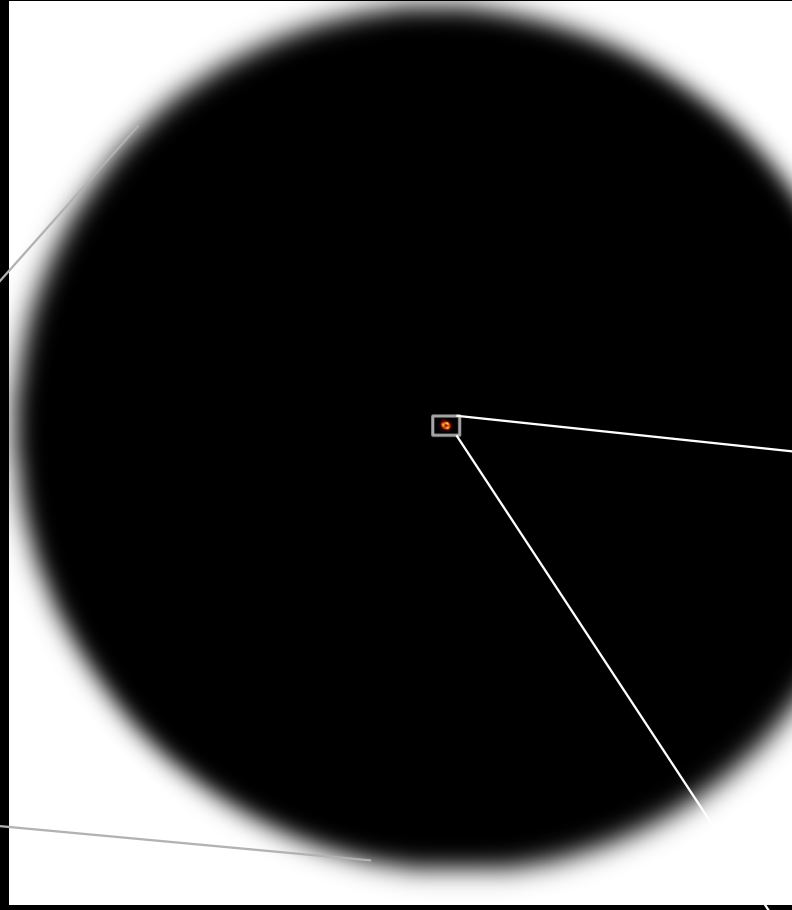
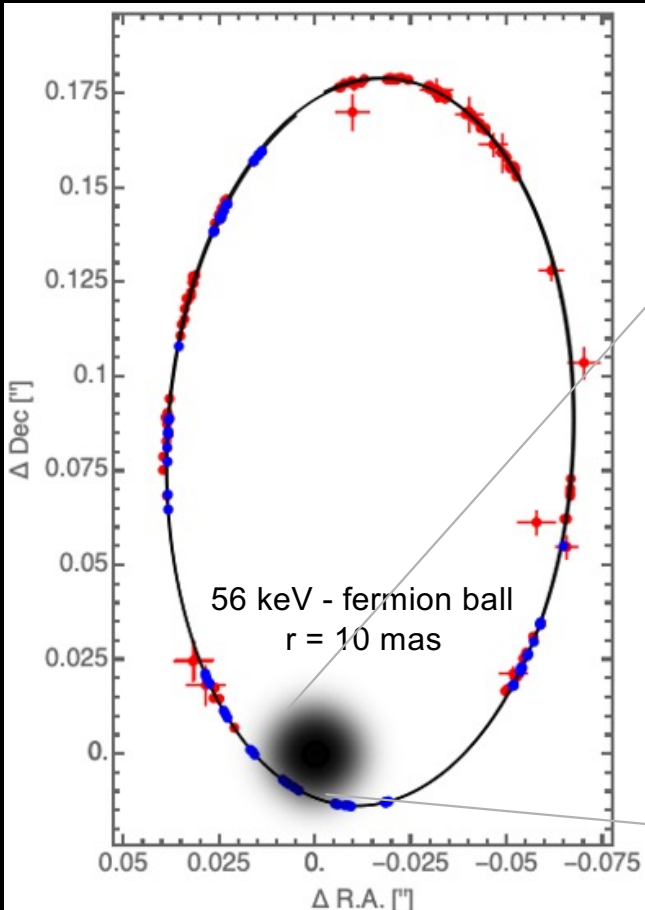


99.9% of the mass of Sgr A* is
inside of S29's apo = 100 AU

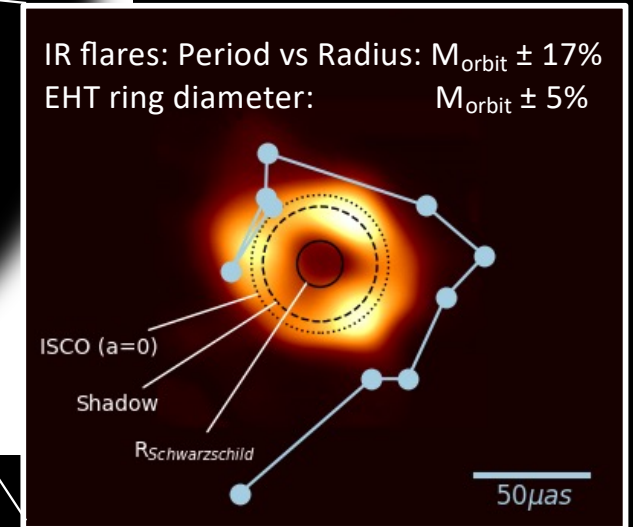
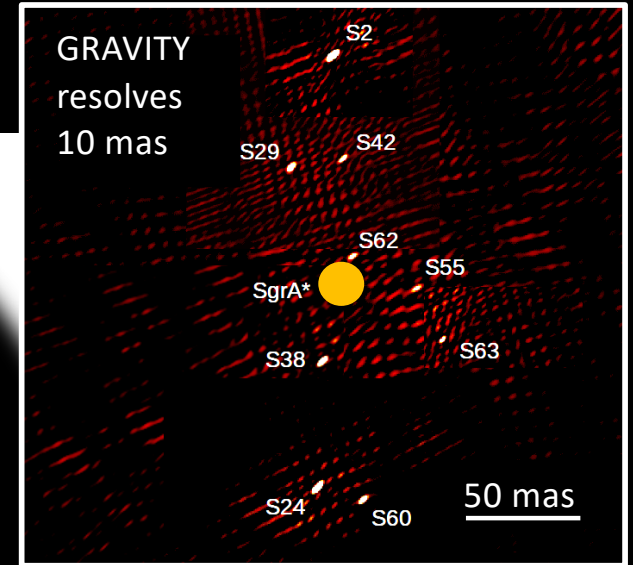
At most
~ 0.1% ~ 4000 M_{\odot}
can be in an extended
configuration



Example of a model completely ruled out: Fermion ball



Argüelles et al. 2019



GRAVITY- starting to set limits

Gravitational redshift f_{RS}



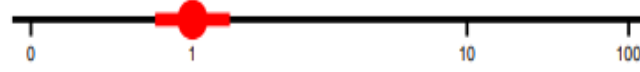
$$f_{RS} = 1.04 \pm 0.04$$

Equivalence principle $1+\Delta\beta$



$$\Delta\beta = 2.4 \pm 5.1 \%$$

Schwarzschild precession f_{SP}



$$f_{SP} = 1.09 \pm 0.13$$

P_{flare}/P_{Kepler}



White hole metric α_{WH}/R_G



Black hole spin X_{spin}

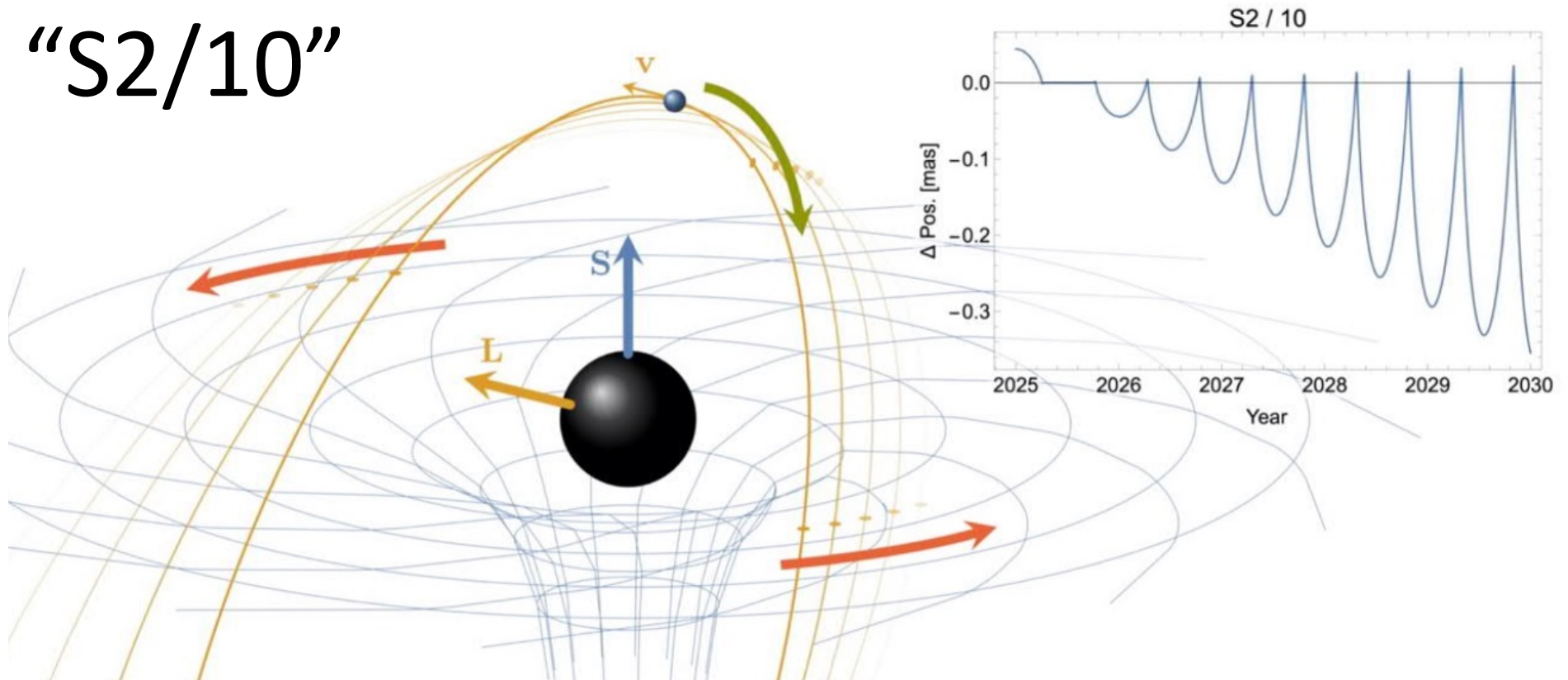


2PN term f_{2PN}

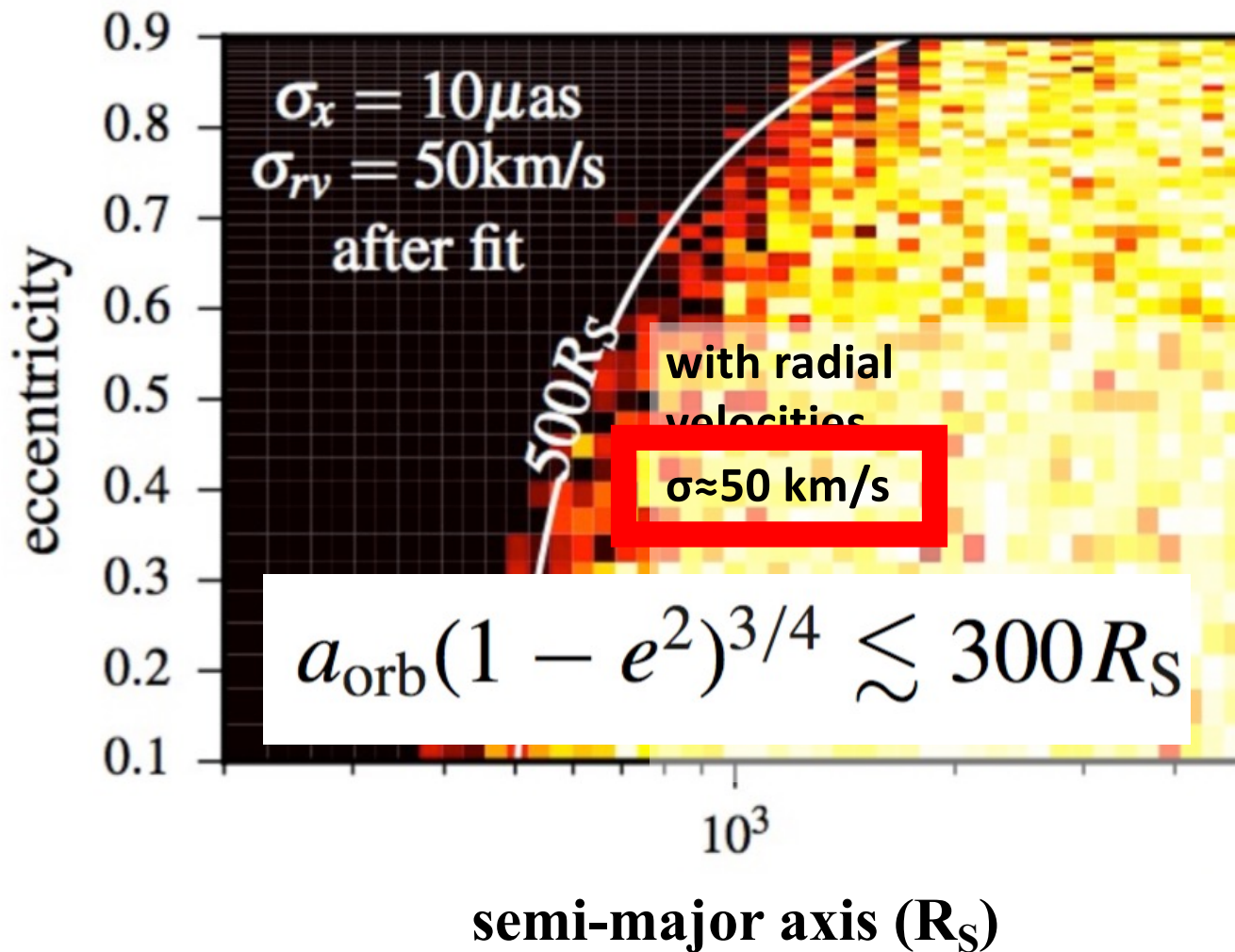


How far can we get in the Galactic Center?

“S2/10”



What star is needed to detect the spin?



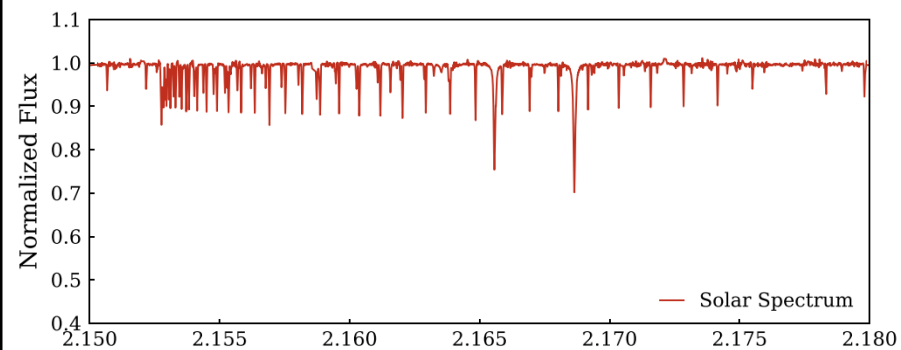
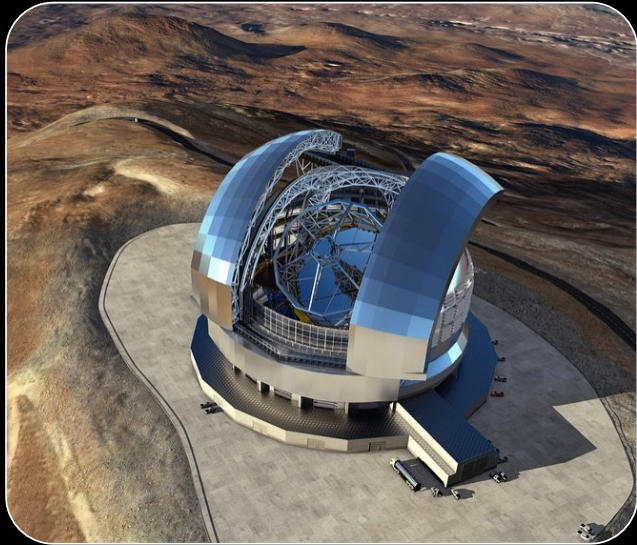
Spin effects: r^{-3}
→ stars further in
→ astrometry harder,
spectroscopy easier

- very conservative radial velocity error
- we don't need to limit us to one star

Waisberg et al. 2018

Medium-term goal: The spin of Sgr A*

ELT Spectroscopy

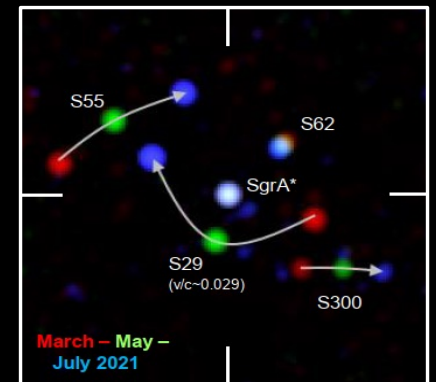


Deep Interferometric Imaging

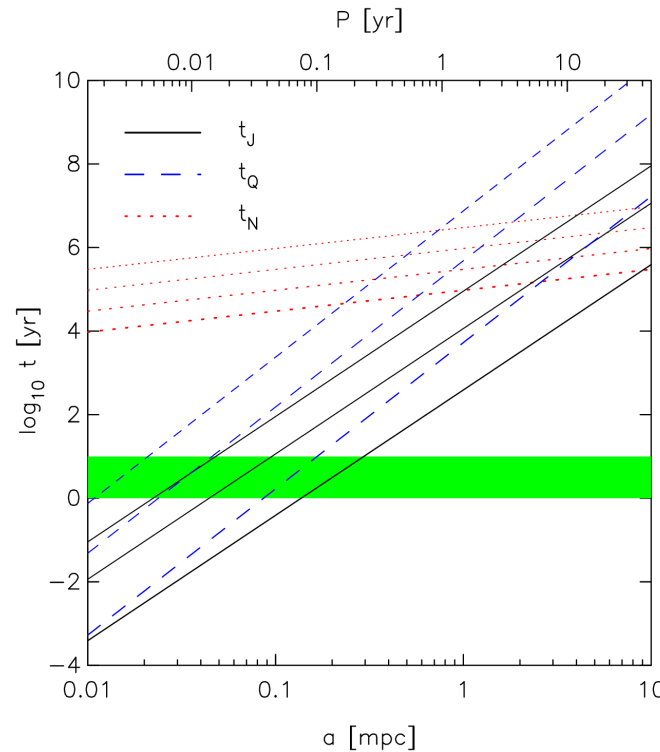
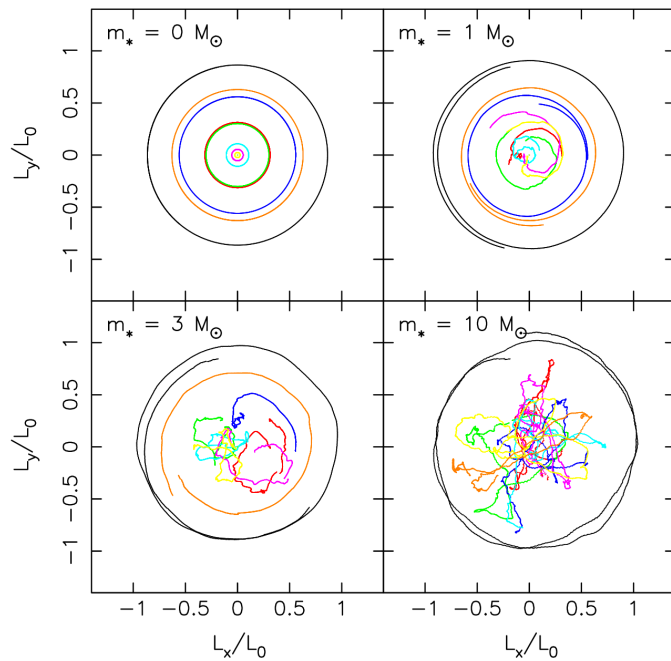


$$\Delta v < 1 \text{ km/s}$$

$$mK > 21$$



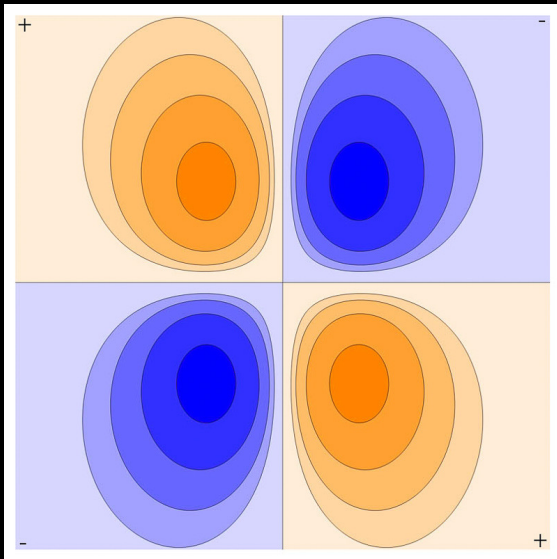
Theoretical arguments were too pessimistic



- arguments based on time-averaged changes
- but observers see “abrupt” changes
- “if a change happens at pericenter, it was the MBH”

Merritt+ 2010

Longer term goal: No-hair theorem



Will 208

$$a = -\frac{Mx}{r^3} + \frac{Mx}{r^3} \left(4 \frac{M}{r} - v^2 \right) + 4 \frac{Mr}{r^2} v$$

$$- \frac{J}{r^3} [2\mathbf{v} \times \hat{\mathbf{J}} - 3i\mathbf{n} \times \hat{\mathbf{J}}]$$

$$+ \frac{3}{2} \frac{Q_2}{r^4} [5\mathbf{n}(\mathbf{n} \cdot \hat{\mathbf{J}})^2 - 2(\mathbf{n} \cdot \hat{\mathbf{J}})\hat{\mathbf{J}} - \mathbf{n}],$$

In general relativity:
 $Q_2 = -J^2/M$

Quadrupole moment:
 observational quantity

Rad err [km/s]	Schwarz- schild	Lense- Thirring	Lensing & Delay	Quad- rupole
10	0.6%	3%	-	-
1	0.4%	0.7%	35%	-
0.1	0.02%	0.04%	4%	11%

State of the art AO

Vibration control

Laser guide stars for all telescopes

Full sky coverage for
adaptive optics

reach $mK=22$

Full sky coverage for
fringe-tracking

Wide field off-axis fringe tracking

Grism upgrade



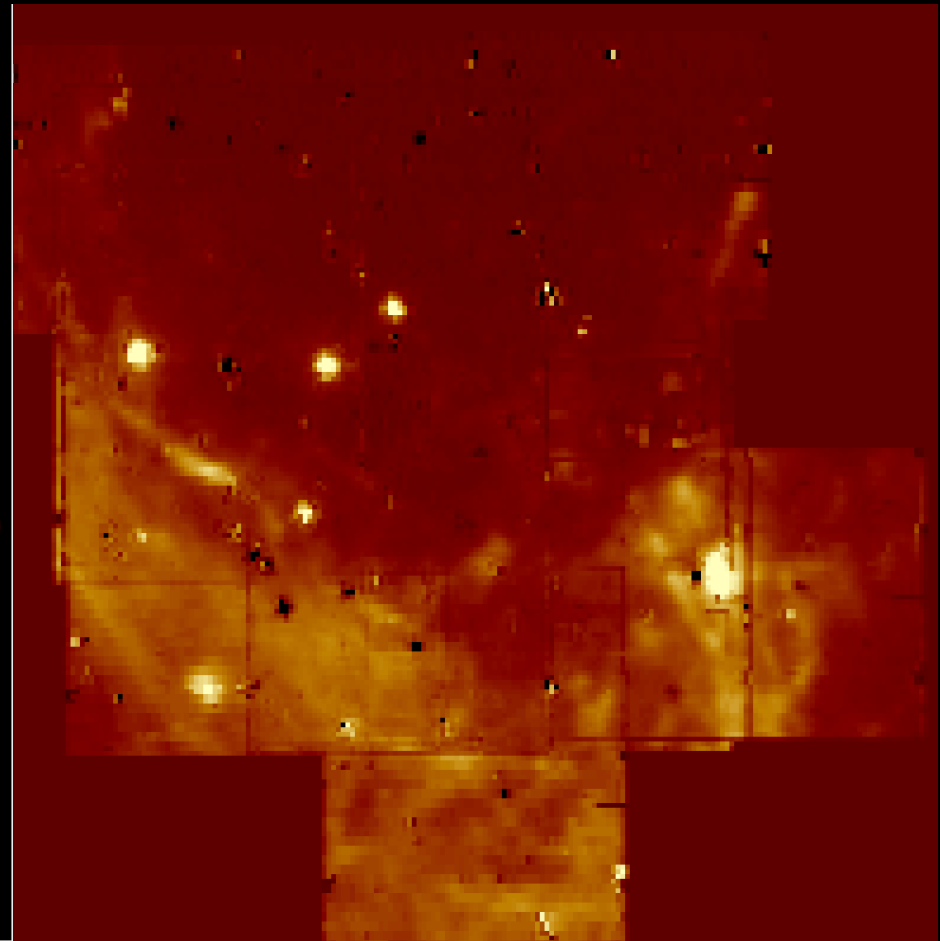
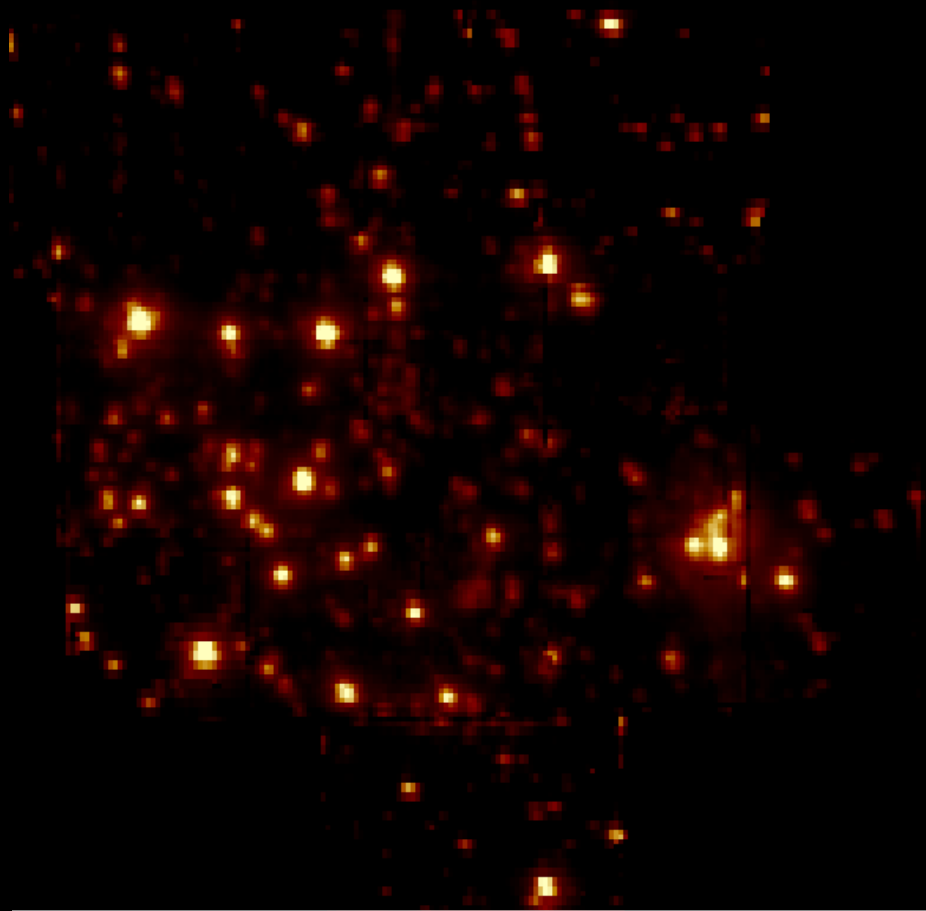
GRAVITY+ Collaboration

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 - SgrA* flares
 - A funny gas cloud

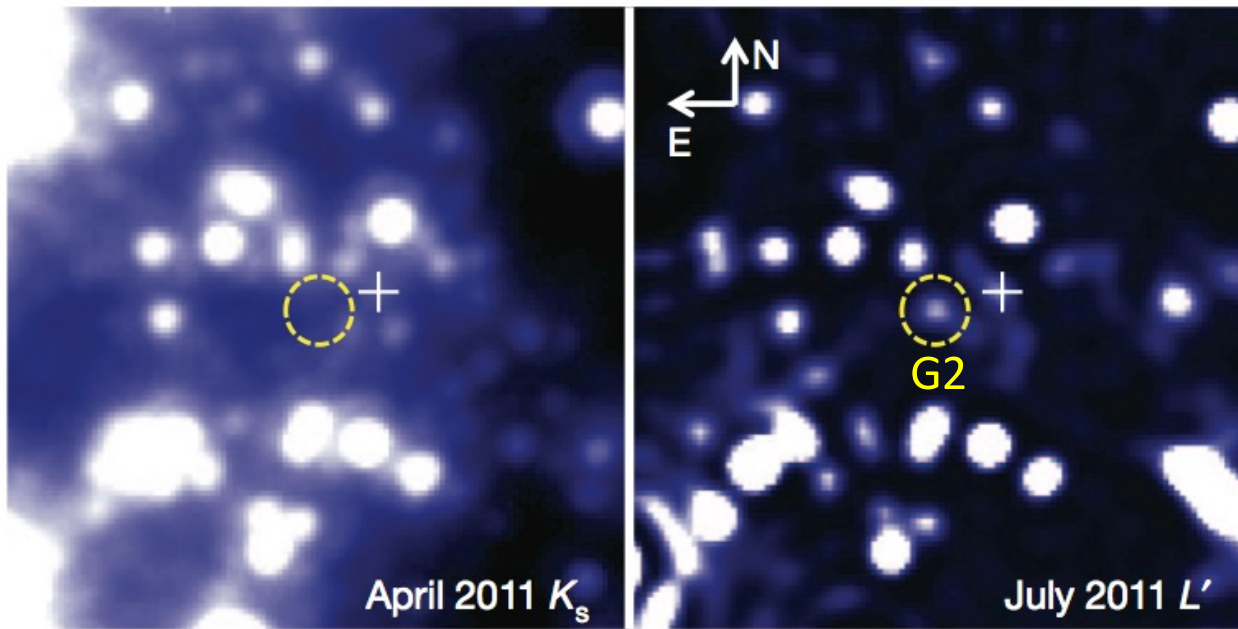
The Galactic Center is gas rich

ERIS, 100mas, 2.166 μ m, R=11200



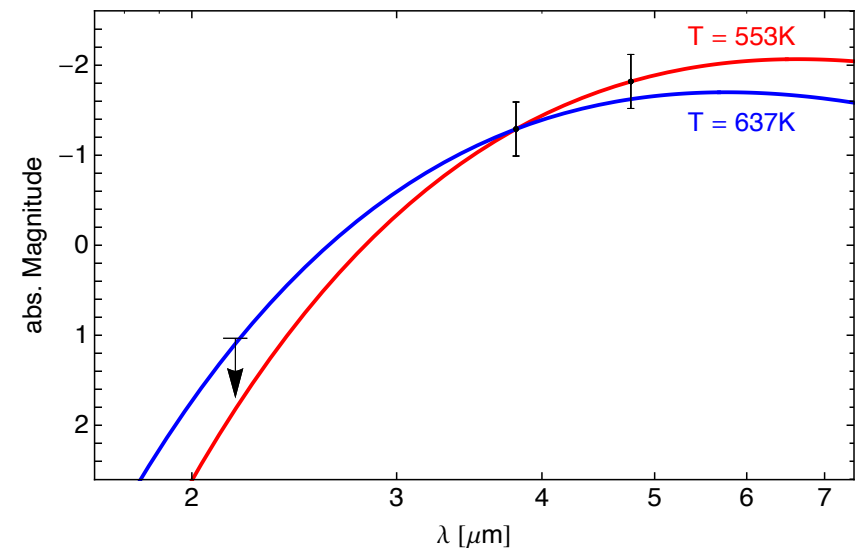
A cool object: G2 - $T = 600$ K

Gillessen et al. 2012



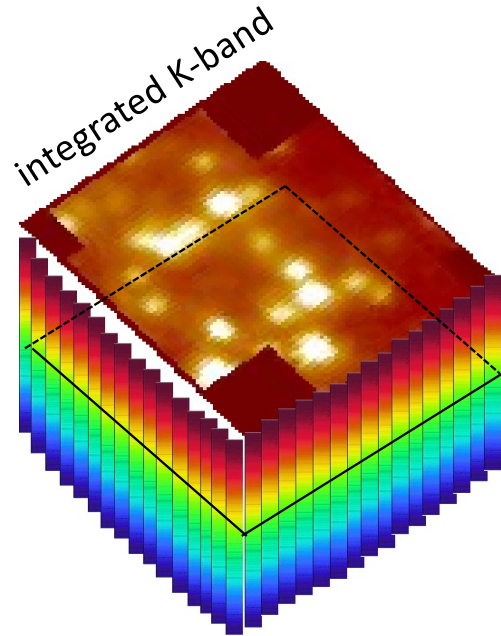
2.15 μm

3.8 μm

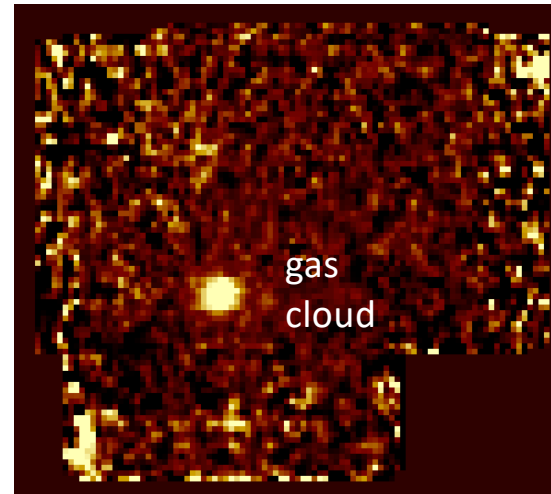


G2 is a dusty, ionized gas cloud

Integral-field spectroscopy: SINFONI

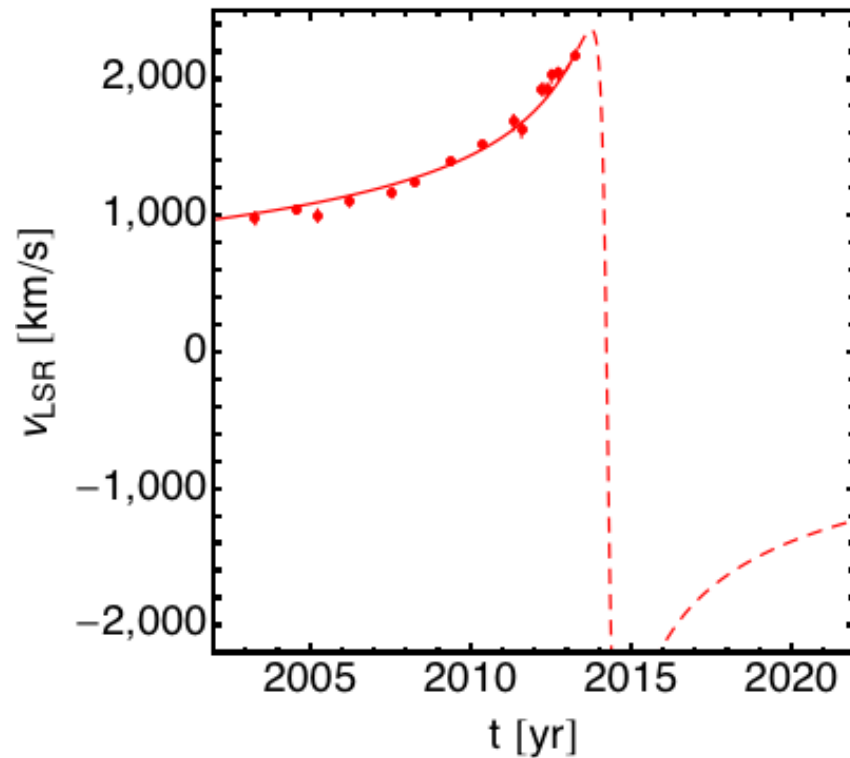
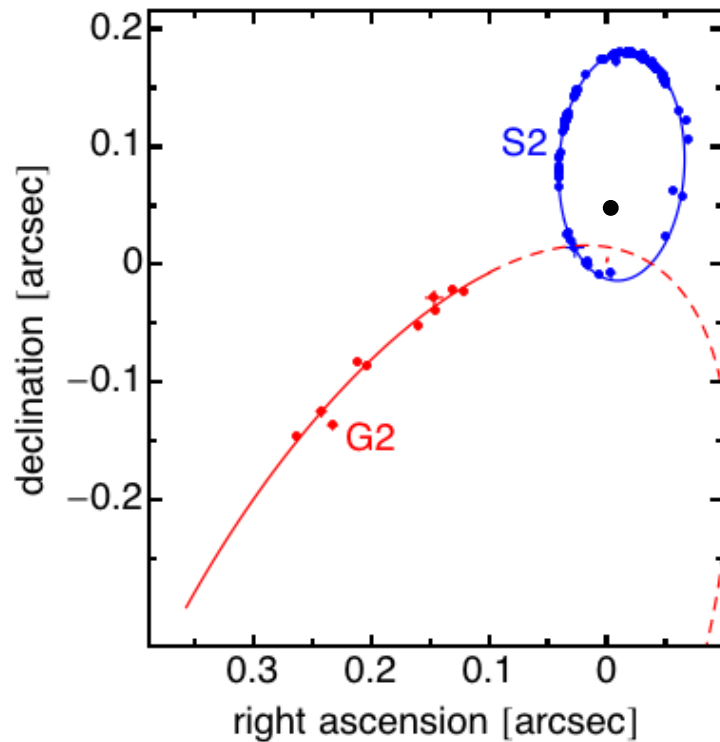


full data cube



channel map at
Br- γ + 1300 km/s
(2.175 μ m)

G2 has an interesting orbit



- pericenter 2014 @ $2000 R_S$
- highly eccentric

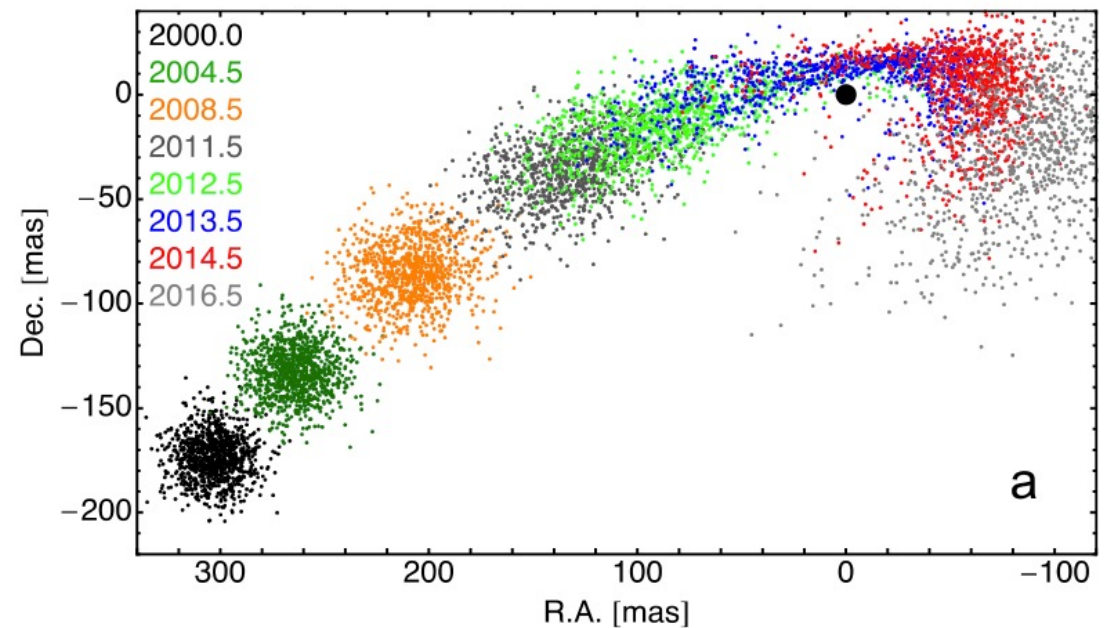
Phifer et al. 2013
Gillessen et al. 2013

For the gas dynamics the source model does not matter

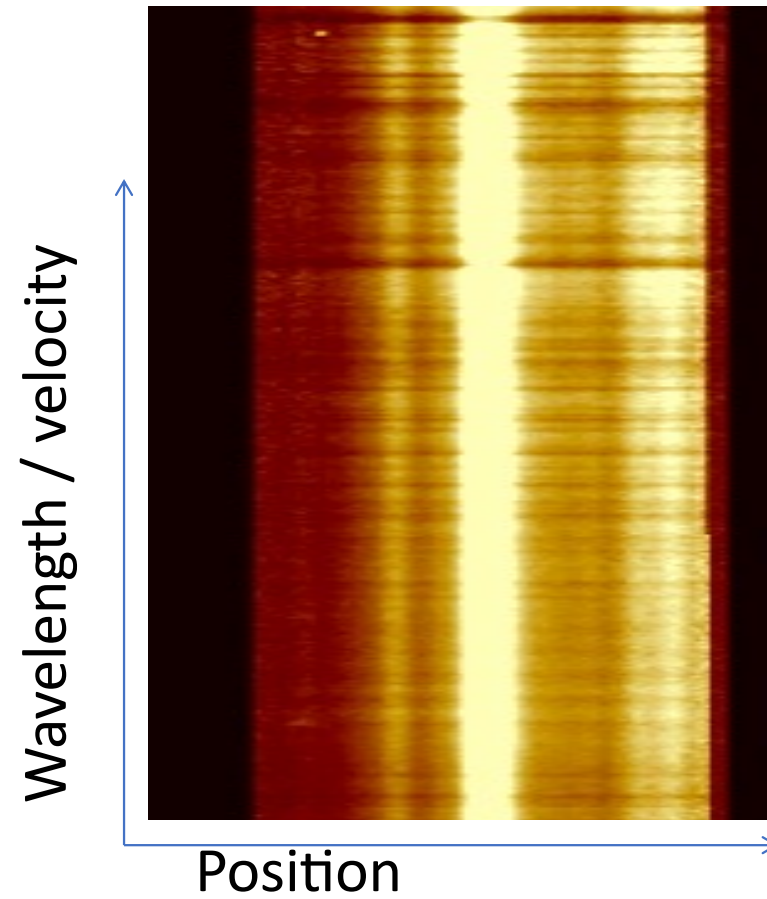
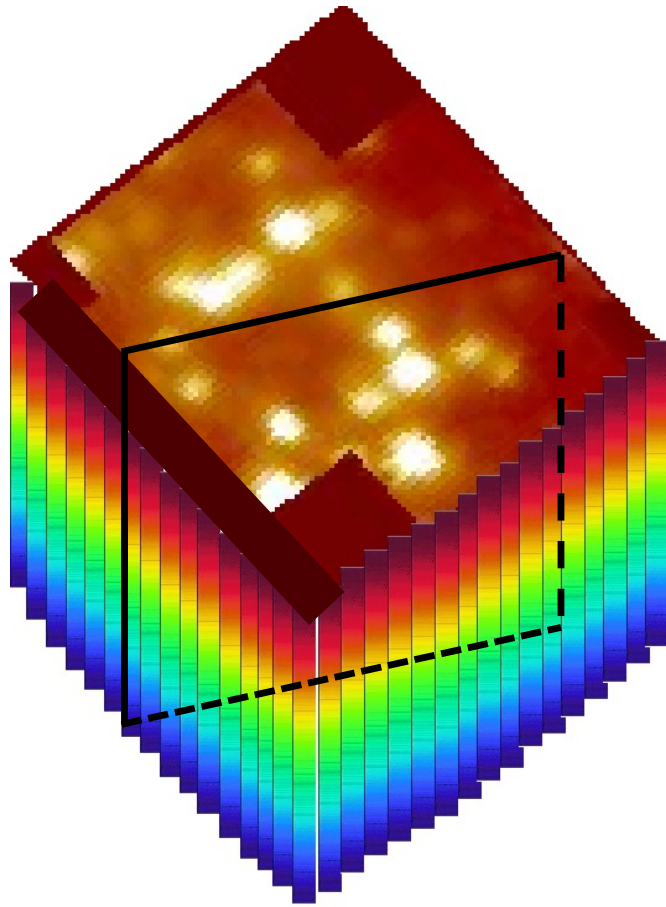
Gravitationally binding G2's
gas in the tidal field of Sgr A*:
 $M \gg 100M_{\odot}$

The gas that we see
cannot be bound to any
central object

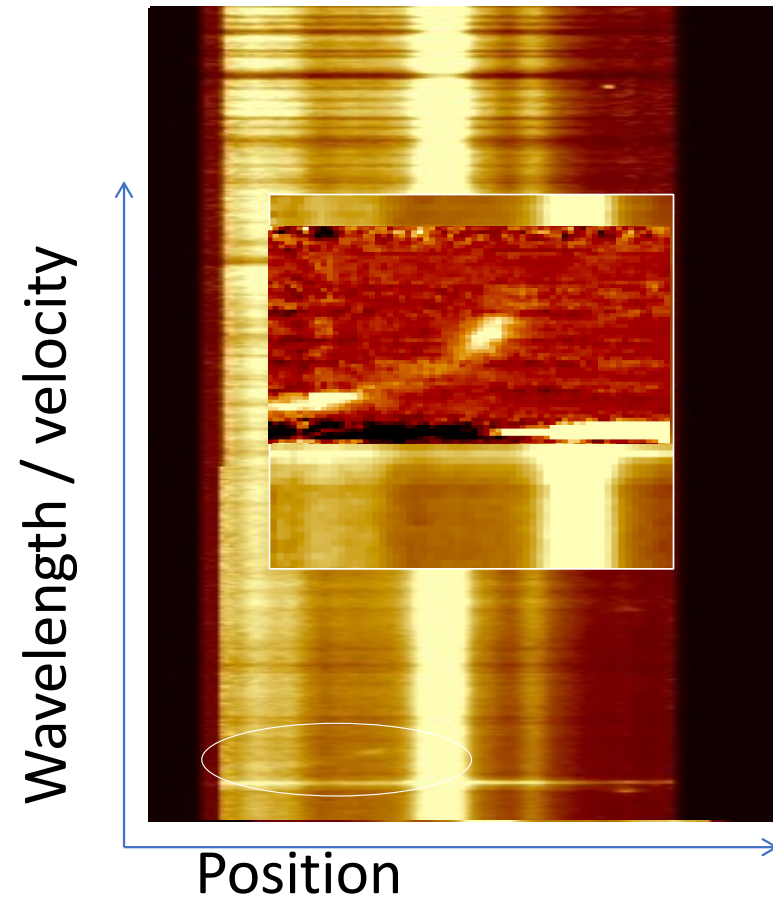
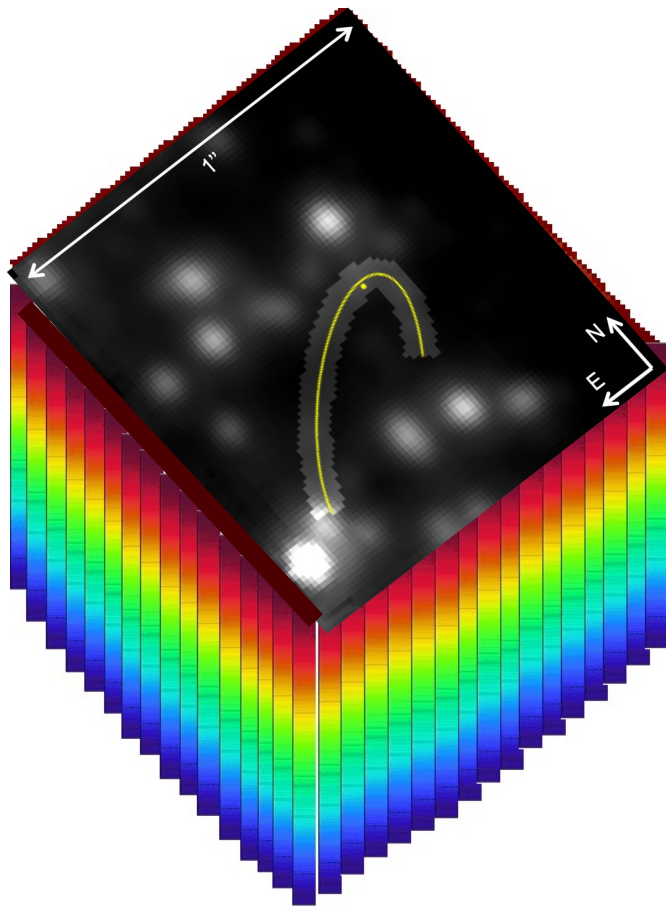
We should see G2
tidally evolve



Position – Velocity - Diagram

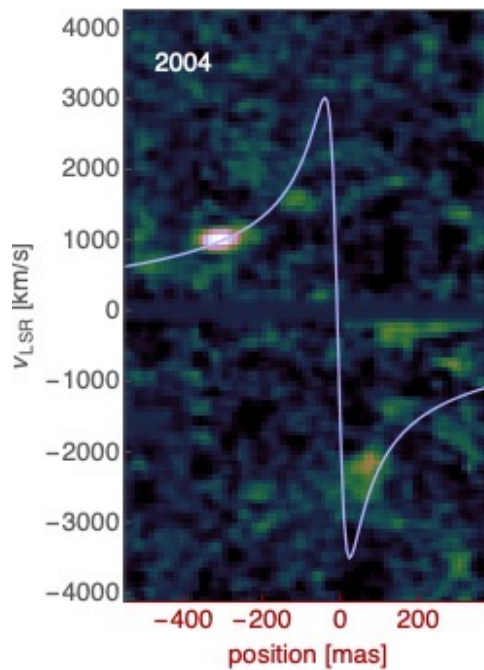


Position – Velocity - Diagram

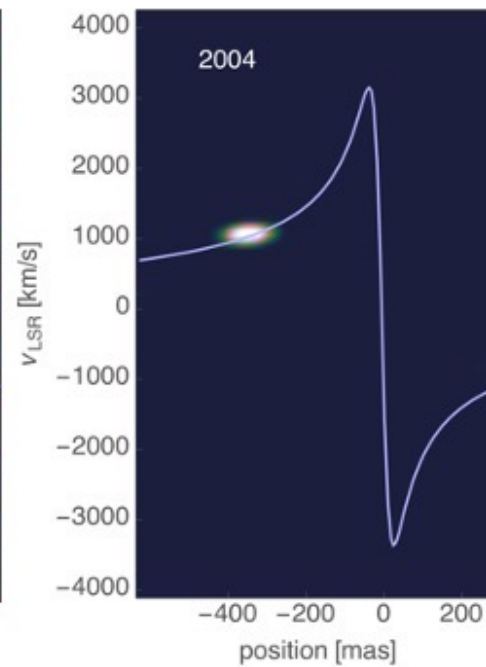


Too beautiful to watch: The tidal evolution of G2

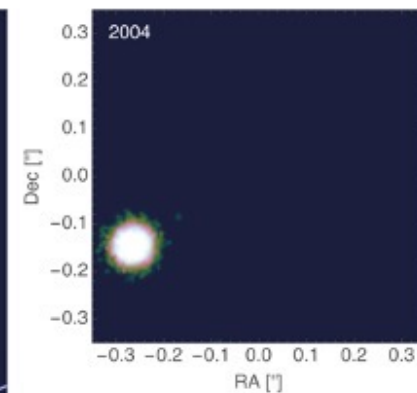
2004



the actual data
PV



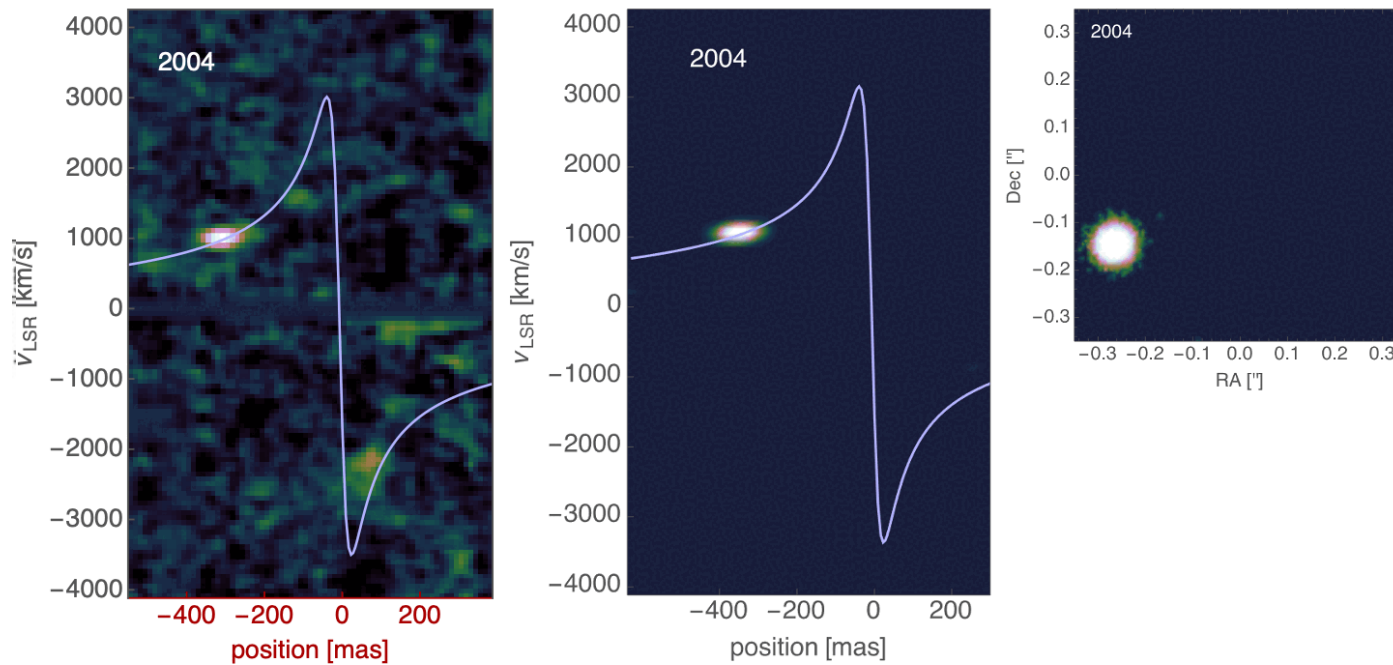
ensemble of
test particles - PV



ensemble of
test particles - RA/Dec

Too beautiful to watch: The tidal evolution of G2

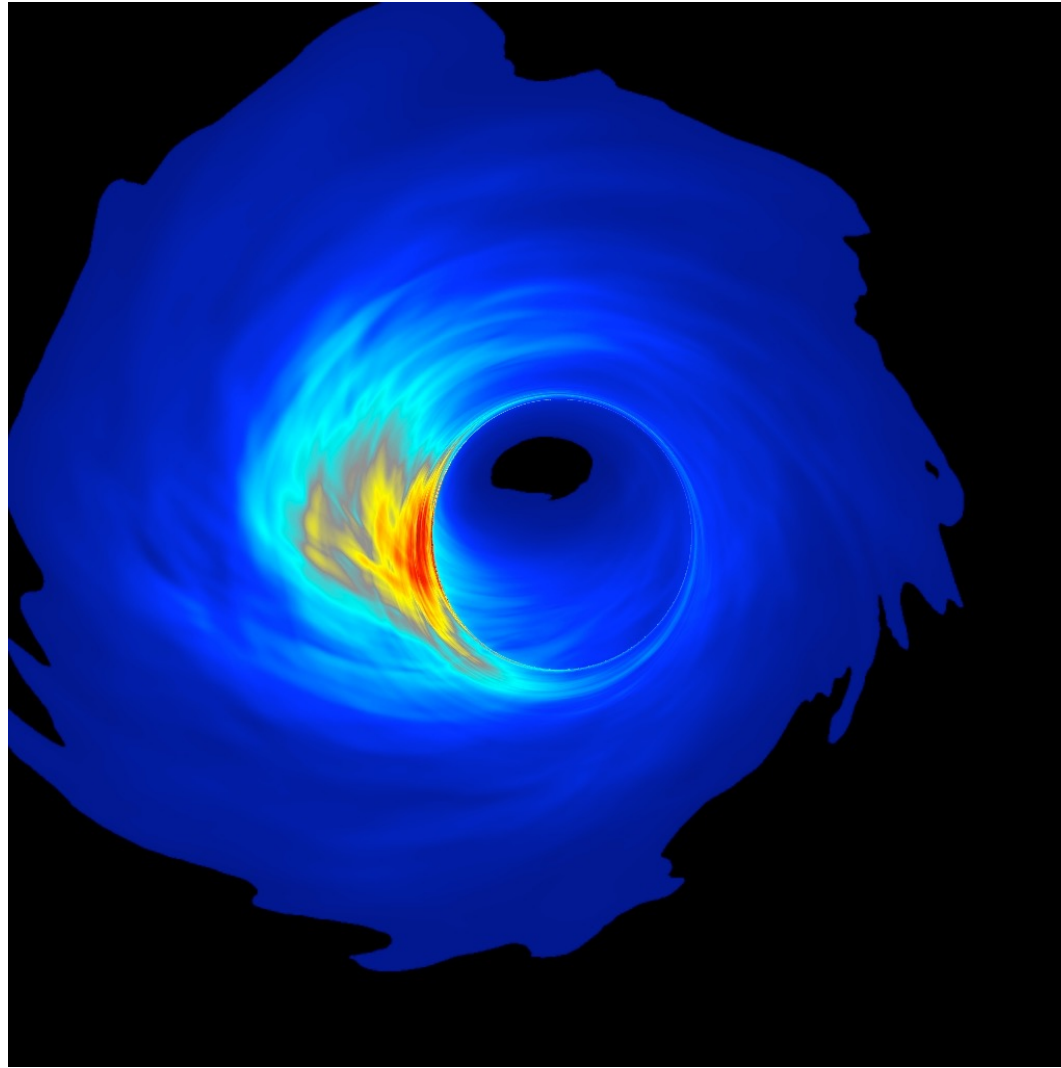
2004



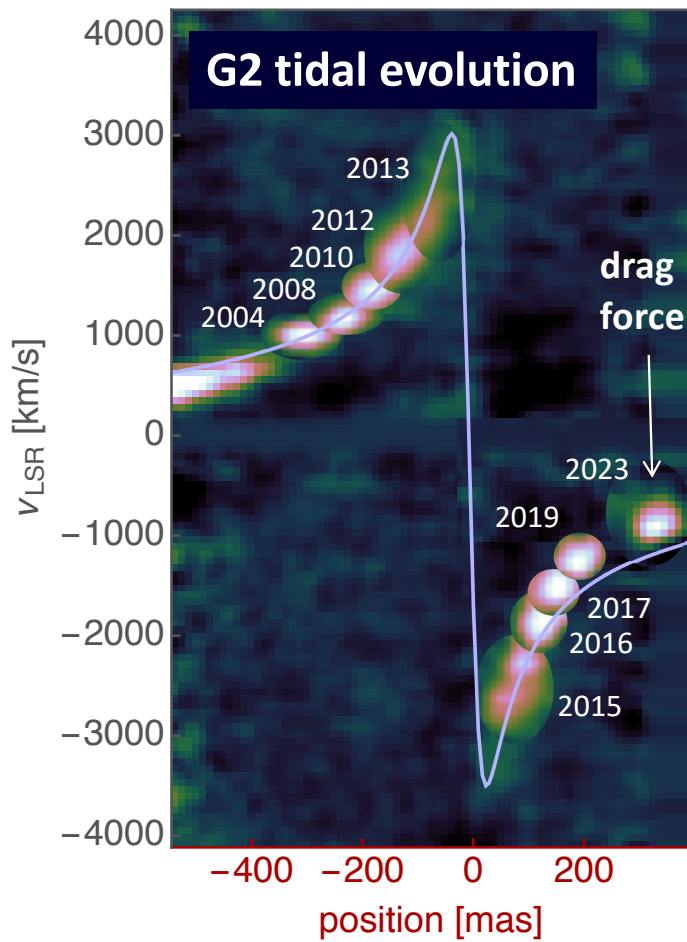
MailOnline



Missing: Interaction with Accretion Flow

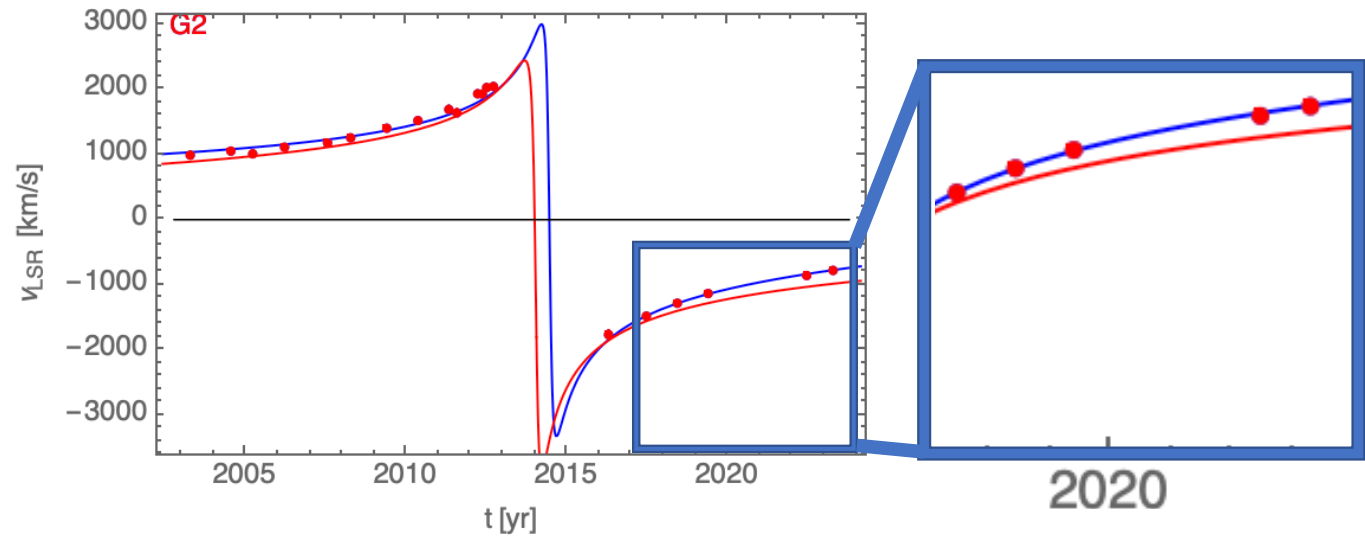
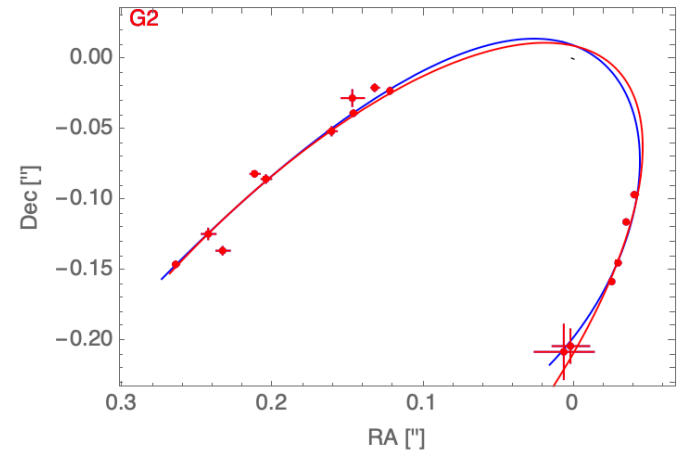


The G2 data since 2017 show a deceleration

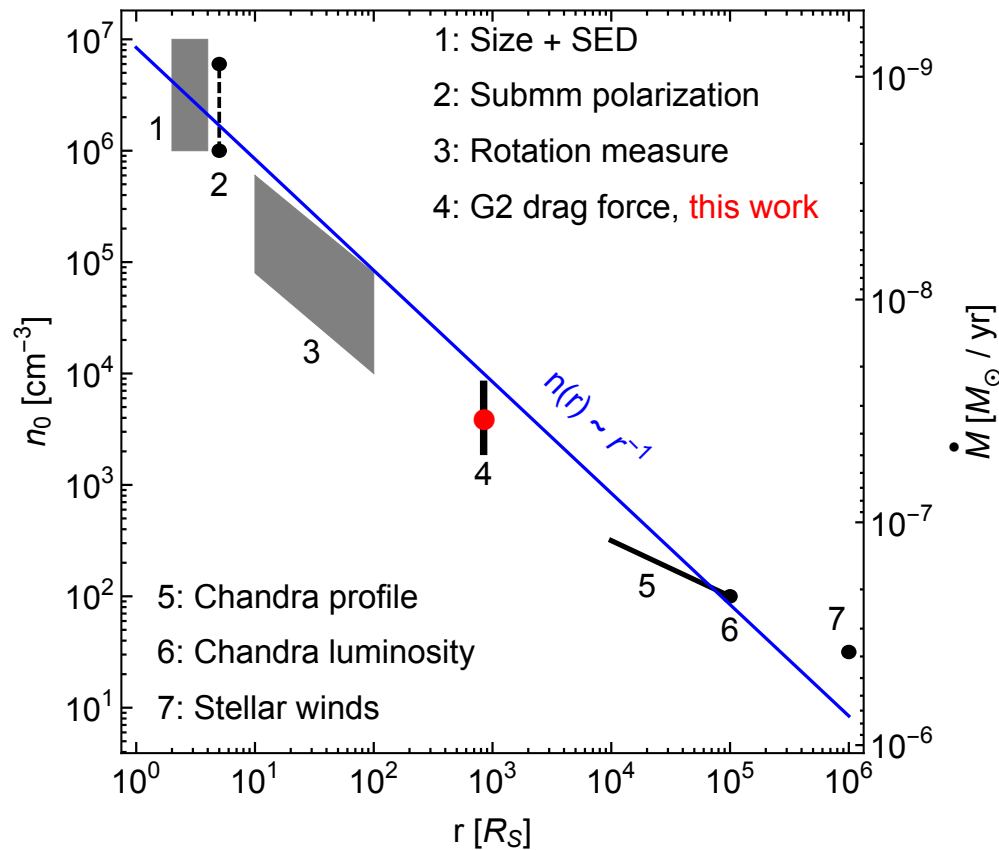


$$F_{\text{drag}} \propto s^2 \rho(r) v^2$$

$$\rho(r) \propto r^{-1}$$



A very valuable constraint of the accretion flow density



- 1: One zone models using size & SED
 (Doeleman et al. 2008,
 Bower et al. 2015,
 von Fellenberg et al. 2018)
- 2: Agol 2000,
 Quataert & Gruzinov 2000a
- 3: Bower et al. 2003,
 Marrone et al. 2006, 2007
- 5: Wang et al. (2013)
- 6: Baganoff et al. (2003)
- 7: Quataert et al. (2004)

Outline

- Part I The mass of Sgr A*
 - Astronomical black holes stellar and massive
 - Measuring mass orbits, Jeans-modeling
 - Infrared observations see through the dust
 - Adaptive Optics see sharp – diffraction limit
 - Stellar Orbits measure the 4.3 million M_{sun}
- Part II Errors, Fitting and all that
- Part III Testing the black hole paradigm
 - The black hole nature of Sgr A* mass, size, position, faintness
 - Interferometry breaks the diffraction limit
 - Testing G.R. in the Gal. Center redshift, precession, ... spin?
 - SgrA* flares accretion flow geometry
 - A funny gas cloud accretion flow density



Thanks to the
GRAVITY team