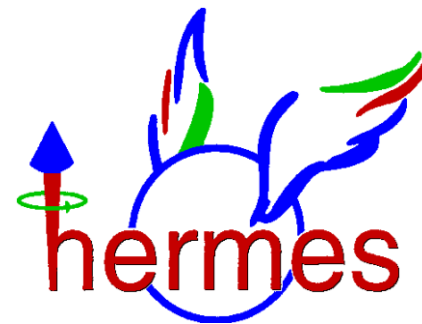


Beam-spin induced polarization of Λ and anti- Λ hyperons in semi-inclusive deep-inelastic scattering

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On behalf of the HERMES collaboration



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Longitudinal spin transfer $D_{LL'}$ in SiDIS

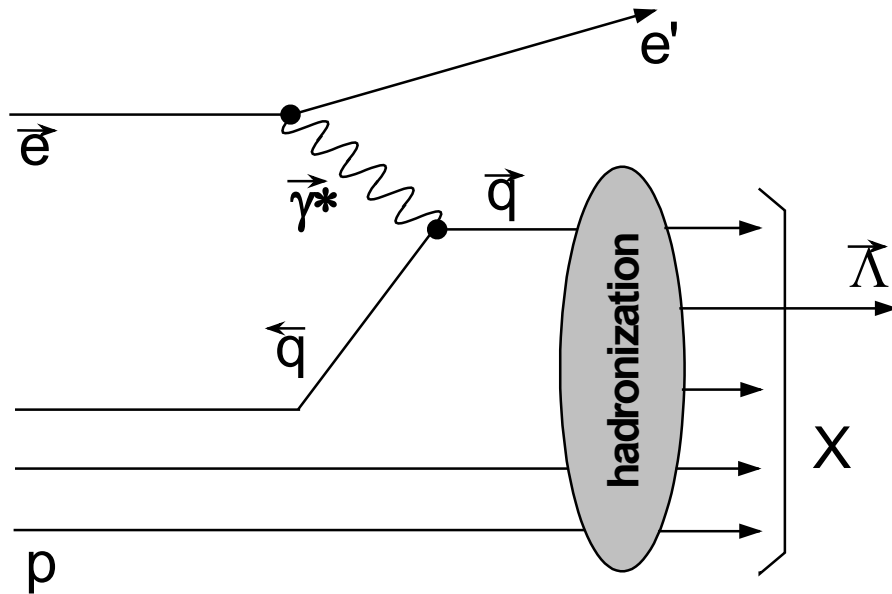
$$\vec{e} + N \rightarrow e' + \vec{\Lambda} + X$$

$$P_{L'}^{\Lambda} = P_{\gamma^*} D_{LL'} = P_b D(y) D_{LL'}$$

Spin transfer $D_{LL'}$ is correlation between beam polarization and final state hadron polarization or in other word double beam spin asymmetry

Why use Λ hyperon as final state hadron?

- Relatively good statistic
- Easy to detect $\Lambda \rightarrow p + \pi$
- Weak decay ("self" analyzing particle), polarization of Λ can be extracted by measuring decay proton angular distribution, no additional scattering needed



$$\frac{dN}{d\Omega_p} = \frac{dN_0}{d\Omega_p} (1 + \alpha P_{L'}^{\Lambda} \cos\theta_{pL'})$$

Angle between proton momentum and Λ spin in Λ rest frame



Spin transfer $D_{LL'}^\Lambda$

P.J. Mulders, R.D. Tangerman

Published in: Nucl.Phys.B 461 (1996), 197-237

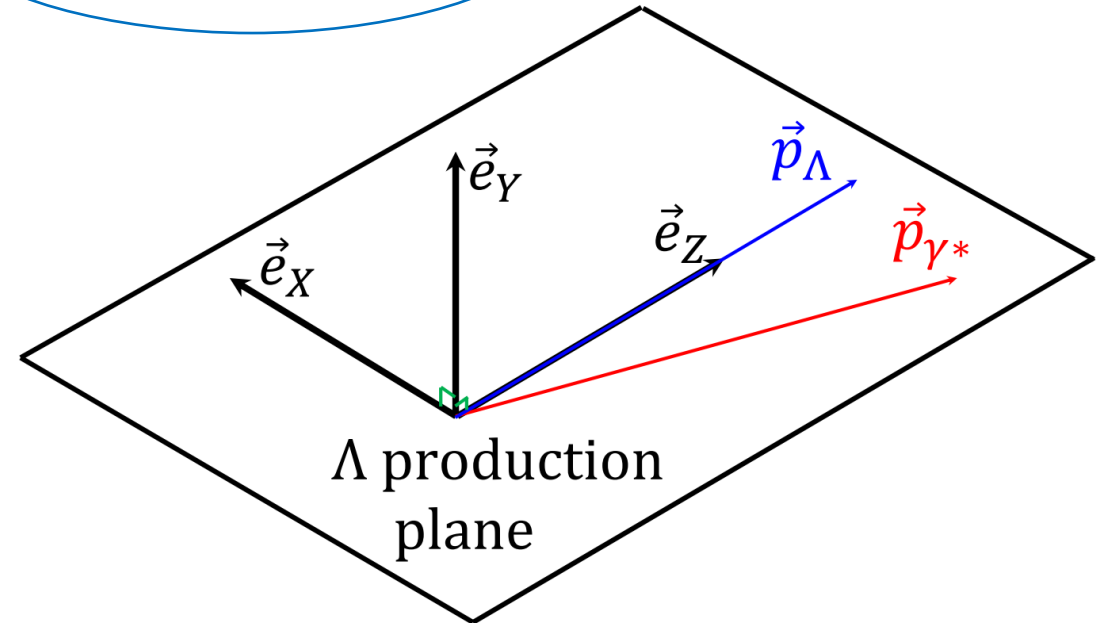
$$P_X^\Lambda = P_B D_X(y) \left\{ \frac{M \sum_q e_q^2 x_B e^q(x_B) H_1^q(z)}{Q \sum_q e_q^2 x_B f_1^q(x_B) D_1^q(z)} - \frac{M^\Lambda \sum_q e_q^2 x_B f_1^q(x_B) \tilde{G}_1^q(z)}{Q \sum_q e_q^2 x_B f_1^q(x_B) D_1^q(z)} \right\} = P_B D_X(y) D_{LX}(x_B, z, Q^2)$$

$$P_Y^\Lambda = D_Y(y) \frac{M \sum_q e_q^2 x_B f_1^q(x_B) D_{1T}^{\perp(1)q}(z)}{Q \sum_q e_q^2 x_B f_1^q(x_B) D_1^q(z)}$$

$$P_Z^\Lambda = P_B D_Z(y) \frac{\sum_q e_q^2 x_B f_1^q(x_B) G_1^q(z)}{\sum_q e_q^2 x_B f_1^q(x_B) D_1^q(z)} = P_B D_Z(y) D_{LZ}(x_B, z, Q^2)$$

$$P_X^\Lambda = P_B D_X(y) D_{LX}(x_B, z, Q^2)$$

$$P_Z^\Lambda = P_B D_Z(y) D_{LZ}(x_B, z, Q^2)$$

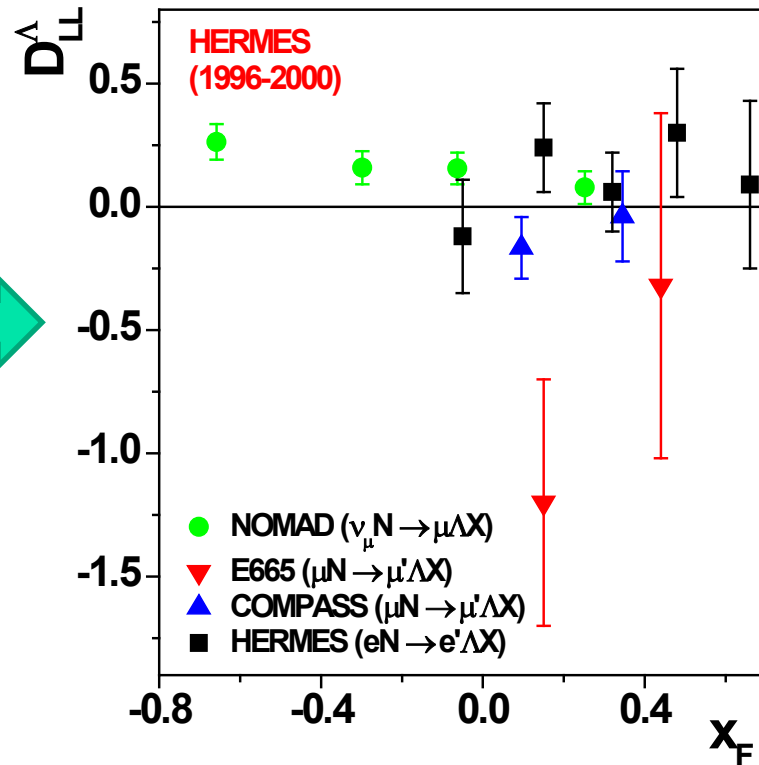
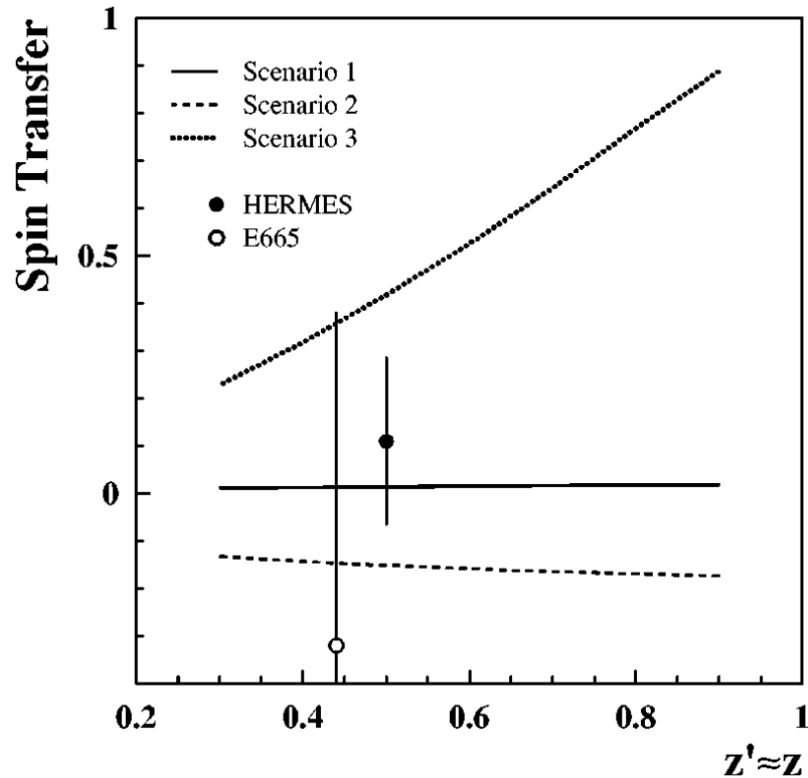


- $D_{LX}(x_B, z)$ and $D_{LZ}(x_B, z)$ called spin transfer coefficients
- P_Y^Λ is independent of beam polarization $P_B \rightarrow$ only 2 components $D_{LX}(x_B, z)$ and $D_{LZ}(x_B, z)$
- Give us access to chiral odd (H_1^q) and even (G_1^q, \tilde{G}_1^q) fragmentation functions



Previous publication

- Phys.Rev. D64 (2001) $\sim 2000 \Lambda s$. $D_{LL} = 0.11 \pm 0.17_{\text{stat}} \pm 0.03_{\text{syst}}$ (1996-1997 data)
- Phys.Rev. D74 (2006), $\sim 8200 \Lambda s$. $D_{LL} = 0.11 \pm 0.10_{\text{stat}} \pm 0.03_{\text{syst}}$ (1996-2000 data)



- For full data set we have $\sim 46000 \Lambda s$ and $\sim 6500 \bar{\Lambda} s$
- More than two times better statistical uncertainty on Λ
- First HERMES result on $\bar{\Lambda}$
- Also possible to extract kinematical dependences for $\bar{\Lambda}$

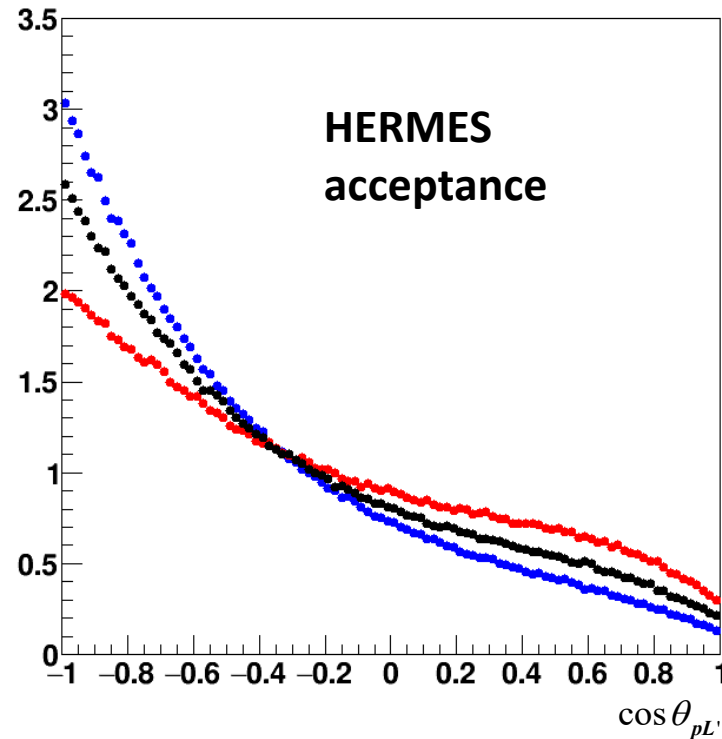
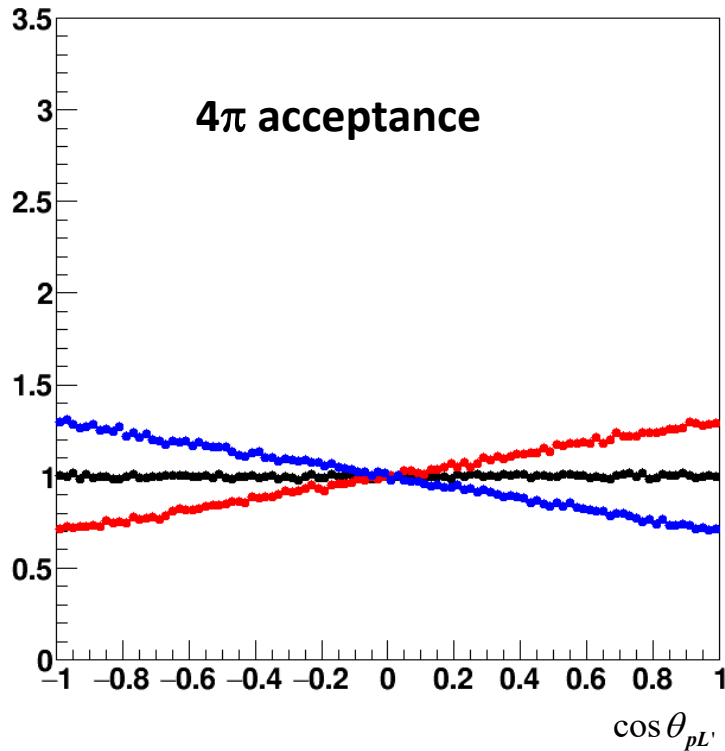


Extraction of Λ polarization

$$\frac{dN^{meas}}{d\Omega_p} = \frac{dN_0}{d\Omega_p} (1 + \alpha P_{L'}^\Lambda \cos\theta_{pL'}) \varepsilon(\theta_{pL'} \dots)$$

Unknown ←

Unpolarized $P^\Lambda = 0$
 Polarized $P^\Lambda = 0.5$
 Polarized $P^\Lambda = -0.5$



How to extract polarization?

1. 4π acceptance (**NOMAD**, " 4π " acceptance)

- P^Λ just a slope from linear fit

2. **Complex acceptance**

- MC simulation of $\varepsilon(\theta_{pL'} \dots)$ (major source of systematic uncertainty!) (**COMPASS**)
- Assume acceptance is stable over time, so we can cancel acceptance effect using positive and negative beam polarization (**HERMES**)



Extraction formalism of $D_{LL'}^\Lambda$

Moment method /Phys.Rev. D64 (2001)/

- calculate $\langle P_B \cos \theta_{pL'} \rangle$ and $\langle \cos^2 \theta_{pL'} \rangle$ with $\frac{d\omega}{d\Omega_p} = \frac{d\omega_0}{d\Omega_p} (1 + \alpha \vec{P}^\Lambda \vec{k}_p)$
- In simple 1D case and helicity balanced data sample $[[P_B]] \equiv \frac{1}{\mathcal{L}_{tot}} \int P_B d\mathcal{L} = 0$

$$\checkmark \langle P_B \cos \theta_{pL'} \rangle = \frac{[[P_B]] \langle \cos \theta_{pL'} \rangle_0 + \alpha D_{LL'} [[P_B^2]] \langle \cos^2 \theta_{pL'} \rangle_0}{1 + \alpha D_{LL'} [[P_B]] \langle \cos \theta_{pL'} \rangle_0} \quad \underline{\underline{[[P_B]] = 0}} \quad \alpha D_{LL'} [[P_B^2]] \langle \cos^2 \theta_{pL'} \rangle_0$$

$$\checkmark \langle \cos^2 \theta_{pL'} \rangle = \frac{\langle \cos^2 \theta_{pL'} \rangle_0 + \alpha D_{LL'} [[P_B]] \langle \cos^3 \theta_{pL'} \rangle_0}{1 + \alpha D_{LL'} [[P_B]] \langle \cos \theta_{pL'} \rangle_0} \quad \underline{\underline{[[P_B]] = 0}} \quad \langle \cos^2 \theta_{pL'} \rangle_0$$

Unpolarized moment
(unknown, but its reduced)

No MC simulation of acceptance needed

$$D_{LL'}^\Lambda = \frac{1}{\alpha [[P_B^2]]} \cdot \frac{\langle P_B \cos \theta_{pL'} \rangle}{\langle \cos^2 \theta_{pL'} \rangle}$$

- Slightly more complicated iteration procedure used in case of unbalanced P_B
- Extended for 3D case



Extraction of spin transfer in 3D and unbalanced case

$$\sum_{k=x,z} D_{Lk} \left\langle \frac{D_k(y)D_i(y)\cos\theta_k\cos\theta_i}{1 + \alpha \llbracket P_B \rrbracket \sum_{j=x,z} D_j(y)D_{Lj} \cos\theta_j} \right\rangle = \frac{1}{\alpha} \frac{\langle P_B D_i(y)\cos\theta_i \rangle - \llbracket P_B \rrbracket \langle D_i(y)\cos\theta_i \rangle}{\llbracket P_B^2 \rrbracket - \llbracket P_B \rrbracket^2}$$

- System of “linear” equations

$$\sum_k D_{Lk} a_{ik} = c_i$$

$$\begin{bmatrix} a_{xx} & a_{xz} \\ a_{zx} & a_{zz} \end{bmatrix} \begin{bmatrix} D_{Lx} \\ D_{Lz} \end{bmatrix} = \begin{bmatrix} c_x \\ c_z \end{bmatrix}$$

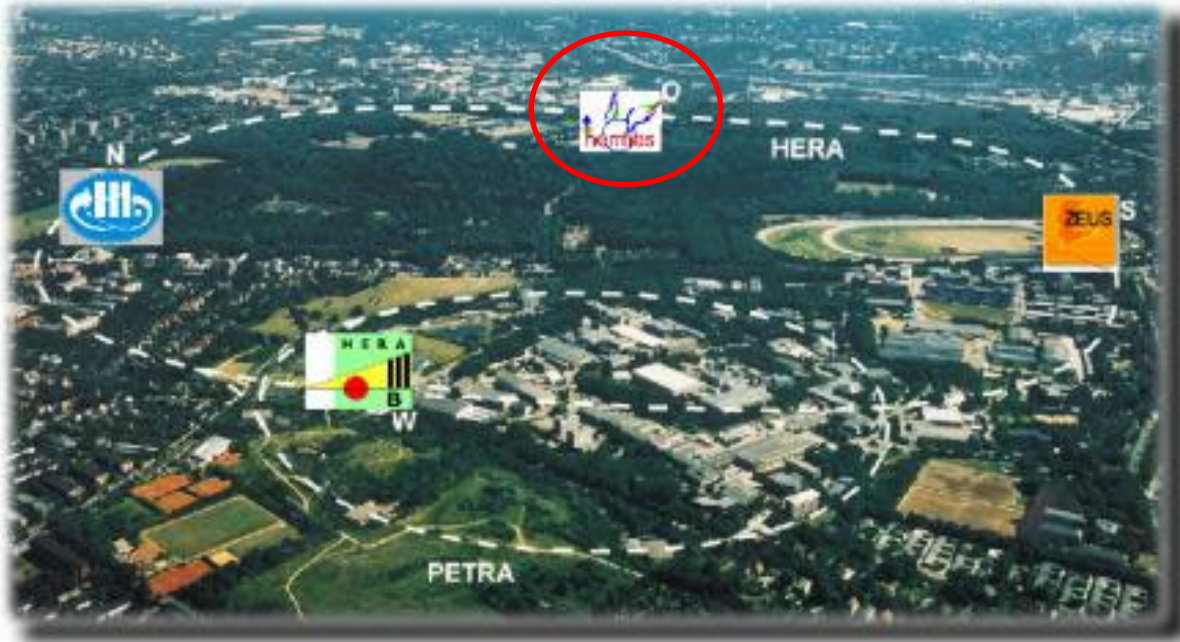
- Coefficients a_{ik} are dependent from D_{Li} (unless helicity balance $\llbracket P_B \rrbracket = 0$)
- Denominator in matrix elements are close to unity so we can use iterative procedure to solve it
- To solve given system, an iteration procedure is used

$$D_{Lk}^{(0)} \xrightarrow{\text{calculate } a_i^k} a_i^k{}^{(0)} \xrightarrow{\text{solve } a_i^k D_{Lk} = c_i} D_{Lk}^{(1)} \rightarrow \dots$$

- 3 steps are sufficient to converge due to $\alpha \llbracket P_B \rrbracket \sum_{j=x,z} D_j(y)D_{Lj} \cos\theta_j \sim 0$

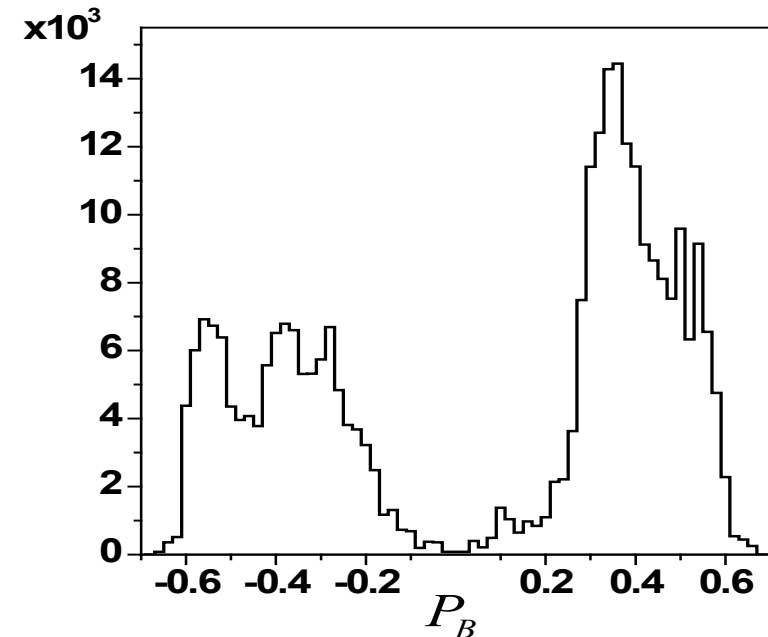


HERMES experiment



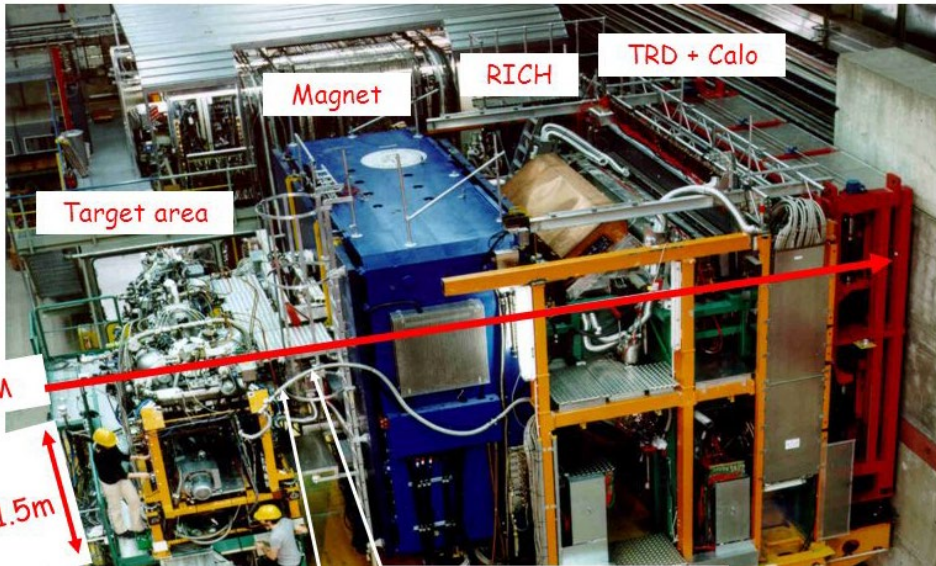
- Located at DESY, Hamburg
- 27.6 GeV e⁺/e⁻ beam and 920 GeV proton beam (HERMES used only lepton beam)
- Data taking end in 2007

- Lepton beam “self” transversely polarized up to 60% due to Sokolov-Ternov effect
- Spin rotators upstream and downstream of HERMES experiment used to get longitudinally polarized beam
- Beam spin flipped every few weeks/months

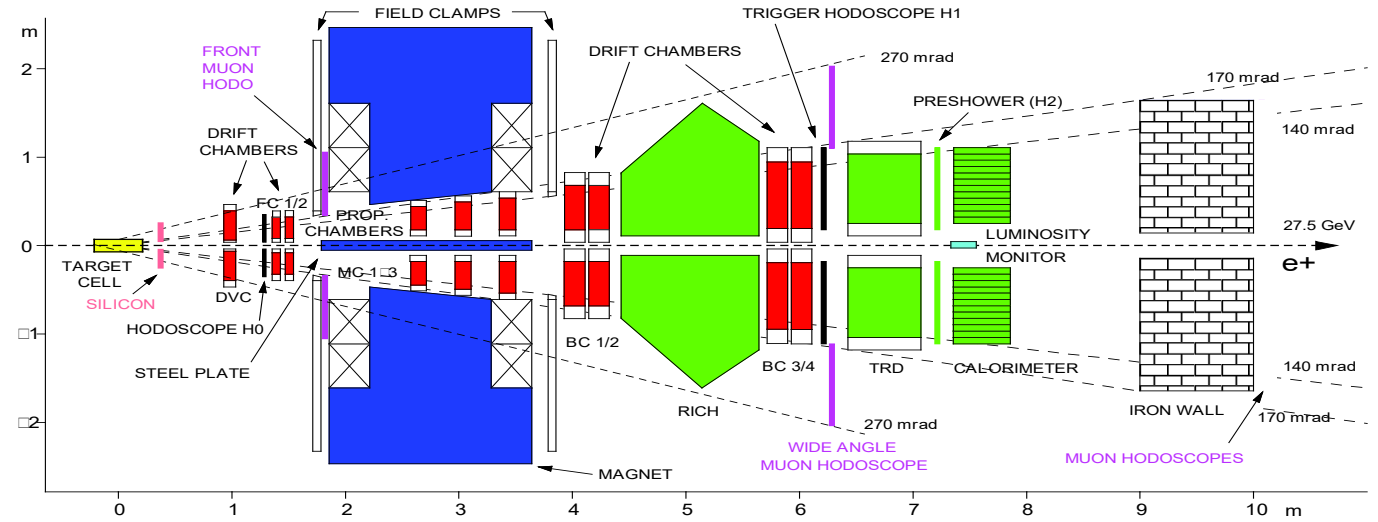




HERMES experiment



- Top/bottom symmetry
- $40 \text{ mrad} < \theta < 220 \text{ mrad}$
- Long. / trans. polarized gas targets H, D, flipped every 60-180 sec, $\langle P_{\text{targ}} \rangle \approx 0$
- Unpolarized targets H, D, He, N, Ne, Kr, Xe



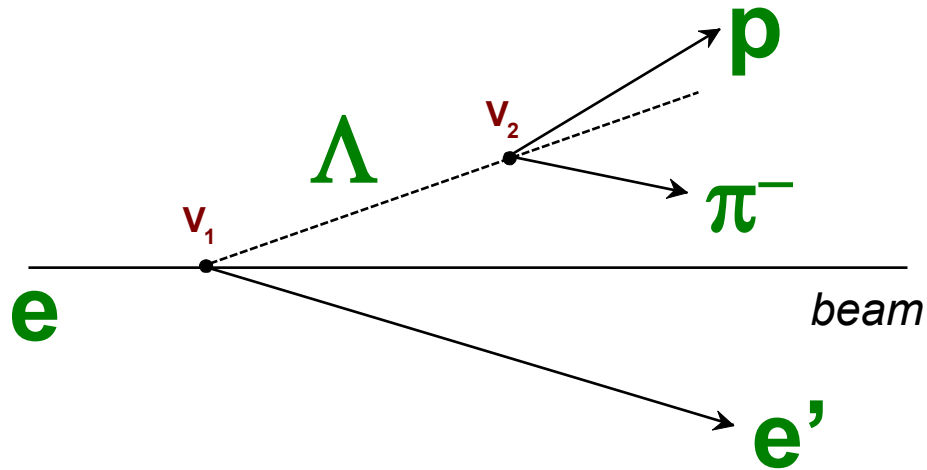
Particle ID:

- lepton/hadron separation
- Ring Cerenkov detector (RICH):
pion/kaon/proton discrimination $2 \text{ GeV} < p < 15 \text{ GeV}$

Good momentum and angular resolution

$$\frac{\Delta p}{p} \leq 1\%, \quad \Delta\theta_x, \Delta\theta_y \leq 1 \text{ mrad}$$

$$x = \frac{Q^2}{2M\nu}, y = \frac{\nu}{E_e} = \frac{E_e - E'_e}{E_e}, Z = \frac{E^\Lambda}{\nu}, x_F = \frac{p_{\parallel}^\Lambda}{p_{max}^\Lambda}$$



DIS cuts

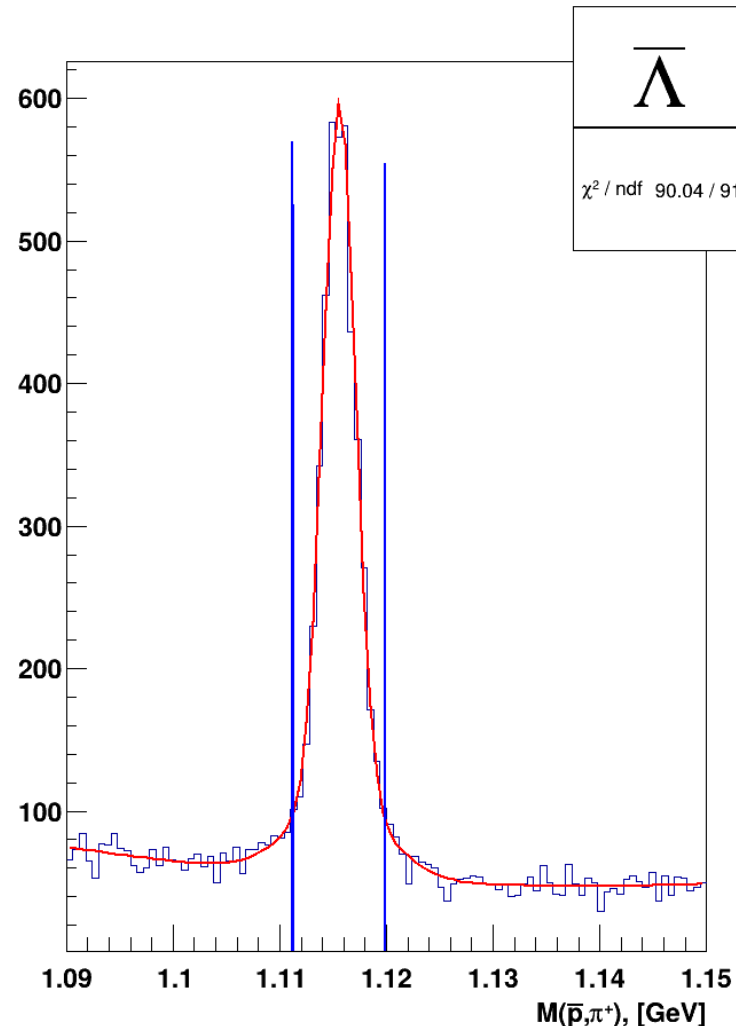
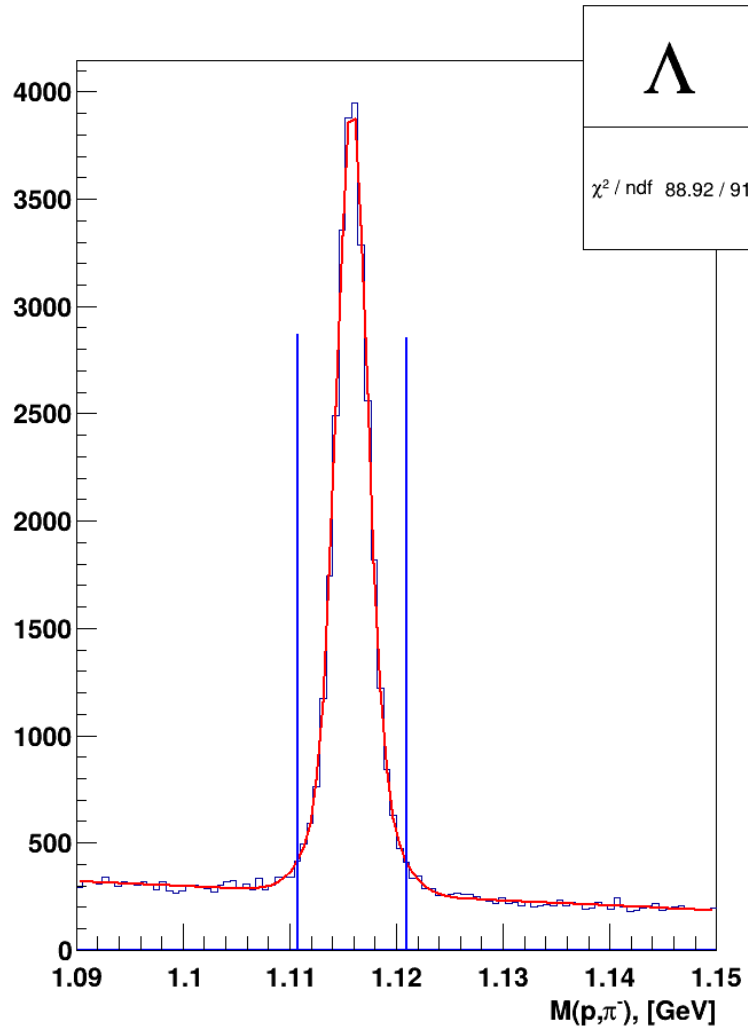
- $0.8 \text{ GeV}^2 < Q^2 < 24 \text{ GeV}^2$
- $W^2 > 10 \text{ GeV}^2$
- $0.2 < y < 0.85$

Background suppression

- leading π rejection (in HERMES kinematics proton is **always leading**):
- h^+h^- pair background (coming from V_1) rejection, vertex separation:
 - $d(V_1, V_2) > 6 \text{ cm}$ for Λ
 - $d(V_1, V_2) > 10 \text{ cm}$ for $\bar{\Lambda}$



$\Lambda (\bar{\Lambda})$ invariant mass



- Just an example for subset of data
- Clear $\Lambda(\bar{\Lambda})$ signal with highly suppressed background
- Background contamination
 - $> 5\%$ for Λ
 - $> 10\%$ for $\bar{\Lambda}$

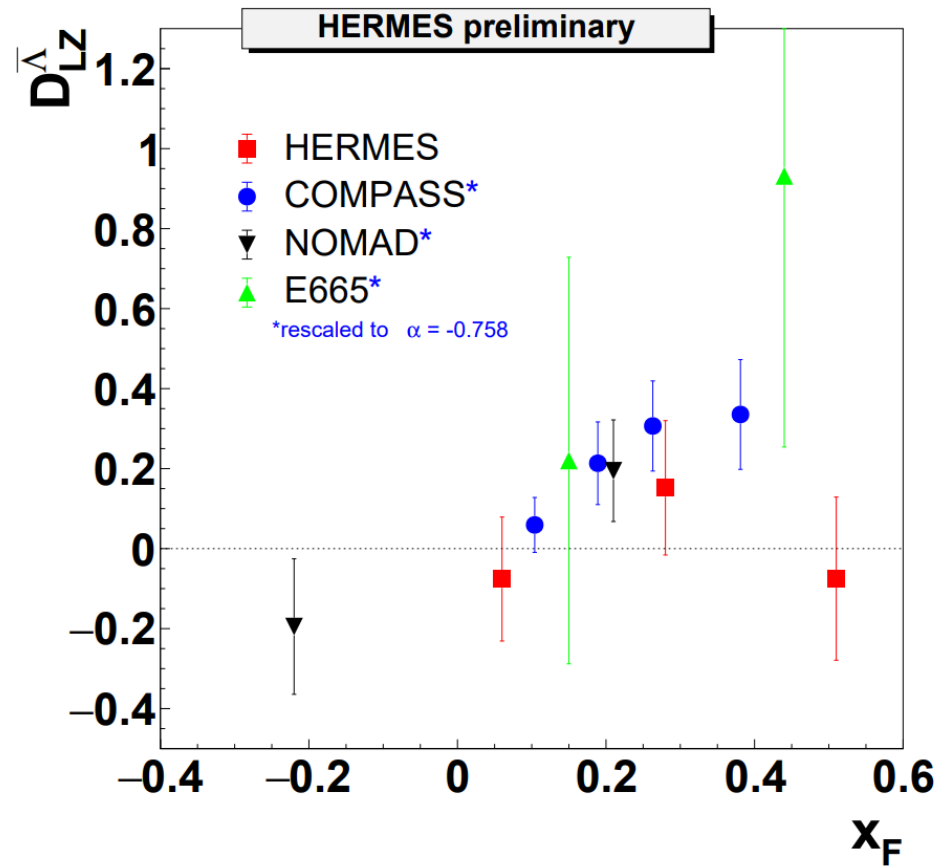
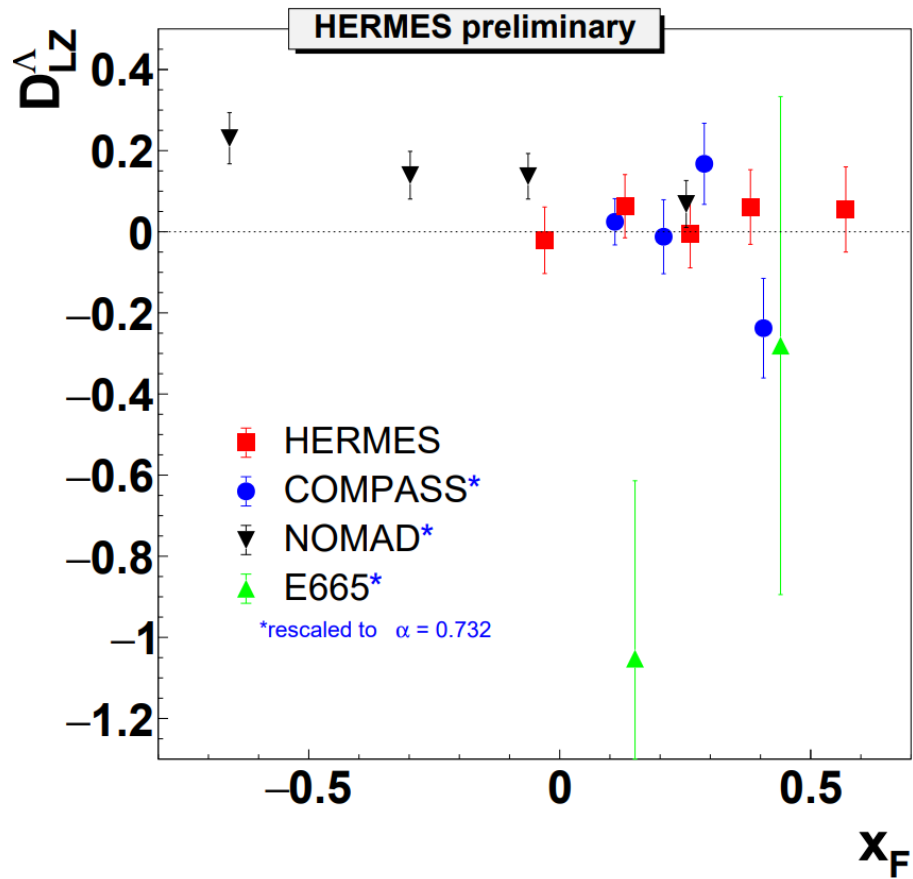
Total numbers of $\Lambda(\bar{\Lambda})$:

$$N_{\Lambda} = 46000$$

$$N_{\bar{\Lambda}} = 6500$$



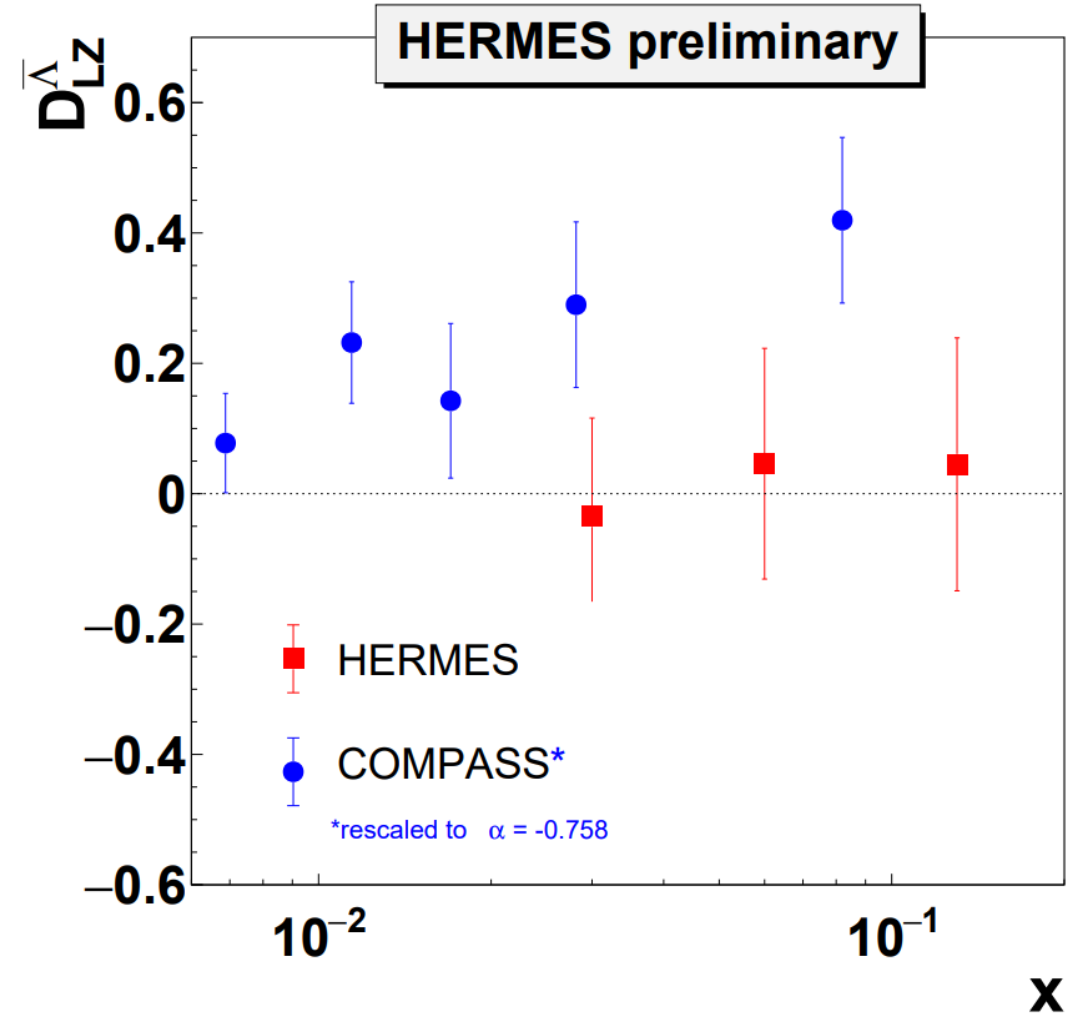
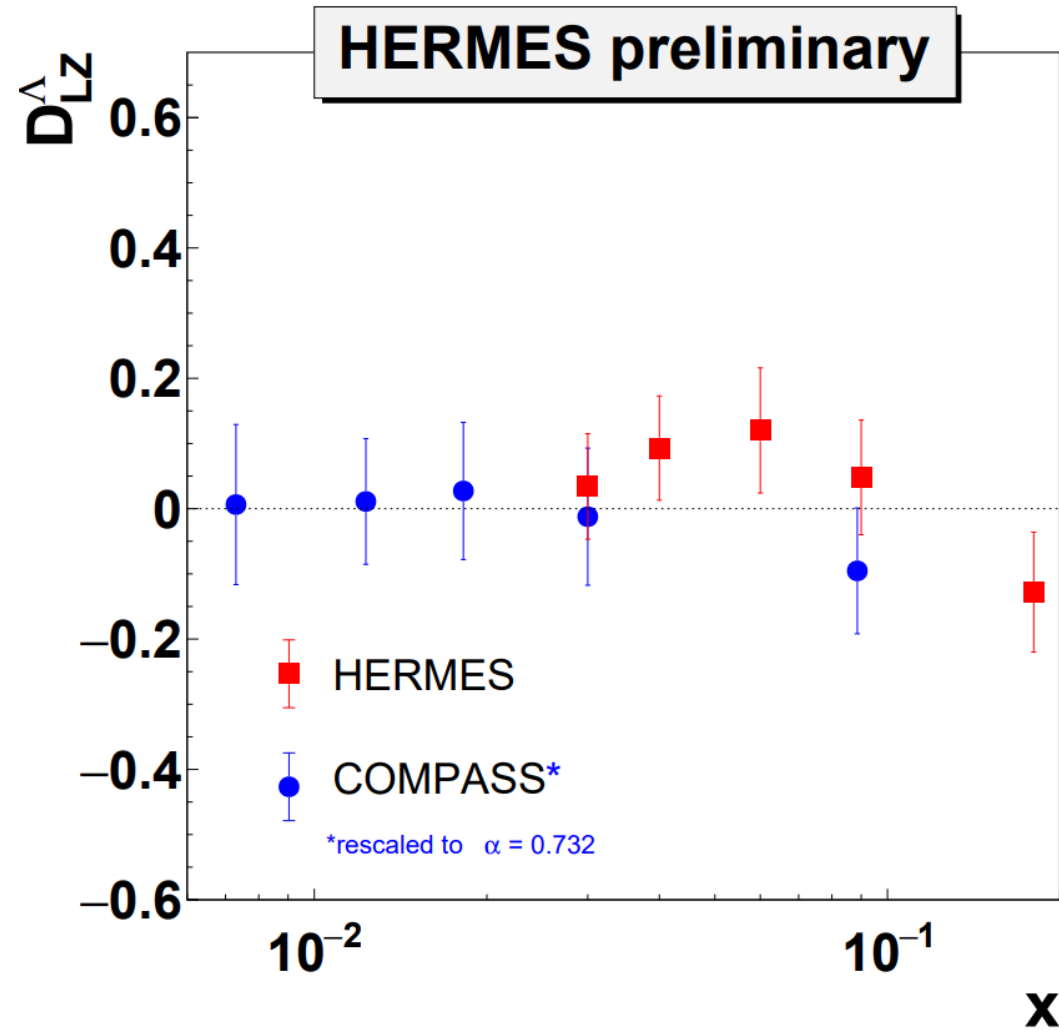
Results



Λ decay constant α has been updated in PDG → use new value and rescale data from other experiment



Results





Conclusion

- "self-analyzing" decay of Λ hyperon provides unique possibilities to study spin effects
- beam-spin transfer to Λ sensitive to various spin-dependent parton distributions and fragmentation functions (e.g. $e(x)$, \tilde{G}_1^q , G_1^q and H_1^q)
- HERMES has a large DIS data set with longitudinal beam polarization
- availability of both beam-helicity states exploited in novel extraction formalism that does not rely on MC simulations
- compared to previous HERMES publications, formalism has been extended to 3D case and to data sets that are not helicity-balanced
- increased data set allows for analysis of also $\bar{\Lambda}$
- analysis in advanced state, final results/publication to come out soon



Kinematical dependences

