

Recent Advances in pre Big Bang String Cosmology: Bouncing solutions in non-perturbative String Theory

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Why String Theory in Cosmology?

- Consistent theory of Quantum Gravity
- Avoid the Classical Big Bang Singularity
- Go beyond EFT approach to Inflation (Transplanckian problem)
- No need for ad-hoc fields to drive Inflation
- Possible observational signature of String Theory

What does String Theory suggest?

- The existence of a low energy massless multiplet $\{\phi, g_{\mu\nu}, B_{\mu\nu}\}$
- Pre Big-Bang scenario instead of Slow-Roll (Bouncing Cosmology)
- A systematic way to implement higher order curvature corrections (α^\prime expansion)
- A systematic way to implement higher order string coupling expansion in a perturbative way $g_s^2 = exp(\phi)$ (Genus Expansion)
- Additional simmetry in space-time with d abelian isometries, O(d,d) invariance in the field space (Continuous generalisation of T-duality)

Main Topics

• General criteria to have bouncing solutions using Hohm Zwiebach (HZ) action (all order α' action) starting from perturbative vacuum of String Theory

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- Dilaton stabilisation, FLRW and de-Sitter attractor from non-perturbative dilaton potential

Main Topics

- General criteria to have bouncing solutions using Hohm Zwiebach (HZ) action (all order α' action) starting from perturbative vacuum of String Theory
- Dilaton stabilisation, FLRW and de-Sitter attractor from non-perturbative dilaton potential
- Isotropisation mechanism via α^\prime corrections and Dilaton potential

$$S = \int Ldt = -\frac{1}{2} \int dt N e^{-\bar{\phi}} \left(N^{-2} \dot{\bar{\phi}}^2 + F(N^{-1} \dot{\beta}_i) + 2(\alpha')^{(d-1)/2} V(\phi) \right)$$

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$$\begin{split} ds^2 &= N^2(t)dt^2 - \sum_i \delta_{ij} e^{2\beta_i} dx^i dx^j \\ F &= -N^{-2} \sum_i \dot{\beta_i} + \mathcal{O}(\alpha') = -\sum_i H_i^2 + \mathcal{O}(\alpha') \quad \text{Even function of the Hubble functions} \\ \bar{\phi} &= \phi - \sum_i \beta_i \quad \text{Shifted Dilaton} \end{split}$$

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Low Energy isotropic solution (Asymptotic Past Triviality)

$$\phi \sim -\left(1+\sqrt{d}\right)\ln\left(-t\right)$$

$$H = \dot{\beta} = \frac{1}{\sqrt{d}(-t)}$$

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$$\bar{\phi} = \phi - \sum_{i} \beta_{i} \quad \text{Shifted Dilaton} \quad \text{Expanding} \quad \text{Expanding}$$

$$\frac{\text{Expanding}}{\text{pre-Big Bang}} \quad \text{Expanding}$$

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The Routhian formalism and the all order $\alpha^\prime~$ Hamiltonian.

Legendre transform only a subset of the variables of the Lagrangian β_i

$$\mathcal{R}(N,\bar{\phi},\pi_i) = \sum_i \pi_i \dot{\beta}_i - L = N e^{\bar{\phi}} \left[\frac{1}{2} N^{-2} \dot{\phi}^2 + h(z_i) + V(\bar{\phi} + \sum_i \beta_i) \right]$$
$$h(z_i) \equiv \frac{1}{2} \left(F - \sum_i \dot{\beta}_i \frac{\partial F}{\partial H_i} \right) = \frac{1}{2} \sum_i z_i^2 + \mathcal{O}(\alpha') + \dots \qquad \frac{\partial h}{\partial z_i} = H_i$$

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$$\begin{aligned} \mathcal{R}(N,\bar{\phi},\pi_i) &= \sum_i \pi_i \dot{\beta}_i - L = N e^{\bar{\phi}} \begin{bmatrix} \frac{1}{2}N^{-2} \dot{\bar{\phi}}^2 + h(z_i) + V(\bar{\phi} + \sum_i \beta_i) \end{bmatrix} \\ h(z_i) &\equiv \frac{1}{2} \left(F - \sum_i \dot{\beta}_i \frac{\partial F}{\partial H_i} \right) = \frac{1}{2} \sum z_i^2 + \mathcal{O}(\alpha') + \dots \qquad \frac{\partial h}{\partial z_i} = H_i \\ & \text{EOM} \\ \\ \dot{\bar{\phi}}^2 &= 2h(z_i) + 2V, \qquad \dot{z}_i = z_i \dot{\bar{\phi}} - \frac{\partial V}{\partial \phi}, \qquad \ddot{\bar{\phi}} = \sum_i z_i \frac{\partial h}{\partial z_i} + \frac{\partial V}{\partial \phi} \end{aligned}$$

The Routhian formalism and the all order $\alpha^\prime~$ Hamiltonian.

Legendre transform only a subset of the variables of the Lagrangian β_i $\mathcal{R}(N,\bar{\phi},\pi_i) = \sum_i \pi_i \dot{\beta}_i - L = N e^{\bar{\phi}} \left| \frac{1}{2} N^{-2} \dot{\bar{\phi}}^2 + h(z_i) + V(\bar{\phi} + \sum_i \beta_i) \right|$ $h(z_i) \equiv \frac{1}{2} \left(F - \sum_i \dot{\beta}_i \frac{\partial F}{\partial H_i} \right) = \frac{1}{2} \sum z_i^2 + \mathcal{O}(\alpha') + \dots \qquad \frac{\partial h}{\partial z_i} = H_i$ EoM $\dot{\bar{\phi}}^2 = 2h(z_i) + 2V, \qquad \dot{z}_i = z_i \dot{\bar{\phi}} - \frac{\partial V}{\partial \phi}, \qquad \ddot{\bar{\phi}} = \sum_i z_i \frac{\partial h}{\partial z_i} + \frac{\partial V}{\partial \phi},$

Advantages:

- Specify only two functions to have a non-perturbative description of the system.
- Simple 'hamiltonians' capture all α' corrections and generate bouncing solutions they have another zero except the trivial one. In general they come from non-holomorphic F. In the isotropic case we used: $d = \frac{2^4}{2^4}$

$$h(z) = \frac{d}{2}(z^2 - \alpha' \frac{z^4}{2})$$

Dilaton Potential with Istantonic behaviour not captured by string coupling expansion (Non-perturbative Potential)



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$$V(\phi) = A e^{-B(\phi)/\beta} \left[\left(c^2 - B(\phi) \right)^2 + \delta B(\phi) \right] \left[1 - q B^{-1}(\phi) \right]$$
$$B(\phi) = \frac{1 + \alpha g_s^2}{\alpha g_s^2} = \frac{1 + \alpha e^{\phi}}{\alpha e^{\phi}}$$



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Is it possible to reach a final state with a frozen dilaton?

- Stabilisation of the dilaton implies stabilisation of the Newton's constant
- What is the final geometry?
- Is the final configuration an attractor? ¹⁷

Dilaton Stabilisation and FLRW attractor

Asymptotic Matter FLRW geometry $\delta=0$

 $V_0 = 0$

Dilaton Stabilisation and FLRW attractor



φ

Dilaton Stabilisation and FLRW attractor



φ

100

Asymptotic de-Sitter geometry $\delta \neq 0$

$$V_0 = V(\phi_m) > 0$$

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$$\phi_0 \simeq \phi_m - \frac{2V_0}{m^2(d-1)}$$

$$H_0 \simeq \left[\frac{2V_0}{d(d-1)}\right]^{1/2}$$



·-3.60

-1.0

-0.5

√dH

-2

0.5

Asymptotic de-Sitter geometry $\delta \neq 0$

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Minimum of the dilaton potential in the E-frame (The dilaton is a minimally coupled scalar in the E-frame)



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Minimum of the dilaton potential in the E-frame (The dilaton is a minimally coupled scalar in the E-frame)



Is this solution a stable attractor?

Numerical analysis of the attraction basin: Isotropic case

Variation of the height of the first peak (β) and the amplitude A

β controls the height of the first peak.

• Conserved quantity from the EoM used to fix the initial condition of the dilaton.

 $e^{-\overline{\phi}}z = \kappa^{-1}$



Numerical analysis of the attraction basin: Isotropic case

 $\alpha = 10$

 $z_{in} = 0.01$

d = 3

β controls the height of the first peak.

 Conserved quantity from the EoM used to fix the initial condition of the dilaton.

 $e^{-\overline{\phi}}z = \kappa^{-1}$

- **q** controls the asymptotic behaviour of the dilaton potential.
- **q**=1: Asymptotically vanishing potential

de Sitter, stabilised dilaton

: Growing dilaton, Runaway

: Time-reversal post Bounce,

decreasing dilaton

 0 ≤ q < 1: Asymptotically constant potential



Variation of the height of the first peak (β) and the amplitude A

Whenever there is a late-time attractor with constant ϕ and z_i the attractor must be isotropic, i.e $z_i = z = z_0$, and consequently $H_i = H = H_0$.

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$$\left(\sum_{j}H_{j}\right)z_{i}+\frac{\partial V}{\partial\phi}=0$$

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Case study: Bianchi I geometry with two isotropic *d* and *n* dimensional subspaces

$$h(z_i) = rac{1}{2}\sum_i z_i^2 - rac{lpha'}{4}\sum_i z_i^4$$

Anisotropic "Hamiltonian" consistent at the first order in α' with the perturbative result obtained for F in heterotic string theory

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From the 2° EoM $\left(\sum_{i} H_{j}\right) z_{i} + \frac{\partial V}{\partial \phi} = 0$ Independent from i! Consistent with the other two EoM Case study: Bianchi I geometry with two isotropic d and n dimensional subspaces H_1 0.06 $h(z_i) = rac{1}{2}\sum_i z_i^2 - rac{lpha'}{4}\sum_i z_i^4$ -1.00.04 -1.5 0.02 Anisotropic "Hamiltonian" 0.00 consistent at the first order in α' -2.0 -0.02 with the perturbative result $d = 1 \quad n = 2$ -2.5obtained for F in heterotic string theory

• Constant quantity

$$\kappa = e^{\phi} \left[\frac{d H_1 + n H_2}{\sqrt{d+n}} \right]^{-(1+d \gamma_1 + n \gamma_2)}$$

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Past-asymptotic low curvature solutions.

• Scale factors:

$$a_i \sim (-t)^{-\gamma_i}$$

• Hubble functions:

$$H_i = \frac{\gamma_i}{t}$$

• Kasner condition

 $\sum \gamma_i^2 = 1$

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$$\gamma_2 = \sqrt{(1 + d\epsilon)/(d + n)}$$

 ϵ : Anisotropy parameter

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de Sitter attraction basin in the $\{\epsilon,\kappa\}$ space



Future research and developments

- Stabilizing mechanism for the internal compact dimensions
- Add the Kalb-Ramond $B_{\mu\nu}$ field
- Compute perturbations (scalar and tensor) to bound the model with observational data
- Can the Lagrangians, Hamiltonians, or Routhians that implement a regular bounce correspond to the dimensional reduction of some general covariant action.

Thanks for your attention!

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This presentation was based on:

P. Conzinu, G.Fanizza, M.Gasperini, E. Pavone, L. Tedesco, G. Veneziano, *From the string vacuum to FLRW or de Sitter via* α *corrections*, JCAP **12**, 019 (2023).