

# *$\eta$ – $\eta'$ Mixing-From Electromagnetic Transitions to Weak Decays of Charm and Beauty Hadrons*

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Work in collab. with C. Di Donato and I. Bigi

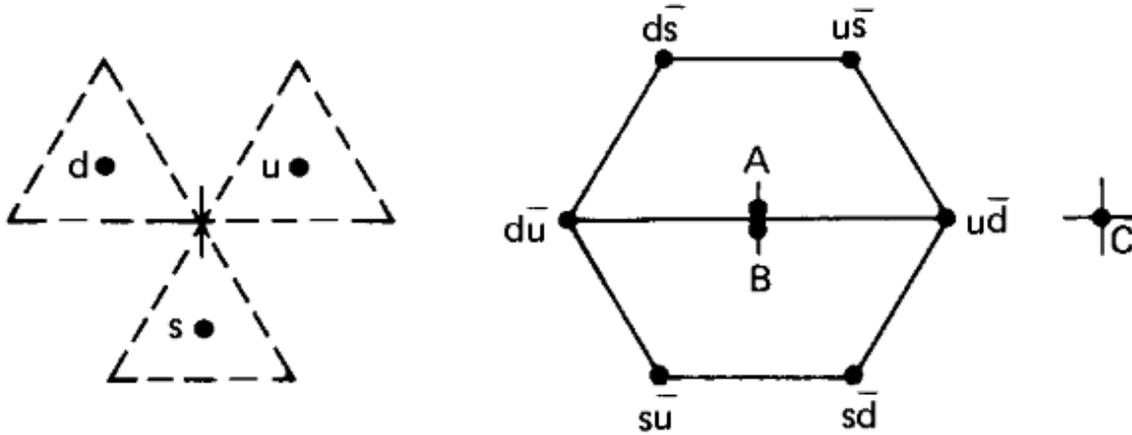
e-Print: [arXiv:1105.3557](https://arxiv.org/abs/1105.3557) [hep-ph]

LNF, Frascati, June 9<sup>th</sup>, 2011

## Outline of the talk

1. Review on theoretical and experimental progresses in  $\eta$ – $\eta'$  mixing—probing their gluonic content
2. Analysis one by one of relevant processes at different energy scales
  - ❖ electromagnetic and strong decays
  - ❖ electroweak D and B decays
3. Conclusions and future prospects

# $SU(3)_{\text{flavor}}$



$$B = \eta_8 = \frac{1}{\sqrt{3}} (u\bar{u} + d\bar{d} + s\bar{s})$$

$$C = \eta_0 = \frac{1}{\sqrt{6}} (u\bar{u} + d\bar{d} - 2s\bar{s})$$

$$3 \otimes \bar{3} = 8 \oplus 1$$

$SU(3)$  decomposition in the  $I_3, Y$  plane

$q\bar{q}$ Orbital Ang. Mom.	$q\bar{q}$ Spin	$J^{PC}$	Observed Nonet			Typical Mass (MeV)
			$I = 1$	$I = \frac{1}{2}$	$I = 0$	
$L = 0$	$S = 0$	$0^{-+}$	$\pi$	K	$\eta, \eta'$	500
	$S = 1$	$1^{--}$	$\rho$	$K^*$	$\omega, \phi$	800

$\eta$ - $\eta'$  physical states originated by the mixing of neutral  $I = 0$  octet  $\eta_8$  and singlet  $\eta_0$  states

# Motivations

- knowing the  $\eta$ – $\eta'$  wave functions probes our understanding of non-perturbative QCD dynamics---**Showing there is a purely gluonic component would establish for the first time that gluons play an independent role also in hadronic spectroscopy.**
  - Also essential to disentangle SM hadronic uncertainties vs. New Physics
- Great effort has been given to this challenge { yet no clear picture has emerged }
- We point out which measurements and theoretical analysis would be most helpful for a more definite conclusion.

## $\eta-\eta'$ Mixing

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos \theta_P & -\sin \theta_P \\ \sin \theta_P & \cos \theta_P \end{pmatrix} \begin{pmatrix} \eta_8 \\ \eta_0 \end{pmatrix}$$

Just for orientation:

quadratic (linear) Gell-Mann-Okubo

mass formula

$$\theta_P \approx -10^\circ \quad (\theta_P \approx -23^\circ)$$

quark-flavor basis:

$$\begin{cases} \eta_{NS} = \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) \\ \eta_S = s\bar{s} \end{cases}$$

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos \varphi_P & -\sin \varphi_P \\ \sin \varphi_P & \cos \varphi_P \end{pmatrix} \begin{pmatrix} \eta_{NS} \\ \eta_S \end{pmatrix}$$

$$\theta_P = \varphi_P - \arctan \sqrt{2} = \varphi_P - 54.7^\circ$$



$$\varphi_P \approx 45^\circ \quad (\varphi_P \approx 32^\circ)$$

## Two Mixing Angles Scenario: octet-singlet basis

- From the late 90's (*Leutwyler, Kraiser, Kroll, Stech, Feldmann etc.*) been shown the mixing cannot be adequately described by a single angle; the fact that the decay constants follow the pattern of state mixing is an a-priori assumption
- Due to SU(3) breaking ( $f_K \neq f_\pi$ ), mixing of decay constants does not follow the same pattern of state mixing

$$J_{\mu L}^a = \bar{q} \frac{\lambda^a}{\sqrt{2}} \gamma_\mu \frac{1 - \gamma_5}{2} q$$

$$J_{\mu R}^a = \bar{q} \frac{\lambda^a}{\sqrt{2}} \gamma_\mu \frac{1 + \gamma_5}{2} q$$

$$J_\mu^a = J_{\mu R}^a + J_{\mu L}^a$$

$$f_\eta^8 \neq f_\eta^0 \neq f_{\eta'}^8 \neq f_{\eta'}^0$$

$$\begin{pmatrix} f_\eta^8 & f_\eta^0 \\ f_{\eta'}^8 & f_{\eta'}^0 \end{pmatrix} = \begin{pmatrix} f_8 \cos \theta_8 & -f_0 \sin \theta_0 \\ f_8 \sin \theta_8 & f_0 \cos \theta_0 \end{pmatrix}$$

$$\langle 0 | J_{\mu 5}^a(0) | P(q) \rangle = i f_p^a q_\mu$$

The estimated difference  $\theta_8 - \theta_0$  can be large  $[-12^\circ - 19^\circ]$

# Quark flavor basis

- The smallness of the mixing angles is consistent with the OZI-rule, i.e. amplitudes that involve quark-antiquark annihilation into gluons are suppressed

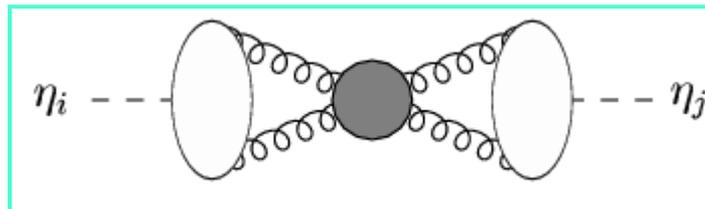
– vector meson sector: mixing angle  $\approx 3^\circ$

$$\omega \approx q\bar{q} \quad \phi \approx s\bar{s}$$

$$\begin{cases} \eta_{NS} = \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) \\ \eta_S = s\bar{s} \end{cases}$$

- In the pseudoscalar sector,  $U(1)_A$  anomaly induces a significant mixing between the fields  $\eta_{NS}$  and  $\eta_S$ .

$$\partial^\mu J_{\mu 5}^u = \partial^\mu (\bar{u} \gamma_\mu \gamma_5 u) = 2 m_u (\bar{u} i \gamma_5 u) + \frac{\alpha_s}{4\pi} G \tilde{G}$$

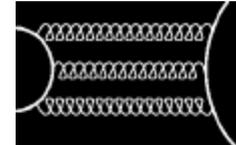


# Two mixing angle scenario: flavor basis

In principle we have

$$\begin{pmatrix} f_{\eta}^q & f_{\eta}^s \\ f_{\eta'}^q & f_{\eta'}^s \end{pmatrix} = \begin{pmatrix} f_q \cos \phi_q & -f_s \sin \phi_s \\ f_q \sin \phi_q & f_s \cos \phi_s \end{pmatrix}$$

- However, while the mixing is large ( $\approx 40^\circ$ ), the difference between the two mixing angles is determined by OZI-rule violating contribution
- OZI-rules rigorous when  $N_c \rightarrow \infty$  or  $\alpha \rightarrow 0$



$\phi_q$  and  $\phi_s$  nearly coincide

# Mix with Gluonium

- The  $\eta'$  meson is a good candidate to have a sizeable gluonic content, (while the  $\eta$  meson is well understood as an SU(3)-flavor octet with a small singlet admixture)

$$\left\{ \begin{array}{l} \eta' = X_{\eta'} \frac{1}{\sqrt{2}} |u\bar{u} + d\bar{d}\rangle + Y_{\eta'} |s\bar{s}\rangle + Z_{\eta'} |glue\rangle \\ \eta = X_{\eta} \frac{1}{\sqrt{2}} |u\bar{u} + d\bar{d}\rangle + Y_{\eta} |s\bar{s}\rangle \end{array} \right.$$

$$\left\{ \begin{array}{l} X_{\eta'} = \cos \phi_G \sin \phi_P; Y_{\eta'} = \cos \phi_G \cos \phi_P; Z_{\eta'} = \sin \phi_G \\ X_{\eta} = \cos \phi_P; Y_{\eta} = -\sin \phi_P \end{array} \right.$$

Mixing with heavier pseudoscalar mesons is ignored

# *Electromagnetic and strong transitions*

## 1. Radiative vector and pseudoscalar meson decays

$$\psi', \psi, \phi \rightarrow \gamma \eta' \quad \text{vs.} \quad \gamma \eta$$

$$\rho, \omega \rightarrow \gamma \eta$$

$$\eta' \rightarrow \gamma \omega, \gamma \rho$$

## 2. Decays into two photons or production in collisions:

$$\eta' \rightarrow 2\gamma \quad \text{vs.} \quad \eta \rightarrow 2\gamma$$

$$\gamma\gamma \rightarrow \eta \quad \text{vs.} \quad \gamma\gamma \rightarrow \eta'$$

## 3. Decays of $\psi$ into PV final states with the vector meson acting as a 'flavor filter':

$$\psi \rightarrow \rho/\omega/\phi + \eta \quad \text{vs.} \quad \eta'$$

# Radiative $\rho/\omega/\phi$ Decays: first modern analyses

no gluonium hypothesis

KLOE 07:  $\phi_P = (41.3 \pm 0.3_{stat} \pm 0.7_{sys} \pm 0.6_{th})^\circ$

Escribano,  
Nadal 07

$$\phi_P = (42.7 \pm 0.7)^\circ$$

Thomas 07:

no form factors

$$\phi_P = (41.7 \pm 0.5)^\circ$$

with form factors

$$\exp\left(-|p|^2/(8\beta^2)\right)$$

$$\phi_P = (42.8 \pm 0.8)^\circ$$

allowing for gluonium

$$\begin{cases} \phi_P = (39.7 \pm 0.7)^\circ \\ Z_{\eta'}^2 = 0.14 \pm 0.04 \end{cases}$$

$$\begin{cases} \phi_P = (42.6 \pm 1.1)^\circ \\ Z_{\eta'}^2 = 0.01 \pm 0.07 \end{cases}$$

$$\begin{cases} \phi_P = (41.7 \pm 0.5)^\circ \\ Z_{\eta'}^2 = 0.04 \pm 0.04 \end{cases}$$

$$\begin{cases} \phi_P = (41.9 \pm 0.7)^\circ \\ Z_{\eta'}^2 = 0.10 \pm 0.08 \end{cases}$$

# Radiative Decays

## KLOE vs. th analyses:

OZI-rule reduces considerably possible transitions and their respective VP wave-functions overlaps ( $P \rightarrow V \gamma$ ,  $V \rightarrow P \gamma$ )

KLOE: VP-overlap parameters without gluonium, (from *Bramon, Escribano and Scadron 01*)

Theory: VP-overlap parameters determined allowing for gluonium  $\Rightarrow$  new values

Differently from the th analyses, Kloe use also constraints from  $\eta' \rightarrow \gamma\gamma$ , where no VP-overlap parameters enter in the analysis: this guarantee “independence” from the parameters in a reasonable range.

KLOE

$$\frac{\Gamma(\phi \rightarrow \eta' \gamma)}{\Gamma(\phi \rightarrow \eta \gamma)}$$

$$\frac{\Gamma(\eta' \rightarrow \rho \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)}$$

$$\frac{\Gamma(\eta' \rightarrow \gamma \gamma)}{\Gamma(\pi^0 \rightarrow \gamma \gamma)}$$

no VP overlap

$$\frac{\Gamma(\eta' \rightarrow \omega \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)}$$

Escribano, Nadal

$$\Gamma(\phi \rightarrow \eta \gamma)$$

$$\Gamma(\phi \rightarrow \eta' \gamma)$$

$$\Gamma(\eta' \rightarrow \rho \gamma)$$

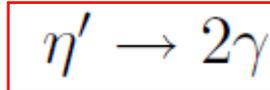
$$\Gamma(\eta' \rightarrow \omega \gamma)$$

# *New KLOE data (2009)*

- Message not truly inconsistent (considering the stated uncertainties), but ambivalent
  - Some studies (KLOE) point to a significant gluonic component, others not
- the 10th discussion has prompted a new KLOE update (2009)
  - Results confirmed
    - no gluons  $\phi = (41.4 \pm 0.5)^\circ$
    - allowing gluons  $\begin{cases} \phi = (40.4 \pm 0.6)^\circ \\ Z_{\eta'}^2 = 0.115 \pm 0.036 \end{cases}$
- the actual difference with KLOE values appears not due to a wrong set of variables, but to the inclusion in the analysis of  $\eta' \rightarrow 2\gamma$

# $\eta'$ decays

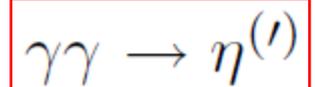
- In lowest order all possible  $\eta'$  strong decays are forbidden by C, CP invariance and G-parity conservation.
- First order electromagnetic decays are forbidden as well, or occur at a suppressed rate because involving an anomaly.
  - The first allowed decay is therefore the second-order electromagnetic transition



- key role in the mixing parameters determination
  - MD-1@Novosibirsk (1985), ASP@SLAC 85

$$\phi_P = 34.9^\circ \pm 2.2^\circ$$

- can be exploited also looking at the inverse processes, namely
  - Crystal Ball Collab@DESY (1983)



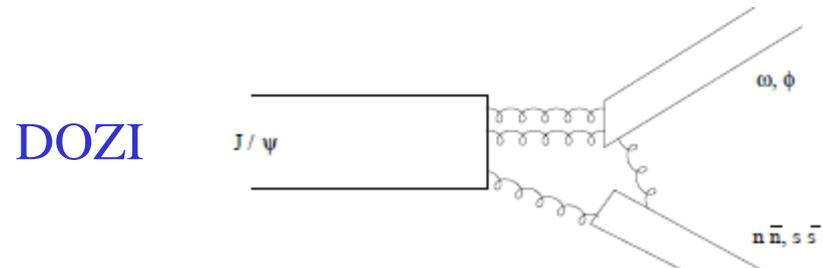
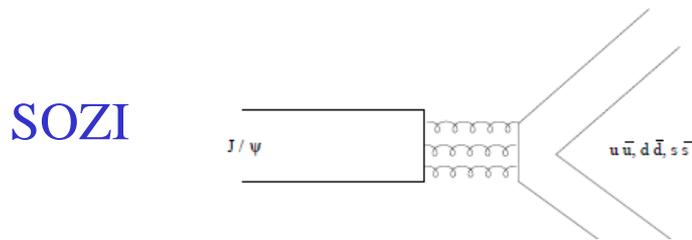
$$\phi_P = 32.3^\circ \pm 1.2^\circ$$

Significantly below previous results

- NEW Results BaBar (2011) disagree with theoretical prediction: (*Bakulev et al 2001*): admixture of the two-gluon component?

## $\psi \rightarrow PV$ : also ambivalent results

- Mark III (1985) :  $(35 \pm 18)\%$  of the  $\eta'$  wave function can be attributed to gluonium or radial excitation
  - analysis is based on the assumption that decays proceed via singly disconnected diagram (SOZI), omitting the doubly disconnected (DOZI) diagram
- Mark III (1988) : any gluonium contribution to the  $\eta'$  wave function is ruled out
  - including DOZI diagrams



- More recent re-analyses of the hadronic  $J/\psi$  and  $\psi'$  decays (including DOZI) furnish a consistent description in terms of one mixing angle with a suggestion of some gluonic component of the  $\eta'$ . F.i.

– no form factors  $\phi_P = (45 \pm 4)^\circ$   $\phi_{glue} = (33 \pm 13)^\circ$

– form factors  $\phi_P = (46^{+4}_{-5})^\circ$   $\phi_{glue} = (44 \pm 9)^\circ$  (*Thomas 2007*)

# Possible future experimental scenarios

Processes	$(\delta\Gamma/\Gamma)_I$	$(\delta\Gamma/\Gamma)_{II}$	$(\delta\Gamma/\Gamma)_{III}$	$(\delta\Gamma/\Gamma)_{IV}$	$(\delta\Gamma/\Gamma)_V$
$\phi \rightarrow \eta'\gamma$	3.5%	3.5%	1.7%	1.7%	1%
$\phi \rightarrow \eta\gamma$	2%	2%	1%	1%	1%
$\eta' \rightarrow \omega\gamma$	9%	4.5%	9%	4.5%	1.7%
$\eta' \rightarrow \rho\gamma$	5%	5%	5%	2.5%	1.7%
$\rho \rightarrow \eta\gamma$	7%	7%	7%	3.4%	7%
$\omega \rightarrow \eta\gamma$	9%	9%	9%	4.5%	9%
$\phi_P$	$(40.6 \pm 0.9)^\circ$	$(40.1^{+1.0}_{-0.8})^\circ$	$(40.6^{+0.6}_{-0.5})^\circ$	$(40.4 \pm 0.5)^\circ$	$(40.1^{+0.3}_{-0.4})^\circ$
$Z_{\eta'}^2$	$(0.09 \pm 0.05)$	$(0.13 \pm 0.06)$	$(0.09 \pm 0.03)$	$(0.10 \pm 0.03)$	$(0.13 \pm 0.02)$
$\chi^2/Ndof$	6.11/4	11.12/4	6.12/4	21.44/4	82.36/4

- "I" : actual uncertainties in the exp input values (PDG 2010)
- "II" : improvement by studying  $\eta' \rightarrow \omega\gamma$  using 1) 20 fb<sup>-1</sup> (KLOE2) 2) selection efficiency 20% 3) neglecting background subtraction—
  - limiting factor : uncertainty in the total  $\eta'$  width
- "III" : improvement in determination of the partial widths for  $\phi \rightarrow \eta^{(\prime)}\gamma$
- "IV" : improvement in determination of the partial widths for  $\eta^{(\prime)} \rightarrow \rho\gamma$
- "V" : an uncertainty of 1% on the measure of branching ratios for  $\eta'$  decays and of 1.4% for the  $\eta'$  full width.

# Experimental status and prospects

Conclusions from previous table :

*both  $BR(\eta^{(\prime)} \rightarrow \rho\gamma)$  and  $\Gamma_{\eta^{(\prime)}}$  to be measured more accurately.*

- $\eta' \rightarrow \omega\gamma$  partial width with relative error 9% (PDG 2010) Relevant experiment (ANL-E-397, 1977) bases on 68 events.
- $\phi \rightarrow \eta^{(\prime)}\gamma$  partial width mainly due to KLOE07; error dominated by systematics due to the secondary  $\eta'$  branching ratio
- $\eta' \rightarrow \rho\gamma$  partial width with relative error 5% (PDG 2010) Relevant experiment (1969) bases on 298 events
  
- $\Gamma_{\eta'} = 0.194 \pm 0.009$  (PDG 2010);  $\Gamma_{\eta'} = 0.30 \pm 0.09$  direct meas.(1996,1979)
  - Crystal Ball@MAMI (started 2009)
  - New insight could come from production in  $\gamma\gamma$  fusion (KLOE)

# Weak Decays of Charm and Beauty Hadrons

- $\eta/\eta'$  wave functions important input for several weak  $D$  and  $B$  decays
  - CP asymmetries involving  $\eta/\eta'$  in the final states
  - Control NP vs SM hadronic uncertainties
- Phenomenological approach (while waiting for lattice)
  - pioneering calculation RBC-UKQCD of  $\eta/\eta'$  masses and mixing using  $N_f = 2+1$  (2010)

$$\theta_P = -14.1(2.8)^\circ \quad (\varphi \approx 40.7^\circ)$$

# Light Flavour Spectroscopy in Semileptonic Decays

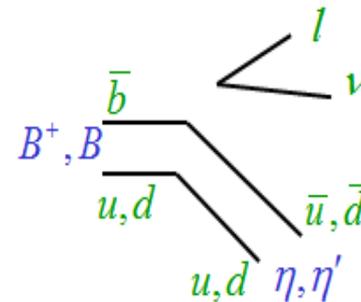
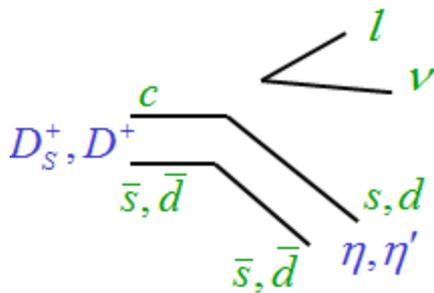
$$D_s^+ \rightarrow \eta^{(\prime)} l^+ \nu$$

$$D^+ \rightarrow \eta^{(\prime)} l^+ \nu$$

$$B^+ \rightarrow \eta^{(\prime)} l^+ \nu$$

- Spectator diagram dominance**

- Cabibbo allowed  $c \rightarrow s$ , suppressed  $c \rightarrow d$  and CKM suppressed  $b \rightarrow u$



- If  $\eta/\eta'$  without gluonic content and pole ansatz for form factors

$$\Gamma(D_s \rightarrow \eta' e^+ \nu) / \Gamma(D_s \rightarrow \eta e^+ \nu) \propto \cot^2 \phi$$

$$\phi = (41.3 \pm 5.3)^\circ \quad \text{Feldmann, Kroll, Stech 98}$$

- Allowing gluonic content

$$|\eta'\rangle \simeq \cos \phi_G \sin \phi_P |\eta_q\rangle + \cos \phi_G \cos \phi_P |\eta_s\rangle + \sin \phi_G |gg\rangle$$

$$|\eta\rangle \simeq \cos \phi_P |\eta_q\rangle - \sin \phi_P |\eta_s\rangle$$

$$\Gamma(D_s \rightarrow \eta' e^+ \nu) / \Gamma(D_s \rightarrow \eta e^+ \nu) \propto \cot^2 \phi \cos^2 \phi$$

Given  $\phi = (37.74 \pm 2.6)^\circ \longrightarrow \phi_G \approx 20.3 \quad \text{Anisovic et al. 97}$

• new (compatible)  
CLEO-c 09 data

$$\frac{\Gamma(D_s^+ \rightarrow \eta' l^+ \nu) / \Gamma(D_s^+ \rightarrow \eta l^+ \nu)}{\Gamma(D^+ \rightarrow \eta' l^+ \nu) / \Gamma(D^+ \rightarrow \eta l^+ \nu)} \simeq \cot^4 \phi_P$$

$$\phi_P \approx 40^\circ$$

BESIII expects errors on  $\phi$  going down to 2%

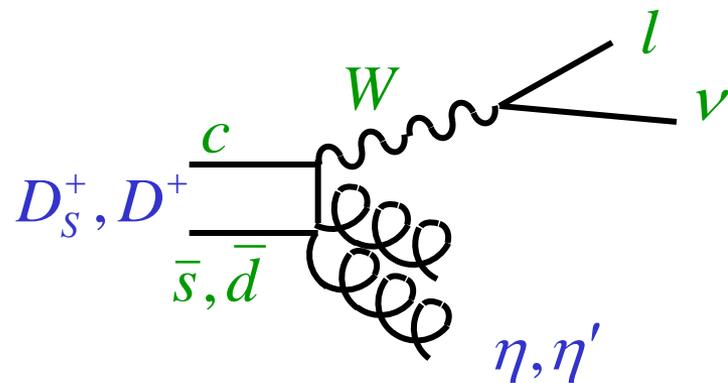
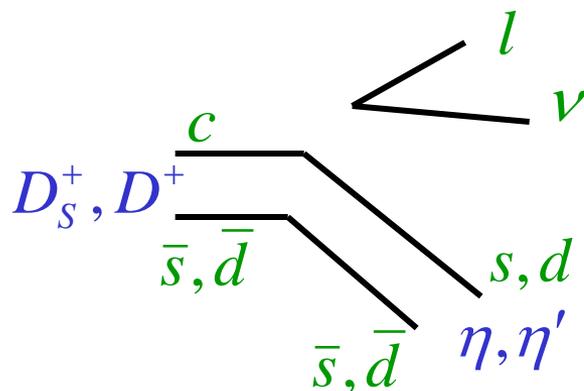
# Weak annihilation (WA) diagrams

Spectator

WA (gluonic)

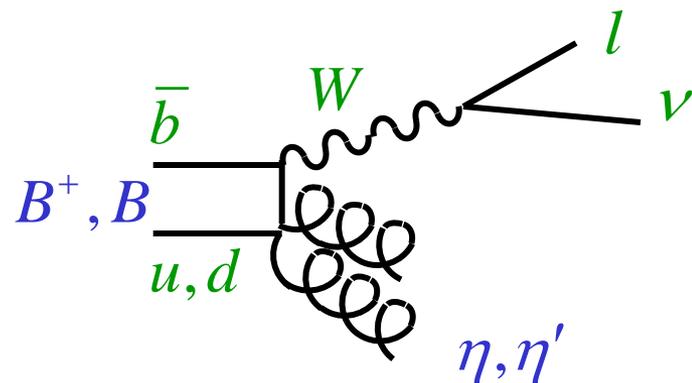
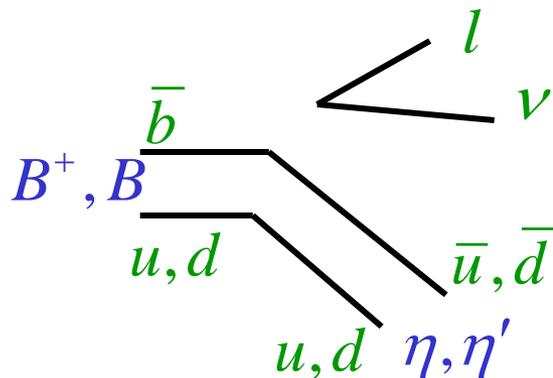
$$D_S^+ \rightarrow \eta(\eta') l^+ \nu$$

$$D^+ \rightarrow \eta(\eta') l^+ \nu$$



$$B^+ \rightarrow \eta(\eta') l^+ \nu$$

$$B \rightarrow \eta(\eta') l^+ \nu$$



# *WA for precision studies*

- WA no more than a nonleading contribution to inclusive rates, BUT could affect exclusive modes considerably
- strength depends on:
  - size of the gg component in the wave functions
  - how much gg radiation one can expect in each semileptonic channel
  - might come from the interference with the spectator amplitude, it can a priori enhance or reduce those
- A recent analysis based on inclusive semileptonic D decays, which considers both the widths and the lepton energy moments, shows no clear evidence of WA effects *Gambino, Kamenik 2010*
- No extensive exclusive theoretical analysis yet

# $B^\pm$ semileptonic decays

Same than  $D^\pm$  only large  $q^2$  and CKM suppressed by  $|V_{ub}|^2$

- first evidence of  $B^+ \rightarrow \eta' l^+ \nu$  by CLEO in 2008
- newest BaBar results (2011) (with a significance of  $3.0 \sigma$ ).
  - an order of magnitude smaller than the CLEO result

$$\mathcal{B}(B^+ \rightarrow \eta' l^+ \nu) / \mathcal{B}(B^+ \rightarrow \eta l^+ \nu) = 0.67 \pm 0.24_{stat} \pm 0.11_{syst}$$

seems to allow a large gluonic singlet contribution

- $\mathcal{B}(B_s \rightarrow \eta' l^+ l^-)$  potentially informative on the gluonic content, exp challenging. In SM BR  $10^{-7}$ - $10^{-8}$  (Super-flavour factories)

# Charmless Hadronic $B$ Decays

- Charmless hadronic  $B$  decays provide valuable tests for the pattern of CP violation in the CKM framework
- CP asymmetries

$$A(t) \equiv \frac{\Gamma(\bar{B}^0(t) \rightarrow \eta' K_S) - \Gamma(B^0(t) \rightarrow \eta' K_S)}{\Gamma(\bar{B}^0(t) \rightarrow \eta' K_S) + \Gamma(B^0(t) \rightarrow \eta' K_S)} = -C_{\eta' K} \cos(\Delta mt) + S_{\eta' K} \sin(\Delta mt)$$

$$S_{\eta' K} \equiv \frac{2\text{Im}(\lambda_{\eta' K})}{1 + |\lambda_{\eta' K}|^2}, \quad C_{\eta' K} \equiv \frac{1 - |\lambda_{\eta' K}|^2}{1 + |\lambda_{\eta' K}|^2}, \quad \lambda_{\eta' K} \equiv -e^{-2i\beta} \frac{A(\bar{B}^0 \rightarrow \eta' \bar{K}^0)}{A(B^0 \rightarrow \eta' K^0)}$$

in the SM, if one single dominant amplitude

$$C_{\eta' K_S} = 0$$

$$S_{\eta' K_S} = \sin 2\beta$$

# *CP violation*

- Decays dominated by single  $b \rightarrow s$  penguin amplitude; In the SM

$$C_{\eta' K_S} = 0$$

$$S_{\eta' K_S} = \sin 2\beta$$

Corrections for suppressed diagrams  $\Delta S_f = S_f - \sin 2\beta$

Estimates in the ranges  $(-0.03, 0.03)$  (QCDF and SCET) and in the range  $(-0.05, 0.09)$  from SU(3) symmetry bounds

*Beneke 05, Williamson, Zupan 06,  
Cheng, Chia, Soni 05, Gronau, Rosner,  
Zupan, 06*

$$C_{\eta' K_S} = 0.03 \pm 0.07$$

$$S_{\eta' K_S} = +0.67 \pm 0.11$$

*Belle 11*

- Compatible with SM and  $\sin 2\beta = 0.657 \pm 0.036 \pm 0.012$  (*Babar 11*) from

$$B^0 \rightarrow J/\psi K_s^0$$

# *Experimental puzzle*

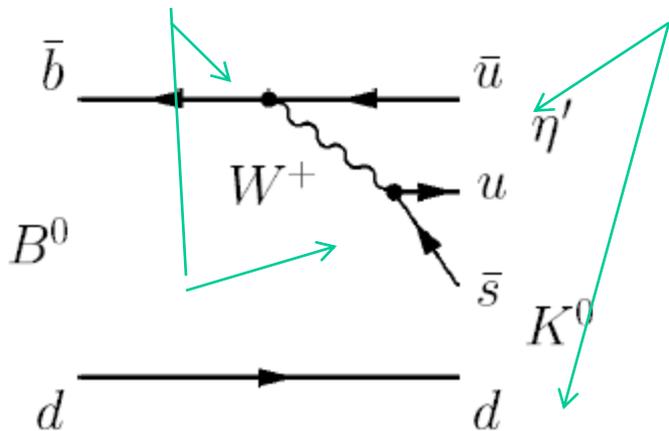
$$\mathcal{B}(B^0 \rightarrow K^0 \eta') = (6.6 \pm 0.4) \times 10^{-5} \gg \mathcal{B}(B^0 \rightarrow K^0 \eta) = (1.1 \pm 0.4) \times 10^{-6}$$

$$\mathcal{B}(B^\pm \rightarrow \eta' K^\pm) = (71.5 \pm 1.3 \pm 3.2) 10^{-6} \gg \mathcal{B}(B^\pm \rightarrow \pi^0 K^\pm) = (13.6 \pm 0.6 \pm 0.7) 10^{-6}$$

- SU(3)<sub>f</sub> singlet penguin ( $B_i$  meson triplet,  $M_j^1$  meson nonet), including gluonic contributions
- Anomaly effects, large charm content, NP?

# Penguin dominated decays

- Tree level CKM and color suppressed

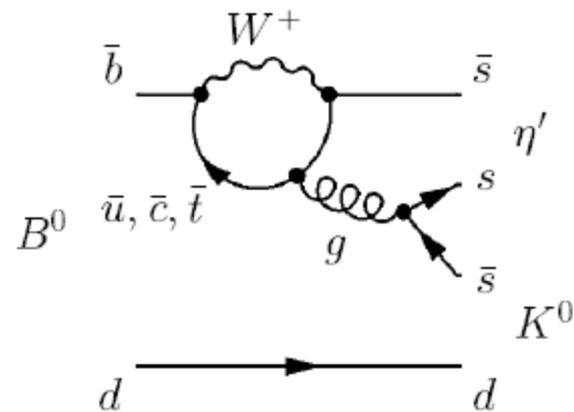
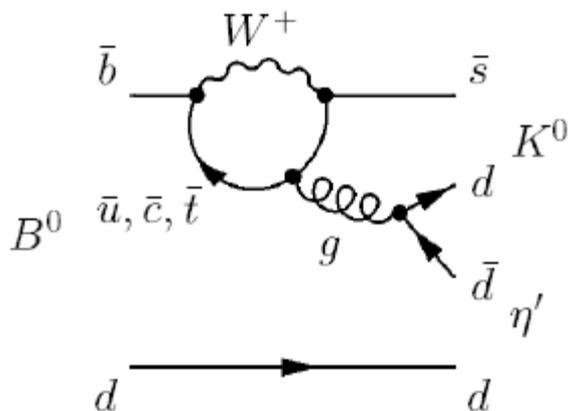


$$\lambda \approx 0.22$$

$$|V_{ub} V_{us}| \approx \lambda^4$$

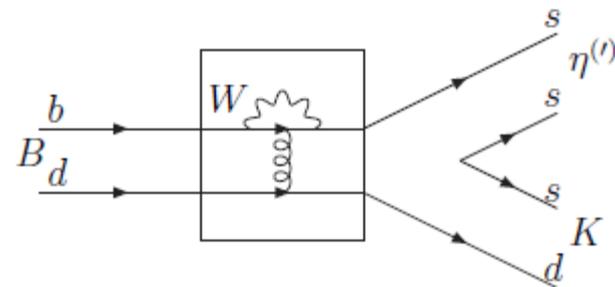
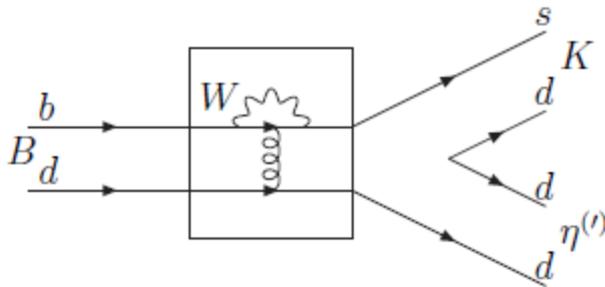
- Penguins CKM not suppressed

$$|V_{tb} V_{ts}| \approx \lambda^2$$

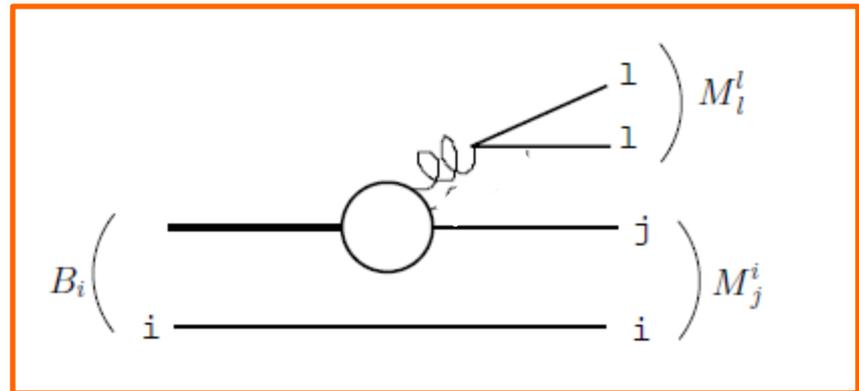
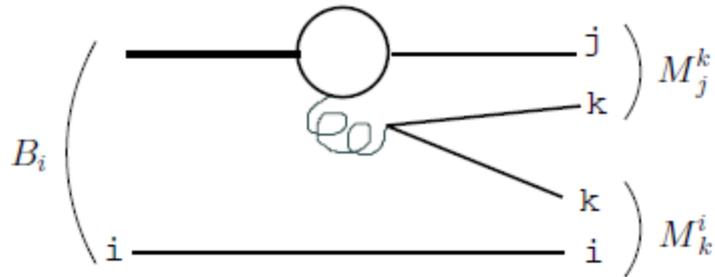


# Other contributing penguins

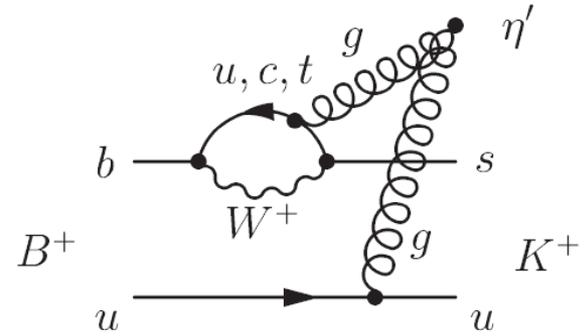
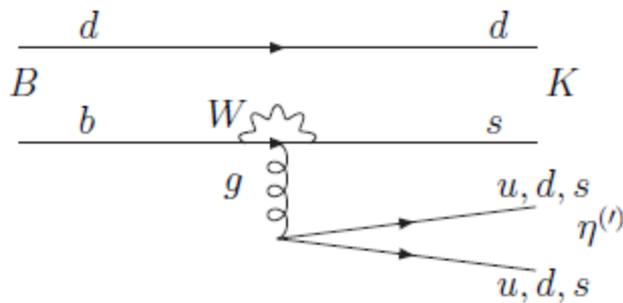
- Annihilation penguin contribution



- SU(3)fl singlet penguin ( $B_i$  meson triplet,  $M_j^l$  meson nonet)



# Singlet contribution



- Additional contribution to the SU(3) singlet contribution: fusion of gluons, one gluon from  $b \rightarrow sg$  process and another one from spectator.

A sizeable gluonium contribution to the  $\eta'$  meson could play an important role: the contribution of the diagram in which two gluons are directly attached to gluonium in  $\eta'$  in principle important

# *Theoretical approaches*

Within large errors BR compatible with interference among different contributions

- QCDF and SCET : sizable gluonic contributions to the  $B \rightarrow \eta'$  form factor  
*Beneke, Neubert 03;*  
*Williamson, Zupan 06*
- pQCD : impact of the gluonic component numerically very small  
*Y.-Y. Charng, T. Kurimoto, H.-n. Li 06*

Major analyses prior to 06, relying on old experimental data

## *More recent data*

2008-2010 CLEO/BaBar

B semileptonic  $\mathcal{B}(B^+ \rightarrow \eta' l^+ \nu) < \mathcal{B}(B^+ \rightarrow \eta l^+ \nu)$

- 2010 BaBar

$$\Gamma(B \rightarrow K^* \eta') < \Gamma(B \rightarrow K^* \eta)$$

- 2007 CLEO

$$D_s^+ \rightarrow K^+ \eta' \approx D_s^+ \rightarrow K^+ \pi^0$$

## Conclusions

Radiative $\rho\omega\phi$	$\theta_P$	$Z_{\eta'}^2$
KLOE ( $+\eta' \rightarrow \gamma\gamma$ )	$(39.7 \pm 0.7)^\circ$	$0.14 \pm 0.04$
Escribano	$(41.4 \pm 1.3)^\circ$	$0.04 \pm 0.09$
Thomas	$(41.7 \pm 0.5)^\circ$	$0.04 \pm 0.04$
Thomas F.F.	$(41.9 \pm 0.7)^\circ$	$0.10 \pm 0.04$

$J/\psi \rightarrow VP$	$\theta_P$	$Z_{\eta'}^2$
Thomas	$(45 \pm 4)^\circ$	$0.30 \pm 0.21$
Thomas F.F.	$(46 + 4 / - 5)^\circ$	$0.48 \pm 0.16$
Escribano	$(44.6 \pm 4.4)^\circ$	$0.29 + 0.28 / - 0.26$

Semileptonic decay	$\theta_P$	$Z_{\eta'}^2$
$\frac{\Gamma(D_s^+ \rightarrow \eta' l^+ \nu)}{\Gamma(D_s^+ \rightarrow \eta l^+ \nu)}$	$(37.7 \pm 2.6)^\circ$	0.12

- The different determinations of  $\eta$ – $\eta'$  mixing angle are generally consistent, but show relevant model and mode dependence
- the message concerning  $\eta'$  gluonium content remains ambivalent
- More dedicated studies are necessary (including theoretical updates with new data) while waiting for lattice (but wait seated...)
- Also  $D^+$ ,  $D_s^+$  and  $B^+$  decays must be included in traditional investigations to check  $\eta'$  gluonium role