

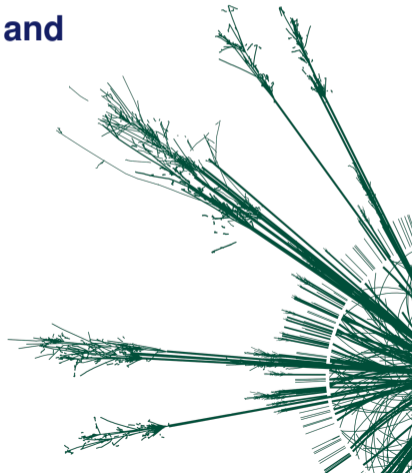
La Thuile 2024

Search for **dark mesons** decaying to top and bottom quarks with the ATLAS detector

Jochen Jens Heinrich for the ATLAS collaboration

08 March, 2024

[See ATLAS-CONF-2023-021](#)

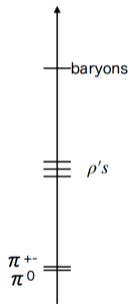


Imagine...

- ... a new strongly-coupled "dark" sector of vector-like fermions
- ... which transform under both the new dark group and the EW part of the SM group
- ... and also permits Higgs interactions

Mesons in QCD

- $SU(3)_{\text{colour}}$
- Exact $SU(2)$ isospin symmetry
- Fermions allowed to transform under EW and QCD
- Triplet of pions (π^- , π^0 , π^+),...



"Dark" mesons

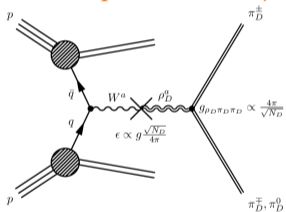
- $SU(N_{\text{dark}})$ (Here: $N_{\text{dark}} = 4$)
- Exact $SU(2)$ custodial symmetry
- Fermions allowed to transform under EW and "dark colour"
- Triplet of dark pions (π_D^- , π_D^0 , π_D^+),...

In [arXiv:1809.10183](https://arxiv.org/abs/1809.10183) and [arXiv:1809.10184](https://arxiv.org/abs/1809.10184) models containing a strongly-coupled, $SU(2)$ dark flavour symmetry preserving dark sector is developed and explored \Rightarrow Stealth Dark Matter [[arXiv:1503.04203](https://arxiv.org/abs/1503.04203)]

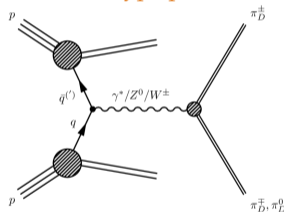
Production of Dark Pions

- Dark pions are always **pair-produced** and decay to pure SM states
- The dark sector is **free of constraints** from precision **EW observables** and **Higgs coupling** measurements

Resonant production via ρ_D



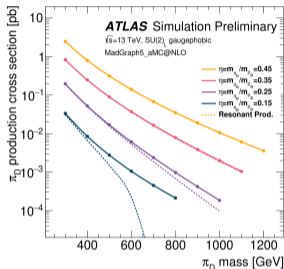
Drell-Yan-type production



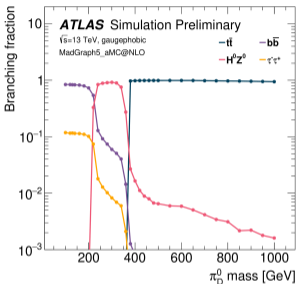
- Dark pion production trivially depends on $m_{\pi_D}/m_{\rho_D} = \eta$ (important parameter!)
- In this work focus on **gaugephobic** dark pion decay behaviour

Dark Pion Decays and Analysis Signature

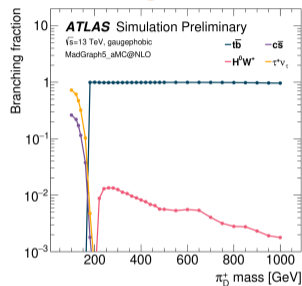
Production cross-section



Neutral π_D



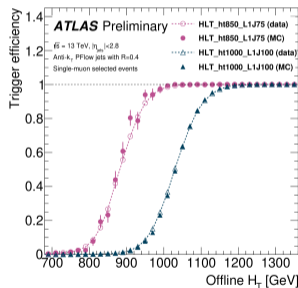
Charged π_D



- Dark pions **decay promptly** back to SM states
- Branching fractions only depend on "gaugephobic-ness" and the mass of the dark pion
- Decays to t 's and b 's dominate once kinematically open \Rightarrow **Consider $t\bar{t}b\bar{b}$ and $t\bar{t}tb$ final states**
- Conduct search in the **all-hadronic** channel (search in **1-lepton channel** is forthcoming)

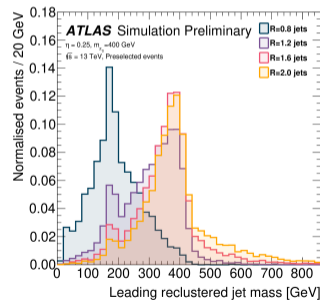
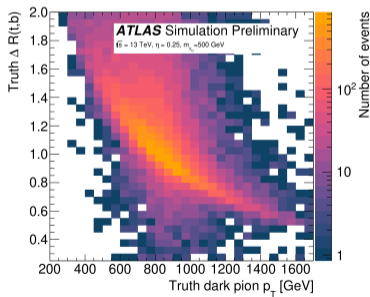
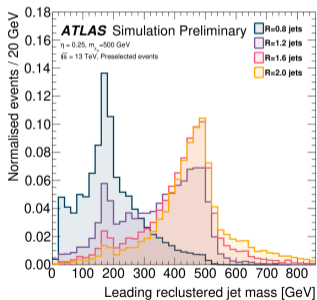
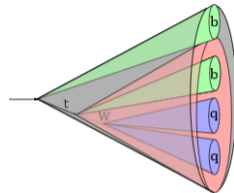
Premise and Online Event Selection

- This work is the **first direct collider search** for this type of model
- Existing limits from reinterpretation of ATLAS and CMS results are fairly weak
- Scanning **2-dimensional parameter space** depending on η and m_{π_D}
- High jet-multiplicity + dark pions from heavy ρ_D decay
 ⇒ **Use lowest unprescaled H_T triggers**
- H_T triggers reach **full efficiency** for large m_{π_D}
- Good agreement between simulated onset and real onset
 ⇒ **Onset has minimal impact on systematics**
- **All-hadronic channel** sensitive to models with low η and low to medium dark pion masses
- **1-lepton channel** more sensitive to higher η values



Dark Pion Reconstruction

- Dark pion reconstruction using **reclustering** of anti- k_t $R = 0.4$ PFlow jets
- R-parameter needs to strike balance between π_D reconstruction and QCD susceptibility
- **$R = 1.2$ has best overall performance** for targeted signal points



Preselection

- **Event quality cuts** that ensure operational detector
- Veto events containing any loose electrons or muons
- Require $H_T > 1150$ GeV
- At least six jets with $p_T > 25$ GeV **to suppress QCD**
- At least 3 b-tagged jets
- At least two reclustered jets with mass > 190 GeV

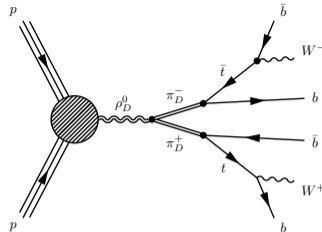
Sample	Yield	Percent SM MC sum
$t\bar{t}$	8756.9	78.0
V + jets	1688.3	15.0
$t\bar{t} + X$	405.6	3.6
SingleTop	375.0	3.3
Multiboson	0.8	< 1
SM sum	11226.5	100
SU2L-45-400	804.7	
SU2L-35-500	867.3	
SU2L-25-500	549.5	
Data	67339	

Process	Generator	PDF	Showering	Tune	Cross section
$t\bar{t}$	POWHEGBOX v2	NNPDF3.0nlo	PYTHIA8	A14	NNLO+NNLL
$t\bar{t}$ +HF	POWHEG BOX RES	NNPDF3.0nlo	PYTHIA8	A14	NNLO
V + jets	SHERPA v2.2.11	NNPDF3.0nnlo	SHERPA	Def.	NLO
Single top	POWHEGBOX v2	NNPDF3.0nlo	PYTHIA8	A14	NLO+NNLL
$t\bar{t}t\bar{t}$	MADGRAPH5_aMC@NLO v2.4.3	NNPDF3.1nlo	PYTHIA8	A14	NLO
$t\bar{t}V$	MADGRAPH5_aMC@NLO v2.3.3	NNPDF3.0nlo	PYTHIA8	A14	NLO
$t\bar{t}H$	POWHEGBOX v2	NNPDF3.0nlo	PYTHIA8	A14	NLO
Other $t\bar{t} + X$	MADGRAPH5_aMC@NLO	NNPDF2.3lo	PYTHIA8	A14	NLO
Multiboson	SHERPA v2.2.1/v2.2.2	NNPDF3.0nnlo	SHERPA	Def.	NLO

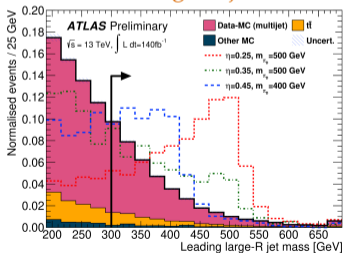
- Dominant background from QCD
 ⇒ **Data-driven background estimate using ABCD method**

Define kinematic variables

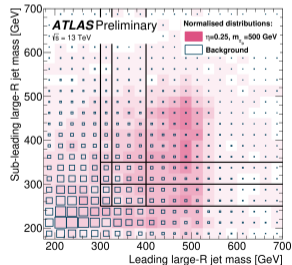
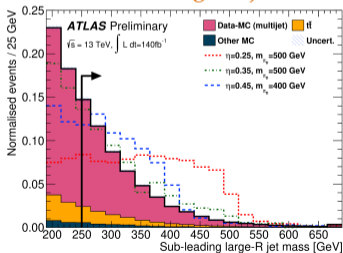
- Define a set of variables motivated by signal kinematics
 - $\Delta R(j, b_2)$: Distance between a RC jet and the second closest b-jet
 - $m_{bb}/p_{T,bb}$: For the closest b-jet pair to an RC jet to suppress QCD
 - $m_{\text{jet},R=1.2}$: Mass of the RC jet, main discriminant
- Select ABCD regions based on $m_{bb}/p_{T,bb}$ and $m_{\text{jet},R=1.2}$
- Subdivided into nine bins based on jet masses



Leading RC jet



Sub-leading RC jet



QCD Multijet Estimation

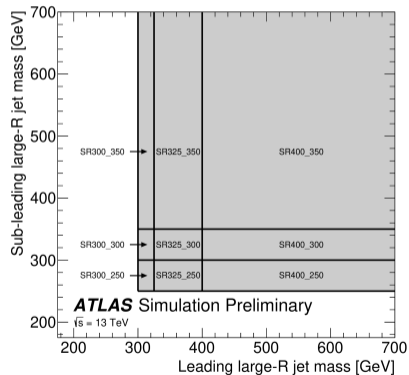
- Use a 4-variable generalisation of the **ABCD method** developed and used in [arXiv:1801.02052](https://arxiv.org/abs/1801.02052)
- Two tags and their slashed "anti-tags" are introduced:
 - bb_i -tag**: i -th $R = 1.2$ jet has 2 b-tagged jets within $R = 1.0$
 - $\pi_{D,i}$ -tag**: i -th $R = 1.2$ jet has a mass within a certain window
- Permutations of all 4 tag states define **16 regions**

		Leading large-R jet			
		$\pi_{D,1}bb_1$	$\pi_{D,1}bb_1$	$\pi_{D,1}bb_1$	$\pi_{D,1}bb_1$
Sub-leading large-R jet	$\pi_{D,2}bb_2$	J	K	L	S
	$\pi_{D,2}bb_2$	B	D	H	N
	$\pi_{D,2}bb_2$	E	F	G	M
	$\pi_{D,2}bb_2$	A	C	I	O

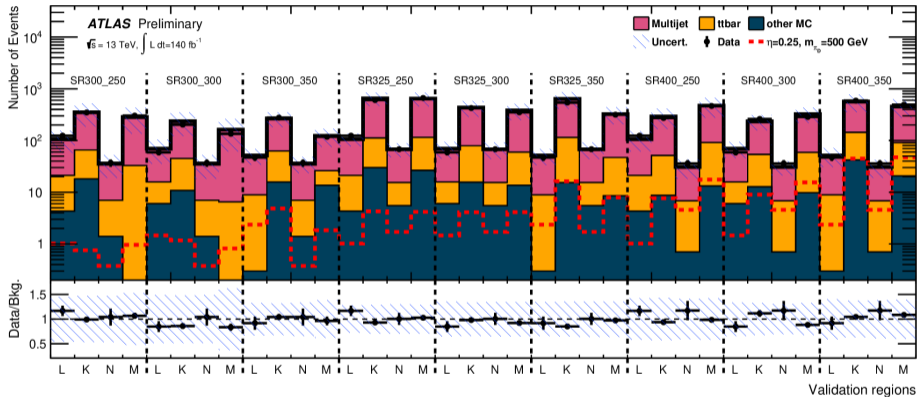
- Extrapolate from orange regions and use 2-tag correction factors for 1-tag correlations
- Signal region S** (red) has all four tags applied
- 3-tag regions K, L, M and N (pink) are validation regions
- Perform independent estimates for each SR bin

Tag selections

	Tag	Variable	Tag selection	Anti-tag selection
Both large- R jets		$m_{bb}/p_{T,M}$		> 0.25
Leading large- R jet	bb_1	$\Delta R(j, b_2)$	< 1.0	≥ 1.0
Sub-leading large- R jet	bb_2	$\Delta R(j, b_2)$	< 1.0	≥ 1.0
Leading large- R jet	$\pi_{D,1}$	$m_{\text{jet},R=1.2}$	$[300 - 325 \text{ GeV}, 325 - 400 \text{ GeV}, > 400 \text{ GeV}]$	$\leq 300 \text{ GeV}$
Sub-leading large- R jet	$\pi_{D,2}$	$m_{\text{jet},R=1.2}$	$[250 - 300 \text{ GeV}, 300 - 350 \text{ GeV}, > 350 \text{ GeV}]$	$\leq 250 \text{ GeV}$



Validation of Multijet Estimate



- Have good agreement in all 36 VRs
 ⇒ Method works very well with no uncontrolled correlations

Systematic Uncertainties

- Main source of systematic uncertainty is the **data-driven multijet estimation**
- Study closure in all VR bins independently via k -factors
 $k_{VR} = (\text{data-MC})/\text{QCD}$
- Also consider statistical uncertainty on k -factors σ_k

$$\text{Systematic: } \sigma_{ABCD} = \sqrt{(1 - \prod_{VR} k_{VR})^2 + \sum_{VR} \sigma_{k_{VR}}^2}$$

	SR300_250	SR300_300	SR300_350
Non-closure uncertainty	40%	45%	1.4%
Stat. uncert. on k -factors	37%	35%	39%
Total Multijet Uncertainty	55%	57%	39%

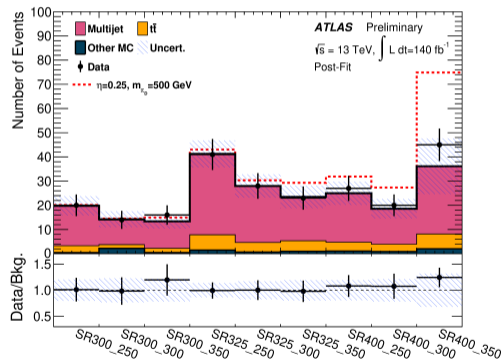
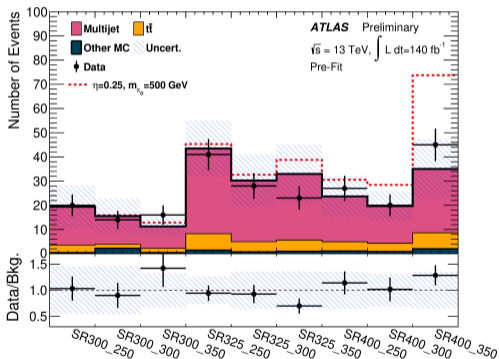
	SR325_250	SR325_300	SR325_350
Non-closure uncertainty	16%	28%	28%
Stat. uncert. on k -factors	29%	29%	29%
Total Multijet Uncertainty	33%	40%	41%

	SR400_250	SR400_300	SR400_350
Non-closure uncertainty	34%	3.2%	29%
Stat. uncert. on k -factors	37%	38%	38%
Total Multijet Uncertainty	51%	38%	48%

- All remaining systematic uncertainties propagated to multijet estimate
- With exception of $t\bar{t}$ modeling uncertainty (up to 6%) and some jet systematics (up to 4%), all systematics are small
⇒ Expected since multijet estimate constrains to data

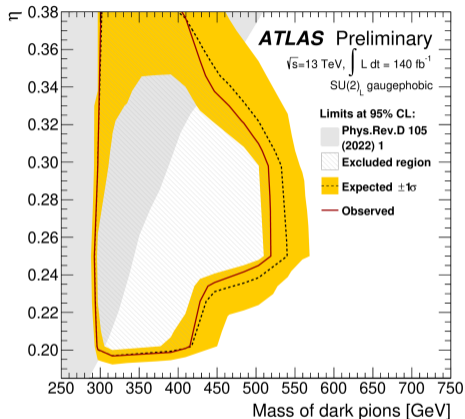
Statistical Analysis

- Perform a **profile likelihood fit** over all signal regions
- Nuisance parameters below 1% impact are pruned away
⇒ **Excellent agreement of data and predicted background**



Results

- Setting limits on dark pion masses given predicted xsections
 ⇒ Observed limits very similar to expected
- Huge gain in coverage over existing limits (grey area) based on reinterpretation
- First direct constraints to this type of model by a collider experiment
- See the full results [here](#)
- Expect 1-lepton channel to increase coverage extensively



Summary

- Searching for **gaugephobic dark pion decays** with signatures **$t\bar{t}b\bar{b}$ or $t\bar{t}t\bar{b}$** in the all-hadronic channel
- Reconstruction of dark pions from reclustered jets
- Cut-based selection with SR binned in large-R jet masses
- **Data-driven estimate of QCD multijet** background that passes validation tests
- Limits **significantly extend existing coverage**

- [Link to publication](#)



Postdoc contract ending in October this year!
Looking for a new member of your group?
⇒ [linkedin.com/in/jochen-jens-heinrich](https://www.linkedin.com/in/jochen-jens-heinrich)

BACKUP MATERIAL

[Link to publication](#)

The Dark Meson Sector of Stealth Dark Matter [\[arXiv:1503.04203\]](https://arxiv.org/abs/1503.04203)

We can explore the phenomenology of such models which has never been done in ATLAS or CMS!

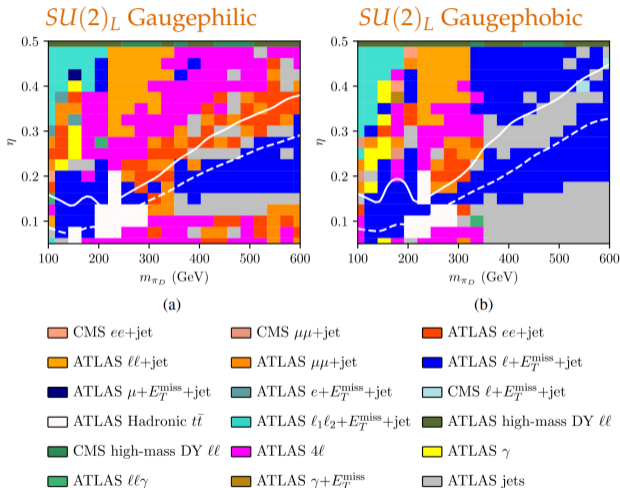
- Higgs interactions break the global (species) symmetries of the dark sector
 \Rightarrow Allows dark pions to decay to SM
- The dark sector is free of constraints from precision EW observables and Higgs coupling measurements
 \Rightarrow Very weak existing limits

Consider models with two flavours of dark fermions \Rightarrow dark pions and one set of dark vector mesons

- Vector-like nature of dark sector permits two possibilities of gauging global flavour symmetry leading to two distinct models of kinetic mixing

$$\begin{aligned}
 SU(2) \text{ kinetic mixing:} & \quad SU(2)_{\text{global flavour}} \leftrightarrow SU(2)_L \\
 U(1) \text{ kinetic mixing:} & \quad SU(2)_{\text{global flavour}} \leftrightarrow SU(2)_R
 \end{aligned}$$

Existing constraints



- [Phys. Rev. D 105, 015008 \(2022\)](#) tested the sensitivity of several ATLAS and CMS analyses
- Solid line is 95% CL exclusion, dashed corresponds to 68% CL
- Even *weaker limits for SU(2)_R*
⇒ *Vast parameter space still open!*
- Other analyses fail to be sensitive because they are optimised for high target masses, expect single production or require large E_T^{miss}
- Strong constraints in models with $\eta > 0.5$ from dilepton searches
⇒ *Only consider models with $\eta < 0.5$ for analysis!*

Object Definitions

Small-R Jets

Small-R Jets:

- $R = 0.4$ Anti- k_t EMPFlow Jets
- $p_T > 20$ GeV
- $|\eta| < 2.8$
- JVT working point: tight
- ForwardJVT working point: tight

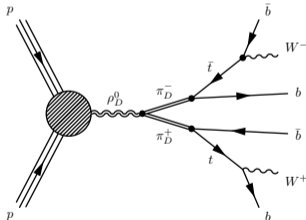
B-tagged Jets:

- $|\eta| < 2.5$
- Only for small-R jets
- DL1r tagger with 77% efficiency working point

Large-R Jets

Reclustered Jets:

- $R = 1.2$ and recluster
- $R = 0.4$ Anti- k_T EMPFlow Jets
- $m > 60$ GeV



Leptons

Baseline electron (for veto):

- ID: MediumLH
- Isolation: None
- $p_T > 10$ GeV
- $|\eta| < 1.37$ OR $1.52 < |\eta| < 2.47$

Baseline muon (for veto):

- ID: Medium
- Isolation: None
- $p_T > 10$ GeV
- $\eta < 2.7$

Details of the multijet estimate

- Consider estimate \hat{K} in region K with standard ABCD

$$\hat{K} = \frac{I \cdot D}{B}$$

Region labels

	$\bar{\pi}_{d,1}bb_1$	$\pi_{d,1}bb_1$	$\bar{\pi}_{d,1}bb_1$	$\pi_{d,1}bb_1$
2nd large-R jet	$\pi_{d,2}bb_2$	J	K	L
$\bar{\pi}_{d,2}bb_2$	B	D	H	N
$\pi_{d,2}bb_2$	E	F	G	M
$\bar{\pi}_{d,2}bb_2$	A	C	I	O
	1st large-R jet			

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$\bar{\pi}_{d,2}bb_2$	B	D	H	N
$\pi_{d,2}bb_2$	E	F	G	M
$\bar{\pi}_{d,2}bb_2$	A	C	I	O
	1st large-R jet			

- This is correct if $\pi_{D,1}$ and $\pi_{D,2}$ uncorrelated, otherwise correct with $k_{\pi_{D,1},\pi_{D,2}} = \frac{F \cdot A}{C \cdot E}$

$$\hat{K} = \frac{I \cdot D}{B} \cdot \frac{F \cdot A}{C \cdot E} = \frac{I \cdot D}{B} \cdot k_{\pi_{D,1},\pi_{D,2}}$$

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Region labels

	$\pi_{d,1}bb_1$	$\pi_{d,1}bb_1$	$\pi_{d,1}bb_1$	$\pi_{d,1}bb_1$	
2nd large-R jet	$\pi_{d,2}bb_2$	J	K	L	S
$\pi_{d,2}bb_2$	B	D	H	N	
$\pi_{d,2}bb_2$	E	F	G	M	
$\pi_{d,2}bb_2$	A	C	I	O	
	1st large-R jet				

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$$\hat{K} = \frac{I \cdot D}{B} \cdot \frac{F \cdot A}{C \cdot E} = \frac{I \cdot D}{B} \cdot k_{\pi_{D,1},\pi_{D,2}}$$

- In a three-variable ABCD estimate with two correlation factors \hat{K} would be computed via

$$\hat{K} = \frac{I \cdot C}{A} \cdot k_{\pi_{D,1},\pi_{D,2}} \cdot k_{\pi_{D,1},bb_2} = \frac{I \cdot C}{A} \cdot \frac{F \cdot A}{C \cdot E} \cdot \frac{D \cdot A}{B \cdot C} = \frac{I \cdot D \cdot F \cdot A}{B \cdot C \cdot E}$$

Details of the multijet estimate

- Consider estimate \hat{K} in region K with standard ABCD

$$\hat{K} = \frac{I \cdot D}{B}$$

		Region labels			
		$\overline{\pi_{d,1}bb_1}$	$\pi_{d,1}bb_1$	$\overline{\pi_{d,1}bb_1}$	$\pi_{d,1}bb_1$
2nd large-R jet	$\overline{\pi_{d,2}bb_2}$	J	K	L	S
	$\pi_{d,2}bb_2$	B	D	H	N
	$\overline{\pi_{d,2}bb_2}$	E	F	G	M
	$\pi_{d,2}bb_2$	A	C	I	O
		1st large-R jet			

- This is correct if $\pi_{D,1}$ and $\pi_{D,2}$ uncorrelated, otherwise correct with $k_{\pi_{D,1},\pi_{D,2}} = \frac{F \cdot A}{C \cdot E}$

$$\hat{K} = \frac{I \cdot D}{B} \cdot \frac{F \cdot A}{C \cdot E} = \frac{I \cdot D}{B} \cdot k_{\pi_{D,1},\pi_{D,2}}$$

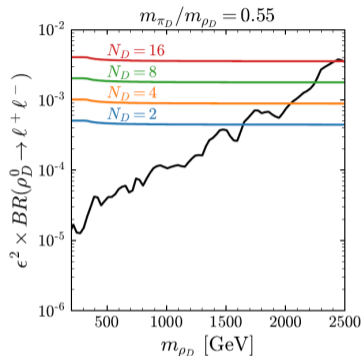
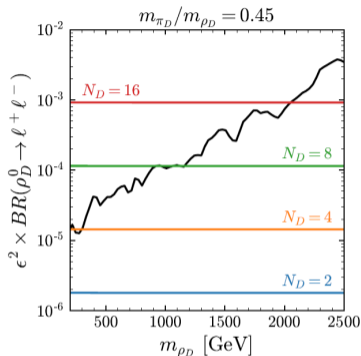
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$$\hat{K} = \frac{I \cdot C}{A} \cdot k_{\pi_{D,1},\pi_{D,2}} \cdot k_{\pi_{D,1},bb_2} = \frac{I \cdot C}{A} \cdot \frac{F \cdot A}{C \cdot E} \cdot \frac{D \cdot A}{B \cdot C} = \frac{I \cdot D \cdot F \cdot A}{B \cdot C \cdot E}$$

- For a four-variable ABCD estimate \hat{S} one needs to consider six correlation factors correcting the traditional estimate \hat{S}'

$$\hat{S} = \hat{S}' \cdot k_{\pi_{D,1},bb_1} \cdot k_{\pi_{D,2},bb_2} \cdot k_{\pi_{D,1},bb_2} \cdot k_{\pi_{D,2},bb_1} \cdot k_{\pi_{D,1},\pi_{D,2}} \cdot k_{bb_1,bb_2}$$

Constraints on ρ_D (see [arXiv:1809.10184](https://arxiv.org/abs/1809.10184))



- Plot shows the **constraints on the kinetic mixing** between SM and ρ_D through absence of a dilepton resonance
- Coloured lines correspond to the $SU(2)_L$ model varying the number of colours N_D
- For $\eta < 0.5$, there is virtually **no constraint on ρ_D when $N_D \leq 4$**