

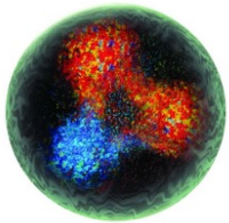
Exclusive Mesons Production with CLAS12 At Jefferson Lab



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Generalized Parton Distributions (GPDs)



Wigner function:
full phase space parton
distribution of the nucleon

$$\rho(x, \vec{k}_T, \vec{b}_T) \quad 5D$$

quark pol.

N/q	U	L	T
nucleon pol.	U	H	\bar{E}_T
L		\tilde{H}	\tilde{E}_T
T	E	\tilde{E}	H_T, \tilde{H}_T

$\bar{E}_T = 2\tilde{H}_T + E_T$

$$\int d^2 k_T$$



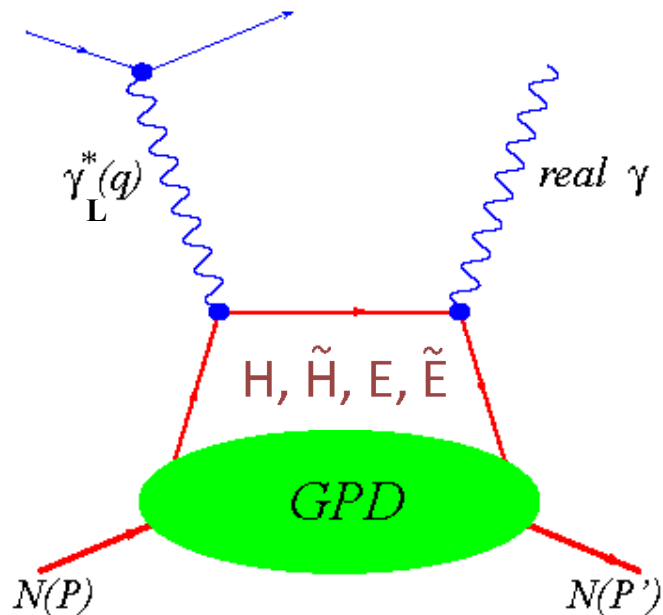
Generalized Parton Distributions
(GPD)



3-D nucleon images in the
transverse coordinate and
longitudinal momentum space

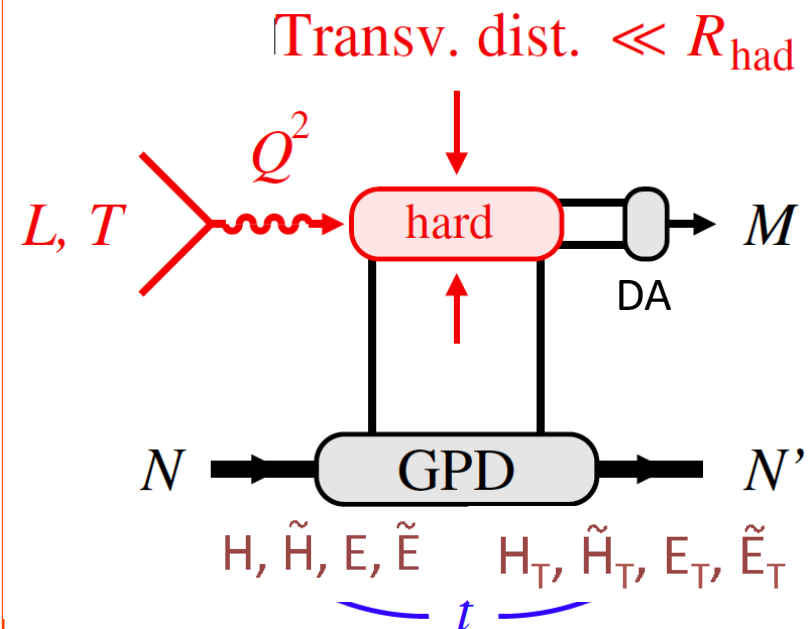
Study GPDs: Deeply Exclusive Processes

Deeply Virtual Compton Scattering (DVCS)



- + Clean process
- Only sensitive to chiral even GPDs

Deeply Virtual Meson Production (DVMP)



- + Access to transversity degrees of freedom described by chiral-odd GPDs
- Distribution Amplitude (DA) is involved as additional soft non pert. quantity

Deeply Virtual Meson Production

	Meson	Flavor
$\mathcal{H}_T, \mathcal{E}_T$ \tilde{H}, \tilde{E}	π^+	$\Delta u - \Delta d$
	π^0	$2\Delta u + \Delta d$
	η	$2\Delta u - \Delta d + 2\Delta s$
\mathcal{H}, \mathcal{E}	ρ^+	$u - d$
	ρ^0	$2u + d$
	ω	$2u - d$
	ϕ	g

H_T is related to the protons tensor charge

$$\delta_T^{u,d} = \int dx H_T^{u,d}(x, \xi = 0, t = 0)$$

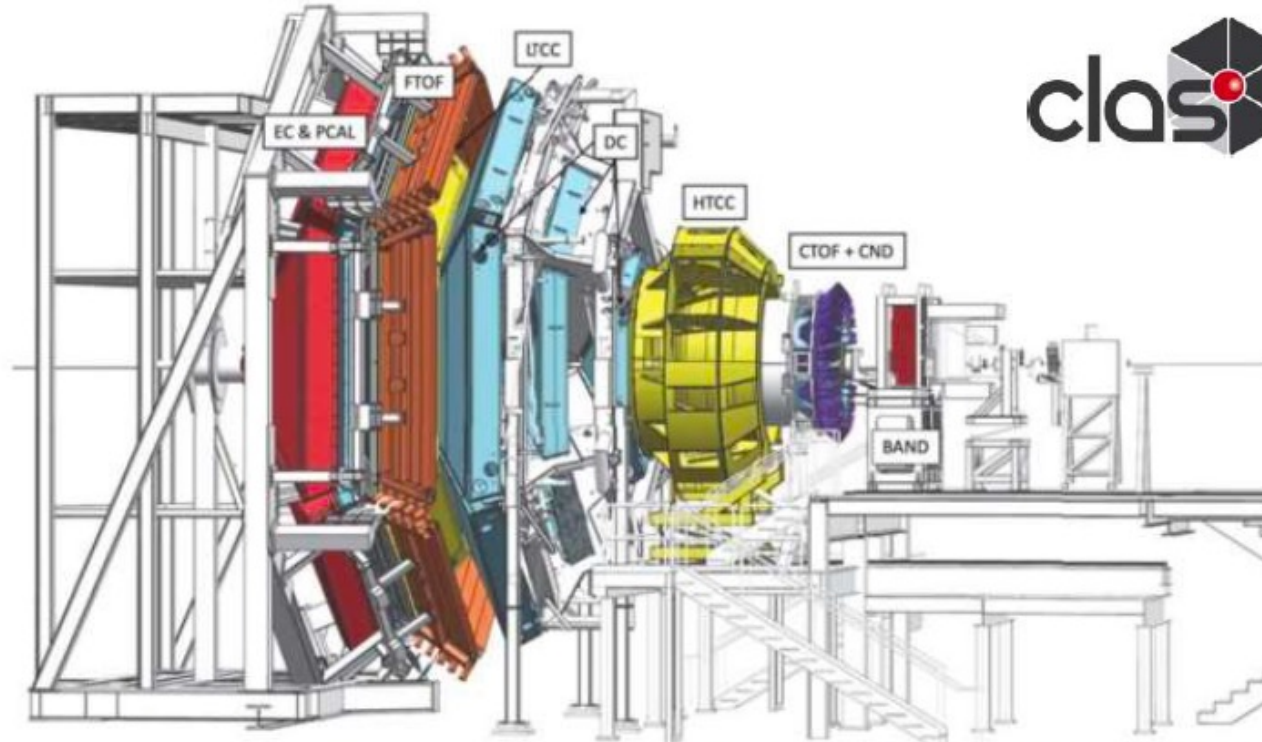
→ Absolute magnitude of transversely polarized valence quarks inside a transv. polarized nucleon

\bar{E}_T is related to the protons anomalous tensor magnetic moment

$$k_T^{u,d} = \int dx \bar{E}_T^{u,d}(x, \xi = 0, t = 0)$$

$$\bar{E}_T = 2\tilde{H}_T + E_T$$

CLAS12 at JLAB

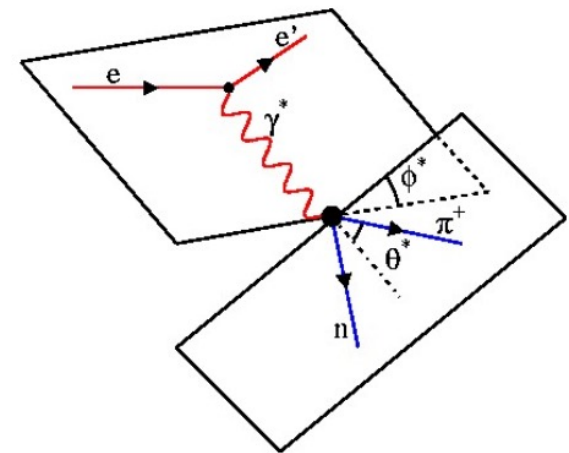
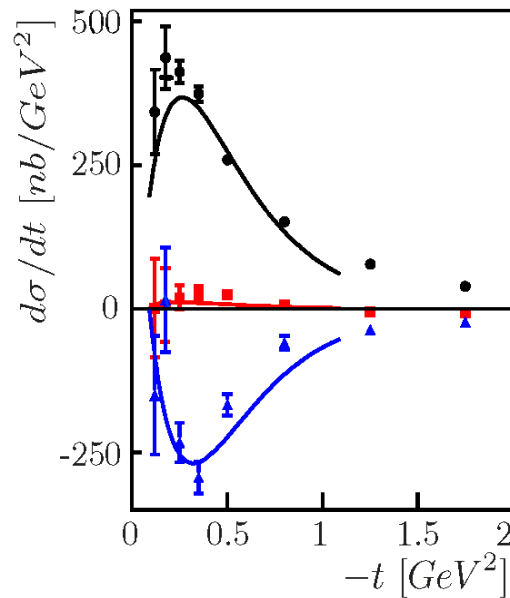
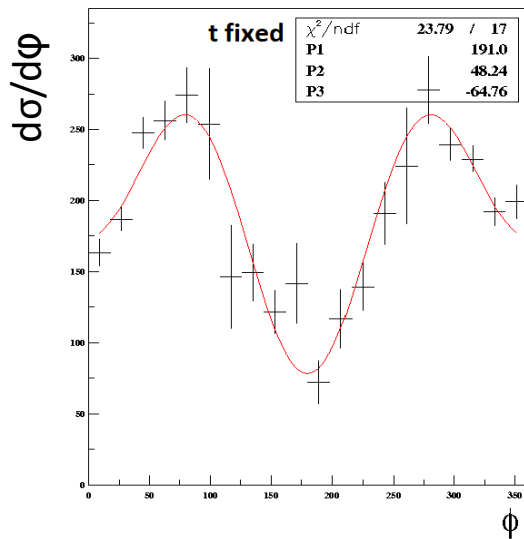


V. Burkert et al., Nucl. Instr. Meth. A 959, 163419 (2020)

- Data recorded with CLAS12 during fall 2018 and spring 2019 (RG-A)
 - 10.6 GeV / 10.2 GeV electron beam ~ 86 % average polarization
 - liquid H₂ target

Differential Cross Section of DVMP (π^0)

$$\frac{d^4\sigma}{dQ^2 dx_B d\phi dt} \propto \boxed{\sigma_T + \epsilon\sigma_L} + \epsilon\boxed{\sigma_{TT}} \cdot \cos(2\phi) + \sqrt{2\epsilon(1+\epsilon)} \cdot \boxed{\sigma_{LT}} \cdot \cos(\phi)$$



CLAS collaboration. I Bedlinskiy et al.
Phys.Rev.Lett. 109 (2012) 112001

$$\sigma_{TT} \propto \frac{t'}{Q^4} |\bar{E}_T|^2$$

$$\sigma_{LT} \propto \frac{\sqrt{-t'}}{Q^4} \xi \sqrt{1-\xi^2} \text{Re} \left[\langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

Unpolarized π^0 Cross Section and Theory Predictions

2 theoretical models:

Goloskokov, Kroll (GK)

Eur. Phys. J. A. 47: 112 (2011)

→ Chiral odd GPDs parameterized using latest results from lattice QCD and transversity parton distribution functions with emphasis on H_T and \bar{E}_T .

Goldstein, Hernandez, Liuti (GGL)

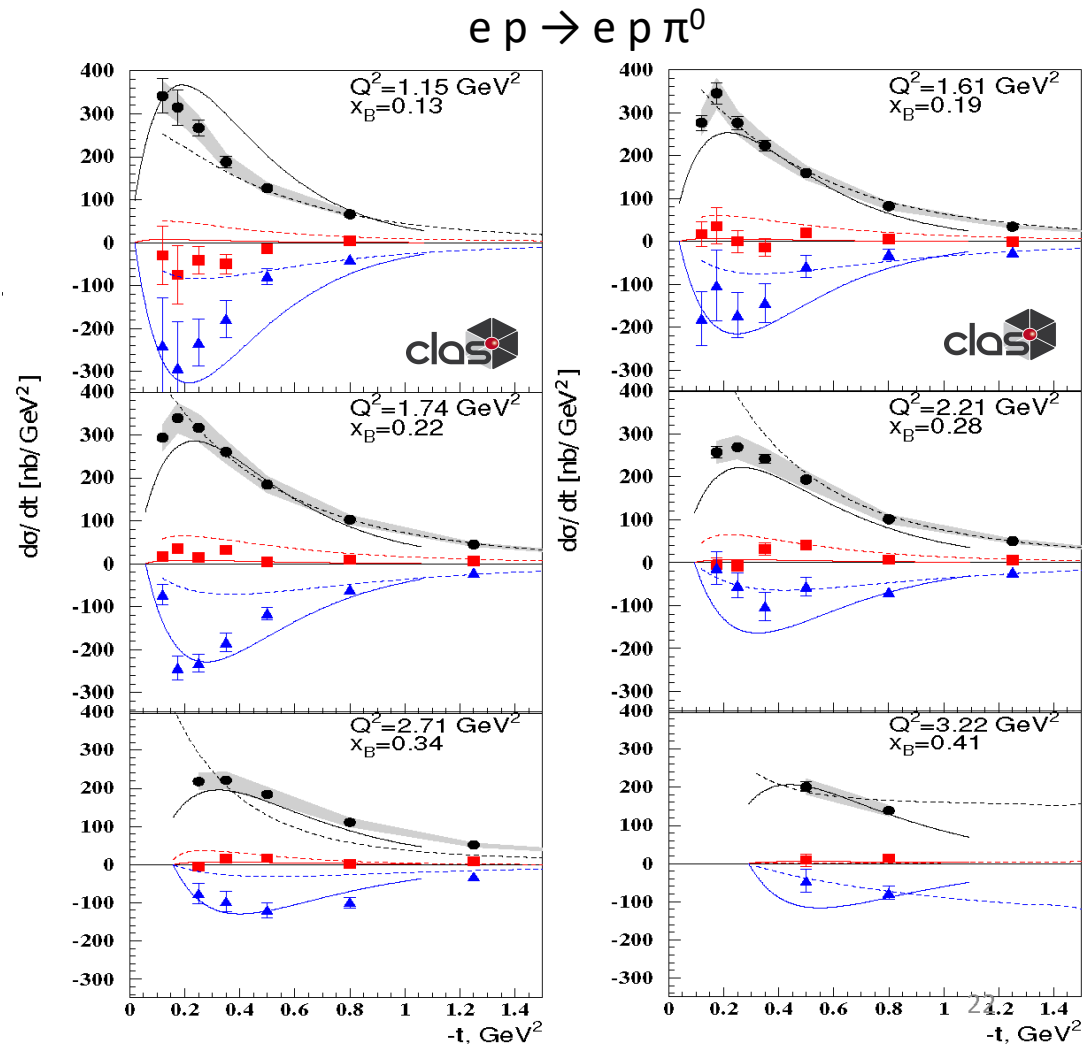
Phys. Rev. D 84, 034007 (2011)

→ Chiral odd GPDs parameterized via linear relations to chiral even GPDs under parity and charge conjugation symmetries in Reggied diquark model

■ σ_0

■ $\sigma_{TT} \rightarrow \cos(2\varphi)$

■ $\sigma_{LT} \rightarrow \cos(\varphi)$



Solid: S. Goloskokov and P. Kroll
Dots: S. Liuti and G. Goldstein

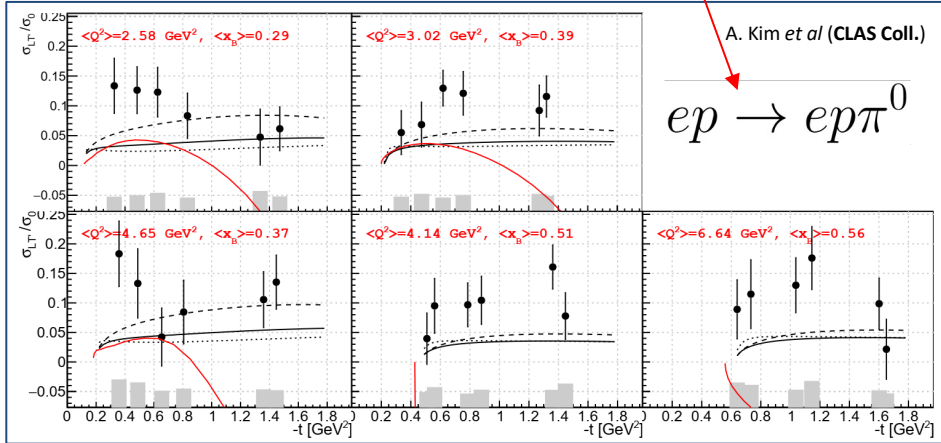
CLAS collaboration. I Bedlinskiy et al.
Phys.Rev.Lett. 109 (2012) 112001

Pseudoscalar meson electroproduction with CLAS12

$$\frac{d^4\sigma}{dQ^2 dx_B d\phi dt} \propto \sigma_T + \epsilon\sigma_L + \epsilon\sigma_{TT} \cdot \cos(2\phi) + \sqrt{2\epsilon(1+\epsilon)} \cdot \sigma_{LT} \cdot \cos(\phi) + P_b \cdot \sqrt{2\epsilon(1-\epsilon)} \cdot \sigma_{LT'} \cdot \sin(\phi)$$

$$\sigma_{LT'} = \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \times \text{Im} \left[\langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \tilde{E}_T \rangle^* \langle \tilde{H} \rangle \right]$$

E=10.6 GeV

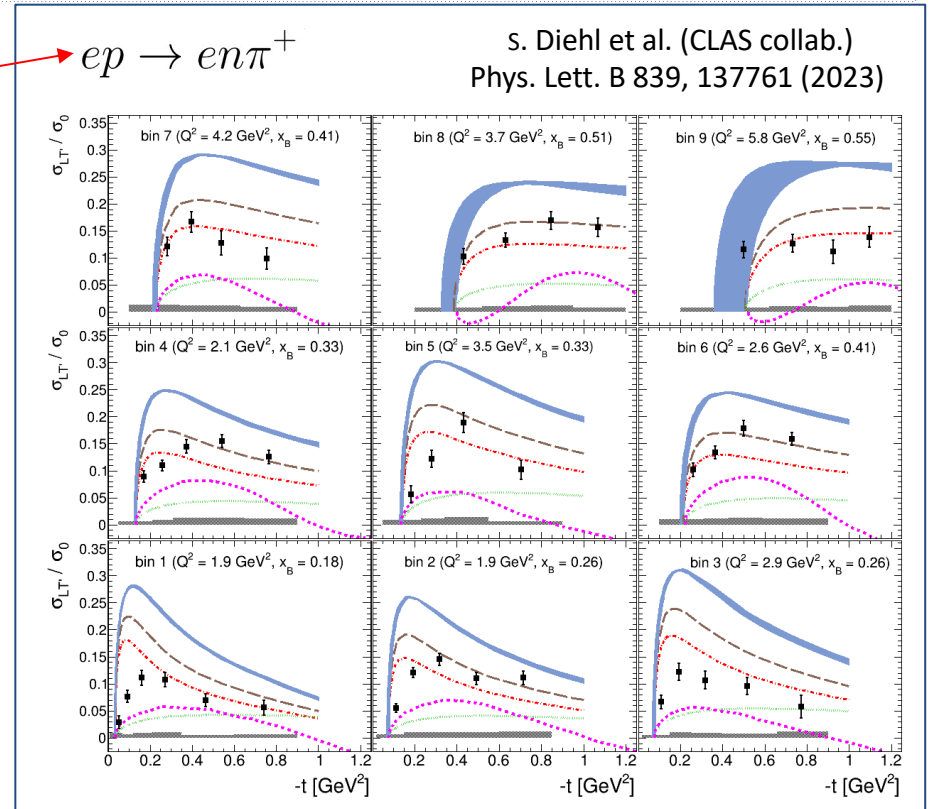


A. Kim et al. (CLAS collab.)
Phys. Lett. B 849, 138459 (2024)

— GK model
..... JML model

\tilde{E}_T is related to the proton's anomalous tensor magnetic moment.

H_T is related to the proton's tensor charge.



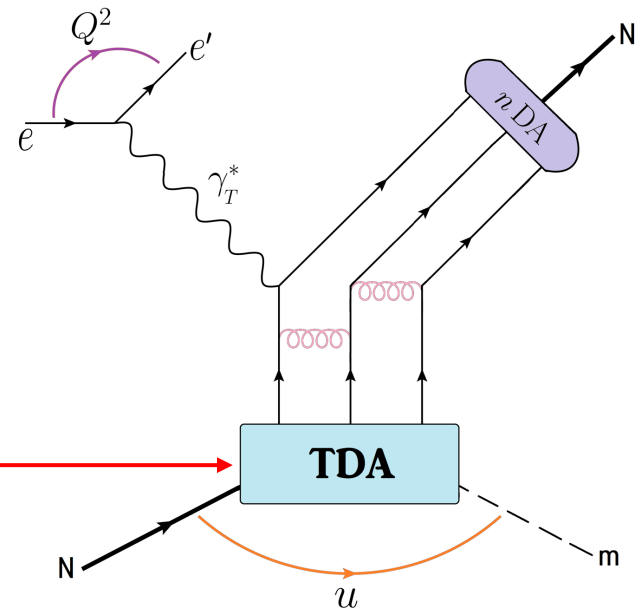
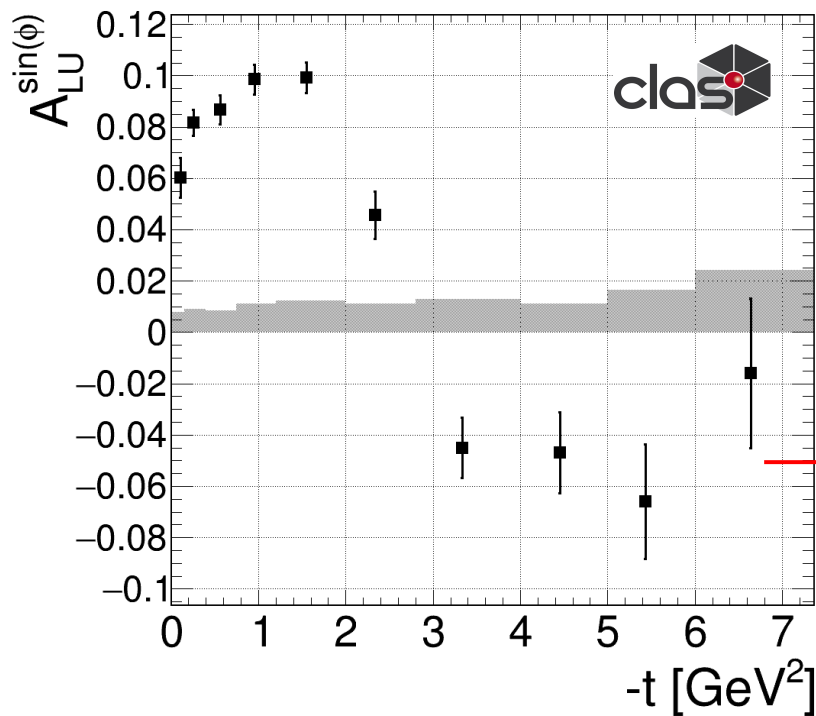
$$\kappa_T^u = \int dx \bar{E}_T^u(x, \xi, t=0) \quad \delta_T^u = \int dx H_T^u(x, \xi, t=0)$$

$$\kappa_T^d = \int dx \bar{E}_T^d(x, \xi, t=0) \quad \delta_T^d = \int dx H_T^d(x, \xi, t=0)$$

From GPDs to Transition Distribution Amplitudes (TDAs)

$$ep \rightarrow en\pi^+ \quad A_{LU}^{\sin\phi} = \sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_T + \epsilon\sigma_L}$$

CLAS data $E = 5.4 \text{ GeV}$
 $W > 2 \text{ GeV}$ $Q^2 > 1 \text{ GeV}^2$



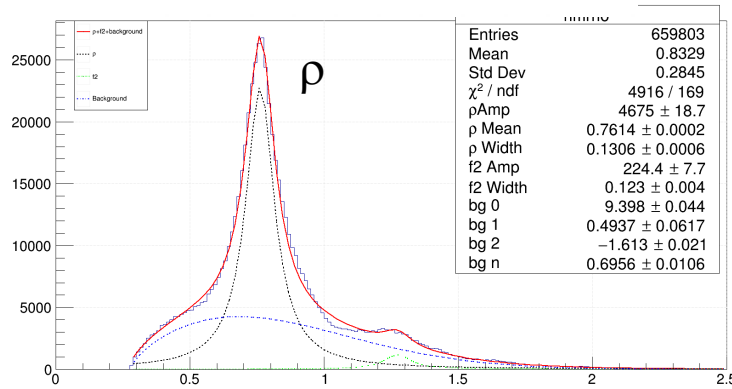
S. Diehl et al. (CLAS collaboration),
 Phys. Rev. Lett. 125, 182001 (2020)

➔ "Backward physics" opens a new window to the 3D nucleon structure!

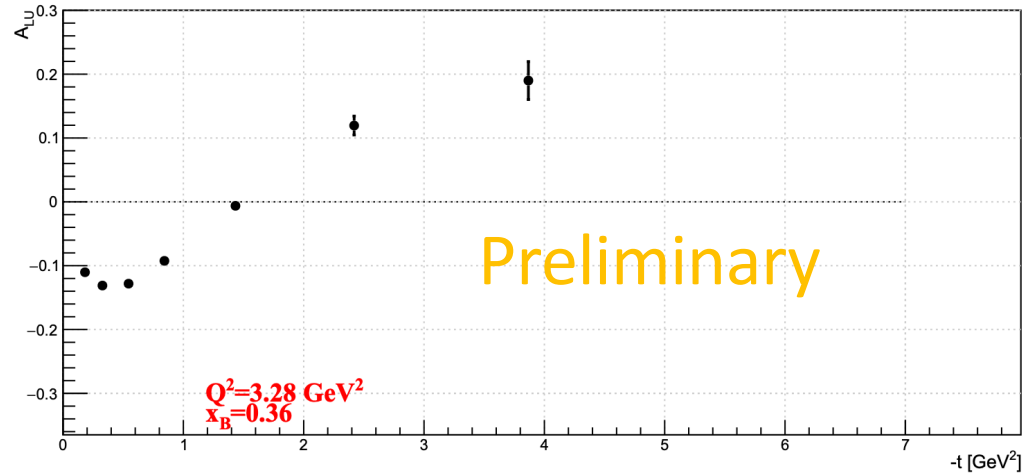
Exclusive ρ/ω production with CLAS12, $ep \rightarrow ep(\rho/\omega)$

$$\sigma_{LT'} \sim r_{00}^8 \sim \text{Im} [\langle H_T \rangle^* \langle E \rangle + \langle \bar{E}_T \rangle^* \langle H \rangle]$$

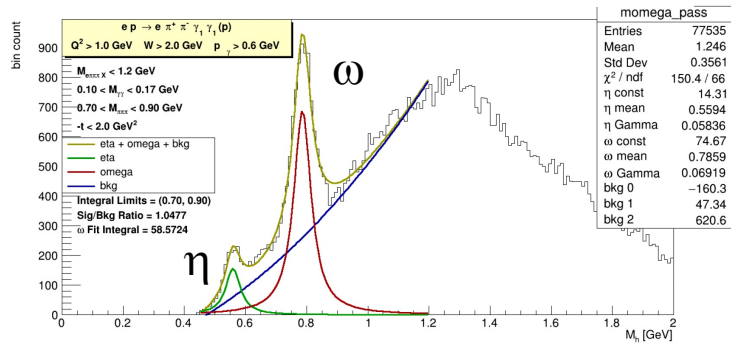
Invariant Mass: $\pi^+ + \pi^-$



$ep \rightarrow ep\rho$

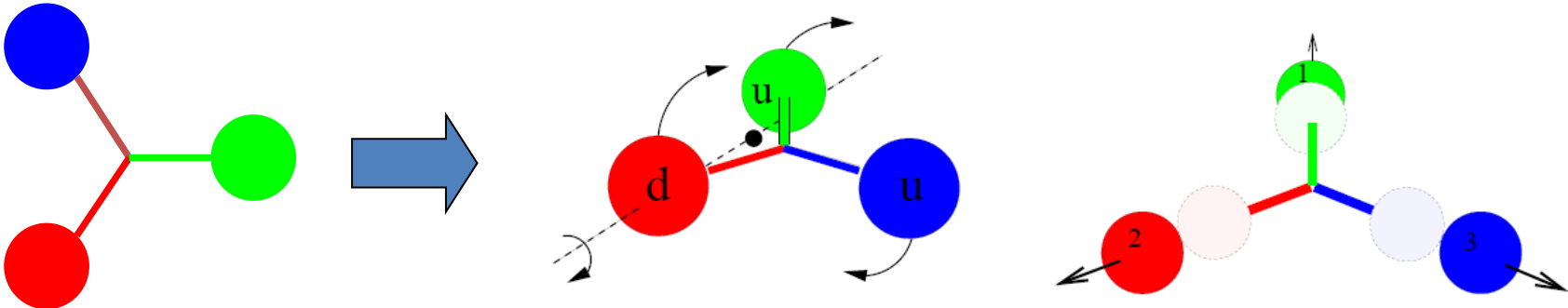


Invariant Mass: $\pi^+ + \pi^- + \pi^0$



N. Trota *et al* (UCONN)

From the ground state nucleon to resonances



How does the excitation affect the 3D structure of the Nucleon?

→ Pressure distributions, tensor charge, ... of resonances?

Traditional way: Study of transition form factors (**2D picture** of transv. position)

3D picture of the excitation process: Encoded in **transition GPDs**

Simplest case: $N \rightarrow \Delta$ transition → **16 transition GPDs**

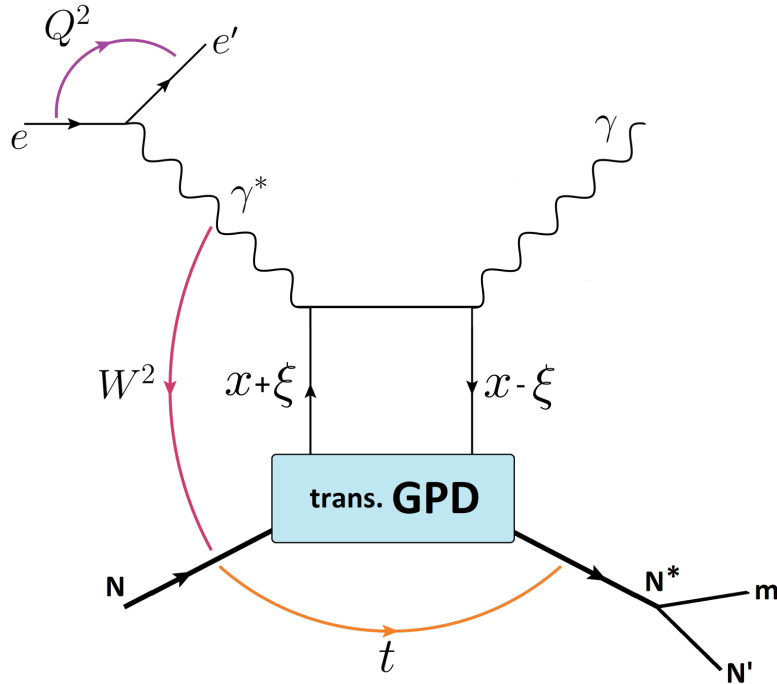
P. Kroll and K. Passek-Kumericki, *Phys. Rev. D* 107, 054009 (2023).

K. Semenov, M. Vanderhaeghen, *arXiv:2303.00119* (2023).

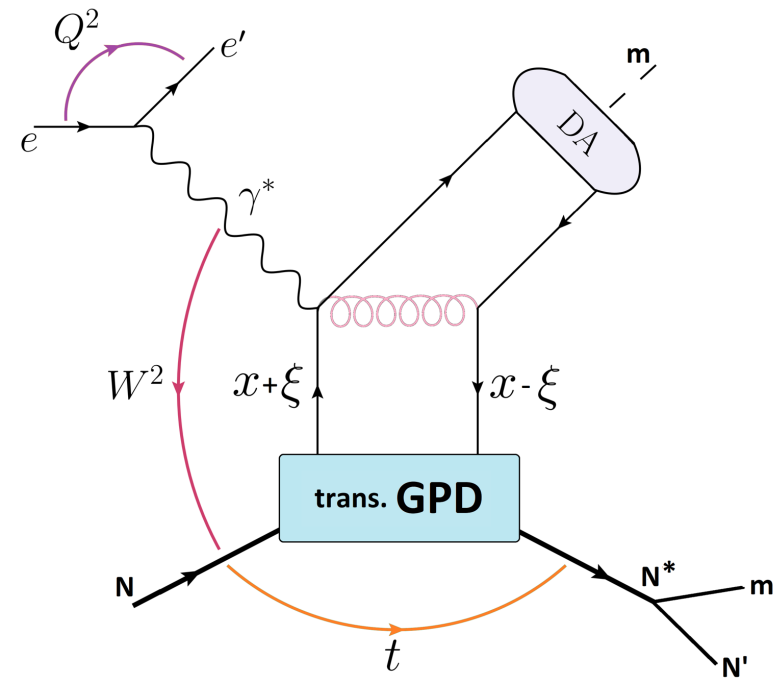
- **8 helicity non-flip transition GPDs (twist 2)**
 - Related to the Jones-Scardon and Adler EM FF for the $N \rightarrow \Delta$ transition
- **8 helicity flip transition GPDs (transversity)**

Non-diagonal DVCS / DVMP

non-diagonal DVCS



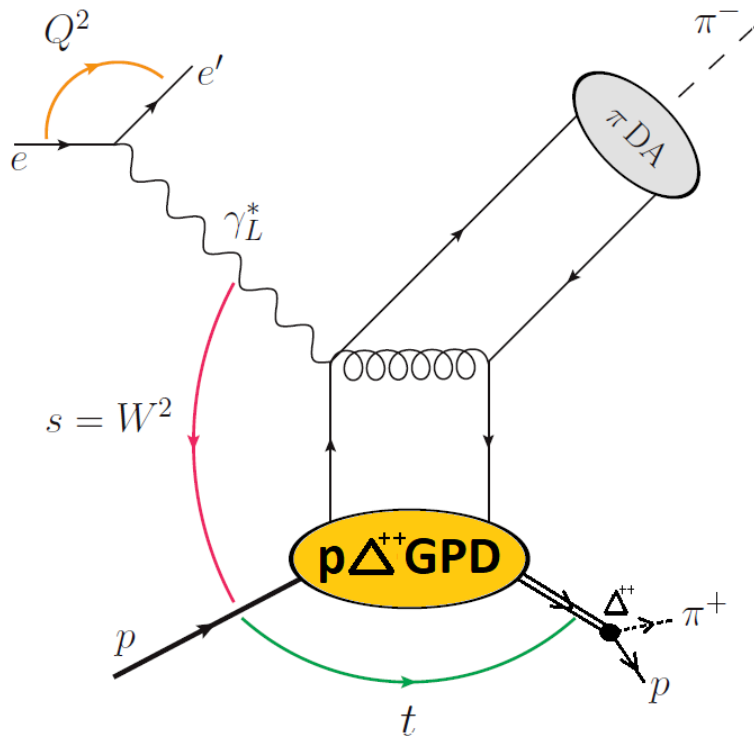
non-diagonal DVMP



factorization expected for: $-t/Q^2$ small, $Q^2 > M_{N^*}^2$ x_B fixed

N- \rightarrow $\Delta(1232)$ transition GPDs: 8 twist-2 GPDs: 4 unpolarized, 4 polarized. [K. Semenov, M. Vanderhaeghen, arXiv:2303.00119 \(2023\)](https://arxiv.org/abs/2303.00119)

$$ep \rightarrow e\Delta^{++}\pi^{-} \rightarrow ep\pi^{+}\pi^{-}$$



Factorization expected for:

$$-t / Q^2 \ll 1, x_B \text{ fixed, and } Q^2 > M_{\Delta}^2$$

- Provides access to p- Δ transition GPDs

$$ep \rightarrow e\Delta^{++}\pi^{-} \rightarrow ep\pi^{+}\pi^{-}$$

$$I_z = +3/2$$

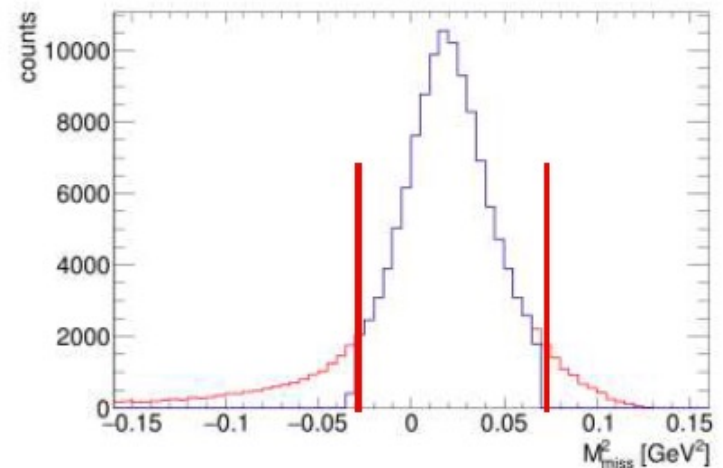
- The $p\pi^+$ final state can **only** be populated by **Δ -resonances** -> Large gap between $\Delta(1232)$ and higher resonances

Event Selection and Kinematic Cuts

Event selection: $ep \rightarrow ep\pi^- X$

$$X = \pi^+$$

→ 2 sigma cut around the missing π^+



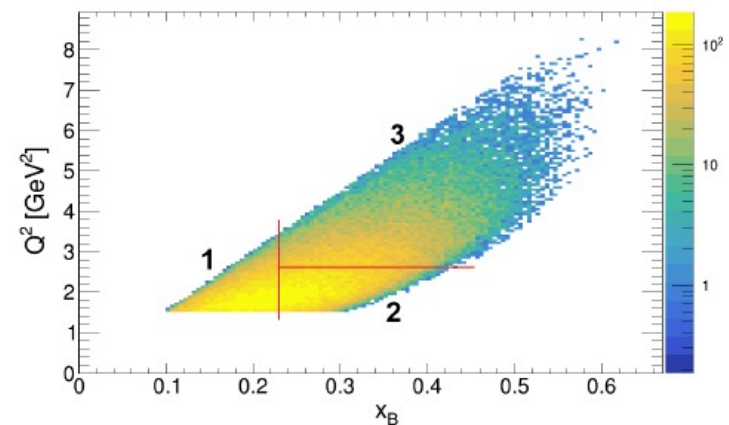
Kinematic cuts:

$$Q^2 > 1.5 \text{ GeV}^2$$

$$W > 2 \text{ GeV}$$

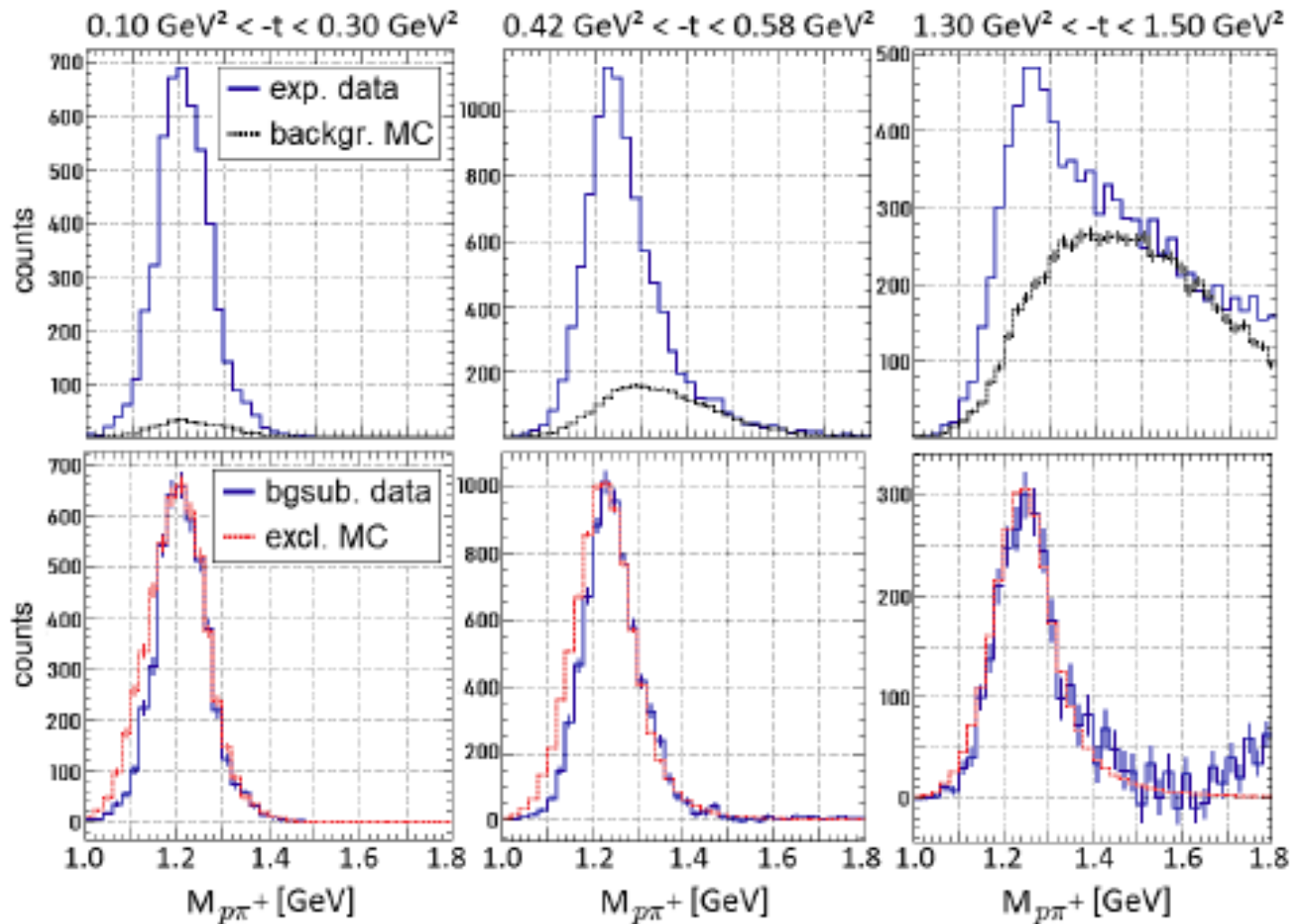
$$y < 0.75$$

$$-t < 1.5 \text{ GeV}^2$$

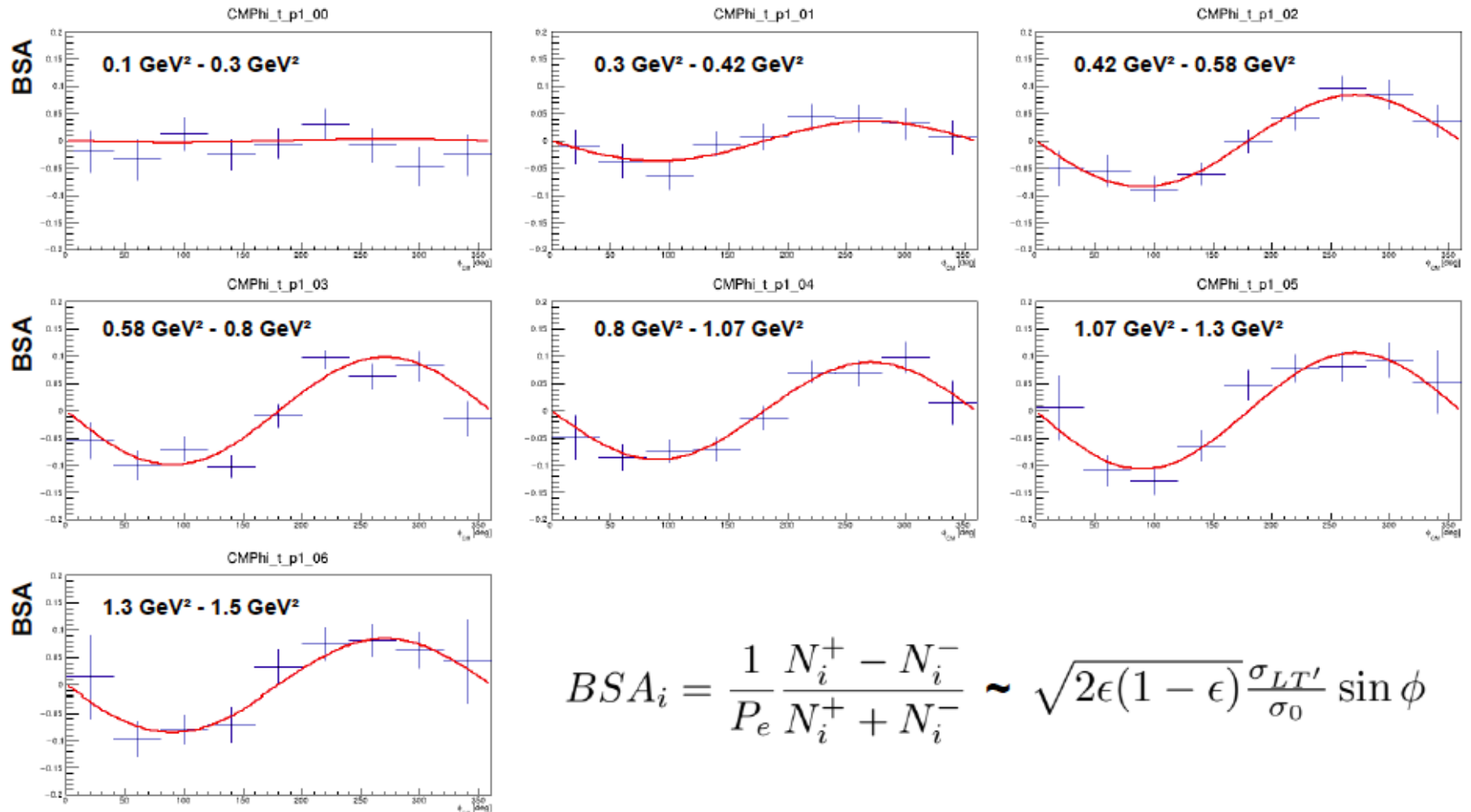


Signal and Background Separation

$$ep \rightarrow e\Delta^{++}\pi^{-} \rightarrow ep\pi^{+}\pi^{-}$$



Resulting Beam Spin Asymmetries (Q^2 - x_B integrated)



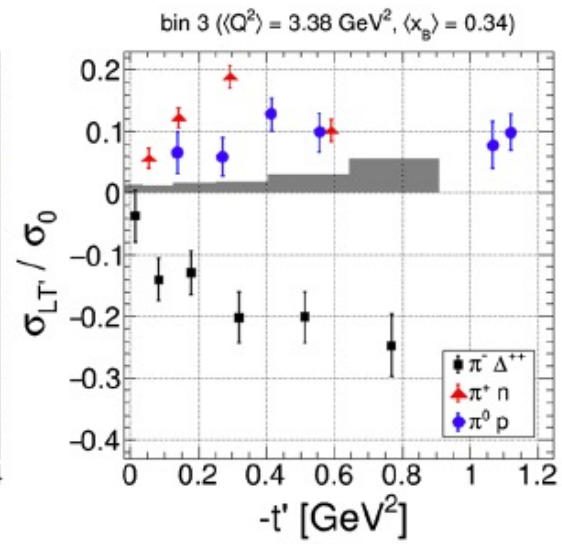
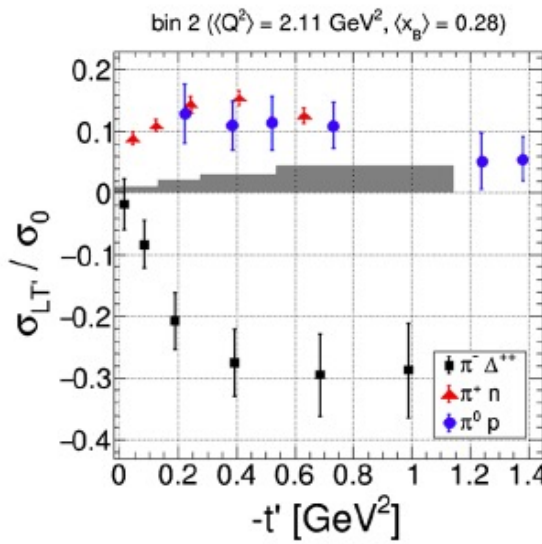
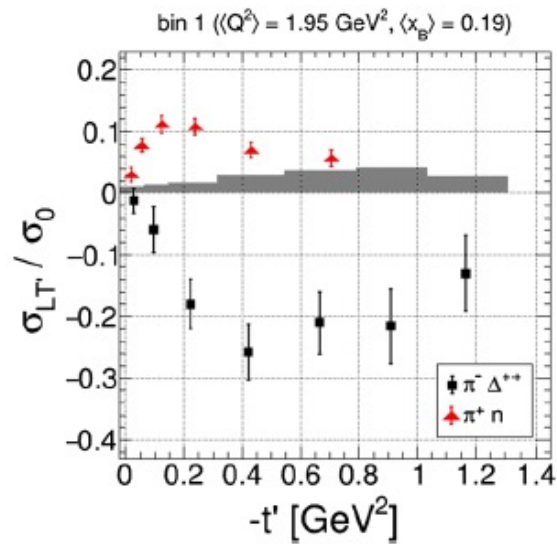
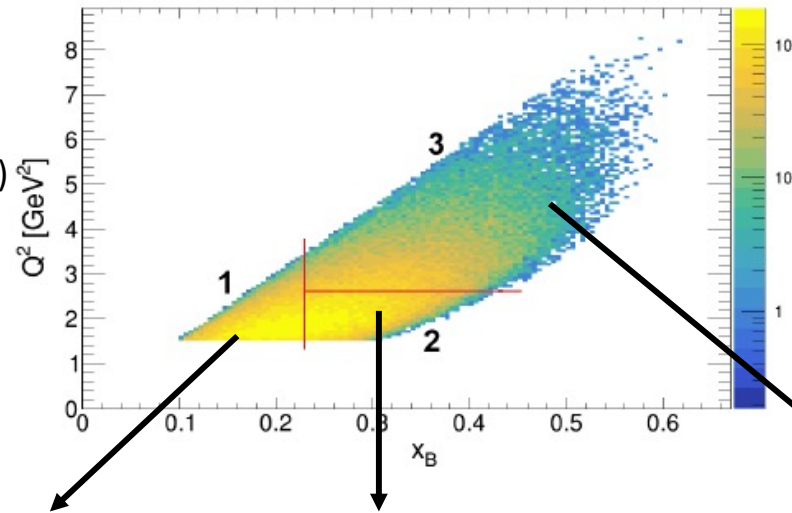
$$BSA_i = \frac{1}{P_e} \frac{N_i^+ - N_i^-}{N_i^+ + N_i^-} \sim \sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_0} \sin \phi$$

Results

S. Diehl et al. (CLAS collab.),
Phys. Rev. Lett. 131, 021901 (2023)

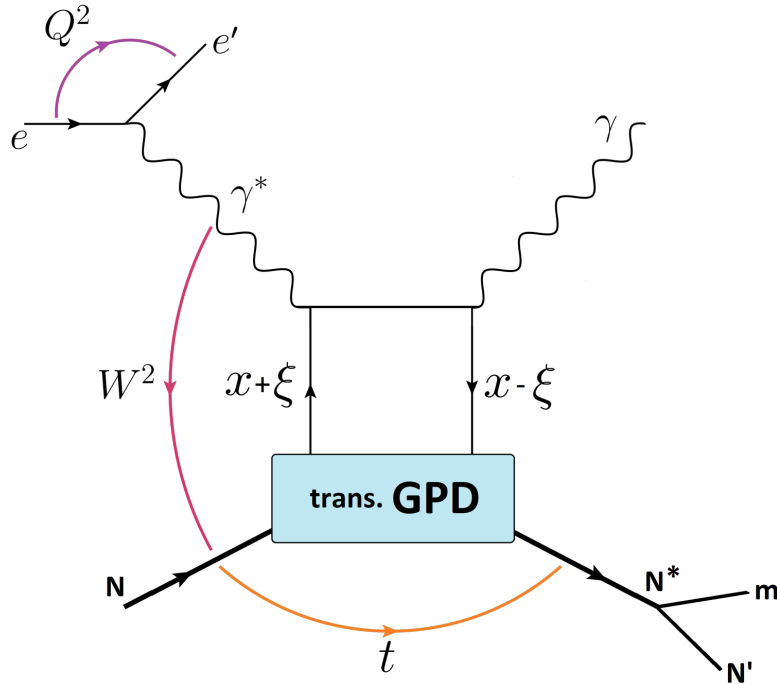
S. Diehl et al. (CLAS collab.)
Phys. Lett. B 839, 137761 (2023)

A. Kim et al. (CLAS collab.)
Phys. Lett. B 849, 138459 (2024)

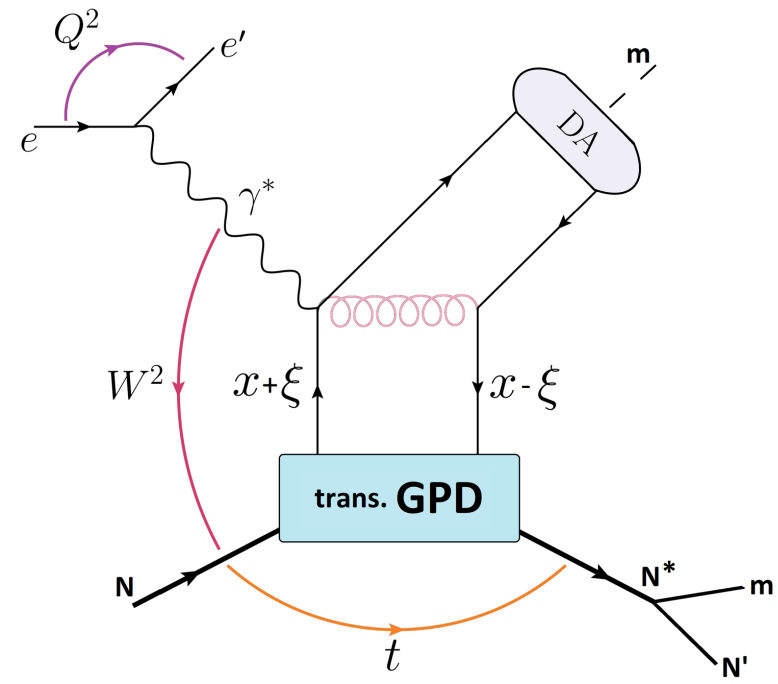


Non-diagonal DVCS / DVMP

non-diagonal DVCS



non-diagonal DVMP



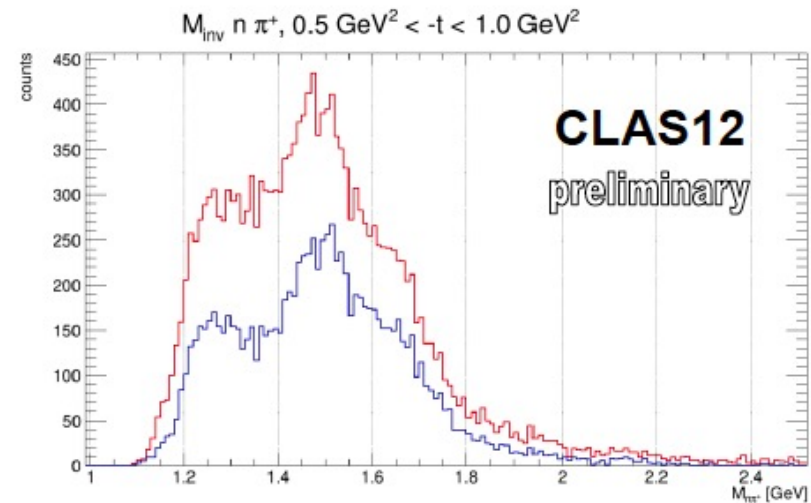
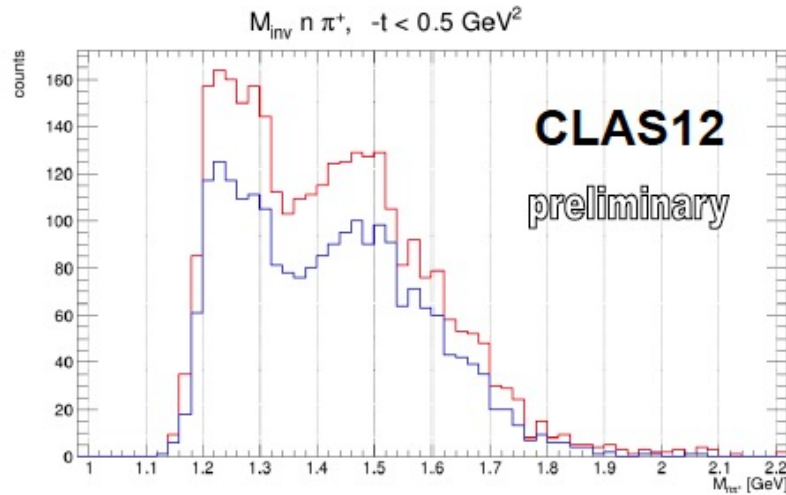
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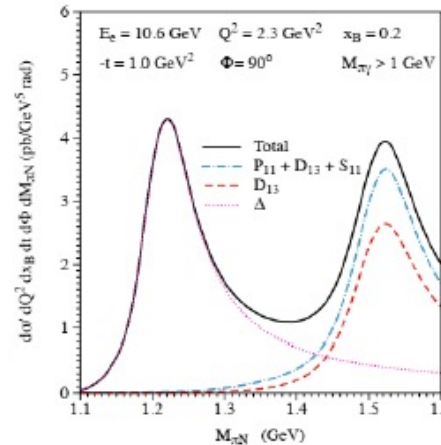
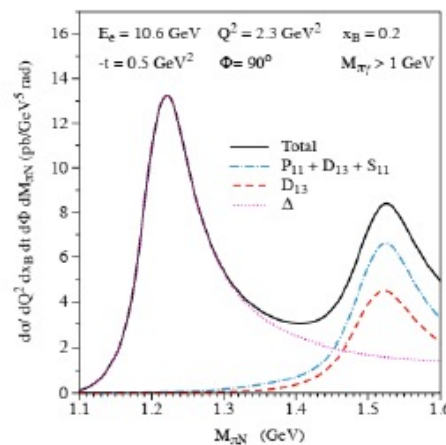
$e p \rightarrow e' \Delta^+ \gamma \rightarrow e' n \pi^+ \gamma$

$e p \rightarrow e' \Delta^+ \gamma \rightarrow e' n \pi^+ \gamma$

— raw — $M(\pi^+ \gamma) > 1.0 \text{ GeV}$



Semenov-Tian-Shansky,
Vanderhaeghen,
arXiv:2303.00119 (2023)

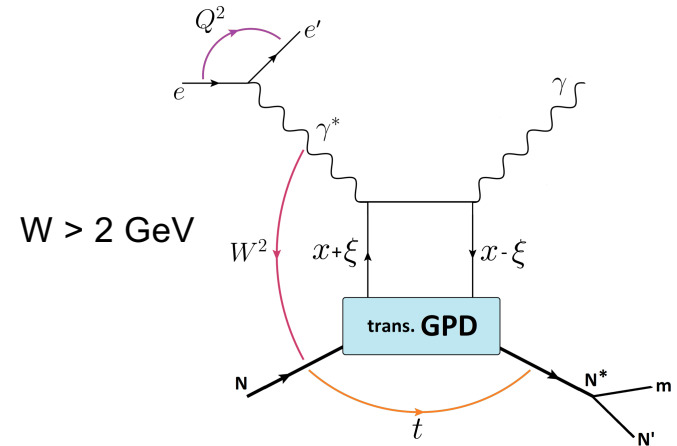
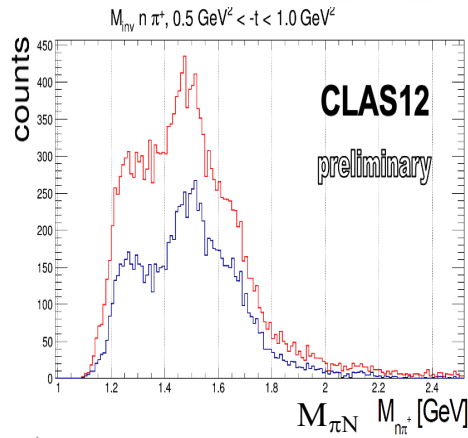
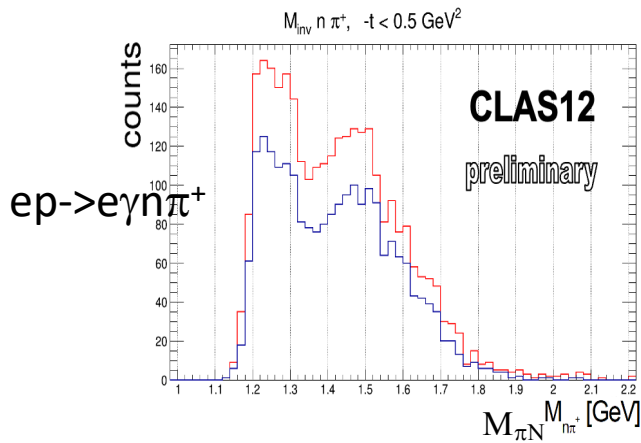
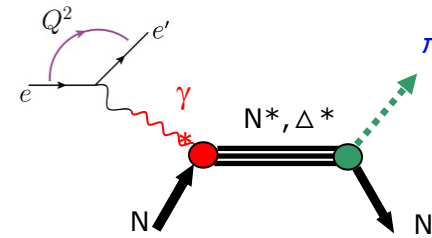
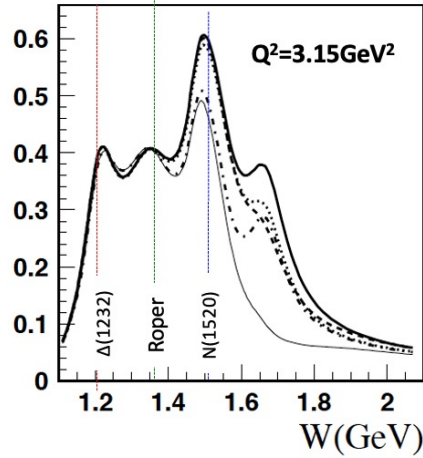
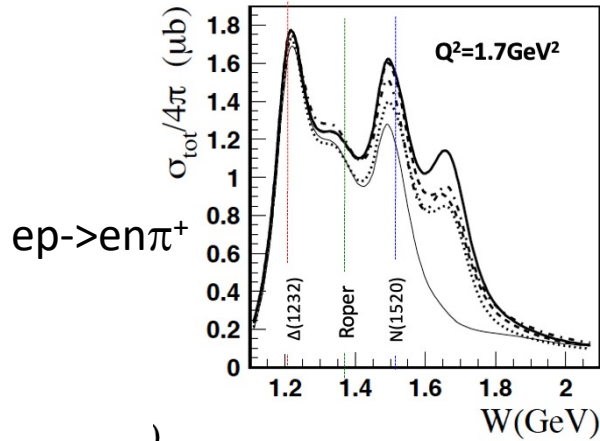


- study in progress
- π^+ mainly in CD and low momentum
- Good agreement of BSA
- awaiting pass 2

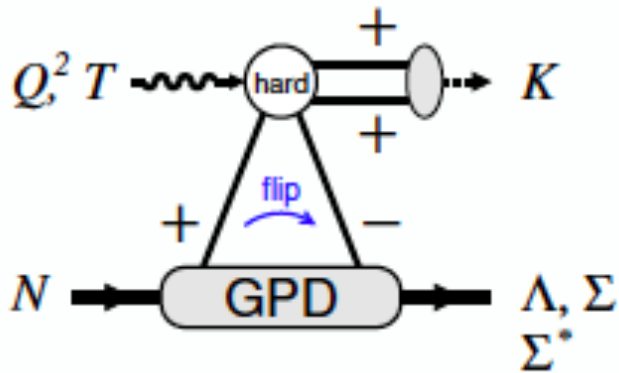
$ep \rightarrow en\pi^+$

vs.

$ep \rightarrow e\gamma n\pi^+$



N → Λ, Σ, Σ* GPDs in K production with CLAS12



Production mechanism

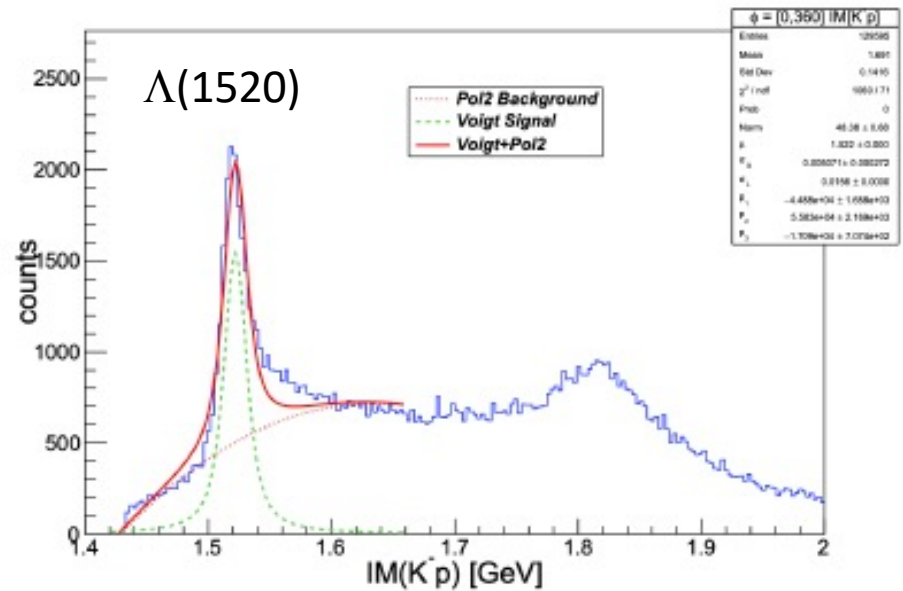
Same twist-3 mechanism with chiral-odd structures as π, η production

Symmetry relations for strange chiral-odd GPDs

$N \rightarrow \Lambda, \Sigma$ related to $N \rightarrow N$
by conventional SU(3) flavor symmetry

$N \rightarrow \Sigma^*$ related to $N \rightarrow N, \Lambda, \Sigma$
by SU(6) spin-flavor symmetry in large- N_c limit

Predictive power; quantitative predictions possible



Invariant mass distribution of pK^-
after $ep \rightarrow e'p'K^+K^-$ events are selected.

U. Shrestha (UConn)

Electron Scattering Binning Scheme

	Resonance Region	DIS Region
Inclusive Scattering	Q^2, W	Q^2, x_B
Exclusive Process ($\gamma, \pi, \rho, \phi, \dots$)	$Q^2, W, \cos\theta^*, \phi$	$Q^2, x_B, -t, \phi$
Off-diagonal DVCS or DVMP	$Q^2, x_B, -t, \phi, M_{\pi N}, \cos\theta^*, \phi^*$	

Summary

- GPDs provide a unifying framework to study the 3-D quark and gluon structure of the nucleon
- 3-D imaging of nucleons uncovers the rich dynamics of QCD.
- CLAS12 allows high precision measurements of Exclusive channels to study ground nucleon GPDs and transition GPDs with large kinematic coverages in the valence quark regime!
- The nucleon-to-resonance ($N \rightarrow N^*$) transition GPDs may provide a unique tool for exploring the 3D structure of baryon resonances.