



Theory and Phenomenology of GPDs - Overview

(Can't be comprehensive in 30 minutes)

- ❑ Inclusive vs. Exclusive – Explore Hadron's Partonic Structure without Breaking it!
- ❑ GPDs, 3D Tomographic Images & Beyond
- ❑ Exclusive Processes for Extracting GPDs
- ❑ QCD Factorization, Angular Modulations, ...
- ❑ Summary and Outlook

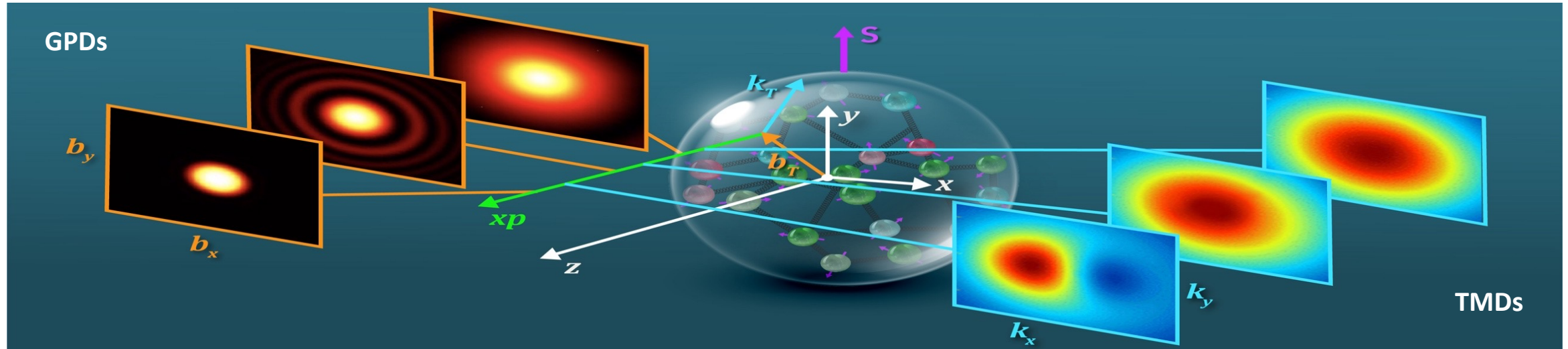


See talks by Nicole D'Hose, Charlotte Van Hulse, Valerio Bertone, Cedric Mezrag, Silvia Niccolai, ...
See also talk by Herve Dutrioux @ QCD Evolution

“See” Internal Structure of Hadron without seeing quarks/gluons directly?

□ 3D hadron structure:

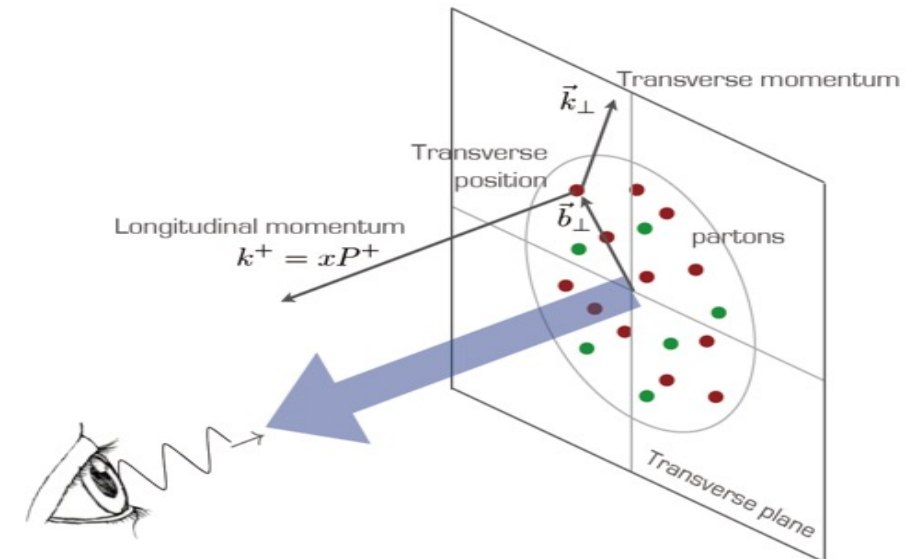
NO quarks and gluons can be seen in isolation!



□ Need new observables with two distinctive scales:

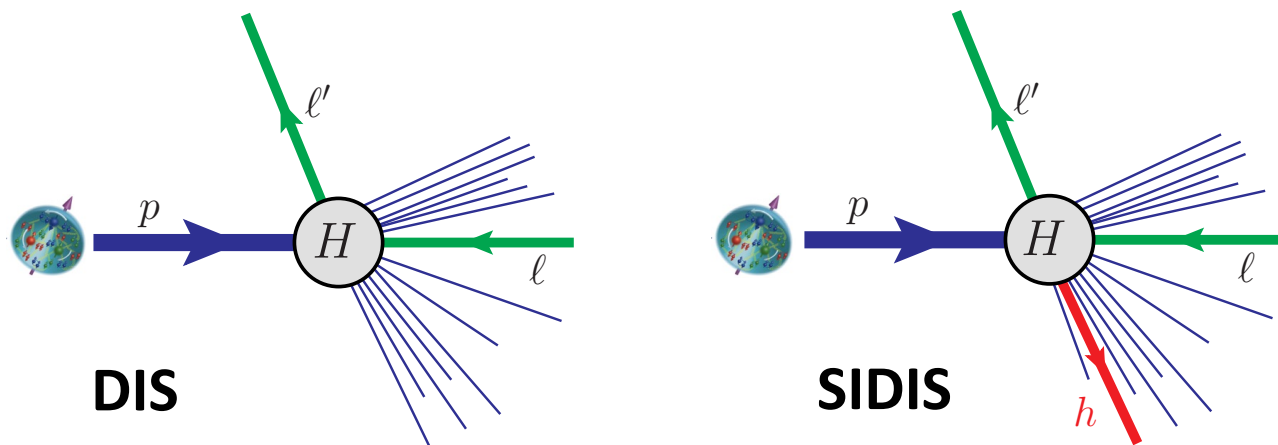
$$Q_1 \gg Q_2 \sim 1/R \sim \Lambda_{\text{QCD}}$$

- **Hard scale:** Q_1 to localize the probe to see the particle nature of quarks/gluons
- **“Soft” scale:** Q_2 to be more sensitive to the emergent regime of hadron structure $\sim 1/\text{fm}$

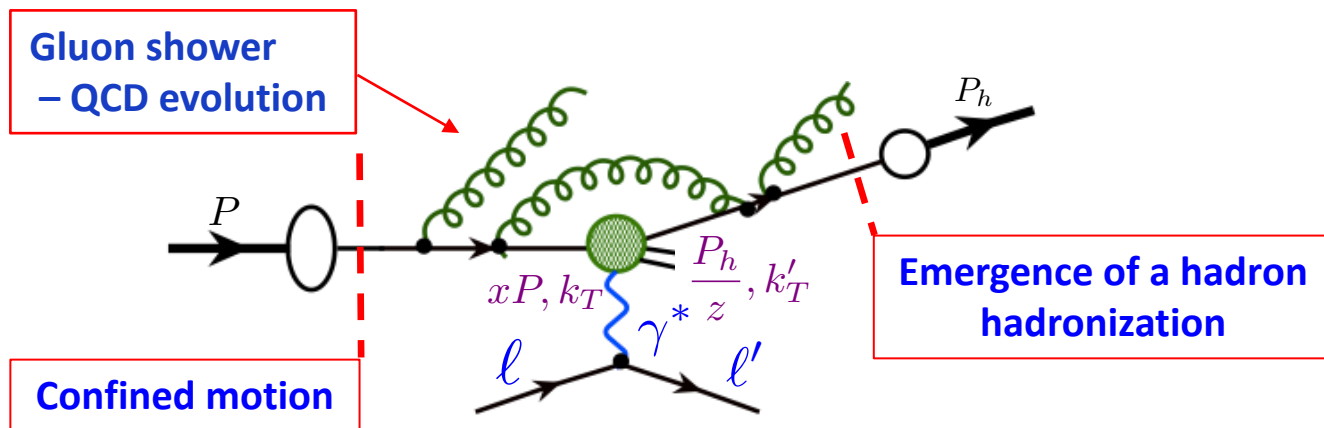
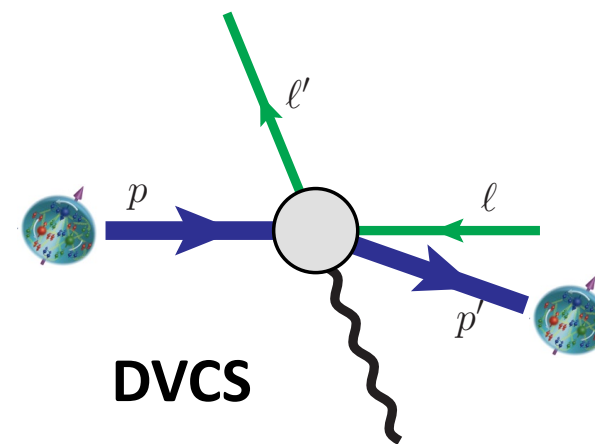


Inclusive vs. Exclusive – Partonic structure without breaking the hadron!

Inclusive scattering



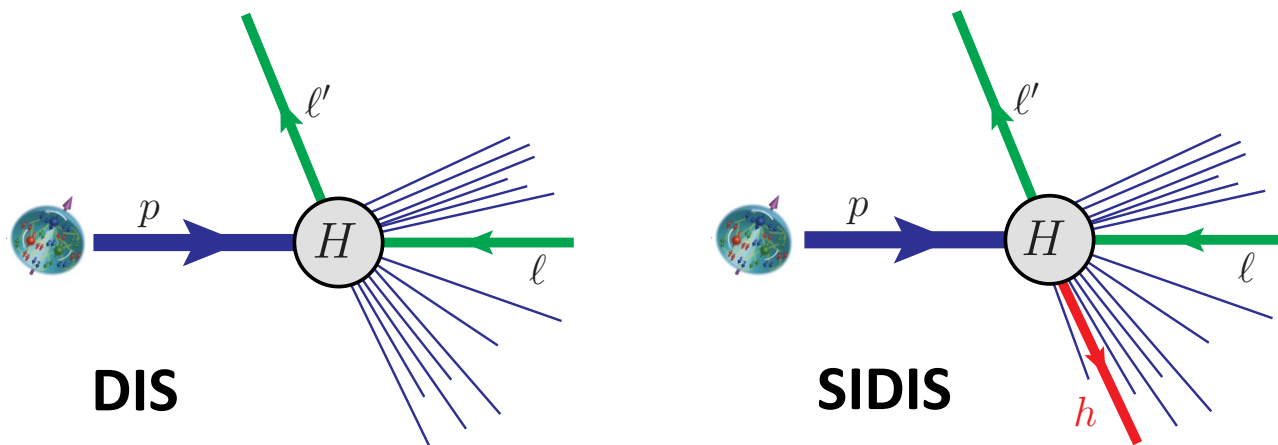
Exclusive diffraction



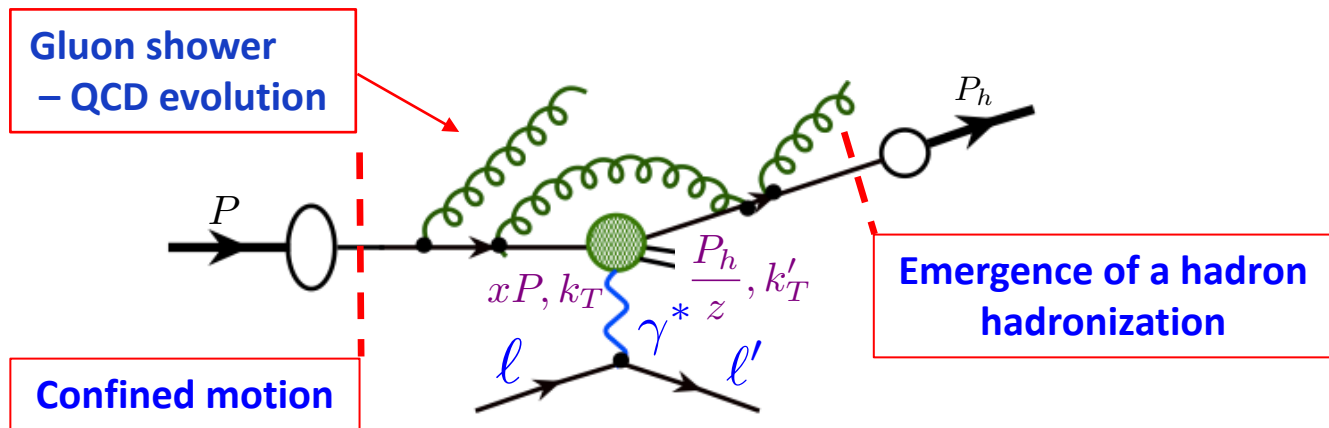
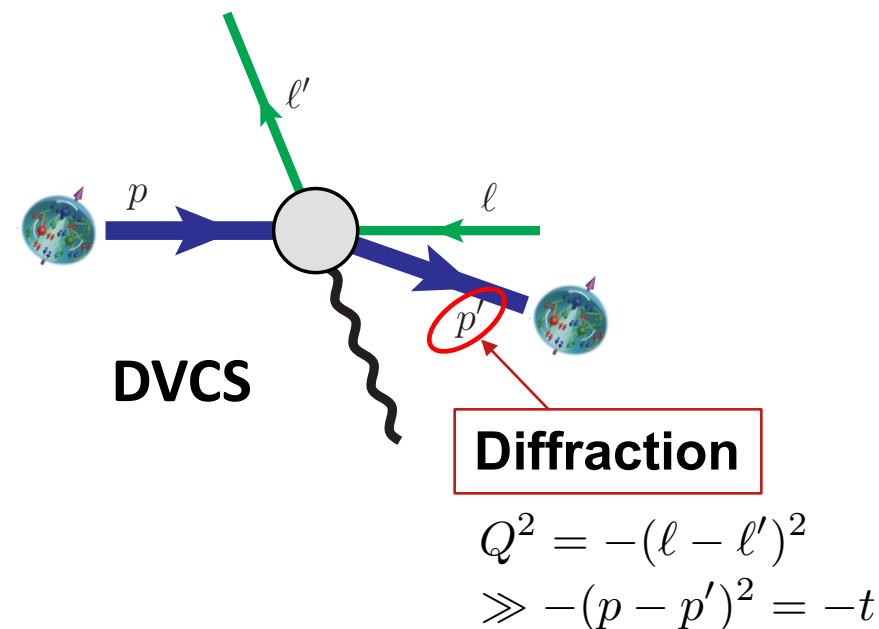
$$e + P \rightarrow e + h + X$$

Inclusive vs. Exclusive – Partonic structure without breaking the hadron!

Inclusive scattering



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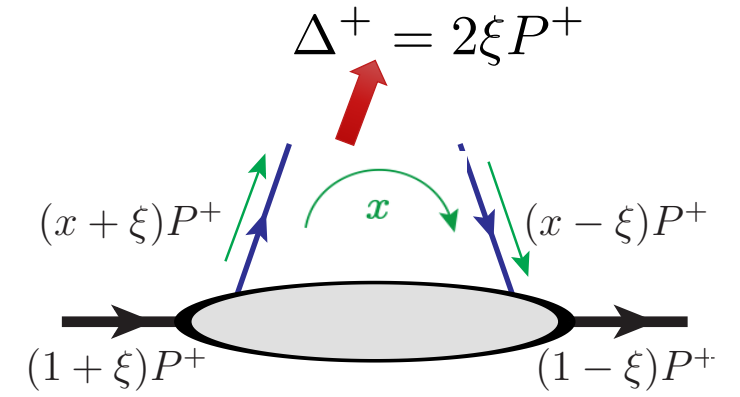
$$e + P \rightarrow e + h + X$$

Generalized Parton Distributions (GPDs)

D. Müller, D. Robaschik, B. Geyer, F.-M. Dittes, J. Hořejši,
Fortsch. Phys. 42 (1994) 101

Definition:

$$\begin{aligned}
 F^q(x, \xi, t) &= \int \frac{dz^-}{4\pi} e^{-ixP^+z^-} \langle p' | \bar{q}(z^-/2) \gamma^+ q(-z^-/2) | p \rangle \\
 &= \frac{1}{2P^+} \left[H^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) - E^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2m} u(p) \right], \\
 \tilde{F}^q(x, \xi, t) &= \int \frac{dz^-}{4\pi} e^{-ixP^+z^-} \langle p' | \bar{q}(z^-/2) \gamma^+ \gamma_5 q(-z^-/2) | p \rangle \\
 &= \frac{1}{2P^+} \left[\tilde{H}^q(x, \xi, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) - \tilde{E}^q(x, \xi, t) \bar{u}(p') \frac{\gamma_5 \Delta^+}{2m} u(p) \right].
 \end{aligned}$$

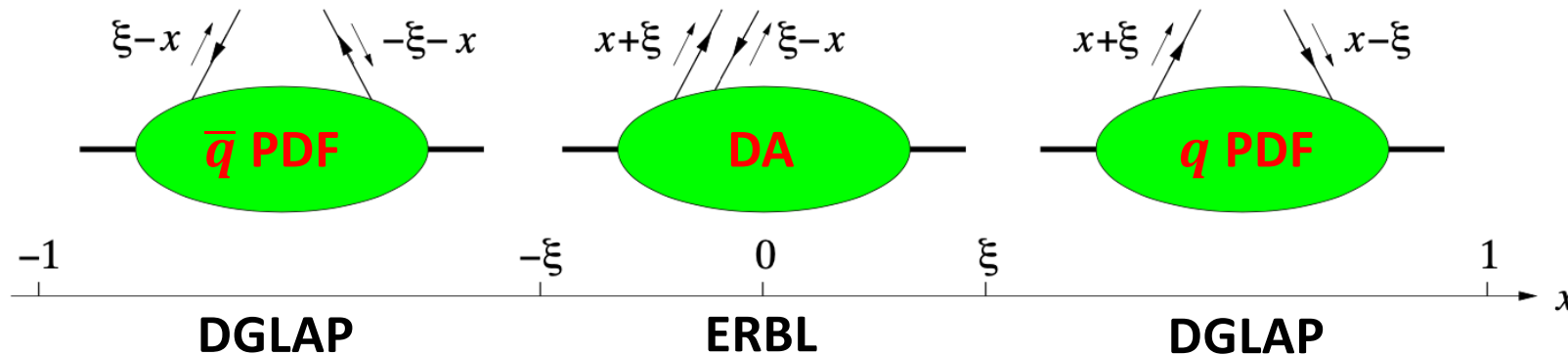


Combine PDF and Distribution Amplitude (DA):

Forward limit $\xi = t = 0$: $H^q(x, 0, 0) = q(x)$, $\tilde{H}^q(x, 0, 0) = \Delta q(x)$

$$\begin{aligned}
 P^+ &= \frac{p^+ + p'^+}{2} \\
 \Delta &= p - p' \quad t = \Delta^2
 \end{aligned}$$

Similar definition
 for gluon GPDs

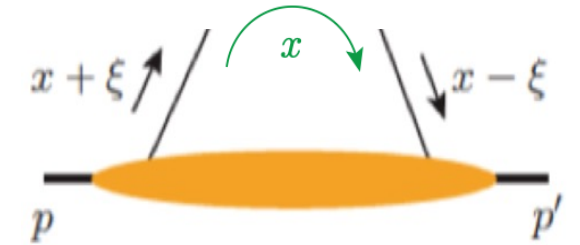


Properties of GPDs – Partonic

□ Impact parameter dependent parton density distribution:

$$q(x, b_{\perp}, Q) = \int d^2 \Delta_{\perp} e^{-i \Delta_{\perp} \cdot b_{\perp}} H_q(x, \xi = 0, t = -\Delta_{\perp}^2, Q)$$

➔ Quark density in $dx d^2 b_T$



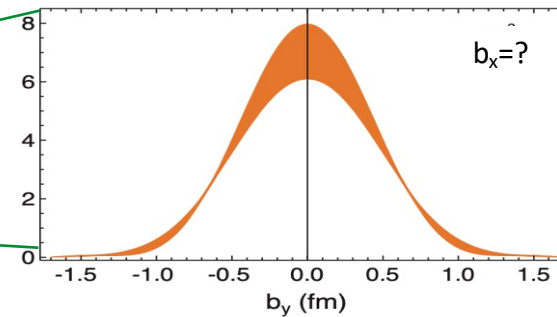
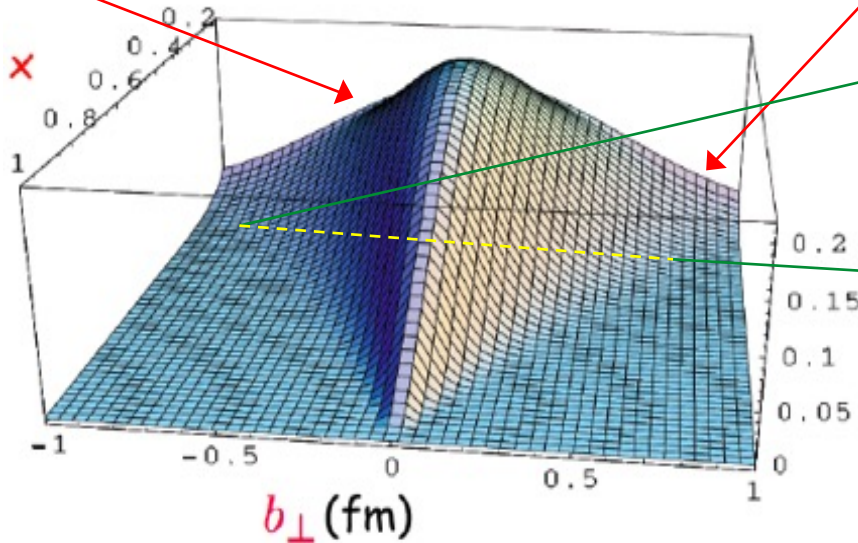
Measurement of p' fixes (t, ξ)
 x = momentum flow between the pair

How fast does glue density fall?

Tomographic image of hadron in slice of x

How far does glue density spread?

➔



Slice in (x, Q)

Modeled by M. Burkardt, PRD 2000

$$\langle q_{\perp}^N \rangle \equiv \int db_{\perp} b_{\perp}^N q(x, b_{\perp}, Q)$$

➔ Proton radii from quark and gluon spatial density distribution, $r_q(x)$ & $r_g(x)$

Properties of GPDs – Partonic

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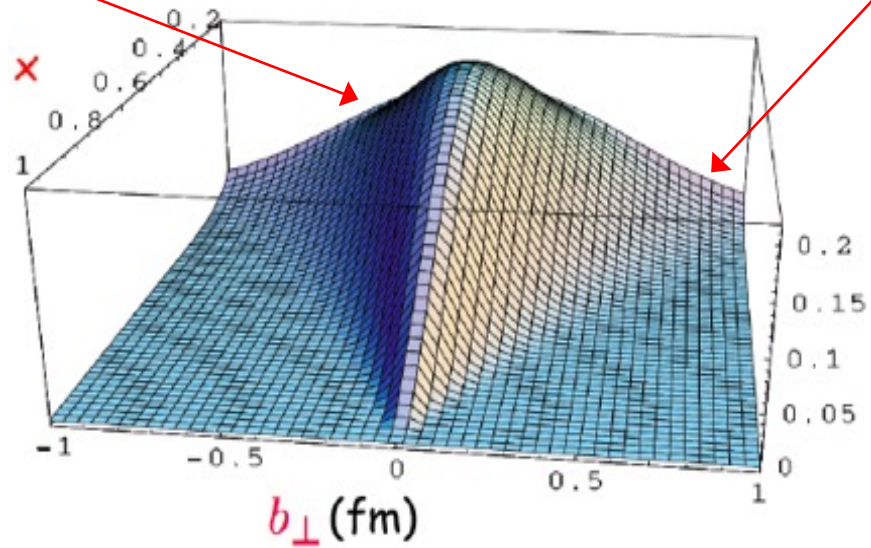
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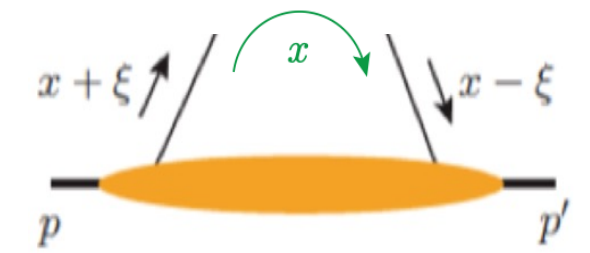
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➔

➔ Proton radii from quark and gluon spatial density distribution, $r_q(x)$ & $r_g(x)$



Measurement of p' fixes (t, ξ)
 x = momentum flow between the pair

- Should $r_q(x) > r_g(x)$, or vice versa?
- Could $r_g(x)$ saturates as $x \rightarrow 0$
- How do they compare with known radius (EM charge radius, mass radius, ...), & why?
- How the image correlate to hadron spin, ... ?
- ...

Properties of GPDs – Hadronic = Moments of GPDs

QCD energy-momentum tensor:

Ji, PRL78, 1997
V. D. Burkert, et al. RMP 95 (2023) 041002

$$T^{\mu\nu} = \sum_{i=q,g} T_i^{\mu\nu} \quad \text{with} \quad T_q^{\mu\nu} = \bar{\psi}_q i\gamma^{(\mu} \overleftrightarrow{D}^{\nu)} \psi_q - g^{\mu\nu} \bar{\psi}_q \left(i\gamma \cdot \overleftrightarrow{D} - m_q \right) \psi_q \quad \text{and} \quad T_g^{\mu\nu} = F^{a,\mu\eta} F^{a,\eta\nu} + \frac{1}{4} g^{\mu\nu} (F_{\rho\eta}^a)^2$$

“Gravitational” form factors:

$$\langle p' | T_i^{\mu\nu} | p \rangle = \bar{u}(p') \left[A_i(t) \frac{P^\mu P^\nu}{m} + J_i(t) \frac{iP^{(\mu} \sigma^{\nu)\Delta}}{2m} + D_i(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{4m} + m \bar{c}_i(t) g^{\mu\nu} \right] u(p)$$

Connection to GPD moments:

$$\int_{-1}^1 dx x F_i(x, \xi, t) \propto \langle p' | T_i^{++} | p \rangle \propto \bar{u}(p') \left[\underbrace{(A_i + \xi^2 D_i)}_{\int_{-1}^1 dx x H_i(x, \xi, t)} \gamma^+ + \underbrace{(B_i - \xi^2 D_i)}_{\int_{-1}^1 dx x E_i(x, \xi, t)} \frac{i\sigma^{+\Delta}}{2m} \right] u(p)$$

$$C_i(t) \leftrightarrow D_i(t)/4$$

**Related to pressure
& stress force inside h**

Polyakov, Schweitzer,
Inntt. J. Mod. Phys.
A33, 1830025 (2018)
Burkert, Elouadrhiri, Girod
Nature 557, 396 (2018)

Angular momentum sum rule:

$$J_i = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H_i(x, \xi, t) + E_i(x, \xi, t)]$$

$i = q, g$

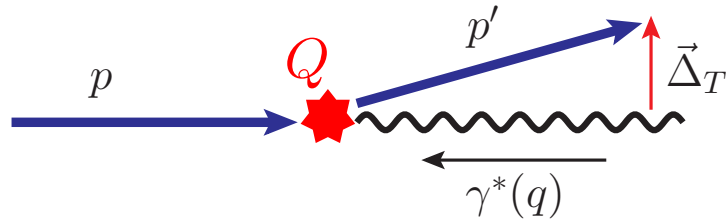
3D tomography
Relation to GFFs
Angular Momentum

**x-dependence
of GPDs!**

Need to know the x-dependence of GPDs to construct the proper moments!

Known Physical Processes for Extracting GPDs

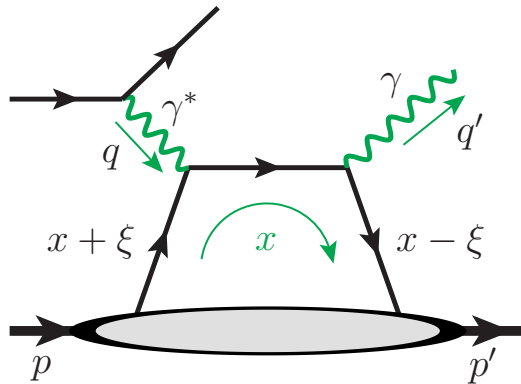
- Hit the proton hard without breaking it \Rightarrow Diffractive scattering to keep proton intact



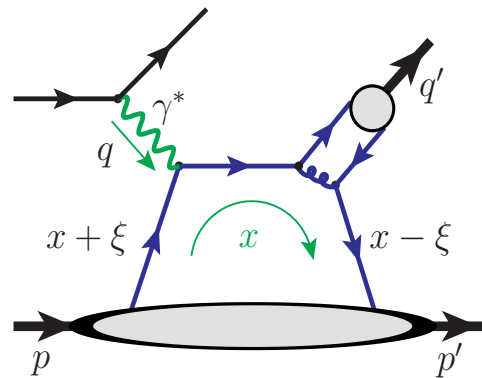
HERA discovery:

\sim 10-15% of HERA events with the Proton stayed intact

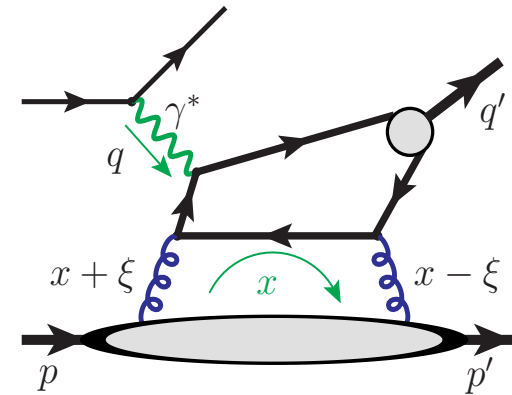
- Known exclusive processes for extracting GPDs:



DVCS: $Q^2 \gg |t|$



DVMP



DVQP

+ DDVCS, ...

Feature: Two-scale observables

$$Q^2 \gg |t| \quad t = (p - p')^2$$

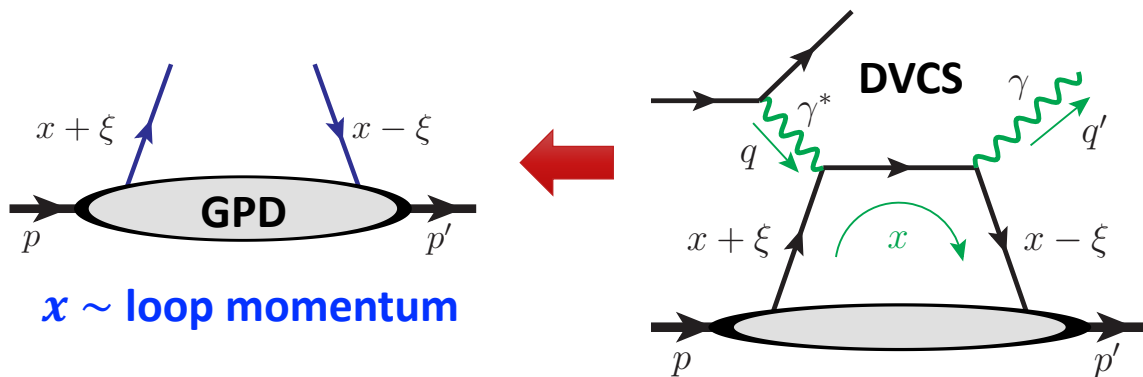
- Hard scale Q : allows pQCD, factorization
- Low scale t : probes non-pert. hadron structure

\rightarrow
Factorization

GPDs: $f_{i/h}(x, \xi, t; \mu)$

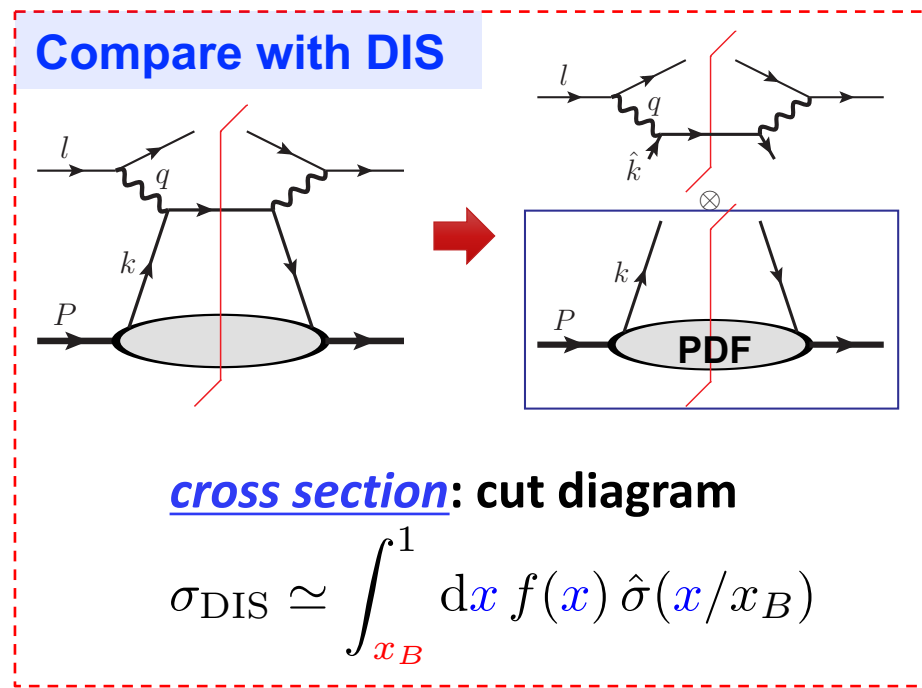
Why is the GPD's x -dependence so *difficult* to measure?

□ **Amplitude** nature: exclusive processes



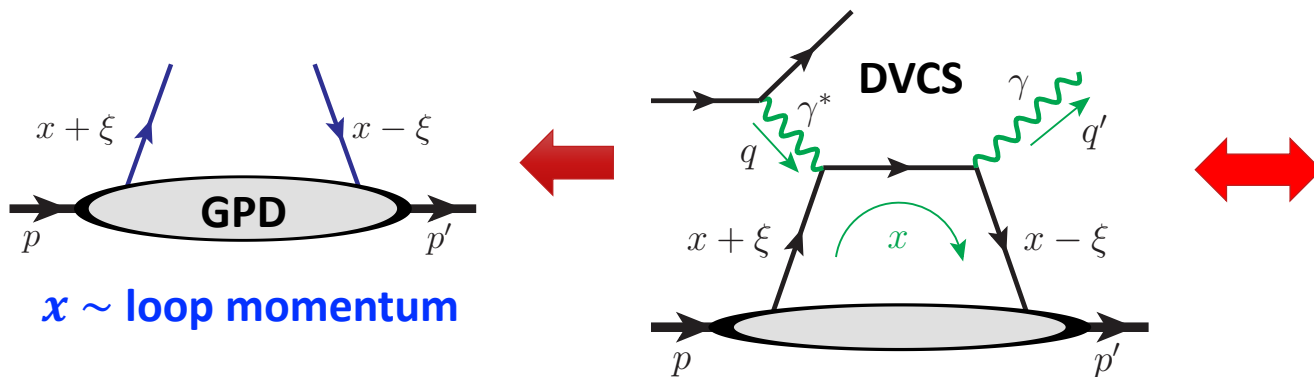
$$i\mathcal{M} \sim \int_{-1}^1 dx F(x, \xi, t) \cdot C(x, \xi; Q/\mu)$$

Full range of x , including $x = 0$; $x = \pm\xi$



Why is the GPD's x -dependence so *difficult* to measure?

Amplitude nature: exclusive processes



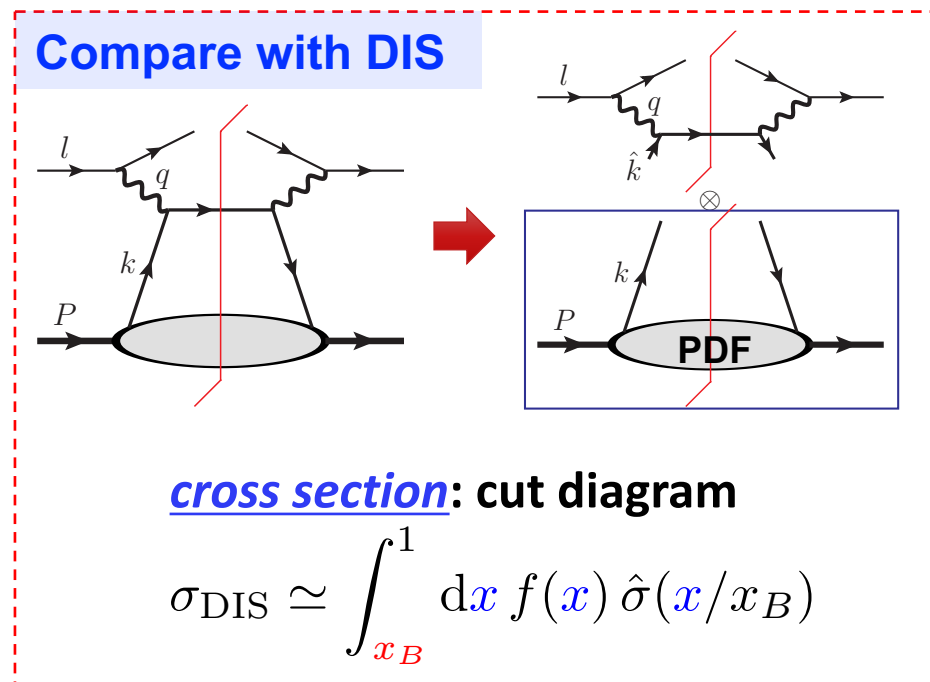
$$i\mathcal{M} \sim \int_{-1}^1 dx F(x, \xi, t) \cdot C(x, \xi; Q/\mu)$$

Full range of x , including $x = 0$; $x = \pm\xi$

Sensitivity to x : comes from $C(x, \xi; Q/\mu)$

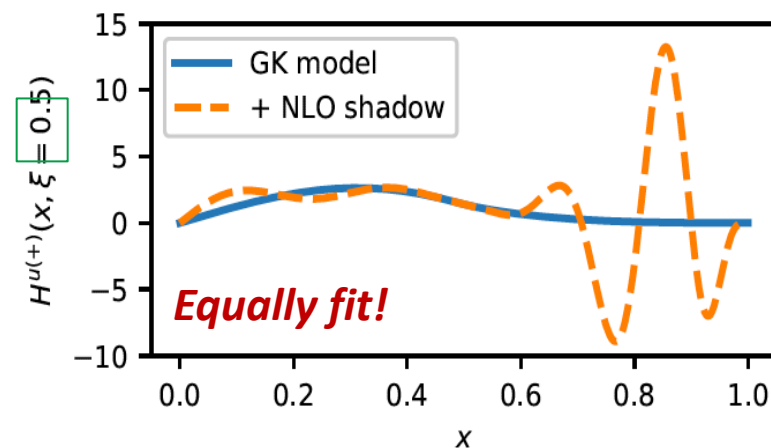
$$C(x, \xi; Q/\mu) = T(Q/\mu) \cdot G(x, \xi) \propto \frac{1}{x - \xi + i\epsilon} \dots$$

$$\Rightarrow i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - \xi + i\epsilon} \equiv "F_0(\xi, t)" \quad \text{"moment"}$$



cross section: cut diagram

$$\sigma_{\text{DIS}} \simeq \int_{x_B}^1 dx f(x) \hat{\sigma}(x/x_B)$$



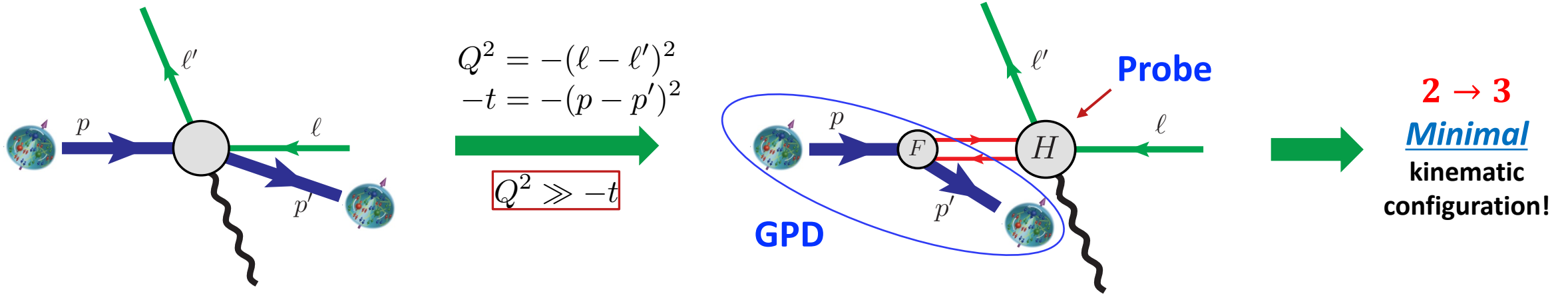
[Bertone et al. PRD '21]

How to Find Physical Processes to be Sensitive to GPDs?

□ Single diffractive hard exclusive processes (SDHEPs):

DVCS in Lab Frame: $h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$

Qiu & Yu, JHEP 08 (2022) 103
PRD 107 (2023) 014007
PRL 131 (2023) 161902



Classification of SDHEPs

Electro-production (JLab, EIC, ...)

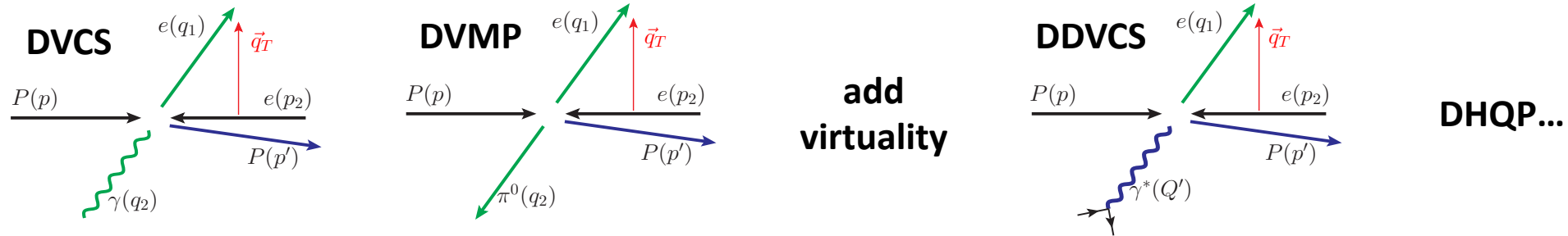
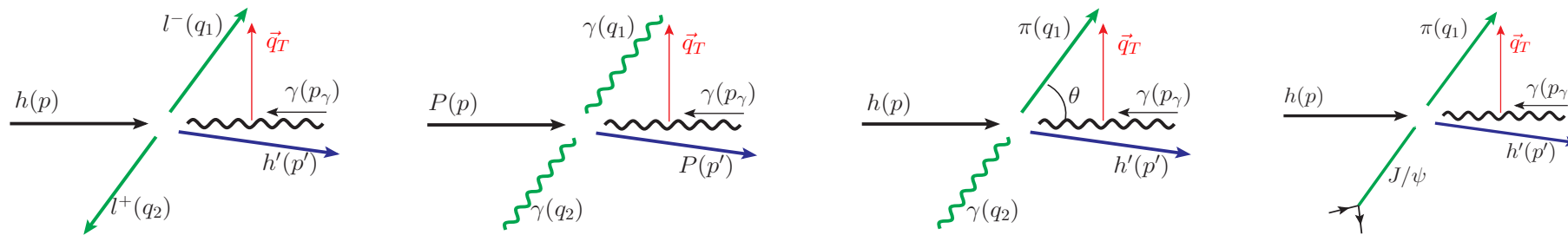
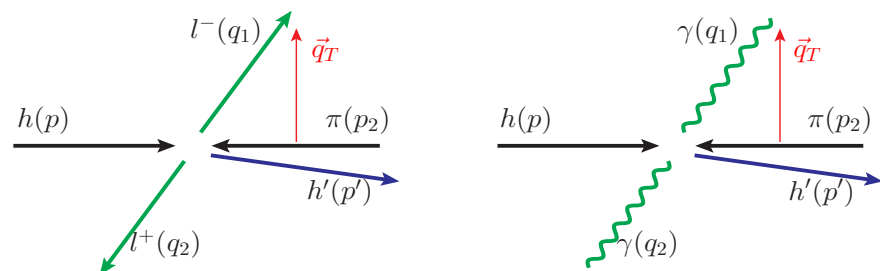


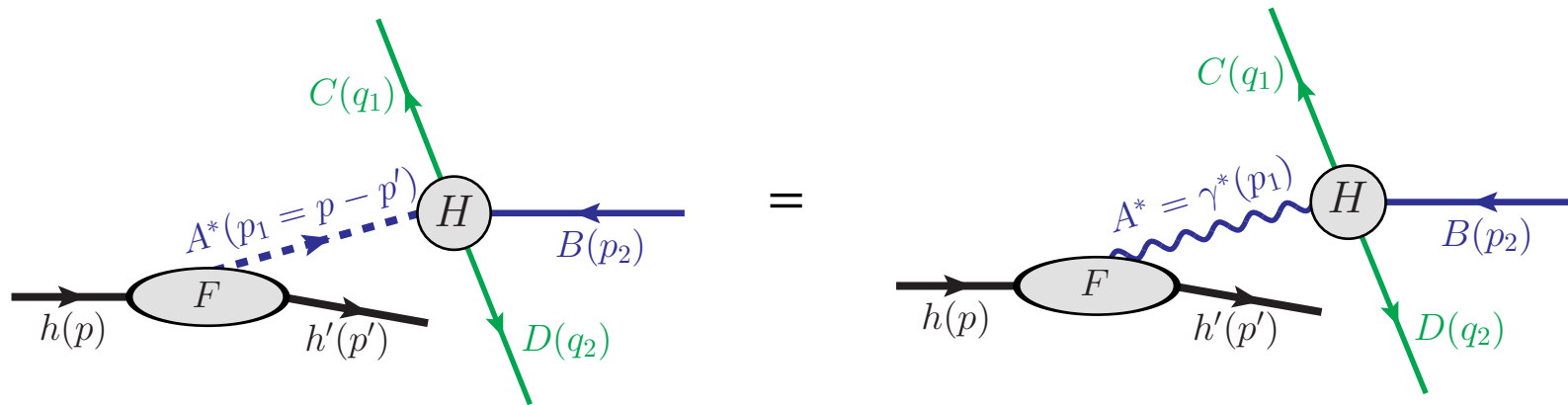
Photo-production (JLab, EIC, ...)



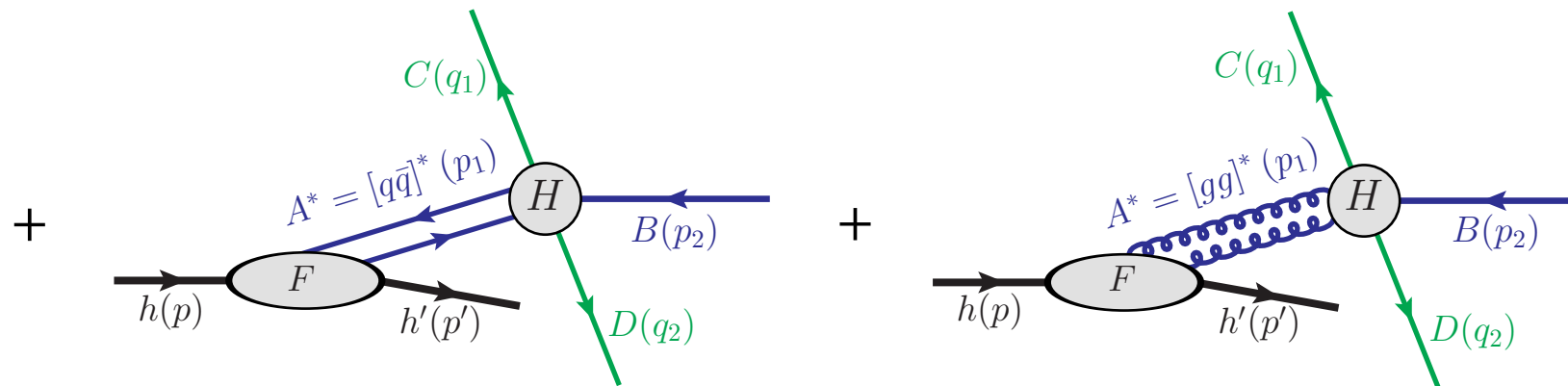
Meso-production (AMBER, J-PARC, ...)



SDHEP: Two-stage Paradigm and Power Expansion



Leading Power in $1/Q$

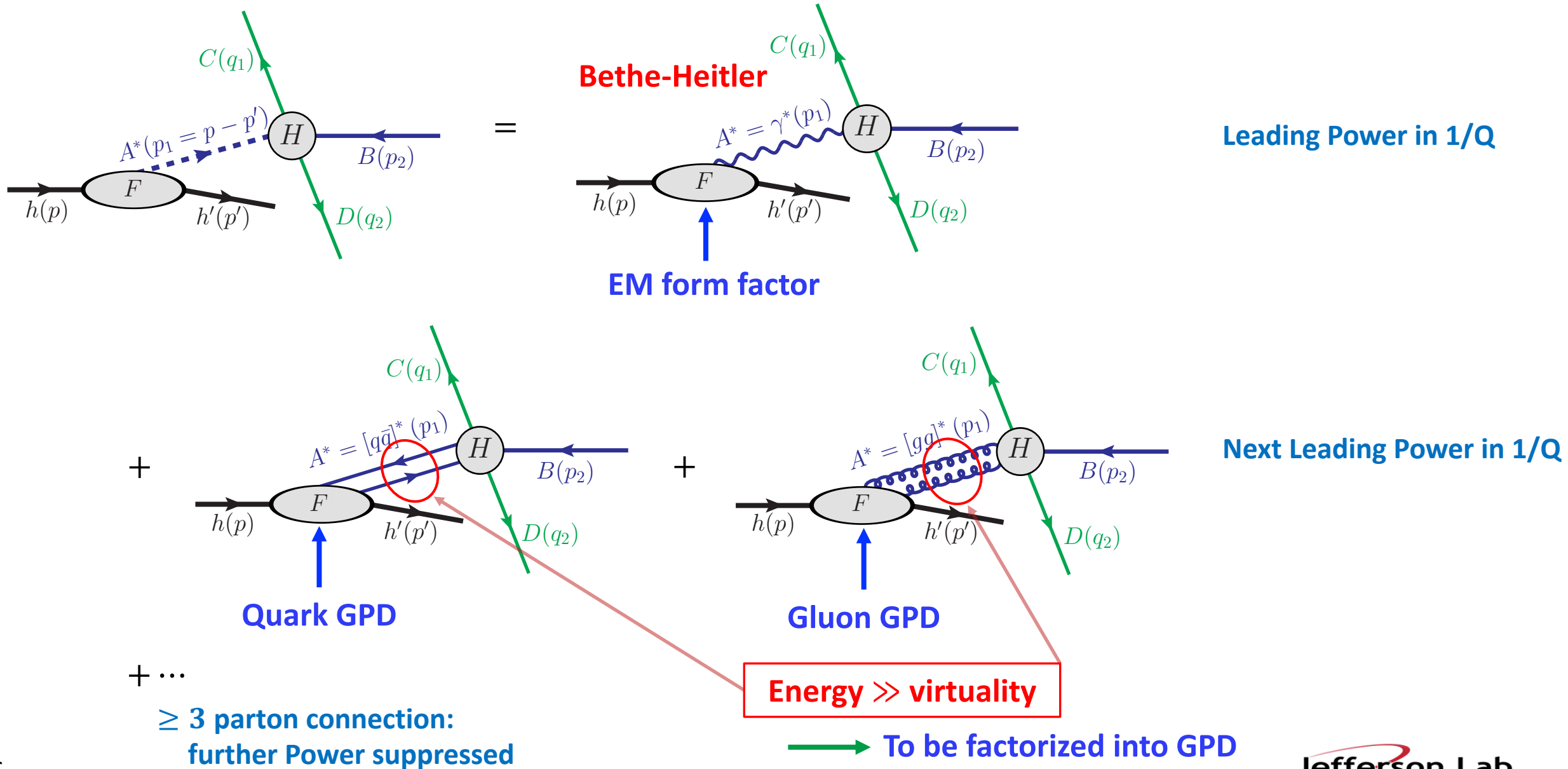


Next Leading Power in $1/Q$

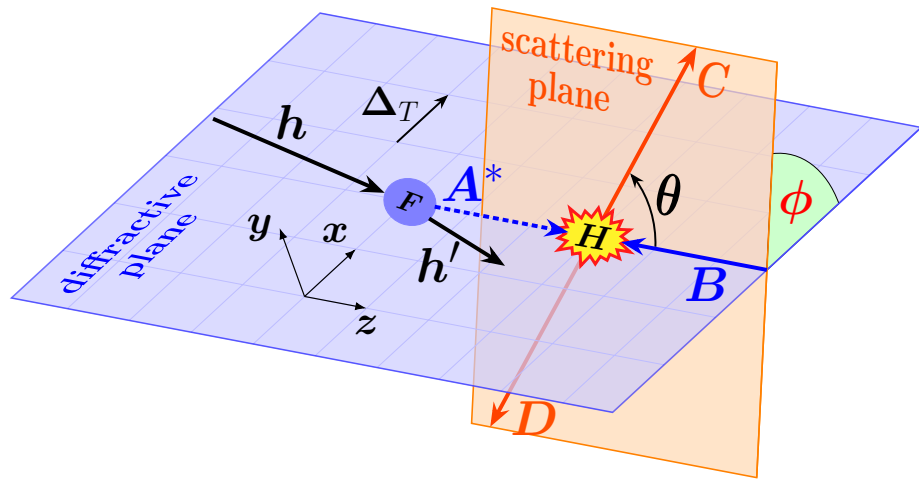
+ ...

≥ 3 parton connection:
further Power suppressed

SDHEP: Two-stage Paradigm and Power Expansion



Where does the ***x*-sensitivity** come from?



□ ***x*-sensitivity** $\Leftrightarrow 2 \rightarrow 2$ hard scattering:

Kinematics:

1. $\hat{s} = 2 \xi s / (1 + \xi)$ ← ξ

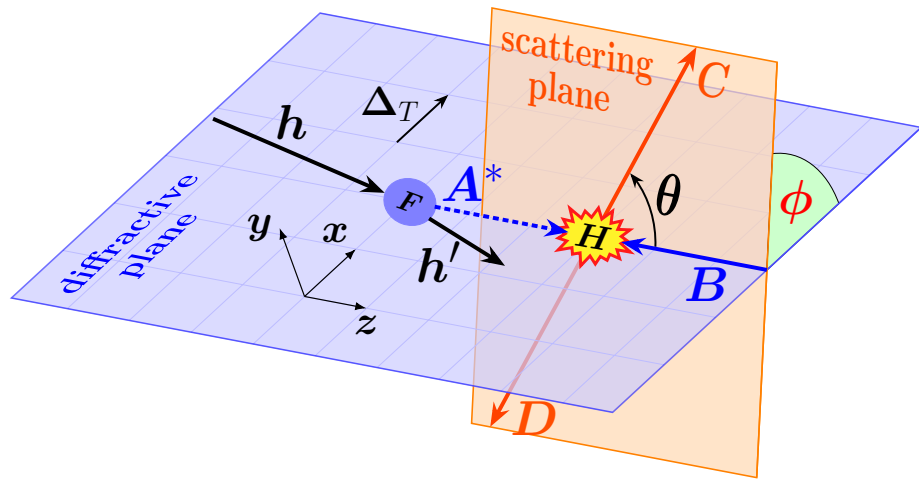
2. θ or $q_T = (\sqrt{\hat{s}}/2) \sin\theta$ ↔ x

3. ϕ ← (A^*B) spin states

$$\mathcal{M}(Q, \phi) \simeq \sum_A e^{i(\lambda_A - \lambda_B)\phi} \cdot \int_{-1}^1 dx F_A(x) C_A(x; Q) \quad (Q = \theta \text{ or } q_T)$$

[suppressing t and ξ dependence]

Where does the **x -sensitivity** come from?



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[suppressing t and ξ dependence]

■ **Moment-type sensitivity:** $C(x; Q) = G(x) \cdot T(Q)$ → $F_G = \int_{-1}^1 dx G(x) F(x, \xi, t)$ **Independent of Q**
Scaling for F_G

→ **Inversion problem: shadow GPD** $S_G = \int_{-1}^1 dx G(x) S(x, \xi) = 0$ [Bertone et al. PRD '21]

■ **Enhanced sensitivity:** $C(x; Q) \neq G(x) \cdot T(Q)$ → $d\sigma/dQ \sim |C(x; Q) \otimes_x F(x, \xi, t)|^2$

Moment-type Sensitivity: $h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$

PRD56 (1997) 5524; PRD58 (1998) 094018; PRD59 (1999) 074009

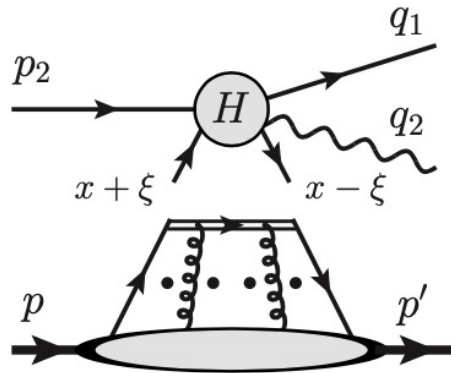
□ DVCS:

$h(p) = \text{Proton}(p)$, $h'(p') = \text{Proton}(p')$, $B(p_2) = \text{electron}(p_2)$, $C(q_1) = \text{electron}(q_1)$, $D(q_2) = \text{photon}(q_2)$

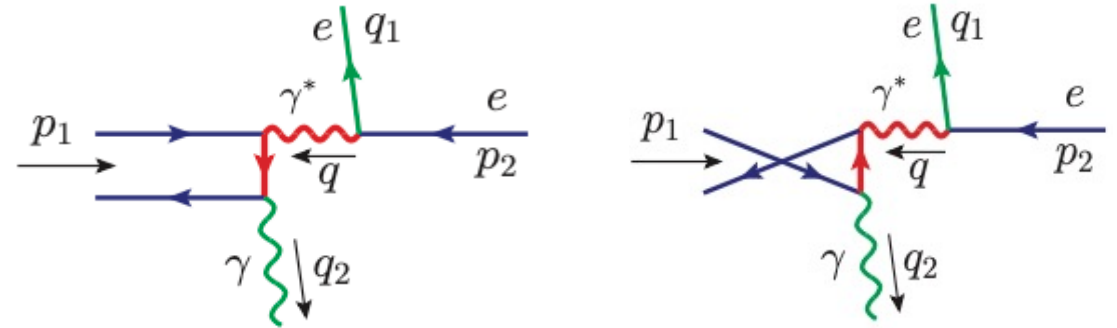
Factorization:

$$\xi = \frac{(p - p')^+}{(p + p')^+}$$

$$t = (p - p')^2$$



LO:



$$C^{(0)} \propto \frac{1}{x - \xi + i\varepsilon} - \frac{1}{x + \xi - i\varepsilon}$$



$$\mathcal{M}_{he \rightarrow h'e\gamma}^{(2)} = \sum_i \int_{-1}^1 dx F_i^h(x, \xi, t) C_{ie \rightarrow e\gamma}(x, \xi, q_T),$$

The x -integration is NOT sensitive to externally measured hard scale, q_T or Q^2 !
Need a very large range of Q^2 , but, cross section is strongly suppressed!

Moment-type Sensitivity: $h(p) + B(p_2) \rightarrow h'(p') + C(q_1) + D(q_2)$

PRD56 (1997) 5524; PRD58 (1998) 094018; PRD59 (1999) 074009

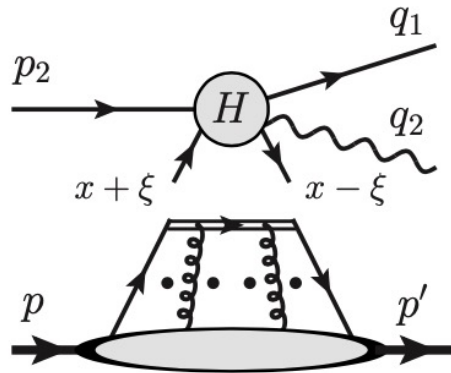
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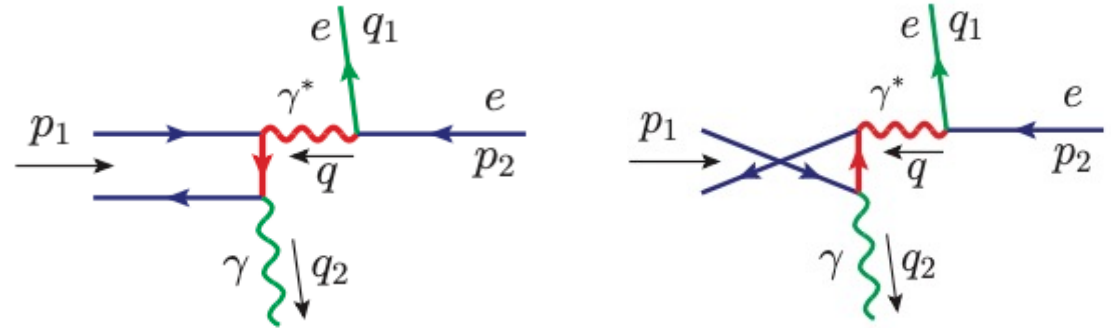
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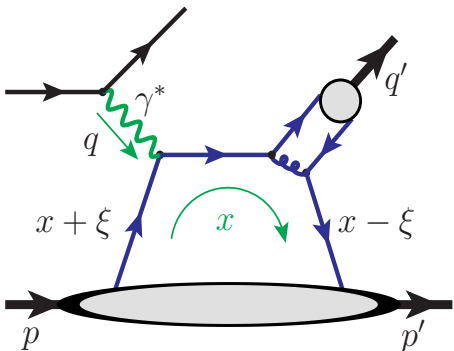


➔ $C^{(0)} \propto \frac{1}{x - \xi + i\epsilon} - \frac{1}{x + \xi - i\epsilon}$

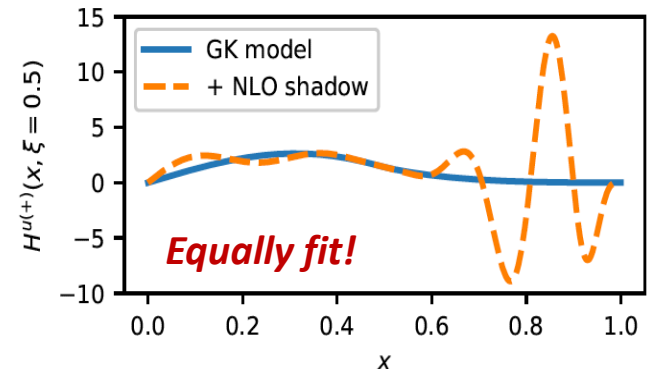
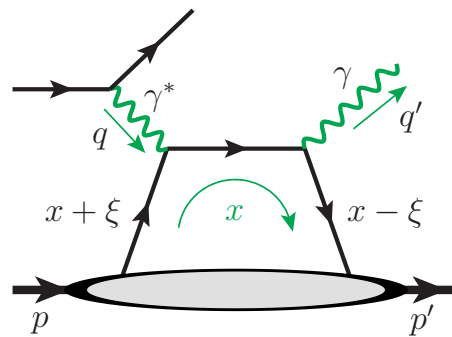
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□ DVMP:



Similar to



[Bertone et al. PRD '21]

What Kind of Process Could be Sensitive to the x -Dependence?

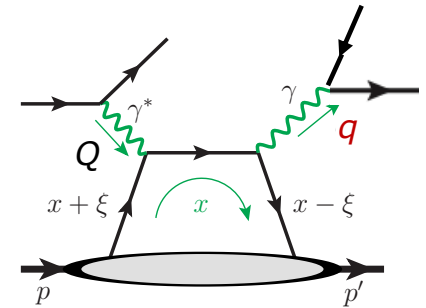
- Create an entanglement between the internal x and an externally measured variable?

$$i\mathcal{M} \propto \int_{-1}^1 dx \frac{F(x, \xi, t)}{x - x_p(\xi, q) + i\varepsilon}$$

Change external q to sample different part of x .

- Double DVCS (two scales):

$$x_p(\xi, q) = \xi \left(\frac{1 - q^2/Q^2}{1 + q^2/Q^2} \right) \rightarrow \xi \text{ same as DVCS if } q \rightarrow 0$$

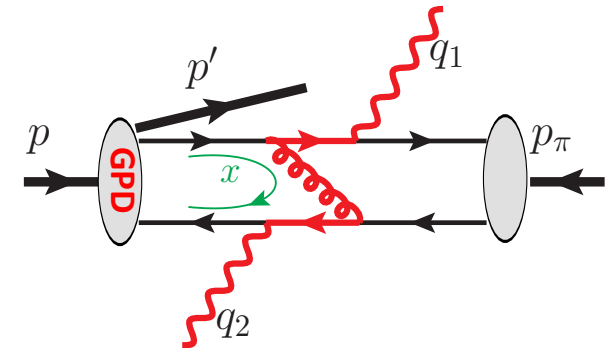


- Production of two back-to-back high p_T particles (say, two photons):

$$\pi^-(p_\pi) + P(p) \rightarrow \gamma(q_1) + \gamma(q_2) + N(p')$$

Hard scale: $q_T \gg \Lambda_{\text{QCD}}$ Soft scale: $t \sim \Lambda_{\text{QCD}}^2$

Qiu & Yu
JHEP 08 (2022) 103



- Factorization:

$$\mathcal{M}(t, \xi, q_T) = \int_{-1}^1 dx F(x, \xi, t; \mu) \cdot C(x, \xi; q_T/\mu) + \mathcal{O}(\Lambda_{\text{QCD}}/q_T) \quad \text{[suppressing pion DA factor]}$$

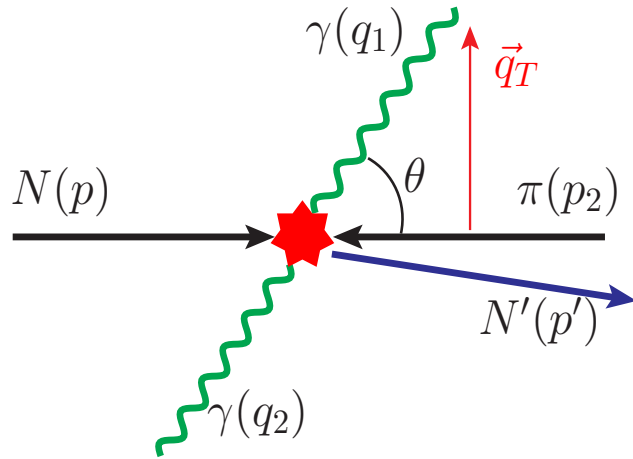
$$\frac{d\sigma}{dt d\xi dq_T} \sim |\mathcal{M}(t, \xi, q_T)|^2$$

q_T distribution is "conjugate" to x distribution

$$x \leftrightarrow q_T$$

Enhanced x -Sensitivity: (1) Diphoton Meso-production

Qiu & Yu, PRD 109 (2024) 074023



In addition to

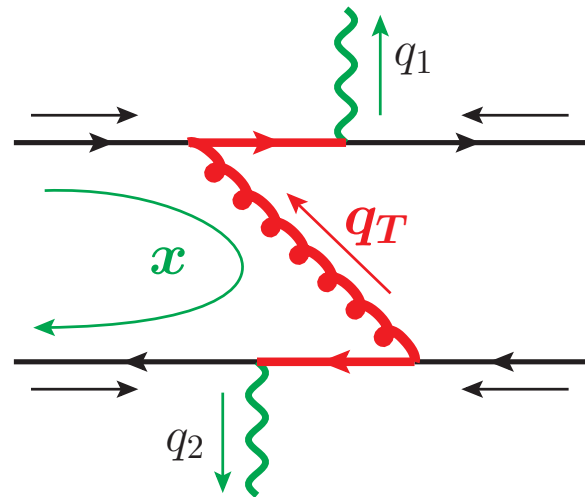
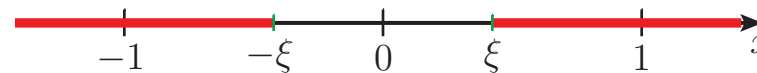
$$F_0(\xi, t) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \xi + i\epsilon}$$

When two photons are radiated from the same charged line

$i\mathcal{M}$ also contains

$$I(t, \xi; z, \theta) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \rho(z; \theta) + i\epsilon \operatorname{sgn}[\cos^2(\theta/2) - z]}$$

$$\rho(z; \theta) = \xi \cdot \left[\frac{1 - z + \tan^2(\theta/2) z}{1 - z - \tan^2(\theta/2) z} \right] \in (-\infty, -\xi] \cup [\xi, \infty)$$



Enhanced x -Sensitivity: (1) Diphoton Meso-production

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□ **Diphoton process:** $N\pi \rightarrow N'\gamma\gamma$: (1) $p\pi^- \rightarrow n\gamma\gamma$; (2) $n\pi^+ \rightarrow p\gamma\gamma$

$$\frac{d\sigma}{d|t|d\xi d\cos\theta} = 2\pi \left(\alpha_e \alpha_s \frac{C_F}{N_c} \right)^2 \frac{1}{\xi^2 s^3} \cdot \left[(1 - \xi^2) \sum_{\alpha=\pm} \left(|\mathcal{M}_\alpha^{[\tilde{H}]}|^2 + |\tilde{\mathcal{M}}_\alpha^{[H]}|^2 \right) - \left(\xi^2 + \frac{t}{4m^2} \right) \sum_{\alpha=\pm} |\tilde{\mathcal{M}}_\alpha^{[E]}|^2 - \frac{\xi^2 t}{4m^2} \sum_{\alpha=\pm} |\mathcal{M}_\alpha^{[\tilde{E}]}|^2 - 2\xi^2 \sum_{\alpha=\pm} \text{Re} \left(\tilde{\mathcal{M}}_\alpha^{[H]} \tilde{\mathcal{M}}_\alpha^{[E]*} + \mathcal{M}_\alpha^{[\tilde{H}]} \mathcal{M}_\alpha^{[\tilde{E}]*} \right) \right]$$

Nucleon transition GPDs

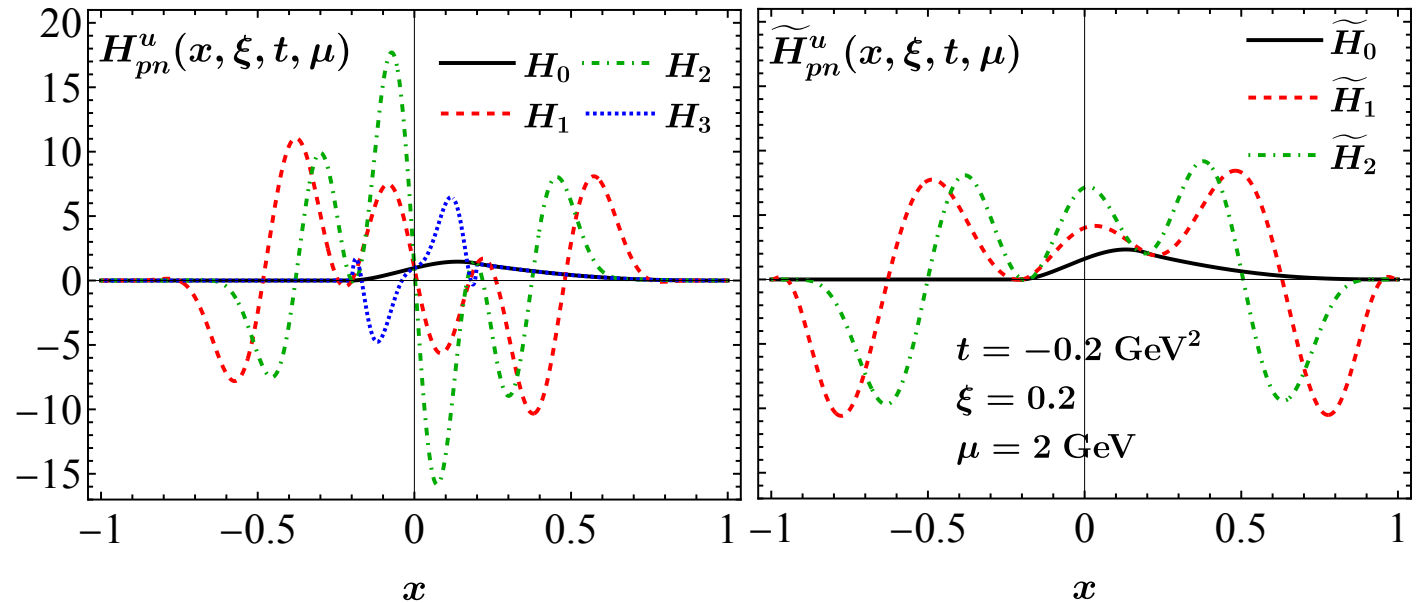
$$H_{pn}^u = H_p^u - H_p^d, \text{ etc.}$$

GPD models = GK model + shadow GPDs

Goloskokov & Kroll, '05, '07, '09

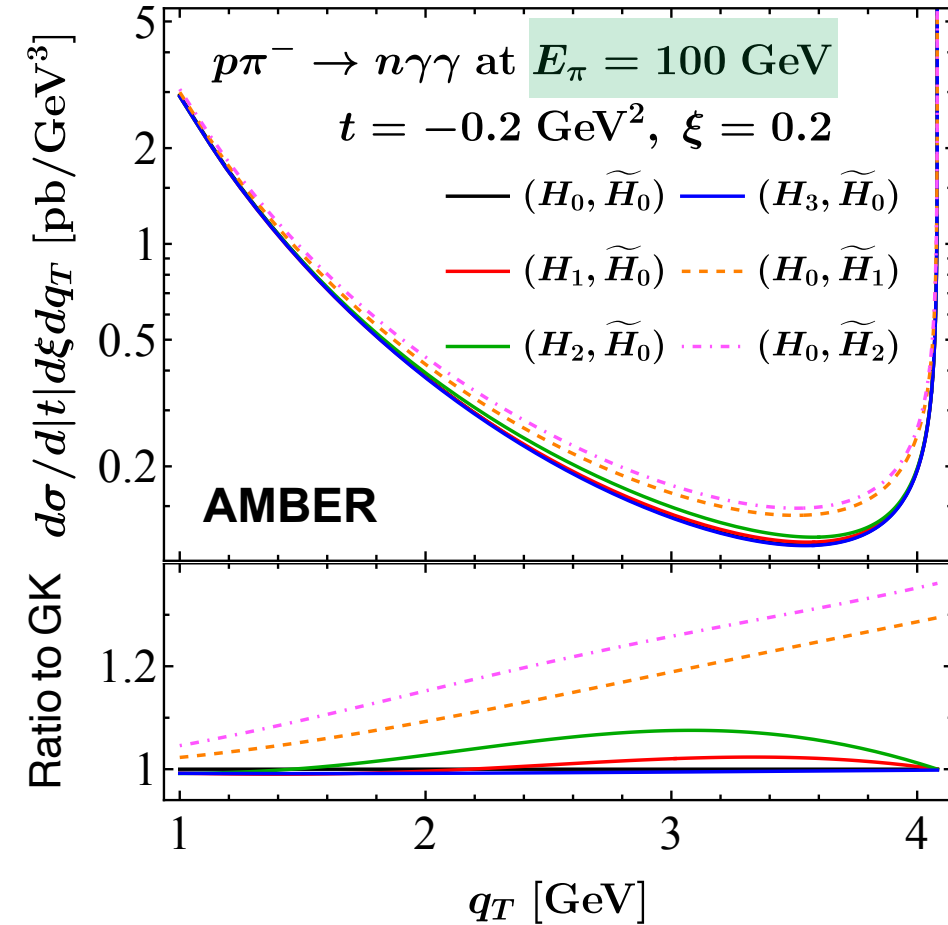
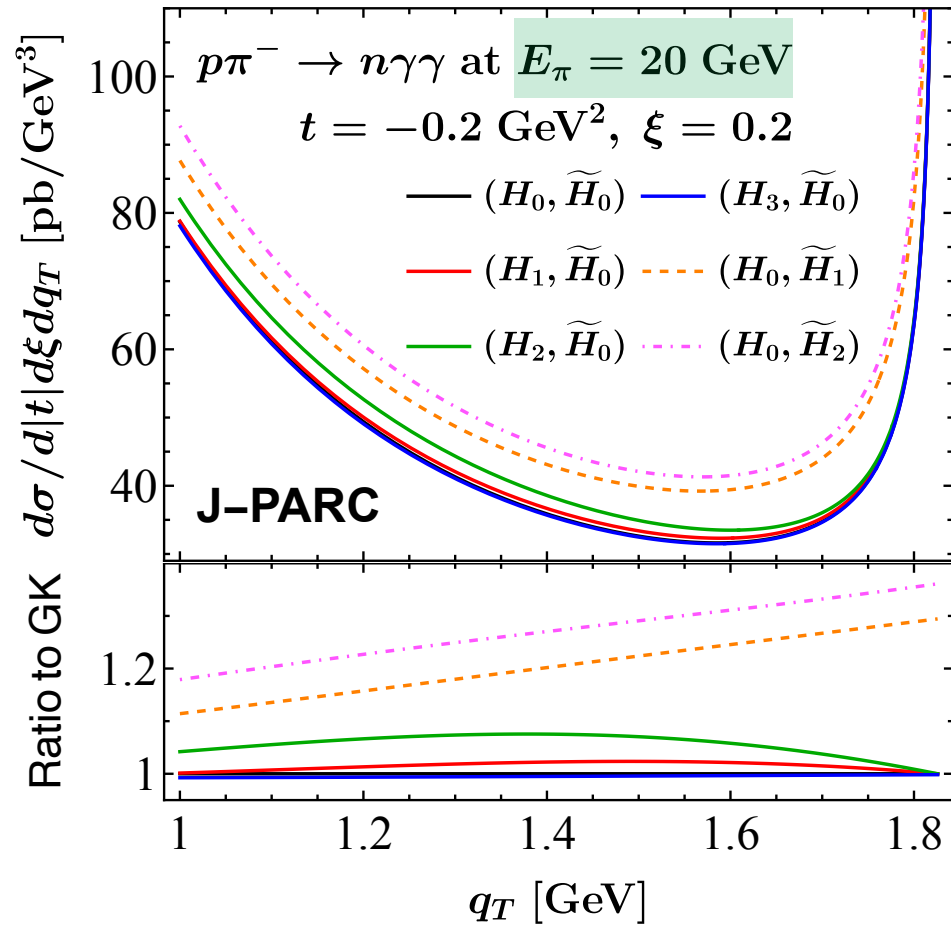
Bertone et al. '21
Moffat et al. '23

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$



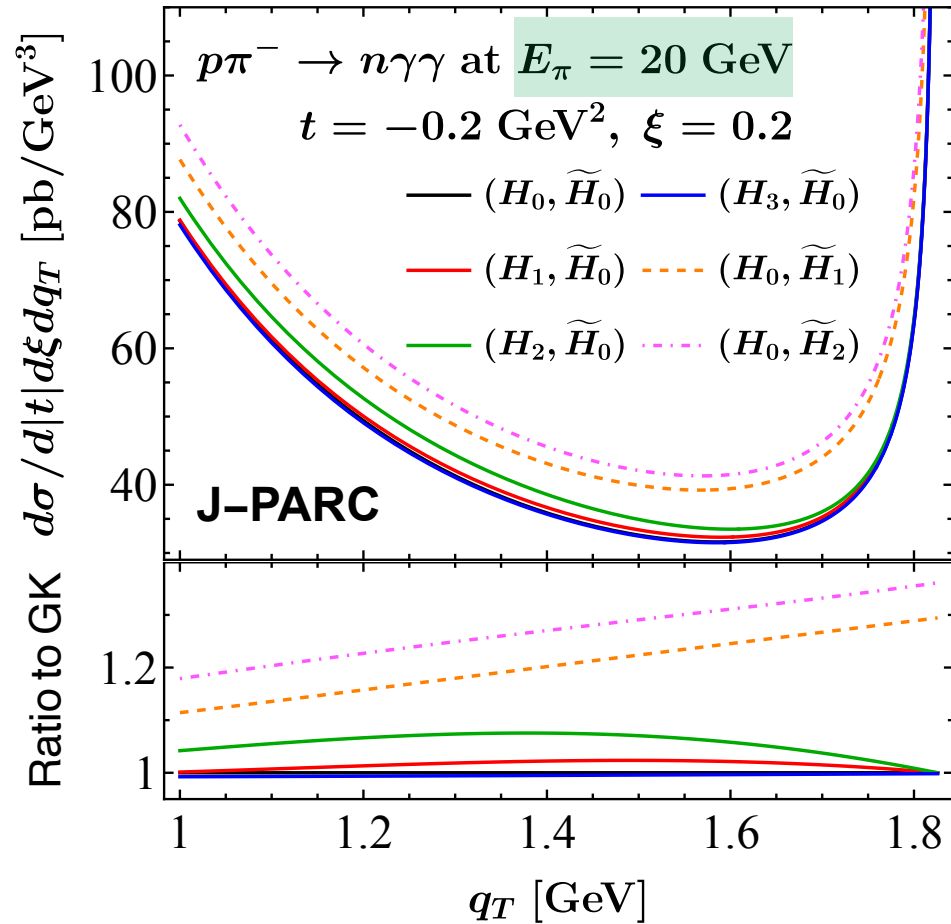
Enhanced x -Sensitivity: (1) Diphoton Meso-production

Qiu & Yu, PRD 109 (2024) 074023



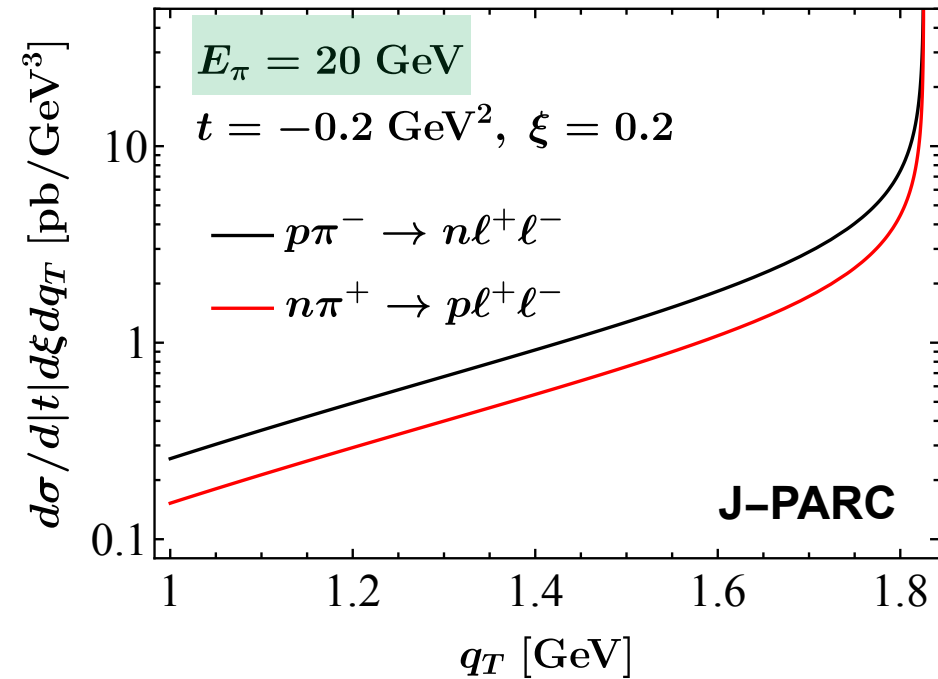
Enhanced x -Sensitivity: (1) Diphoton Meso-production

Qiu & Yu, PRD 109 (2024) 074023



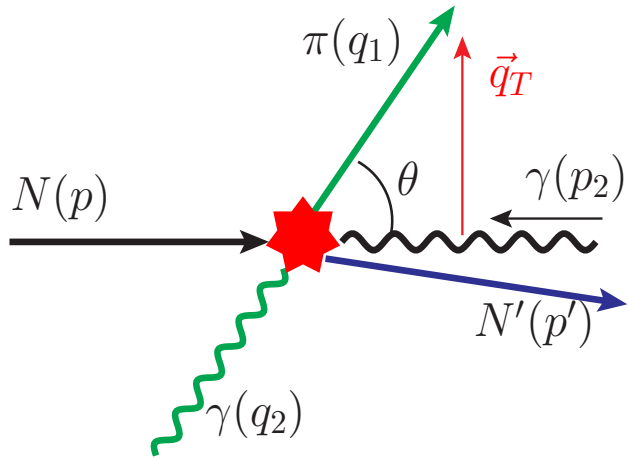
Exclusive Drell-Yan dilepton production

$$N + \pi \rightarrow N' + \gamma^* [\rightarrow \ell^+ + \ell^-]$$



- Lower rate
- Blind to shadow GPDs

Enhanced x -Sensitivity: (2) γ - π Pair Photoproduction



$i\mathcal{M}$ also contains the special integral:

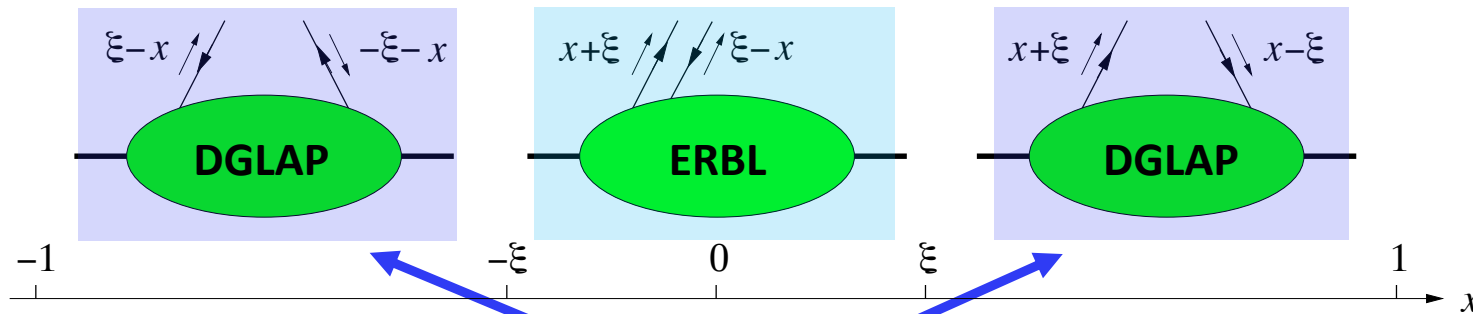
$$I'(t, \xi; z, \theta) = \int_{-1}^1 \frac{dx F(x, \xi, t)}{x - \rho'(z; \theta) + i\epsilon}$$

$$\rho'(z; \theta) = \xi \cdot \left[\frac{\cos^2(\theta/2) (1 - z) - z}{\cos^2(\theta/2) (1 - z) + z} \right] \in [-\xi, \xi]$$



G. Duplancic et al., JHEP 11 (2018) 179
 G. Duplancic et al., JHEP 03 (2023) 241
 G. Duplancic et al., PRD 107 (2023), 094023
 Qiu & Yu, PRL 131 (2023), 161902

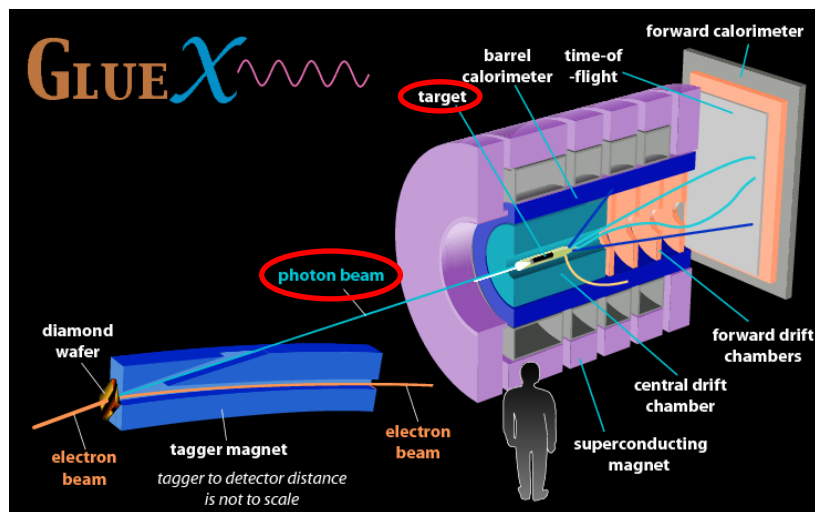
Complementary sensitivity:



$$N \pi \rightarrow N' \gamma \gamma$$

Enhanced x -Sensitivity: (2) γ - π Pair Photoproduction (at JLab Hall D)

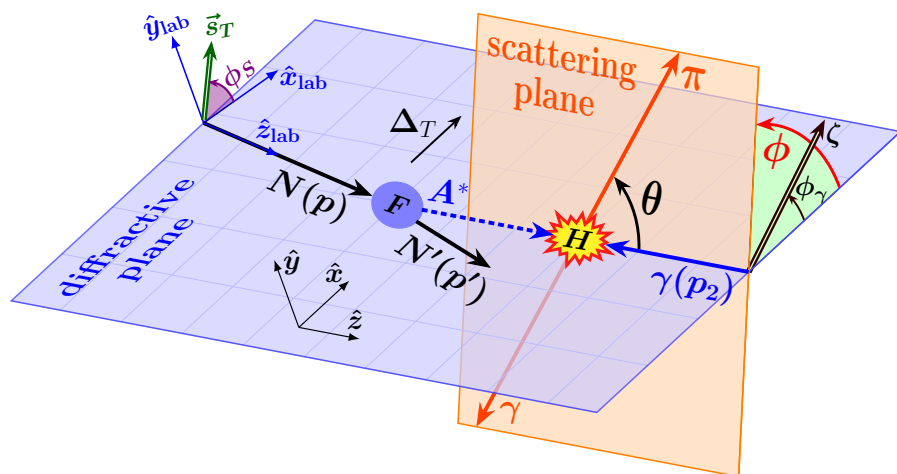
Qiu & Yu, PRL 131 (2023), 161902



□ Polarization asymmetries:

$$\frac{d\sigma}{d|t| d\xi d \cos \theta d\phi} = \frac{1}{2\pi} \frac{d\sigma}{d|t| d\xi d \cos \theta} \cdot [1 + \lambda_N \lambda_\gamma A_{LL} + \zeta A_{UT} \cos 2(\phi - \phi_\gamma) + \lambda_N \zeta A_{LT} \sin 2(\phi - \phi_\gamma)]$$

$$\frac{d\sigma}{d|t| d\xi d \cos \theta} = \pi (\alpha_e \alpha_s)^2 \left(\frac{C_F}{N_c} \right)^2 \frac{1 - \xi^2}{\xi^2 s^3} \Sigma_{UU}$$



$$\begin{aligned} \Sigma_{UU} &= |\mathcal{M}_+^{[\tilde{H}]}|^2 + |\mathcal{M}_-^{[\tilde{H}]}|^2 + |\tilde{\mathcal{M}}_+^{[H]}|^2 + |\tilde{\mathcal{M}}_-^{[H]}|^2, \\ A_{LL} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} \right], \\ A_{UT} &= 2 \Sigma_{UU}^{-1} \text{Re} \left[\tilde{\mathcal{M}}_+^{[H]} \tilde{\mathcal{M}}_-^{[H]*} - \mathcal{M}_+^{[\tilde{H}]} \mathcal{M}_-^{[\tilde{H}]*} \right], \\ A_{LT} &= 2 \Sigma_{UU}^{-1} \text{Im} \left[\mathcal{M}_+^{[\tilde{H}]} \tilde{\mathcal{M}}_-^{[H]*} + \mathcal{M}_-^{[\tilde{H}]} \tilde{\mathcal{M}}_+^{[H]*} \right]. \end{aligned}$$

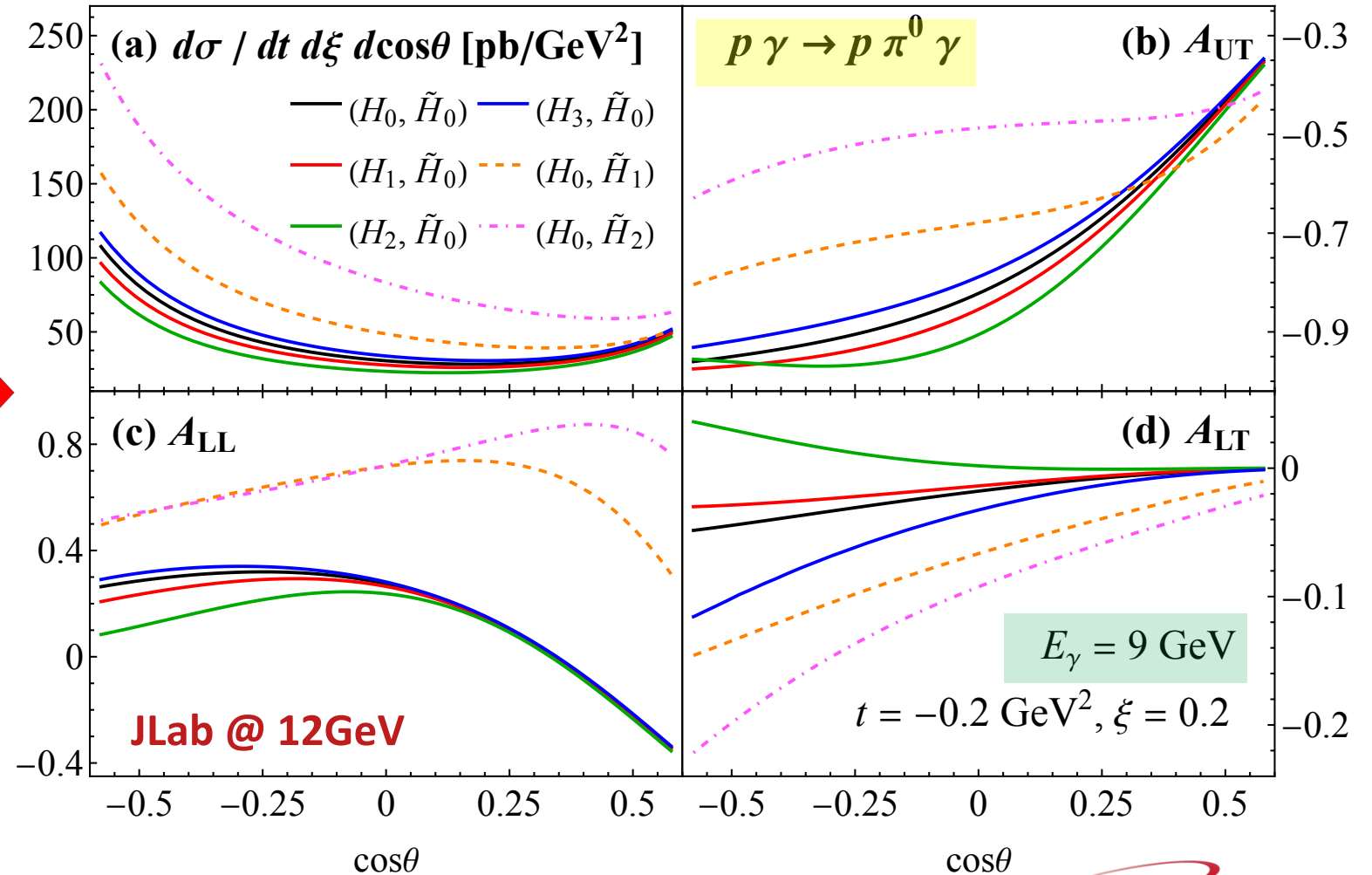
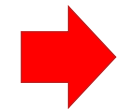
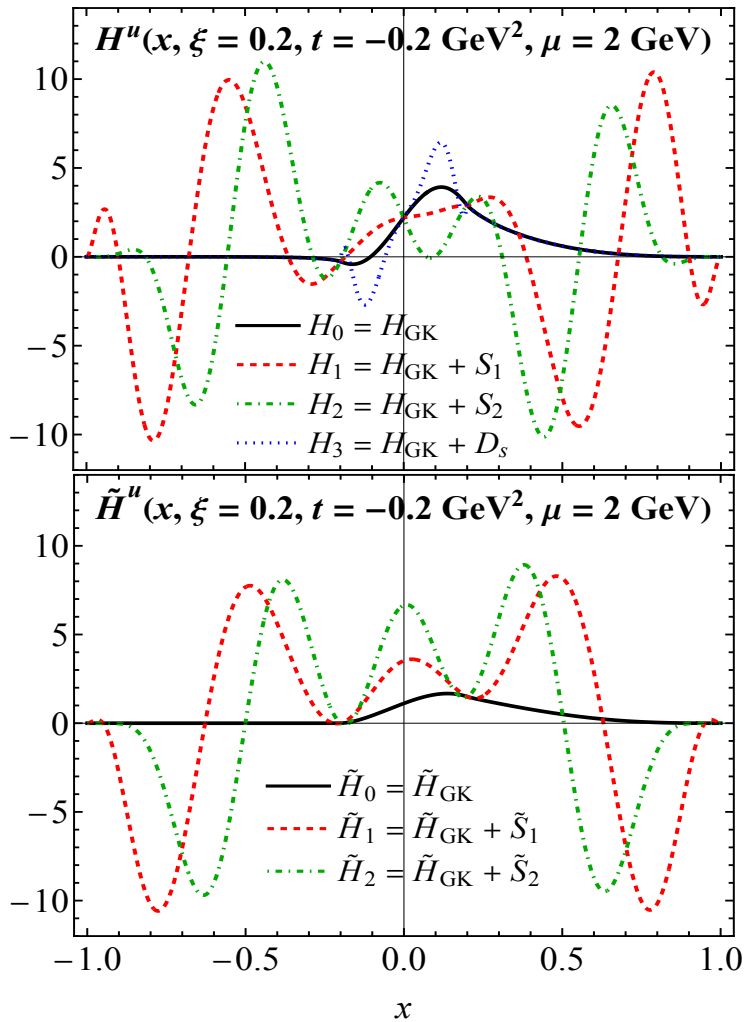
Neglecting: (1) E and \tilde{E} ; (2) gluon channel

Enhanced x -sensitivity: (2) γ - π pair photoproduction (at JLab Hall D)

Goloskokov, Kroll, '05, '07, '09
 Bertone et al. '21
 Moffat et al. '23

GPD models = GK model + shadow GPDs

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

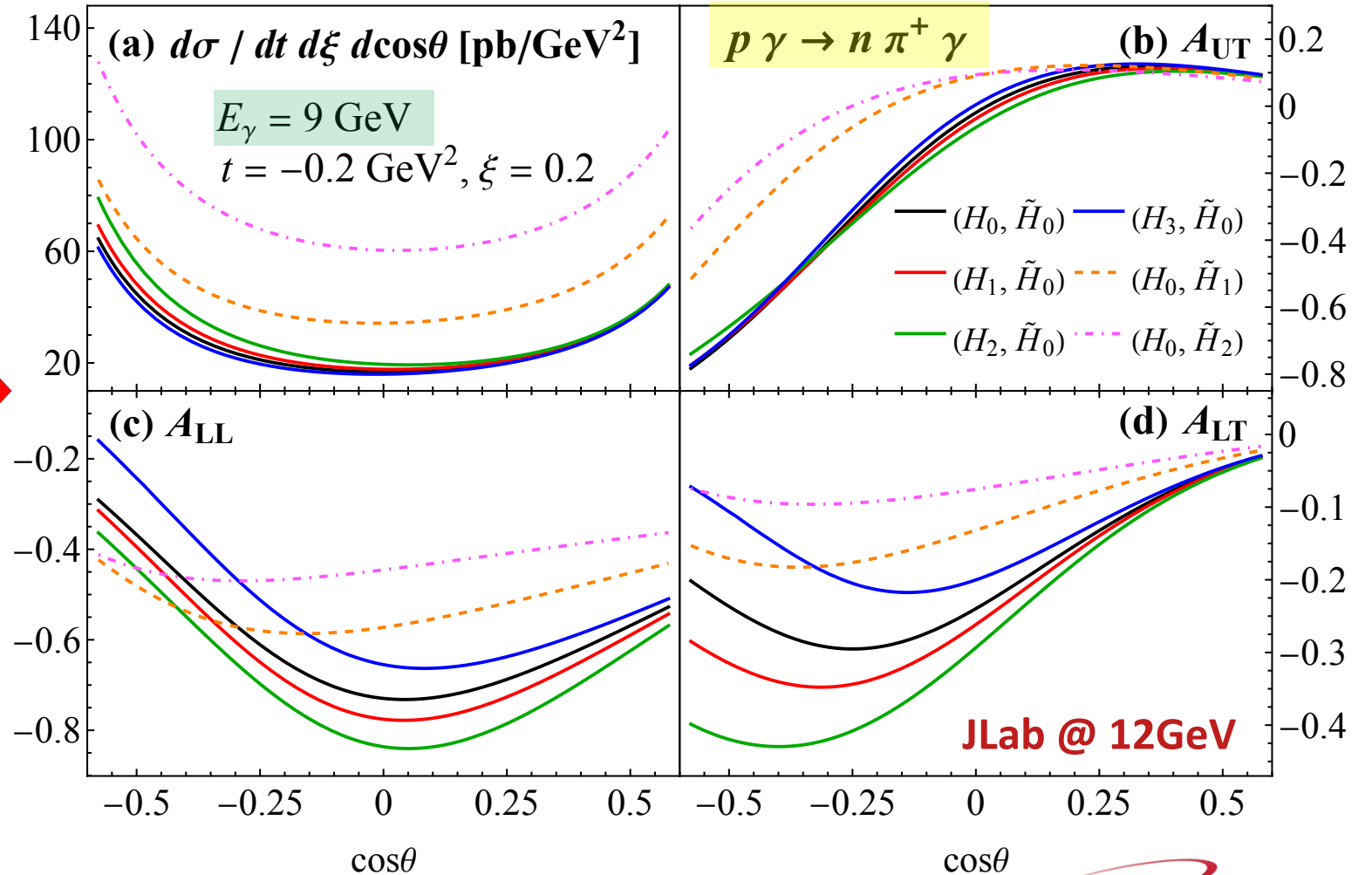
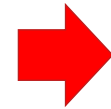
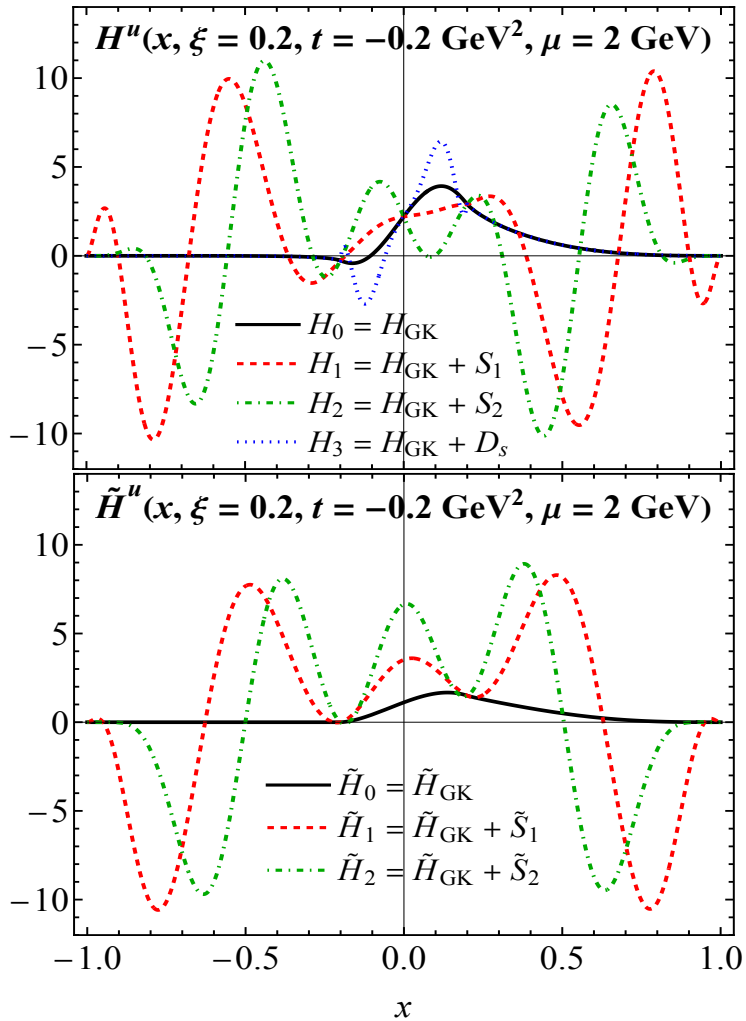


Enhanced x -sensitivity: (2) γ - π pair photoproduction (at JLab Hall D)

GPD models = GK model + shadow GPDs

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
Bertone et al. '21
Moffat et al. '23

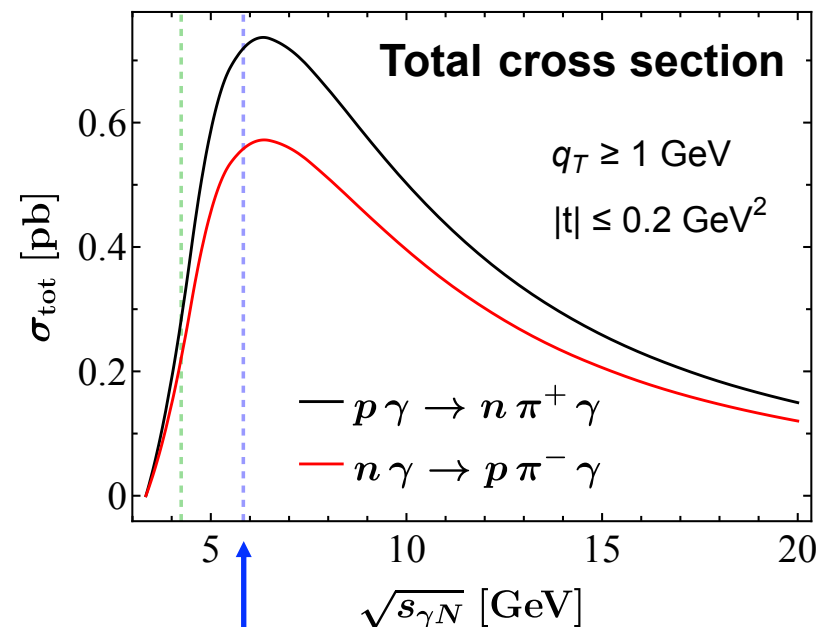
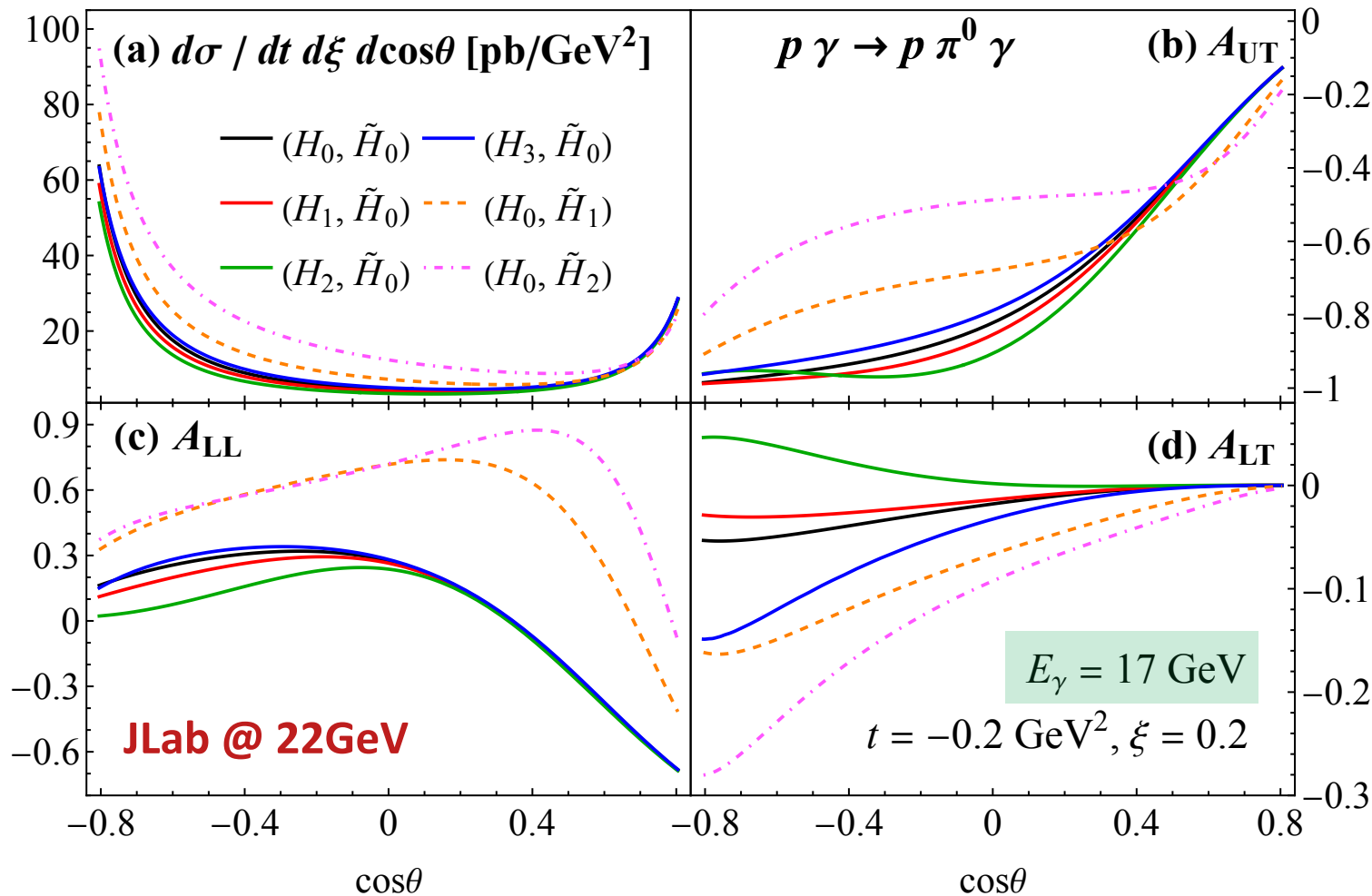


Enhanced x -sensitivity: (2) γ - π pair photoproduction (at upgraded energy)

GPD models = GK model + shadow GPDs

$$\int_{-1}^1 \frac{dx S(x, \xi)}{x - \xi \pm i\epsilon} = 0$$

Goloskokov, Kroll, '05, '07, '09
 Bertone et al. '21
 Moffat et al. '23
 Qiu & Yu, '23

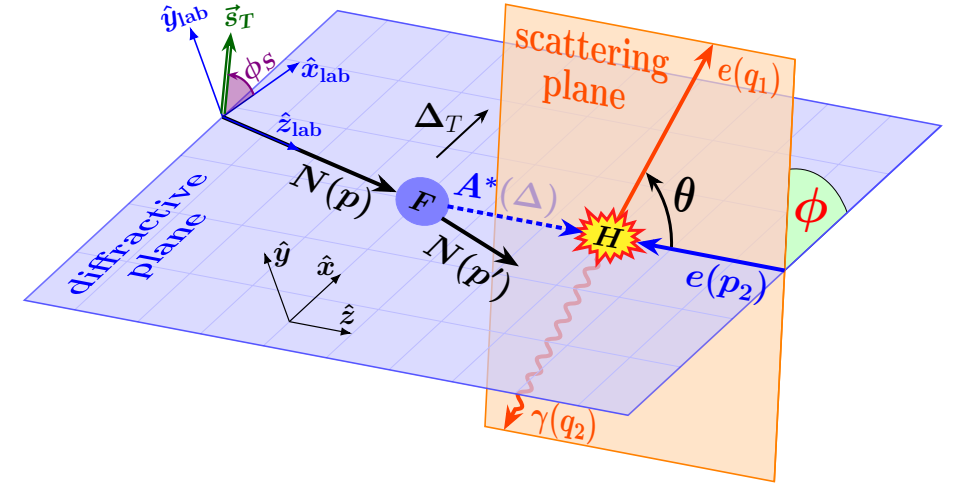
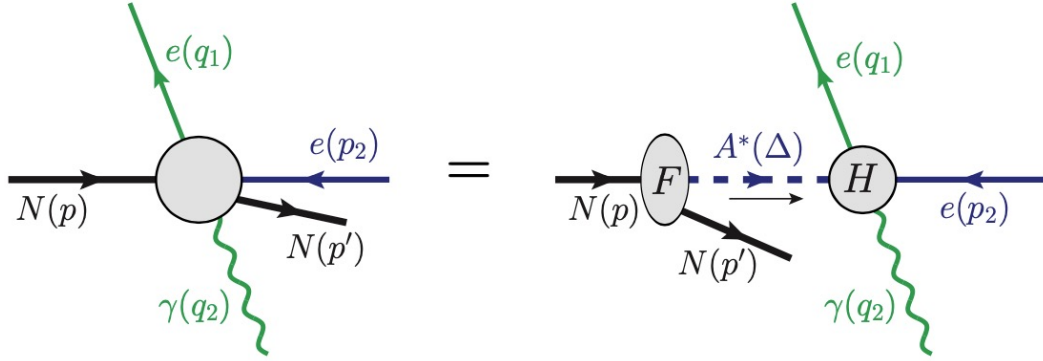


JLab @ 22GeV

A. Accardi et al.
 [arXiv:2306.09360]

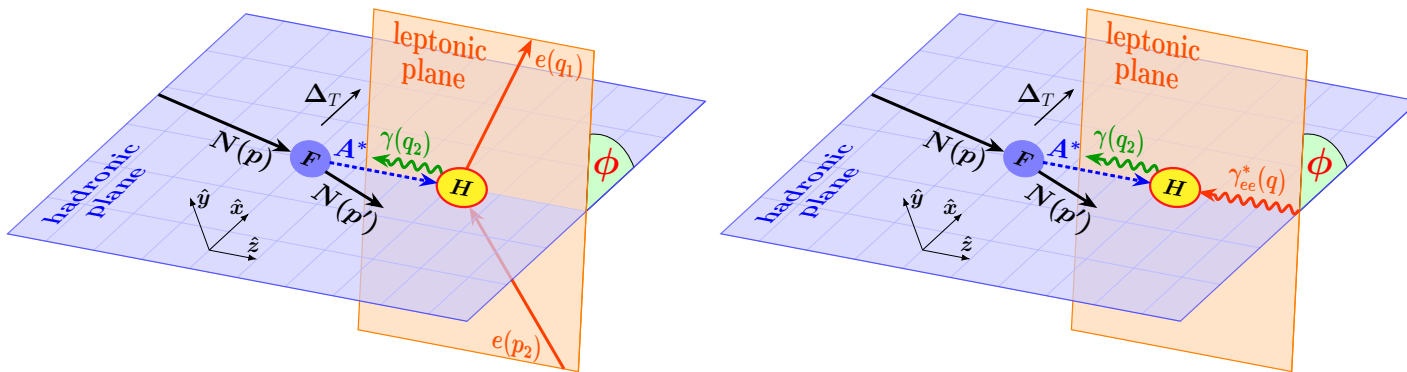
Angular Modulations – Separation of Different GPDs & Global Analyses

- GPDs depend on the choice of Frame (light-cone n-vector to define the “+” component):



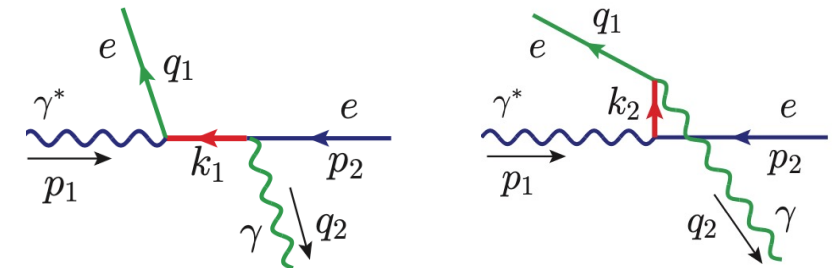
Angular modulation between “diffractive” and “scattering” planes to select the spin-state of A^* - different GPDs

- Experimental Breit frame is not ideal:



Angular modulation between “leptonic” and “hadronic” planes do not necessarily select the definite spin-state of A^* - different GPDs!

BH is not a “t”-channel process:

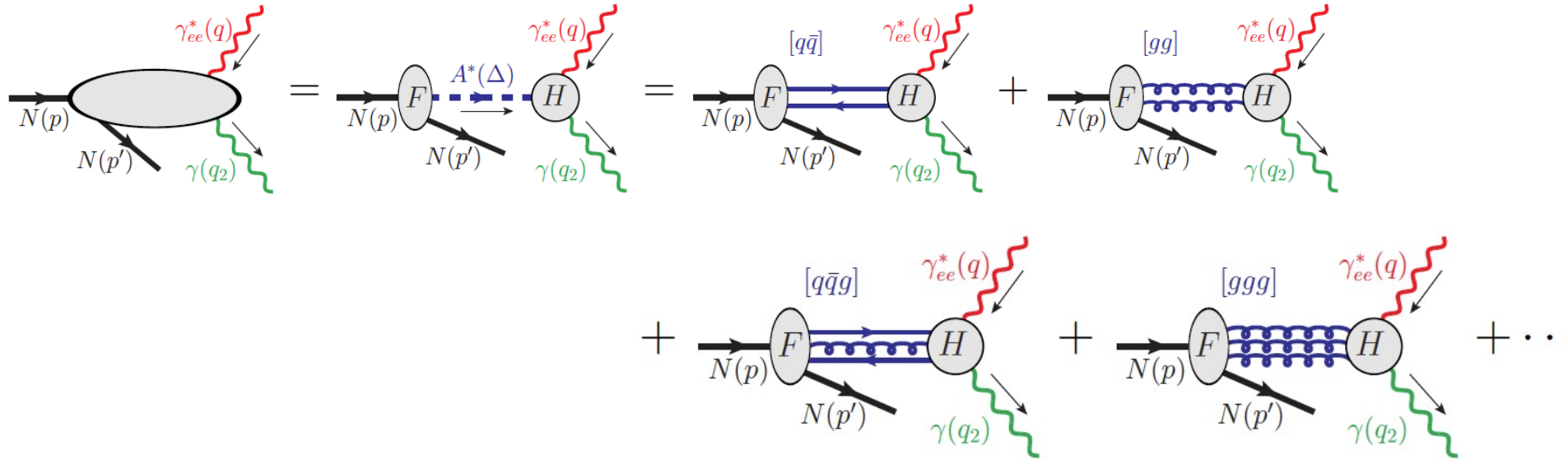


Propagators of k_1 & k_2 have ϕ -dependence!

DVCS Compton Form Factors (CFFs)

Qiu, Sato & Yu in preparation

□ Like structure functions, CFFs include all power contributions:



In terms of EM gauge invariance, this virtual photon Compton scattering amplitude can be decomposed into **18** scalar Compton Form Factors.

□ Extraction of GPDs from QCD global analyses:

- 1) Universality of GPDs – unlike PDFs, the choice of the light-cone vector to define “+” component is not unique! When considering the power corrections, it is critical to carefully define the GPDs!
- 2) Angular modulation is critically important for separating different GPDs, but, consistent choice of the “angle” to define the modulation is crucial!

Summary and Outlook

- GPDs are fundamental, carrying rich information on:
 - Tomographic images of confined quarks and gluons
 - Underline dynamics of hadronic properties
- It is challenging to extract the x -dependence of GPDs
 - The relative momentum fraction x needs to be entangled with externally measured hard scales!

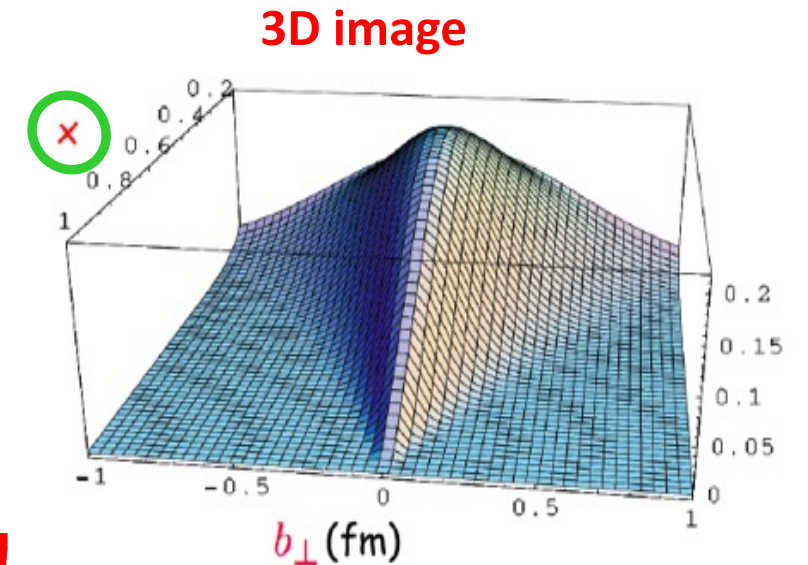
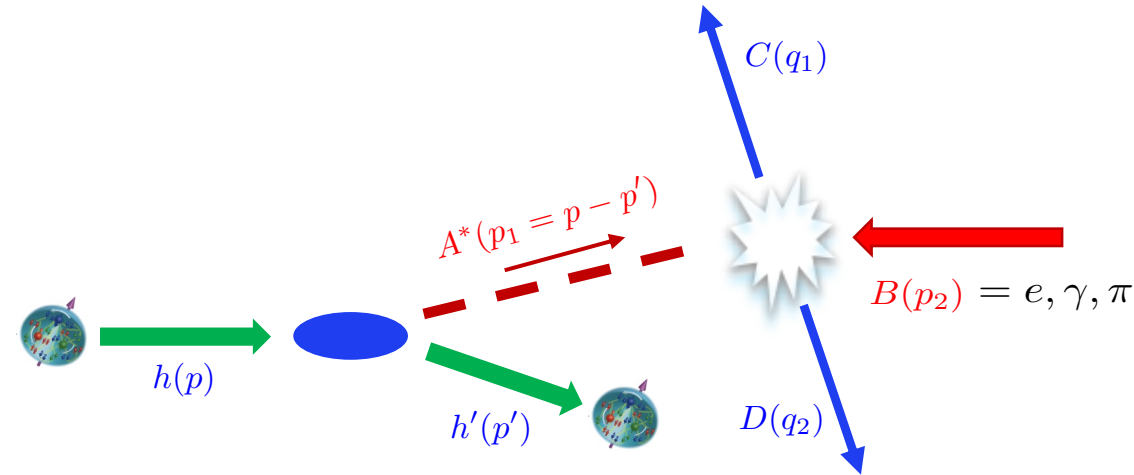
□ QCD Global analyses to extract GPDs:

- Need more processes, more observables (see Silvia's talk)
- Need input from lattice QCD (see Herve's talk at QCD Evolution)
- Controlled GPD evolution (see Valerio's talk)
- Need to solve the end point issues of exclusive processes (momentum of active parton goes to zero)
- ...

A long but challenging & exciting way to go!

Many on-going efforts:

PARTON Collab, QuantOM Collab,
QGT Collab, FemtoNET, ...



Thanks!