

# NNLO global analysis of helicity PDFs

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$$\Delta q(x) = \text{---} \quad \begin{array}{c} \text{---} \\ \text{---} \end{array}$$

A diagram illustrating the operator definition of the quark distribution  $\Delta q(x)$ . It consists of two red circles representing quarks. The left quark contains a white center with a yellow arrow pointing right, and a green arrow pointing right from its side. The right quark contains a white center with a yellow arrow pointing left, and a green arrow pointing right from its side. A minus sign between the circles indicates the difference between these two contributions.

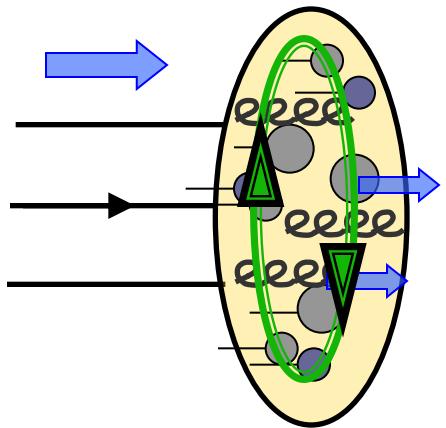
$$\Delta g(x) = \text{---} \quad \begin{array}{c} \text{---} \\ \text{---} \end{array}$$

A diagram illustrating the operator definition of the gluon distribution  $\Delta g(x)$ . It consists of two red circles representing gluons. Both circles contain a white center with a yellow arrow pointing right, and a green arrow pointing right from its side. The left quark is labeled "eee" above it. A minus sign between the circles indicates the difference between these two contributions.

- in QCD: operator definition  
→ dependence on “resolution” scale  $\mu$

# proton spin:

Jaffe, Manohar; Chen et al;  
Wakamatsu; Hatta; ...



$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_q + L_g$$

$$\Delta\Sigma = \int_0^1 dx [\Delta u + \Delta \bar{u} + \Delta d + \Delta \bar{d} + \Delta s + \Delta \bar{s}](x) \ll 1$$

$$\Delta G = \int_0^1 dx \Delta g(x)$$

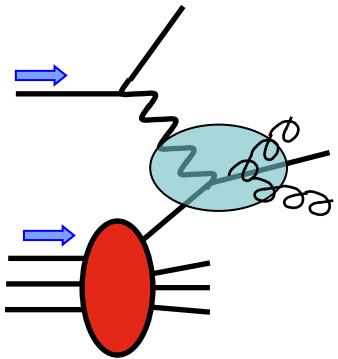
# Outline:

- Helicity PDFs at NNLO: framework
- Results
- Concluding remarks

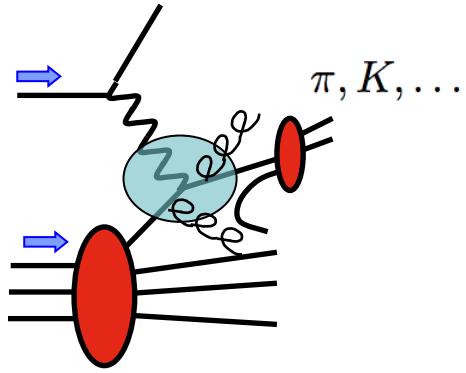
Thanks to my collaborators

I. Borsa, D. de Florian, R. Sassot, M. Stratmann

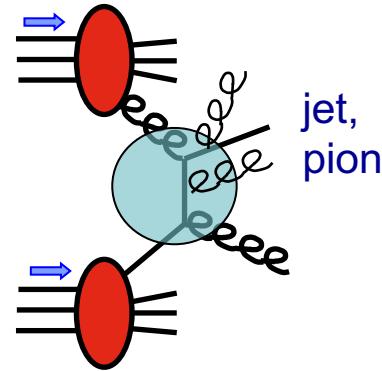
# Helicity PDFs at NNLO: framework



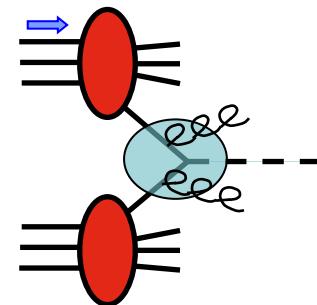
DIS



SIDIS



pp high- $p_T$



W bosons

- Partonic hard scattering:  $\Delta\hat{\sigma}_{ab} = \Delta\hat{\sigma}_{ab}^{\text{LO}} + \frac{\alpha_s}{\pi} \Delta\hat{\sigma}_{ab}^{\text{NLO}} + \left(\frac{\alpha_s}{\pi}\right)^2 \Delta\hat{\sigma}_{ab}^{\text{NNLO}} + \dots$
- PDF evolution:  $\Delta\mathcal{P}_{ij} = \frac{\alpha_s}{2\pi} \Delta P_{ij}^{\text{LO}} + \left(\frac{\alpha_s}{2\pi}\right)^2 \Delta P_{ij}^{\text{NLO}} + \left(\frac{\alpha_s}{2\pi}\right)^3 \Delta P_{ij}^{\text{NNLO}} + \dots$

# First NLO “global” analysis of nucleon helicity (GRSV, 1996)

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## Next-to-leading-order radiative parton model analysis of polarized deep inelastic lepton-nucleon scattering

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(Received 28 August 1995; revised manuscript received 10 January 1996)

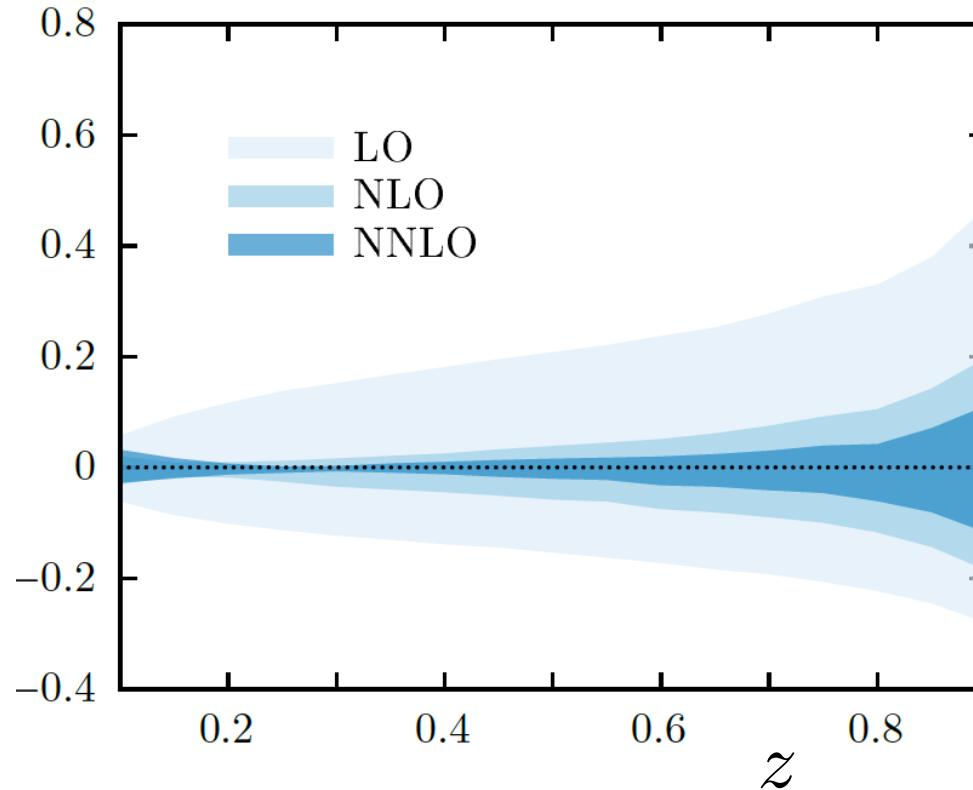
A next-to-leading-order QCD analysis of spin asymmetries and structure functions in polarized deep inelastic lepton-nucleon scattering is presented within the framework of the radiative parton model. A consistent NLO formulation of the  $Q^2$  evolution of polarized parton distributions yields two sets of plausible NLO spin-dependent parton distributions in the conventional  $\overline{\text{MS}}$  factorization scheme. They respect the fundamental positivity constraints down to the low resolution scale  $Q^2 = \mu_{\text{NLO}}^2 = 0.34 \text{ GeV}^2$ . The  $Q^2$  dependence of the spin asymmetries  $A^{p,n,d}(x, Q^2)$  is similar to the leading-order (LO) one in the range  $1 \leq Q^2 \leq 20 \text{ GeV}^2$  and is

# Why NNLO ?

- need per cent accuracy for EIC and JLab (cf. LHC experience)
- reduce theory uncertainty

$$\frac{\sigma(\mu) - \sigma(Q)}{\sigma(Q)}$$

$$Q/2 \leq \mu_{R,F} \leq 2Q$$



$Q^2 > 5 \text{ GeV}^2$

SIDIS@EIC

Abele, De Florian, WV

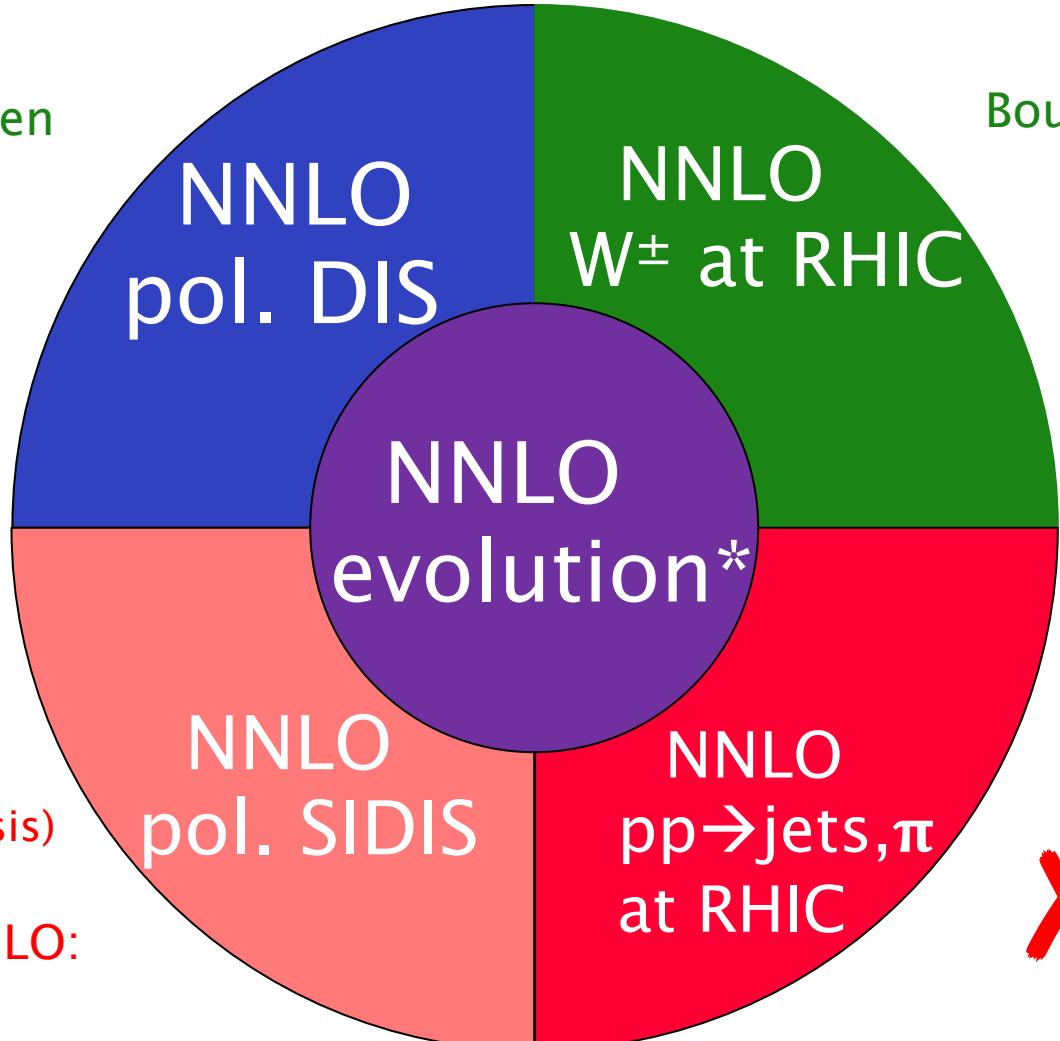
- progress in lattice computations

hard scattering:  $\Delta\hat{\sigma}_{ab} = \Delta\hat{\sigma}_{ab}^{\text{LO}} + \frac{\alpha_s}{\pi} \Delta\hat{\sigma}_{ab}^{\text{NLO}} + \left(\frac{\alpha_s}{\pi}\right)^2 \Delta\hat{\sigma}_{ab}^{\text{NNLO}} + \dots$

evolution:  $\Delta\mathcal{P}_{ij} = \frac{\alpha_s}{2\pi} \Delta P_{ij}^{\text{LO}} + \left(\frac{\alpha_s}{2\pi}\right)^2 \Delta P_{ij}^{\text{NLO}} + \left(\frac{\alpha_s}{2\pi}\right)^3 \Delta P_{ij}^{\text{NNLO}} + \dots$

Zijlstra, van Neerven  
1994

Boughezal, Li, Petriello 2021



Bonino, Gehrmann, Löchner,  
Schönwald, Stagnitto 2024  
Goyal, Lee, Moch, Pathak,  
Rana, Ravindran 2024

(not suitable yet for glob. analysis)

soft-gluon approximate NNLO:

Anderle, Ringer, WV 2012  
Abele, de Florian, WV 2021

X

use soft-gluon approx.

\* Moch, Vogt, Vermaseren  
Blümlein, Marquard,  
Schneider, Schönwald  
**QCD Pegasus:** A. Vogt

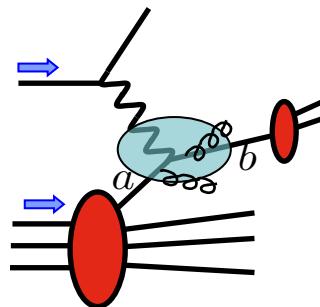
# NNLO helicity PDFs:

- DIS-only analysis Taghavi-Shahri et al., 2016
- DIS and (approximate) SIDIS MAP: Bertone, Chiefa, Nocera 2024
- **Fully global analysis with approximate NNLO for SIDIS and pp**  
BDSSV: Borsa, de Florian, Sassot,Stratmann, WV (to appear)

$$\Delta\hat{\sigma}_{qq}^{\text{N}^k\text{LO}}(\hat{x}, \hat{z}) \sim \alpha_s^k \left[ \delta(1-\hat{x}) \left( \frac{\ln^{2k-1}(1-\hat{z})}{1-\hat{z}} \right)_+ + \delta(1-\hat{z}) \left( \frac{\ln^{2k-1}(1-\hat{x})}{1-\hat{x}} \right)_+ \right.$$


$$\left. + \frac{1}{(1-\hat{x})_+} \left( \frac{\ln^{2k-2}(1-\hat{z})}{1-\hat{z}} \right)_+ + \frac{1}{(1-\hat{z})_+} \left( \frac{\ln^{2k-2}(1-\hat{x})}{1-\hat{x}} \right)_+ + \dots \right]$$

$$\hat{x} = \frac{Q^2}{2p_a \cdot q}$$



- logs can be resummed to all orders: threshold resummation

Anderle, Ringer, WV  
Abele, de Florian, WV

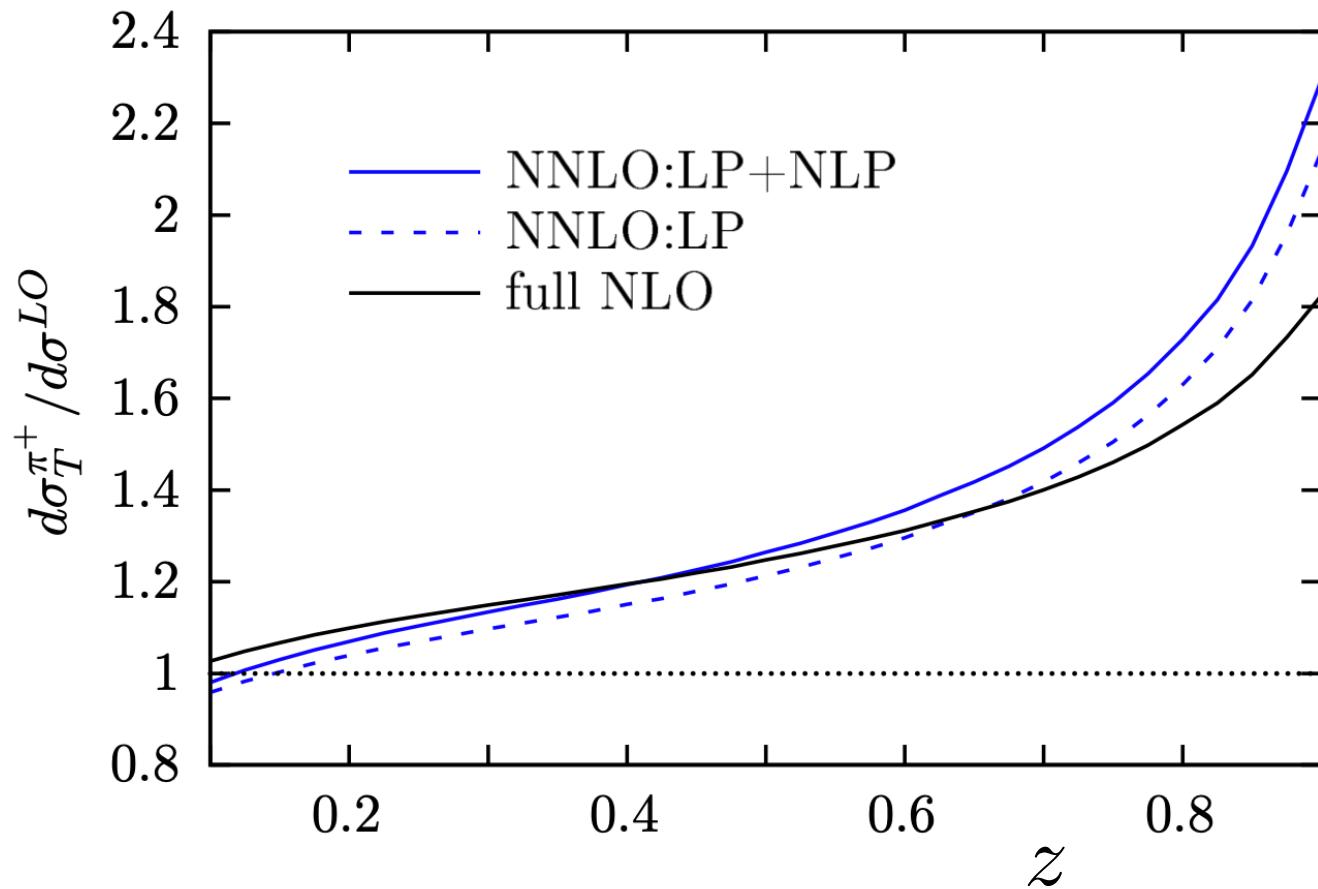
Fixed Order							
Resummation	LO	1					
	NLO	$\alpha_s L^2$	$\alpha_s L$	$\alpha_s$			
NNLO	$\alpha_s^2 L^4$	$\alpha_s^2 L^3$	$\alpha_s^2 L^2$	$\alpha_s^2 L$	$\alpha_s^2$		
...	...	...	...	...	...	...	
N <sup>k</sup> LO	$\alpha_s^k L^{2k}$	$\alpha_s^k L^{2k-1}$	$\alpha_s^k L^{2k-2}$	$\alpha_s^k L^{2k-3}$	$\alpha_s^k L^{2k-4}$		...

↓                    ↓                    ↓

LL                NLL                NNLL

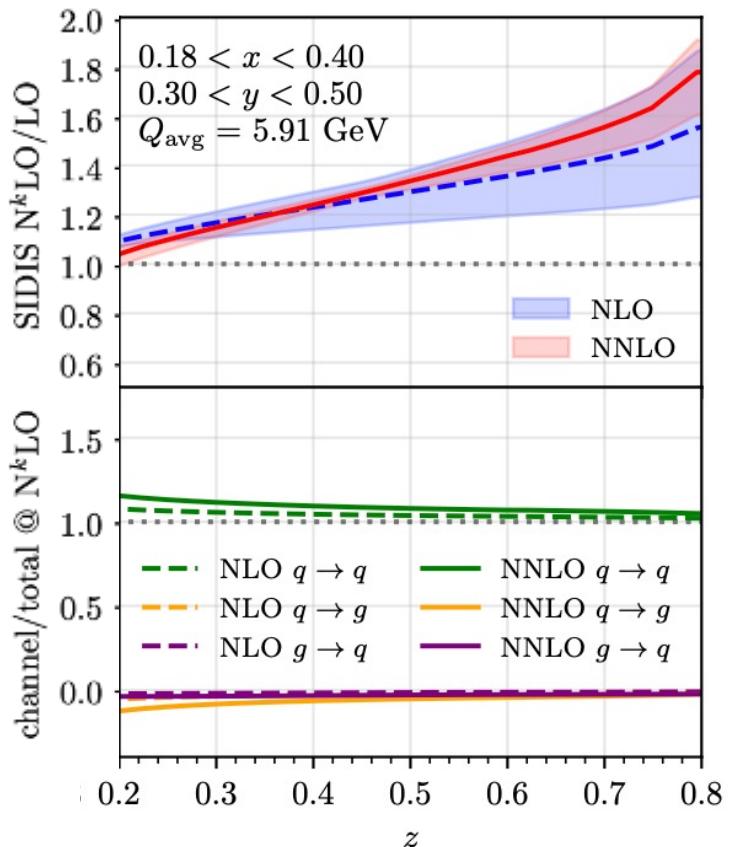
$\mu p \rightarrow \mu \pi^+ X$ 

COMPASS

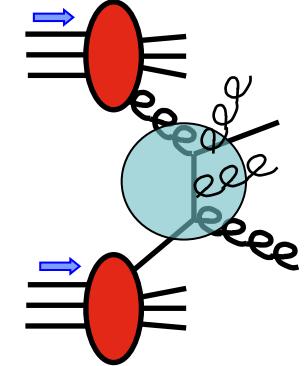
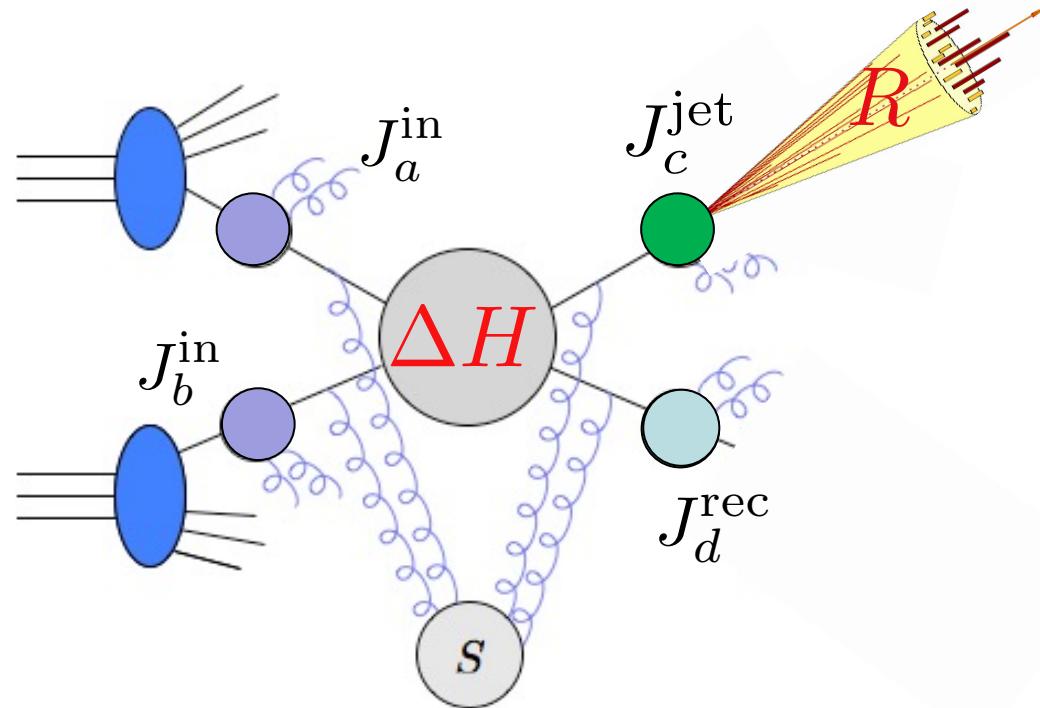
**Full NNLO unpolarized:**

Bonino, Gehrmann, Stagnitto

Goyal, Moch, Pathak, Rana, Ravindran



## Approximate NNLO corrections for $pp \rightarrow \text{jet}+X$ at RHIC:



Kidonakis, Oderda, Sterman  
de Florian, WV  
Hinderer, Ringer, Sterman, WV, ....

$$\text{threshold logs } \ln \left( 1 - \frac{s_{\text{rad}}}{s} \right)$$

$$\Delta \hat{\sigma}^{ab \rightarrow cd} \sim J_a^{\text{in}} \times J_b^{\text{in}} \times J_c^{\text{jet}} \times J_d^{\text{rec}} \times \text{Tr} [\Delta H S^\dagger S S]_{ab \rightarrow cd}$$

# Global NNLO analysis: results

# Parameters & data selection:

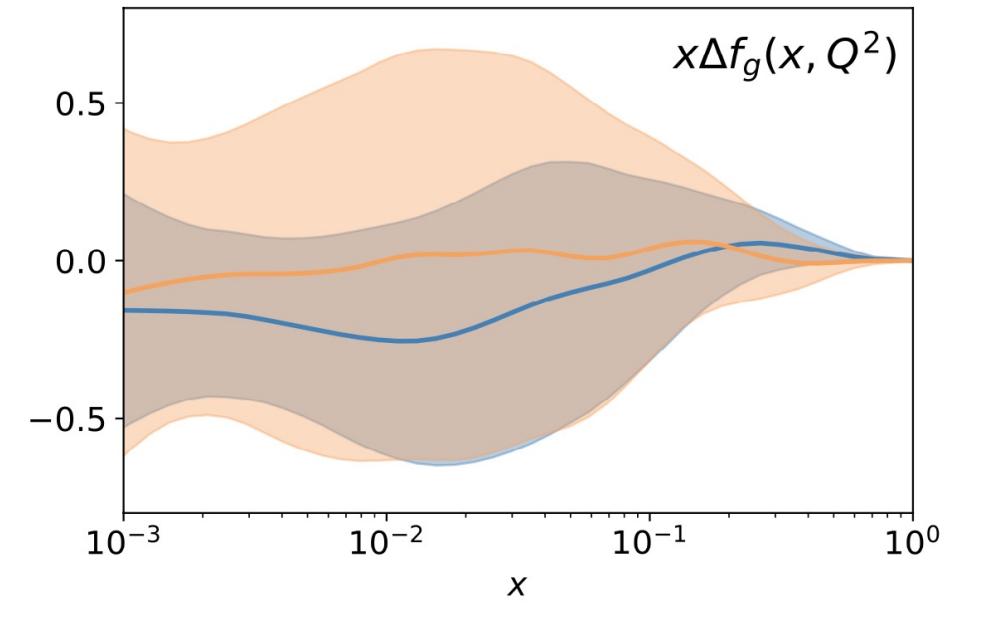
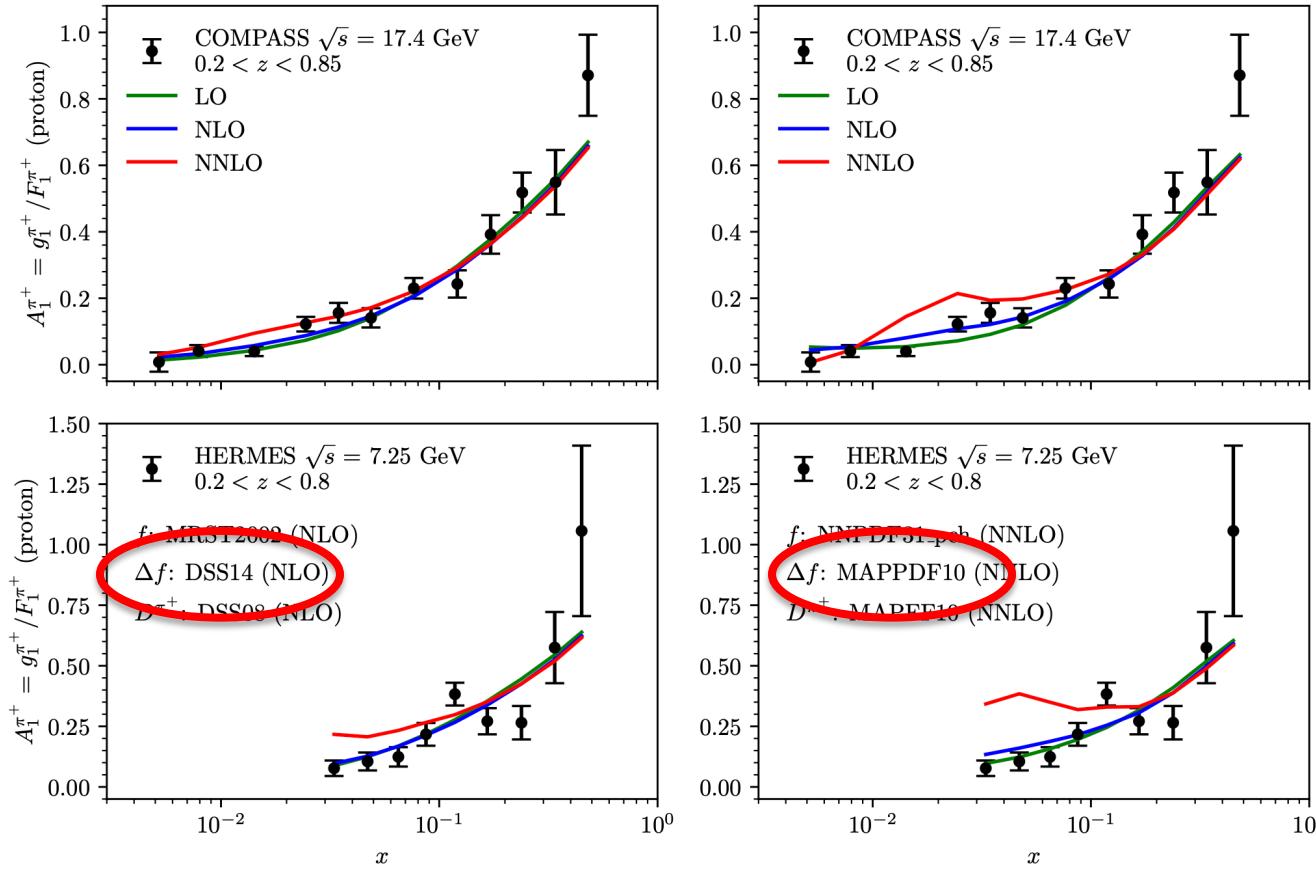
- unpol. PDFs: MSHT20–NNLO
- unpol. FFs: Borsa–de Florian–Sassot–Stratmann NLO 2022,2024  
(BDSSV NNLO fit does not have  $p\bar{p} \rightarrow$  hadron data)
- data:

<b>DIS:</b> EMC,SMC,E142,E143,E154,E155, HERMES, COMPASS, HALL-A,CLAS (p,n,d,He)	378
<b>SIDIS:</b> HERMES, COMPASS ( $p-\pi^\pm, d-\pi^\pm$ )	<b>277</b>
<b>PP-JETS:</b> STAR run 5,6,9,12,13,15 $(\sqrt{s} = 200, 510 \text{ GeV})$ (no dijets yet)	91
<b>PP-<math>\pi^0/\pi^\pm</math>:</b> PHENIX, STAR	82
<b>PP <math>W^\pm</math>:</b> PHENIX, STAR	22
<b>Total:</b>	<b>850</b>

however...

- interesting feature for SIDIS:

Bonino et al.

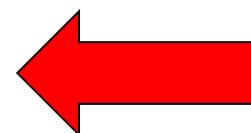


- “bump” not present in threshold approximation for SIDIS  
→ breakdown of approximation?

- in any case, lower-x SIDIS data involve FFs (and unpol. PDFs) outside regimes where they have been determined (and don't work so well)
- we find:

	NLO	NNLO
DIS	304.72	308.74
all SIDIS	276.08	322.46
all pp	199.72	193.91

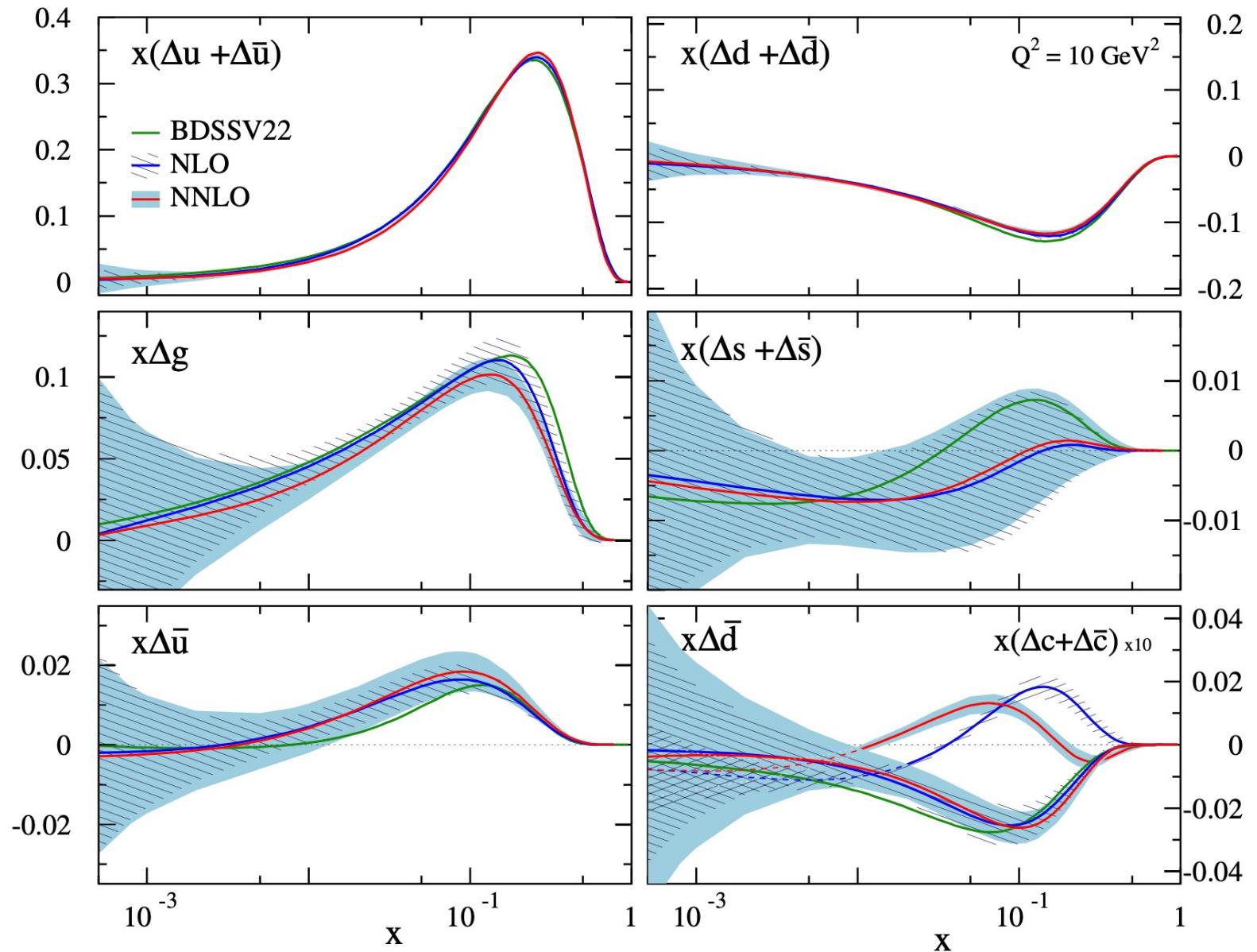
- introducing cut  $x > 0.12$  (and  $p_T^h > 2$  GeV):

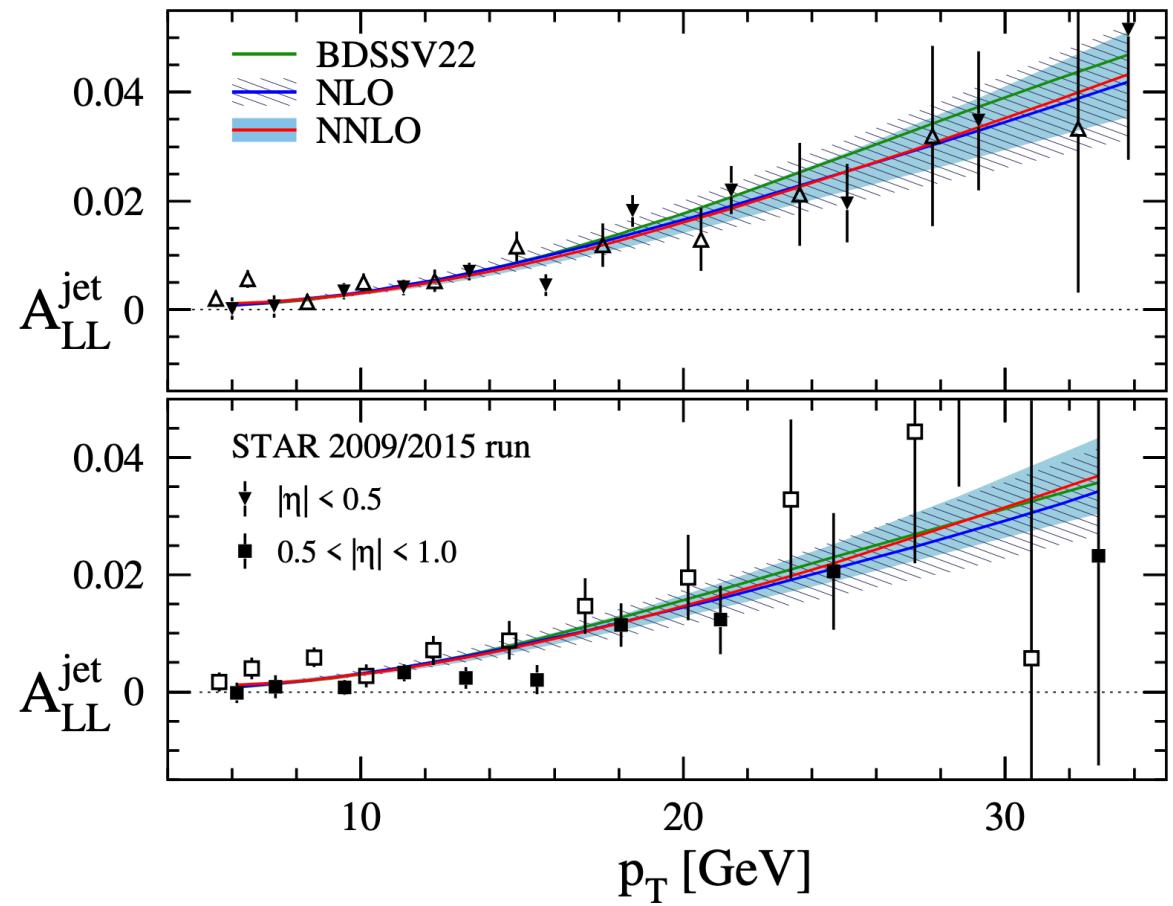
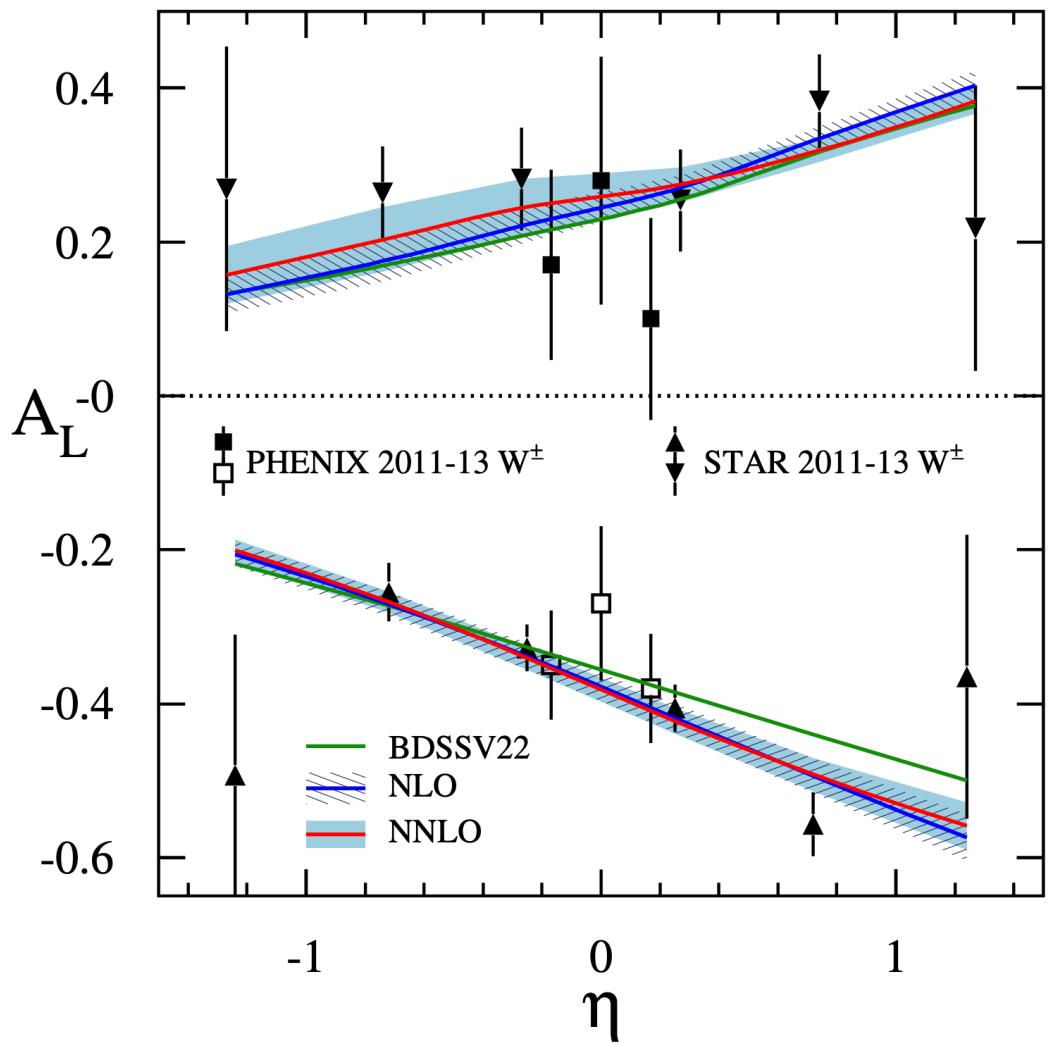


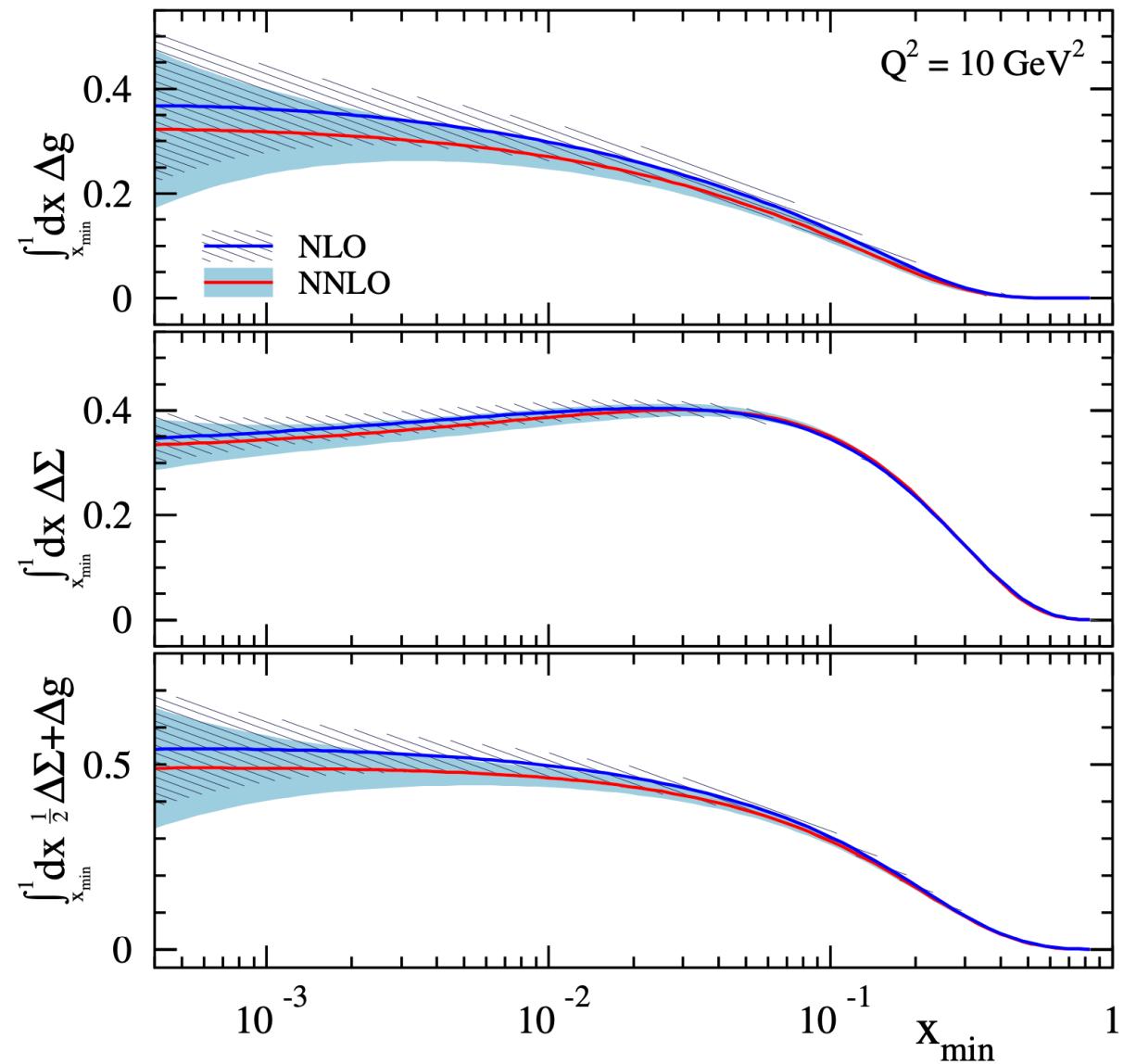
implemented  
in our analysis  
to be conservative

	NLO	NNLO
DIS	302.76	294.48
all SIDIS	126.62	124.42
all pp	196.93	189.90

- similar observation: FF fits and MAP analysis







## Concluding remarks:

- qualitative step forward: first global NNLO analysis of helicity PDFs
- pQCD analysis “in good shape”
- further improvements will come: full SIDIS, improved NNLO for high- $p_T$
- SIDIS crisis?
- numerous outstanding issues: low- $x$ , power corrections, synergies with lattice