

Measurement-induced phase transitions in quantum circuits

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This talk

LP, Y. Li, R. Vasseur, A. Nahum, PRB **107**, 224303 (2023)

M. Fava, LP, T. Swann, D. Bernard, A. Nahum, arXiv:2302.12820
[Accepted PRX]

Main interests

Many-body physics out of equilibrium

Quantum information & quantum simulation

Tensor networks

Quantum circuits & cellular automata

Low-dimensional quantum field theory (Integrability, CFT)

Cold atoms

Quantum chaos & random-Matrix theory

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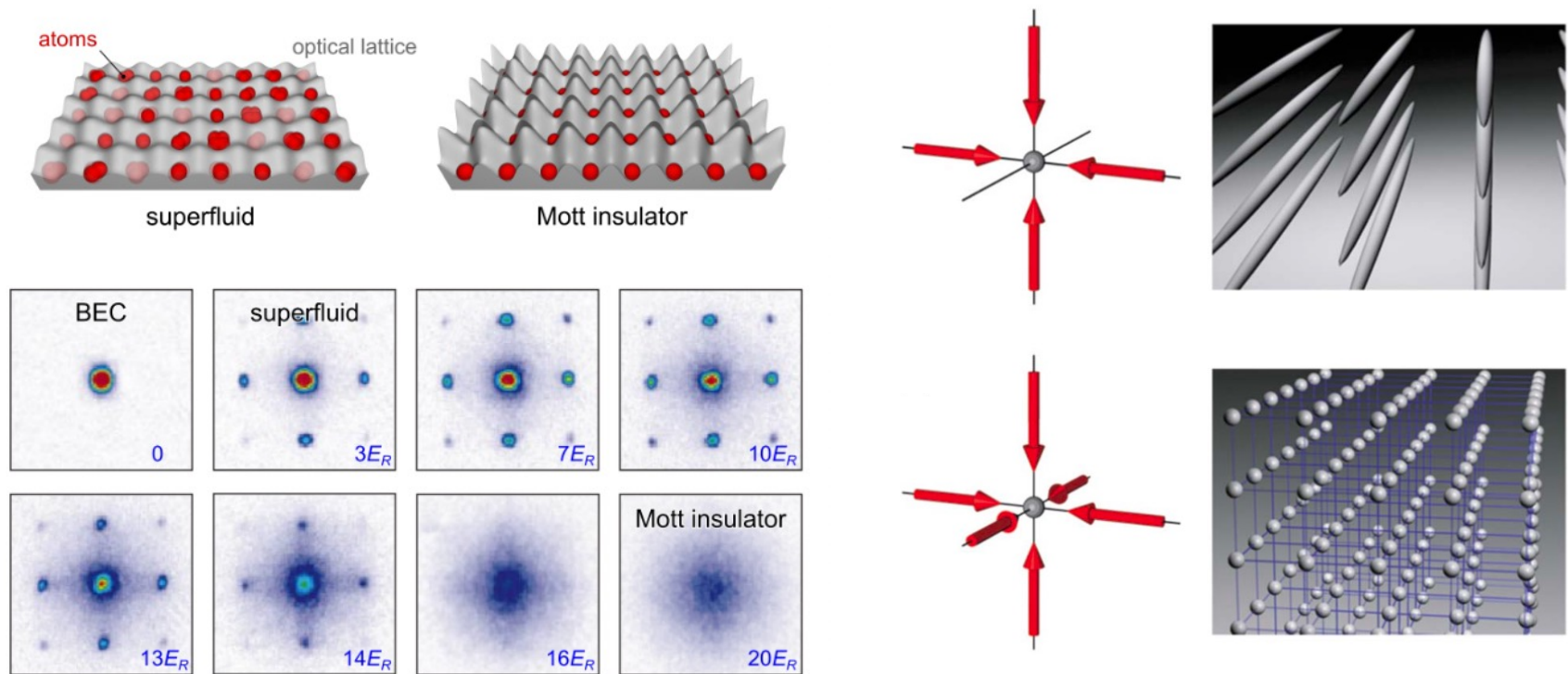
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- Renewed interest from physics of **ultra-cold atoms**:
 - large-number of atoms confined in arbitrary **geometries** and **effective dimensions**
 - ideal **isolation** and **control** of the system parameters
 - real-time dynamics** accessible to high precision



- Non-perturbative problems & hard to study numerically

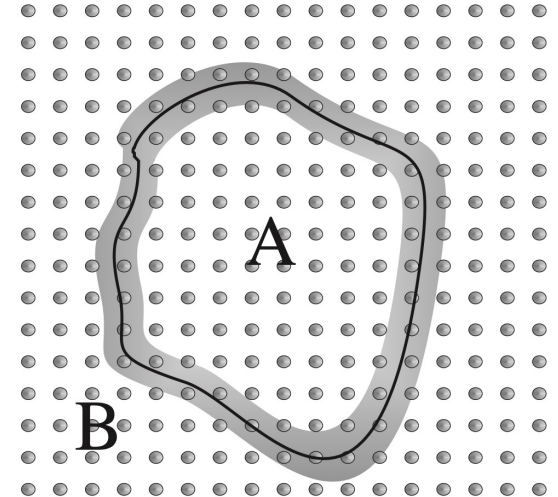
- **Non-perturbative** problems & **hard** to study numerically
- Entanglement emerged as **organizing principle**

- Define **reduced state** $\rho_A = \text{tr}_B [|\psi\rangle\langle\psi|]$

- Entanglement **entropies**

$$S_A^{(n)}[\rho] = \frac{1}{1-n} \log[\text{tr}(\rho^n)]$$

$$S_A[\rho] = -\text{tr}(\rho \log[\rho])$$



$$|\psi\rangle \in \mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$$

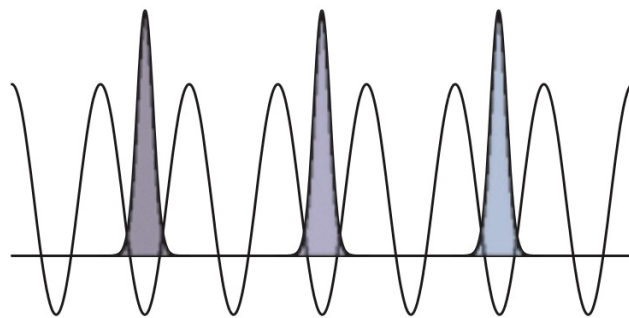
- They quantify **quantum correlations**, and answer the question:
“how far is $|\psi\rangle$ from a (classical) product state”?

$$|\psi_0\rangle = |\phi_0\rangle \otimes \cdots \otimes |\phi_0\rangle$$

- Entanglement key ingredient for **exotic** quantum effects
- It makes quantum dynamics **hard** to simulate classically

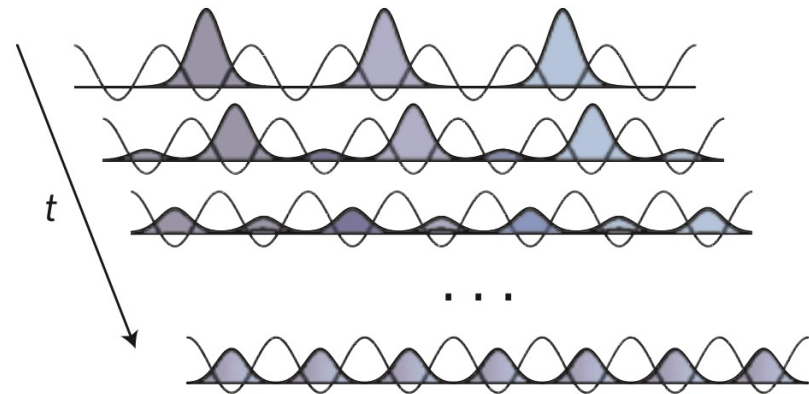
- Entanglement key ingredient for **exotic** quantum effects
- It makes quantum dynamics **hard** to simulate classically
- Ex: “**quantum quench**”

$$\hat{H} = \sum_j \left[-J \left(\hat{a}_j^\dagger \hat{a}_{j+1} + \text{h.c.} \right) + \frac{U}{2} \hat{n}_j (\hat{n}_j - 1) + \frac{K}{2} \hat{n}_j j^2 \right]$$



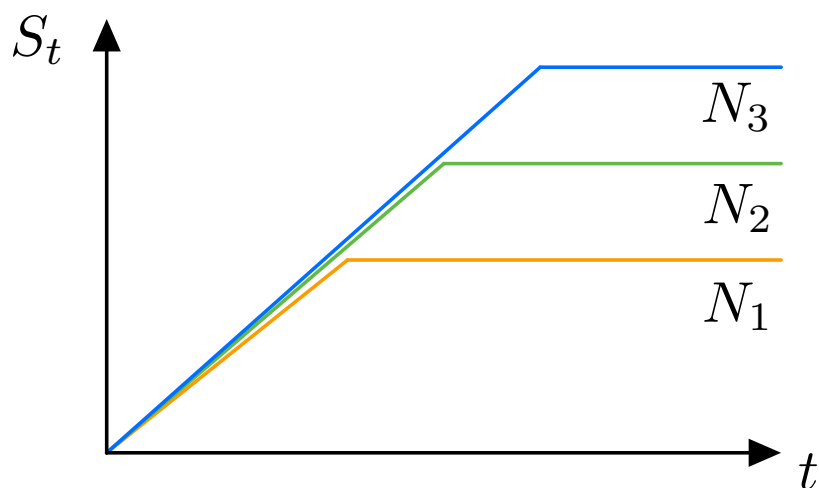
(1) Preparation

$$|\psi(t=0)\rangle = |\dots 10101\dots\rangle$$



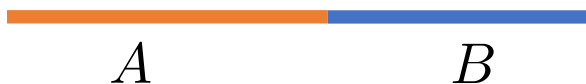
(2) Evolution

- Entanglement key ingredient for **exotic** quantum effects
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$$S(|\Psi\rangle) = -\text{Tr} [\rho_A \log \rho_A]$$

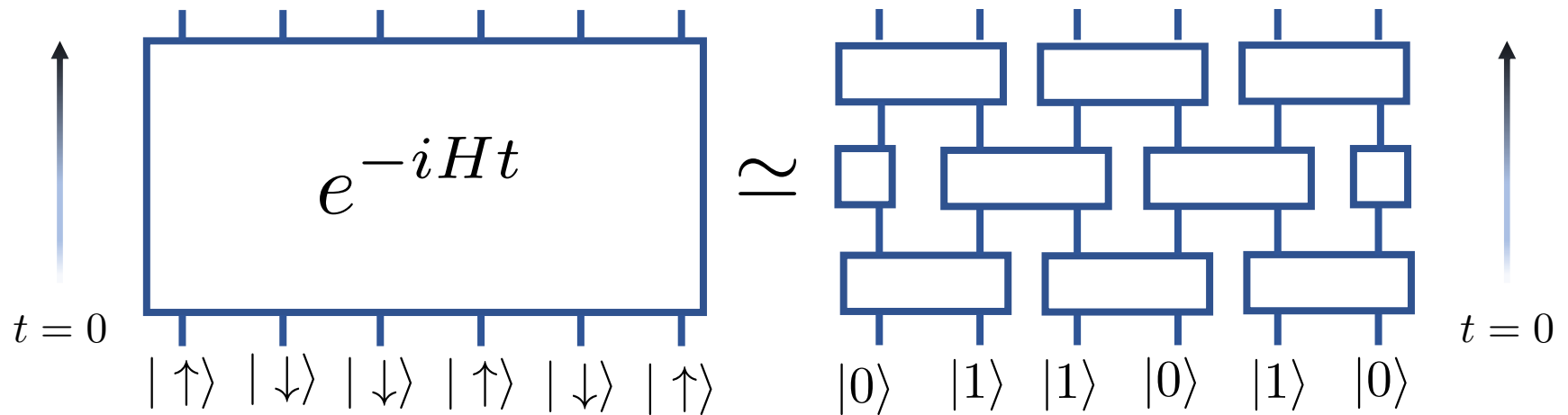
$$\rho_A = \text{Tr}_B [|\Psi\rangle\langle\Psi|]$$



$$|\Psi_0\rangle \xrightarrow{e^{-iHt}} |\Psi_t\rangle$$

- Large-scale dynamics simulable for low entanglement via **tensor-network** methods
- Bottleneck:
“**Entanglement barrier**”

- What if we had a **quantum computer**?
- Direct access to quantum dynamics via **Trotter-Suzuki decomposition**



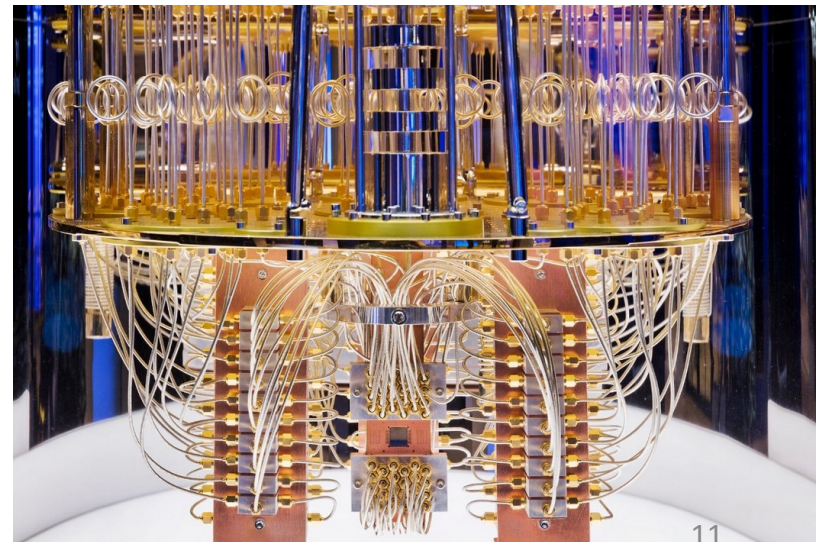
Lloyd (1996)

$$H = \sum_j h_{j,j+1} \Rightarrow e^{-iHt} \simeq \left[e^{-iH_e \delta} e^{-iH_o \delta} \right]^{t/\delta}$$

- Implementation of **Digital Quantum Simulation**

Feynman, 1982

- Remarkable recent progress in platforms for quantum simulation:
cold & Rydberg atoms, trapped ions, superconducting circuits
- Most advanced experiments in **analog quantum simulation**
J. Daley *et al.*, Nature (2022)
- **Digital quantum simulation** \Rightarrow greater control and versatility,
but experimental research is at an **early stage**
- Noisy-Intermediate-Scale-Quantum
(**NISQ**) era
J. Preskill, Quantum 2 (2018)



NISQ physics already opens new avenues for **fundamental** research

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1. Discrete dynamics with new features compared to Hamiltonian evolution:

*What **new phases** and **universal behavior** in NISQ platforms?*

2. **New paradigm** for many-body physics out of equilibrium
 - **Minimally-structured** models to study hard questions in many-body physics
 - Implementation **feasible** in NISQ devices

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1. Discrete dynamics with new features compared to Hamiltonian evolution:

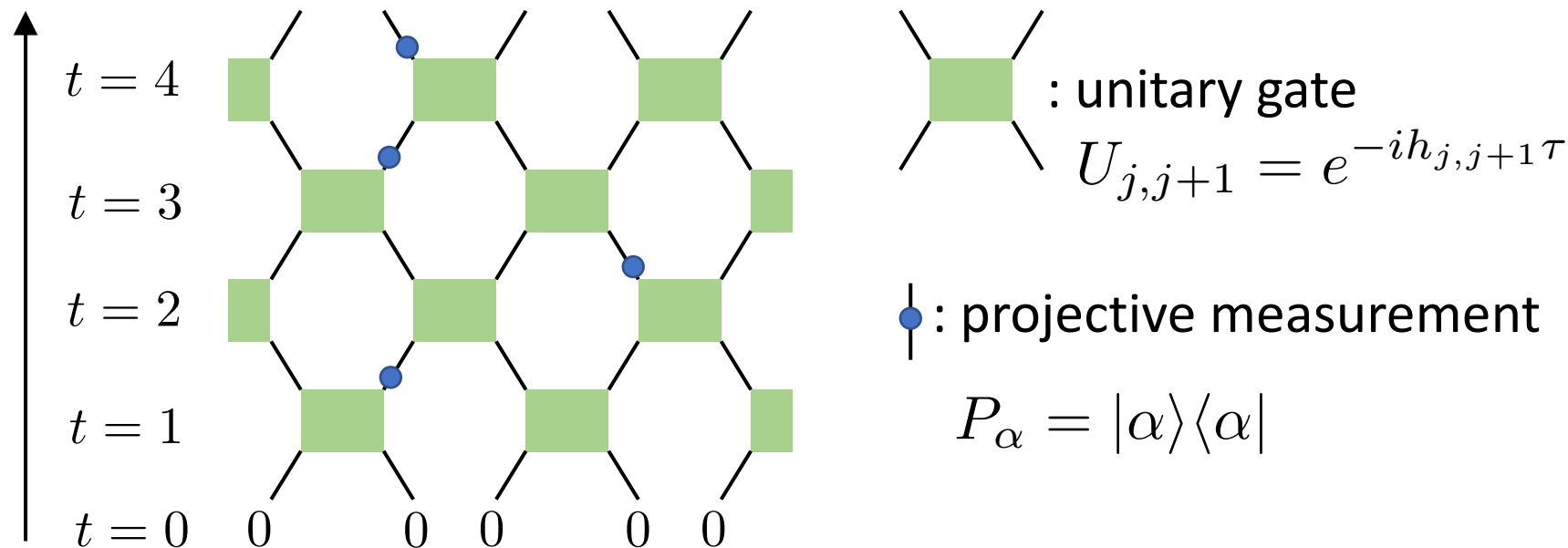
*What **new phases** and **universal behavior** in NISQ platforms?*

2. **New paradigm** for many-body physics out of equilibrium

- **Minimally-structured** models to study hard questions in many-body physics
- Implementation **feasible** in NISQ devices

⇒ **This talk**: focus on class of **stat.-mech. problems** arising in this context

Measurement-induced phase transitions

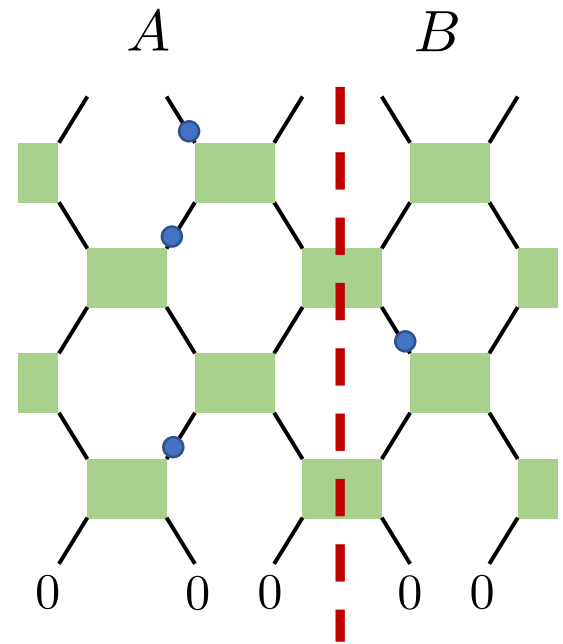


- $\mathcal{H} = \mathbb{C}^q \otimes \dots \otimes \mathbb{C}^q$

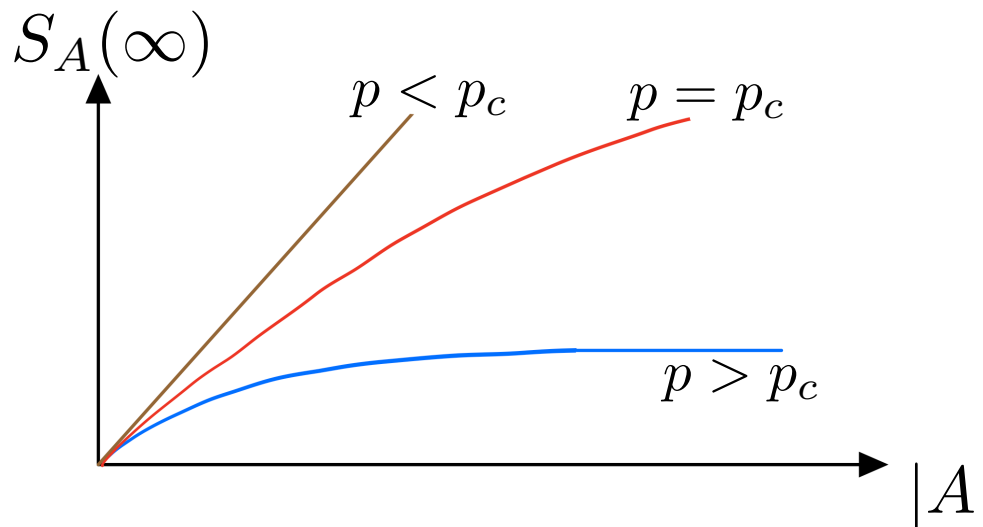
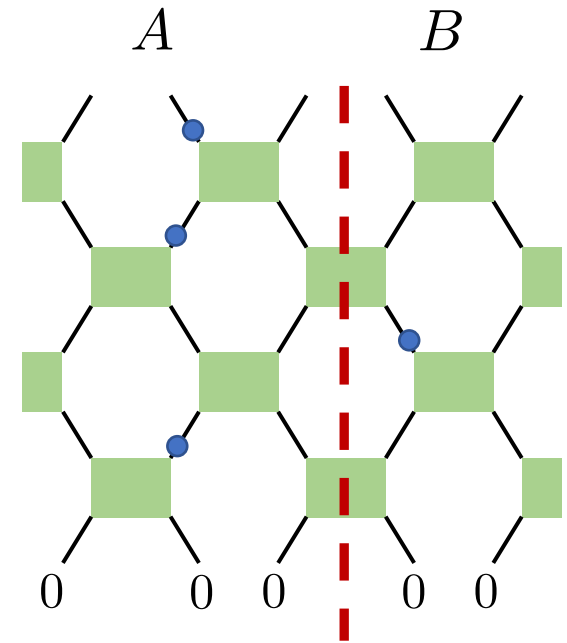
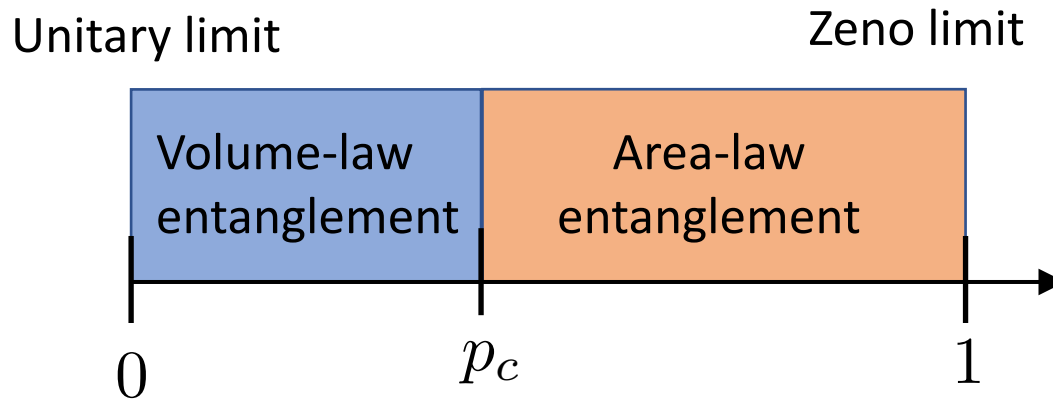
$$\mathbb{C}^{q+1} = \text{span}(\{|0\rangle, |1\rangle, \dots, |q\rangle\})$$

- Born's rule** $\text{prob}(\alpha) = \langle\psi(t)|(|\alpha\rangle\langle\alpha|_j)|\psi(t)\rangle$

Each qubit measured **randomly**, with
a rate (probability) p



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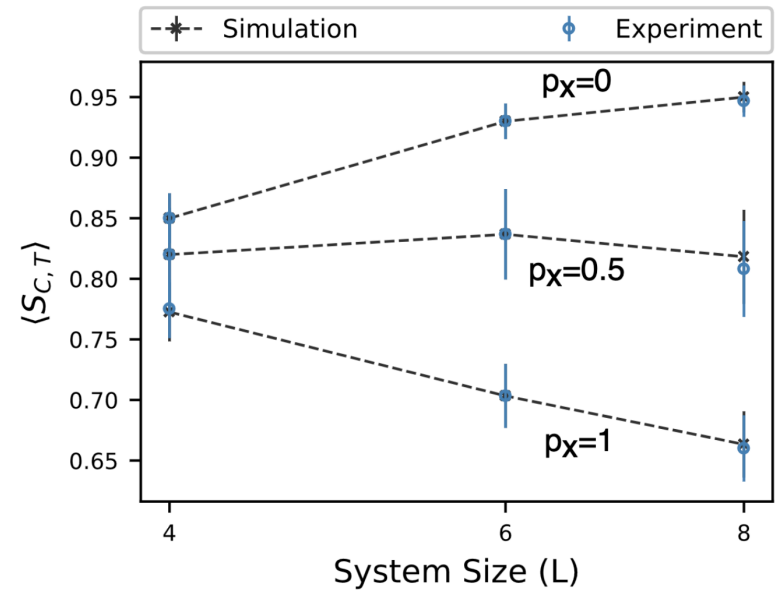
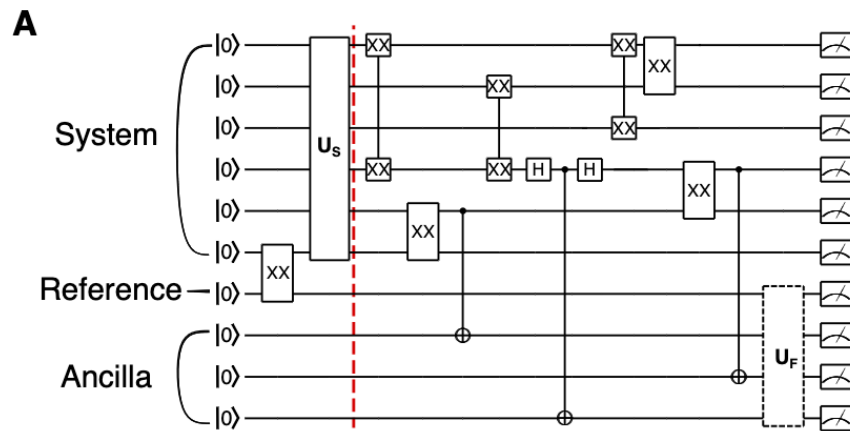


$$S(t) = -\text{Tr} \{ \rho_A(t) \log [\rho_A(t)] \}$$

$$\rho_A(t) = \text{Tr}_B [|\Psi(t)\rangle \langle \Psi(t)|]$$

Skinner, Ruhman, Nahum, PRX (2019)
Li, Chen, Fisher, PRB 98, (2018)

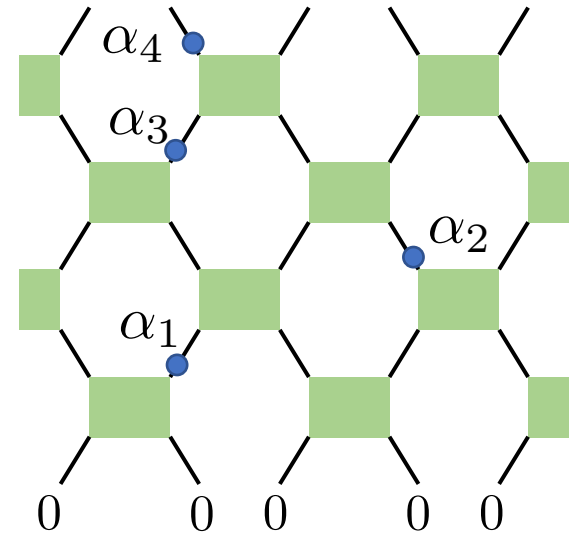
Experimentally probed in point-of-principle experiments



C. Noel et al., Nature Phys. 18, 760 (2022)
 J. Ming Koh et al. arXiv (2022)

- Measurement histories define **quantum trajectories**

$$\mathcal{E} = \{|\psi(t, \{\alpha_j\}_j)\rangle, \text{prob}(\{\alpha_j\}_j)\}$$



- Dynamics of averaged density matrix is **dephasing**

$$\rho \mapsto \mathcal{E}_D^{(j)}[\rho] = \sum_{\alpha} {}_j\langle \alpha | \rho | \alpha \rangle_j \otimes |\alpha\rangle \langle \alpha|_j$$

- Drives locally system towards infinite-temperature
 \Rightarrow **no competition** with unitary dynamics
 \Rightarrow transition not detected by local observables

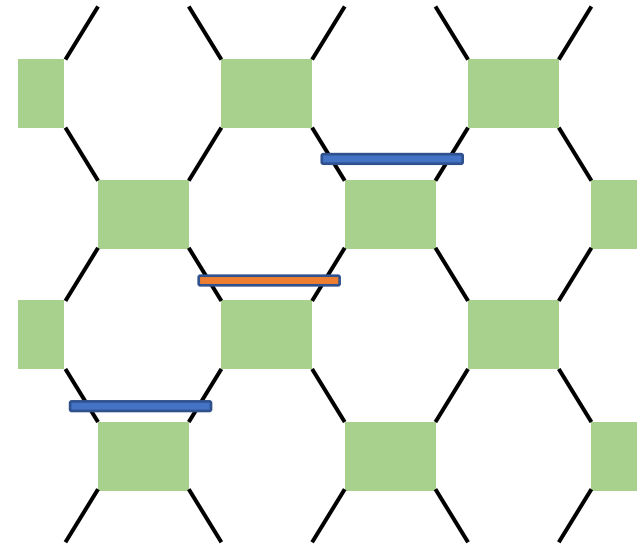
- Example of **new universal behavior** beyond conventional dynamical phase transitions
- Recent work aimed at characterizing phenomenology of the transition (critical exponents, order parameters, etc.)
- This talk:
 - **analytic predictions** for simplest type of transition
[Fava, LP, Swann, Bernard, Nahum, PRX (2023)]
 - **enriched phenomenology** via adaptive measurements
[LP, Li, Vasseur, Nahum, PRB (2023)]

- Non-interacting dynamics

$$\text{Diagram} \propto \exp \left(\sum_{ij} J_{ij}(t) \hat{\gamma}_i \hat{\gamma}_j \right)$$

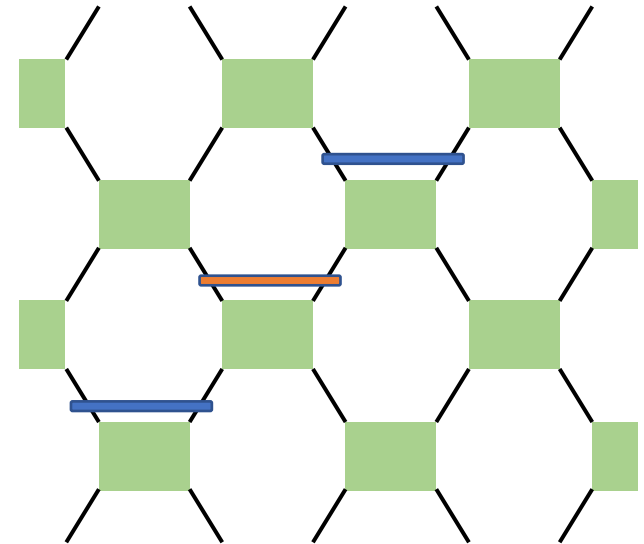
$$\{\hat{\gamma}_i, \hat{\gamma}_j\} = \delta_{ij}$$

$$\mathbb{E} [J_{i,j}(t) J_{k,l}(t')] = J^2 \delta(t - t') \delta_{ik} \delta_{jl}$$

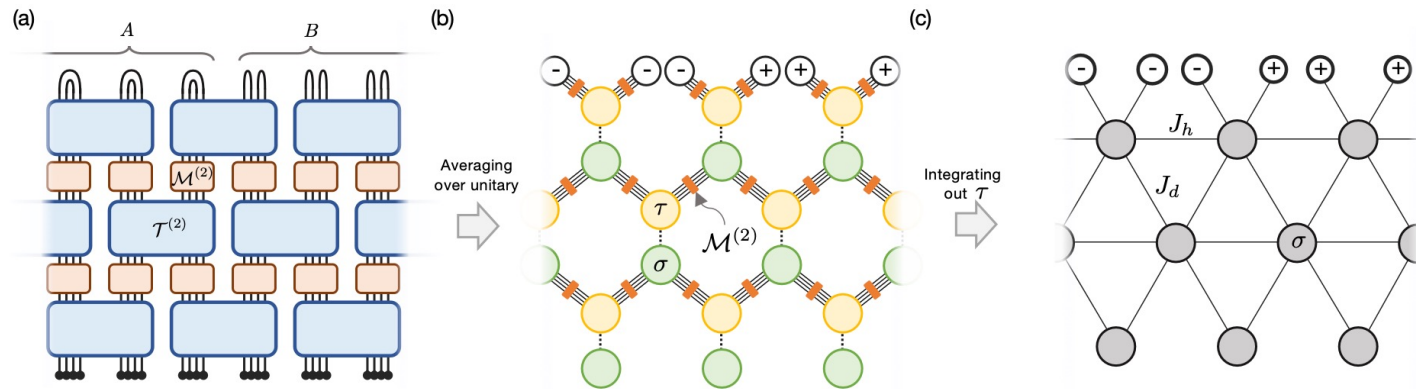


- Measure fermionic local parity $\hat{\gamma}_{2j} \hat{\gamma}_{2j+1}, \hat{\gamma}_{2j+1} \hat{\gamma}_{2j+2}$ with rates $\Gamma_1, \Gamma_2 = \Gamma(1 \pm \Delta)$
- Simple model, but **monitored** dynamics non-trivial
- Numerical evidence suggests entanglement transition from **area-law** phase to **non-trivial phase**

- We provide analytic predictions based on mapping to stat. mech. problem [Bao, et al. Ann. Phys. (2021); Jian, et al. arXiv:2302.09094]
- Basic idea:
 - map space-time pattern to 2D static disordered fermions
 - the problem can be tackled using **replica trick** developed for Anderson localization problem
 - for large space-time scales, develop a field theoretical description in terms of **non-linear sigma model**



- We provide analytic predictions based on mapping to stat. mech. problem [Bao, et al. Ann. Phys. (2021); Jian, et al. arXiv:2302.09094]



- Replica limit arising due to

$$S_{n,A}[|\psi\rangle] = \frac{1}{1-n} \log \text{Tr}_A (\rho_A^n)$$

$$\log x = \lim_{m \rightarrow 0} (x^m - 1) / m$$

- Bipartite **entanglement entropy** mapped to partition function for a matrix field $Q \in SO(N)$

$$\mathcal{S}[Q] = \frac{1}{2g_B} \int dx \, dt \, \text{Tr} \left[\frac{1}{v} \partial_t Q^T \partial_t Q + v \partial_x Q^T \partial_x Q \right]$$

N : number of replicas.

- Same description as **Anderson transition**
- Physical predictions obtained in replica limit $N \rightarrow 1$ because of **Born probabilities**
- We obtain asymptotic exact prediction based on **RG analysis**

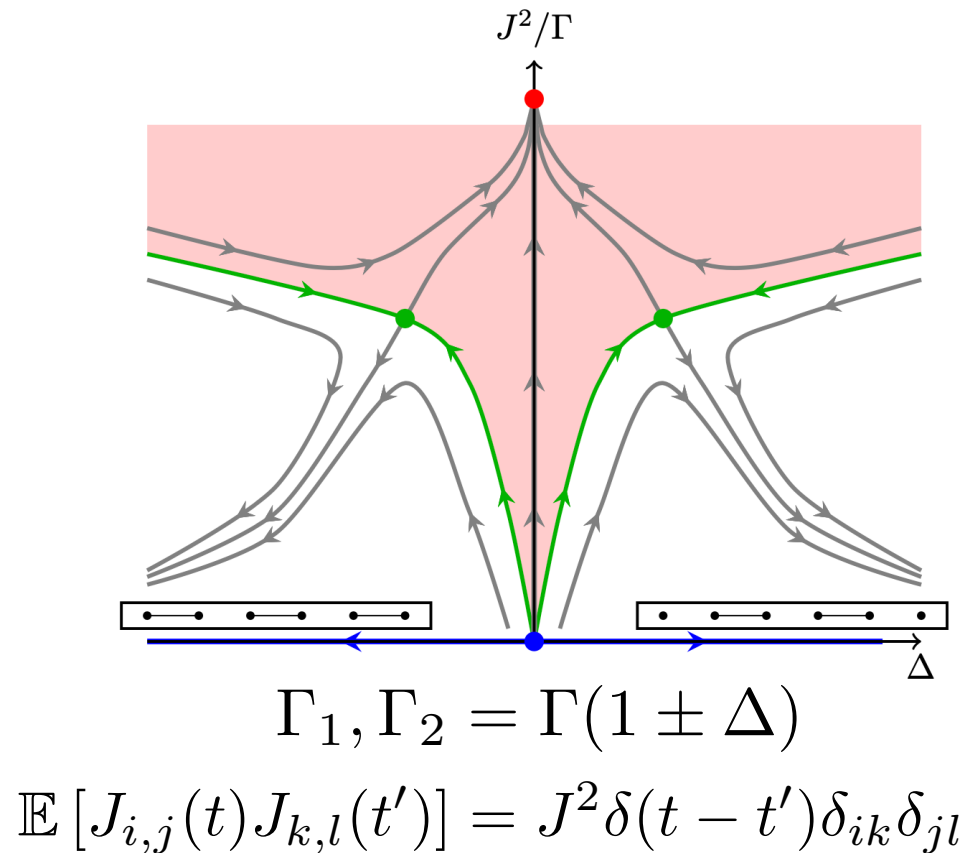
$$\frac{dg}{d \ln L} = \frac{1}{8\pi} (N - 2) g^2 + O(g^3)$$

- Non-trivial phase:

$$S_L(\infty) \sim \frac{1}{48} (\ln L)^2$$

- Trivial phase:
entanglement **area law**

$$S_L(\infty) \sim O(1)$$

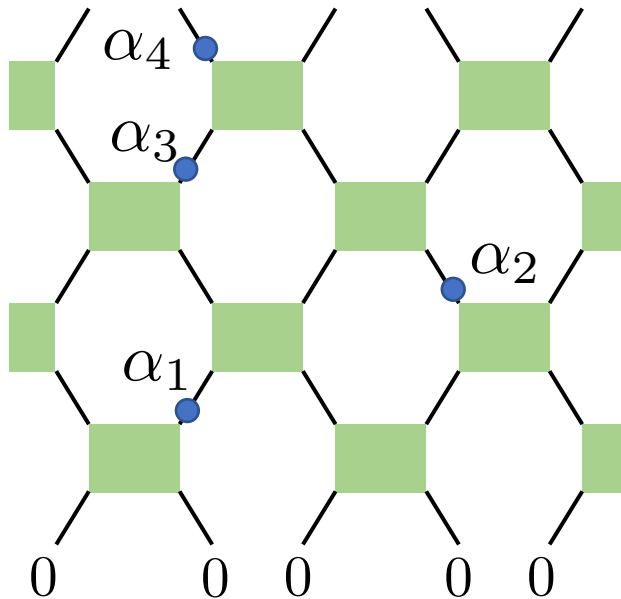


- Transition out of the non-trivial phase driven by **proliferation of vortices**, but different from Kosterlitz–Thouless

\Rightarrow **new** universality class!

- Simplest example of universal behavior made possible by quantum-circuit setting
- Wide phenomenology, depending on unitary operations and measurement process
- Our work: generalize measurements via **feedback** operations, depending on measurement outcome
- Simplest example: local **resetting** measurements
 - measure qubit k
 - if outcome $\alpha = 0$, do nothing
 - if outcome $\alpha \neq 0$, flip the spin

- Resetting described by **Kraus operators** $M_\alpha = |0\rangle\langle\alpha|$



$$\Rightarrow |\psi\rangle \mapsto \frac{M_\alpha |\psi\rangle}{\sqrt{\text{prob}(\alpha)}}$$

$$\text{prob}(\alpha) = \langle\psi|M_\alpha^\dagger M_\alpha|\psi\rangle$$

$$\bullet = \frac{\overline{0}}{\overline{\alpha}}$$

- Choose block-diagonal gates

$$U_{j,j+1} = \begin{pmatrix} 1 & 0 \\ 0 & W \end{pmatrix} \Rightarrow U_{j,j+1} |00\rangle_{j,j+1} = |00\rangle_{j,j+1}$$

- The dynamics features **absorbing-state**

$$|\mathbf{0}\rangle = |0\rangle \otimes |0\rangle \otimes \cdots \otimes |0\rangle$$

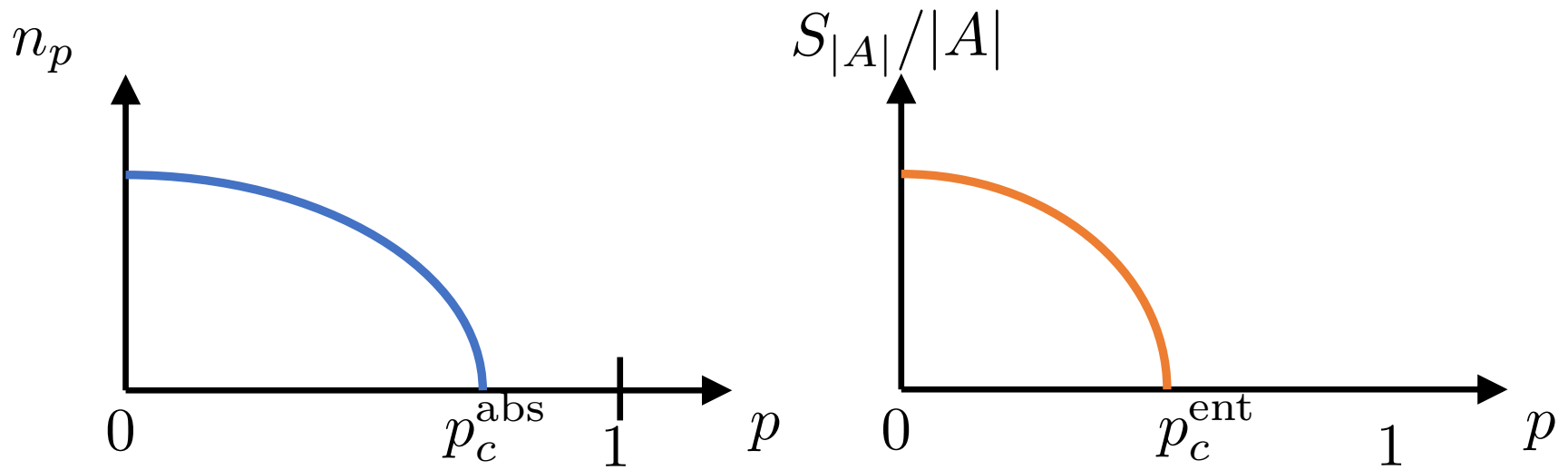
fixed by gates and measurements

- We expect absorbing-state transition detected by order-parameter

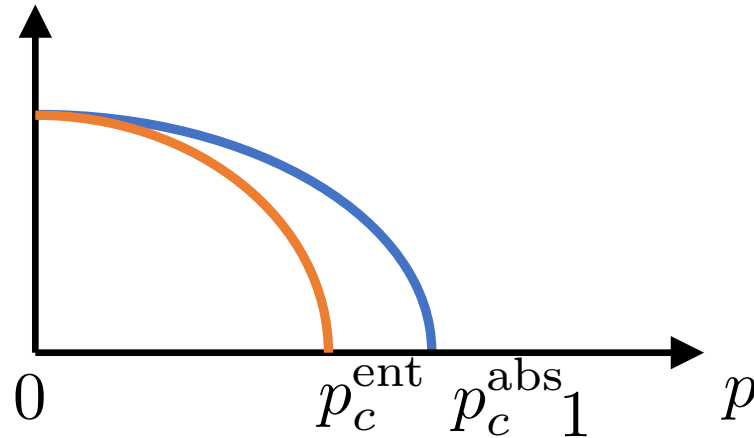
$$n(t) = \frac{1}{L} \sum_{j=1}^L \text{Tr} \{ \mathbb{E}[\rho(t)] \mathcal{P}_j \}$$

$$\mathcal{P}_j = 1 - |0\rangle\langle 0|_j$$

$$\left\{ \begin{array}{ll} p < p_c : \text{active phase} & n_p \equiv \lim_{t \rightarrow \infty} \lim_{L \rightarrow \infty} n(t) > 0 \\ p > p_c : \text{inactive phase} & n_p = 0 \end{array} \right.$$

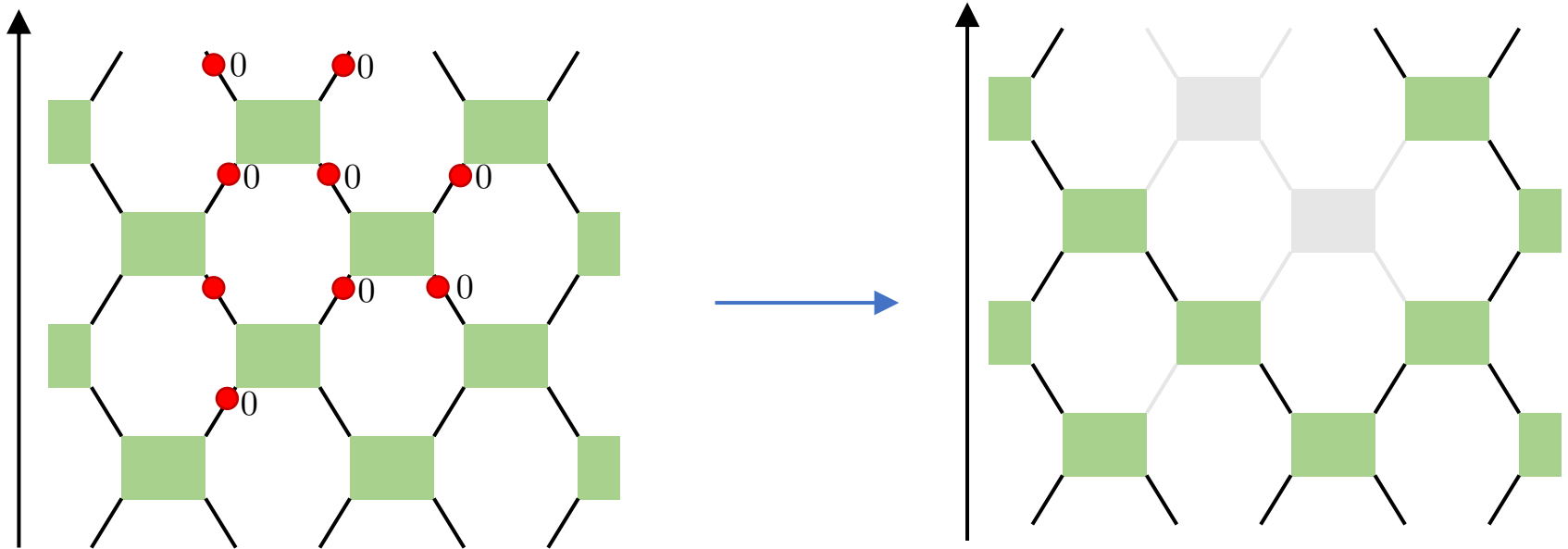


- Clearly $p_c^{\text{ent}} \leq p_c^{\text{abs}}$; can the transitions coincide?
- Does feedback modify the universality class of entanglement MIPT?

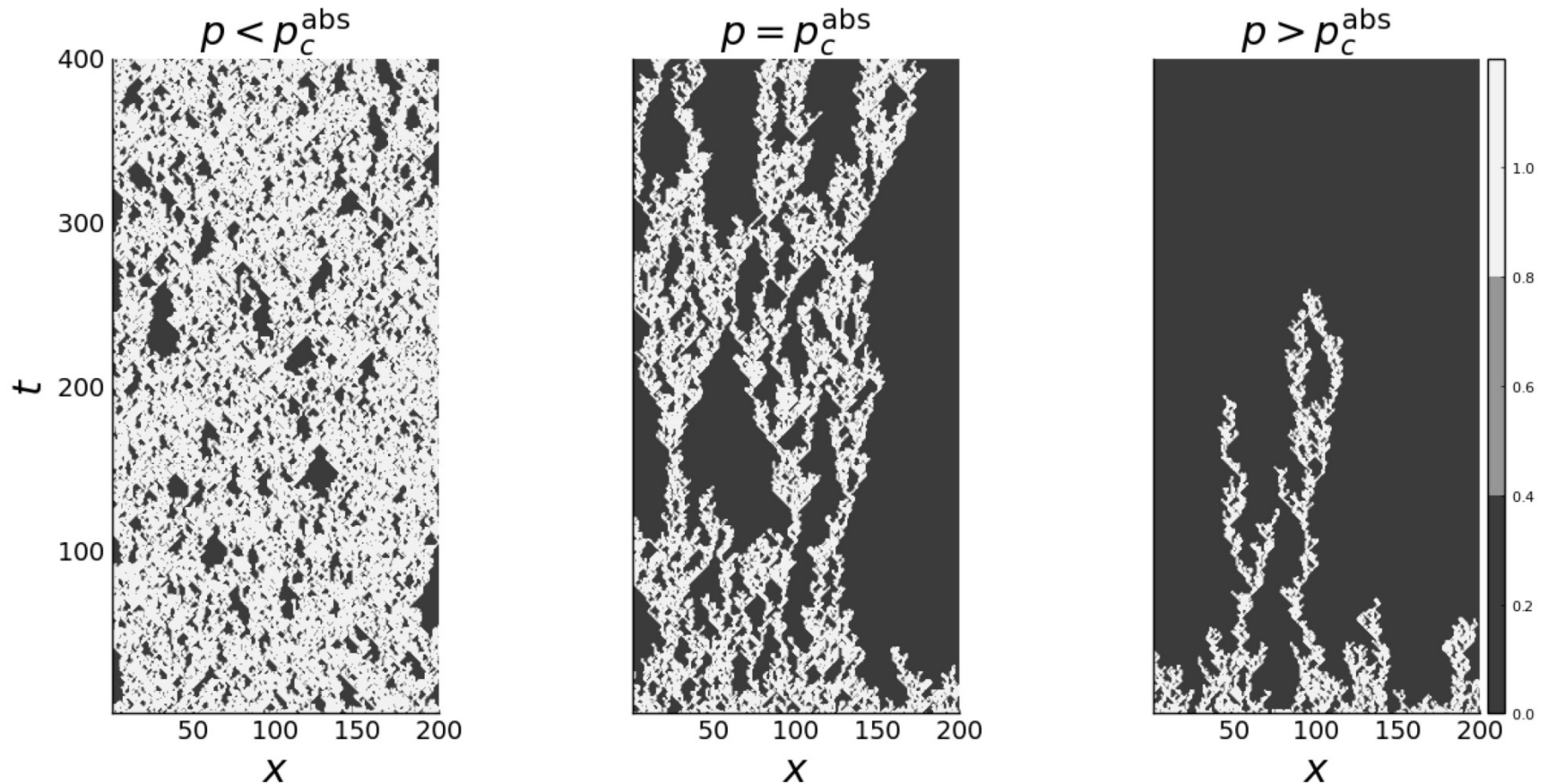


- Clearly $p_c^{\text{ent}} \leq p_c^{\text{abs}}$; can the transitions coincide?
- Does feedback modify the universality class of entanglement MIPT?
- Our work:
 - transitions generically **distinct** and **unrelated**
 - they only coincide in ``**semiclassical limit**''

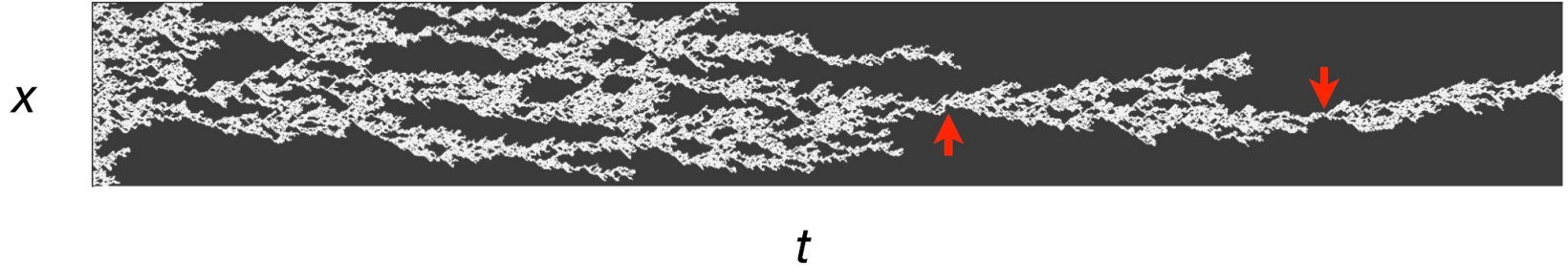
- Basic idea: resetting measurements act as effective cuts in 2D space-time network



- Absorbing-state transition coincides with **connectivity transition**



- **Entanglement** \sim number of links connecting the cluster
- Can be studied exploiting standard results in stat. mech.



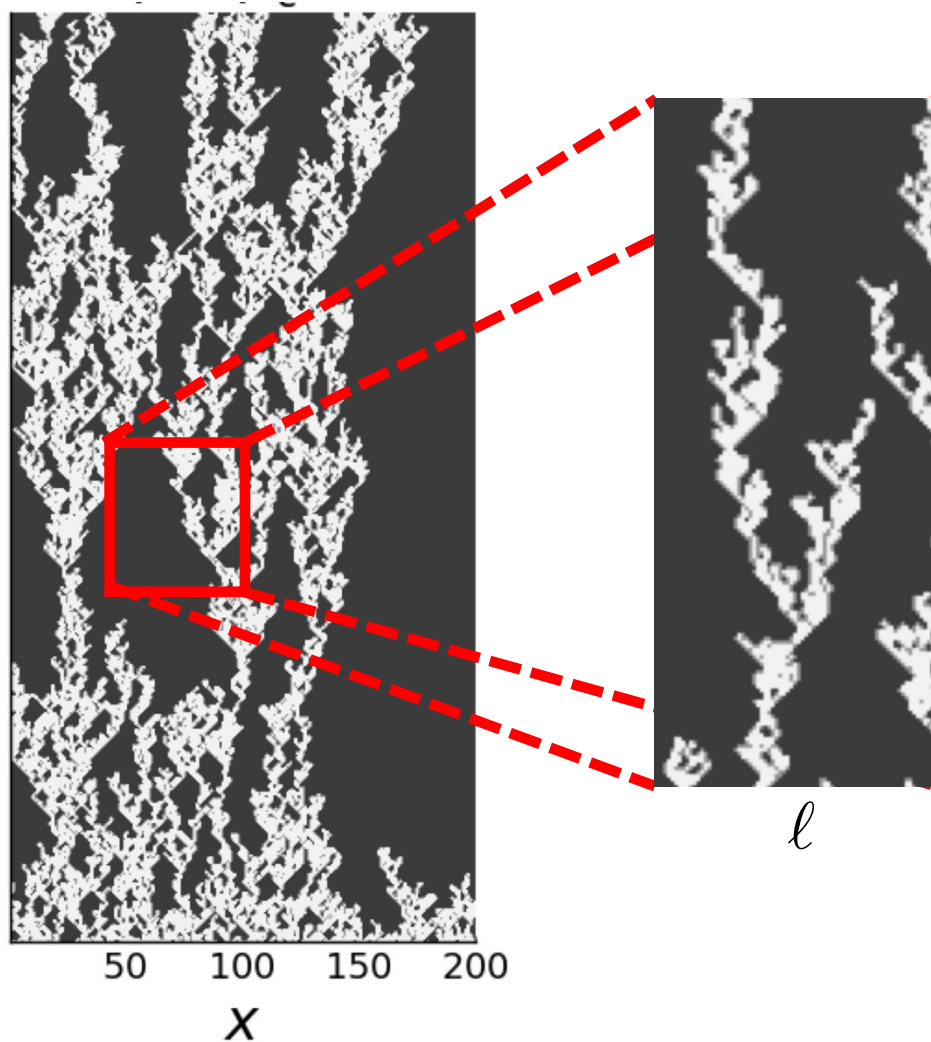
- Define **red bonds** as link that disconnect a percolating cluster
- At criticality

$$N_{\text{red}}(\tau) \sim \tau^{1/\nu_{\parallel}}$$

Huber, Jensen, Sneppen, PRE (1995)

$$\xi_{\perp} \sim |p - p_c^{\text{abs}}|^{-\nu_{\perp}} \quad \nu_{\perp} \simeq 1.097$$

$$\xi_{\parallel} \sim |p - p_c^{\text{abs}}|^{-\nu_{\parallel}} \quad \nu_{\parallel} \simeq 1.73$$



τ

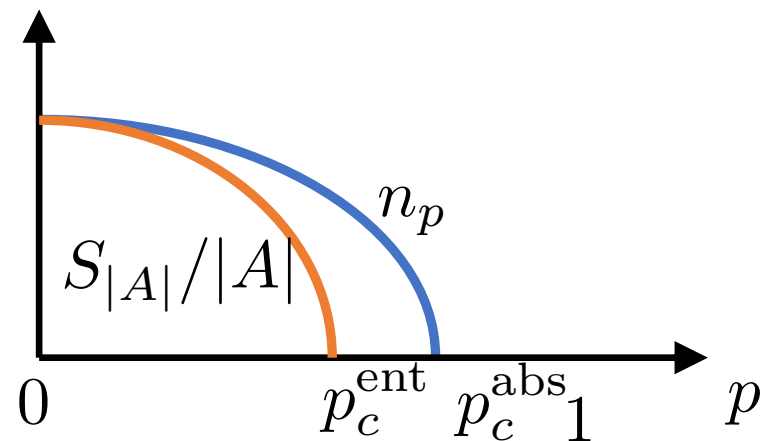
$$N_{\text{red}}(\tau) \sim \tau^{1/\nu_{\parallel}}$$

- Near criticality, we can apply this result for a space-time region of dimensions

$$\ell = \xi_{\perp} \sim |p - p_c^{\text{abs}}|^{-\nu_{\perp}} \quad \tau = \xi_{\parallel} \sim |p - p_c^{\text{abs}}|^{-\nu_{\parallel}}$$

- Using simple estimates: if

$$|p - p_c^{\text{abs}}| \ll 1/q,$$



the system is in the **area-law** phase

⇒ close to DP transition, quantum trajectories are **disentangled**

- We recover that two transitions only coincide in the limit $q \rightarrow \infty$
- Results supported by extensive numerics in simple models

Outlook

- Entanglement phase transitions as new instances of universality out of equilibrium
- Many **qualitative** and **quantitative** questions remain open:
 - Characterization of the transition in **interacting models**
 - Can we efficiently probe the transition?
- More generally, many fundamentally interesting questions are emerging from “digital quantum physics”
- Important new synergies between groups in quantum-information theory, stat.-mech. and many-body physics

My project **QUANTHEM** awarded an ERC starting grant!

Research areas:

- *Quantum information theory*
- *Quantum simulation*
- *Many-body physics out of equilibrium*



European Research Council

Established by the European Commission

New postdoc positions at the University of Bologna starting 2024!



QUANTHEM hosted by and interacting with **Quantum group** in Bologna: E. Ercolessi, P. Pieri, T. Calarco

Thank you for your attention!