Neutrinos as a Window Beyond the SM

Enrique Fernández-Martínez



HIDDE Hunting Invisibles: Dark sectors, Dark matter and Neutrinos Asymmetry Essential Asymmetries of Nature

v masses require BSM physics

See talk by Mikhail Shaposhnikov

v masses require BSM physics Simplest option to add v_R to the SM content

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ν masses require BSM physics Simplest option to add ν_R to the SM content



See posters by Xaver Stribl, Alexandre Sousa, Kevin Urquia, Matheus Hostert, Julia Book Motzkin

Probed in fixed target including ND of oscillation experiments: NuTeV, T2K, NA62, ProtoDUNE, SHIP, DUNE, ICARUS, SBND, μ**BooNE**... Or from atmospheric: SK, IceCube, HK ESSvSB, INO-ICAL, KM3NeT-ORCA,... Also in nuclear decay kinematics: KATRIN/Tristan, HUNTER... Collider searches: ATLAS, CMS, Faser, Belle II...

v masses require BSM physics Simplest option to add v_R to the SM content



If they are too heavy to be produced: indirect searches from PMNS non-unitarity: electroweak precision and flavour observables

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v masses require BSM physics Simplest option to add v_R to the SM content



If they are very light they participate in oscillations See talks at sessions 2,14 and 15 and posters by Alexandre Sousa, Tetiana Kozynets Adapted from: P. Coloma, L. Koerner, I. Shoemaker and J. Yu Snowmass report see arXiv:2209.10362 for summary and links to dedicated analyses

v masses require BSM physics Simplest option to add v_R to the SM content v mass, BAU v mass, BAU

non-Unitarity EWPO, flavour oscillations...
HNLs
sterile v
fixed target,
colliders...

If they are very light they participate in oscillations

See talks at sessions 2,14 and 15 and poster by Alexandre Sousa, Tetiana Kozynets

But also "zero distance" effect in averaged-out oscillations: solar, reactors, MINOS/MINOS+, $NO_{V}A$, T2K, IceCube, HK, ESS_vSB, INO-ICAL, KM3NeT-ORCA, DUNE, JUNO, TAO, SUPERCHOOZ, CLOUD...

Possible connections to other open problems: LED may address the hierarchy problem and ν masses



Similar pheno to steriles but with characteristic masses and mixings: solar, reactors, MINOS/MINOS+, NOvA, T2K, IceCube, HK, ESSvSB, INO-ICAL, KM3NeT-ORCA, DUNE, JUNO/TAO, SUPERCHOOZ/CLOUD...

Also searches for non-standard ν properties: NSI: affect oscillations solars, MINOS/MINOS+, NOvA, T2K, IceCube, HK, ESSvSB, INO-ICAL, KM3NeT-ORCA, DUNE IsoDAR... and directly probed through CEvNs: COHERENT, CONNIE, CONUS...



Also searches for non-standard ν properties: Longer range forces or interactions with DM \rightarrow modified matter potentials Self-interactions \rightarrow impact cosmological abundance and distort SN fluxes See talk by Stefano Gariazzo, posters by Diyaselis Delgado, Caroline Fengler



see arXiv:2209.10362 for summary and links to dedicated analyses

Also searches for non-standard ν properties: Neutrino decay or decoherence would also impact oscillations: solar, MINOS/MINOS+, NO_vA, T2K, IceCube, HK, ESS_vSB, INO-ICAL, KM3NeT-ORCA, DUNE, JUNO,... See posters by Anil Kumar, George Parker, Alicia Pérez García, Anthony Mard Calatayud Cadenillas







See posters by Gray Putnam, Daniel Mayer, Austin Schneider, Matheus Hostert



See posters by Lukas Hennig, Supriya Pan, Alicia Pérez García



See posters by Natsumi Taniuchi, Joshua Barrow, Daisy Kalra, Roxanne Guenette, Cailian Jiang

Simplest option add v_R and acquire Dirac masses via Yukawas

$$Y_{\nu}\bar{\nu}_{R}\phi\nu_{L} \xrightarrow{\text{SSB}} \frac{Y_{\nu}\nu}{\langle\phi\rangle} = \frac{Y_{f}\nu}{\sqrt{2}} \quad \frac{Y_{\nu}\nu}{\sqrt{2}}\bar{\nu}_{R}\nu_{L} \quad m_{D} = \frac{Y_{\nu}\nu}{\sqrt{2}}$$

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but a Majorana mass is also allowed...

 $M_N \bar{\nu}_R \nu_R^c$

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Fermion number violation \rightarrow Baryogenesis via Leptogenesis

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The first mass scale not related to the EW scale and the Higgs To be searched for at experiments!!

A new physics scale



 M_N could be anywhere...

Very different phenomenology at different scales

A new physics scale

eV	keV	MeV	GeV	TeV 📏
				Precision electroweak and flavour violation

Looking for v_R : Non-Unitarity

$$U^{t}\begin{pmatrix} 0 & m_{D}^{t} \\ m_{D} & M_{N} \end{pmatrix}U \approx \begin{pmatrix} N^{t} & -\Theta^{*} \\ \Theta^{t} & X^{t} \end{pmatrix}\begin{pmatrix} 0 & m_{D}^{t} \\ m_{D} & M_{N} \end{pmatrix}\begin{pmatrix} N & \Theta \\ -\Theta^{\dagger} & X \end{pmatrix} = \begin{pmatrix} m & 0 \\ 0 & M \end{pmatrix}$$

The 3×3 submatrix *N* of active neutrinos will not be unitary



 $\sum_{i} \sum_{j=1}^{V_i} (N^{\dagger}N)_{ij}$

Effects in weak interactions...

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Effects in weak interactions...

When the W and Z are integrated out to obtain the Fermi theory NSI are recovered!

see e.g. M. Blennow, P.Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon arXiv:1609.08637 for the dictionary

Just replace *U* by *N*

$$P_{\alpha\beta}(L) = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-\Delta m_{ij}^2 L}{2E}}$$

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At L=0, $P_{\alpha\beta} \neq \delta_{\alpha\beta}$ this "zero distance effect" can be striking and is usually the source of the most stringent constraints

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For instance, if the prediction for $P_{\mu e}$ comes from near detector data on $P_{\mu\mu}$: $\hat{P}_{\mu e}(L) = \frac{P_{\mu e}(L)}{P_{\mu\mu}(0)} = \frac{\sum_{i,j} N_{ei} N_{\mu i}^* N_{\mu j} N_{ej}^* e^{\frac{-\Delta m_{ij}^2 L}{2E}}}{\left| \left(N N^{\dagger} \right)_{\mu \mu} \right|^2}$

Also, no zero distance effect in disappearance channles!! $\widehat{P}_{\mu\mu}(L) = \frac{P_{\mu\mu}(L)}{P_{\mu\mu}(0)} = \frac{\sum_{i,j} N_{\mu i} N_{\mu i}^* N_{\mu j} N_{\mu j}^* e^{\frac{-\Delta m_{ij}^2 L}{2E}}}{\left| \left(N N^{\dagger} \right)_{\mu\mu} \right|^2}$

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These are often thought to be the strongest bounds, but the effect cancels (together with the systematics) when using actual data involving v to predict the unoscillated events

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And data involving neutrinos is always necessary: If I know π flux from hadroproduction need $\pi \rightarrow \mu \nu$ Br Even if computing from "first principles" need $G_F(\mu \text{ decay})$ and V_{ud} (β decay)

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These are often thought to be the strongest bounds, but the effect cancels (together with the systematics) when using actual data involving v to predict the unoscillated events

At most, if the prediction comes from a different channel, one may constrain the ratio

$$\frac{\left|\left(NN^{\dagger}\right)_{\alpha\alpha}\right|^{2}}{\left|\left(NN^{\dagger}\right)_{\beta\beta}\right|^{2}}$$

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But these are more efficiently constraint from LFU bounds, from instance π decay ratios, no need to also detect the v...
G_F from μ decay is affected!



 G_F from μ decay is affected!



But this agrees at ~10⁻³ with G_F from M_W (modulo CDF), measurents of $\sin \theta_W$ from LEP, Tevatron and LHC and β and K decays

ratios:

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LFU also strong bounds on

 $\frac{\left(NN^{\dagger}\right)_{\alpha\alpha}}{\left(NN^{\dagger}\right)_{\beta\beta}}$

From ratios of π , *K*, and lepton decays

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Also the invisible width of the Z since NC are also affected

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And LFV processes such as $\mu \rightarrow e \gamma$ since the GIM cancellation is lost

Bounds from a global fit to flavour and Electroweak precision data

		~	<u> </u>	
	$(1 - \alpha_{ee})$	0	0	
N =	$-\alpha_{\mu e}$	$1 - \alpha_{\mu\mu}$	0	
	$\langle -\alpha_{\tau e} \rangle$	$-\alpha_{\tau\mu}$	$1 - \alpha_{\tau\tau}$	/
	Zz. Xing ()709.2220 a	nd 1110.00	83.
F. J. E	Scrihuela, D). V. Forero,	O. G. Miran	da, 70
111		J. W. I. Vall		/ 5.
M. B	Blennow, E	EFM, J. He	ernandez-	•
Nare	do-Tuero	rcano ano	D. Dez-Pavo	n
2306	5.01040			
	N =	$N = \begin{pmatrix} 1 - \alpha_{ee} \\ -\alpha_{\mu e} \\ -\alpha_{\tau e} \\ \text{Zz. Xing 0} \\ \text{F. J. Escrihuela, D} \\ \text{M. Tortola, and} \\ \text{M. Blennow, I} \\ \text{Garcia, X. Ma} \\ \text{Naredo-Tuero} \\ 2306.01040 \\ \end{pmatrix}$	$N = \begin{pmatrix} 1 - \alpha_{ee} & 0 \\ -\alpha_{\mu e} & 1 - \alpha_{\mu \mu} \\ -\alpha_{\tau e} & -\alpha_{\tau \mu} \\ \text{Zz. Xing 0709.2220 a} \\ \text{F. J. Escrihuela, D. V. Forero,} \\ \text{M. Tortola, and J. W. F. Vall} \\ \text{M. Blennow, EFM, J. He} \\ \text{Garcia, X. Marcano and} \\ \text{Naredo-Tuero and J. Lo} \\ 2306.01040 \end{pmatrix}$	$N = \begin{pmatrix} 1 - \alpha_{ee} & 0 & 0 \\ -\alpha_{\mu e} & 1 - \alpha_{\mu\mu} & 0 \\ -\alpha_{\tau e} & -\alpha_{\tau\mu} & 1 - \alpha_{\tau\tau} \\ Zz. Xing 0709.2220 and 1110.00 \\ F. J. Escrihuela, D. V. Forero, O. G. Miran M. Tortola, and J. W. F. Valle 1503.088 \\ M. Blennow, EFM, J. Hernandez-Garcia, X. Marcano and D. Naredo-Tuero and J. Lopez-Pavo 2306.01040 \\ \end{bmatrix}$

See also R. E. Shrock 1980, 1981; P. Langaker and D. London 1988; S. M. Bilenky and C. Giunti hep-ph/9211269; E. Nardi, E. Roulet and D. Tommasini hep-ph/9503228; D. Tommasini, G. Barenboim, J. Bernabeu and C. Jarlskog hep-ph/9503228; S. Antusch, C. Biggio, EFM, B. Gavela and J. López Pavón hep-ph/0607020; S. Antusch, J. Baumann and EFM 0807.1003; D. V. Forero, S. Morisi, M. Tortola, and J. W. F. Valle 1107.6009; S. Antusch and O. Fischer 1407.6607; F. J. Escrihuela, D. V. Forero, O. G. Miranda, M. Tortola, and J. W. F. Valle 1503.08879; F.J. Escrihuela, D.V. Forero, O.G. Miranda, M. Tórtola, J.W.F. Valle 1612.07377; EFM, J. Hernandez-Garcia and J. Lopez-Pavon 1605.08774; R. Coy, M. Frigerio, 1812.0316, 2110.09126...

Bounds from a global fit to flavour and Electroweak precision data



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Bounds from a global fit to flavour and Electroweak precision data

	"flavor+electroweak" $m > EW \ (2\sigma \ limit)$	Oscillations (from zero distance effects in disappearance, 90%)
α_{ee}	$[0.081, 1.4] \cdot 10^{-3}$	$8.4 \cdot 10^{-3}$ [55]
$\alpha_{\mu\mu}$	$1.4 \cdot 10^{-4}$	$5.0 \cdot 10^{-3}$ [15]
$\alpha_{\tau\tau}$	$8.9\cdot10^{-4}$	$6.5 \cdot 10^{-2}$ [56]
$ \alpha_{\mu e} $	$6.8 \cdot 10^{-4} \ (2.4 \cdot 10^{-5})$	$9.2 \cdot 10^{-3}$
$\left \alpha_{\tau e} \right $	$1.8\cdot10^{-3}$	$1.4 \cdot 10^{-2}$
$ \alpha_{\tau\mu} $	$3.6 \cdot 10^{-4}$	$1.1 \cdot 10^{-2}$

From C. Argüelles et al Snowmass Whitepaper 2203.10811,

M. Blennow, P. Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon 1609.08637 and M. Blennow, EFM, J. Hernandez-Garcia, X. Marcano and D. Naredo-Tuero and J. Lopez-Pavon 2306.01040

A new physics scale

Short and long baseline v oscillations				
eV	keV	MeV	GeV	TeV
-	•			

Precision

electroweak

and flavour

violation

The way out: lighter Steriles

For very light (< keV) extra neutrinos the strong constraints from EW and flavour are lost and ν oscillations dominate



S. Parke and M. Ross-Lonergan arXiv:1508.05095

$$U = \begin{pmatrix} N & \Theta \\ -\Theta^{\dagger} & X \end{pmatrix}$$

"Heavy v'' Non-Unitarity $P_{\alpha\beta}$

$$P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$$

M. Blennow, P. Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon arXiv:1609.08637 C. S. Fong, H. Minakata and H. Nunokawa arXiv:1609.08623

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"Heavy v" Non-Unitarity $P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$
"Light v" Steriles $P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$
 $+ \sum_{I,J} \Theta_{\beta I} \Theta_{\alpha I}^* \Theta_{\alpha J} \Theta_{\beta J}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$
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M. Blennow, P. Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon arXiv:1609.08637 C. S. Fong, H. Minakata and H. Nunokawa arXiv:1609.08623

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"Light v'' Steriles

If $\frac{\Delta m_{ij}^2 L}{2E} \gg 1$ oscillations too fast to resolve and only see average effect

$$P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}L}{2E}}$$

$$+ \sum_{I,J} \Theta_{\beta I} \Theta_{\alpha I}^* \Theta_{\alpha J} \Theta_{\beta J}^* e^{\frac{-i\Delta m_{IJ}^2 L}{2E}}$$

$$+ \sum_{i,J} N_{\beta i} N_{\alpha i}^* \Theta_{\alpha J} \Theta_{\beta J}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$$

 \cdot 2 τ

M. Blennow, P. Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon 1609.08637 C. S. Fong, H. Minakata and H. Nunokawa 1609.08623

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 $+ \sum_{I,J} \Theta_{\beta I} \Theta_{\alpha J}^* \Theta_{\alpha J} \Theta_{\beta J}^* e^{\frac{i\Delta m_{ij}^2 L}{2E}}$
At leading order "heavy" non-unitarity and avergaed-out
"light" steriles have the same impact in oscillations

M. Blennow, P. Coloma, EFM, J. Hernandez-Garcia and J. Lopez-Pavon arXiv:1609.08637

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$$U = \begin{pmatrix} N & \Theta \\ -\Theta^{\dagger} & X \end{pmatrix}$$

Heavy v" Non-Unitarity $P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$
`Light v" Steriles $P_{\alpha\beta} = \sum_{i,j} N_{\beta i} N_{\alpha i}^* N_{\alpha j} N_{\beta j}^* e^{\frac{-i\Delta m_{ij}^2 L}{2E}}$

If $\frac{\Delta m_{ij}^2 L}{2E} \ll 1$ at the near detector or in the data to estimate the flux and cross section, the zero distance effect is recovered and bounds apply

A new physics scale



See talk by Albert de Roeck

A new physics scale



Looking for v_R : Beam Dumps



A. M. Abdullahi, P. Barham Alzas et al. arXiv:2203.08039

See also P. Coloma, EFM, M. González-López, J. Hernández-García arXiv:2007.03701 for a FeynRules file with interactions between mesons and N_R (HNLs) and J. L. Feng, A. Hewitt, F. Kling and D. La Rocco 2405.07330 for a python library

Looking for v_R : Beam Dumps



P. Coloma, J. Lopez-Pavon, L. Molina-Bueno and S. Urrea 2304.06765

A new physics scale





EFM, M. González-López, J. Hernández-García, M. Hostert, J. López-Pavón arXiv:2304.06772 https://github.com/mhostert/Heavy-Neutrino-Limits See also: P. D. Bolton, F. F. Deppisch and P. S. B. Dev arXiv:1912.03058



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Looking for v_R

All together:



EFM, M. González-López, J. Hernández-García, M. Hostert, J. López-Pavón arXiv:2304.06772 https://github.com/mhostert/Heavy-Neutrino-Limits See also: P. D. Bolton, F. F. Deppisch and P. S. B. Dev arXiv:1912.03058

A new physics scale



See talks by Maria Archidiacono and Stefano Gariazzo

Cosmology



A. C Vincent, EFM, P. Hernandez, M. Lattanzi and O. Mena arXiv:1408.1956 See also K. Langhoff, N. J. Outmezguine, and N. L. Rodd arXiv:2209.06216

A new physics scale



See talks in Sesion 6

Conclusions

- Neutrino masses and mixings imply new BSM physics
- The simplest extension, right-handed neutrinos, already imply a lot of new phenomenology to search for:
 - Non-unitarity, searches at colliders, fixed targets, cosmology, *0νββ*,...
- Also offers conexions to other open problems of the SM
 - Baryogenesis, Dark Matter, Flavour puzzle...
- Neutrino detectors can also probe for other BSM physics
- Neutrino physics is an excellent window BSM!!

Non-unitarity and *M*_W from CDF



M. Blennow, P. Coloma, EFM, M-González-Lopez 2204.04559 M. Blennow, EFM, J. Hernandez-Garcia, X. Marcano and D. Naredo-Tuero and J. Lopez-Pavon 2306.01040

It has become common to call them:

"Indirect" or "charged leptons"		"Direct" or "neutrinos"	
	"flavor+electroweak"	Oscillations (from zero distance	
	$m > EW \ (2\sigma \ limit)$	effects in disappearance, 90%)	
α_{ee}	$[0.081, 1.4] \cdot 10^{-3}$	$8.4 \cdot 10^{-3}$ [55]	
$lpha_{\mu\mu}$	$1.4\cdot 10^{-4}$	$5.0 \cdot 10^{-3}$ [15]	
$lpha_{ au au}$	$8.9\cdot10^{-4}$	$6.5 \cdot 10^{-2}$ [56]	
$ \alpha_{\mu e} $	$6.8 \cdot 10^{-4} \ (2.4 \cdot 10^{-5})$	$9.2 \cdot 10^{-3}$	
$ \alpha_{ au e} $	$1.8 \cdot 10^{-3}$	$1.4 \cdot 10^{-2}$	
$ \alpha_{\tau\mu} $	$3.6 \cdot 10^{-4}$	$1.1 \cdot 10^{-2}$	

From C. Argüelles et al Snowmass Whitepaper arXiv:2203.10811 and M. Blennow, EFM, J. Hernandez-Garcia, X. Marcano and D. Naredo-Tuero and J. Lopez-Pavon 2306.01040

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But they all involve



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"Indirect" or "charged leptons"

But they all involve



"Direct" or "neutrinos"

it's where the sensitivity comes from...

So they are all equally "direct" and they all have a neutrino and a charged lepton...

Which one is more robust/model-independent?

"Indirect" or "charged leptons" "Direct" or "neutrinos"

Which one is more robust/model-independent?

Looking for v_R : Non-Unitarity

"Indirect" or "charged leptons"

Introducing an NSI operator with u and d quarks the zero distance effect could be cancelled

"Direct" or "neutrinos"

Which one is more robust/model-independent?

Looking for v_R : Non-Unitarity

"Indirect" or "charged leptons"

Introducing an NSI operator with u and d quarks the zero distance effect could be cancelled They also come from zero-distance effect...

"Direct" or "neutrinos"

Which one is more robust/model-independent?

Looking for v_R : Non-Unitarity

 \downarrow $G_F \text{ from } \mu \text{ decay}$ compared to from M_W ,
measurents of $\sin \theta_W$ at
different energies
(Moller, colliders) and β and *K* decays. Very
different physics! Not
easy to cancel all...

"Indirect" or "charged leptons"

Introducing an NSI operator with u and d quarks the zero distance effect could be cancelled They also induce a zero-distance effect...

"Direct" or "neutrinos"

But the "neutrino" bounds are often assumed to be more robust... why??

v mass from right-handed neutrinos

All SM fermions acquire Dirac masses via Yukawa couplings

$$Y_f \bar{f}_R \phi f_L \xrightarrow{\text{SSB}} \frac{Y_f v}{\langle \phi \rangle} = \frac{Y_f v}{\sqrt{2}} \quad \frac{Y_f v}{\sqrt{2}} \quad \bar{f}_R f_L \quad m_D = \frac{Y_f v}{\sqrt{2}}$$

Simplest option add N_R : a Majorana mass is also allowed

$$m_N N_R N_R^*$$

$$m_V = \begin{pmatrix} 0 & m_D^t \\ m_D & M_N \end{pmatrix} \longrightarrow U^t \begin{pmatrix} 0 & m_D^t \\ m_D & M_N \end{pmatrix} U = \begin{pmatrix} m & 0 \\ 0 & M \end{pmatrix}$$
v mass from right-handed neutrinos

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Simplest option add N_R : a Majorana mass is also allowed

$$m_{N}N_{R}N_{R}^{c}$$

$$m_{\nu} = \begin{pmatrix} 0 & m_{D}^{t} \\ m_{D} & M_{N} \end{pmatrix} \xrightarrow{} U^{t} \begin{pmatrix} 0 & m_{D}^{t} \\ m_{D} & M_{N} \end{pmatrix} U = \begin{pmatrix} m & 0 \\ 0 & M \end{pmatrix}$$
Seesaw

If $M_N \gg m_D$ then $M \approx M_N$ and $m \approx m_D^t M_N^{-1} m_D \rightarrow$ lightness of ν

v mass from right-handed neutrinos



If $M_N \gg m_D$ then $M \approx M_N$ and $m \approx m_D^t M_N^{-1} m_D \rightarrow \text{lightness of } V$

But a very high M_N worsens the Higgs hierarchy problem

Lightness of v masses could also come naturally from an approximate symmetry (B-L)

eV	keV	MeV	GeV	TeV

 M_N could be anywhere...

But a very high M_N worsens the Higgs hierarchy problem

Lightness of ν masses could also come naturally from an approximate symmetry (B-L)

But a very high M_N worsens the Higgs hierarchy problem

Lightness of v masses could also come naturally from an approximate symmetry (B-L)

$$m_D \overline{N}_R \nu_L + M_N \overline{N}_R N_L$$

$$\begin{pmatrix} 0 & m_D^t & 0 \\ m_D & 0 & M_N \\ 0 & M_N & 0 \end{pmatrix}$$

So that $m_{\nu} = 0$ even if $\Theta \approx m_D^{\dagger} M_N^{-1}$ is large

But a very high M_N worsens the Higgs hierarchy problem

Lightness of ν masses could also come naturally from an approximate symmetry (B-L)

$$\begin{split} m_{D}\overline{N}_{R}\nu_{L} + M_{N} \ \overline{N}_{R}N_{L} + \mu \overline{N}_{L}^{c} \ N_{L} \\ \begin{pmatrix} 0 & m_{D}^{t} & 0 \\ m_{D} & 0 & M_{N} \\ 0 & M_{N} & \mu \end{pmatrix} \quad \text{``inverse Seesaw''} \\ \text{R. Mohapatra and J. Valle 1986} \end{split}$$

Small
$$m_{\nu} \approx \mu \frac{m_D^2}{M_N^2}$$
 even if $\Theta \approx m_D^{\dagger} M_N^{-1}$ is large and M_N low



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