 "la Caixa" Foundation
Junior Leader
Fellowship
LCF/BQ/PI23/11970034

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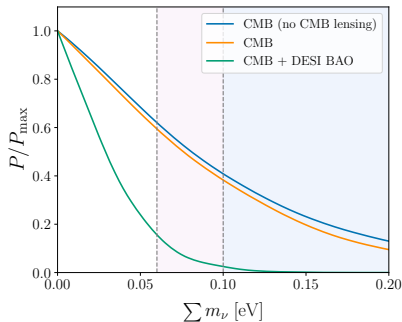
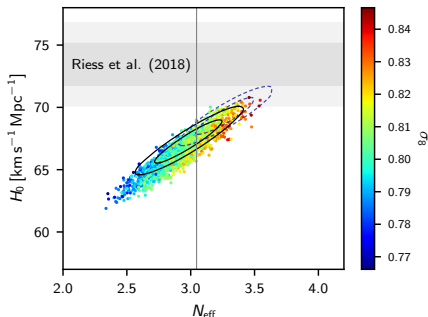
ν non-standard scenarios and cosmology

Don't be tricked!

cosmological ν bounds are not what they seem...

XXXI International Conference on Neutrino Physics and Astrophysics,
Milan, Italy, 16-22/06/2024

Selected neutrino constraints from cosmology



Expectation: $N_{\text{eff}} = 3.044$

[Akita+, 2020] [Froustey+, 2020]

[Bennett, SG+, 2020] [Drewes+, 2024]

Measurement:

$$N_{\text{eff}} = 2.99 \pm 0.17$$

(CMB+lens+BAO)

[Planck Collaboration, 2018]

Oscillations:

$$\Sigma m_\nu \gtrsim 60 \text{ meV (normal ord.)}$$
$$\Sigma m_\nu \gtrsim 100 \text{ meV (inverted ord.)}$$

[any global fit]

Measurement:

$$\Sigma m_\nu < 72 \text{ meV}$$

(95%, CMB+DESI BAO)

[DESI 2024]

C

Clarifications

about the cosmological presence of neutrinos

What is N_{eff} ?

radiation density:

$$\rho_r = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma$$

ρ_γ photon energy density, $7/8$ for fermions, $(4/11)^{4/3}$ due to photon reheating after neutrino decoupling

$$N_{\text{eff}}^\nu = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_\nu}{\rho_\gamma} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{1}{\rho_\gamma} \sum_i g_i \int \frac{d^3 p}{(2\pi)^3} E(p) f_{\nu,i}(p)$$

It is a measurement of the energy density of **relativistic neutrinos!**

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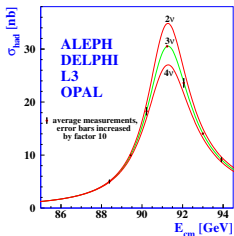
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Nothing to do with [LEP (2006)]

$$N_\nu^{(Z)} = 2.9840 \pm 0.0082$$



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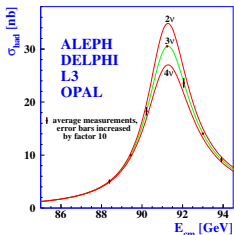
$$N_\nu^{(Z)} = 2.9840 \pm 0.0082$$

instantaneous decoupling:

$$N_{\text{eff}}^\nu = 1 \text{ for each } \nu \text{ family}$$

non-instantaneous decoupling:

$N_{\text{eff}}^\nu > 3$ because of entropy transfer to photons and neutrinos when electrons become non-relativistic



Other relativistic species?

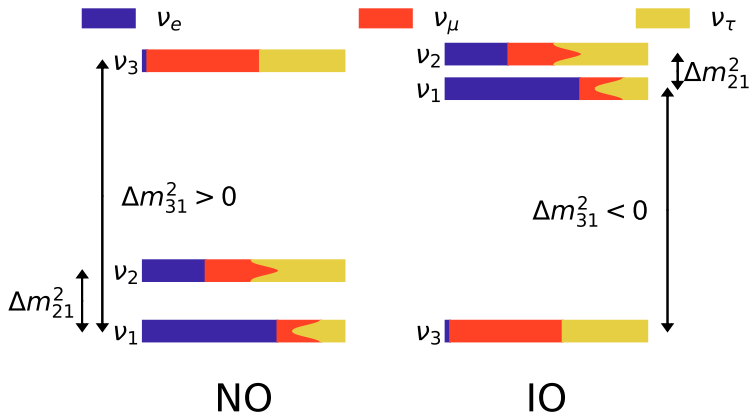
What is Σm_ν ?

normal ordering (NO):

$$m_2 \gtrsim 9 \text{ meV}, m_3 \gtrsim 50 \text{ meV}$$

inverted ordering (IO):

$$m_{1,2} \gtrsim 50 \text{ meV}$$



[Valencia global fit]

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relic neutrinos have a Fermi-Dirac distribution with $T_\nu \approx \mathcal{O}(0.1) \text{ meV}$!

many relic neutrinos are non-relativistic today!

$$\text{energy density } \rho_{\nu, \text{non-rel}} = \sum_i m_i n_i$$

$$\text{fractional energy density } \omega_\nu = \Omega_\nu h^2 = \frac{\sum_i m_i n_i}{\rho_{\text{cr}}} = \frac{\Sigma m_\nu}{94.1 \text{ eV}}$$

Background measurements are sensitive to ω_ν , not Σm_ν !

“Cosmological” Σm_ν measures the energy density of non-relativistic ν s

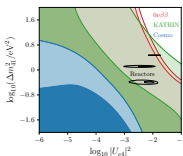
Note: free-streaming scale directly depends on m_ν :

$$\lambda_{\text{FS}} \propto v_{\text{th}}/H \propto \langle p \rangle / (m_\nu H) \propto \sim 3T_\nu / (m_\nu H)$$

N

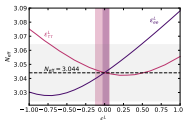
N_{eff}

\Rightarrow energy density of relativistic neutrinos



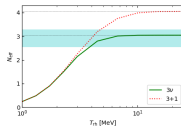
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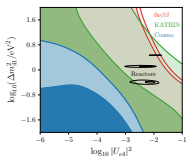
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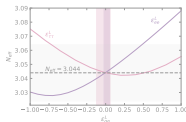
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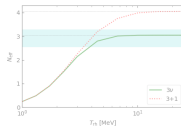
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Four neutrinos \rightarrow new oscillations in the early Universe

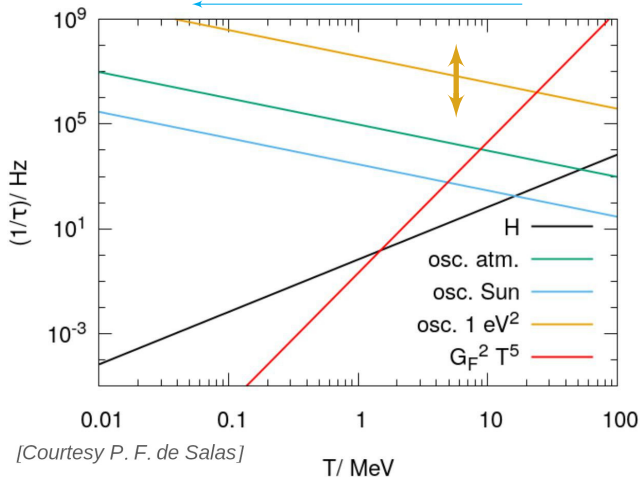
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Sterile neutrino in the early universe

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need to produce it through oscillations, but matter effects may block them
time



beginning of oscillations depends on Δm_{41}^2

later oscillations
 \downarrow
less time before ν decoupling!

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when are they enough to allow full equilibrium of active-sterile states?

$$0 \longleftarrow \Delta N_{\text{eff}} = N_{\text{eff}}^{4\nu} - N_{\text{eff}}^{3\nu} \longrightarrow \simeq 1$$

no sterile production active&sterile in equilibrium

$$\frac{\Delta m_{as}^2}{\text{eV}^2} \sin^4(2\vartheta_{as}) \simeq 10^{-5} \ln^2(1 - \Delta N_{\text{eff}}) \quad (1+1 \text{ approx.})$$

[Dolgov&Villante, 2004]

$$\text{e.g.: } \Delta m_{as}^2 = 1 \text{ eV}^2, \sin^2(2\vartheta_{as}) \simeq 10^{-3} \implies \Delta N_{\text{eff}} \simeq 1$$

$$N_{\text{eff}}^{3\nu} = 3.044 \quad [\text{JCAP 2021}]$$

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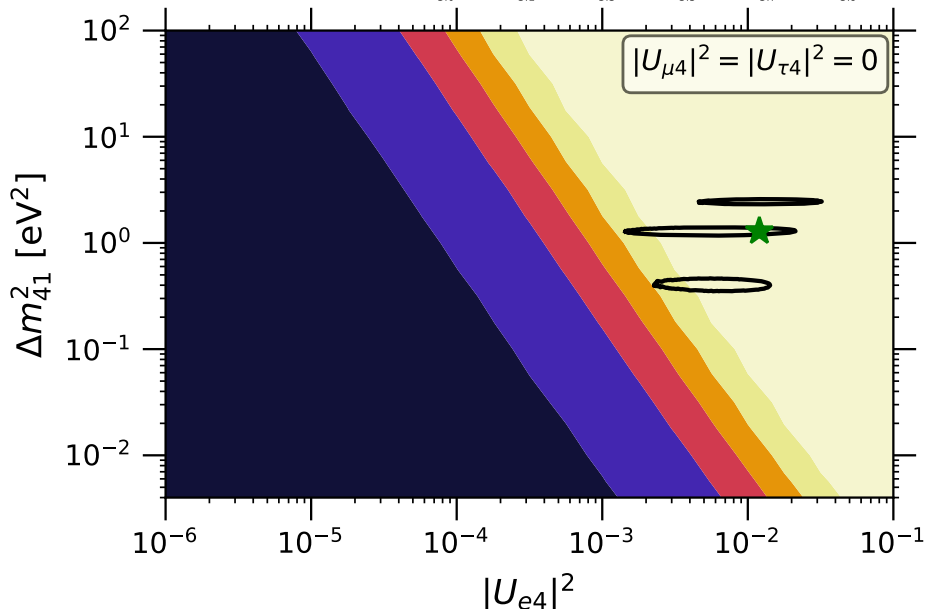
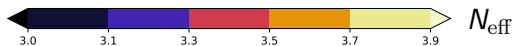
Full calculation: use numerical code!

FORTran-Evolved Primordial Neutrino Oscillations
(FortEPiano)

https://bitbucket.org/ahep_cosmo/fortepiano_public

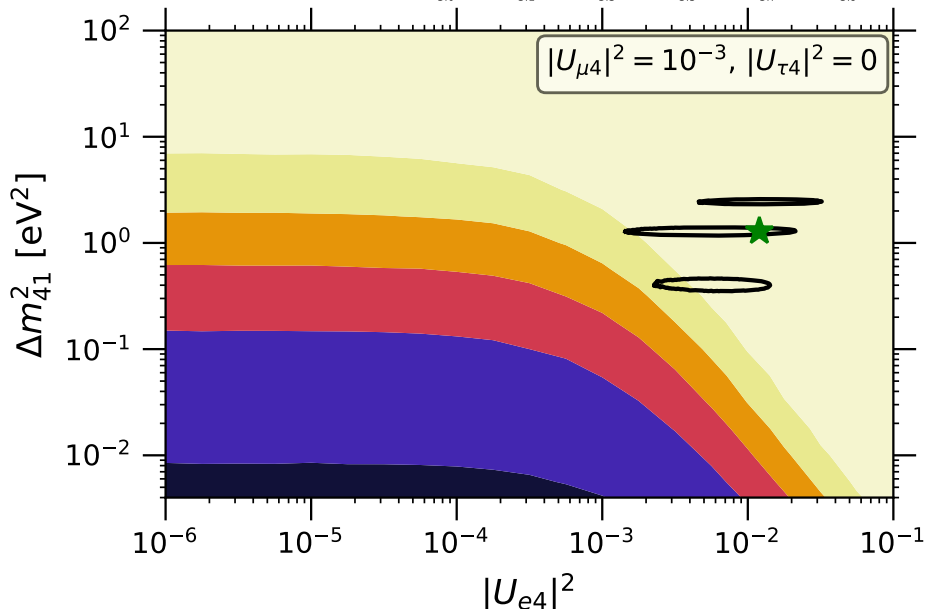
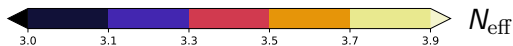
N_{eff} and the new mixing parameters

We can vary more than one angle:



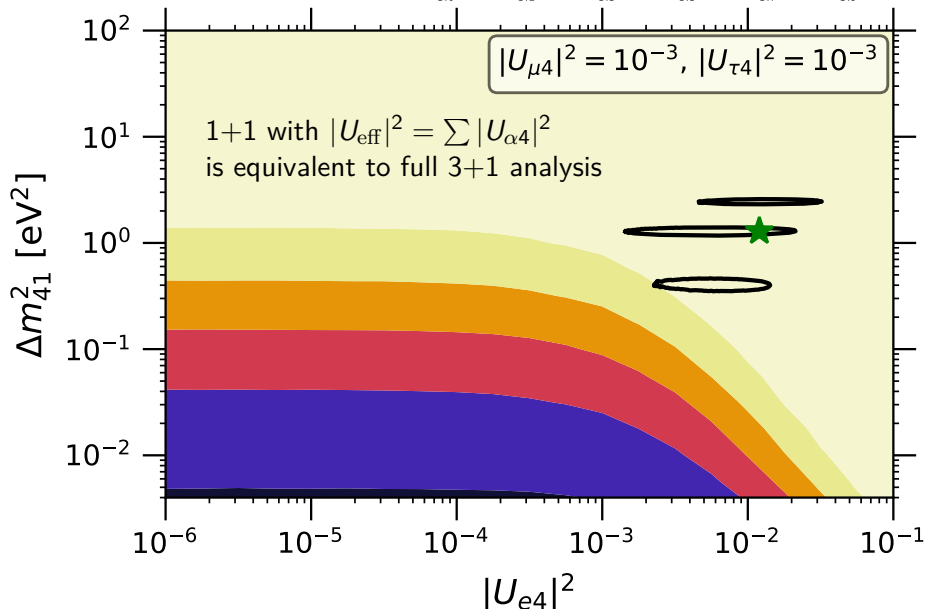
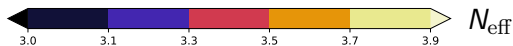
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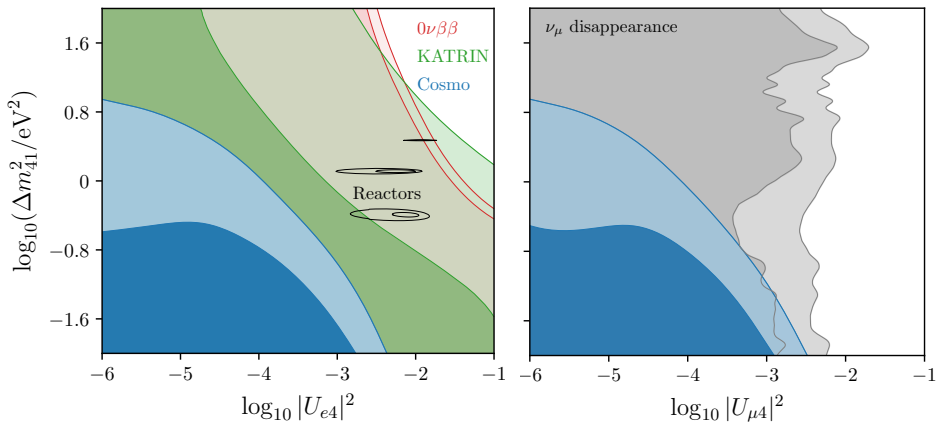
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Comparing constraints

Cosmological constraints are stronger than most other probes

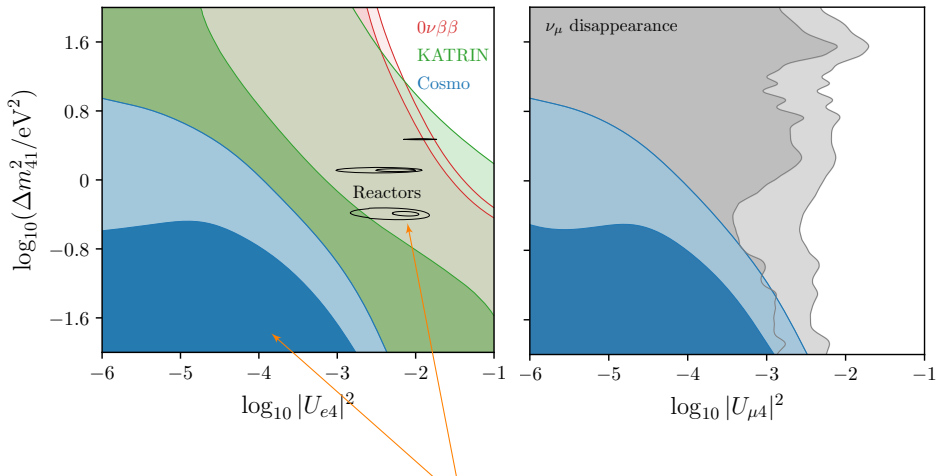
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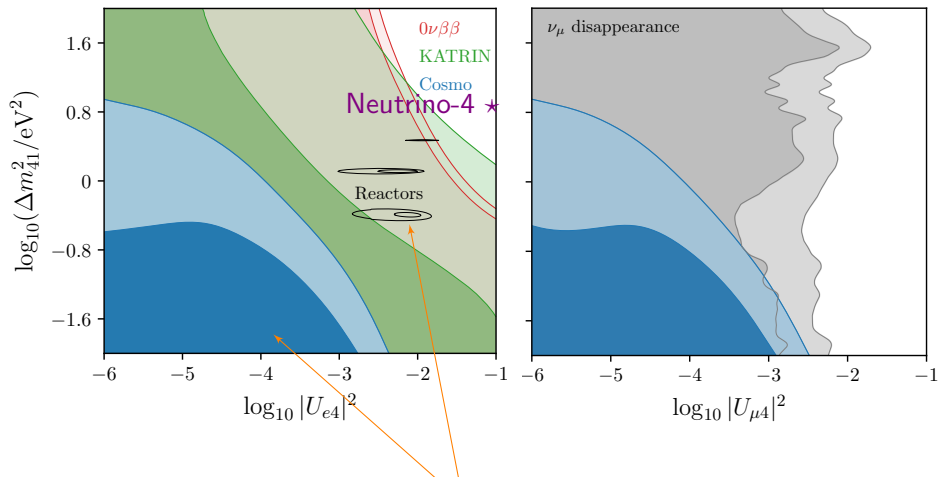


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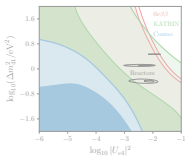


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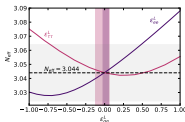
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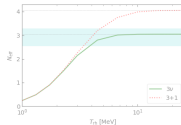
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$N_{\text{eff}} < 3?$

e.g. low-reheating [PRD 92 (2015) 123534], [update in prep.]

Can neutrinos have interactions beyond the SM ones?

e.g.: $\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{NSIe}}$, with $\mathcal{L}_{\text{NSIe}} \propto G_F \sum_{\alpha,\beta} \epsilon_{\alpha\beta}^{L,R} (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{e} \gamma_\mu P_{L,R} e)$
 see e.g. [Farzan+, 2018]

coupling strength governed by the $\epsilon_{\alpha\beta}^{L,R}$ coefficients ($\alpha = e, \mu, \tau$)

new interactions **affect all phenomena** involving neutrinos and electrons
 including neutrino decoupling:

collision terms

$$G_{\text{SM}}^L = \text{diag}(g_L, \tilde{g}_L, \tilde{g}_L)$$

$$G_{\text{SM}}^R = \text{diag}(g_R, g_R, g_R)$$

$g_R = \sin^2 \theta_W$, $\tilde{g}_L = g_R + 1/2$, $\tilde{g}_L = g_R - 1/2$

$$G^{L,R} = G_{\text{SM}}^{L,R} + \begin{pmatrix} \epsilon_{ee}^{L,R} & \epsilon_{e\mu}^{L,R} & \epsilon_{e\tau}^{L,R} & \cdots \\ \epsilon_{e\mu}^{L,R} & \epsilon_{\mu\mu}^{L,R} & \epsilon_{\mu\tau}^{L,R} & \cdots \\ \epsilon_{e\tau}^{L,R} & \epsilon_{\mu\tau}^{L,R} & \epsilon_{\tau\tau}^{L,R} & \cdots \\ \vdots & & & \ddots \end{pmatrix}$$

matter effects in oscillations
 (subdominant!)

$$\mathcal{H}_{\text{eff,SM}} \supset k \cdot \text{diag}(\rho_e + P_e, 0, 0)$$

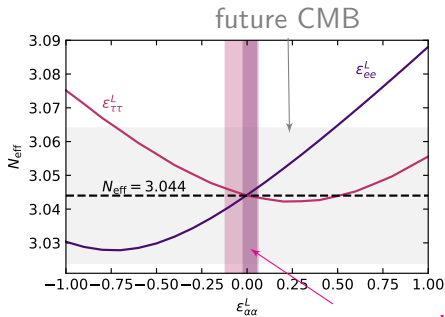
$$\mathcal{H}_{\text{eff}} \supset k(\rho_e + P_e) \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu} & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau} & \epsilon_{\mu\tau} & \epsilon_{\tau\tau} \end{pmatrix}$$

with $\epsilon_{\alpha\beta} = \epsilon_{\alpha\beta}^L + \epsilon_{\alpha\beta}^R$

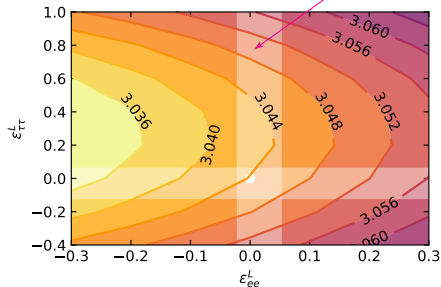
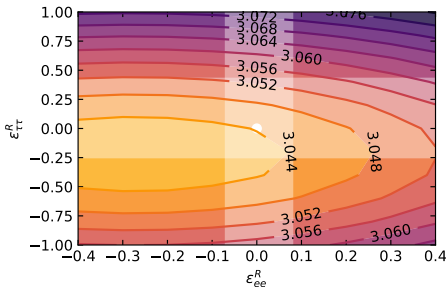
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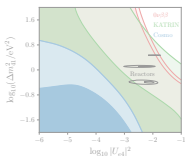
e.g.:

$$\begin{aligned} G_{ee}^L &\rightarrow 0.727 + \epsilon_{ee}^L \\ G_{\tau\tau}^L &\rightarrow -0.273 + \epsilon_{\tau\tau}^L \\ G_{\alpha\alpha}^R &\rightarrow 0.233 + \epsilon_{\alpha\alpha}^R \end{aligned}$$

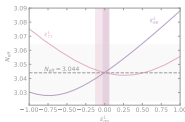


current terrestrial

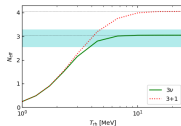


 N_{eff} \Rightarrow energy density of relativistic neutrinos $N_{\text{eff}} > 3?$

e.g. sterile neutrino [JCAP 07 (2019) 014]

 $N_{\text{eff}} \simeq 3?$

e.g. NSI [PLB 820 (2021)]

 $N_{\text{eff}} < 3?$

e.g. low-reheating [PRD 92 (2015) 123534], [update in prep.]

Reheating: phase ending inflation

during inflation, the inflaton (non-rel. scalar) dominates the energy density

during reheating: inflaton decays into standard model particles

⇒ photons, electrons, ... are populated directly

radiation domination begins after reheating

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Low reheating temperature: when reheating occurs at $T_{\text{rh}} \lesssim 20$ MeV

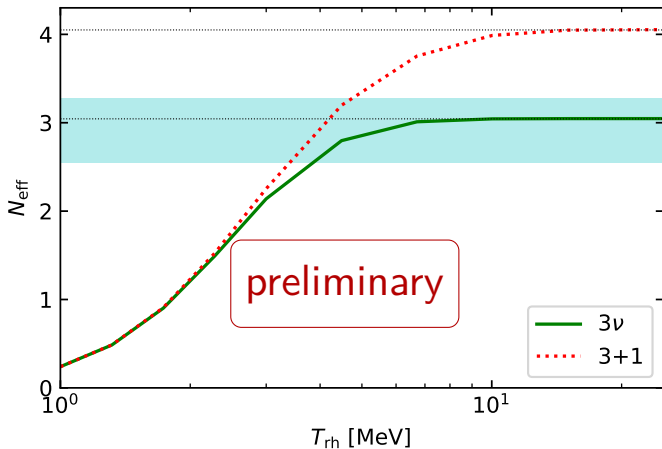
notice: if $T_{\text{rh}} \lesssim 3$ MeV, BBN is broken!

3 neutrino oscillations start to be affected when $T_{\text{rh}} \lesssim 8$ MeV

what about sterile neutrinos?

N_{eff} with low reheating

N_{eff} as a function of T_{rh} (3 or 3+1 neutrinos):

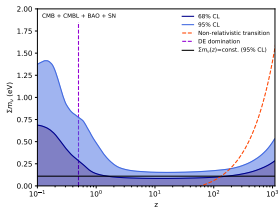


Planck constraint: $N_{\text{eff}} = 2.92^{+0.36}_{-0.37}$ (95%, TT,TE,EE+lowE)



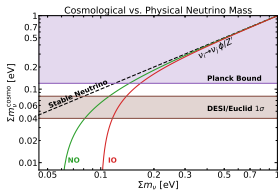
$$\Sigma m_\nu^{\text{cosmo}}$$

⇒ energy density of non-relativistic neutrinos



$\Sigma m_\nu^{\text{cosmo}} < \Sigma m_\nu$ by time-varying masses?

e.g. [Lorenz+, PRD 104 (2021) 123518]



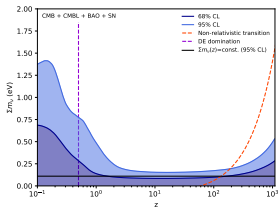
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e.g. [Escudero+, JHEP 12 (2020) 119]



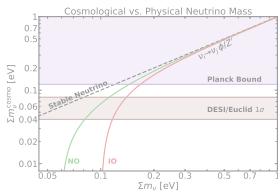
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Time-varying neutrino masses

How to make neutrino masses non-constant in time?

new coupling!

a scale factor

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To new scalar field
(early) dark energy?

e.g. [Ayaita+, PRD 2016]

$m_\nu = 0$ at early times, then

phase transition generates
 m_ν at late times ($a \gtrsim 0.2$)

Masses can grow up to
 $m_\nu = 0.6$ eV at $a = 1$

ν may decay to radiation later

if supercooled transition, data
prefer it at $a \sim 1$ [Lorenz+, PRD 2019]

Massless ν at all times!

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↓
Massless ν at all times!

To **dark radiation**
e.g. [Dvali+, PRD 2016]

assume anomalous symmetry
("axial" neutrino lepton number)

broken at scale Λ_G

s.b. generates effective m_ν

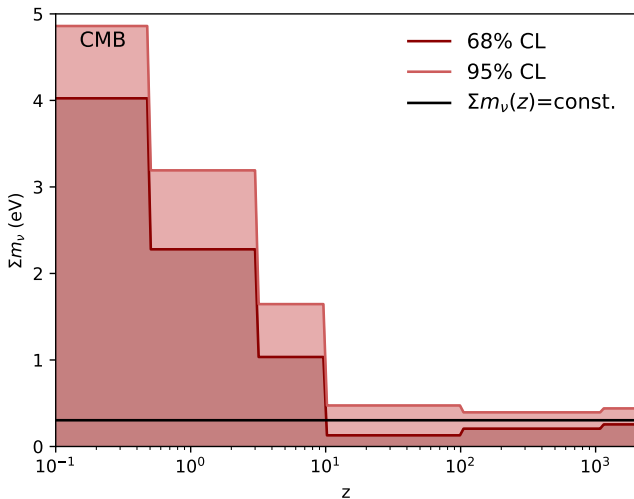
$m_\nu = 0$ at early times, then
generation of masses at $T_\Lambda \lesssim m_\nu$
($T_\Lambda < T_{\text{CMB}}$ in order to preserve recombination)

Massive ν may decay into lightest ν
and annihilate into Goldstone bosons

↓
cosmological mass bound is avoided!

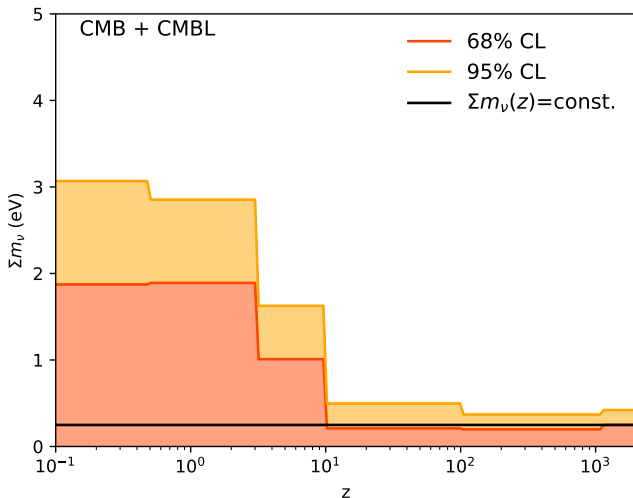
a scale factor

Phenomenological approach: $\Sigma m_\nu(z) = \text{const}$ in 6 redshift bins



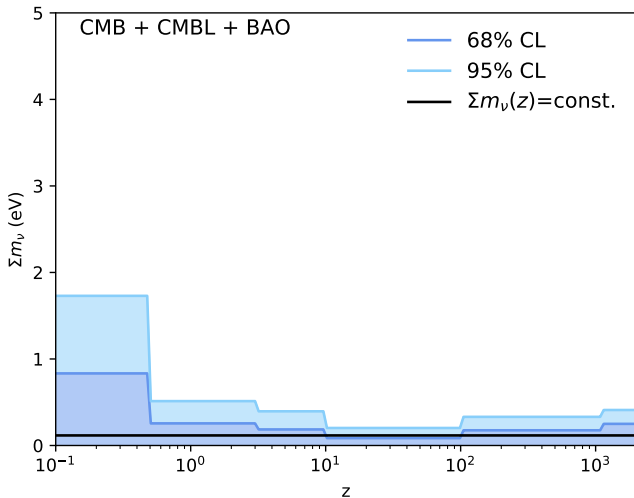
CMB: very high redshift, poor constraints on late universe

Phenomenological approach: $\Sigma m_\nu(z) = \text{const}$ in 6 redshift bins



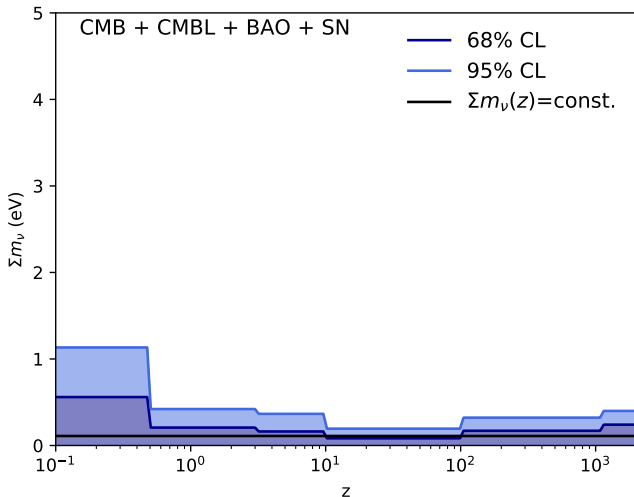
CMBL: lensing affects CMB evolution until today

Phenomenological approach: $\Sigma m_\nu(z) = \text{const}$ in 6 redshift bins



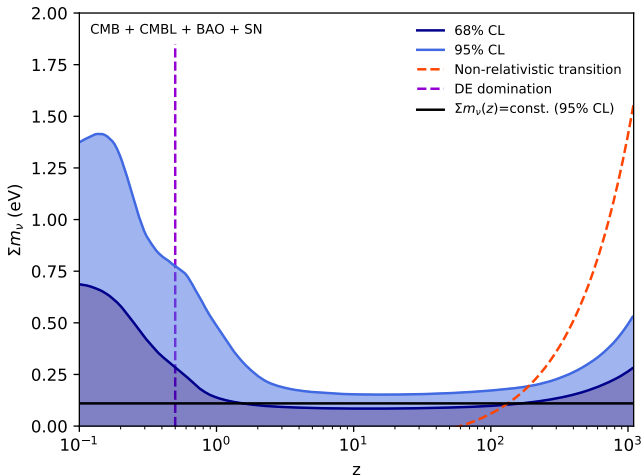
BAO: SDSS BAO+Lyman- α +quasars measurements at $z < 2.3$

Phenomenological approach: $\Sigma m_\nu(z) = \text{const}$ in 6 redshift bins



SN: Pantheon SN samples at $z < 2.3$ from SDSS, SNLS, HST

Phenomenological approach: $\Sigma m_\nu(z) = \text{const}$ in 6 redshift bins

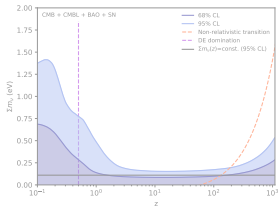


ν energy density is strongly constrained in the range $1 \lesssim z \lesssim \mathcal{O}(100)$, corresponding to matter domination while ν are non-relativistic



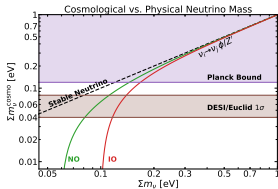
$$\Sigma m_\nu^{\text{cosmo}}$$

⇒ energy density of non-relativistic neutrinos



$\Sigma m_\nu^{\text{cosmo}} < \Sigma m_\nu$ by time-varying masses?

e.g. [Lorenz+, PRD 104 (2021) 123518]



$\Sigma m_\nu^{\text{cosmo}} < \Sigma m_\nu$ by invisible decay?

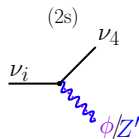
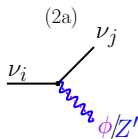
e.g. [Escudero+, JHEP 12 (2020) 119]

Can a neutrino decay? Is the decay lifetime τ_ν constrained?

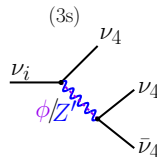
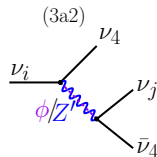
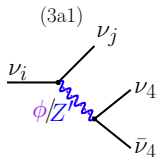
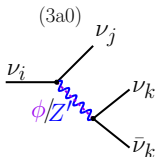
Decay into **visible (e.m.)** particles is constrained to $\tau_\nu > 10^2 t_U$

Invisible decay is much less constrained! A few examples:

2-body Decays:



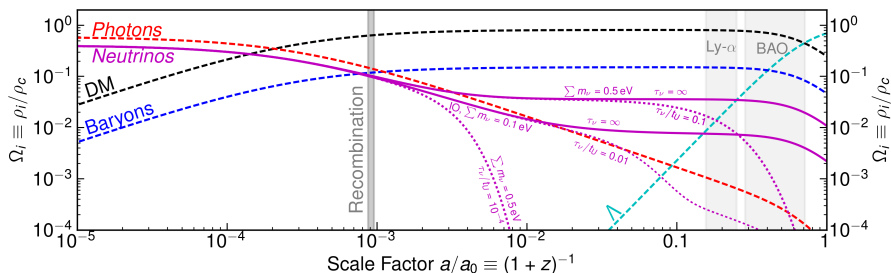
3-body Decays:



if $m_4 = m_\phi = m_{Z'} = 0$, neutrinos decay to radiation and the mass bounds are avoided if $\tau_\nu \lesssim t_U$ is short enough

Can a neutrino decay? Is the decay lifetime τ_ν constrained?

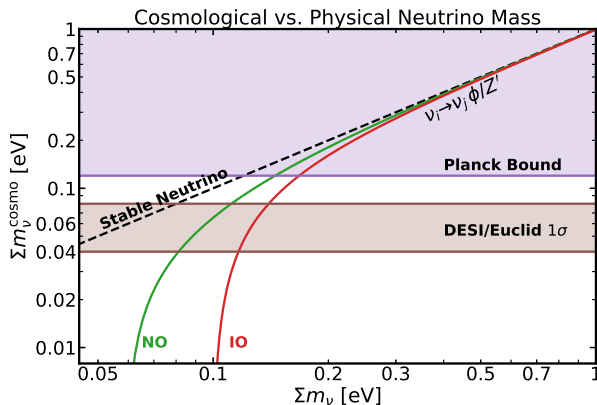
Depending on τ_ν , the standard energy density of neutrinos is suppressed:



if $\tau_\nu \simeq 10^{-4} t_U$, decay occurs at recombination
(if $\Sigma m_\nu < 0.6$ eV, neutrinos are still relativistic)

for large τ_ν , suppression may be partial

Consider 2-body decay channels:



$$\Sigma m_\nu^{\text{cosmo}} = 3m_{\text{lightest}} \text{ if all except lightest neutrino decay to } \nu_{\text{lightest}} \phi / Z'$$

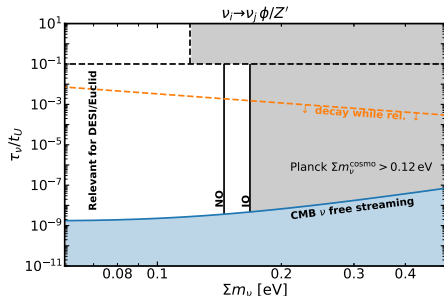
Bounds obtained with $\tau_\nu \lesssim t_U/10$

DESI/Euclid bound is forecast with $\Sigma m_\nu = 60 \text{ meV}$

Consider 2-body decay channels:

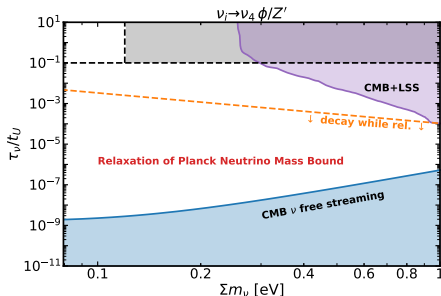
Constraints from **Planck legacy** or **late time clustering+CMB lensing**

Active neutrino + boson:



late time constraints not yet considered, may provide additional constraints

Sterile neutrino + boson:

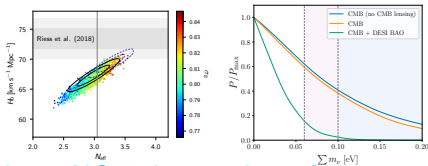


for appropriate τ_ν and Σm_ν , the mass bounds do not apply because all goes into radiation and $\Sigma m_\nu^{\text{cosmo}} = 0$

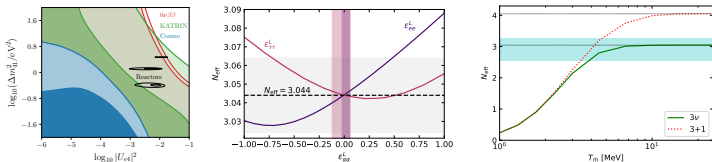
E Conclusions

What do we learn about non-standard ν scenarios?

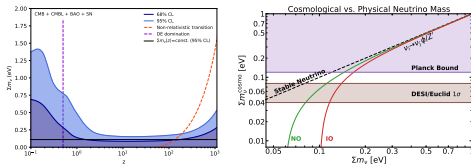
C Cosmology measures (mostly) neutrino energy densities!



N N_{eff} is NOT the number of neutrinos!

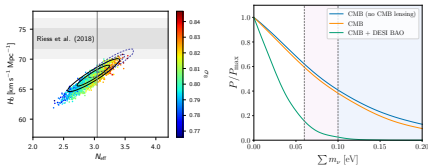


M The “cosmological” Σm_ν is NOT the sum of neutrino masses!

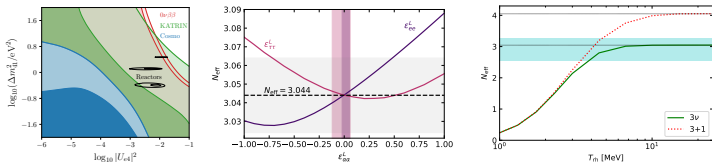


What do we learn about non-standard ν scenarios?

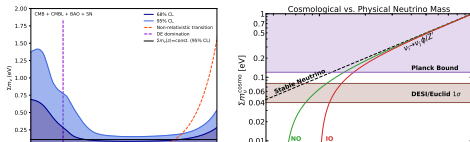
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Thanks for your attention!