

## Introduction

- Neutrino oscillation is a quantum mechanical phenomenon that arises due to the *coherent superposition of neutrino mass states*. However, if neutrino as a quantum system is coupled to an environment, the coherence between two or more propagating states may be lost leading to the suppression of flavour oscillations.
- Such type of **environmentally induced quantum decoherence (QD)** in neutrino states might emerge *from quantum gravity effects or space-time "foam" which acts as dissipative sources and can modify the  $\nu$ -oscillation probability in various ways [1]*.

## QD: Formalism

The time evolution of neutrinos in an open quantum system is given by

$$\frac{\partial \rho(t)}{\partial t} = -i[H, \rho(t)] + \mathcal{D}[\rho(t)], \quad (1)$$

where  $H$  is the neutrino Hamiltonian which can be written as

$$H = \frac{1}{2E} \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + V_{CC} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}. \quad (2)$$

The effect of decoherence is given by the dissipator matrix  $\mathcal{D}$  which on imposing relevant physical conditions can be parametrized as

$$\mathcal{D} = -\text{diag}(\Gamma_{21}, \Gamma_{21}, 0, \Gamma_{31}, \Gamma_{31}, \Gamma_{32}, \Gamma_{32}, 0) \quad (3)$$

The solution to the Eq.1 is given by

$$\rho_{ij}^\alpha(x) = \tilde{U}_{\alpha i}^* \tilde{U}_{\alpha j} e^{-(\Gamma_{ij} + i\tilde{\Delta}_{ij})x} \quad (4)$$

and the oscillation probabilities in the presence of decoherence read as

$$P(\nu_\alpha \rightarrow \nu_\beta) = \text{Tr}[\rho^\alpha(0)\rho^\beta(x)] \quad (5)$$

Here,  $\tilde{U}$  is the modified PMNS matrix in matter and  $\tilde{\Delta}_{ij} = \frac{\Delta \tilde{m}_{ij}^2 L}{4E}$ , with  $\Delta \tilde{m}_{ij}^2$  being the mass squared differences in matter.

- $L = 360$  km, *Water Cherenkov detector of fiducial volume 538 kt and 5 years run-time of neutrino + 5 years of antineutrino* [2].

## Oscillation Probability including Decoherence

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} - 2 \sum_{i>j} \text{Re}[\tilde{U}_{\alpha i}^* \tilde{U}_{\beta i} \tilde{U}_{\beta j} \tilde{U}_{\alpha j}^*] [1 - \cos(2\tilde{\Delta}_{ij}) e^{-\Gamma_{ij}L}] + 2 \sum_{i>j} \text{Im}[\tilde{U}_{\alpha i}^* \tilde{U}_{\beta i} \tilde{U}_{\beta j} \tilde{U}_{\alpha j}^*] \sin(2\tilde{\Delta}_{ij}) e^{-\Gamma_{ij}L}$$

*QD induces terms similar to damping phenomena of the form  $e^{-\Gamma_{ij}L}$  in the oscillation probability.*

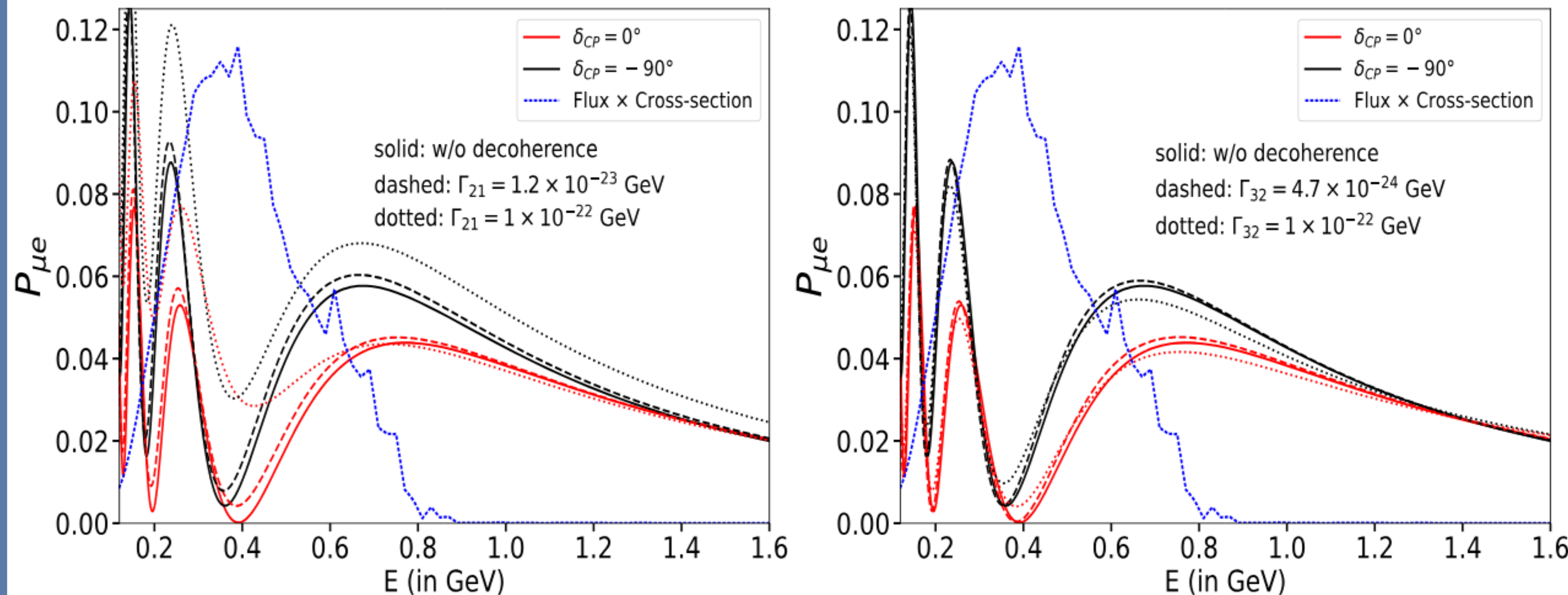


Fig.: Effect of  $\Gamma_{21}$  and  $\Gamma_{32}$  on the appearance probability for the ESSnuSB.

- The  $P_{\mu e}$  probability is **mostly affected by  $\Gamma_{21}$**  in comparison to  $\Gamma_{32}$ .
- The effect of  $\Gamma_{21}$  is significant when  $\delta_{CP}$  is  $-90^\circ$ .

## Important Analytical Formulae

For maximal CPV, the relevant contributing terms for  $\Gamma_{21}$  and  $\Gamma_{32}$  are

$$|P_{\mu e}^{\text{CP-odd}}| \propto |2\alpha\Delta_{31}(\Gamma_{21}L - 2\sin^2\Delta_{31}) - \Gamma_{21}L\sin 2\Delta_{31}|$$

$$|P_{\mu e}^{\text{CP-odd}}| \propto |\Gamma_{32}L\cos 2\Delta_{31} + 2\sin^2\Delta_{31}|$$

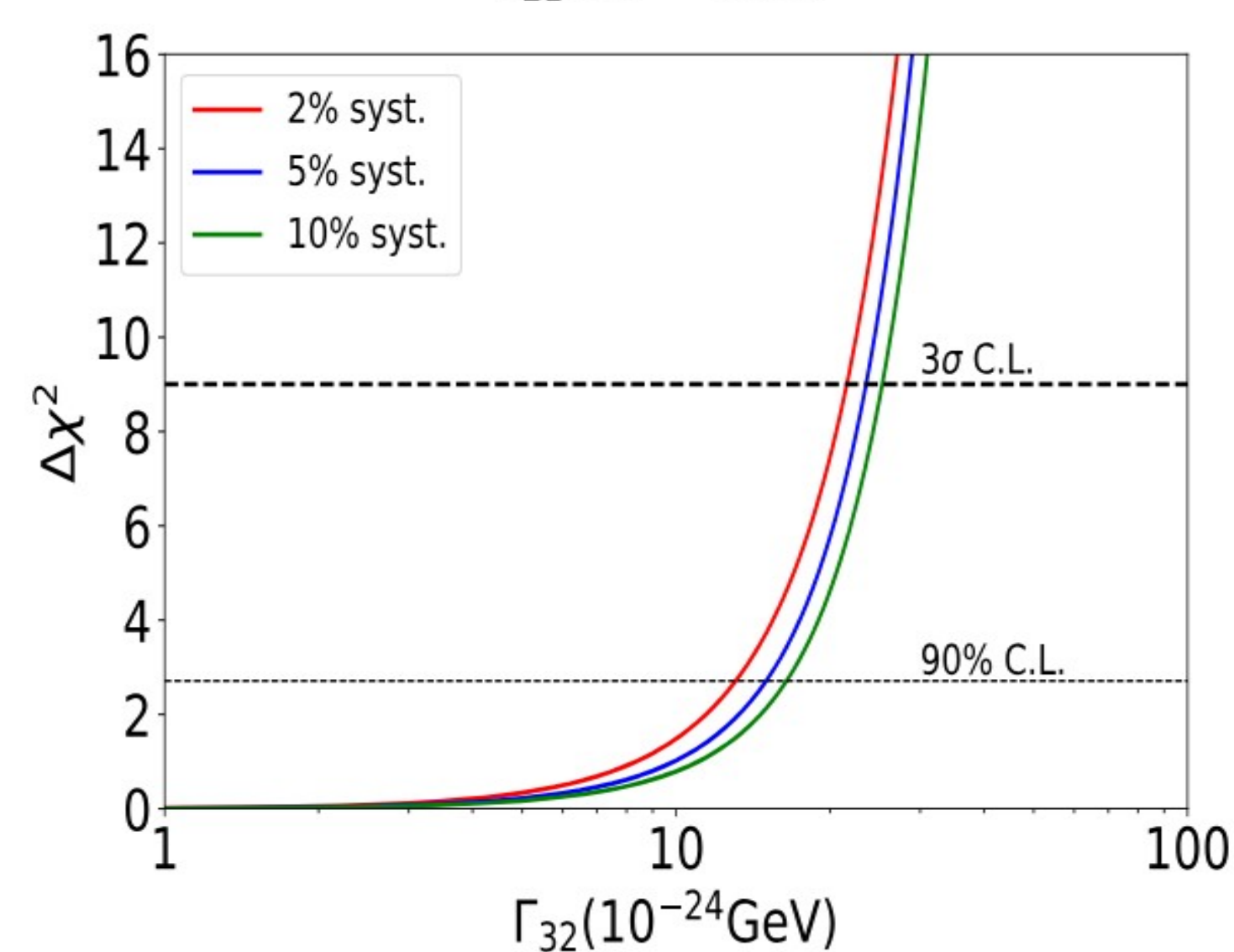
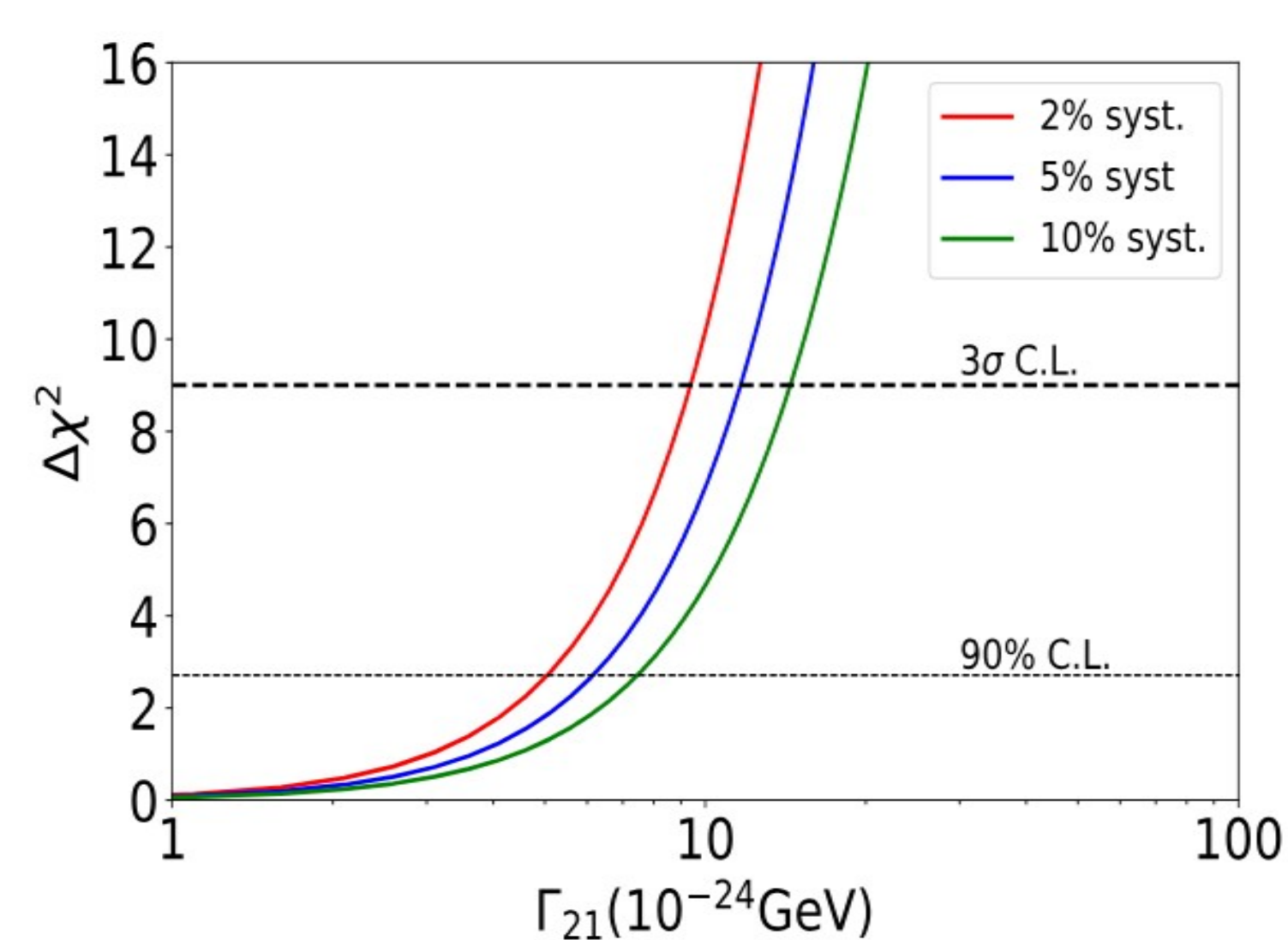
For  $\delta_{CP} = 0, 180^\circ$ , expression for CPV precision is given by:

$$\Delta\delta_{CP} \propto \frac{1}{|2\alpha\Delta_{31}(\Gamma_{21}L - 2\sin^2\Delta_{31}) - \Gamma_{21}L\sin 2\Delta_{31}|}; \frac{1}{|\alpha\Delta_{31}\Gamma_{32}L\cos 2\Delta_{31} + 2\alpha\Delta_{31}\sin^2\Delta_{31}|}$$

On the other hand, for maximal CP violation, we obtain

$$\Delta\delta_{CP} \propto \frac{1}{|\Gamma_{21}L(\cos 2\Delta_{31} + \cos 2\theta_{12}) + 2\alpha\Delta_{31}\sin 2\Delta_{31}|}; \frac{1}{|\alpha\Delta_{31}(1 - \Gamma_{32}L)\sin \Delta_{31}|}$$

## Bounds on $\Gamma_{21}$ and $\Gamma_{32}$



Decoherence Parameter	90% C.L. (in GeV)		
	2% syst.	5% syst.	10% syst.
$\Gamma_{21}$ when $\Gamma_{32} = 0$	$5.06 \times 10^{-24}$	$6.15 \times 10^{-24}$	$7.45 \times 10^{-24}$
$\Gamma_{32}$ when $\Gamma_{21} = 0$	$1.31 \times 10^{-23}$	$1.50 \times 10^{-23}$	$1.64 \times 10^{-23}$
$\Gamma_{21} = \Gamma_{32}$	$4.23 \times 10^{-24}$	$4.99 \times 10^{-24}$	$5.64 \times 10^{-24}$

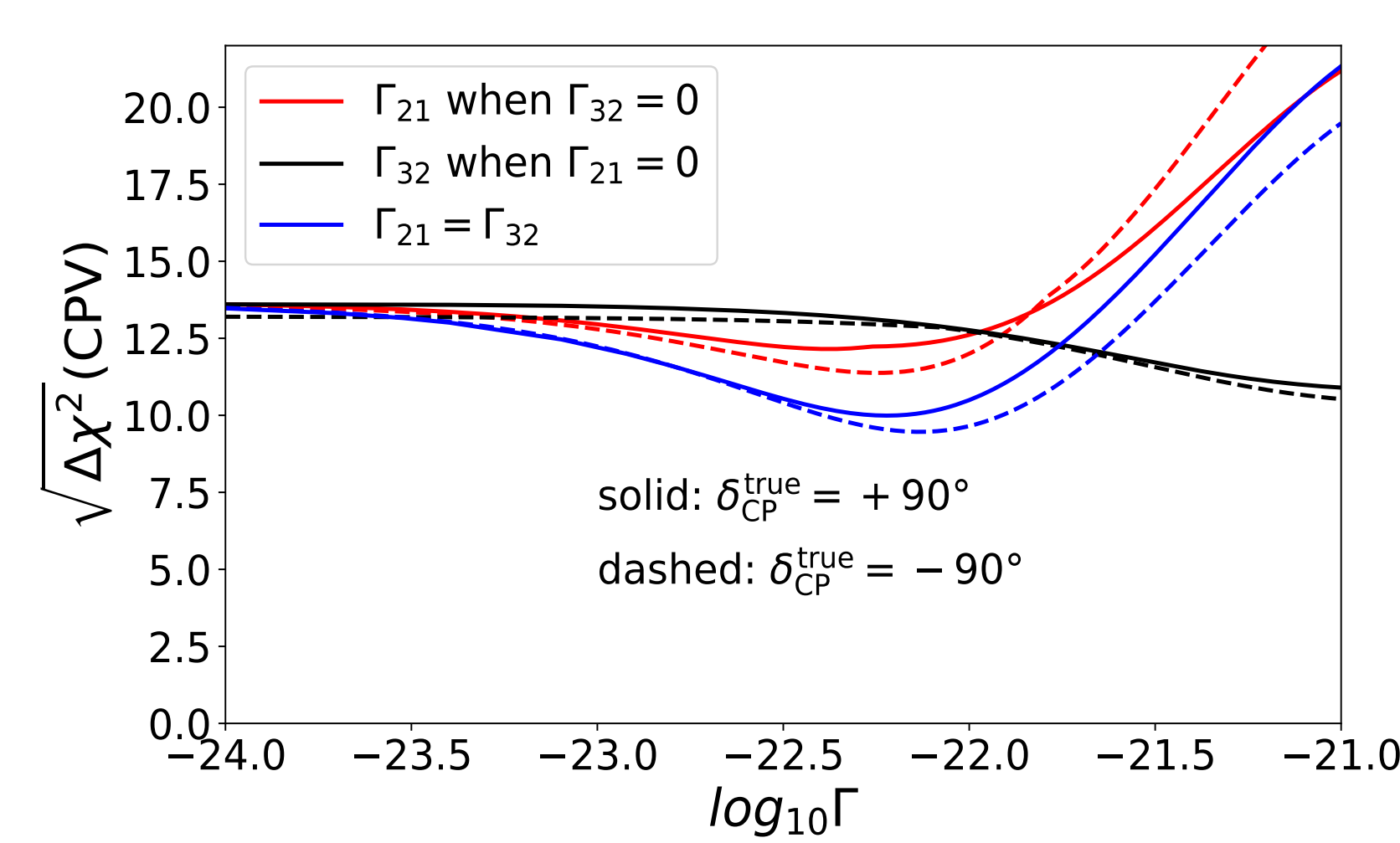
Other experimental bounds at 90% C.L. are [3, 4]

$$\Gamma_{32} = \Gamma_{21} < 9.4 \times 10^{-24} \text{ GeV [MINOS/MINOS+]}$$

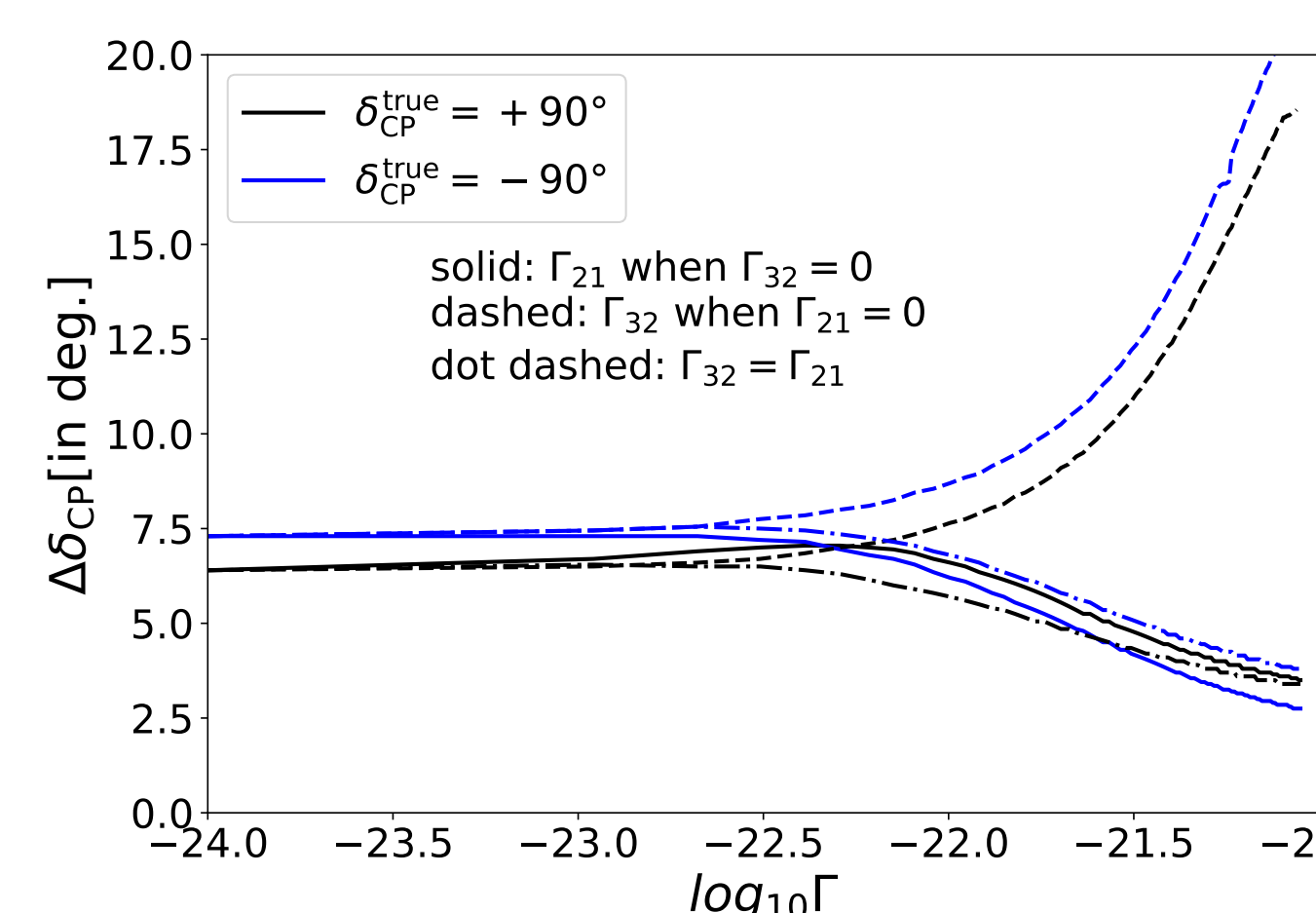
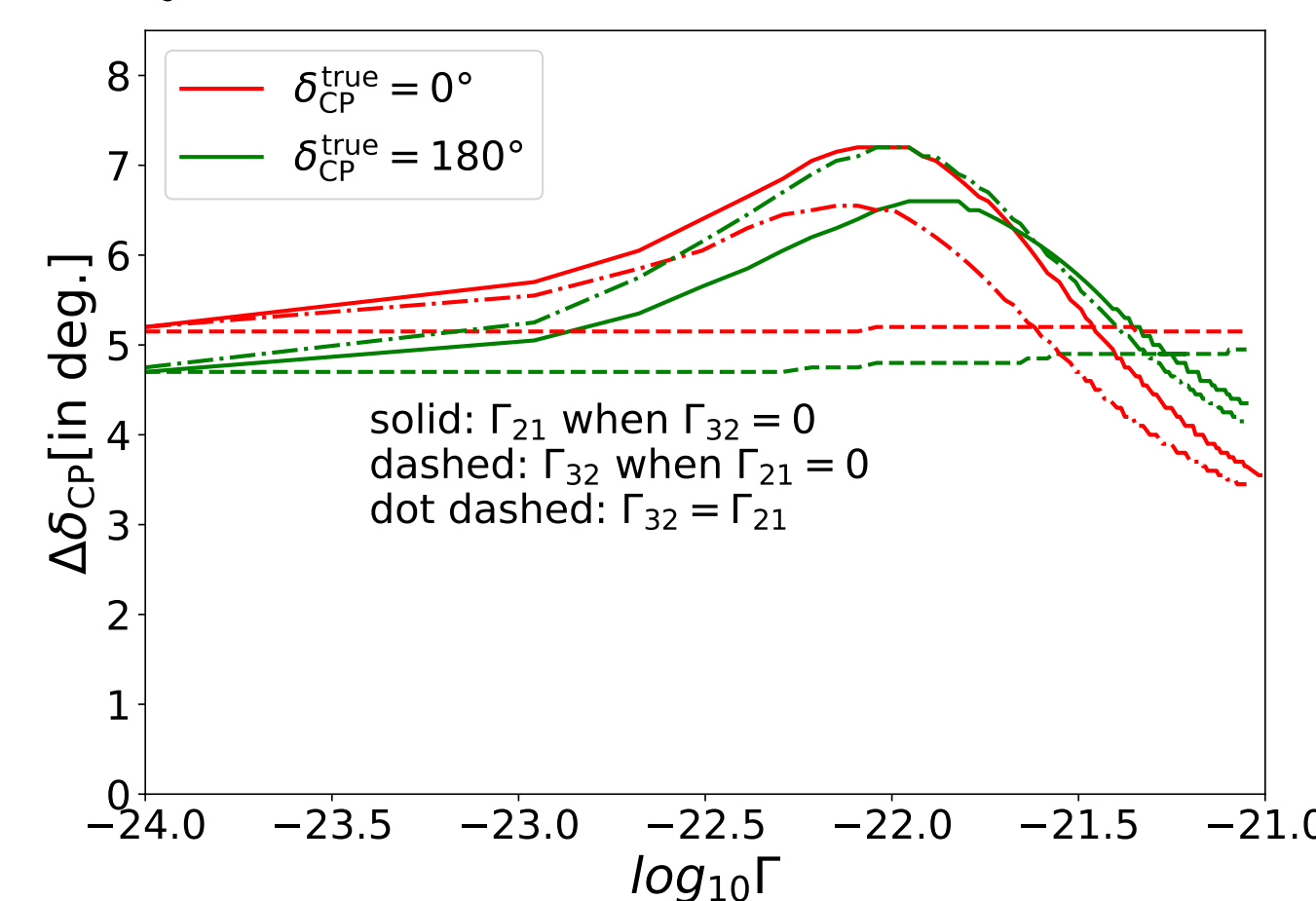
$$\Gamma_{21} < 1.2 \times 10^{-23} \text{ GeV [DUNE]},$$

$$\Gamma_{32} < 4.7 \times 10^{-24} \text{ GeV [DUNE]}.$$

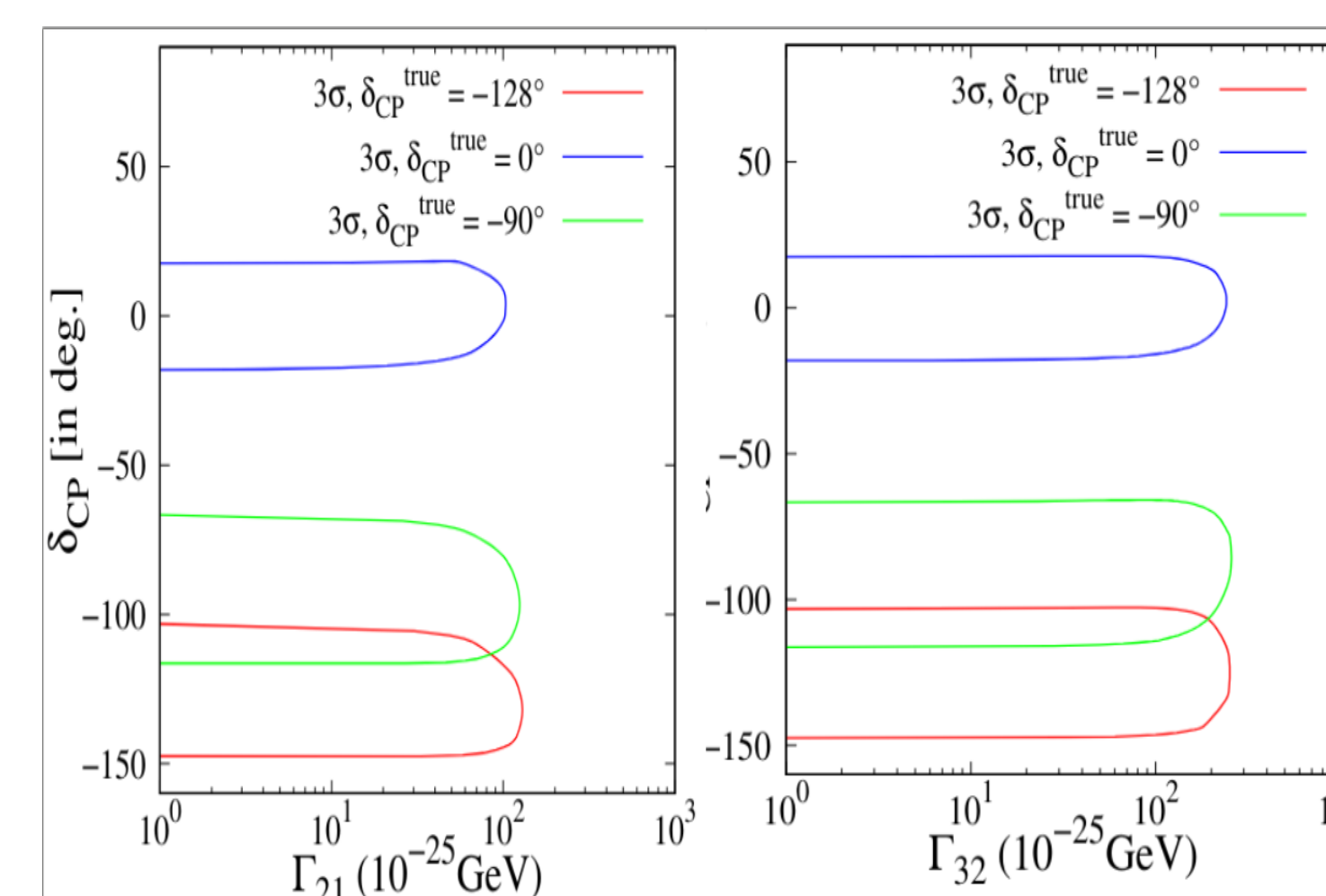
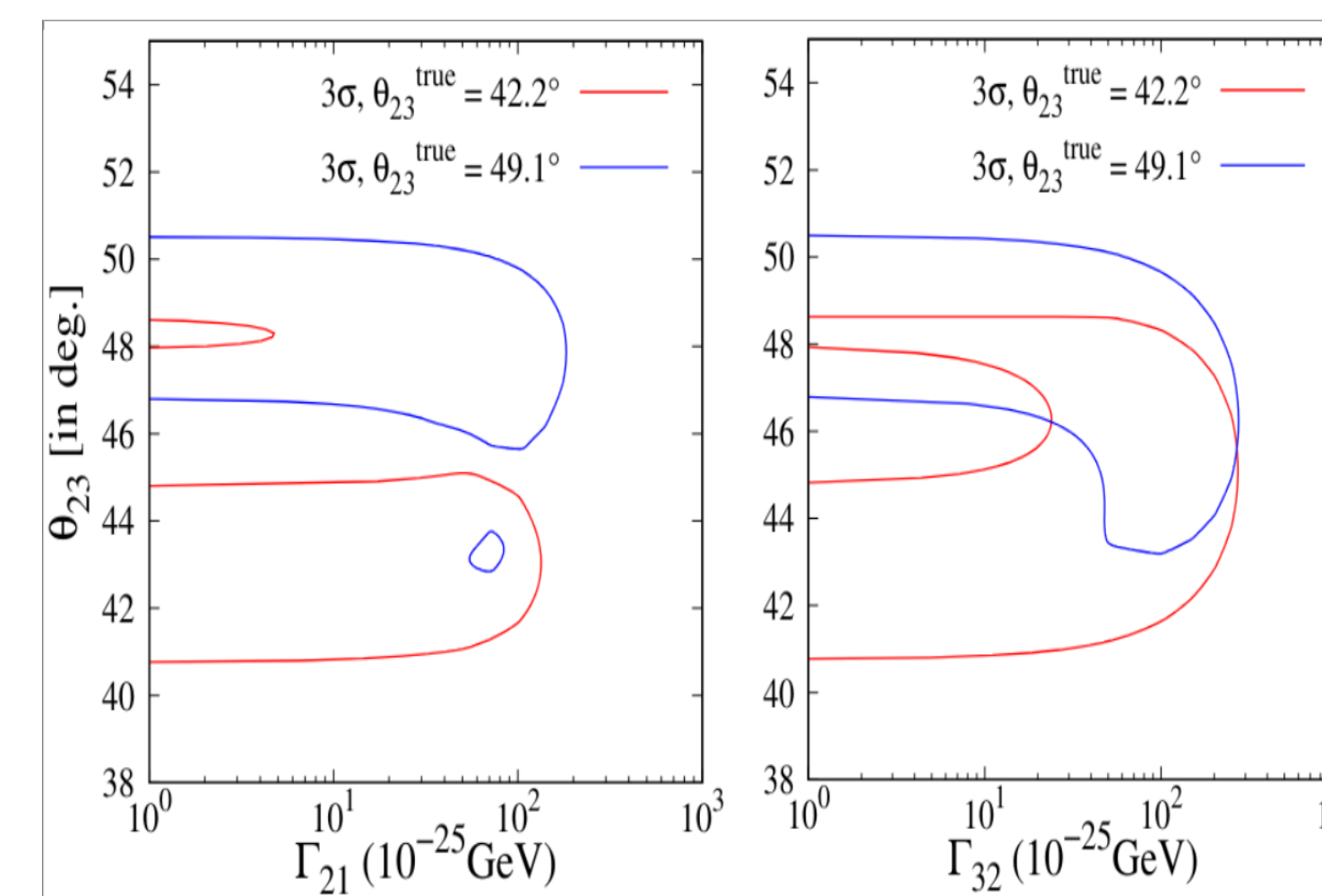
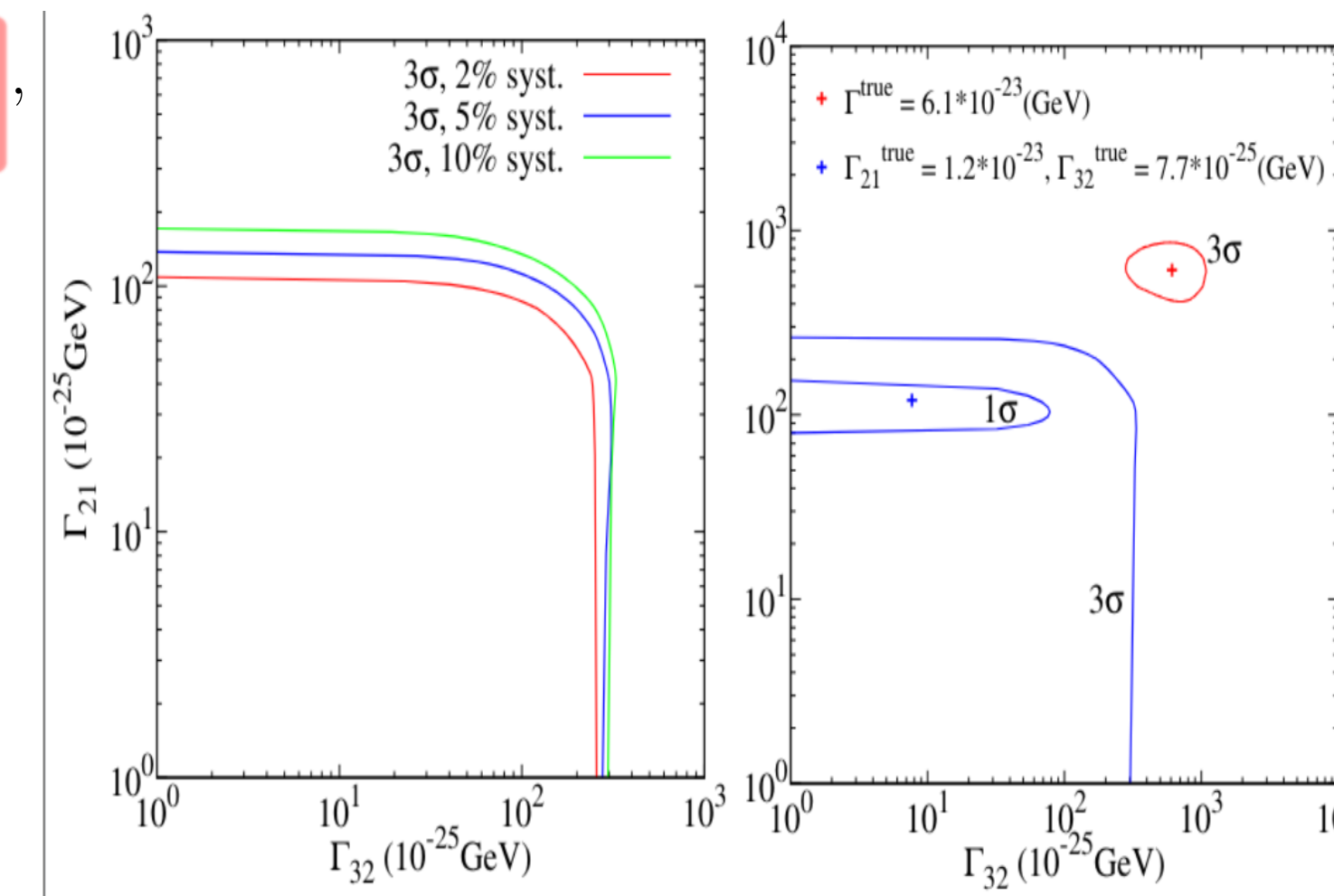
## CPV Sensitivity & Precision



For small  $\Gamma_{21}$ , the sensitivity first slightly decreases and when the decoherence term becomes dominant, the relevant probability increases along with  $\Gamma_{21}$ , improving the sensitivity.



## Correlations



## Key Takeaways

- For the first time, we explore the sensitivity of **ESSnuSB** to constrain  $\Gamma_{21}$  and  $\Gamma_{32}$  [5].
- We find that, *the bounds on  $\Gamma_{21}$  are better than MINOS/MINOS+ and DUNE, while constraint on  $\Gamma_{32}$  is competitive.*
- ESSnuSB measurement of  $\delta_{CP}$  remains robust for  $\Gamma_{ij}$  in the range  $[10^{-24}, 10^{-21}]$  GeV.**
- For the case of maximal CP violation, an uncertainty below  $10^\circ$  can be maintained for  $\Gamma_{ij} \gtrsim 10^{-22}$  GeV.*
- Interesting correlations have been observed among  $\theta_{23}$  and  $\Gamma_{ij}$ .*

## References

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