

Dynamical friction on compact binary systems

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[VD, Nusser, Buehler ApJ, arXiv: 2111.07366]

[Buehler & VD PRD, arXiv: 2207.13740]

[Buehler, Kolyada & VD, arXiv: 2310.05244]

GWFPF workshop, Rome, 15/2/24

Outline

Dynamical Friction (DF) in structure formation

Solving circular DF with a multipole expansion

Perfect fluid vs. ULA dark matter

Including eccentricity and density profiles

Dynamical Friction (DF)





this results in an enhanced region of density behind the mass, with a drag force, F_d known as dynamical friction



- The density wake is global (far-field matters)
 - The medium can be anything (stars, DM, gas...)



Figure @ Luke Leisman

DF in structure formation

DF is a key mechanism which controls

- The merging of DM halos, galaxies, Globular clusters
- The slow-down of spinning galactic bars
- The coalescence of compact objects

Fornax Globular clusters



coalescing AGNs



Signature of DF in GW data



 $\bigcirc \boxtimes \oslash$





[Ginat, Glanz, Perets, Grishin, VD, MNRAS, arXiv: 1903.11072]

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Analytical approaches to DF

Standard assumptions:

- Linear response theory
- The perturber is a point mass in straight-line motion at constant velocity
- Infinite, homogeneous and isotropic medium
- Self-gravity of the medium is ignored

DF for perturbers in straight-line motion



[Chandrasekhar 43; Dokuchaev 64; Ruderman+ 71; Rephaeli+ 80; Just+ 91; Ostriker 99, ...]

Compact binary coalescence: DF on bound eccentric orbits

DF for perturbers on circular orbits



DF for perturbers on circular orbits



Circular DF in a perfect fluid

[VD, Nusser, Buehler ApJ, arXiv: 2111.07366]

• Linear response theory:

$$\ddot{\delta}\rho - c_s^2 \nabla^2 \delta\rho = 4\pi G \bar{\rho} \,\rho_p$$

$$G_{\rm ret}(\vec{r},\tau) = \frac{1}{r} \,\delta_D \left(\tau - \frac{r}{c_s}\right)$$
$$\delta\rho \sim \frac{GM}{c_s^2 r} \bar{\rho}$$

• Multipole expansion:

$$I(\mathcal{M}) = \sum_{\ell m} I_{\ell m}(\mathcal{M})$$
$$I_{\ell m}(\mathcal{M}) \sim j_{\ell}(m\mathcal{M}) h_{\ell}^{(1)}(m\mathcal{M})$$

$$\vec{F}_{\rm DF}(t) = -4\pi \left(\frac{GM}{v_p}\right)^2 \bar{\rho} \left(\Re(I)\hat{r}(t) + \Im(I)\hat{\varphi}(t)\right) \frac{\vec{v}_p}{v_p} \qquad (v_p = \Omega r_0)$$

Circular DF in a perfect fluid

[VD, Nusser, Buehler ApJ, arXiv: 2111.07366]



Circular DF in a perfect fluid

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Compact circular binary



[Kim+ 08]

Compact circular binary: friction coefficient

[VD, Nusser, Buehler ApJ, arXiv: 2111.07366]



Ultra-light dark matter (ULDM)

ULDM is a Bose-Einstein condensate characterized by a wave function

• Wave-like description (Gross-Pitaevskii-Poisson):

$$i\partial_t \psi = -\frac{\hbar}{2m_a} \nabla_{\mathbf{r}}^2 \psi + \frac{m_a}{\hbar} \Phi \psi$$
$$\nabla_{\mathbf{r}}^2 \Phi = 4\pi G\rho \qquad (\rho \equiv |\psi|^2)$$

• Hydrodynamic description (Madelung):

$$\psi = \sqrt{\rho} e^{i\theta} \qquad \qquad \partial_t \rho + \nabla_r (\rho \mathbf{v}) = 0$$
$$\mathbf{v} = \frac{\hbar}{m_a} \nabla_r \theta \qquad \qquad \partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla_r) \mathbf{v} = -\nabla_r (Q + \Phi)$$

[Ruffini+ 68; Baldeschi+83; Sin 92; Hu+ 99; Chavanis 11; Hui+ 17]

Circular DF in ULDM

[Buehler & VD PRD, arXiv: 2207.13740]

$$\begin{split} I_{\ell m} \propto & \int_{0}^{\infty} dz \, z \frac{j_{\ell}(z) j_{\ell-1}(z)}{z^{4}/R_{\Omega}^{2} - (m+i\epsilon)^{2}} \\ & - \frac{R_{\Omega}}{2} e^{imt} \int_{0}^{\infty} \frac{dz}{z} j_{\ell}(z) j_{\ell-1}(z) \left(\frac{e^{-i(z^{2}/R_{\Omega} - i\epsilon)t}}{z^{2}/R_{\Omega} - m - i\epsilon} + \frac{e^{i(z^{2}/R_{\Omega} + i\epsilon)t}}{z^{2}/R_{\Omega} + m + i\epsilon} \right) \end{split}$$



Circular DF in ULDM

[Buehler & VD PRD, arXiv: 2207.13740]

$$G_{\rm ret}(\vec{r},\tau) \propto \Im \left[\operatorname{erf} \left(\frac{1+i}{2\sqrt{2}} \sqrt{R_\Omega} \frac{r}{\sqrt{\tau}} \right) \right] \qquad (R_\Omega = 2m_a \Omega r_0^2 / \hbar)$$

Friction coefficient for a compact circular binary











Circular DF in isothermal density profiles

w/ Gali Eytan and Robin Buehler

 $P = c_s^2 \rho$

Singular Isothermal Sphere (SIS):

$$\rho_0(r) = \frac{c_s^2}{2\pi G} r^{-2}$$

$$\phi_0(r) = 2c_s^2 \ln(r) + \text{const}$$



Outlook

- Linear response theory DF can be solved analytically in the circular case
- Include eccentric orbits, inhomogeneous backgrounds, other matter "excitations"
- Provide (i) useful physical insights and (ii) a versatile tool to explore DF for a wide range of media and dynamical systems (single perturber, binaries ...)