Hadronic contribution to to the muon g-2: theoretical challenges with the MUonE experiment

F. Piccinini



INFN, Sezione di Pavia (Italy)

QCD@Work, Trani, 18-21 June 2024









$$\vec{\mu}_{\mu} = -\frac{g_{\mu}}{2m_{\mu}}\vec{s}$$

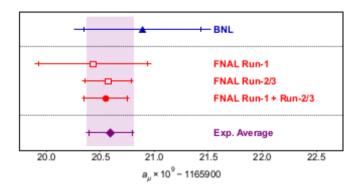
• in QED without quantum fluctuations

$$g_{\mu}=2$$

• with quantum fluctuations: $g_{\mu} \neq 2$ at the 0.1% level

$${f a}_{\mu}\equiv rac{{f g}_{\mu}-{f 2}}{2}$$

Experimental status



 $\hbox{Muon $g-2$ Coll., Phys. Rev. Lett. 126 (2021) 120801; Phys. Rev. Lett. 131 (2023) 161802; arXiv:2402.15410 } \\$

present world average: $\mathbf{a}_{\mu}^{\mathrm{exp}} = 116592059(22) \cdot 10^{-11}$ (0.19 ppm)

M. Knecht, MITP Workshop 3-7/06/2024

 $\Longrightarrow \sim 0.14$ ppm in the (very) near future

Theoretical contributions to a_{μ}

• QED: it accounts for more than 99.99% of the total, with negligible uncertainty at the present precision

Theoretical contributions to \mathbf{a}_{μ}

- QED: it accounts for more than 99.99% of the total, with negligible uncertainty at the present precision
- **ElectroWeak**: calculated up to three loops, with negligible uncertainty $(\sim 153(1) \cdot 10^{-11})$

Theoretical contributions to \mathbf{a}_{μ}

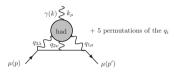
- QED: it accounts for more than 99.99% of the total, with negligible uncertainty at the present precision
- **ElectroWeak**: calculated up to three loops, with negligible uncertainty ($\sim 153(1) \cdot 10^{-11}$)
- QCD: the largest source of uncertainty, due to non-perturbative effects

Theoretical contributions to ${f a}_{\mu}$

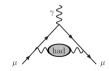
- QED: it accounts for more than 99.99% of the total, with negligible uncertainty at the present precision
- **ElectroWeak**: calculated up to three loops, with negligible uncertainty ($\sim 153(1) \cdot 10^{-11}$)
- QCD: the largest source of uncertainty, due to non-perturbative effects
- possible New Physics?

QCD contributions

Hadronic Light-by-Light (HLxL)



Hadronic Vacuum Polarization (HVP)

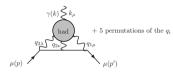


F. Jegerlehner, arXiv:0902.3360

F. Jegerlehner, arXiv:0902.3360

QCD contributions

Hadronic Light-by-Light (HLxL)



Hadronic Vacuum Polarization (HVP)

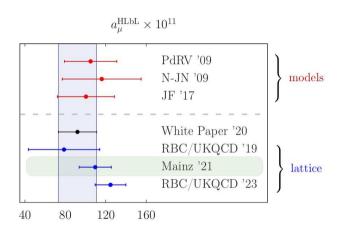


F. Jegerlehner, arXiv:0902.3360

F. Jegerlehner, arXiv:0902.3360

- two approaches for both contributions:
 - first principle calculations with LQCD
 - dispersion relations

HLxL estimates



M. Knecht, MITP Workshop 3-7/06/2024

$${f a}_{\mu}^{f HLxL}=91(19)\cdot 10^{-11}~({\sf WP~2020})$$
 T. Aoyama et al., Phys. Rept. 887 (2020) 1

The dispersive approach for HVP in a nutshell

$$a_{\mu}^{\rm HLO} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^{2} \int_{m_{\pi}^{2}}^{\infty} ds \frac{K(s)R(s)}{s^{2}}$$

$$K(s) = \int_{0}^{1} dx \frac{x^{2}(1-x)}{x^{2}+(1-x)(s/m_{\mu}^{2})}$$

$$R(s) = \frac{\sigma^{0}(e^{+}e^{-} \to \text{hadrons} + \gamma))}{\sigma_{\rm pt}}$$

$$\sigma_{\rm pt} = \frac{4\pi\alpha^{2}}{3s}$$

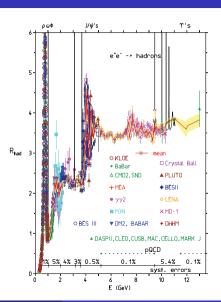
The dispersive approach for HVP in a nutshell

$$a_{\mu}^{\rm HLO} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^{2} \int_{m_{\pi}^{2}}^{\infty} ds \frac{K(s)R(s)}{s^{2}}$$

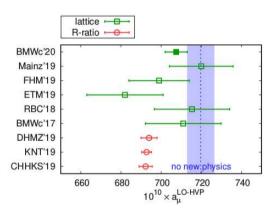
$$K(s) = \int_{0}^{1} dx \frac{x^{2}(1-x)}{x^{2}+(1-x)(s/m_{\mu}^{2})}$$

$$R(s) = \frac{\sigma^{0}(e^{+}e^{-} \to \text{hadrons} + \gamma))}{\sigma_{\text{pt}}}$$

$$\sigma_{\text{pt}} = \frac{4\pi\alpha^{2}}{3s}$$



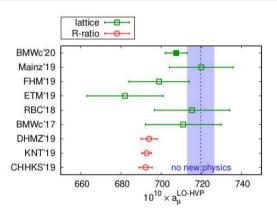
HVP-LO estimates



M. Knecht, MITP Workshop 3-7/06/2024

$${f a}_{\mu}^{HVP-LO}=6931(40)\cdot 10^{-11}$$
 (WP 2020, without BMWc) T. Aoyama et al., Phys. Rept. 887 (2020) 1

HVP-LO estimates



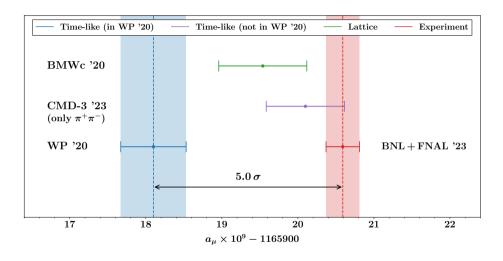
M. Knecht, MITP Workshop 3-7/06/2024

$$\mathbf{a}_{\mu}^{HVP-LO}=6931(40)\cdot 10^{-11}$$
 (WP 2020, without BMWc) T. Aoyama et al., Phys. Rept. 887 (2020) 1

the new CMD-3 result on pion form-factor introduced an additional puzzle

F.V. Ignatov et al, arXiv:2302.08834

Summary of the comparison data-theory



courtesy of A. Gurgone



- * G. Abbiendi, C.M. Carloni Calame, U. Marconi, C. Matteuzzi, G. Montagna, O. Nicrosini, M. Passera, F. Piccinini, R. Tenchini, L. Trentadue, G. Venanzoni,

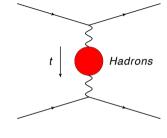
 Measuring the leading hadronic contribution to the muon g-2 via μe scattering

 Eur. Phys. J. C 77 (2017) no.3, 139 arXiv:1609.08987 [hep-ph]
- ★ C. M. Carloni Calame, M. Passera, L. Trentadue and G. Venanzoni, A new approach to evaluate the leading hadronic corrections to the muon g-2 Phys. Lett. B 746 (2015) 325 - arXiv:1504.02228 [hep-ph]

Master formula

$$a_{\mu}^{\mathrm{HLO}} = \frac{\alpha}{\pi} \int_{0}^{1} dx \left(1 - x\right) \Delta \alpha_{\mathrm{had}}[t(x)]$$

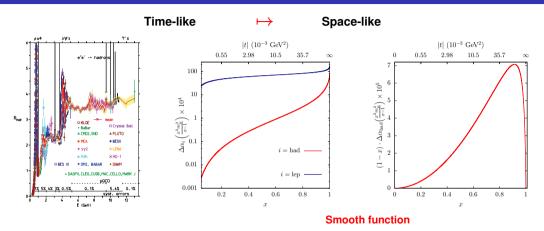
$$t(x) = \frac{x^{2} m_{\mu}^{2}}{x - 1} < 0$$



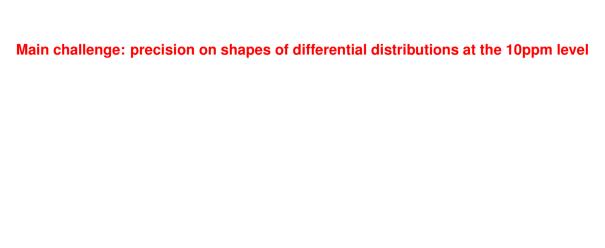
e.a. Lautrup, Peterman, De Rafael, Phys. Rept. 3 (1972) 193

- $\star \Delta \alpha_{\rm had}(t)$ can be directly measured in a (single) experiment involving a space-like scattering process and $\mathbf{a}_u^{\mathrm{HLO}}$ obtained through numerical integration Carloni Calame, Passera, Trentadue, Venanzoni PLB 746 (2015) 325
- * A data-driven, inclusive evaluation of a_u^{HLO} , but with space-like data

From time-like to space-like evaluation of a_{μ}^{HLO}



- \mapsto Time-like: combination of many experimental data sets, control of RCs better than $\mathcal{O}(1\%)$ on hadronic channels required
- → Space-like: in principle, one single experiment, it's a one-loop effect, very high accuracy needed



Main challenge: precision on shapes of differential distributions at the 10ppm level Main sources of systematics on the theory side

Main challenge: precision on shapes of differential distributions at the 10ppm level Main sources of systematics on the theory side

Radiative corrections to the signal

Predictions for Background processes

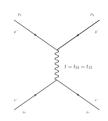
Main challenge: precision on shapes of differential distributions at the 10ppm level Main sources of systematics on the theory side

Radiative corrections to the signal

Predictions for Background processes

High precision Monte Carlo simulation tools required

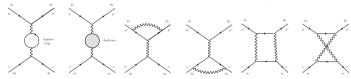
First step towards precision: QED NLO and MC (2018)



analytical expression for tree level

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2}{\lambda(s, m_{\mu}^2, m_e^2)} \left[\frac{(s - m_{\mu}^2 - m_e^2)^2}{t^2} + \frac{s}{t} + \frac{1}{2} \right]$$

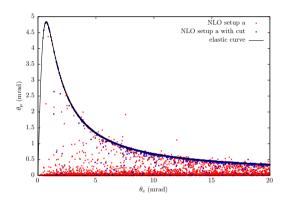
- VP gauge invariant subset of NLO rad. corr.
- factorized over tree-level: $\alpha \rightarrow \alpha(t)$
- QED NLO virtual diagrams and real emission diagrams with exact finite m_e and m_u effects

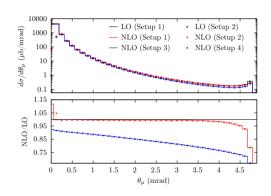


- tree-level Z-exchange important at the 10^{-5} level ($\sim tG_u/4\pi\alpha\sqrt{2}$ in the Fermi theory)
- SM weak RCs at most at a few 10^{-6} level, negligible

Alacevich et al. JHEP 02 (2019) 155

First realistic description of scattering events



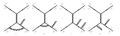


- many points fall out of the $2 \to 2$ correlation curve $\theta_{\mu} \theta_{e}$ because of the radiative events
- NLO QED radiative corrections at the % level, enhanced by exclusive event selections

Second step, towards photonic radiative corrections at NNLO (2020)

exact calculation of corrections along one lepton line with all finite mass effects





two independent calculations, with different IR singularities handling procedures (slicing and subtraction)

Carloni Calame et al., JHEP 11 (2020) 028,

P. Banerjee, T. Engel, A. Signer, Y. Ulrich, SciPost Phys. 9 (2020) 027

- implemented in Mesmer and McMule, perfect numerical agreement
- NNLO with finite mass effects and approximate up-down interference in Mesmer
 - interference of LO $\mu e
 ightarrow \mu e$ amplitude with



NNLO double-virtual amplitudes where at least 2 photons connect the e and μ lines are approximated according to the Yennie-Frautschi-Suura ('61) formalism to catch the IR divergent structure

Second step, photonic radiative corrections at NNLO (2023)

- ullet complete calculation of the amplitude $f^+f^- o F^+F^-$ with $m_f=0$, $m_F
 eq 0$
- "massification" to recover the leading m_e terms, i.e. neglecting powers of m_e^2/Q^2

T. Engel, C. Gnendiger, A. Signer and Y. Ulrich, JHEP 02 (2019) 118

Y. Ulrich, PoS RADCOR2023 (2024) 077

FKS^ℓ subtraction scheme

T. Engel, A. Signer, Y. Ulrich, JHEP 01 (2020) 085

 Next-to-soft stabilisation, to obtain numerical stability in real-virtual corrections with soft and/or collinear photon configurations

T. Engel, A. Signer, Y. Ulrich, JHEP 04 (2022) 097; T. Engel, JHEP 07 (2023) 177

- with the above ingredients
 - ullet NNLO calculation neglecting terms of $\mathcal{O}(m_e^2/Q^2)$ in <code>McMule</code>

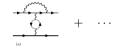
A. Broggio et al., JHEP 01 (2023) 112

NNLO virtual leptonic pairs (vacuum polarization insertion) (2021)

- any lepton (and hadron) in the VP blobs
- interfered with $\mu e \rightarrow \mu e$ or $\mu e \rightarrow \mu e \gamma$ amplitudes









• interfered with $\mu e
ightarrow \mu e$ amplitude



2-loop integral evaluated with dispersion relation techniques in Mesmer

used e.g. in the past for Bhabha: Actis et al., Phys. Rev. Lett. 100 (2008) 131602; Carloni Calame et al., JHEP 07 (2011) 126

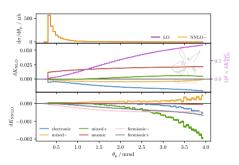
$$\frac{g_{\mu\nu}}{q^2+i\epsilon} \rightarrow g_{\mu\nu} \frac{\alpha}{3\pi} \int_{4m_\ell^2}^{\infty} \frac{dz}{z} \frac{R_\ell(z)}{q^2-z+i\epsilon} = g_{\mu\nu} \frac{\alpha}{3\pi} \int_{4m_\ell^2}^{\infty} \frac{dz}{z} \frac{1}{q^2-z+i\epsilon} \left(1+\frac{4m_\ell^2}{2z}\right) \sqrt{1-\frac{4m_\ell^2}{z}}$$

• 2-loop integral evaluated (also) with hyperspherical method in McMule

M. Fael, JHEP02 (2019) 027

NNLO order of magnitude

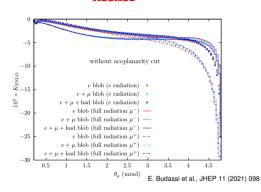
McMule



A. Broggio et al., JHEP 01 (2023) 112

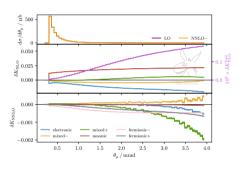
• NNLO corrections at the $10^{-4} - 10^{-3}$ level

Mesmer

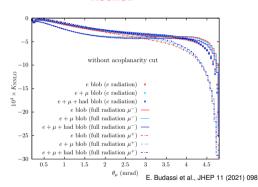


NNLO order of magnitude

McMule



Mesmer

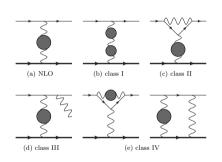


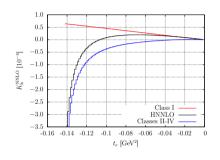
A. Broggio et al., JHEP 01 (2023) 112

- NNLO corrections at the $10^{-4} 10^{-3}$ level
- eventually fixed order calculations need to be matched to resummation of higher order corrections, through PS techniques (e.g. BaBaYaga) or YFS techniques (e.g. KKMC/SHERPA)

NNLO hadronic contributions (2019)

using the dispersion relation approach





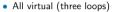
Fael, Passera, Phys. Rev. Lett. 122 (2019) 192001

- corrections of the order of 10⁻⁴
- hyperspherical integration method to calculate hadronic NNLO corrections, where the hadronic vacuum
 polarization is employed in the space-like region (used in McMule)

 M. Fael, JHEP02 (2019) 027

Towards N³LO on the electron line

Y. Ulrich, N³LO kick-off workstop/thinkstart, Durham, 3-5 August 2022





• Single real emission (two loops)



• Double real emission (one loops)



Triple real



M. Fael, MUonE Collaboration Meeting, 16/05/2023, CERN

 this contribution will allow improved perturbative predictions and more reliable theoretical uncertainty estimates

Recent progress

the three-loop form factor with finite fermion mass is now available

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev. Lett 128 (2022) 172003

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev.D 106 (2022) 034029

M. Fael, F. Lange, K. Schönwald, M. Steinhauser, Phys. Rev.D 107 (2023) 094017

All order subtraction scheme FKS^ℓ availale

T. Engel, A. Signer, Y. Ulrich, JHEP 01 (2020) 085

 very recent generalisation of the LBK theorem to multi-photon emission ⇒ extension of next-to-soft stabilisation to multiple radiation

T. Engel, JHEP 03 (2024) 004

ullet real-virtual-virtual corrections recently recalculated with $m_e o 0$

S. Badger, J. Krys, R. Moodle, S. Zoia, JHEP 11 (2023) 041

V.S. Fadin, R.N. Lee, JHEP 11 (2023) 148

Fixed target experiment \Longrightarrow bound electron effects

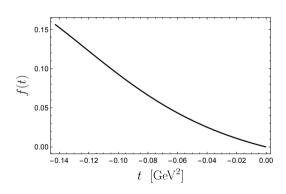
Fixed target experiment ⇒ bound electron effects

very recently estimated

R. Plestid and M.B. Wise, arXiv:2403.12184

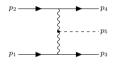
• for C $\frac{1}{\sigma}\frac{d\sigma}{dt}=\frac{1}{\sigma^0}\frac{d\sigma^0}{dt}\left(1-Kf(t)\right)$

• $K = 4.5 \cdot 10^{-4}$, scaling as $1/Z_A$



Backgrounds

- pion pair production forbidden kinematically with the available \sqrt{s}
- single π^0 production possible

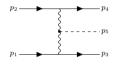


• π^0 production calculated and shown to be well below 10^{-5} w.r.t. $\mu e o \mu e$

E. Budassi et al., PLB 829 (2022) 137138

Backgrounds

- pion pair production forbidden kinematically with the available \sqrt{s}
- single π^0 production possible

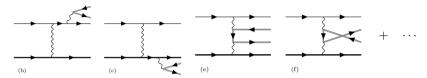


• π^0 production calculated and shown to be well below 10^{-5} w.r.t. $\mu e \to \mu e$

E. Budassi et al., PLB 829 (2022) 137138

- lepton pair production
 - $\mu^{\pm}e^{-} \rightarrow \mu^{\pm}e^{-}\ell^{+}\ell^{-}$
 - $\mu^{\pm}N \rightarrow \mu^{\pm}N\ell^{+}\ell^{-}$

it also contributes at NNLO accuracy



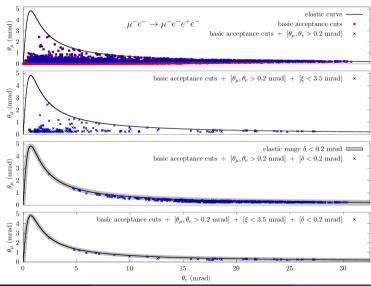
• the emission of an extra electron pair $\mu e \to \mu e \ e^+ e^-$ is potentially a dramatically large background, because of the presence of "peripheral" diagrams which develop powers of collinear logarithms upon integration

G. Racah, Il Nuovo Cimento 14 (1937) 83-113; L.D. Landau, E.M. Lifschitz, Phys. Z. Sowjetunion 6 (1934) 244; H.J. Bhabha, Proc. Roy. Soc. Lond. A152 (1935) 559;

R.N. Lee, A.A. Lyubyakin, V.A. Smirnov, Phys. Lett. B 848 (2024) 138408

• $\mu^\pm e^- o \mu^\pm e^- \ell^+ \ell^-$ calculated with finite mass effects and implemented in Mesmer

simulation of $5\cdot 10^5$ points of $\mu^\pm e^- o \mu^\pm e^- \ell^+ \ell^-$



Real pair emission from scattering on nucleus: $\mu^{\pm}N ightarrow \mu^{\pm}N\ell^{+}\ell^{-}$

G. Abbiendi et al., Phys. Lett B854 (2024) 138720

- it can mimic the signal if one particle is not reconstructed or two tracks overlap within resolution
- ullet cross section scaling $\sim Z^2$

ullet GEANT4: "for the process of e^+e^- pair production the muon deflection is neglected"

A.G. Bogdanov et al., IEEE transactions on nuclear science, 53, n. 2, April 2006

 \implies a dedicated calculation implemented in the Monte Carlo generator ${ t Mesmer}$

- approximation: scattering on the external nucleus field
- finite extension of the nucleus through a form factor

$$F_Z(q) = \frac{1}{Ze} \int_0^\infty dr \, r^2 \rho_Z(r) \frac{\sin(qr)}{qr}$$

- q: momentum transferred to the nucleus
- ρ_Z : nuclear charged density
- different models for charge density

J. Heeck, R. Szafron, Y. Uesaka, PRD 105 (2022) 053006

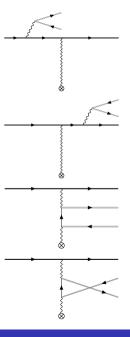
- $F_Z(q) = 1$ (conservative)
- 1 parameter Fermi model (1pF)

$$\rho_Z(r) = \frac{\rho_0}{1 + \exp\frac{r - c}{z}}$$

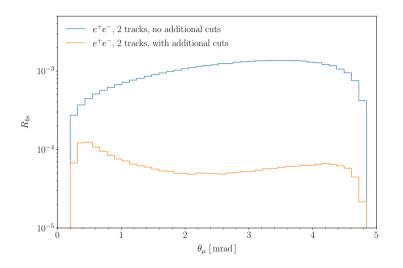
Fourier Bessel expansion (FB)

$$\rho_Z(r) = \sum_{k=0}^{n} a_k j_0 \left(\frac{k\pi r}{R}\right), \quad r \ge R$$
$$= 0 > R$$

modified-harmonic oscillator model



Background/signal ratio



G. Abbiendi, E. Budassi, C.M. Carloni Calame, A. Gurgone, F.P., Phys.Lett.B 854 (2024) 138720

Possible New Physics contamination in the $\Delta \alpha(t)$ determination?

A. Masiero, P. Paradisi and M. Passera, Phys. Rev. D102 (2020) 075013

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, JHEP 05 (2020) 053

- Effects of **heavy** $(M_{NP} \gg 1 \text{ GeV})$ NP mediators investigated through EFT with dim-6 operators
 - excluded (at the 10^{-5} level) by existing data
- Effects of **light** $(M_{NP} \le 1 \text{ GeV})$ NP mediators investigated with spin-dependent general models
 - spin—0 NP mediators (ALPs)
 - spin-1 NP mediators (Dark Photons, light Z' vector bosons)
 - excluded (at the 10^{-5} level) by existing data

HVP determination with MUonE data will be robust against New Physics

Possible New Physics studies with MUonE (in complementary regions to $\Delta \alpha_h$)

- ullet interesting proposals for NP searches at MUonE (new light mediators) in 2 o 3 processes
 - invisibly decaying light Z' in $\mu e \to \mu e Z'$

Asai et al., Phys. Rev. D106 (2022) 5

- a relevant background can be $\mu e \to \mu e \pi^0$, in addition to $\mu e \to \mu e \gamma$
- long-lived mediators with displaced vertex signatures $\mu e \to \mu e A' \to \mu e e^+ e^-$

Galon et al., Phys.Rev.D 107 (2023) 095003

• through scattering off the target nuclei $\mu N \to \mu N X \to \mu N e^+ e^-$

Grilli di Cortona and E. Nardi, Phys. Rev. D105 (2022) L111701

Summary

- Given its precision requirements, MUonE represents a challenge for
 - QFD corrections
 - background calculation
- at present we have two independent Monte Carlo tools, Mesmer and McMule featuring
 - NLO QED corrections
 - NNLO QED corections from single lepton legs
 - YFS inspired approximation to the full NNLO QED in Mesmer
 - full NNLO QED with electron "massification" in McMule
 - pair production in Mesmer
 - $\mu^{\pm}e^{-} \rightarrow \mu^{\pm}e^{-}\ell^{+}\ell^{-}$
 - $\mu^{\pm}N \rightarrow \mu^{\pm}N\ell^{+}\ell^{-}$
- efforts for N³LO started
- work in progress for matching with higher order QED corrections

Past MUonE topical meetings

- MUonE theory workshops
 - Theory Kickoff Workshop, Padova, 4-5 September 2017
 - MITP Workshop, Mainz 19-23 February 2018
 - 2nd Workstop/ThinkStart, Zürich, 4-7 February 2019
 - N³LO kick-off workstop/thinkstart IPPP Durham, 3-5 August 2022
 - MITP Workshop, Mainz 14-18 November 2022
 - MITP Workshop, Mainz 03-07 June 2024
- Five General MUonE Collaboration Meetings

A collection of references on calculation developments

- → Carloni Calame et al., PLB 746 (2015), 325
- → Abbiendi et al., EPJ C77 (2017), 139
- → Mastrolia et al., JHEP 11 (2017) 198
- → Di Vita et al., JHEP 09 (2018) 016
- → Alacevich et al., JHEP 02 (2019) 155
- --- Fael and Passera, PRL 122 (2019) 19, 192001
- → Fael, JHEP 02 (2019) 027
- → Engel et al., JHEP 02 (2019) 118
- → Engel et al., JHEP 01 (2020) 085
- → Carloni Calame et al., JHEP 11 (2020) 028
- → Banerjee et al., SciPost Phys. 9 (2020), 027
- → Banerjee et al., EPJC 80 (2020) 6, 591
- → Budassi et al., JHEP 11 (2021) 098
- → Balzani et al., PLB 834 (2022) 137462

- → Bonciani et al., PRL 128 (2022) 2, 022002
- Budassi et al., PLB 829 (2022) 137138
- → Engel et al., JHEP 04 (2022) 097
- → Fael et al., PRL 128 (2022) 172003
- → Fael et al., PRD 106 (2022) 034029
- → Broggio et al., JHEP 01 (2023) 112
- → Fael et al., PRD 107 (2023) 094017
- → Engel, JHEP 07 (2023) 177
- → Badger et al., JHEP 11 (2023) 041
- → Fadin and Lee., JHEP 11 (2023) 148
- → Ahmed et al., JHEP 01 (2024) 010
- → Engel, JHEP 03 (2024) 004
- → Abbiendi et al., PLB 854 (2024) 138720
- → Plestid and Wise, arXiv:2403.12184

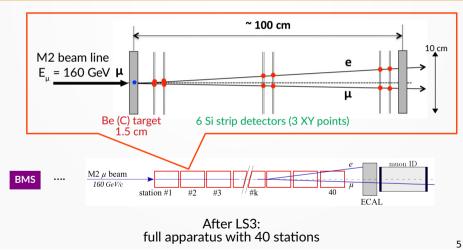
THANK YOU!

Few slides from R. Pilato for an exp. update

MITP Workshop, 03-07/06/2024, Mainz, Germany

The experimental apparatus





Achievable accuracy



40 stations (60 cm Be) + 3 years of data taking =
$$({}^{\sim}4x10^{12} \text{ events} \atop E_e > 1 \text{ GeV})$$
 = $({}^{\sim}0.3\% \text{ statistical} \atop \text{accuracy on } a^{\text{HLO}}$

Main challenge: keep systematic accuracy at the same level of the statistical one

Systematic uncertainty of 10 ppm in the signal region

Main systematic effects:

Longitudinal alignment (<10 μm)

theoretical predictions

- Knowledge of the beam energy (few MeV)
- Multiple scattering (<1%)
- Angular intrinsic resolution
- Non-uniform detector response

6

Staged approach towards the full experiment



- 2017: dedicated test beam to study multiple scattering.
- 2018: test beam to study elastic scattering properties and event selection.
- 2021: first joint test CMS-MUonE with a few 2S modules prototypes (parasitic).
- 2022:
 - test with 1 tracking station.
 - · test the calorimeter.
- 2023: test with 2 tracking stations + calorimeter.
- 2025: run with a scaled version of the complete apparatus:
 - 3 tracking stations;
 - Calorimeter;
 - Muon ID;
 - Beam Momentum Spectrometer (BMS).

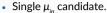
7

Fulvio Piccinini (INFN, Pavia)

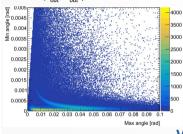
TB 2023 μ -e elastic scattering event selection



Max angle [rad]



• μ_{out} , e_{out} pair candidate.



Initial selection

• Loose χ^2_{vtx} cut. • $|Z_{vtx} - Z_{target}| < 3$ cm. • Acoplanarity cut (elastic events are planar). • 0.003

Work in progress:

- Exploit dedicated MC generators to study the backgrounds.
- Study the main sources of systematic error using tracker data:
 - Angular intrinsic resolution;
 - Beam energy scale.

15

Run 2025



- MUonE recently submitted a proposal for a phase 1 of the experiment to the SPSC, concerning a small scale version of the final apparatus.
- If approved, MUonE will request 4 weeks of data taking in 2025.

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

MUonE Phase 1 Experiment Proposal



April 25, 2024

Proposal for phase 1 of the MUonE Experiment

The MUonE Collaboration

16

Fulvio Piccinini (INFN, Pavia)