

## Introduction

The QCD phase diagram features a variety of phases in the  $\mu_B - T$  plane. Nonetheless, determining the equations of state and phase-transition lines from the first-principles QCD is difficult.



## Holographic Hardwall Model

$$S_{0} = \frac{N_{c}^{2}}{8\pi^{2}L^{3}} \int d^{5}x \sqrt{g} \left(R - 2\Lambda\right) - \frac{\theta N_{f}N_{c}}{24\pi^{2}L} \int d^{5}x \sqrt{g} \frac{d^{5}x}{4\pi^{2}L} \frac{d^{5}x}{4\pi^{2}L} \int d^{5}x \sqrt{g} \frac{d^{5}x}{4\pi^{2}L} \frac{d^{5}x}{4\pi^$$

L is the radius of AdS.

Parameter  $\theta$ : Dilaton coupling and different choices of compactification. The cosmological constant  $\Lambda$ : Asymptotical  $AdS_5$  geometries.

## Methodology

- An IR cutoff  $z_0$  is implemented to break the conformal symmetry.
- The equations of motion are solved with suitable **boundary conditions**.
- The holographic renormalization procedure is applied to determine the on-shell action.
- The phase with the lowest free energy and the equations of state are identified.
- Phenomenological methods are employed to establish the phase diagram in physical units.

## Solutions/Phases

The diagonal metric ansatz:

 $ds^{2} = -rac{g(z)}{h(z)}L^{2}dt^{2} + rac{L^{2}}{z^{2}}d\vec{x}^{2} + rac{L^{2}}{g(z)}dz^{2}$ 

and the gauge field ansatz  $A_0 = \phi(z)$ . There are four boundary conditions: one each for g and h, and two for  $\phi$ . There are three possible geometries:

- **Thermal AdS:** zero denisty, confined phase.
- 2. Charged Black Hole: finite density, deconfined phase.
- 3. Charged AdS: finite density, confined phase.

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## **Charged AdS**

vanishing densities

The charged AdS is a **horizonless geometry** in the bulk that represents a finite-density confined phase of the boundary theory.

$$g = z^2 \left( 1 - \left( 1 - \frac{g_0}{z_0^2} \right) \frac{z^4}{z_0^4} + \frac{\theta N_f}{9N_c} Q^2 (z^6 - z^4 z_0^2) \right);$$

Scalar Glueball mass:  $\implies z_0^{-1} = 290 \text{ MeV}$ 

Two free parameters: Q and  $g_0 = g(z_0)$ 

 $\rho$ -meson mass:  $\implies \bar{g}_0 \simeq 6$ .

Perturbation of the Vector field, with boundary conditions given by  $a_{\nu}(k,0) =$  $\partial_z a_\nu(k, z_0) = 0.$ 

$$z\partial_z(rac{g}{z^3}\partial_z V(z))+k^2V(z)$$

The equation of state relating charge density Q with  $\mu$  is determined by:

• Simple boundary condition:  $\phi(z_0) = 0 \implies 0$ The deconfinement phase transition at T = 0

$$\left(\frac{\mu\alpha_g}{M_g}\right)^2 > \frac{6N_c}{\theta N_f} \left(1 + \sqrt{4 + 3}\right)$$

 $arepsilon = p - rac{N_c^2}{4\pi^2 z_c^6} (z_0^2 + g_0)$ The equation of state: Physically motivated boundary condition:

$$\rho = \frac{\theta N_f N_c}{12\pi^2} Q \qquad g_0 = \bar{g}_0 + \frac{8\pi^2}{N_c^2} \left( p - \frac{2\pi^2 z_0^2}{\theta N_f N_c} \rho^2 \right)$$

- Nambu-Jona-Lasinio: Temperature and chemical potential dependent constituent quark mass.
- Van der Waals: Isospin-symmetric nuclear matter (in-medium ChPT). The equation of state:  $\varepsilon = \alpha + \beta \sqrt{p} + \gamma p$

# Phase diagram



A similar analysis with a softwall model suggests low density confined phase.

$$h=z^4; \quad \phi(z)=\pm\mu+Qz^2;$$

Rinaldi: 2018

- = 0

$$Q = \frac{\mu}{z_0^2}$$

$$B\alpha_\rho \left(\frac{M_g}{M_\rho}\right)^2$$

Klevansky:1992 Fiorilla:2012

## **Condensate solutions**

A complex scalar field  $\psi$  is included to break  $U(1)_B$  symmetry spontaneously through the action.

$$S = S_0 - \lambda_s \int d^5 x \sqrt{g} \left( |D\psi| \right)$$

- Scaling dimension  $\Delta \quad m^2 L^2 = \Delta (\Delta d)$
- Asymptotic analysis  $\psi \sim \psi_+ z^{\Delta_+} + \psi_- z^{\Delta_-} + \dots$
- Breaking symmetry spontaneously  $\psi_{-}=0$ ;  $\psi_+ \neq 0$

Phase	sign	Onset $\mu$ (MeV)	0
VdW	-ve	1426.8	
VdW	+ve	1235.4	
NJL	-ve	_	
NJL	+ve	1139.7	

## **Conclusions and Discussion**

- A simple 5-D holographic hardwall model provided low density confined phase resulting in a complete phase diagram.
- The probe analysis of baryon condensates in the confined phase at zero temperature identifies a critical chemical potential beyond which the condensate disappears.
- An ongoing study for backreacted condensate solutions with physical boundary conditions is suggesting a similar result for vanishing condensate at high chemical potential.
- Improve the model by incorporating the Running coupling, Isospin, and Chiral condensate.
- D3/D7 hardwall model with backreaction can provide finite density confined phase with baryons as solitons.





## References

[1] Akash Singh and K. P. Yogendran. Confined phases at finite density in the Hardwall model. arxiv, 2407.xxxxx. [2] Akash Singh and K. P. Yogendran. Phases of nuclear matter from AdS Hardwall models. arxiv, 24xx.xxxx.

cf: Costa: 2019