

Collision integrals for first-order phase transitions

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C. Branchina, A. Conaci, S. De Curtis, LDR, A. Guiggiani, A. Gil Muyor, G. Panico based on: JHEP 03 (2022) 163, JHEP 05 (2023) 194, JHEP 05 (2024)

Thermal History of the Universe

Phase transitions are crucial events in the evolution of the Universe

- the SM predicts two of them (the two phases are smoothly connected (cross over)) no strong breaking of thermal equilibrium
 - no distinctive experimental signatures
- New physics may change the nature of these PhTs, or add new ones



First-order EWPhT

New physics may provide first order phase transitions New Physics at finite temperature

- a barrier in the potential may be generated from tree-level deformations, thermal or quantum effects
- the field tunnels from false to true minimum at $T = T_n < T_c$
- the transition proceeds through bubble nucleation



• interesting experimental signatures (eg. gravitational waves)



Bubble nucleation

Bubble dynamics can produce gravitational waves and baryogenesys



Key features of a first-order PhT

- the nucleation temperature T_n
- the strength α
- the (inverse) time duration of the transition β/H
- the speed of the bubble wall v_w
- the thickness of the bubble wall L_w

equilibrium quantities

non-equilibrium quantities

Gravitational waves and the efficiency of the EW-baryogenesis crucially depend on them

EWBG is typically efficient for slowly-moving walls. Recent results show efficiency also for fast-moving walls [Dorsch, Huber, Konstandin, 2021]

BG at the EW scale also possibile with $v_w \sim 1$ [Azatov, Vanvlasselaer, Yin, 2021]

GWs are maximised for fast-moving walls

GW from a first-order PhT



Dynamics of the bubble wall

System setup: scalar field + plasma



- The bubble wall drives plasma out of equilibrium
- Interactions between plasma and wall front produce a friction
- If the friction and pressure inside the bubble balance we can realise a steady state regime (terminal velocity reached)

Dynamics of the bubble wall

out-of-equilibrium

For each particle species in the plasma

 $f(p, z) = f_v(p, z) + \delta f(p, z)$ equilibrium

I. Scalar field equation

$$\phi' \Box \phi - V'_T = \sum N_i \frac{dm^2}{dz} \int \frac{d^3p}{(2\pi)^3 2E_p} \delta f(p)$$

2. Boltzmann equation

$$\left(\frac{p_z}{E}\partial_z - \frac{(m^2)'}{2E}\partial_{p_z}\right)(f_v + \delta f) = -\mathcal{C}[f_v + \delta f]$$

two competing effects: $(m^2)' \sim (\phi^2)'$ and C

we assume a planar wall and a steady state regime

Approaches to the Boltzmann equation

To deal with the collision term, previous approaches made assumptions on the shape of $\delta f(p,z)$ in momentum space

- Fluid approximation [1]
- Extended fluid approximation [2]
- New formalism [3]

[1] Moore, Prokopec, 1995[2] Dorsch, Huber, Konstandin, 2022[3] Laurent, Cline, 2020

[1] and [2] dubbed "old formalism" (OF) in the following

1... the $\partial_{p_z} \delta f$ term neglected

2!!! Boltzmann equation integrated with a set of (not unique) weights

Alternative methods

- Expansion of δf in a polynomial basis [4]
- Holographic approach [5]

[4] Laurent, Cline, 2022[5] Bigazzi, Caddeo, Canneti, Cotrone

Full solution to the Boltzmann equation

- * We propose a new method to solve the Boltzmann equation without imposing any ansatz for δf De Curtis, LDR, Guiggiani, Gil Muyor, Panico, 2022
- We developed an algorithm to solve the coupled system of bubble wall and Boltzmann equations, thus getting v_w , L_w , etc.

De Curtis, LDR, Guiggiani, Gil Muyor, Panico, 2023

Key features

- New approach (spectral decomposition) to deal with collision integrals
- No term in the Boltzmann equation is neglected
- Iterative routine where convergence is achieved in few steps

Structure of the collision integral

The collision integral yields two classes of terms:

$$C[f] = \mathcal{Q}(p) \frac{\delta f(p)}{f'_{v}(p)} + f_{v}(p) \langle \delta f \rangle$$

- perturbations do not appear inside the integral: easy to handle
- perturbations are integrated (brackets): very challenging

Structure of the collision integral

 $\langle \delta f \rangle = \mathcal{O}[\delta f] \sim \int \mathcal{O}k \mathcal{K}(p, \cos \theta_p, k, \cos \theta_k) \delta f(k, \cos \theta_k, z)$

Brackets can be seen as the application of Hermitian operators \mathcal{O} on perturbations

Main idea: decompose the operators on the basis of their eigenfunctions

Exploit rotational invariance of the collision integral: the kernel is block diagonal on the angular momentum basis

$$\mathcal{K}(p,\cos\theta_p,k,\cos\theta_k) = \sum_{l} \frac{2l+1}{2} P_l(\cos\theta_p) P_l(\cos\theta_k) \mathcal{G}_l(p,k)$$
$$\mathcal{G}_l(p,k) = \sum_{i} \lambda_i^{(l)} \varphi_i(p) \varphi_i(k)$$

kernels can be (numerically) evaluated only once huge improvement in time performance (~ 2 orders of magnitude)

Integrated friction



Determination of the wall speed



- Important corrections from out-of-equilibrium perturbations
- Sizeable corrections given by the W bosons
- Peak corresponding to the Jouguet velocity

De Curtis, LDR, Guiggiani, Gil Muyor, Panico, 2024

Conclusions and outlook

Conclusions:

- Fully quantitative solution to the Boltzmann equation
- New spectral method based on multipole decomposition of the collision integral
- Computation of v_w
- Quantitative and qualitative differences with previous approaches
- Important impact of out-of-equilibrium friction

Work in progress:

- inclusion of $1 \rightarrow 2$ and $2 \rightarrow 1$ plasma processes in the collision integrals
- improving the description of W bosons in the ultra-soft regime



The **effective kinetic theory** depends on the **momentum** of the particles in the plasma

Hot plasma of weakly interacting particles (separation of scales)



• Point-like particles interacting through local collisions

• Screening and damping effects are important

• Quantum corrections are relevant