



ALICE

# Measurement of multiple (multi-)strange hadron production in small collision systems with ALICE

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on behalf of the ALICE Collaboration



*6° incontro nazionale di Fisica Nucleare  
26-28 Feb.*

1. Università degli Studi di Torino



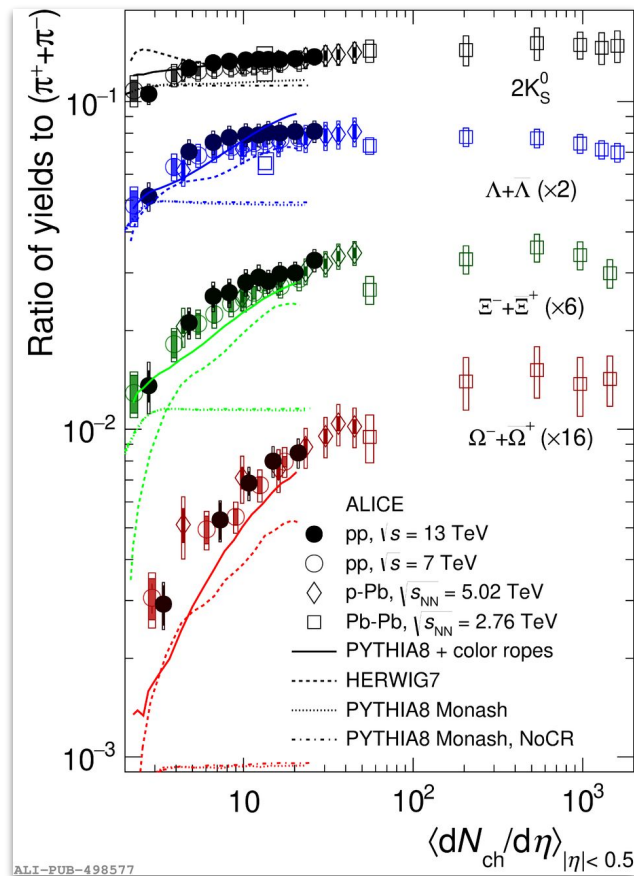
2. INFN Torino





## Strangeness Enhancement (SE):

- at the LHC  $S/\pi$  increases as a function of multiplicity independently on the collision energy and system
- hierarchy: enhancement proportional to the strangeness content in the hadron  
 $\rightarrow \Omega > \Xi > \Lambda$



ALICE-PUB-498577

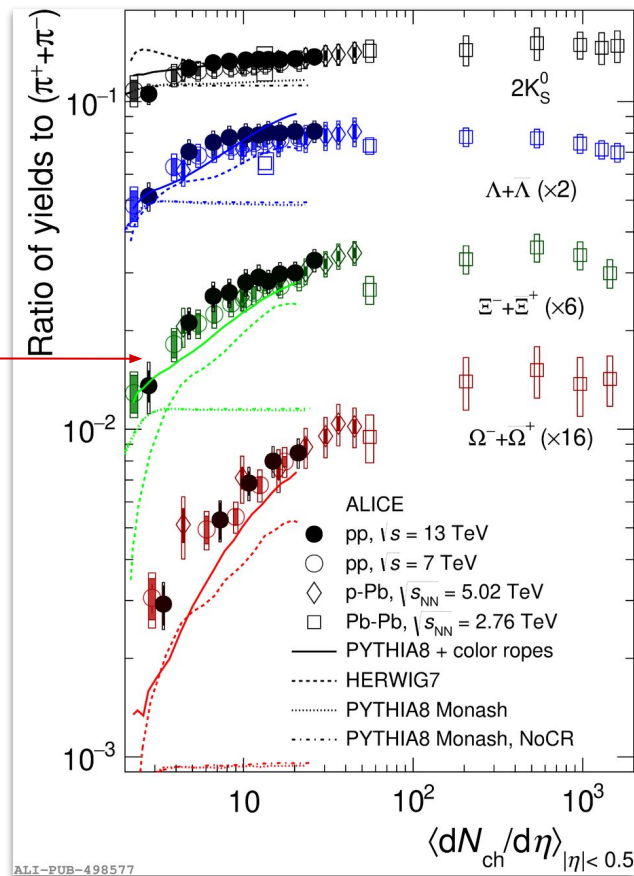


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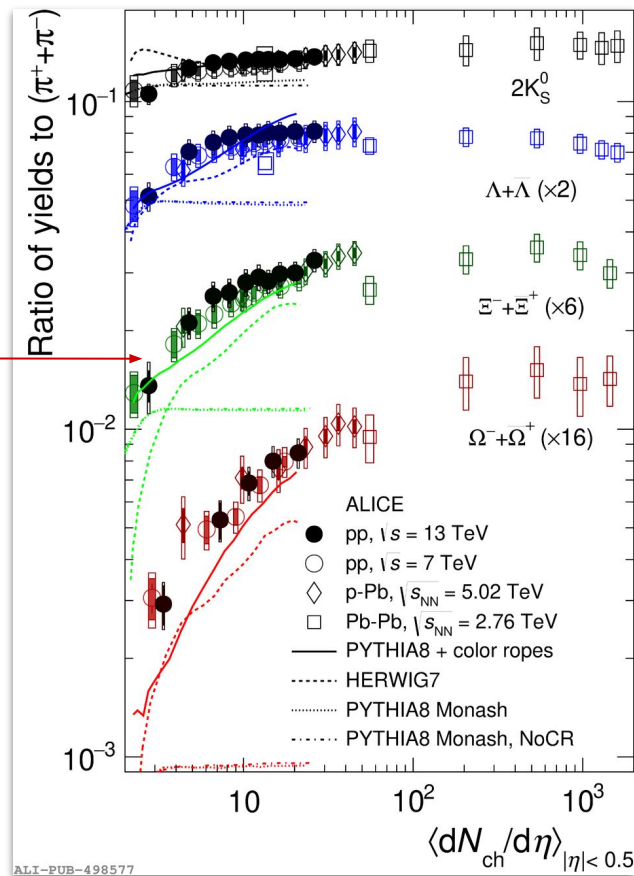
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### New in this talk:

Strange particle multiplicity distribution ( $P(n_s)$ ) for  $K_S^0$ ,  $\Lambda$ ,  $\Xi$ ,  $\Omega$

- extend beyond the average of the distribution
- unique opportunity to test the connection between charged and strange particle multiplicity production
- stronger constraints to models



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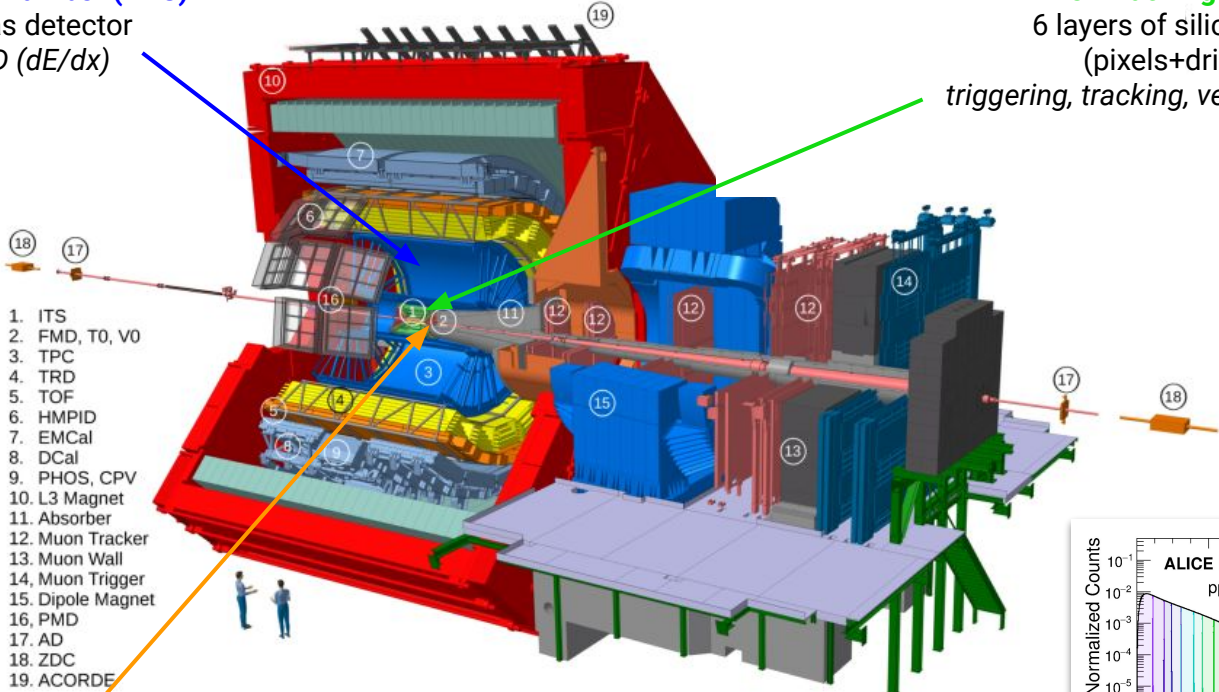


**Time Projection Chamber (TPC)**

Large area gas detector  
tracking, PID ( $dE/dx$ )

**Inner Tracking System (ITS)**

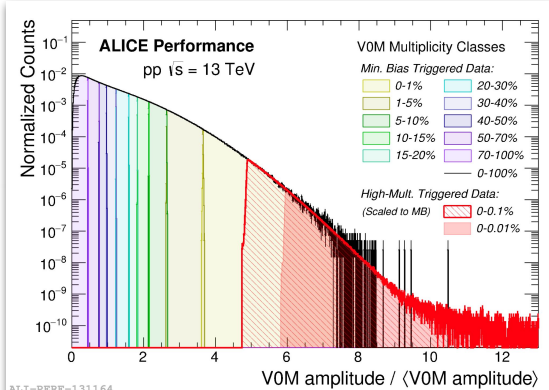
6 layers of silicon detectors  
(pixels+drift+strips)  
triggering, tracking, vertexing, PID ( $dE/dx$ )



**V0 detectors (V0A, V0C)**

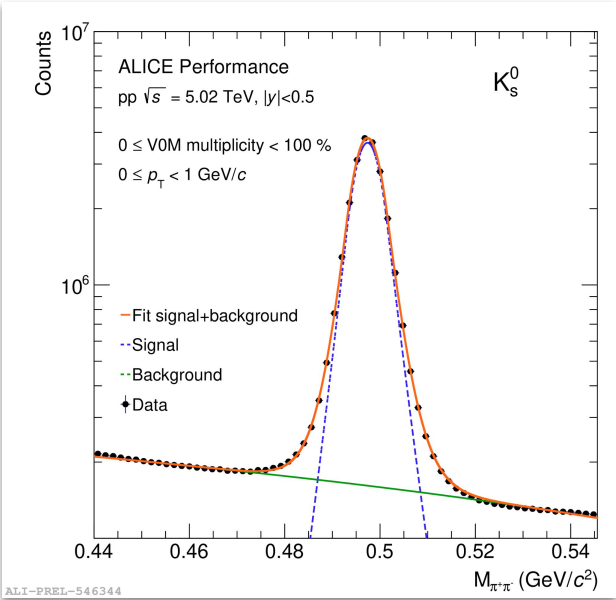
Forward-rapidity arrays of scintillators  
triggering, particle multiplicity  
estimation

- V0M amplitude distribution is divided into percentiles to define multiplicity classes
- The corresponding charged particle multiplicity at mid-rapidity was constructed





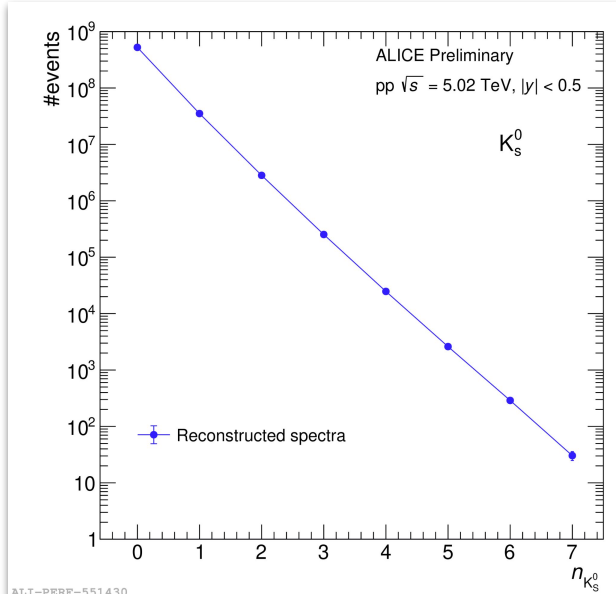
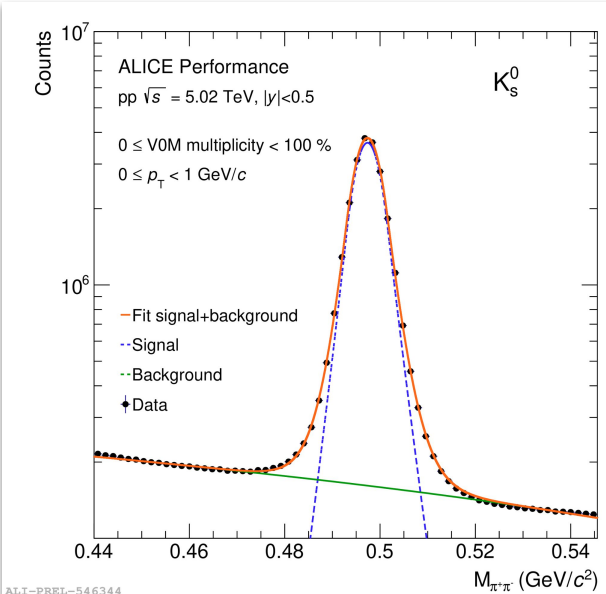
Analysis based on counting the number of strange particles event-by-event in pp collisions at  $\sqrt{s} = 5.02$  TeV



Each candidate **weighted by  $P(\text{Sig})$**  or  **$P(\text{Bkg})$**  estimated by **1D invariant mass fit** in  $p_T$ /multiplicity bins



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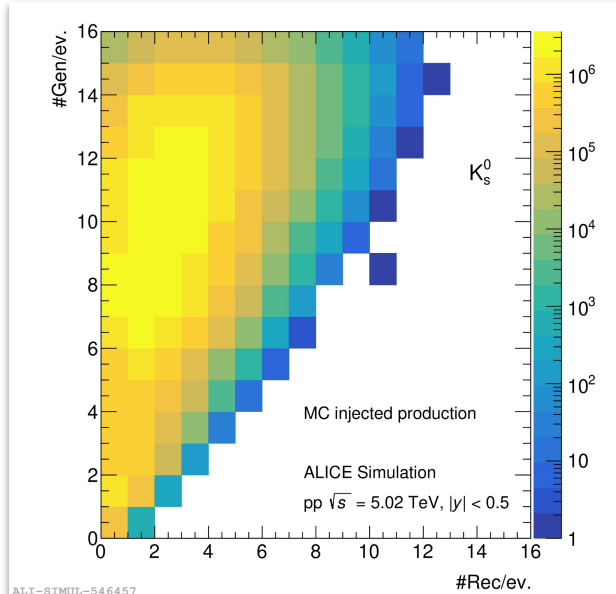
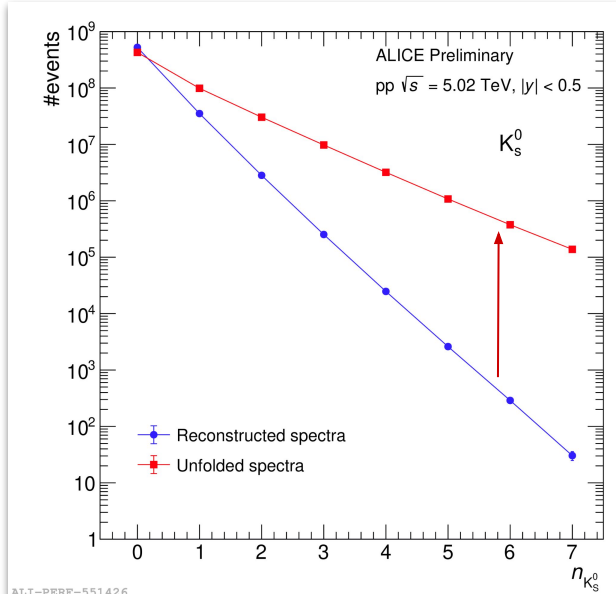
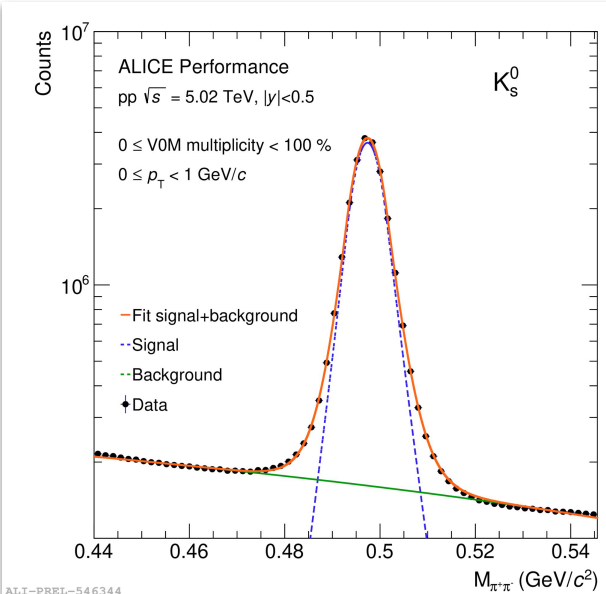
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**Weights** associated to each of the  $N$  candidates **combined** to obtain:  $P(\text{all-sig}), \dots, P(\text{all-bkg})$

For each event: full probability spectrum spanning from 0 to  $N$



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Correction for detector response (MC simulation: measured  $p_T$  distribution)

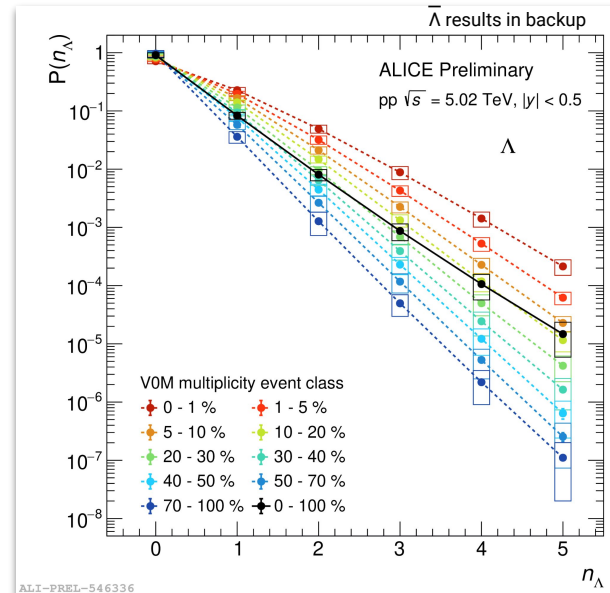
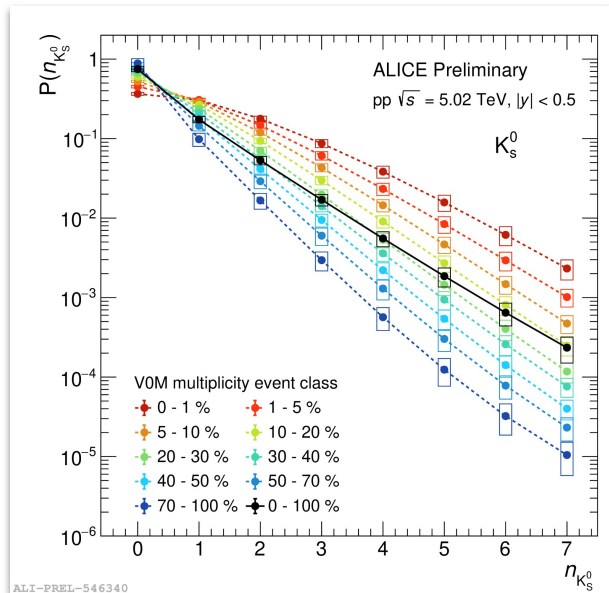
**Bayesian unfolding** procedure applied





Probability to produce  $n$  particle ( $n$  up to 7 for  $K_S^0$ , 5 for  $\Lambda$ ) of a given species per event

- $P(n_{S>1})$  increases with the event charged-particle multiplicity
- Spanning across large ranges of strange/multiplicity variations, all the way to very “extreme” situations (e.g. 7  $K_S^0$  at low average charged-particle multiplicity, 0  $K_S^0$  at high average charged-particle multiplicity)



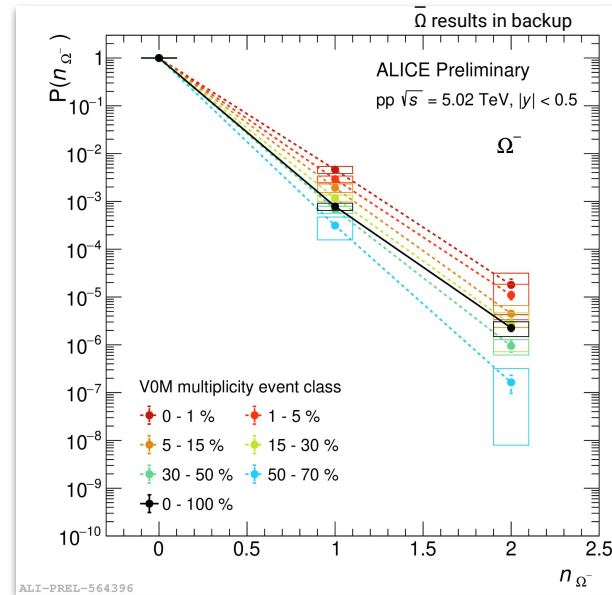
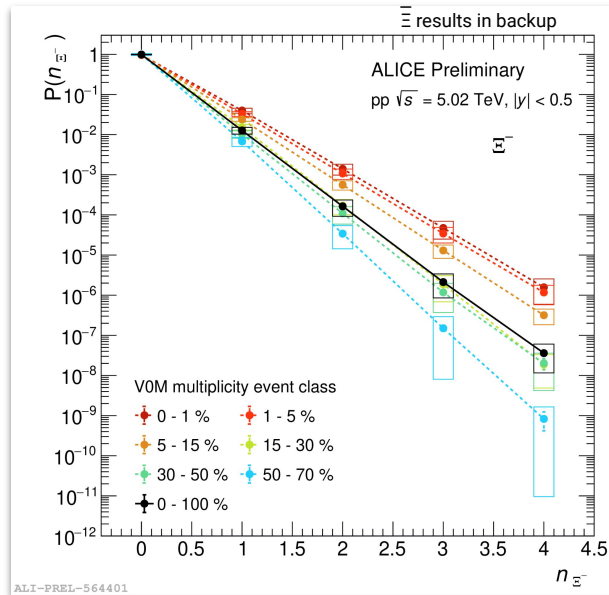
Unique opportunity to test the connection between charged and strange particle multiplicity production

NOTE: multiplicity can fluctuate in each VOM bin and  $\langle dN_{ch}/d\eta \rangle$  can significantly change for events with small/large  $n_S$



Probability to produce  $n$  particle ( $n$  up to 4 for  $\Xi$ , 2 for  $\Omega$ ) of a given species per event

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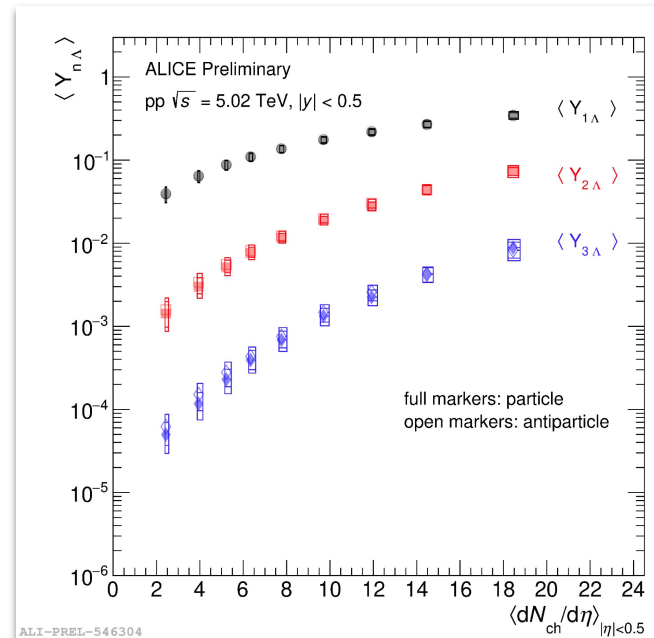
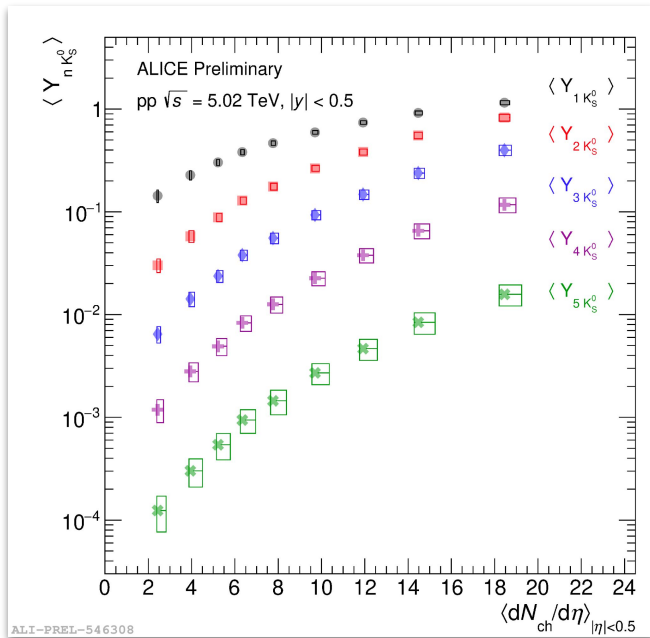
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The distribution allows to calculate the **average probability for the production yield of 1, 2, 3, ... particles/event:**

$$\langle Y_{k-part} \rangle = \sum_{n=k}^{\infty} \frac{n!}{k!(n-k)!} P(n)$$



The increase with multiplicity of the probability for multiple strange hadrons is more than linear

NOTE: very good agreement between  $\langle Y_{1-part} \rangle$  and previous results ([1],[2])

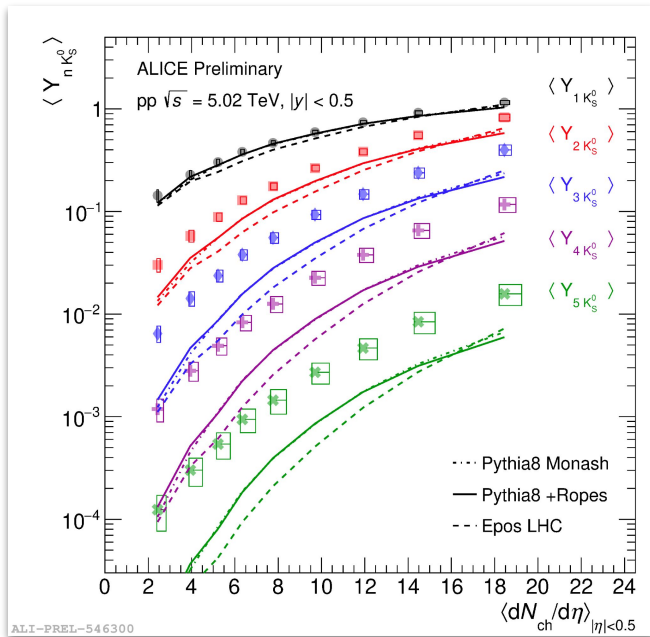
[1] ALICE, *Nature Physics* v13, pages 535–539 (2017)

[2] ALICE, *Eur. Phys. J. C* 80, 167 (2016)

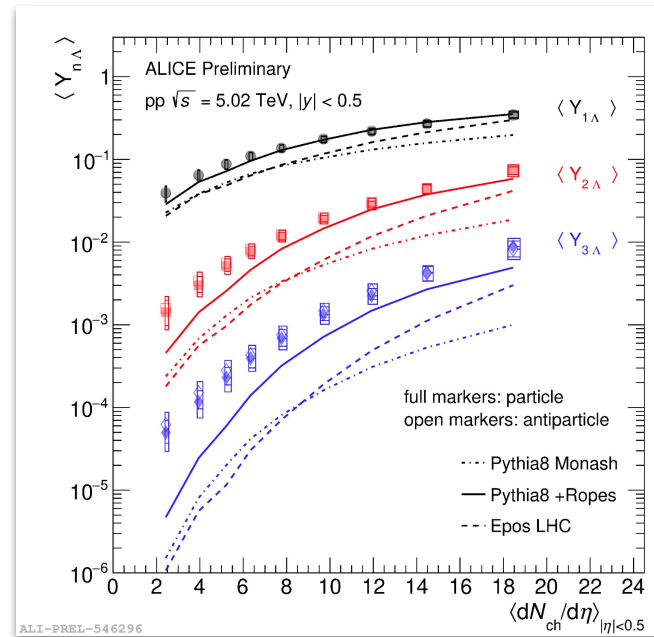


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No difference between [Pythia8 Monash](#) and Ropes for  $K_S^0$ : [Pythia8 + Ropes](#) (with QCD-CR) tends to increase baryons



Ropes approaches the data at high multiplicity for  $\Lambda$

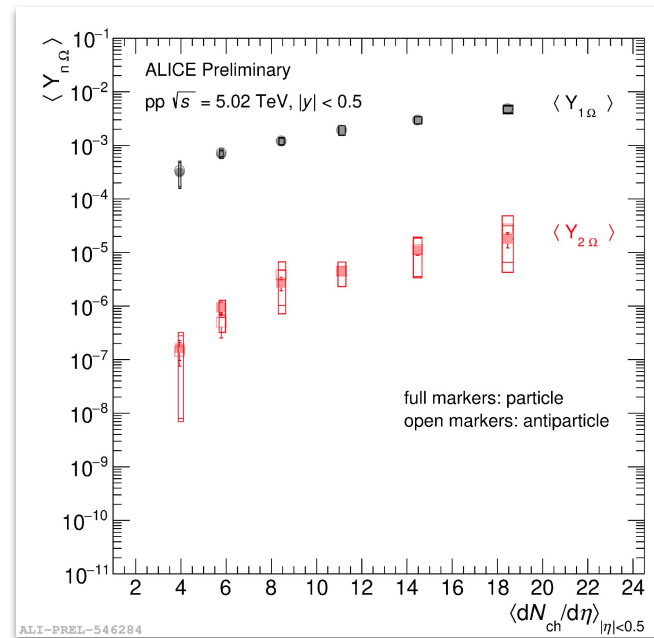
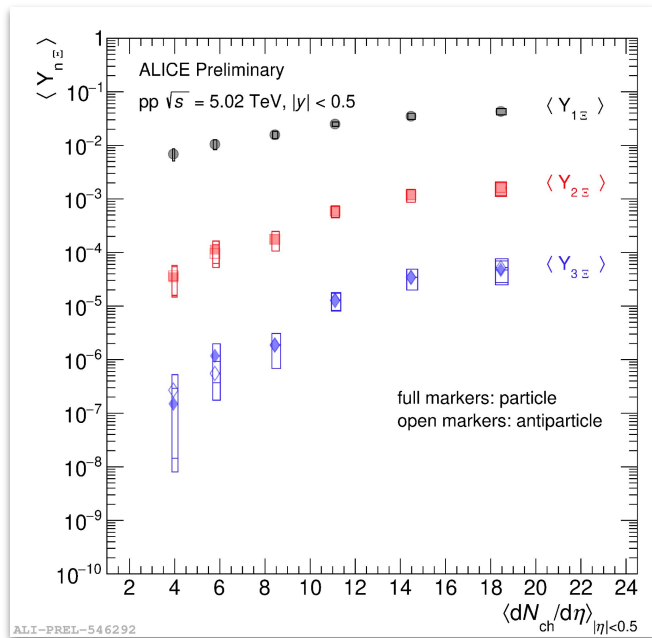
[Epos LHC](#) has a better agreement with the data at high multiplicity, but departs from the trend at low multiplicity

For both particles the agreement gets worse as  $n$  increases



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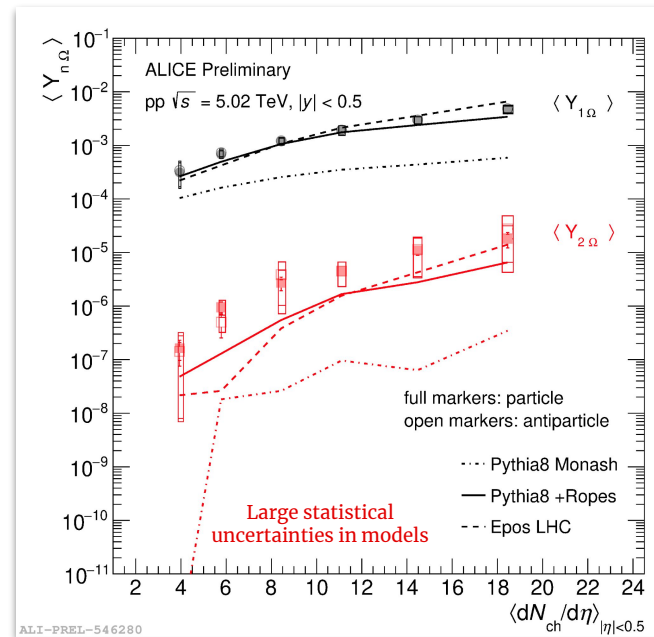
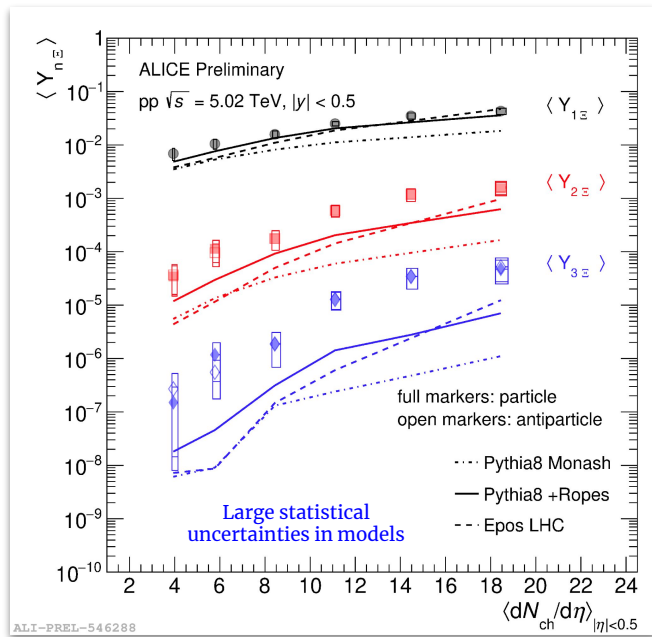
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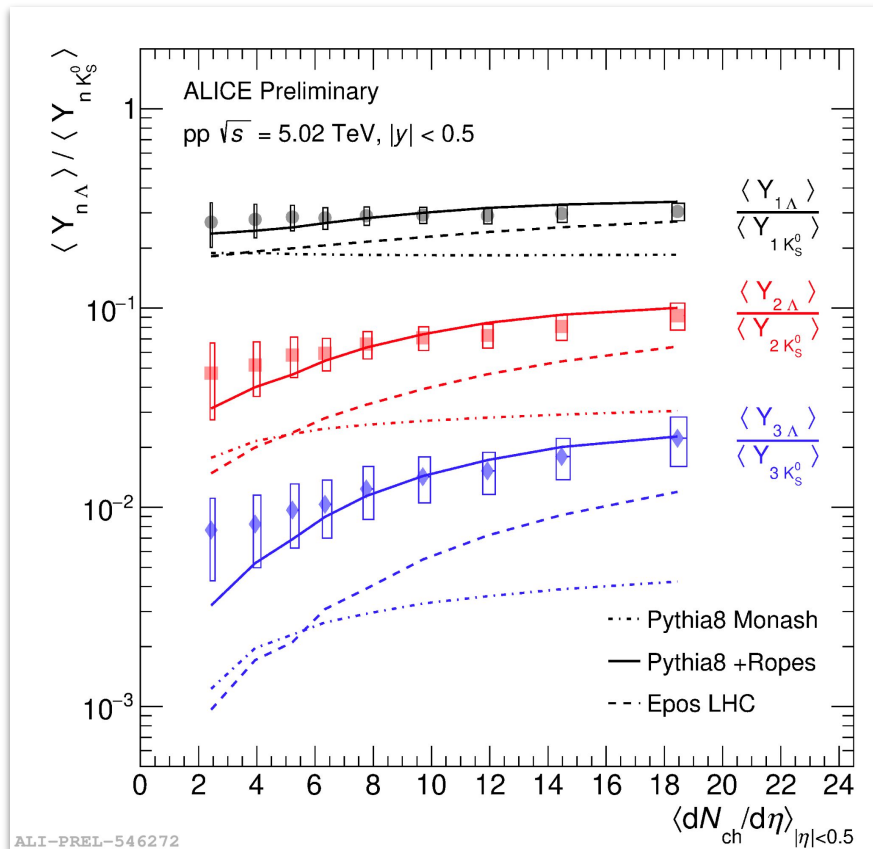
Pythia8 + Ropes approaches the data at high multiplicity

Epos LHC does a rather good job at high multiplicity, but shows larger discrepancy at low multiplicity



Very important to decouple strangeness-related from baryon-related effects!

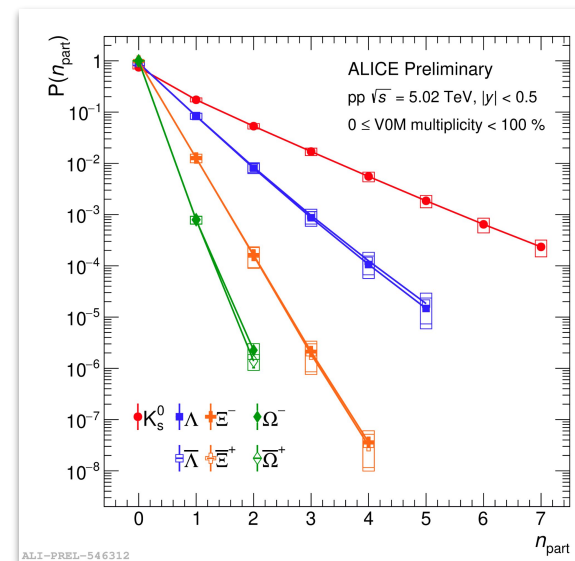
- Increase of  $\Lambda/K_s^0$  VS multiplicity when looking at multiple production!
- Possibly in all strange-hadron/ $\pi$  VS multiplicity plots we have a strangeness-related AND a baryon-related contribution to the enhancement
- Baryon-related effect well reproduced by Ropes (with QCD-CR) at high multiplicity





## First measurements of (multi-)strange particle multiplicity distribution ( $P(n_s)$ )

- perfect benchmark to MC models to test the interplay between charged particle and strange particle multiplicity
- it is a relevant extension of  $\langle dN/dy \rangle$  studies, as it tests at a higher order the connection between global and local characteristics of the analyzed event
- 2- and 3-  $\Lambda/K_S^0$  yield ratios increase with multiplicity (baryon-related effect)
- Comparison to model comparison:
  - for  $K_S^0$  Pythia8 Monash and Ropes are equal across multiplicity
  - for baryons all trends are rather well reproduced by Pythia8 + Ropes: very good improvement in this model except for purely-strangeness effect, all other models fail

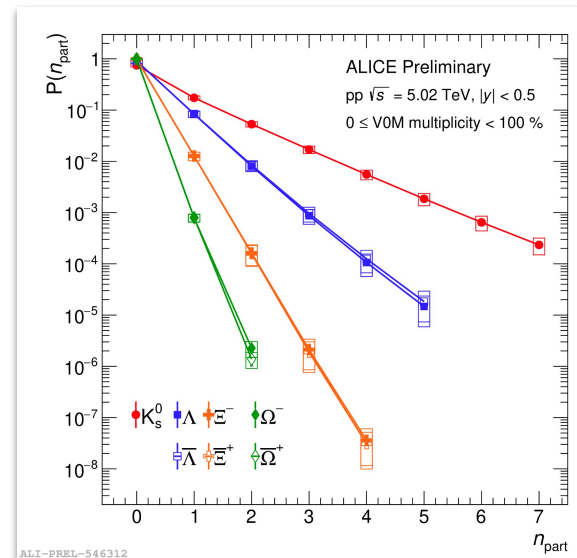






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## Outlook:

- different yield ratios under investigation to show the relative importance of baryon and strangeness effects
- obtained results can be used to study strangeness enhancement at its extremes (yield ratios with  $\Delta S > 3$ )
- Run3 will allow to have larger statistics (3/4 orders of magnitude higher) useful for cascade analyses



# Backup slides



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**iterative procedure based on the Bayes' theorem using a picture of causes C ("true values") and effects E ("observed values")**

$$P(C_i | E_j) = \frac{P(E_j | C_i) \cdot \pi(C_i)}{\sum_{i=1}^{n_C} P(E_j | C_i) \cdot \pi(C_i)}$$

$P(E_j | C_i)$  estimated by using Monte Carlo (response matrix)  
 $P(C_i | E_j) \rightarrow$  probability that different  $C_i$  were responsible for the observed effect  $E_j \rightarrow$  GOAL  
 $\pi(C_i) \rightarrow$  prior probabilities (initially arbitrary, but updated on subsequent iterations)

- Choosing a prior distribution in order to apply Bayes' theorem  $\rightarrow$  posterior probability matrix obtained
- Applied to "observed spectra"  $\rightarrow$  1<sup>st</sup> estimation of the corrected spectra
- The corrected spectra obtained in the previous step becomes the prior probability and the correction proceeds as before
- Procedure is re-iterated until stability is achieved (**regularization parameter:  $n_{iter}$** )

$$\hat{n}(C_i) = \frac{1}{\epsilon_i} \sum_{j=1}^{n_E} n(E_j) \cdot P(C_i | E_j) = \sum_{j=1}^{n_E} M_{ij} \cdot n(E_j)$$

**expected number of events in the cause bin  $i$**

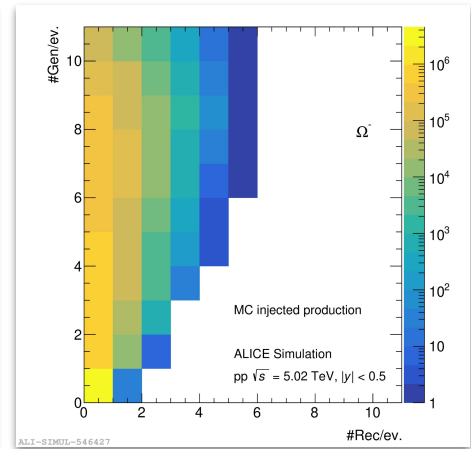
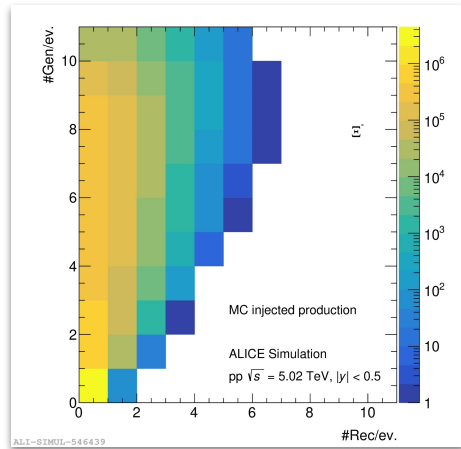
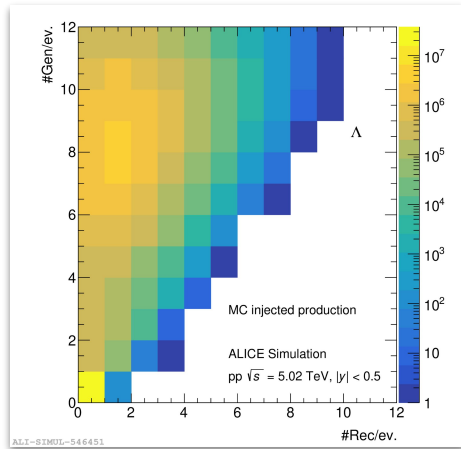
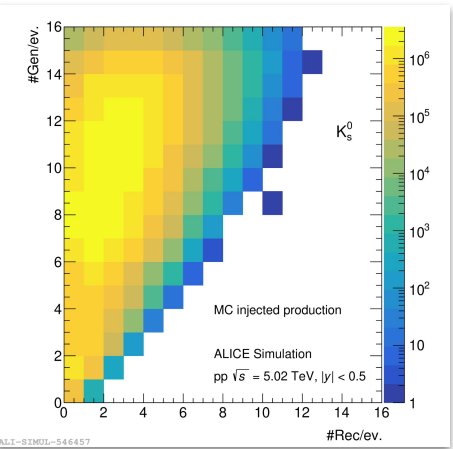
$\rightarrow M_{ij}$  is the unfolding matrix: 
$$M_{ij} = \frac{P(E_j | C_i) \cdot \pi(C_i)}{\epsilon_i \cdot \sum_{l=1}^{n_C} P(E_j | C_l) \cdot \pi(C_l)}$$

$\rightarrow n(E_j)$  measurements (effects)

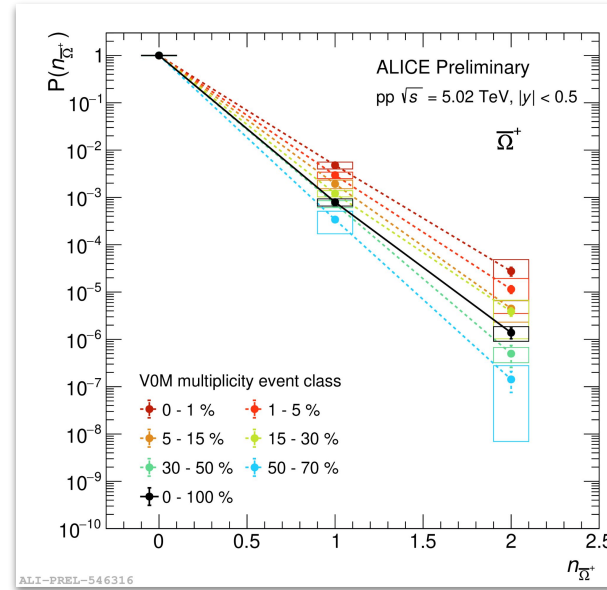
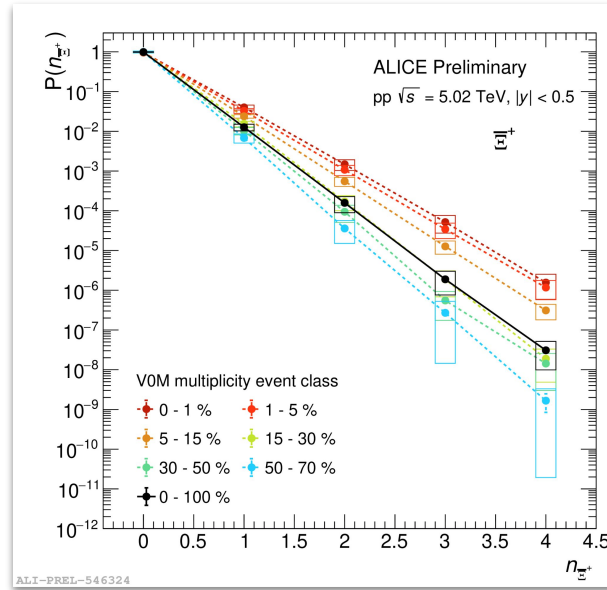
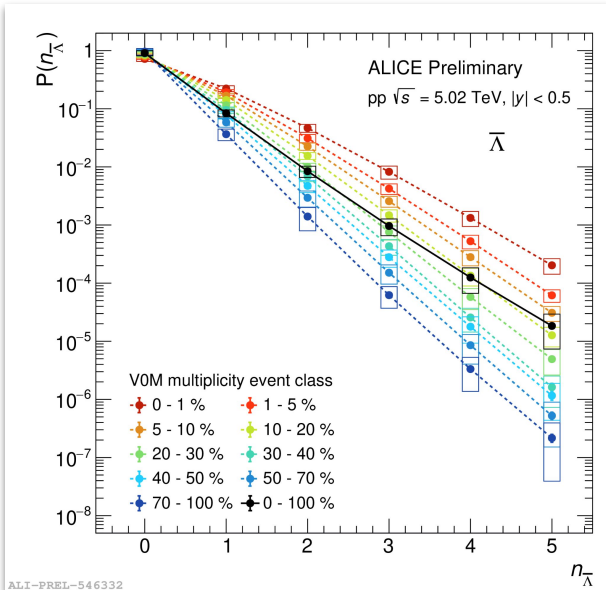
$\rightarrow \epsilon_i$  efficiencies

**unfolding errors: covariance matrix**

$$V(\hat{n}(C_k), \hat{n}(C_l)) = \sum_{i,j=1}^{n_E} \frac{\partial \hat{n}(C_k)}{\partial n(E_i)} V(n(E_i), n(E_j)) \frac{\partial \hat{n}(C_l)}{\partial n(E_j)}$$



Moving from  $K_s^0$  to  $\Omega$  particle the response matrices are increasingly "squeezed" toward a low number of reconstructed particles/event



Probability to produce  $n$  particle ( $n$  up to 5 for  $\bar{\Lambda}$ , 4 for  $\bar{\Xi}$ , 2 for  $\bar{\Omega}$ ) of a given species per event

Spanning across large ranges of strange/multiplicity variations, all the way to very “extreme” situations

Unique opportunity to test the connection between charged and strange particle multiplicity production

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