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DI TRENTO



## Effects of nuclear matter properties in neutron star mergers

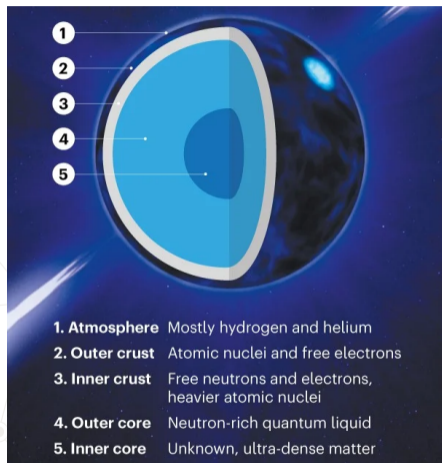
Federico Maria Guercilena (Università di Trento)



with **Maximilian Jacobi** (Friedrich-Schiller-Universität Jena)  
*M. Jacobi, F.M. Guercilena, S. Huth, G. Ricigliano, A. Arcones,  
A. Schwenk, MNRAS, 527, 3, 2023*

Sesto Incontro Nazionale di Fisica Nucleare, Trento - 26/02/2024

# BACKGROUND: NEUTRON STARS



- 1. Atmosphere** Mostly hydrogen and helium
- 2. Outer crust** Atomic nuclei and free electrons
- 3. Inner crust** Free neutrons and electrons, heavier atomic nuclei
- 4. Outer core** Neutron-rich quantum liquid
- 5. Inner core** Unknown, ultra-dense matter

Image from [Mann, 2020]

- ▶ Mass:  $\sim 1$  to  $\sim 2.3 M_{\odot}$
- ▶ Density:  $\sim 1$  to  $\sim 10 n_0$
- ▶ Radius:  $\sim 12$  km
- ▶ Composition: mostly neutrons

The core is well modeled as nuclear matter.

👉 Details of nuclear interaction in these conditions aren't well constrained → Neither is high-density EOS

Note: "isolated" NS too can give us precious information (e.g.  $M_{\max}$ )

$$M_{\max} \simeq 1.97 M_{\odot} [\text{Demorest et al., 2010}]$$

$$M_{\max} \simeq 2.01 M_{\odot} [\text{Antoniadis et al., 2013}]$$

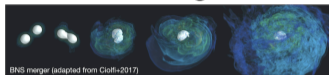
$$M_{\max} \simeq 2.08 M_{\odot} [\text{Fonseca et al., 2021}]$$

$$M_{\max} \simeq 2.35 M_{\odot} [\text{Romani et al., 2022}]$$

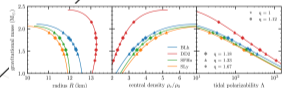
# BACKGROUND: BINARY NEUTRON STARS



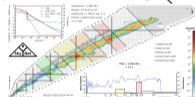
## BNS mergers



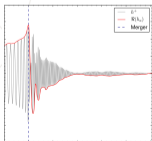
### EOS of dense matter



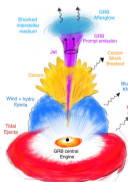
### r-process nucleosynthesis



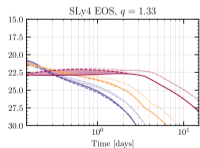
### Gravitational waves



### sGRB



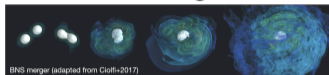
### Kilonovae



# BACKGROUND: BINARY NEUTRON STARS



## BNS mergers

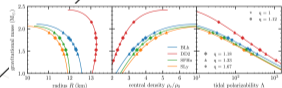


A spectacular confirmation from the events GW170817 (GW signal), GRB170817A (sGRB) and AT2017gfo (kilonova), the first multimessenger detection of a BNS merger

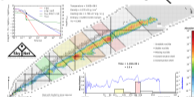


The start of multimessenger astronomy

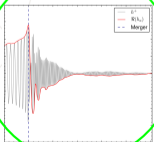
## EOS of dense matter



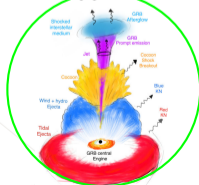
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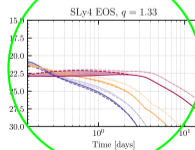
## Gravitational waves



## sGRB



## Kilonovae





# METHODS: SELECTING NUCLEAR PROPERTIES



The problem in a nutshell: can we pinpoint nuclear properties starting from BNS observables? Our approach:

Construct an EOS from a Skyrme parametrization of the energy per nucleon...

$$\frac{E}{A}|_{T=0} = \frac{3\hbar}{10m^*} (3\pi^2 n)^{2/3} [(1-x)^{5/3} + x^{5/3}] + [a + 4bx(1-x)]n + cn^\delta - x\Delta$$
$$x = \frac{n_p}{n}$$

$$\hbar^2/(2m^*) = \hbar^2/(2m) + \alpha n$$

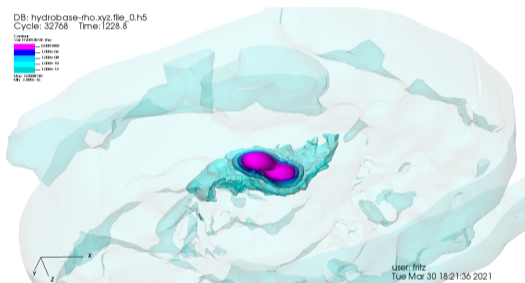
and fix coefficients  $a, b, c$  and  $\delta$  to target specific values of  $B, K, E_{\text{sym}}$  and  $n_0$ .  $m^*$  is varied directly, we cannot control  $L$ .

EOS	$m^*$	$B$	$K$	$E_{\text{sym}}$	$L$	$n_0$
LS175	1.0	16.0	175	29.3	73.7	0.155
LS220	1.0	16.0	220	29.3	73.7	0.155
LS255	1.0	16.0	255	29.3	73.7	0.155
$m_{0.8}^*$	0.8	16.0	220	29.3	79.3	0.155
$m_S^*$	0.634	16.0	220	29.3	86.5	0.155
$(m^*K)_S$	0.634	16.0	281	29.3	86.5	0.155
$(m^*KE)_S$	0.634	16.0	281	36.9	109.3	0.155
SkShen	0.634	16.3	281	36.9	109.4	0.145
Shen	0.634	16.3	281	36.9	110.8	0.145

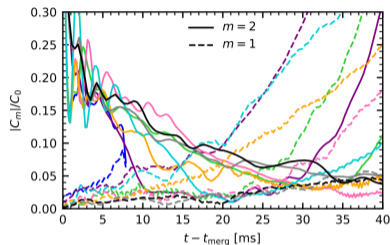
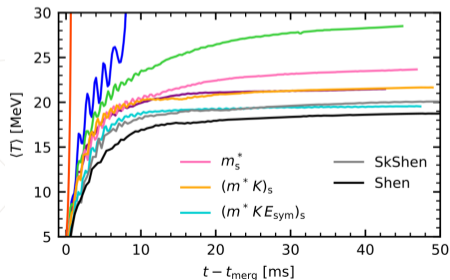
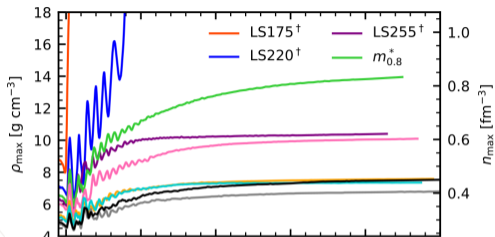
# METHODS: SIMULATING BNS MERGERS



- ▶ One  $1.365 M_{\odot}$  equal mass simulation for each EOS
- ▶ BNS merger simulations performed with the **EinsteinToolkit** [Zlochower et al., 2022]
- ▶ Full GR (BSSN0K formulation of the EfE)
- ▶ Valencia formulation of Euler equations (WhiskyTHC code) [Radice and Rezzolla, 2012]
- ▶ Leakage (emission) + M0 (absorption) neutrino transport module [Galeazzi et al., 2013, Radice et al., 2016]



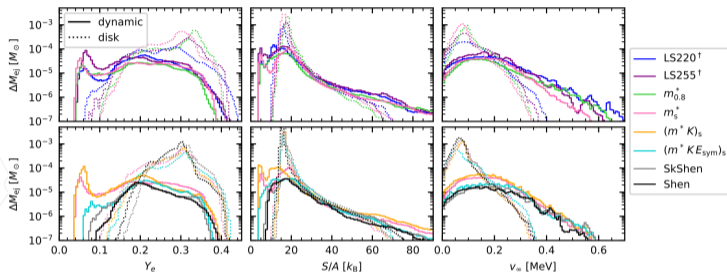
# RESULTS: STRUCTURE OF THE MERGER REMNANT



$$C_m = \int e^{-im\phi} \rho dx dy$$

- ▶ merger remnant structure and fate are tied to  $P|_{n_0}$  and  $dP/dn|_{n_0}$
- ▶ the incompressibility  $K$  has particular influence
- ▶ the final density correlates tightly with pressure at ( $6 \sim 7 \cdot 10^{14}$  g/cm<sup>3</sup>)
- ▶ the crossover from the  $m = 2$  to  $m = 1$  mode is also influenced by the incompressibility

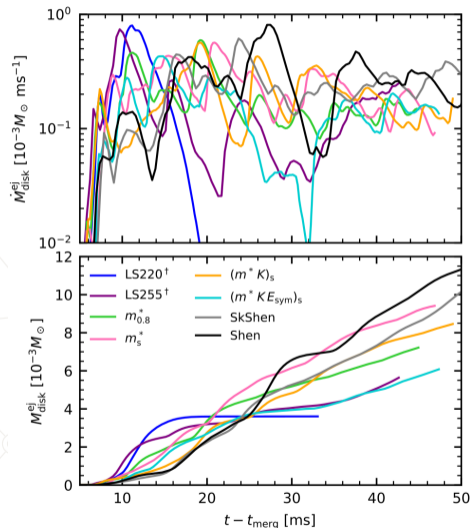
# RESULTS: EJECTA PROPERTIES



- ▶ ejecta are divided in dynamical (satisfying the Geodesic criterion) and disk ejecta (satisfying the Bernoulli criterion)
- ▶ dynamical ejecta are further divided in tidal ( $Y_e < 0.1$ ) and shock-heated ( $Y_e \geq 0.1$ )

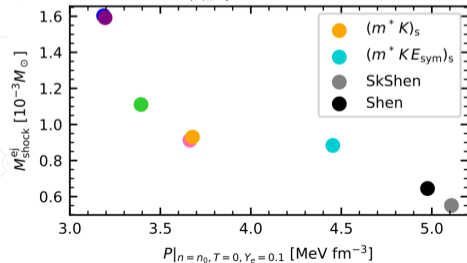
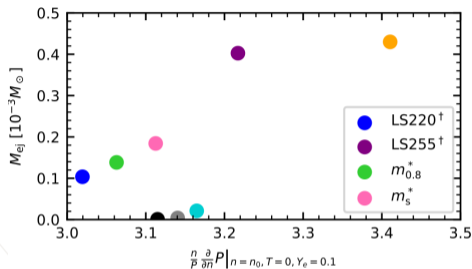
- ▶ ejecta properties are different between dynamical (lower  $Y_e$ , higher  $v_\infty$  and broader  $s$ ) and disk
- ▶ not much difference between EOS (no visible trend)
- ▶ cf. [Bovard et al., 2017]

# RESULTS: DISK EJECTA MASSES

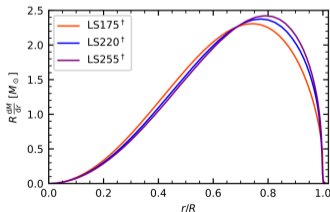


- ▶ of course collapse plays a dominant role
- ▶ dominance of the  $m = 2$  bar mode leads to efficient mass ejection - so called spiral wave wind {see [Nedora et al., 2019]}
- ▶ besides very general observations (e.g. softer EOSs lead to enhanced neutrino winds), no clear trends with EOS

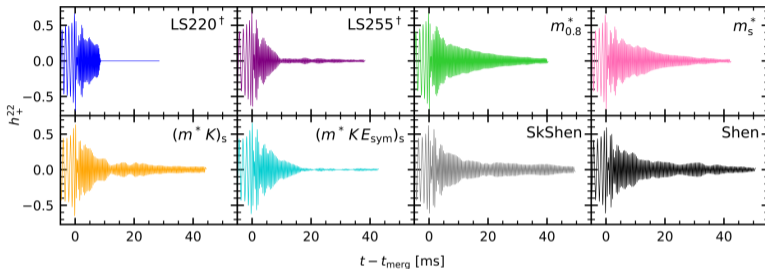
# RESULTS: DYNAMICAL EJECTA MASS VS. NUCLEAR PROPERTIES



- ▶ tidal and shock-heated components have different ejection mechanisms, so they correlate to different nuclear properties
- ▶ shocked ejecta correlate with  $P|_{n_0}$  (therefore  $m^*$  or  $E_{sym}$ )  $\rightarrow$  almost no dependence on  $K$
- ▶ for very stiff EOS, shock-heating dominates and there is no tidal ejecta
- ▶ otherwise, clear correlation between  $M_{ej}^{tid}$  and  $K$  (possibly due to radial matter pile-up)

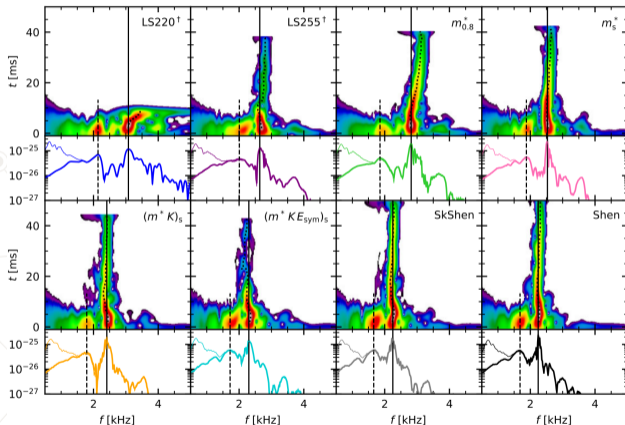


# RESULTS: GRAVITATIONAL WAVES



- ▶ + polarization of the (2, 2)-mode of the GW strain  $h$
- ▶ collapse induced shut-down for models LS175 and LS220
- ▶ “two phases” behaviour in amplitude reduction slope post-merger
  - 👉 linked to transition from  $l = 2$  to  $m = 1$  mode in the remnant

# RESULTS: GRAVITATIONAL WAVES



- ▶ model LS220 shows a “chirp-like” behaviour of the peak frequency before collapse, due to increased rotational velocity as the star contracts
- ▶ same effects is seen in other models (e.g.  $m_{0.8}^*$  and  $m_s^*$ ), but less dramatic
- ▶ decreasing  $m^*$ , increasing  $K$  and  $E_{\text{sym}}$  increases the central pressure and decreases the density, reducing  $f_2$
- ▶ the sequence of models  $m_s^*$ ,  $(m^*K)_s$ ,  $(m^*KE_{\text{sym}})_s$  and SkShen tends toward the Shen EOS spectrum. In particular SkShen and Shen are almost indistinguishable
- ▶ computed GW properties are therefore independent of the EOS microphysical model and mostly sensitive to the EOS around saturation density





- ▶  $K$  has a large impact at high density, i.e. in the core of the merger remnant. Increasing  $K$  decreases the compactness of the remnant and lowers the dominant GW frequency. It also quickly halts the remnant contraction, dampening its oscillations
- ▶  $m^*$  is important for both  $P_{\text{cold}}$  and  $P_{\text{th}}$ . Lowering it increases the  $P_{\text{cold}}$  at all densities, with a similar influence as  $K$ . But: the pressure density dependence is less steep, leading longer oscillations. Decreasing  $m^*$  increases  $\Gamma_{\text{th}}$ , while reducing shock-heating via reduced compactness. So  $P_{\text{th}}$  actually drops
- ▶ The mass of tidal ejecta correlates with  $K$  and the mass of shock-heated ejecta correlates with  $m^*$ . Mass ejection from the disk is more complex and hard to correlate with the EOS.
- ▶ The SkShen EOS is very similar to the original Shen (larger high density  $P$  and a slightly lower  $\Gamma_{\text{th}}$ . They lead to a similar evolution, and the GW spectra are remarkably similar
  - 👍 nuclear matter properties are a useful measure to quantify EOS effects in BNS mergers



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**Future work: run many more models, try to perform actual fits as function of  $m^*$ ,  $K$ , etc.**



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**Future work: run many more models, try to perform actual fits as function of  $m^*$ ,  $K$ , etc.**

Thank you

The background features a diagonal split between a teal upper-left section and a light beige lower-right section. The word 'REFERENCES' is centered in the white area between these two colors.

## REFERENCES

## REFERENCES I

- [Antoniadis et al., 2013] Antoniadis, J., Freire, P. C., Wex, N., Tauris, T. M., Lynch, R. S., et al. (2013).  
A Massive Pulsar in a Compact Relativistic Binary.  
*Science*, 340:6131.
- [Bovard et al., 2017] Bovard, L., Martin, D., Guercilena, F., Arcones, A., Rezzolla, L., and Korobkin, O. (2017).  
 $r$ -process nucleosynthesis from matter ejected in binary neutron star mergers.  
*Phys. Rev.*, D96(12):124005.
- [Demorest et al., 2010] Demorest, P., Pennucci, T., Ransom, S., Roberts, M., and Hessels, J. (2010).  
Shapiro Delay Measurement of A Two Solar Mass Neutron Star.  
*Nature*, 467:1081–1083.
- [Fonseca et al., 2021] Fonseca, E. et al. (2021).  
Refined Mass and Geometric Measurements of the High-mass PSR J0740+6620.  
*Astrophys. J. Lett.*, 915(1):L12.
- [Galeazzi et al., 2013] Galeazzi, F., Kastaun, W., Rezzolla, L., and Font, J. A. (2013).  
Implementation of a simplified approach to radiative transfer in general relativity.  
*Phys.Rev.*, D88:064009.

## REFERENCES II

[Mann, 2020] Mann, A. (2020).

The golden age of neutron-star physics has arrived.

*Nature*, 579.

[Nedora et al., 2019] Nedora, V., Bernuzzi, S., Radice, D., Perego, A., Endrizzi, A., and Ortiz, N. (2019).

Spiral-wave wind for the blue kilonova.

*Astrophys. J.*, 886(2):L30.

[Radice et al., 2016] Radice, D., Galeazzi, F., Lippuner, J., Roberts, L. F., Ott, C. D., and Rezzolla, L. (2016).

Dynamical Mass Ejection from Binary Neutron Star Mergers.

*Mon. Not. Roy. Astron. Soc.*, 460(3):3255–3271.

[Radice and Rezzolla, 2012] Radice, D. and Rezzolla, L. (2012).

THC: a new high-order finite-difference high-resolution shock-capturing code for special-relativistic hydrodynamics.

*Astron. Astrophys.*, 547:A26.

[Romani et al., 2022] Romani, R. W., Kandel, D., Filippenko, A. V., Brink, T. G., and Zheng, W. (2022).

PSR J0952–0607: The Fastest and Heaviest Known Galactic Neutron Star.

*Astrophys. J. Lett.*, 934(2):L17.

## REFERENCES III

[Zlochower et al., 2022] Zlochower, Y., Brandt, S. R., Diener, P., Gabella, W. E., Gracia-Linares, M., Haas, R., Kedia, A., Alcubierre, M., Alic, D., Allen, G., Ansorg, M., Babiuc-Hamilton, M., Baiotti, L., Bengert, W., Bentivegna, E., Bernuzzi, S., Bode, T., Bozzola, G., Brendal, B., Bruegmann, B., Campanelli, M., Cipolletta, F., Corvino, G., Cupp, S., Pietri, R. D., Dimmelman, H., Dooley, R., Dorband, N., Elley, M., Khamra, Y. E., Etienne, Z., Faber, J., Font, T., Friebe, J., Giacomazzo, B., Goodale, T., Gundlach, C., Hawke, I., Hawley, S., Hinder, I., Huerta, E. A., Husa, S., Iyer, S., Johnson, D., Joshi, A. V., Kastaun, W., Kellermann, T., Knapp, A., Koppitz, M., Laguna, P., Lanferman, G., Löffler, F., Masso, J., Menger, L., Merzky, A., Miller, J. M., Miller, M., Moesta, P., Montero, P., Mundim, B., Nelson, P., Nerozzi, A., Noble, S. C., Ott, C., Paruchuri, R., Pollney, D., Radice, D., Radke, T., Reisswig, C., Rezzolla, L., Rideout, D., Ripeanu, M., Sala, L., Schewtschenko, J. A., Schnetter, E., Schutz, B., Seidel, E., Seidel, E., Shalf, J., Sible, K., Sperhake, U., Stergioulas, N., Suen, W.-M., Szilagy, B., Takahashi, R., Thomas, M., Thornburg, J., Tobias, M., Tonita, A., Walker, P., Wan, M.-B., Wardell, B., Werneck, L., Witek, H., Zilhão, M., and Zink, B. (2022).

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