

## 5PN and 6PN completion of the 4PM EOB Hamiltonian

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### **PM-based EOB Hamiltonians**

Standard EOB Hamiltonian, built on PN results:

 $(\hat{H}_{\text{eff}}^{\text{PN}})^2 \equiv A_{\text{PN}}(u) \left[ 1 + p_{\varphi}^2 u^2 + \frac{1}{B_{\text{PN}}(u)} p_r^2 + Q_{\text{PN}}(p_r, u) \right]$ 

$$\hat{H}_{\rm EOB} \equiv \frac{1}{\nu} \sqrt{1 + 2\nu \left(\hat{H}_{\rm eff} - 1\right)}$$

 $\left(D_{\rm PN} = A_{\rm PN} \, B_{\rm PN}\right)$ 

 $\hat{H}_S^2 \equiv A_S \left( 1 + p_{\varphi}^2 u^2 + A_S p_r^2 \right)$ 

DJS gauge, no dependence in  $p_{\varphi}$ 

Alternatives devised to use instead PM information:

### **PS\* Hamiltonian**

$$A = A_{S} + A_{PM},$$
  

$$A_{PM} \equiv u^{2}a_{2PM} + u^{3}a_{3PM} + u^{4}a_{4PM} + \mathcal{O}(u^{5}),$$
  

$$B = B_{S} \equiv A_{S}^{-1}, \quad Q = 0$$
  

$$(\hat{H}_{eff}^{PS*})^{2} \equiv (A_{S} + A_{PM}) \left(1 + p_{\varphi}^{2}u^{2} + A_{S} p_{r}^{2}\right) = \hat{H}_{S}^{2} + A_{PM} \left(1 + p_{\varphi}^{2}u^{2} + A_{S} p_{r}^{2}\right) = \hat{H}_{S}^{2} \left(1 + A_{PM}/A_{S}\right)$$

Here  $(q_{nPM}, a_{nPM})$  are PM-informed coefficients

### **PM-based EOB Hamiltonians**

The coefficients  $(q_{nPM}, a_{nPM})$  are determined from the **PM expansion of the** scattering angle:

$$\chi_{\text{eff}}(\hat{H}_{\text{eff}}^{\text{PS}}) = \chi_{\text{PM}} \rightarrow q_{n\text{PM}} = q_{n\text{PM}}(\gamma) \qquad \qquad \gamma \equiv E_{\text{eff}}/\mu \qquad \qquad \text{[Damour 2017]: up to 2PM]} \\ \chi_{\text{eff}}(\hat{H}_{\text{eff}}^{\text{PS}*}) = \chi_{\text{PM}} \rightarrow a_{n\text{PM}} = a_{n\text{PM}}(\gamma) \qquad \qquad \gamma \equiv E_{\text{eff}}/\mu \qquad \qquad \text{[Khalil et al. 2022]: up to 4PM]}$$

For instance:

$$q_{2\mathrm{PM}}(\gamma) = \frac{3\left(5\gamma^2 - 1\right)(\Gamma - 1)}{2\Gamma}, \qquad a_{2\mathrm{PM}}(\gamma) = \frac{3\left(5\gamma^2 - 1\right)(\Gamma - 1)}{2\Gamma\gamma^2}, \qquad \Gamma \equiv \sqrt{1 + 2\nu\left(\gamma - 1\right)}$$

Here the dependence on the effective energy must be understood **perturbatively in the PM-sense**. At 4PM accuracy:

• 
$$q_{4\text{PM}}(\gamma), q_{3\text{PM}}(\gamma), a_{4\text{PM}}(\gamma), a_{3\text{PM}}(\gamma) \rightarrow \gamma = \hat{H}_{\text{eff}}^{\text{PS}} \Big|_{1\text{PM}} = \hat{H}_S$$

• 
$$q_{2PM}(\gamma) \rightarrow \gamma = \hat{H}_{eff}^{PS} \Big|_{2PM} = \sqrt{\hat{H}_S^2 + (1 - 2u) u^2 q_{2PM}(\hat{H}_S)}$$

• 
$$a_{2PM}(\gamma) \rightarrow \gamma = \hat{H}_{eff}^{PS*} \Big|_{2PM} = \hat{H}_S \sqrt{1 + \frac{u^2 a_{2PM}(\hat{H}_S)}{1 - 2u}}$$

### **Need for a PN completion**



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W.r.t. to the usual 6PN Hamiltonian, the 4PM Hamiltonian make us gain new analytical information but also lose some

How to gain with no pain?

We compute extra **PN complementary terms** for our 4PM Hamiltonian, going beyond the original 4PN result of [Khalil et al. 2022], up to 6PN

### **Computation outline**



### 6PN gauge transformation: DJS to E

General generating function:  $(p^2 = p_r^2 + p_{\varphi}^2 u^2)$ 

$$g_{\rm E}(r,p_r,p) = \frac{r\,p_r}{r^2} \left\{ \frac{f_1}{c^4} + \frac{1}{c^6} \left( g_1\,p^2 + g_2\,p_r^2 + \frac{g_3}{r} \right) + \frac{1}{c^8} \left( h_1\,p^4 + h_2\,p^2p_r^2 + h_3\,\frac{p^2}{r} + h_4\,p_r^4 + h_5\,\frac{p_r^2}{r} + \frac{h_6}{r^2} \right) \right\}$$
$$+ \frac{\log r}{r^2} h_{1,\log} \right\} + \frac{1}{c^{10}} \left[ p^6\,n_1 + \dots + \frac{n_{10}}{r^3} + \frac{\log r}{r^2} \left( n_{1,\log}\,p^2 + n_{2,\log}\,p_r^2 + \frac{n_{3,\log}}{r} \right) \right]$$
$$+ \frac{1}{c^{12}} \left[ p^8\,w_1 + \dots + \frac{w_{15}}{r^4} + \frac{\log r}{r^2} \left( w_{1,\log}\,p^4 + \dots + \frac{w_{6,\log}}{r^2} \right) \right]$$

Corresponding (parameter-dependent) gauge transformation:

$$r \to r - \frac{\partial g_E}{\partial p_r} + \frac{\partial g_E}{\partial p_r} \frac{\partial^2 g_E}{\partial r \, \partial p_r} - \left[ \frac{\partial g_E}{\partial p_r} \left( \frac{\partial^2 g_E}{\partial r \, \partial p_r} \right)^2 + \frac{1}{2} \left( \frac{\partial g_E}{\partial p_r} \right)^2 \frac{\partial^3 g_E}{\partial r^2 \, \partial p_r} \right]$$

$$p_r \to p_r + \frac{\partial g_E}{\partial r} - \frac{\partial g_E}{\partial p_r} \frac{\partial^2 g_E}{\partial r^2} + \left[ \frac{\partial g_E}{\partial p_r} \frac{\partial^2 g_E}{\partial r \partial p_r} \frac{\partial^2 g_E}{\partial r^2} + \frac{1}{2} \left( \frac{\partial g_E}{\partial p_r} \right)^2 \frac{\partial^3 g_E}{\partial r^3} \right]$$

### 6PN gauge transformation: DJS to E

In the energy gauge, the Hamiltonian must be writable as a function of **only**  $\varepsilon = \hat{H}_s^2 - 1$  and *u*, both of order  $1/c^2$ 

$$\varepsilon = p^2(1 - 2u) + 2\left(2u^2 - 3u + 1\right)p_r^2 - 2u \implies p_r^2 = \frac{\varepsilon - p^2(1 - 2u) + 2u}{2 - 6u + 4u^2}$$

If we use this to replace  $p_r^2$  in the energy gauge Hamiltonian, the dependence on  $p^2$  should also disappear

### **Procedure to fix the ansatz parameters in the generating function:**

- We use the ansatz-dependent transformation on the DJS  $\hat{H}_{\text{eff}}^{6\text{PN}}(p^2, p_r^2, u)$
- We replace  $p_r^2$  with  $\varepsilon$  using the relation above, re-expanding at 6PN
- We find the expressions of the parameters that cancel out each term that still depends on  $p^2\,$

This completely determines the transformation and allows us to obtain  $\hat{H}_{\rm eff}^{\rm 6PN}(\varepsilon, u)$ 

### EOB@Work23

## General PN complementary parts in $\hat{H}_{\text{eff}}^{\text{PS}*}$

$$\left(\hat{H}_{\text{eff}}^{\text{PS}*}\right)^2 = (1+\varepsilon) \left[1 + \frac{A_{4\text{PM}}(\varepsilon, u)}{1-2u}\right]$$

### **PN completion:**

$$A_{4\text{PM}} \rightarrow A_{4\text{PM}}^{6\text{PN}} \equiv u^2 a_{2\text{PM}} + u^3 a_{3\text{PM}} + u^4 a_{4\text{PM}} + \Delta_{4\text{PN}}^{\text{gen}} + \Delta_{5\text{PN}}^{\text{gen}} + \Delta_{6\text{PN}}^{\text{gen}}$$

where:

$$\Delta_{4\mathrm{PN}}^{\mathrm{gen}} = \sum_{n=2}^{5} \alpha_n \, u^n \varepsilon^{5-n} + \log u \sum_{n=0}^{1} \alpha_{n+4,\log u} \, u^{n+4} \epsilon^{1-n} + \alpha_{4,\log \varepsilon} \, u^4 \varepsilon \log \varepsilon \,,$$

$$\Delta_{5\mathrm{PN}}^{\mathrm{gen}} = \sum_{n=2}^{6} \beta_n \, u^n \varepsilon^{6-n} + \log u \sum_{n=0}^{2} \beta_{n+4,\log u} \, u^{n+4} \epsilon^{2-n} + \beta_{4,\log \varepsilon} \, u^4 \varepsilon^2 \log \varepsilon,$$

$$\Delta_{6\mathrm{PN}}^{\mathrm{gen}} = \sum_{n=2}^{7} \gamma_n \, u^n \varepsilon^{7-n} + \log u \sum_{n=0}^{3} \gamma_{n+4,\log u} \, u^{n+4} \epsilon^{3-n} + \gamma_{4,\log \varepsilon} \, u^4 \varepsilon^3 \log \varepsilon$$

### Hamiltonian matching

$$\hat{H}_{\text{eff}}^{6\text{PN}}(\varepsilon, u) = \hat{H}_{\text{eff}}^{\text{PS}*}(\varepsilon, u, \Delta_{4\text{PN}}^{\text{gen}}, \Delta_{5\text{PN}}^{\text{gen}}, \Delta_{6\text{PN}}^{\text{gen}})$$

After the PN expansion (simultaneous expansion in  $\varepsilon$  and u) the coefficients of each combination of ( $\varepsilon$ , u) must coincide



Agrees with the result of [Khalil et al. 2022]

## Resulting $\Delta_{5PN}$



## Resulting $\Delta_{5PN}$



## Resulting $\Delta_{6PN}$



## Resulting $\Delta_{6PN}$



## Is there a better PM structure for $\hat{H}_{\rm eff}$ ?

**Issue with the PS Hamiltonian:** 

$$\left(\hat{H}_{\text{eff}}^{\text{PS}}\right)^{2} = \underbrace{(1-2u)}_{\square} \left[1 + p_{\varphi}^{2}u^{2} + (1-2u)p_{r}^{2} + Q_{\text{PM}}\right]$$
  
The effective horizon is locked to its Schwarzschild value,  $u_{\text{EH}} = 1/2$ 

#### **Issue with the PS\* Hamiltonian:**

$$\left(\hat{H}_{\text{eff}}^{\text{PS}*}\right)^2 \equiv \left(1 - 2u + A_{\text{PM}}\right) \left(1 + p_{\varphi}^2 u^2 + (1 - 2u) p_r^2\right)$$

OK effective horizon but in the circular limit we have singularities for  $u \to 1/2$  (i.e.  $\hat{H}_S \to 0$ ) in every PM coefficient  $a_{nPM}$ 

For instance:

$$\hat{H}_{S}^{2} \equiv (1 - 2u) \left[ 1 + p_{\varphi}^{2} u^{2} + (1 - 2u) p_{r}^{2} \right]$$

$$a_{2\text{PM}}^{\text{circ}}(\hat{H}_{S}) = \frac{3\left(5\hat{H}_{S}^{2}-1\right)\left[\sqrt{2\nu(\hat{H}_{S}-1)+1}-1\right]}{2\hat{H}_{S}^{2}\sqrt{2\nu(\hat{H}_{S}-1)+1}} \xrightarrow{\hat{H}_{S}\rightarrow 0} \infty$$

# **BONUS SLIDES**

## Explicit $a_{3PM}$ and $a_{4PM}$

a3PM /. Hs  $\rightarrow\gamma$  // Simplify

$$8 \gamma \left(-25 + 11 \gamma^{2} + 14 \gamma^{4}\right) \nu + 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \left(1 + 2 \left(-1 + \gamma\right) \nu\right) + 48 \sqrt{-1 + \gamma^{2}} \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \left(1 + 2 \left(-1 + \gamma\right) \nu\right) + 48 \sqrt{-1 + \gamma^{2}} \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \left(1 + 2 \left(-1 + \gamma\right) \nu\right) + 48 \sqrt{-1 + \gamma^{2}} \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \left(1 + 2 \left(-1 + \gamma\right) \nu\right) + 48 \sqrt{-1 + \gamma^{2}} \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \left(1 + 2 \left(-1 + \gamma\right) \nu\right) + 48 \sqrt{-1 + \gamma^{2}} \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} - 9 \left(5 - 31 \gamma^{2} + 30 \gamma^{4}\right) \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} + 9 \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} + 9 \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right] \sqrt{1 + 2 \left(-1 + \gamma\right) \nu} + 9 \left(-3 - 12 \gamma^{2} + 4 \gamma^{4}\right) \nu ArcSinh\left[\frac{\sqrt{-1 + \gamma}}{\sqrt{2}}\right]$$

 $6 \gamma^{2} \left(-1+\gamma^{2}\right) \left(1+2 \left(-1+\gamma\right) \nu\right)$ 

a4PM /. Hs  $\rightarrow \gamma$  // Simplify

$$\frac{7}{(169 + 380 \gamma^2)} \times \text{Elliptic} \left[\frac{14\gamma}{1\gamma}\right]^2}{8 (-1+\gamma) \gamma^2 (1+2 (-1+\gamma) \gamma)^{3/2}} - \frac{(1183 + 2929 \gamma + 2660 \gamma^2 + 1200 \gamma^1) \times \text{Elliptic} \left[\frac{14\gamma}{1\gamma}\right]^2}{4 \gamma^2 (-1+\gamma^2) (1+2 (-1+\gamma) \gamma)^{3/2}} + \frac{2 (-852 - 283 \gamma^2 - 140 \gamma^4 + 75 \gamma^5) \times \text{Log}(\gamma)}{3 \gamma (-1+\gamma^2) (1+2 (-1+\gamma) \gamma)^{3/2}} + \frac{1}{144 \gamma^6 (-1+\gamma^2)^2 (1+2 (-1+\gamma) \gamma)^{3/2}} - \frac{(183 + 2929 \gamma + 2660 \gamma^2 + 1200 \gamma^1) \times \text{Log}(\gamma)}{3 \gamma (-1+\gamma^2) (1+2 (-1+\gamma) \gamma)^{3/2}} + \frac{1}{144 \gamma^6 (-1+\gamma^2)^2 (1+2 (-1+\gamma) \gamma)^{3/2}} - \frac{(163 + 203 \gamma^3 + 1260 \gamma^3) \times \text{Log}(\gamma)}{3 \gamma (-1+\gamma^2) (1+2 (-1+\gamma) \gamma)^{3/2}} + \frac{1}{144 \gamma^6 (-1+\gamma^2)^2 (1+2 (-1+\gamma) \gamma)^{3/2}} - \frac{(163 + 203 \gamma^3 + 1260 \gamma^3) \times 1260 \gamma^4 + 7575 \gamma^5 \gamma - 2883 \gamma^2 \gamma + 83864 \gamma^4 + 1842 \gamma^5 \gamma + 9576 \gamma^6 \gamma - 2520 \gamma^2 \gamma^5 \gamma - 104 \gamma^4 + 729 \gamma^5 \gamma + 1205 \gamma^5 \gamma + 1205 \gamma^4 \gamma - 5575 \gamma^2 \gamma - 2883 \gamma^2 \gamma + 83864 \gamma^4 + 1842 \gamma^3 \gamma + 1520 \gamma^2 \gamma^5 \gamma + 104 \gamma^2 \gamma^5 \gamma + 200 \gamma^2 \gamma^5 \gamma + 200 \gamma^2 \gamma^5 \gamma + 120 \gamma^2 \gamma + 12(-1+\gamma) \gamma \sqrt{142 (-1+\gamma) \gamma 142 (-1+\gamma) \gamma$$