

Viability Gravity Mediation

Takemichi Okui

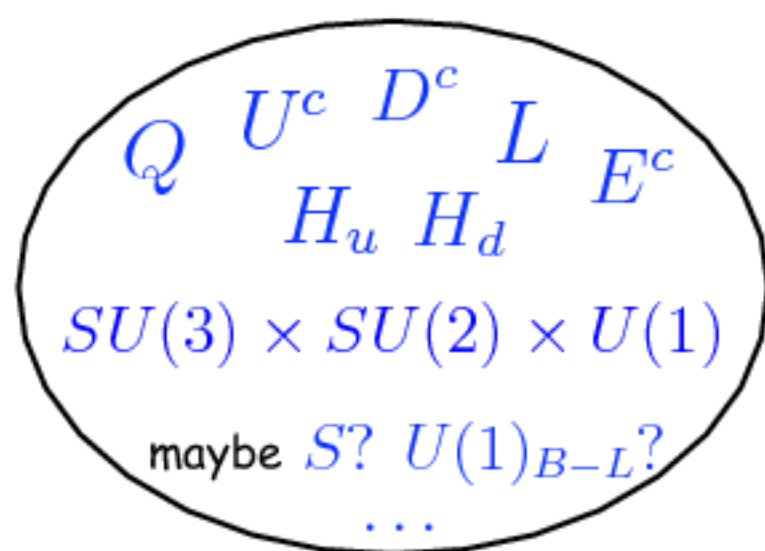
(Florida State Univ.)

PRD 82, 115010 (2010)

With Graham Kribs (U. Oregon, Fermilab) and Tuhin Roy (U. Washington)

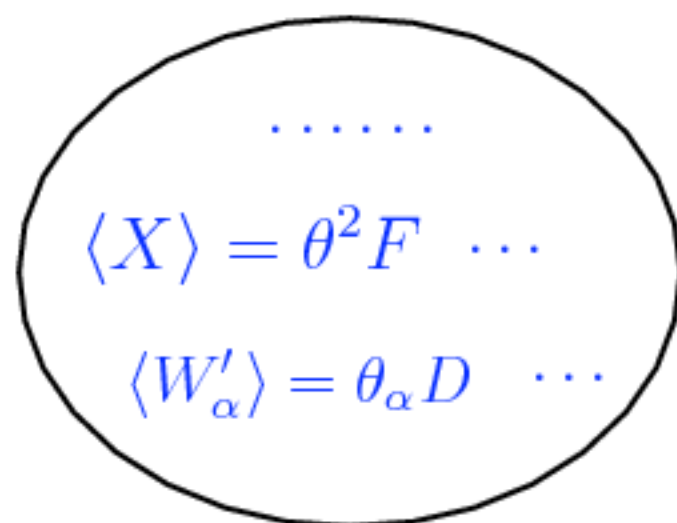
The biggest question in SUSY

"Visible" sector



$$\mathcal{L} = \mathcal{L}_{\text{vis.}}(Q, U^c, \dots)$$

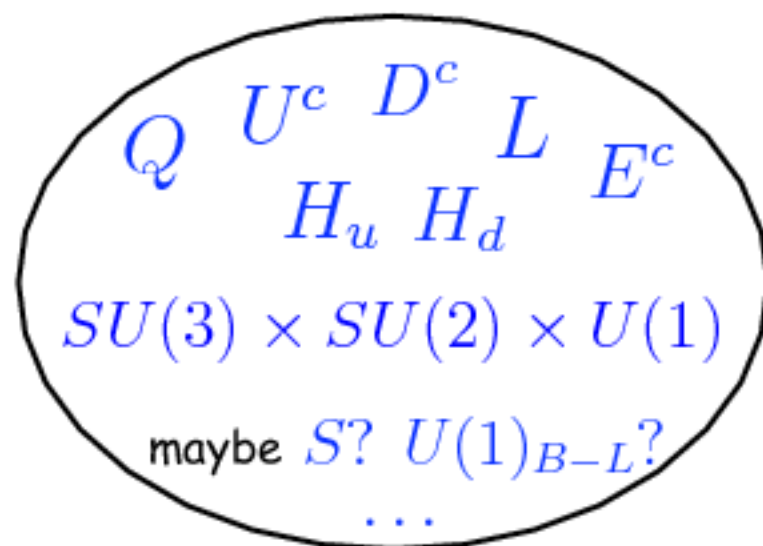
~~SUSY~~ sector



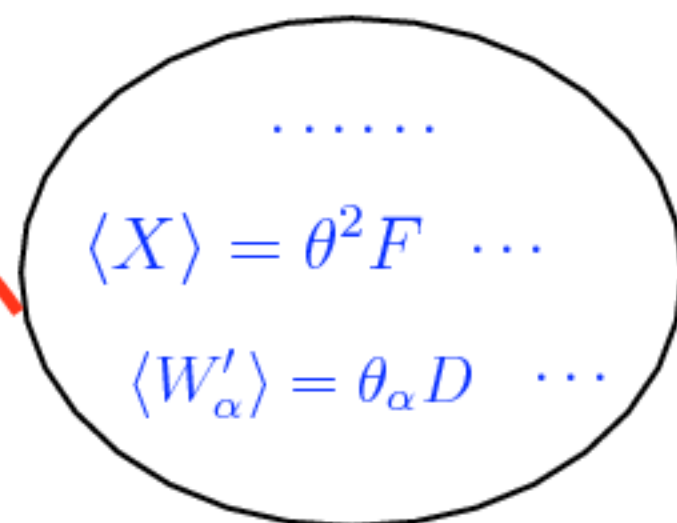
$$+ \mathcal{L}_{\text{S.B.}}(X, W', \dots)$$

The biggest question in SUSY

"Visible" sector



~~SUSY~~ sector



$$\mathcal{L} = \mathcal{L}_{\text{vis.}}(Q, U^c, \dots) + \mathcal{L}_{\text{med.}}(Q, U^c, \dots, X, W') + \mathcal{L}_{\text{S.B.}}(X, W', \dots)$$

How is SUSY breaking mediated to "us"?

How is SUSY breaking mediated to "us"?

At Planck scale, we **always** have

$$\mathcal{L}_{\text{med.}} \sim \int d^4\theta \frac{X^\dagger X Q^\dagger Q}{M_{\text{Pl}}^2} + \dots \xrightarrow{\langle X \rangle = \theta^2 F} m_{\tilde{q}}^2 \sim \frac{F^2}{M_{\text{Pl}}^2}$$

How is SUSY breaking mediated to "us"?

At Planck scale, we **always** have

$$\mathcal{L}_{\text{med.}} \sim \int d^4\theta \frac{X^\dagger X Q^\dagger Q}{M_{\text{Pl}}^2} + \dots \xrightarrow{\langle X \rangle = \theta^2 F} m_{\tilde{q}}^2 \sim \frac{F^2}{M_{\text{Pl}}^2}$$

The simplest scenario: **this is it!** "Gravity Mediation"

→ All mass scales right, no μ problem, neutralino DM, already unified...

How is SUSY breaking mediated to "us"?

At Planck scale, we **always** have

$$\mathcal{L}_{\text{med.}} \sim c_{ij} \int d^4\theta \frac{X^\dagger X Q_i^\dagger Q_j}{M_{\text{Pl}}^2} + \dots \xrightarrow{\langle X \rangle = \theta^2 F} (m_{\tilde{q}}^2)_{ij} \sim c_{ij} \frac{F^2}{M_{\text{Pl}}^2}$$

The simplest scenario: **this is it!** "Gravity Mediation"

→ All mass scales right, no μ problem, neutralino DM, already unified...

Alas, **too much flavor!** Generically, $c_{ij} \not\propto \delta_{ij} \dots$

How is SUSY breaking mediated to "us"?

At Planck scale, we **always** have

$$\mathcal{L}_{\text{med.}} \sim c_{ij} \int d^4\theta \frac{X^\dagger X Q_i^\dagger Q_j}{M_{\text{Pl}}^2} + \dots \xrightarrow{\langle X \rangle = \theta^2 F} (m_{\tilde{q}}^2)_{ij} \sim c_{ij} \frac{F^2}{M_{\text{Pl}}^2}$$

The simplest scenario: **this is it!** "Gravity Mediation"

→ All mass scales right, no μ problem, neutralino DM, already unified...

Alas, **too much flavor!** Generically, $c_{ij} \not\propto \delta_{ij} \dots$

Traditional Ansatz: "flavor universality" $c_{ij} \propto \delta_{ij}$

Self-consistent: $c_{ij} \propto \delta_{ij}$ in UV → $c_{ij} \propto \delta_{ij}$ in IR

How is SUSY breaking mediated to "us"?

At Planck scale, we **always** have

$$\mathcal{L}_{\text{med.}} \sim c_{ij} \int d^4\theta \frac{X^\dagger X Q_i^\dagger Q_j}{M_{\text{Pl}}^2} + \dots \xrightarrow{\langle X \rangle = \theta^2 F} (m_{\tilde{q}}^2)_{ij} \sim c_{ij} \frac{F^2}{M_{\text{Pl}}^2}$$

The simplest scenario: **this is it!** "Gravity Mediation"

————> All mass scales right, no μ problem, neutralino DM, already unified...

Alas, **too much flavor!** Generically, $c_{ij} \not\propto \delta_{ij} \dots$

Traditional Ansatz: "flavor universality" $c_{ij} \propto \delta_{ij}$

Self-consistent: $c_{ij} \propto \delta_{ij}$ in UV $\longrightarrow c_{ij} \propto \delta_{ij}$ in IR

But flavor symmetry **IS** violated by Yukawas!

Then, WHY $c_{ij} \propto \delta_{ij}$? Any mechanism? "SUSY flavor problem"

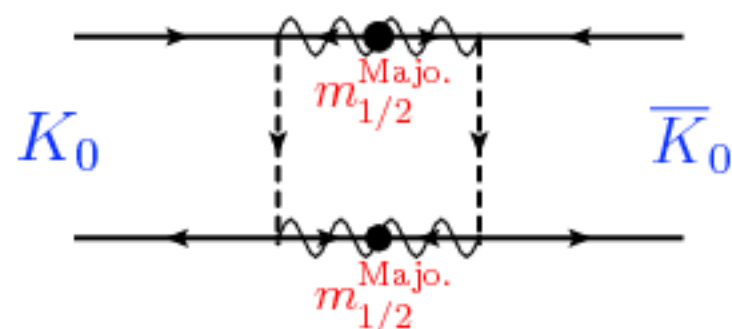
Solutions: gauge med., gaugino med., anomaly med., etc.

Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --

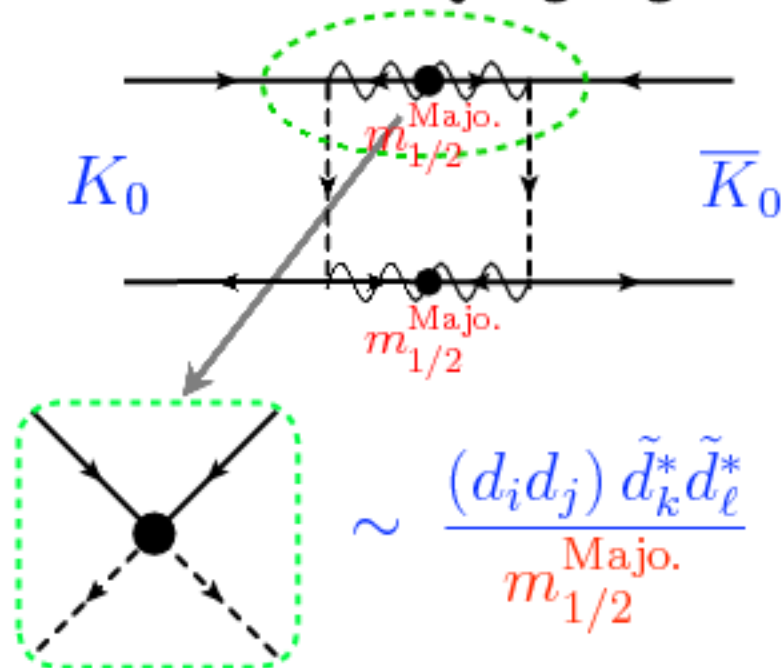


Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --

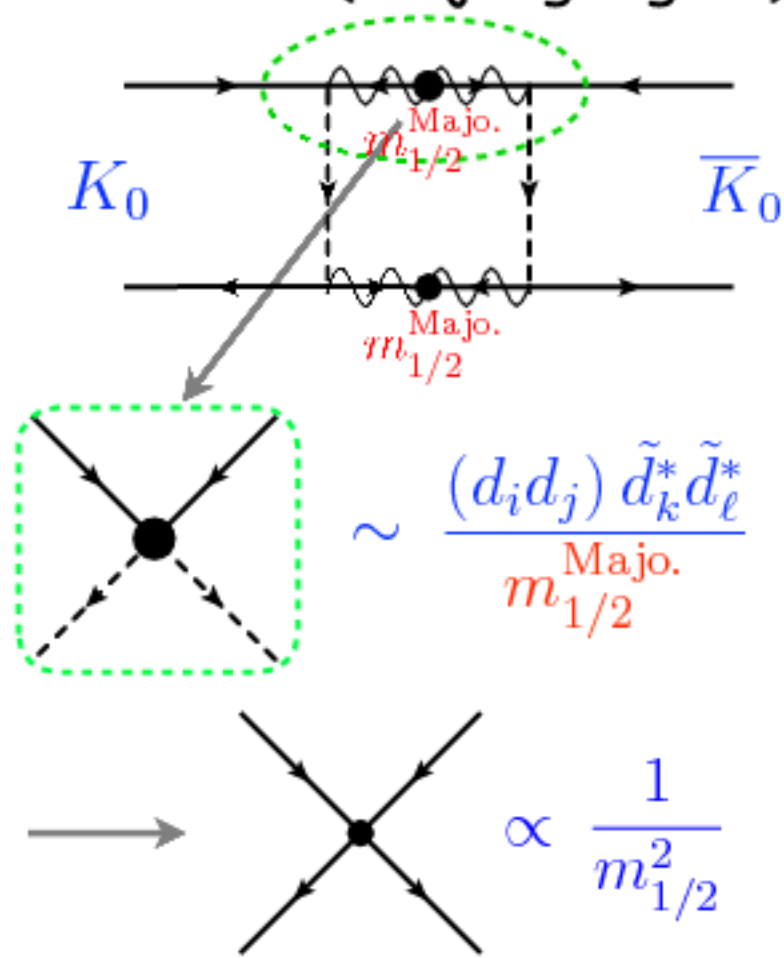


Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --



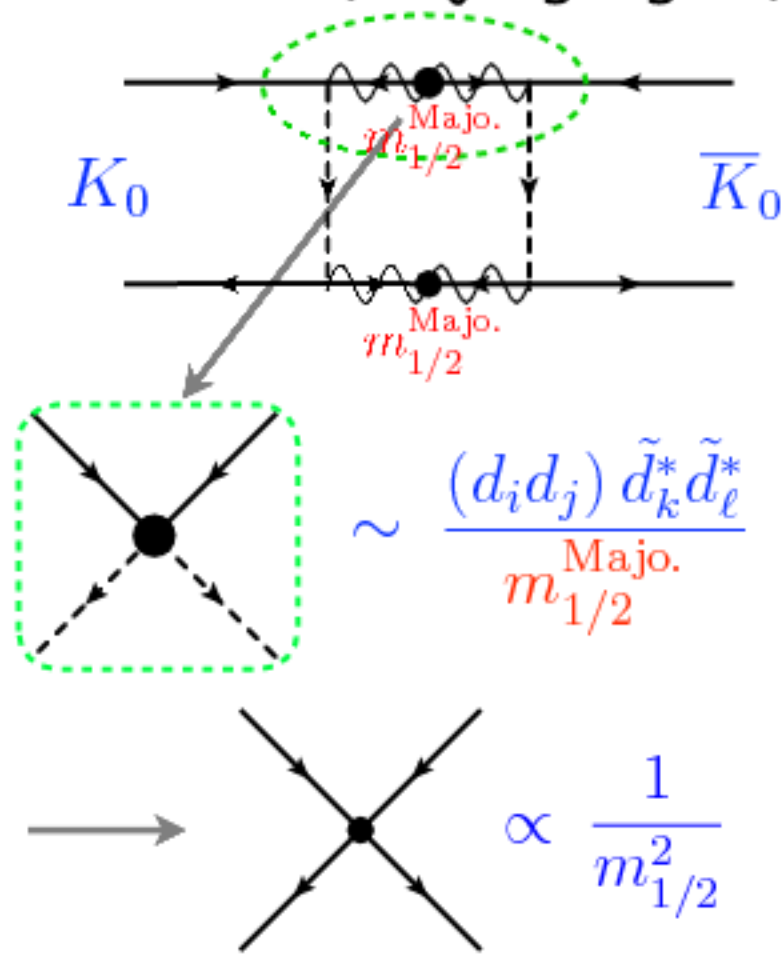
Suppressed but not enough...

Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

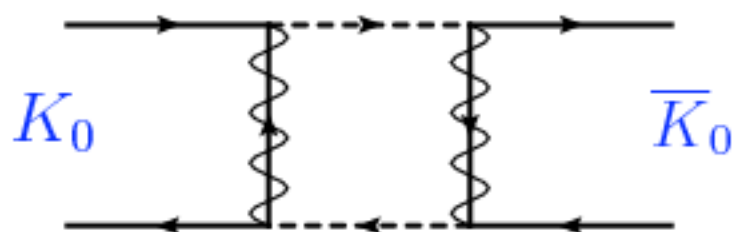
--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --



Suppressed but not enough...

-- Dirac gaugino --

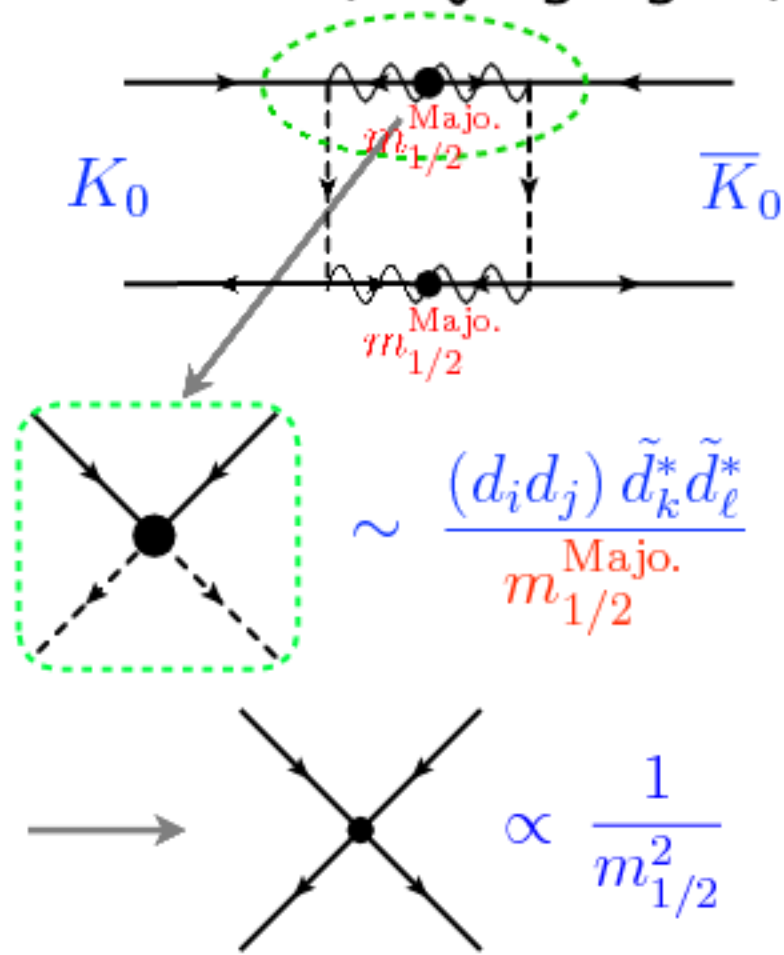


Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

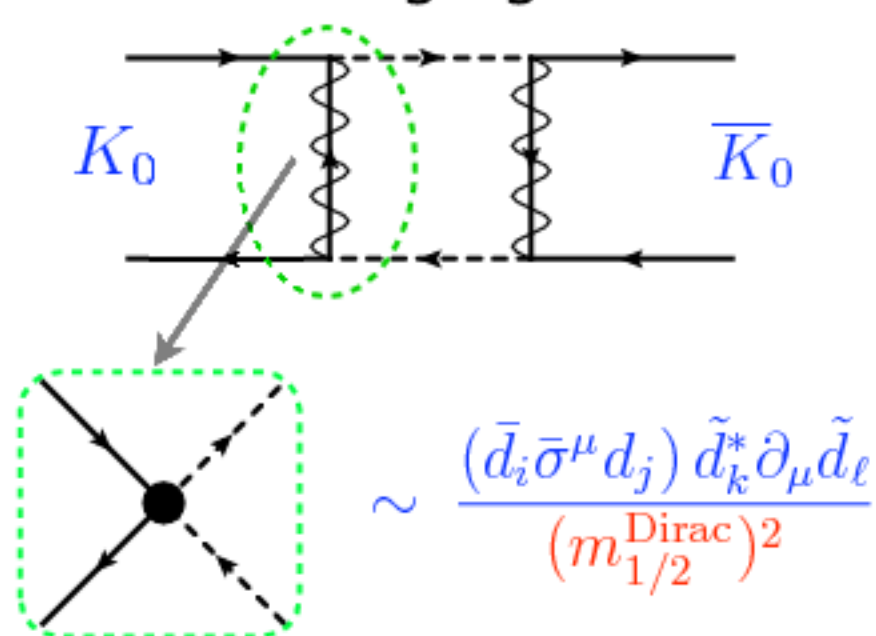
--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --



Suppressed but not enough...

-- Dirac gaugino --

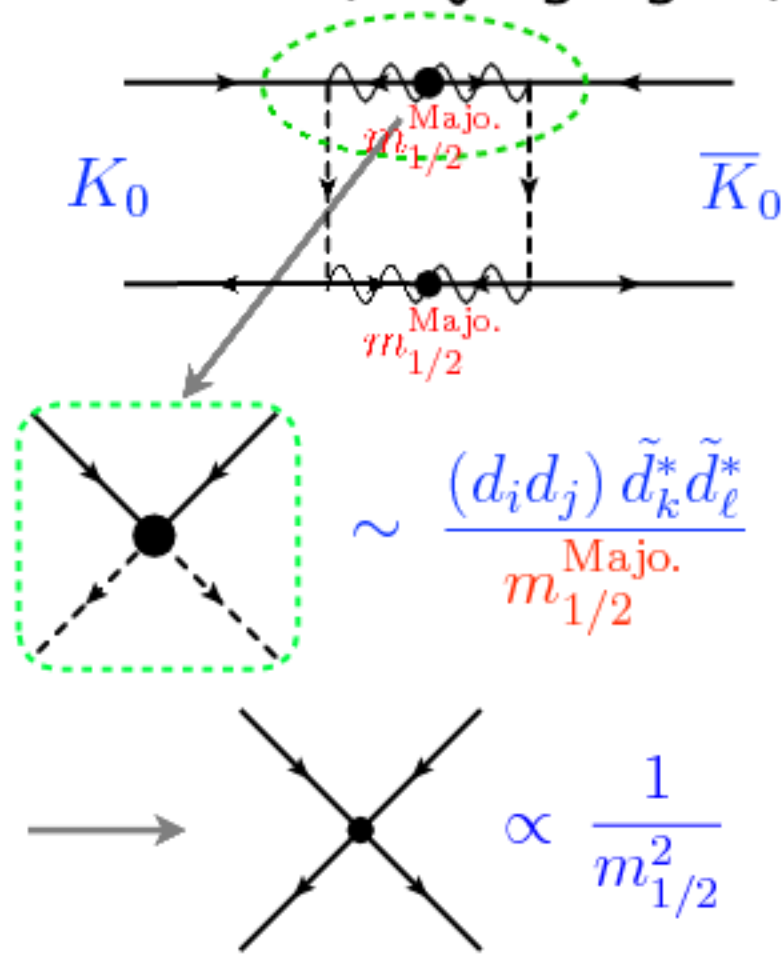


Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

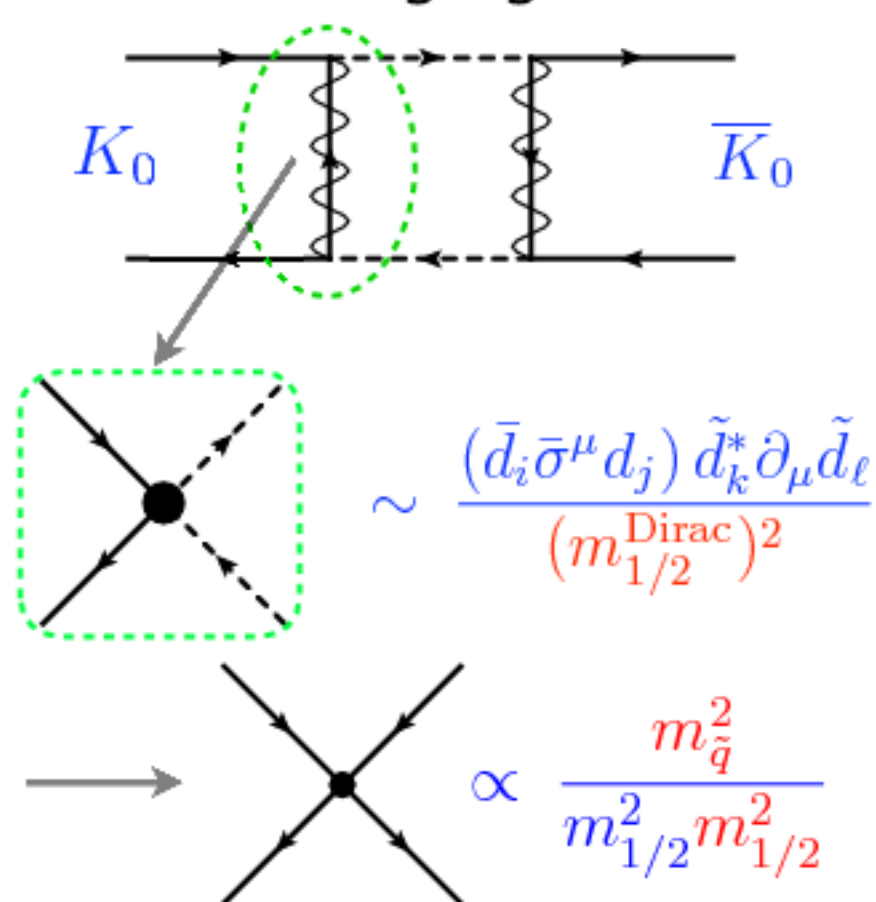
--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --



Suppressed but not enough...

-- Dirac gaugino --

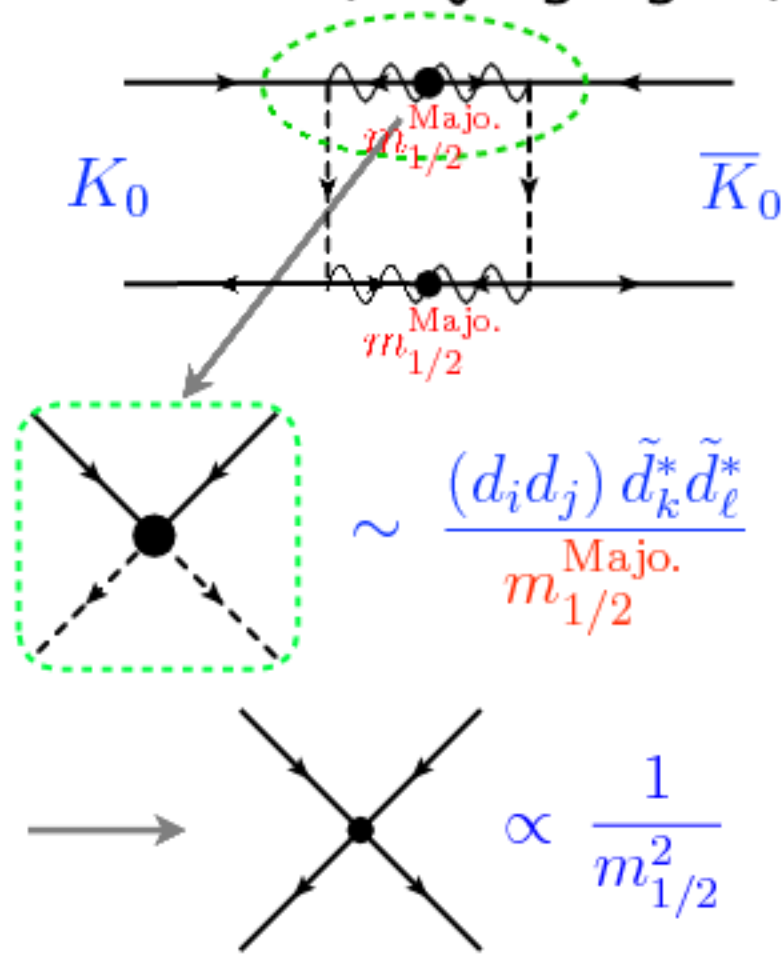


Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

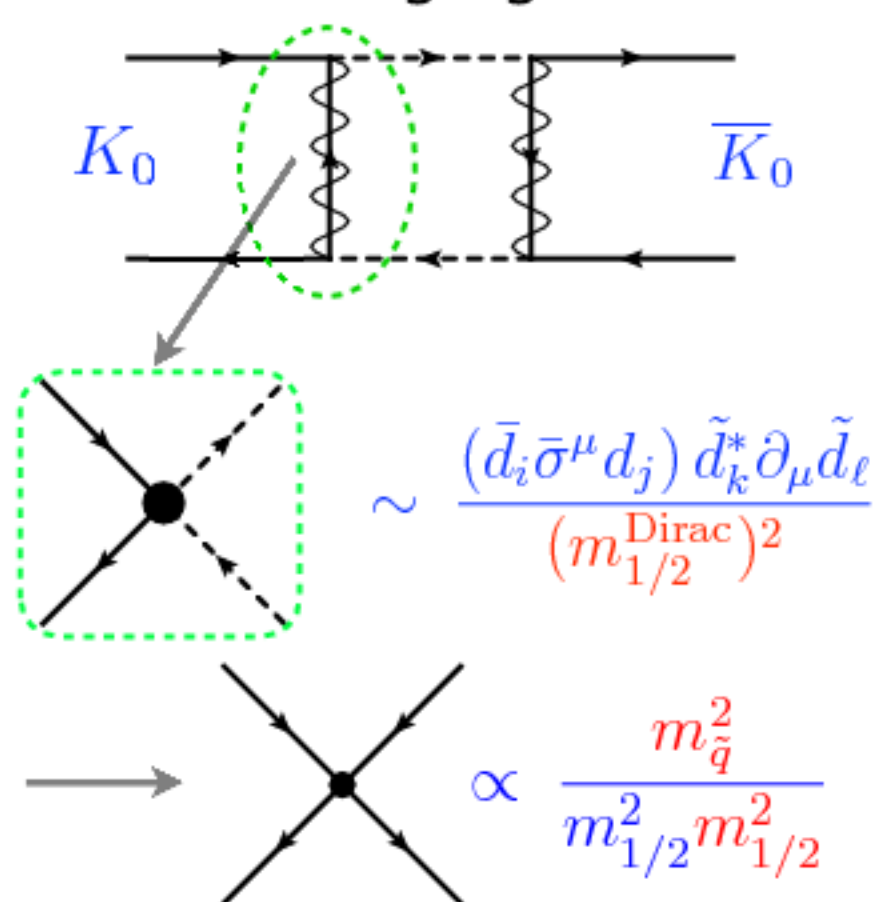
--- Excessive K_0 - $\bar{K}_0$ mixing ---

-- MSSM (Majo. gaugino) --



Suppressed but not enough...

-- Dirac gaugino --



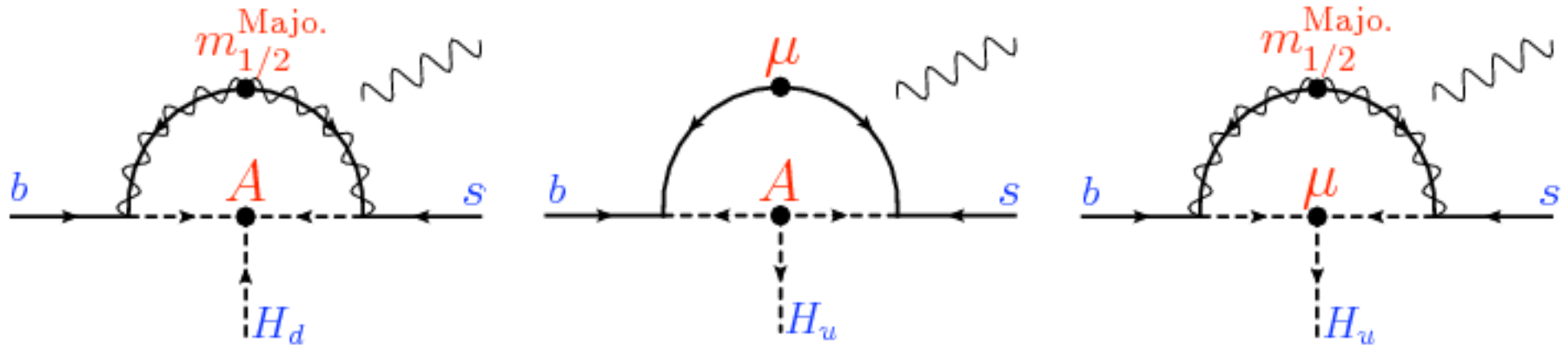
Enough suppression if

$$m_{1/2} \sim 10 m_{\tilde{q}}, \quad c_{12} \sim c_{11}/10 \sim c_{22}/10!$$

Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

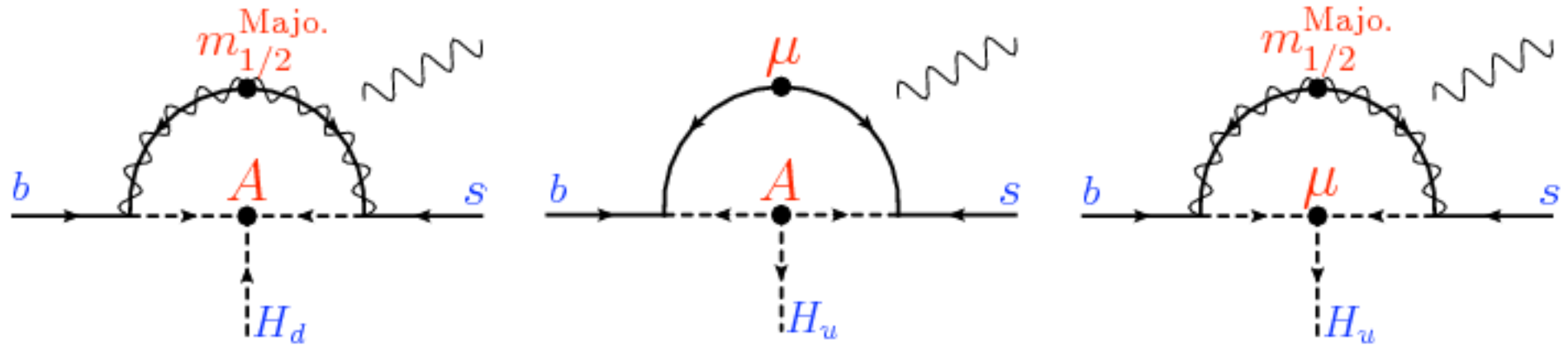
--- Excessive $\mu \rightarrow e\gamma$, $b \rightarrow s\gamma$ etc. ---



Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive $\mu \rightarrow e\gamma$, $b \rightarrow s\gamma$ etc. ---

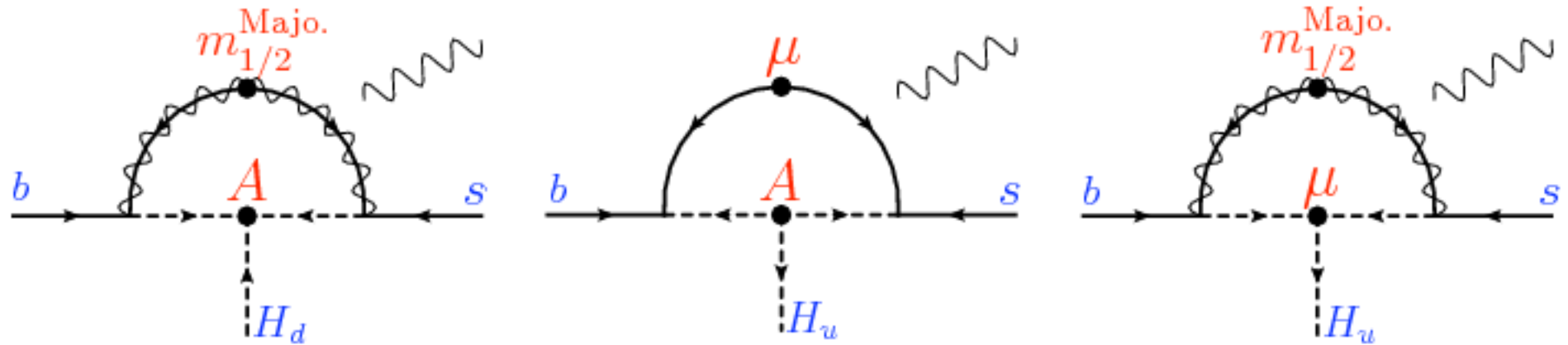


Will be **doubly forbidden** if $A = \mu = m_{1/2}^{\text{Majo.}} = 0!$

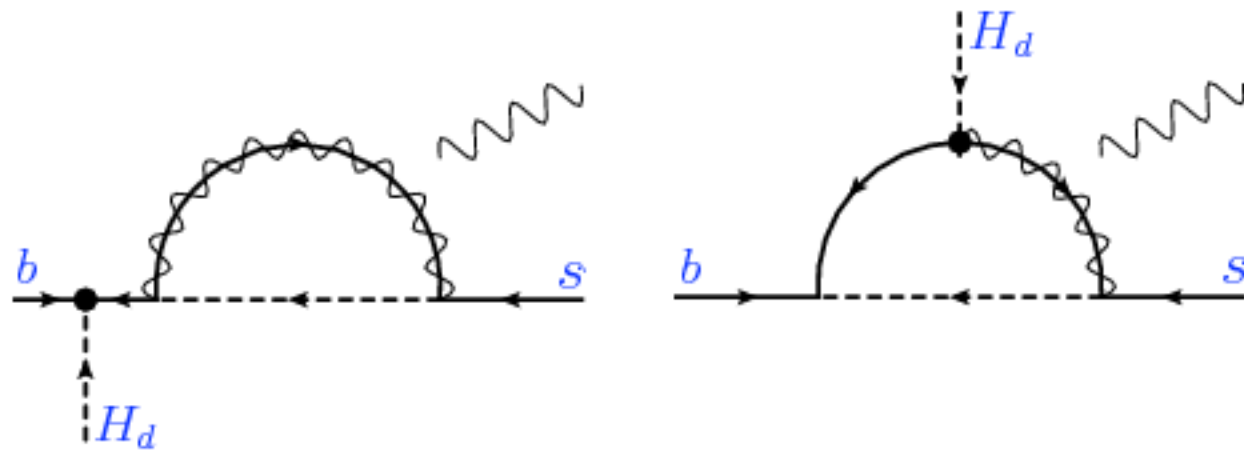
Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive $\mu \rightarrow e\gamma$, $b \rightarrow s\gamma$ etc. ---



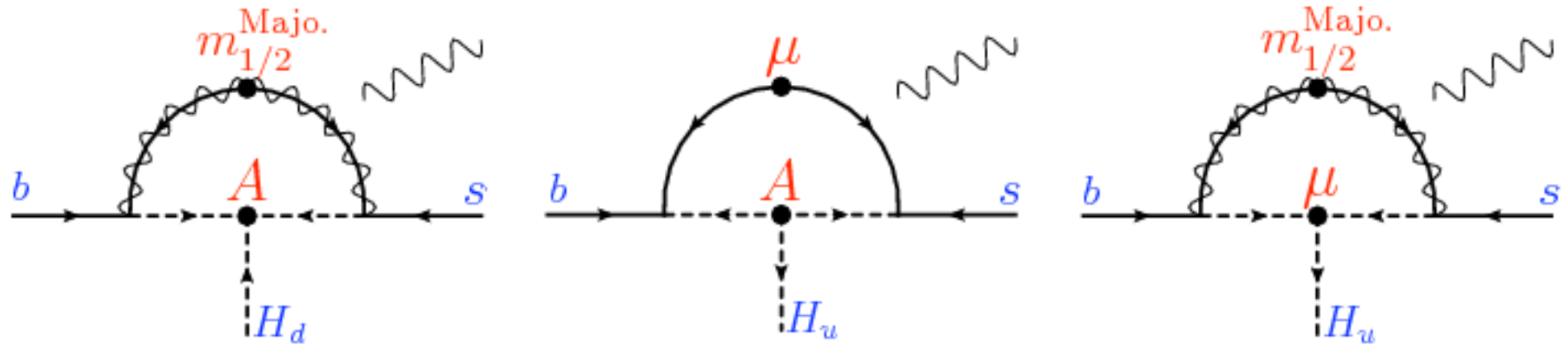
Will be **doubly forbidden** if $A = \mu = m_{1/2}^{\text{Majo.}} = 0!$



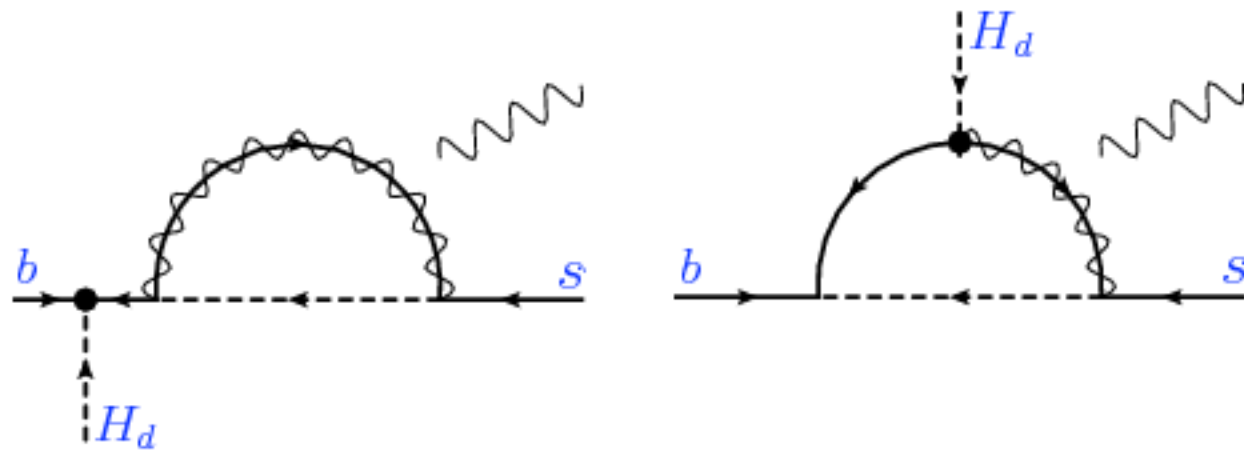
Is there a way to be flavor safe while $c_{ij} \not\propto \delta_{ij}$?

A closer look at SUSY flavor violations:

--- Excessive $\mu \rightarrow e\gamma$, $b \rightarrow s\gamma$ etc. ---



Will be **doubly forbidden** if $A = \mu = m_{1/2}^{\text{Majo.}} = 0$!



Will survive but **small for** $m_{1/2} \sim 10m_{\tilde{q}}$, $c_{12} \sim c_{11}/10 \sim c_{22}/10$!

How do we get $A = \mu = m_{1/2}^{\text{Majo.}} = 0$?

Kribs, Poppitz, Weiner '07: "A $U(1)_R$ can do it!"

$\lambda\lambda^c$	$H_u Q U^c$	$h_u h_d$
$\lambda\lambda$	$h_u \tilde{q} \tilde{u}^c$	$H_u H_d$
$m_{1/2}^{\text{Majo.}}$	A	μ

How do we get $A = \mu = m_{1/2}^{\text{Majo.}} = 0$?

Kribs, Poppitz, Weiner '07: "A $U(1)_R$ can do it!"

$$\left\{ \begin{array}{l} R[\lambda] = 1 = -R[\lambda^c] \\ R[Q] = R[U^c, D^c] = 1 \\ R[H_{u,d}] = 0 \end{array} \right. \longrightarrow \begin{array}{ccc} \lambda\lambda^c & H_u Q U^c & h_u h_d \\ \cancel{\lambda\lambda} & \cancel{h_u \tilde{q} u^c} & \cancel{H_u H_d} \\ \cancel{m_{1/2}^{\text{Majo.}}} & \cancel{A} & \cancel{\mu} \end{array}$$

How do we get $A = \mu = m_{1/2}^{\text{Majo.}} = 0$?

Kribs, Poppitz, Weiner '07: "A $U(1)_R$ can do it!"

$$\begin{cases} R[\lambda] = 1 = -R[\lambda^c] \\ R[Q] = R[U^c, D^c] = 1 \\ R[H_{u,d}] = 0 \end{cases}$$

$\lambda\lambda^c$ $H_u Q U^c$ $h_u h_d$
 ~~$\lambda\lambda$~~ ~~$h_u \tilde{q} u^c$~~ ~~$H_u H_d$~~
 ~~$m_{1/2}^{\text{Majo.}}$~~ ~~A~~ ~~μ~~

Where does it come from?

What about higgsino masses?

How do we get $A = \mu = m_{1/2}^{\text{Majo.}} = 0$?

Kribs, Poppitz, Weiner '07: "A $U(1)_R$ can do it!"

$$\begin{cases} R[\lambda] = 1 = -R[\lambda^c] \\ R[Q] = R[U^c, D^c] = 1 \\ R[H_{u,d}] = 0 \end{cases}$$

$\lambda\lambda^c$ (circled in green)
 $\lambda\lambda$ (crossed out)
 $m_{1/2}^{\text{Majo.}}$ (crossed out)
 $H_u Q U^c$
 $h_u \tilde{q} u^c$ (crossed out)
 A (crossed out)
 $h_u h_d$
 $H_u H_d$ (circled in green)
 μ (crossed out)

Where does it come from?

Introduce: $\begin{cases} \Sigma_3 \supset \tilde{g}^c \sim (\mathbf{8}, \mathbf{1})_0, R[\Sigma_3] = 0 \\ \Sigma_2 \supset \tilde{W}^c \sim (\mathbf{1}, \mathbf{3})_0, R[\Sigma_2] = 0 \\ \Sigma_1 \supset \tilde{B}^c \sim (\mathbf{1}, \mathbf{1})_0, R[\Sigma_1] = 0 \end{cases}$

$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_3 \Sigma_3$ etc.
 with $\langle W'_\alpha \rangle = \theta_\alpha D$

What about higgsino masses?

How do we get $A = \mu = m_{1/2}^{\text{Majo.}} = 0$?

Kribs, Poppitz, Weiner '07: "A $U(1)_R$ can do it!"

$$\begin{cases} R[\lambda] = 1 = -R[\lambda^c] \\ R[Q] = R[U^c, D^c] = 1 \\ R[H_{u,d}] = 0 \end{cases}$$

$$\lambda\lambda^c \quad H_u Q U^c \quad h_u h_d$$

$$\cancel{\lambda\lambda} \quad \cancel{h_u \tilde{q} u^c} \quad \cancel{H_u H_d}$$

$$\cancel{m_{1/2}^{\text{Majo.}}} \quad \cancel{A} \quad \cancel{\mu}$$

Where does it come from?

Introduce: $\begin{cases} \Sigma_3 \supset \tilde{g}^c \sim (\mathbf{8}, \mathbf{1})_0, R[\Sigma_3] = 0 \\ \Sigma_2 \supset \tilde{W}^c \sim (\mathbf{1}, \mathbf{3})_0, R[\Sigma_2] = 0 \\ \Sigma_1 \supset \tilde{B}^c \sim (\mathbf{1}, \mathbf{1})_0, R[\Sigma_1] = 0 \end{cases} \longrightarrow \int d^2\theta \frac{W'}{M_{\text{Pl}}} W_3 \Sigma_3 \text{ etc.}$

with $\langle W'_\alpha \rangle = \theta_\alpha D$

What about higgsino masses?

Introduce: $\begin{cases} R_u \sim (\mathbf{1}, \mathbf{2})_{-1/2}, R[R_u] = 2 \\ R_d \sim (\mathbf{1}, \mathbf{2})_{+1/2}, R[R_d] = 2 \end{cases} \longrightarrow H_u R_u + H_d R_d$

But $U(1)_R$ **cannot** be a fundamental symmetry!

- * No global symmetry in quantum gravity.
- * Gravitino mass requires ~~$U(1)_R$~~ !
- * (cosmo. const.) = ~~(SUSY)~~ + (const. in superpotential)
$$0 = (10^{11} \text{ GeV})^4 + (-(10^{14} \text{ GeV})^6 / M_{\text{Pl}}^2)$$

But $U(1)_R$ **cannot** be a fundamental symmetry!

- * No global symmetry in quantum gravity.

- * Gravitino mass requires ~~$U(1)_R$~~ !

- * (cosmo. const.) = ~~(SUSY)~~ + (const. in superpotential)
$$0 = (10^{11} \text{ GeV})^4 + \left(-(10^{14} \text{ GeV})^6 / M_{\text{Pl}}^2 \right)$$

A **big explicit** $U(1)_R$ breaking!

(Also gives "R-axion" mass)

But $U(1)_R$ **cannot** be a fundamental symmetry!

- * No global symmetry in quantum gravity.

- * Gravitino mass requires ~~$U(1)_R$~~ !

- * (cosmo. const.) = ~~(SUSY)~~ + (const. in superpotential)
 $0 = (10^{11} \text{ GeV})^4 + \boxed{-(10^{14} \text{ GeV})^6 / M_{\text{Pl}}^2}$

A **big explicit** $U(1)_R$ breaking!

(Also gives "R-axion" mass)

[Traditional flavor prob.]

Sfermions: $U(3)^5$ (i.e. $c_{ij} \propto \delta_{ij}$)

Fermions: ~~$U(3)^5$~~

How?

[New flavor prob.]

"Normal" stuff: $U(1)_R$

Gravity stuff: ~~$U(1)_R$~~

How?

But $U(1)_R$ **cannot** be a fundamental symmetry!

* No global symmetry in quantum gravity.

* Gravitino mass requires ~~$U(1)_R$~~ !

* (cosmo. const.) = ~~(SUSY)~~ + (const. in superpotential)

$$0 = (10^{11} \text{ GeV})^4 + \boxed{-(10^{14} \text{ GeV})^6 / M_{\text{Pl}}^2}$$

A **big explicit** $U(1)_R$ breaking!

(Also gives "R-axion" mass)

[Traditional flavor prob.]

Sfermions: $U(3)^5$ (i.e. $c_{ij} \propto \delta_{ij}$)

Fermions: ~~$U(3)^5$~~

How?

Gauge Med.: $U(3)^5$ **"accidental"**

[New flavor prob.]

"Normal" stuff: $U(1)_R$

Gravity stuff: ~~$U(1)_R$~~

How?

$U(1)_R$ **must be accidental!**

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow h_u h_d, r_u r_d, (\sigma_{1,2,3})^2$$

(B-type masses)

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow h_u h_d, r_u r_d, (\sigma_{1,2,3})^2$$

(B-type masses)

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

$$\longrightarrow \text{Dirac gaugino masses}$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow h_u h_d, r_u r_d, (\sigma_{1,2,3})^2$$

(B-type masses)

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

$$\longrightarrow \text{Dirac gaugino masses}$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\longrightarrow \text{Majorana gaugino masses}$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow h_u h_d, r_u r_d, (\sigma_{1,2,3})^2$$

(B-type masses)

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

$$\longrightarrow \text{Dirac gaugino masses}$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\longrightarrow \text{Majorana gaugino masses}$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow H_u H_d, R_u R_d \text{ etc.}$$

(μ -type masses)

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi \quad (\Phi = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow \tilde{q}^\dagger \tilde{q}, h_u^\dagger h_u, \sigma_3^\dagger \sigma_3 \text{ etc.}$$

(Scalar soft masses)

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow h_u h_d, r_u r_d, (\sigma_{1,2,3})^2$$

(B-type masses)

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W_{1,2,3} \Sigma_{1,2,3} \quad (\langle W'_\alpha \rangle = \theta_\alpha D)$$

$$\longrightarrow \text{Dirac gaugino masses}$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\longrightarrow \text{Majorana gaugino masses}$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \quad (\Phi, \Phi' = H_{u,d}, R_{u,d}, \Sigma_{1,2,3})$$

$$\longrightarrow H_u H_d, R_u R_d \text{ etc.}$$

(μ -type masses)

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi'' \quad (\Phi, \Phi', \Phi'' = Q, \dots, H, R, \Sigma)$$

$$\longrightarrow h_u \tilde{q} \tilde{u}^c, \sigma_1 \sigma_3 \sigma_3, h_u \sigma_2 h_d \text{ etc.}$$

(A-type couplings)

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi$$

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi'$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W \Sigma$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} W W$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi''$$

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi$$

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi'$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W \Sigma$$

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} WW$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi''$$

Forbidden if X
is not a singlet.

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi$$

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi'$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W \Sigma$$

Allowed even if X
is not a singlet.

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} W W$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi''$$

Forbidden if X
is not a singlet.

Accidental $U(1)_R$ in visible sector

[What we want]

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi^\dagger \Phi$$

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} \Phi \Phi'$$

$$\int d^2\theta \frac{W'}{M_{\text{Pl}}} W \Sigma$$

Allowed even if X
is not a singlet.

[What we don't want]

$$\int d^2\theta \frac{X}{M_{\text{Pl}}} W W$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$$

$$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi' \Phi''$$

Forbidden if X
is not a singlet.

If no singlets in ~~SUSY~~ sector,

→ Accidental $U(1)_R$ in visible sector!

An Example Hidden Sector

Criteria:

- * No gauge singlet for accidental R symm.
- * $\langle D \rangle \sim \langle F \rangle$ for Dirac gauginos

An Example Hidden Sector

Criteria:

- * No gauge singlet for accidental R symm.
- * $\langle D \rangle \sim \langle F \rangle$ for Dirac gauginos

Very simple example — "4-1 model": (Dine, Nelson, Nir, Shirman, '95)

	$\bar{\Phi}$	S	A	Φ
$SU(4)$	$\bar{4}$	1	6	4
$U(1)$	-1	4	2	-3

$$\mathcal{W} = \lambda S \Phi \bar{\Phi}$$

An Example Hidden Sector

Criteria:

- * No gauge singlet for accidental R symm.
- * $\langle D \rangle \sim \langle F \rangle$ for Dirac gauginos

Very simple example — "4-1 model": (Dine, Nelson, Nir, Shirman, '95)

		$\bar{\Phi}$	S	A	Φ
$SU(5)$	$SU(4)$	$\bar{4}$	1	6	4
	$U(1)$	-1	4	2	-3
		$5\bar{ar}$		10	

$$\mathcal{W} = \lambda S \Phi \bar{\Phi}$$

An Example Hidden Sector

Criteria:

- * No gauge singlet for accidental R symm.
- * $\langle D \rangle \sim \langle F \rangle$ for Dirac gauginos

Very simple example — "4-1 model": (Dine, Nelson, Nir, Shirman, '95)

		$\bar{\Phi}$	S	A	Φ
$SU(5)$	$SU(4)$	$\bar{4}$	1	6	4
	$U(1)$	-1	4	2	-3

5bar 10

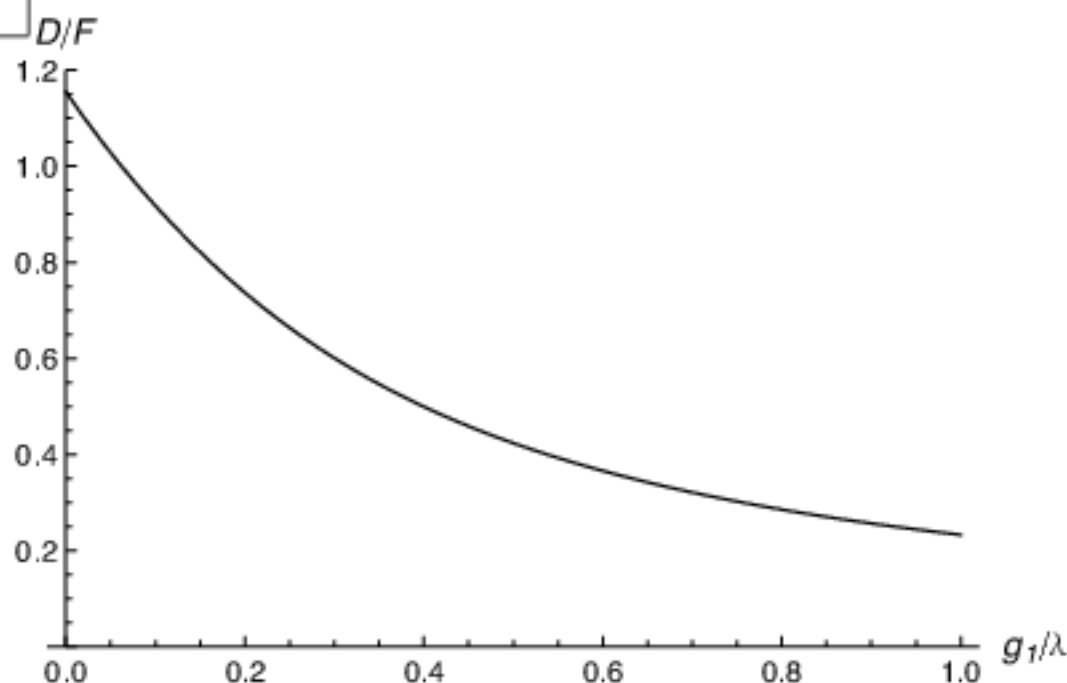
$$\mathcal{W} = \lambda S \Phi \bar{\Phi}$$

$$g_4 \gg \lambda \gtrsim g_1$$

$$\longrightarrow \langle D_1 \rangle \sim \langle F \rangle !$$

(Carpenter, Kaplan, Fox, '05;
Gregoire, Rattazzi, Scrucce, '05)

Simple existence proof!



Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}$?

$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q \quad + \quad \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$

Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}?$

$$\begin{array}{c} 0.01 \\ ? \end{array} \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q \quad + \quad \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$

$m_{\tilde{q}}^2$
 $m_{\tilde{g}}$

Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}$?

$$0.01 \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$

Look at RG effects:

$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ M_{\text{Pl}}$$

$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ 1 \text{ TeV}$$

Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}?$

$$\underset{?}{0.01} \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$

Look at RG effects:

$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \underset{!}{1} \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ M_{\text{Pl}}$$
$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \underset{!}{5.4} \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ 1 \text{ TeV}$$

Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}?$

$$\underset{?}{0.01} \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$



Look at RG effects:

$$\int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \mathbf{1} \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ M_{\text{Pl}}$$



$$\mathbf{0.3} \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \mathbf{5.4} \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ 1 \text{ TeV}$$

Can we get the hierarchy $m_{\tilde{g}} \sim \underline{10} m_{\tilde{q}}?$

$$0.01 \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma$$

$m_{\tilde{q}}^2$
 $m_{\tilde{g}}$

Look at RG effects:

$$0.3 \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + 1 \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ M_{\text{Pl}}$$

$\updownarrow$
 (Dirac gauginos $\rightarrow$ 1-loop
 running **only by Yukawa!**)

$\updownarrow$
RG

$$0.3 \int d^4\theta \frac{\theta^4 \langle F \rangle^2}{M_{\text{Pl}}^2} Q^\dagger Q + \underline{5.4} \int d^2\theta \frac{\theta \langle D \rangle}{M_{\text{Pl}}} W \Sigma \quad @ 1 \text{ TeV}$$

The hierarchy arises naturally!

What about hierarchies inside $m_{\tilde{q}}^2$?

Flavor $\longrightarrow m_{1/2} \sim 10m_{\tilde{q}}, c_{12} \sim c_{11}/10 \sim c_{22}/10$

CP ($\text{Im } K\text{-}\overline{K}$) $\longrightarrow \theta_{\text{CP}} \sim 1/10$ (Blechman, Ng, '08)

What about hierarchies inside $m_{\tilde{q}}^2$?

Flavor $\longrightarrow m_{1/2} \sim 10 m_{\tilde{q}}, c_{12} \sim c_{11}/10 \sim c_{22}/10$

CP ($\text{Im } K\text{-}\overline{K}$) $\longrightarrow \theta_{\text{CP}} \sim 1/10$ (Blechman, Ng, '08)

But c_{ij} here is not the "fundamental" c_{ij} :

$\begin{array}{c} \text{E} \\ \uparrow \\ M_{\text{Pl}} \end{array} (c_{ij}^{\text{UV}}) = \mathcal{O}(0.1) \begin{pmatrix} \sim 1 & \sim 1 & \sim 1 \\ \sim 1 & \sim 1 & \sim 1 \\ \sim 1 & \sim 1 & \sim 1 \end{pmatrix} \quad \Delta\theta_{\text{CP}} \sim 0.1$

What about hierarchies inside $m_{\tilde{q}}^2$?

Flavor $\longrightarrow m_{1/2} \sim 10 m_{\tilde{q}}, c_{12} \sim c_{11}/10 \sim c_{22}/10$

CP ($\text{Im } K\text{-}\overline{K}$) $\longrightarrow \theta_{\text{CP}} \sim 1/10$ (Blechman, Ng, '08)

But c_{ij} here is not the "fundamental" c_{ij} :

$$\begin{array}{c} \text{E} \uparrow \\ M_{\text{Pl}} \text{---} (c_{ij}^{\text{UV}}) = \mathcal{O}(0.1) \begin{pmatrix} \sim 1 & \sim 1 & \sim 1 \\ \sim 1 & \sim 1 & \sim 1 \\ \sim 1 & \sim 1 & \sim 1 \end{pmatrix} \Delta\theta_{\text{CP}} \sim 0.1 \end{array}$$

Flavor universal, no phases!

$m_{\tilde{g}}$ 5 TeV

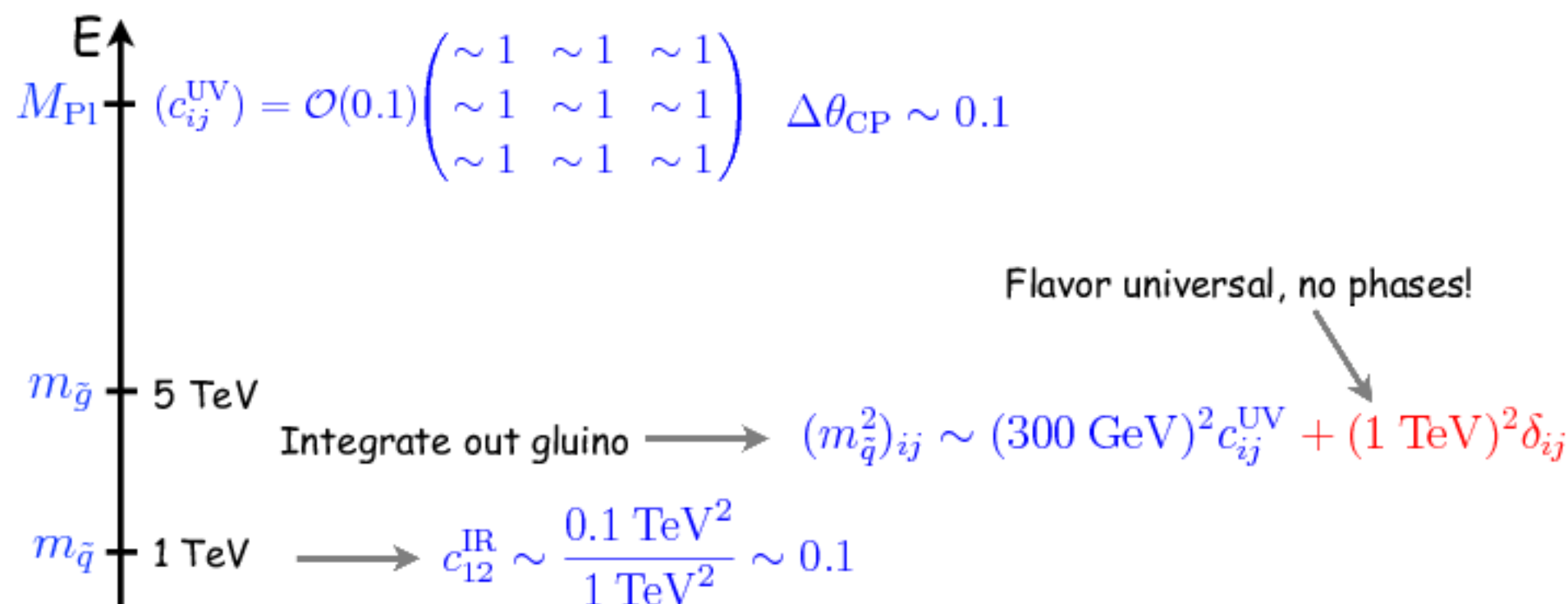
Integrate out gluino $\longrightarrow (m_{\tilde{q}}^2)_{ij} \sim (300 \text{ GeV})^2 c_{ij}^{\text{UV}} + (1 \text{ TeV})^2 \delta_{ij}$

What about hierarchies inside $m_{\tilde{q}}^2$?

Flavor $\longrightarrow m_{1/2} \sim 10 m_{\tilde{q}}, c_{12} \sim c_{11}/10 \sim c_{22}/10$

CP ($\text{Im } K\text{-}\overline{K}$) $\longrightarrow \theta_{\text{CP}} \sim 1/10$ (Blechman, Ng, '08)

But c_{ij} here is not the "fundamental" c_{ij} :



"Fundamental" c_{ij} can be "anarchic"!

A new $U(1)$ gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \quad \text{would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

A new $U(1)$ gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

→ Need a new gauge symmetry $U(1)_{\text{“N”}_{\text{ew}}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

A new $U(1)$ gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

→ Need a new gauge symmetry $U(1)_{\text{“N” ew}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

Can be naturally broken at TeV: $\langle S \rangle \neq 0$ by $m_s^2 < 0$

↖ (Naturally TeV)

A new U(1) gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

→ Need a new gauge symmetry $U(1)_{\text{N ew}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

Can be naturally broken at TeV: $\langle S \rangle \neq 0$ by $m_s^2 < 0$

↖ (Naturally TeV)

The new U(1) solves two additional problems!

(1) No singlets → ~~$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$~~ → ~~$\mu H_u H_d$~~ good!
 ~~$H_u R_u + H_d R_d$~~ oops!

A new U(1) gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

→ Need a new gauge symmetry $U(1)_{\text{“N” ew}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

Can be naturally broken at TeV: $\langle S \rangle \neq 0$ by $m_s^2 < 0$

↖ (Naturally TeV)

The new U(1) solves two additional problems!

(1) No singlets → ~~$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$~~ → ~~$\mu H_u H_d$~~ good!
 ~~$H_u R_u + H_d R_d$~~ oops!

But we have $\langle S \rangle \neq 0$! Just assign

$$N[S_u] = -N[S_d] = -1 \quad \& \quad R[S_{u,d}] = 0$$

$$\longrightarrow \mathcal{W} = S_u H_u R_u + S_d H_d R_d \longrightarrow \langle S_u \rangle H_u R_u + \langle S_d \rangle H_d R_d!$$

A new U(1) gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R!$$

(we want $R[H_{u,d}] = 0$, $R[R_{u,d}] = 2$)

→ Need a new gauge symmetry $U(1)_{\text{N}^{\text{ew}}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

Can be naturally broken at TeV: $\langle S \rangle \neq 0$ by $m_s^2 < 0$
↖ (Naturally TeV)

The new U(1) solves two additional problems!

(1) No singlets → ~~$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$~~ → ~~$\mu H_u H_d$~~ good!
 ~~$H_u R_u + H_d R_d$~~ oops!

But we have $\langle S \rangle \neq 0$! Just assign

$$N[S_u] = -N[S_d] = -1 \quad \& \quad R[S_{u,d}] = 0$$

$$\longrightarrow \mathcal{W} = S_u H_u R_u + S_d H_d R_d \longrightarrow \langle S_u \rangle H_u R_u + \langle S_d \rangle H_d R_d!$$

(2) LEP Higgs bound?

A new U(1) gauge symmetry

A serious problem:

$$\int d^4\theta \frac{X^\dagger X}{M_{\text{Pl}}^2} R_{d,u}^\dagger H_{u,d} \text{ would kill accidental } U(1)_R! \quad (\text{we want } R[H_{u,d}] = 0, R[R_{u,d}] = 2)$$

→ Need a new gauge symmetry $U(1)_{\text{N}^{\text{ew}}}$! (No global symm with gravity)

$$N[R_u] = -N[R_d] = 1, \quad N[H, Q, U, D, L, E] = 0$$

Can be naturally broken at TeV: $\langle S \rangle \neq 0$ by $m_s^2 < 0$

(Naturally TeV)

The new U(1) solves two additional problems!

(1) No singlets → ~~$\int d^4\theta \frac{X^\dagger}{M_{\text{Pl}}} \Phi \Phi'$~~ → ~~$\mu H_u H_d$~~ good!
 ~~$H_u R_u + H_d R_d$~~ oops!

But we have $\langle S \rangle \neq 0$! Just assign

$$N[S_u] = -N[S_d] = -1 \quad \& \quad R[S_{u,d}] = 0$$

$$\longrightarrow \mathcal{W} = S_u H_u R_u + S_d H_d R_d \longrightarrow \langle S_u \rangle H_u R_u + \langle S_d \rangle H_d R_d!$$

(2) LEP Higgs bound? $\mathcal{O}(1) \longrightarrow$ Large H quartic! $\longrightarrow m_H > 114 \text{ GeV}!$

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \rightarrow \begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix} \quad R = -1 \quad R_u \rightarrow \begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix} \quad R = 1 \quad S_u \rightarrow \tilde{s}_u \quad R = -1$$

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \left[\begin{array}{c} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{array} \right] \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \left[\begin{array}{c} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{array} \right] \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{S}_u \quad R = -1$$

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix}$$



H_u



$$\begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix}$$

$$R = -1$$

R_u

$$\begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix}$$

$$R = 1$$

S_u

$\rightarrow \tilde{s}_u$

$$R = -1$$

Charginos good!

1 massless neutralino!?

Dirty Laundries...

* Neutralino masses

$$\left(\begin{array}{l} R[H] = R[S] = 0 \\ R[R] = 2 \end{array} \right) \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \left[\begin{array}{c} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{array} \right] \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \left[\begin{array}{c} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{array} \right] \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{s}_u \text{?} \\ R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{s}_{u,d} \tilde{t}_{u,d}$$

Dirty Laundries...

* Neutralino masses

$$\left(\begin{array}{l} R[H] = R[S] = 0 \\ R[R] = 2 \end{array} \right) \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \left[\begin{array}{c} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{array} \right] \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \left[\begin{array}{c} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{array} \right] \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{s}_u \text{?} \\ R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{s}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix} \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix} \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{S}_u? \quad R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{S}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Embed $U(1)$ s into non-Abelian groups.

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix} \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix} \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{S}_u? \quad R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{S}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Embed $U(1)$ s into non-Abelian groups.

* Unification

$$\text{MSSM} + \Sigma_{2,3} + R_{u,d} + \text{SM singlets}$$

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix} \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix} \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{s}_u? \quad R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{s}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Embed $U(1)$ s into non-Abelian groups.

* Unification

Just perturbative ($b_3 = 0$)!

$$\text{MSSM} + \Sigma_{2,3} + R_{u,d} + \text{SM singlets} + 2 \times [(\mathbf{1}, \mathbf{1})_1 \oplus (\mathbf{1}, \mathbf{1})_{-1}] = \text{unifies!}$$

Dirty Laundries...

* Neutralino masses

$$\begin{pmatrix} R[H] = R[S] = 0 \\ R[R] = 2 \end{pmatrix} \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \begin{bmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{bmatrix} \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \begin{bmatrix} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{bmatrix} \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{S}_u? \quad R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{S}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Embed $U(1)$ s into non-Abelian groups.

* Unification

Just perturbative ($b_3 = 0$)!

$$\text{MSSM} + \Sigma_{2,3} + R_{u,d} + \text{SM singlets} + 2 \times [(\mathbf{1}, \mathbf{1})_1 \oplus (\mathbf{1}, \mathbf{1})_{-1}] = \text{unifies!}$$

Maybe trinification? $[SU(3)_C \otimes SU(3)_L \otimes SU(3)_R] / \mathbb{Z}_3$

$$(\mathbf{8}, \mathbf{1}, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{8}, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{1}, \mathbf{8}) \longrightarrow \Sigma_{1,2,3} \oplus R_{u,d} \oplus S_{u,d} \oplus 2 \times [(\mathbf{1}, \mathbf{1})_1 \oplus (\mathbf{1}, \mathbf{1})_{-1}]$$

Dirty Laundries...

* Neutralino masses

$$\left(\begin{array}{l} R[H] = R[S] = 0 \\ R[R] = 2 \end{array} \right) \longrightarrow H_u \longrightarrow \begin{array}{c} \text{Charginos good!} \\ \left[\begin{array}{c} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{array} \right] \\ R = -1 \end{array} \xrightarrow{R_u} \begin{array}{c} \text{1 massless neutralino!} \\ \left[\begin{array}{c} \tilde{r}_u^0 \\ \tilde{r}_u^- \end{array} \right] \\ R = 1 \end{array} \quad S_u \longrightarrow \tilde{S}_u? \quad R = -1$$

Need $T_{u,d}$ with $R[T_{u,d}] = 2$. Then,

$$\mathcal{W} = S_{u,d} S_{u,d} T_{u,d} \longrightarrow \langle S_{u,d} \rangle \tilde{S}_{u,d} \tilde{t}_{u,d}$$

* Tadpoles for singlets $\Sigma_{1,N}, \sigma_{1,N}$

(chiral "adjoints" for $U(1)_Y, U(1)_N$)

Embed $U(1)$ s into non-Abelian groups.

* Unification

Just perturbative ($b_3 = 0$)!

MSSM + $\Sigma_{2,3} + R_{u,d} + \text{SM singlets} + 2 \times [(\mathbf{1}, \mathbf{1})_1 \oplus (\mathbf{1}, \mathbf{1})_{-1}] = \text{unifies!}$

Maybe trinification? $[SU(3)_C \otimes SU(3)_L \otimes SU(3)_R] / \mathbb{Z}_3$

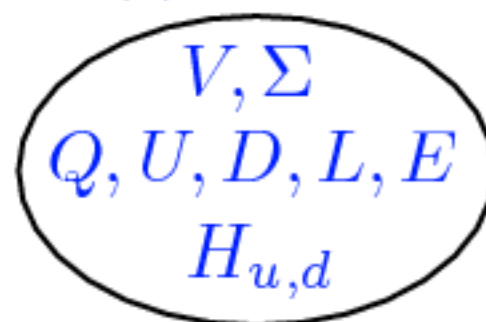
$$(\mathbf{8}, \mathbf{1}, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{8}, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{1}, \mathbf{8}) \longrightarrow \Sigma_{1,2,3} \oplus R_{u,d} \oplus S_{u,d} \oplus 2 \times [(\mathbf{1}, \mathbf{1})_1 \oplus (\mathbf{1}, \mathbf{1})_{-1}]$$

But what about $U(1)_N$? $T_{u,d}$? Quartification $[SU(3)]^4$?

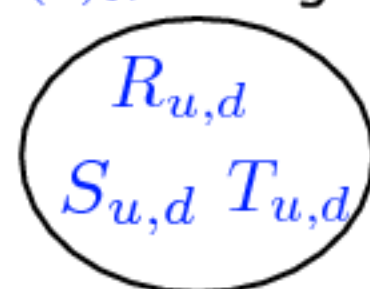
Collider Prospects

$U(1)_N$ will be very interesting!

$U(1)_N$ -neutral



$U(1)_N$ -charged



Collider Prospects

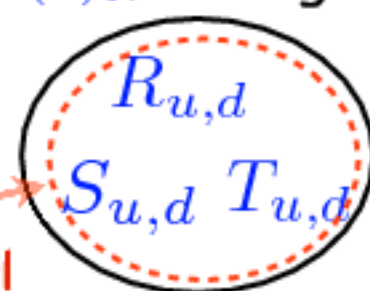
$U(1)_N$ will be very interesting!

A Higgs-philic, fermi-phobic Z' !

$U(1)_N$ -neutral



$U(1)_N$ -charged



Mix!

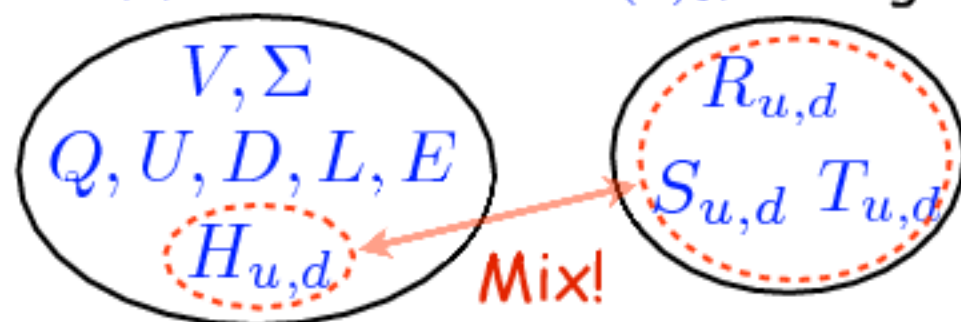
Collider Prospects

$U(1)_N$ will be very interesting!

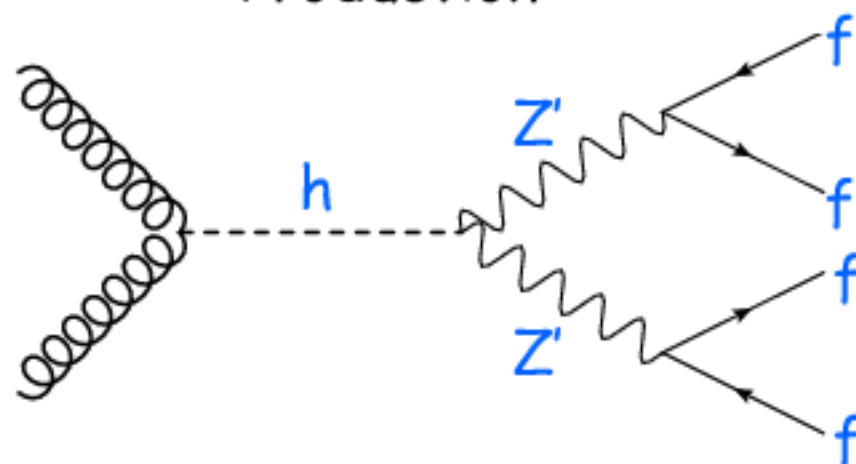
A Higgs-philic, fermi-phobic Z' !

$U(1)_N$ -neutral

$U(1)_N$ -charged



Production:



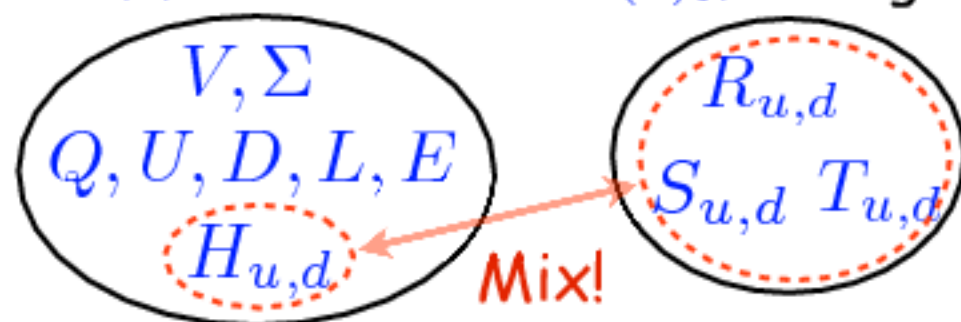
Collider Prospects

$U(1)_N$ will be very interesting!

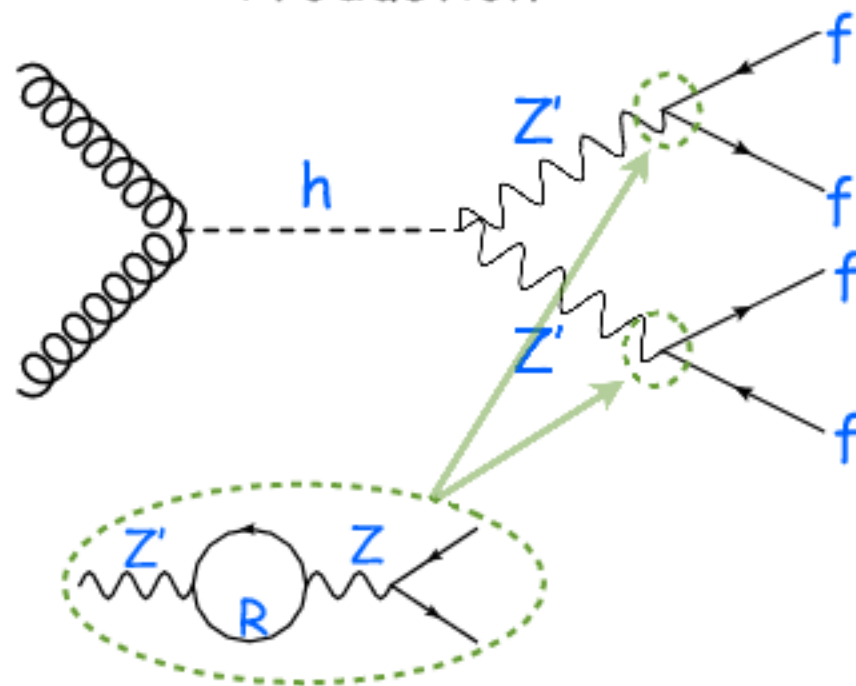
A Higgs-philic, fermi-phobic Z' !

$U(1)_N$ -neutral

$U(1)_N$ -charged



Production:



Tiny, only loop-induced

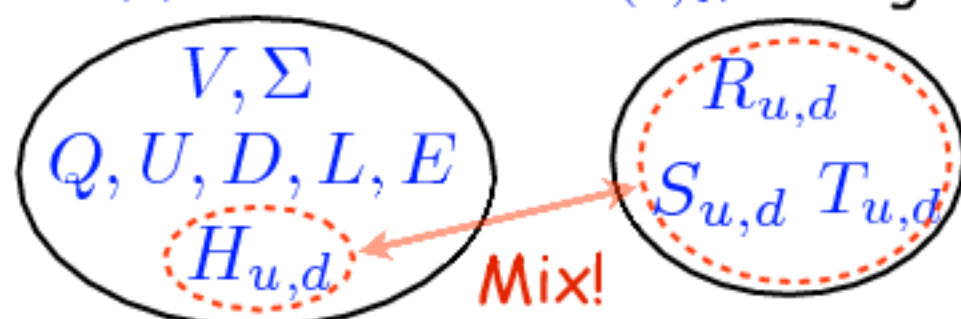
Collider Prospects

$U(1)_N$ will be very interesting!

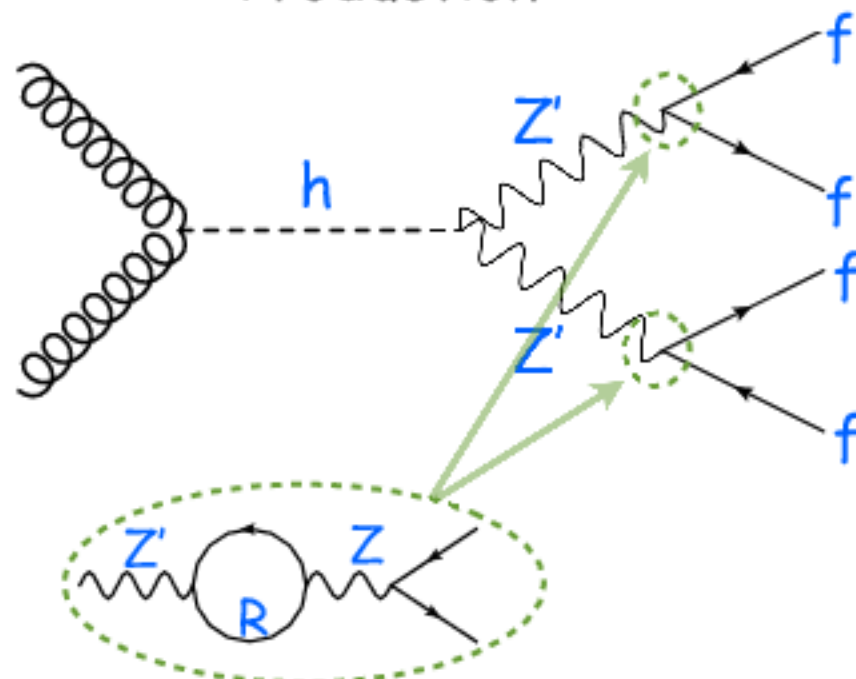
A Higgs-philic, fermi-phobic Z' !

$U(1)_N$ -neutral

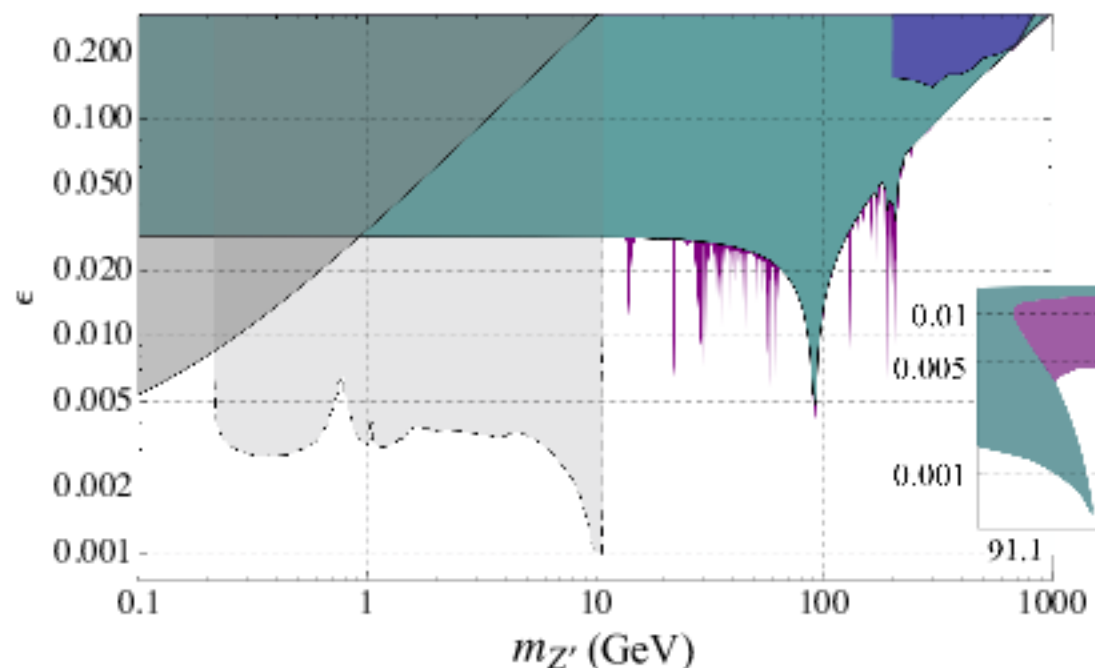
$U(1)_N$ -charged



Production:



Tiny, only loop-induced

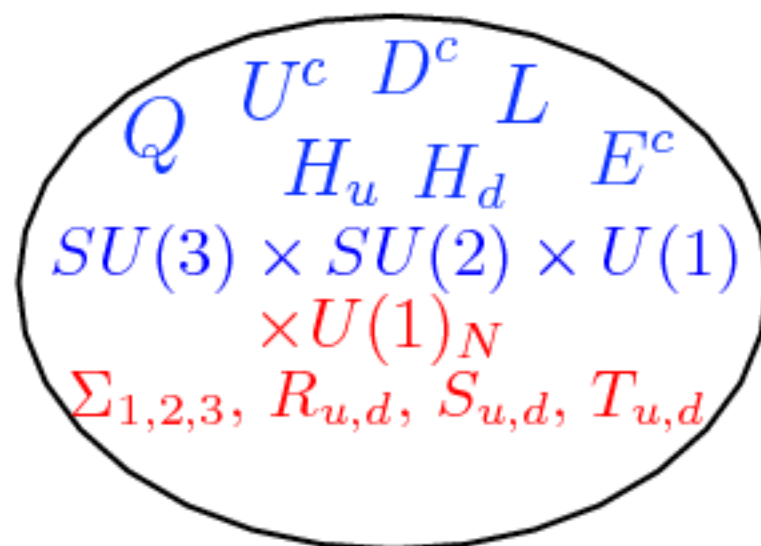


(Hook, Izaguirre, Wacker, 1006.0973)

Z' Can be very light!
Possibly at early LHC!?

Summary

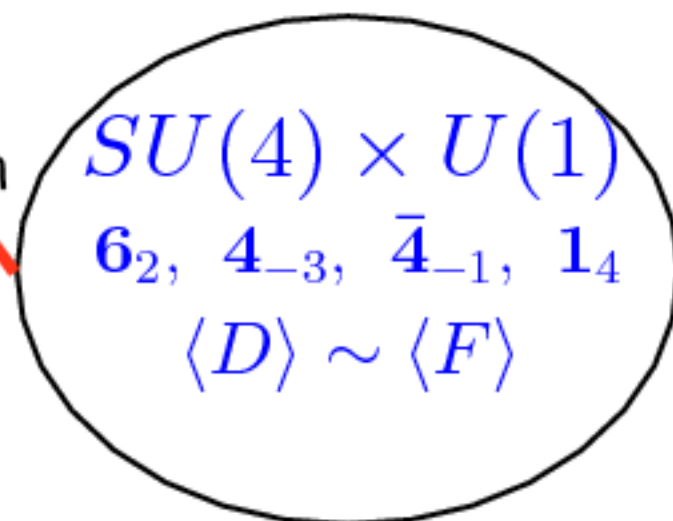
"Visible" sector



gravity mediation

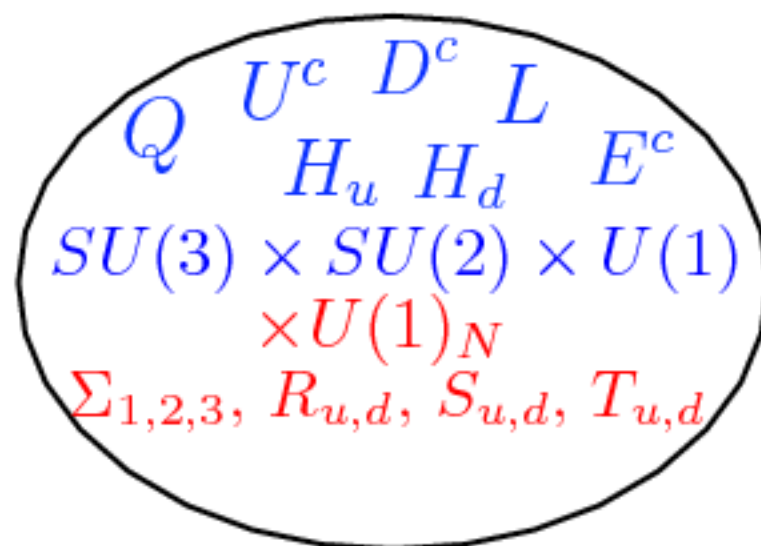
$$\frac{\mathcal{O}_{\text{vis.}} \mathcal{O}_{\text{S.B.}}}{M_{\text{Pl}}^{\#}}$$

~~SUSY~~ sector



Summary

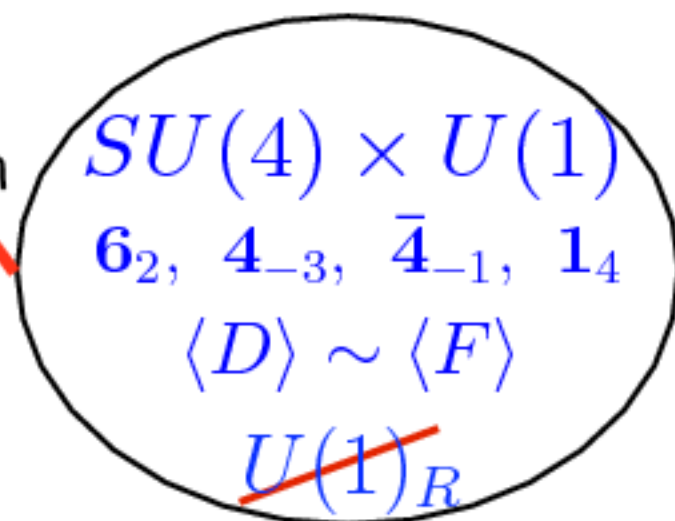
"Visible" sector



gravity mediation

$$\frac{\mathcal{O}_{\text{vis.}} \mathcal{O}_{\text{S.B.}}}{M_{\text{Pl}}^{\#}}$$

~~SUSY~~ sector

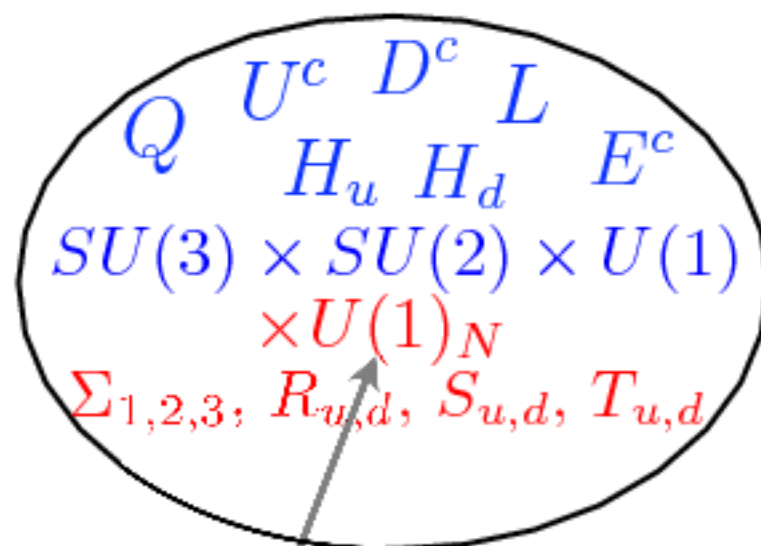


Accidental $U(1)_R$! (SUSY flavor problem solved!)

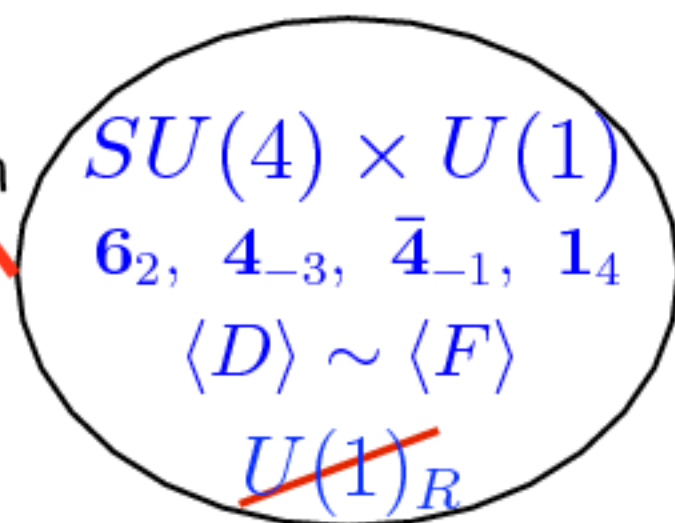
All mass scales right (esp. $m_{\tilde{g}} \sim 10m_{\tilde{q}}$ from RG)
 no mu-Bmu problem, neutralino DM.

Summary

"Visible" sector



~~SUSY~~ sector



gravity mediation

$$\frac{\mathcal{O}_{\text{vis.}} \mathcal{O}_{\text{S.B.}}}{M_{\text{Pl}}^{\#}}$$

Accidental $U(1)_R$! (SUSY flavor problem solved!)

All mass scales right (esp. $m_{\tilde{g}} \sim 10m_{\tilde{q}}$ from RG)
 no mu-Bmu problem, neutralino DM.

Essential for accidental $U(1)_R$! Interesting LHC pheno!