

# Jet Substructure Measurements at STAR

**Diptanil Roy**

*(On behalf of the STAR Collaboration)*

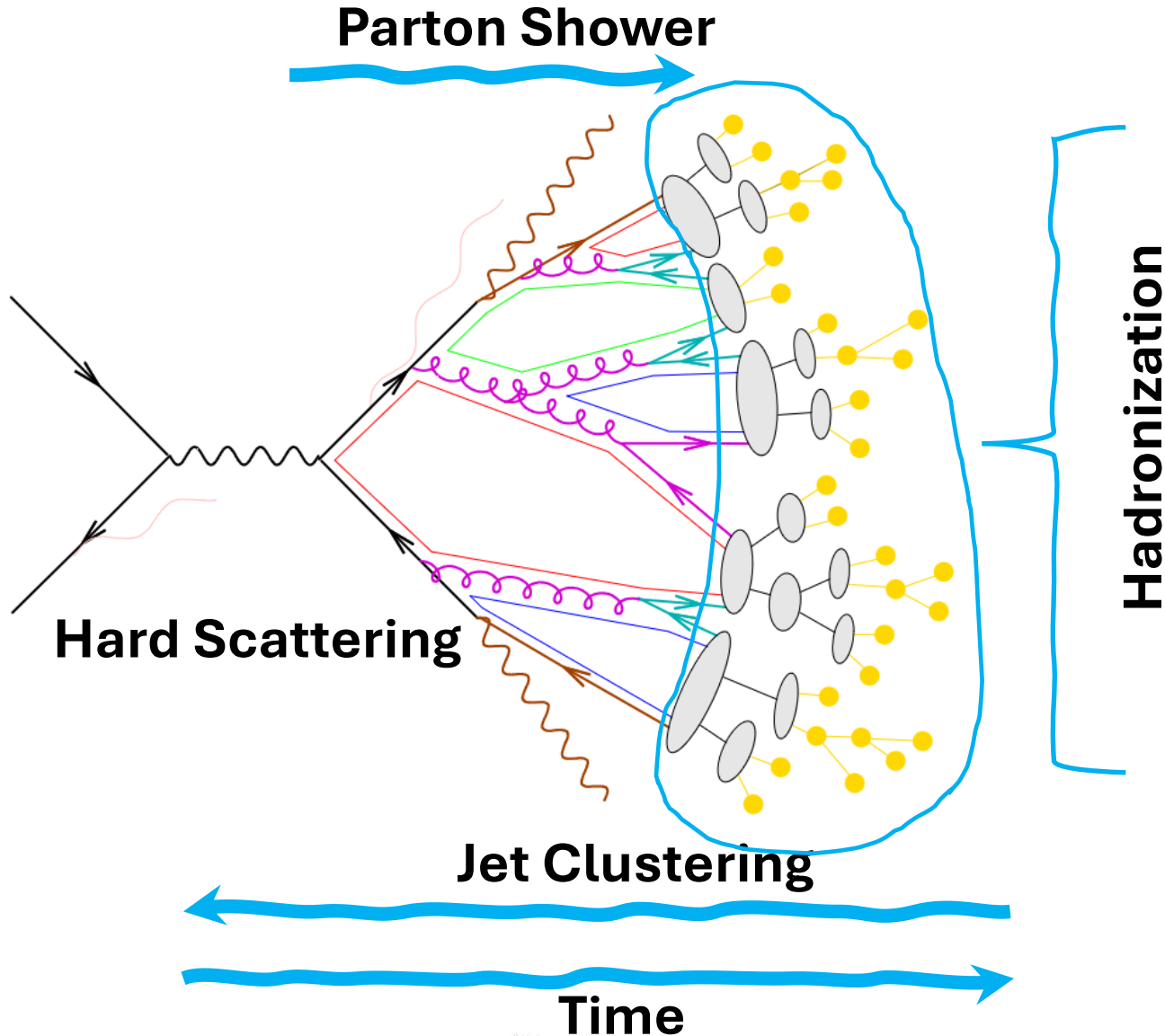
*Rutgers University*

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*July 28 – Aug 2, 2024*

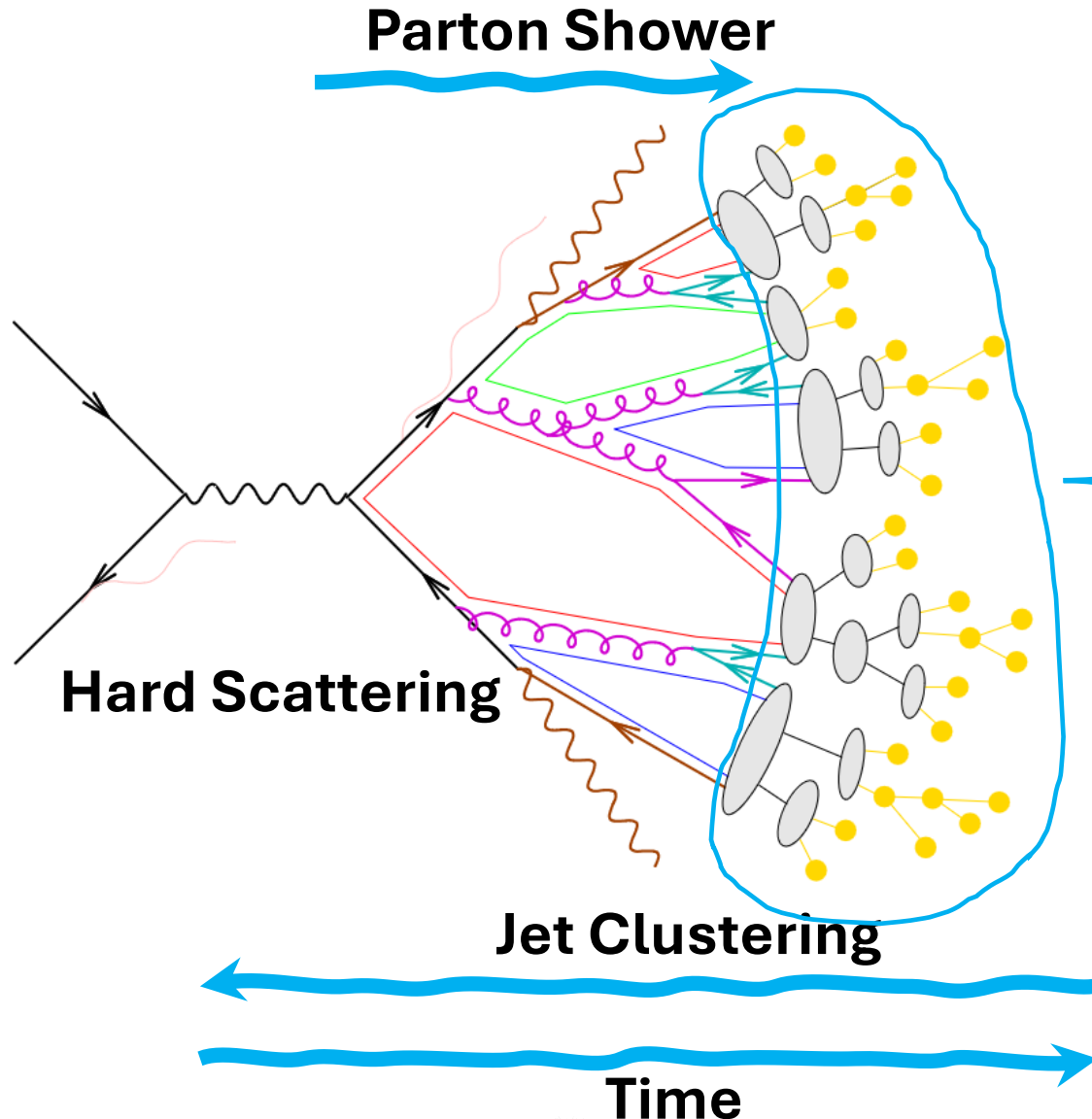


# Jets in vacuum



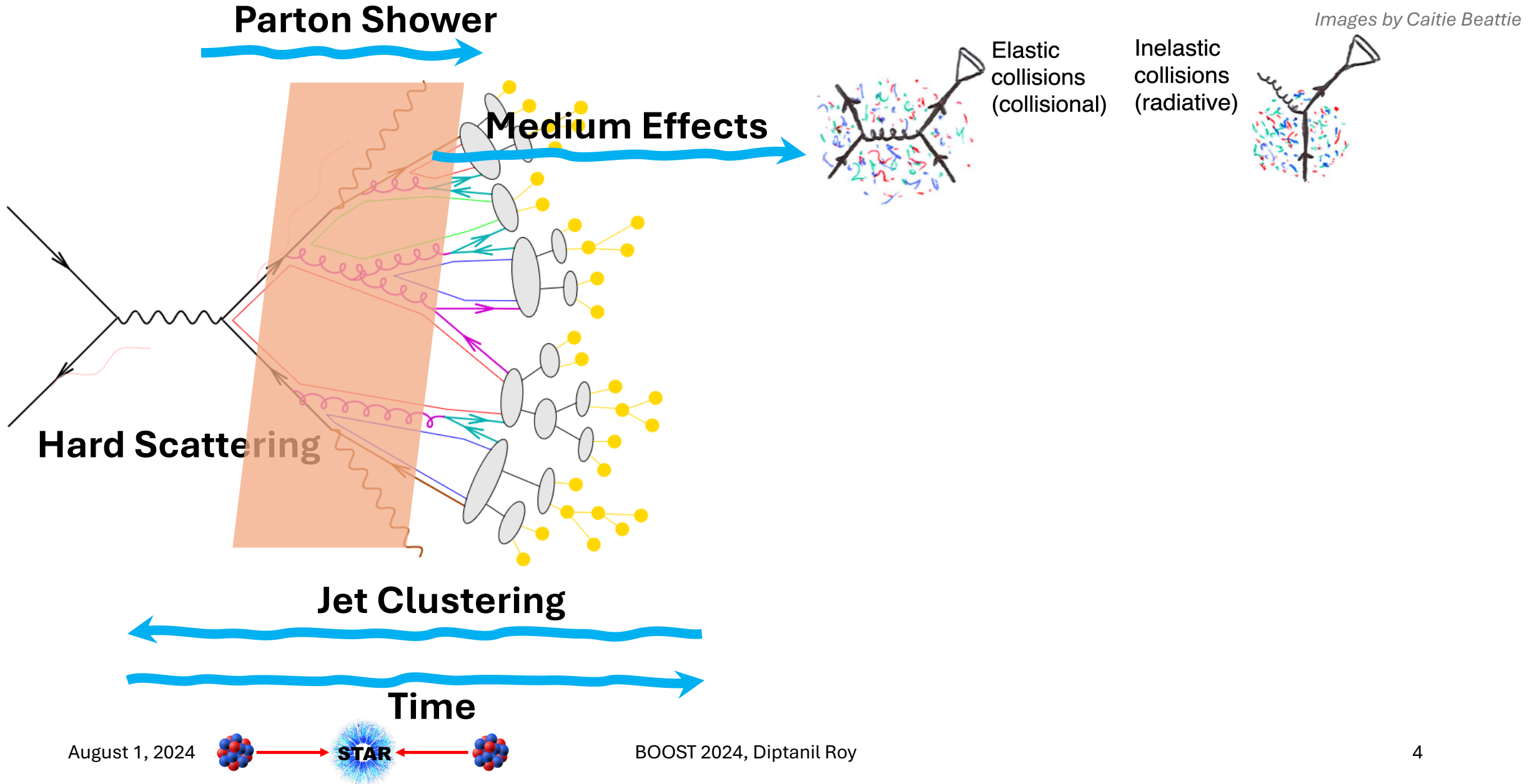
- Proxies for hard scattered partons
- Production explained by pQCD
- Clustering algorithms use final state particles to reconstruct jets
- Jet substructure holds information about fragmentation and hadronization processes

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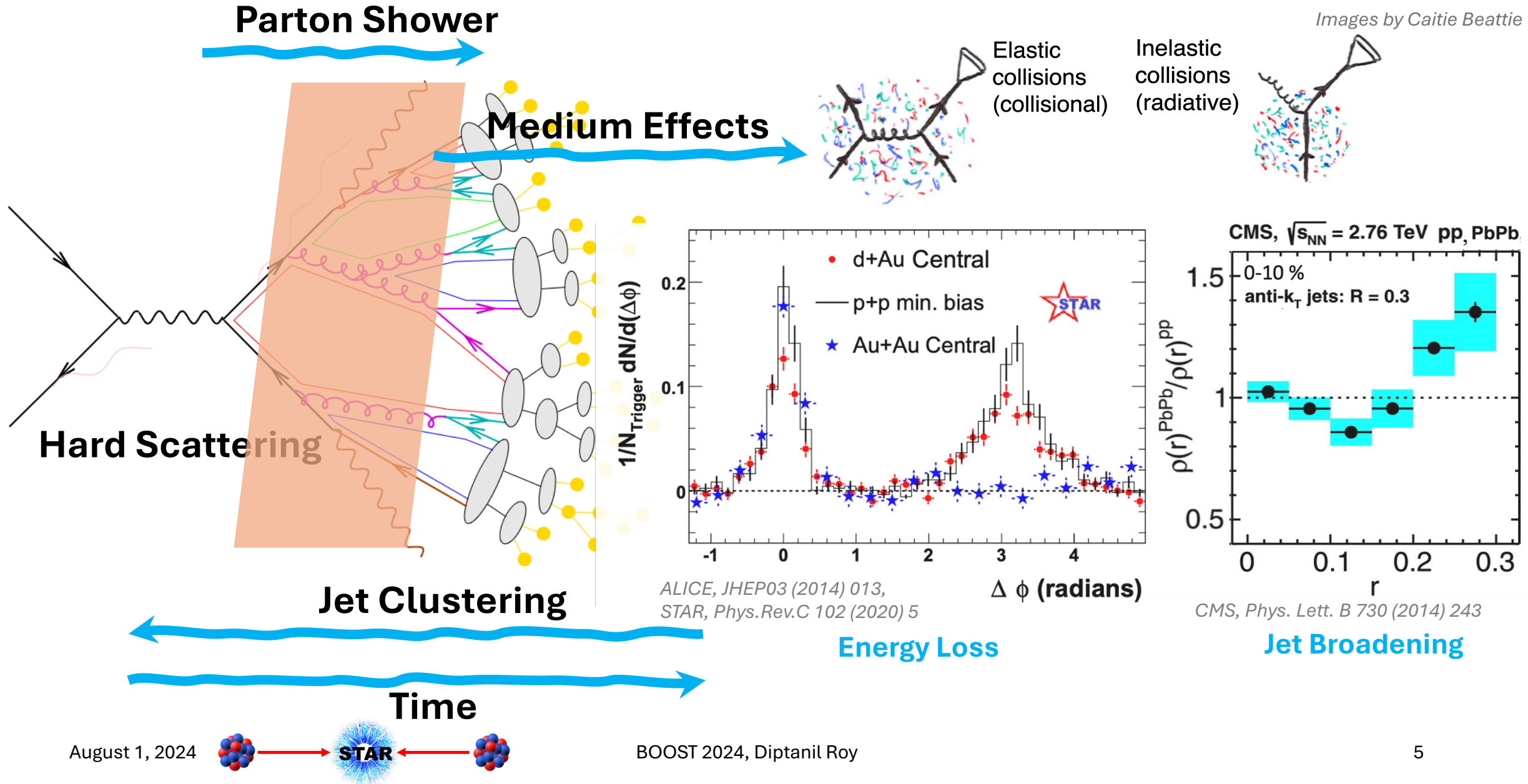


- Proxies for hard scattered partons
  - Production explained by pQCD
  - Clustering algorithms use final state particles to reconstruct jets
  - Jet substructure holds information about fragmentation and hadronization processes
- ✓ *Can we disentangle perturbative and non-perturbative physics in vacuum?*
- ✓ *Can we use jet substructure to understand hadron formation better?*

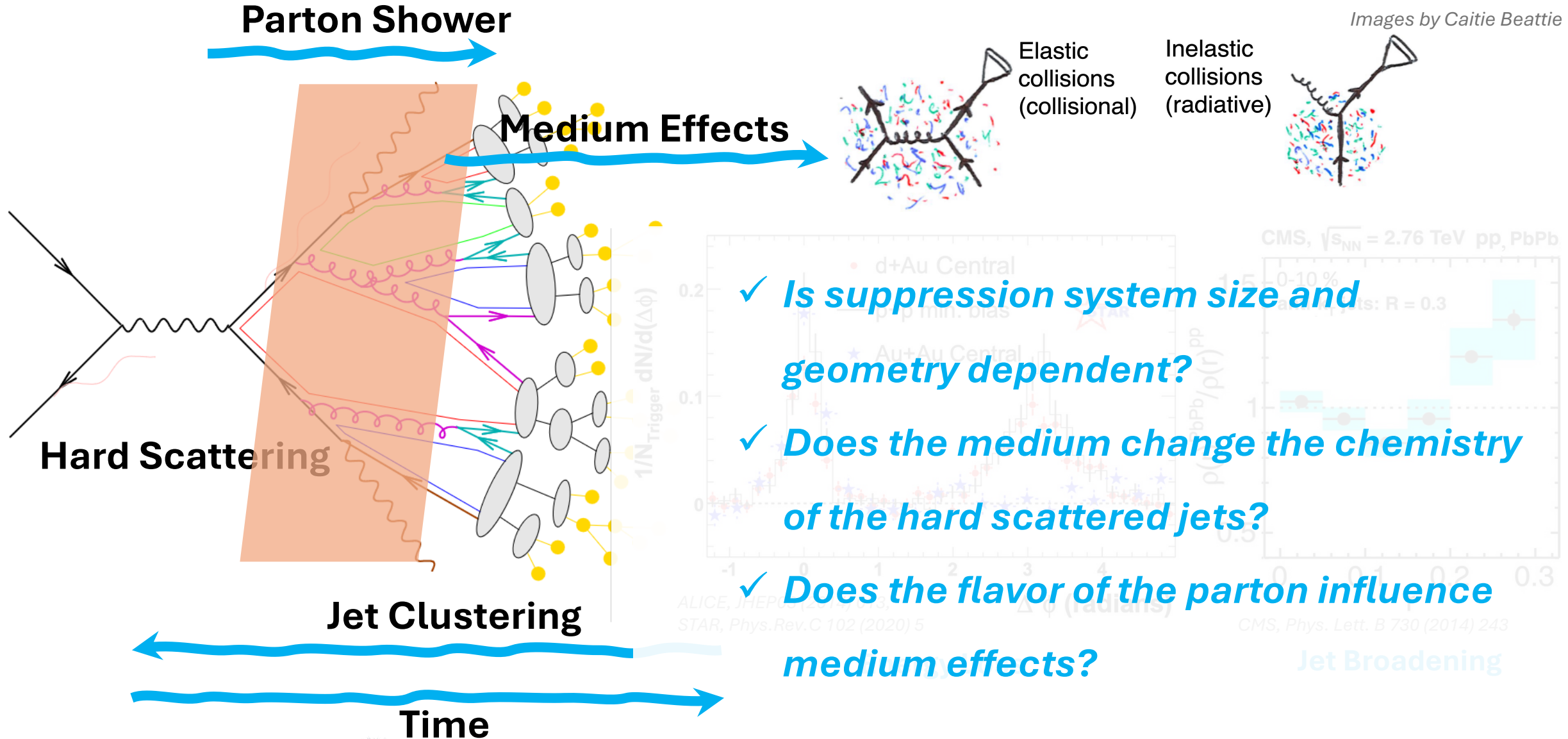
# Jets in medium



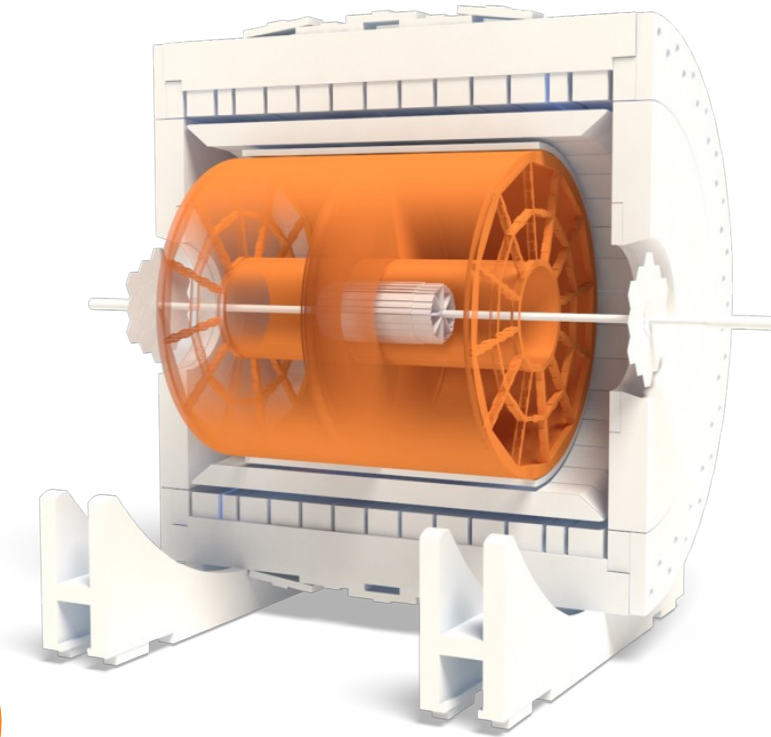
# Jets in medium



# Jets in medium



# STAR at RHIC



## Time Projection Chamber

- ✓ Measures momenta of charged tracks  $[|\eta| < 1, 0 < \phi < 2\pi]$
- ✓ PID using  $dE/dx$

Images: [NSWW](#)



# STAR at RHIC

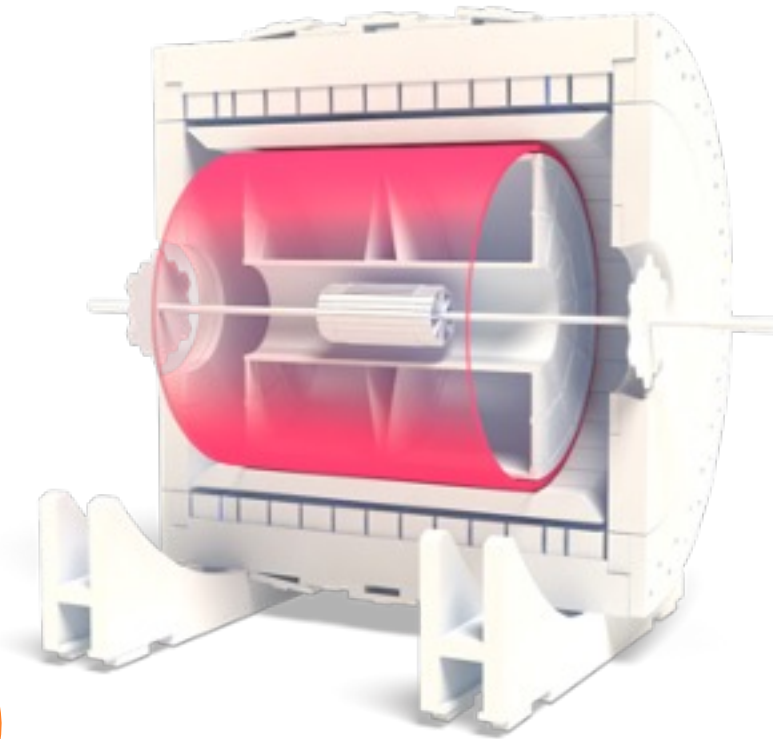
## Time-Of-Flight Detector

- ✓ PID using TOF measurement

$$[|\eta| < 1, 0 < \phi < 2\pi]$$

## Time Projection Chamber

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Images: [NSWW](#)





# STAR at RHIC

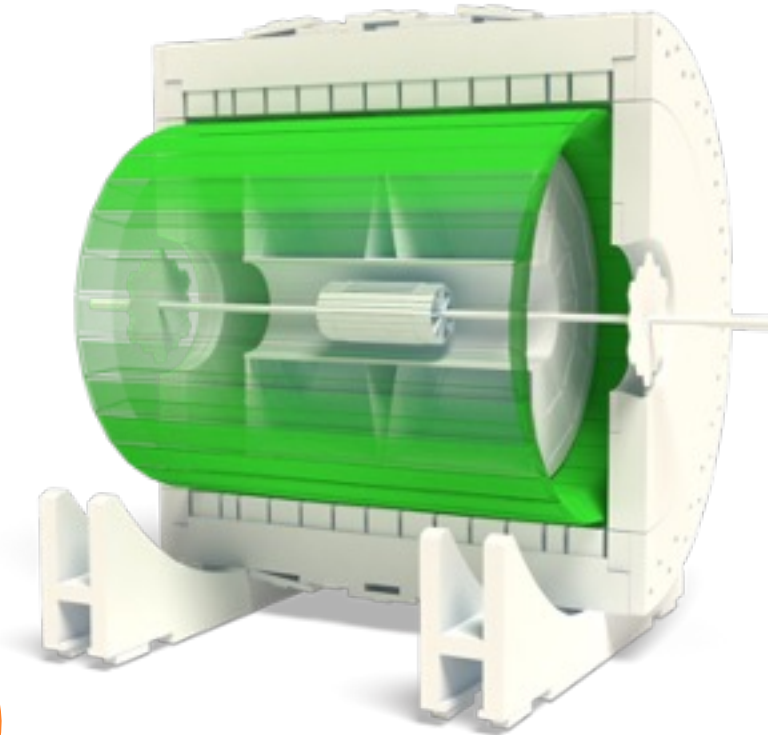
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## Barrel Electromagnetic Calorimeter

- ✓ Measures neutral component of jet energy  $[|\eta| < 1, 0 < \phi < 2\pi]$

Images: [NSWW](#)



# STAR at RHIC

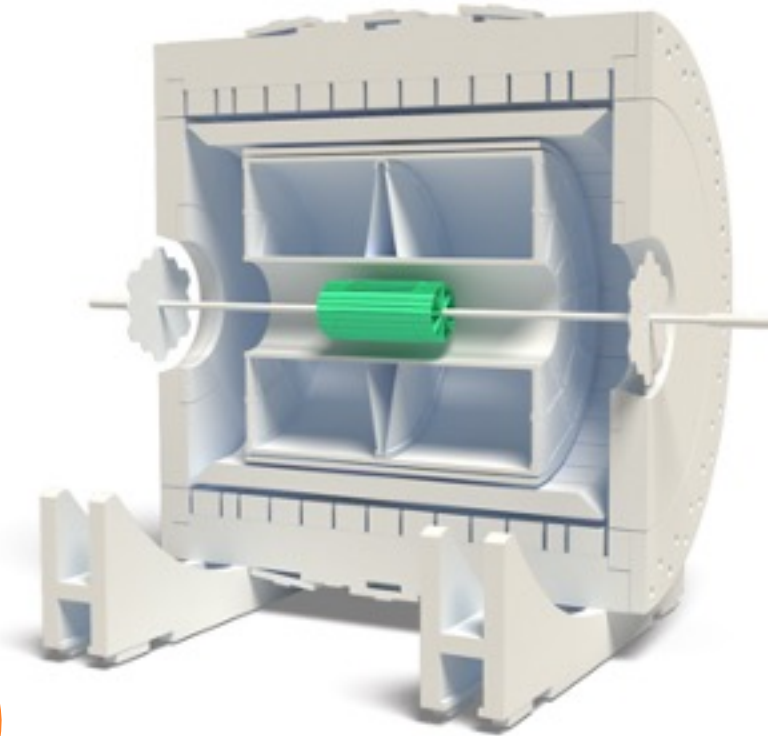
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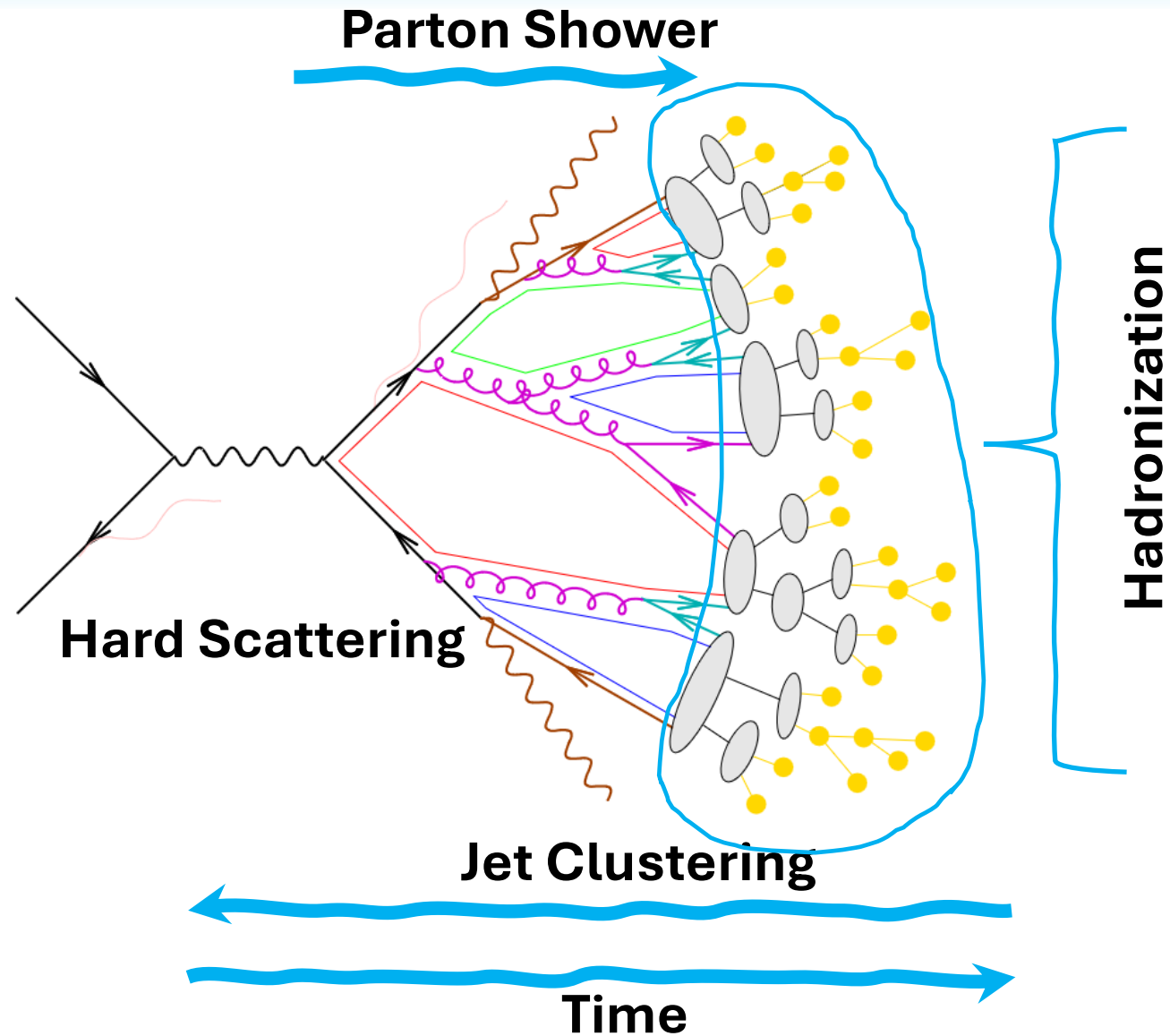
## Heavy Flavor Tracker (2014-2016)

- ✓ Improves position resolution for secondary vertices

Images: [NSWW](#)



# Jets in vacuum



# Isolating pQCD and npQCD in vacuum

$$\Delta M = M - M_g \text{ [GeV]}$$

Image: Larkoski, Marzani, Thaler, Xue, *PRL* 119 (2017) 13, 132003

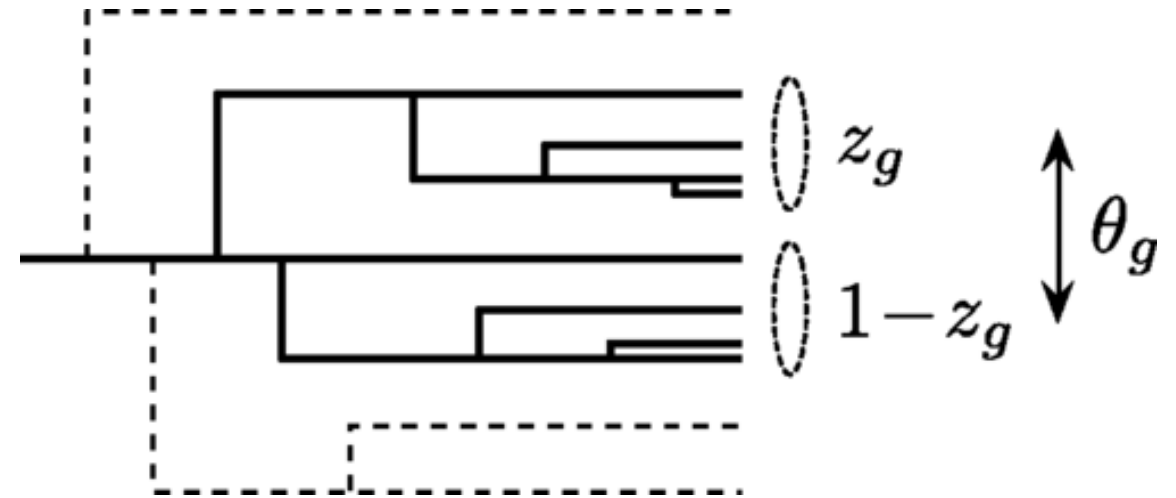
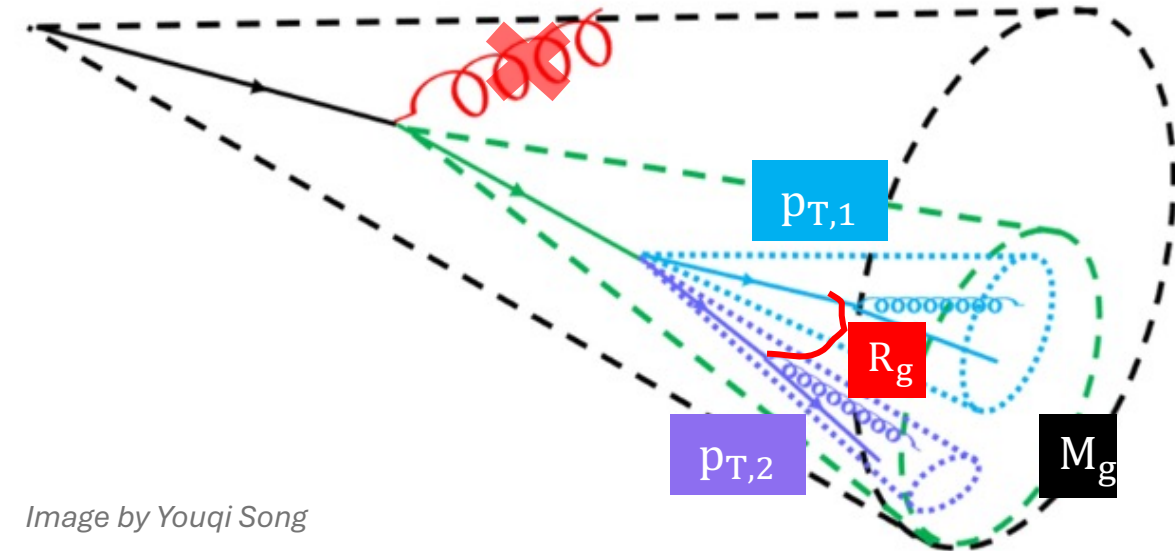


Image by Youqi Song

- **SoftDrop**: Groom a reconstructed jet to remove wide-angle soft radiation
- **CollinearDrop**: Difference of an observable for an ungroomed vs groomed jet  $\rightarrow$  Access to soft component of jet
- **Iterative SoftDrop**: Access to 1<sup>st</sup>, 2<sup>nd</sup>, 3<sup>rd</sup> splits of the shower

$z_g$  = Shared momentum fraction

$R_g$  = Distance of subjets at split

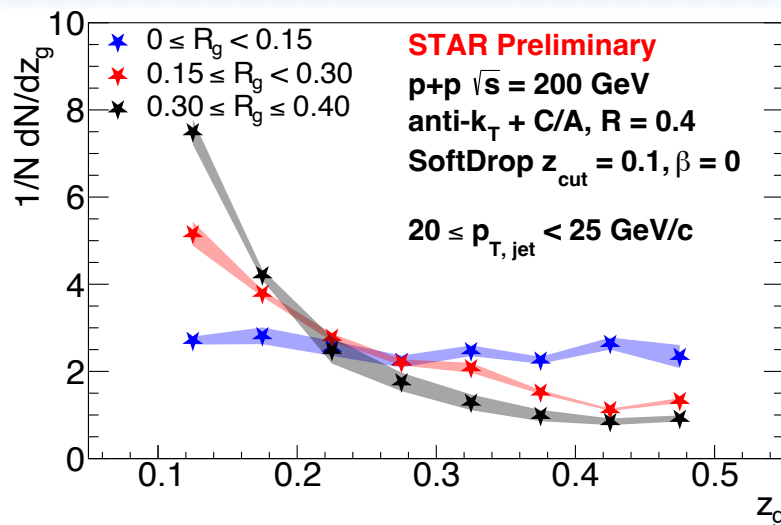
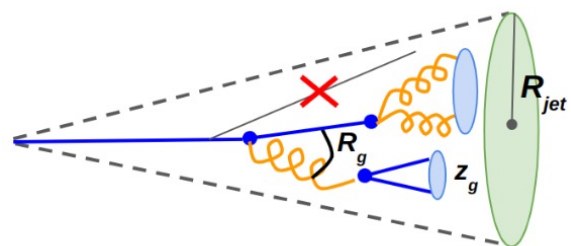
$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{\text{cut}} (R_g / R_{\text{jet}})^\beta \quad \begin{matrix} z_{\text{cut}} = 0.1 \\ \beta = 0 \end{matrix}$$

$$\rightarrow \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > 0.1$$



# Isolating pQCD and npQCD in vacuum

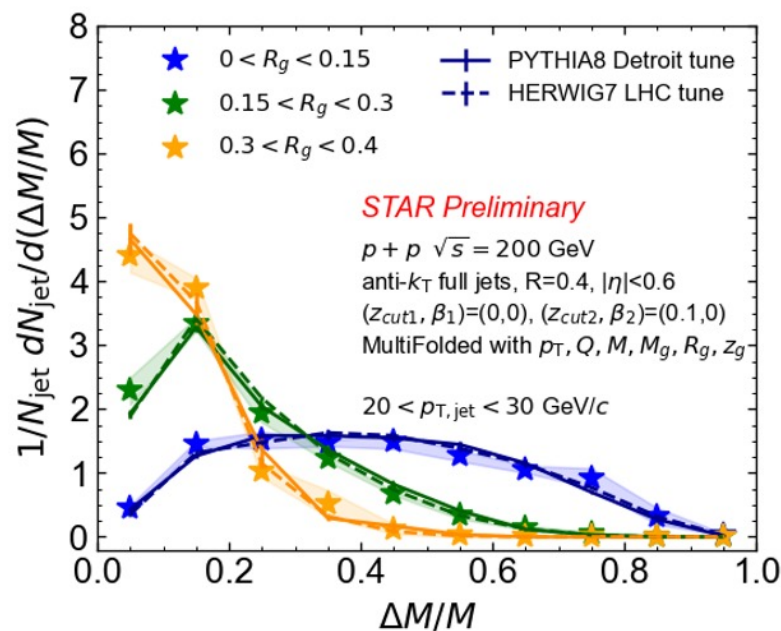
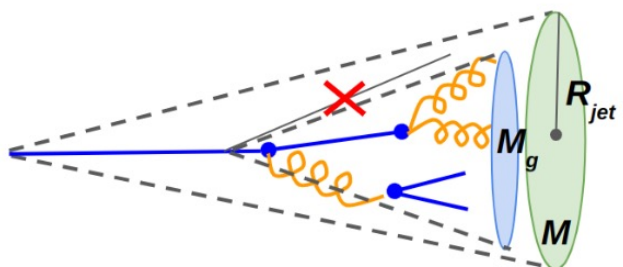
1<sup>st</sup> split



$z_g$  = Shared momentum fraction  
 $R_g$  = Distance of subjects at split  
 $\Delta M = M - M_g$  [GeV]

**Perturbative**

Larger  $R_g \rightarrow$  Smaller  $\langle \Delta M/M \rangle \rightarrow$  Steeper  $z_g$



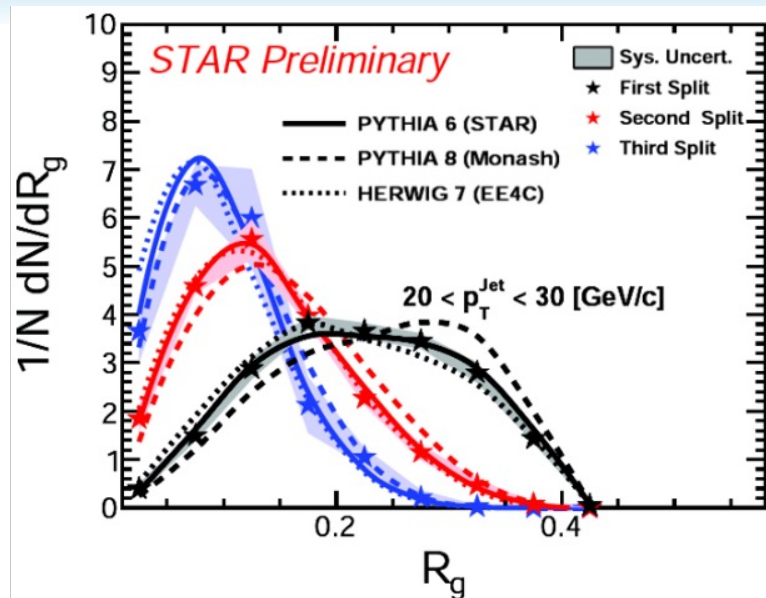
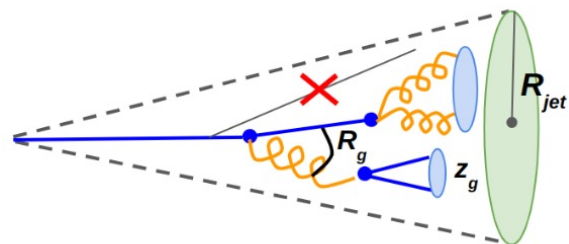
**Non-perturbative**

Smaller  $R_g \rightarrow$  Larger  $\langle \Delta M/M \rangle \rightarrow$  Flatter  $z_g$



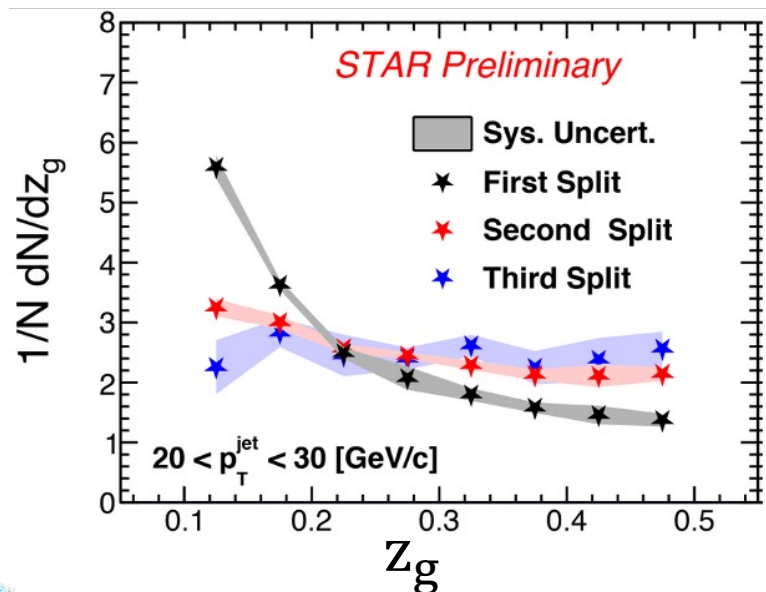
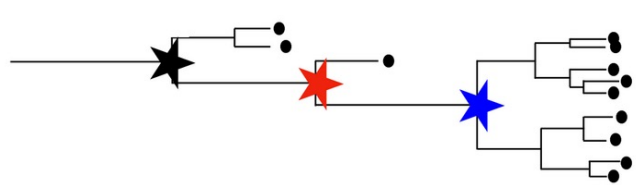
# Isolating pQCD and npQCD in vacuum

1<sup>st</sup>, 2<sup>nd</sup>, 3<sup>rd</sup> splits



$z_g$  = Shared momentum fraction  
 $R_g$  = Distance of subjects at split  
 $\Delta M = M - M_g$  [GeV]

**$R_g$  becomes narrower from 1<sup>st</sup> to 3<sup>rd</sup> split**  
 Change from soft wide-angle to hard collinear  
 splitting



**$z_g$  becomes flatter from 1<sup>st</sup> to 3<sup>rd</sup> split**  
 Perturbative to Non-perturbative transition

*Can we pinpoint a distinct transition region?*

Image by Raghav Kunnawalkam Elayavalli



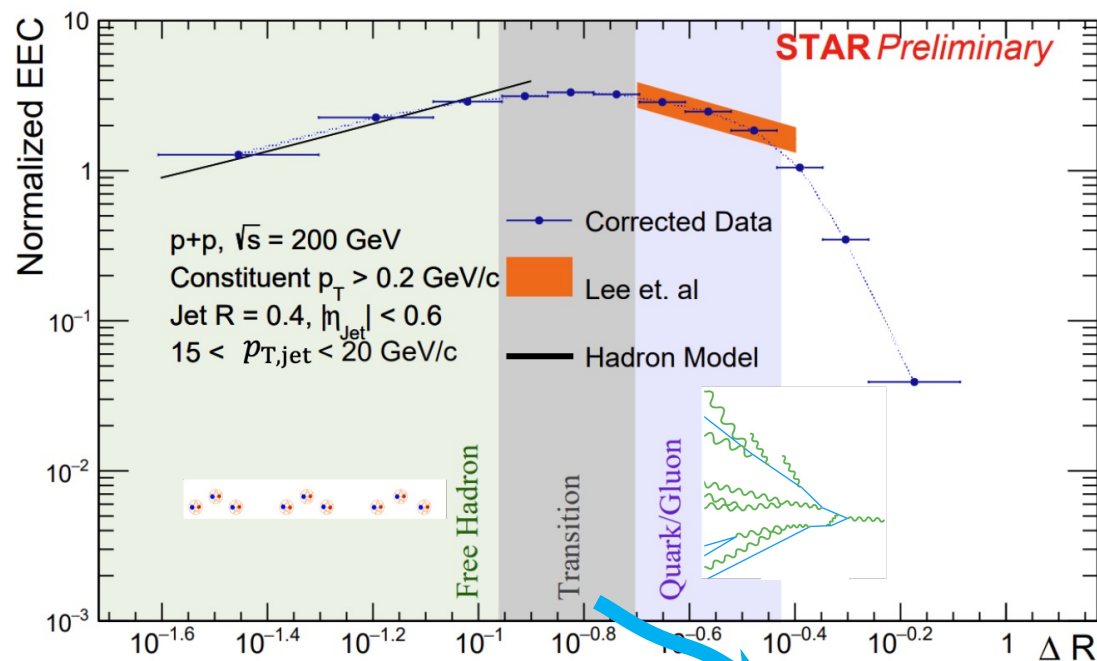
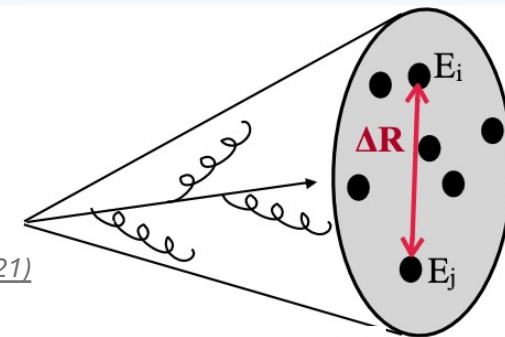
# Time evolution of jets in vacuum

## Energy-Energy Correlators

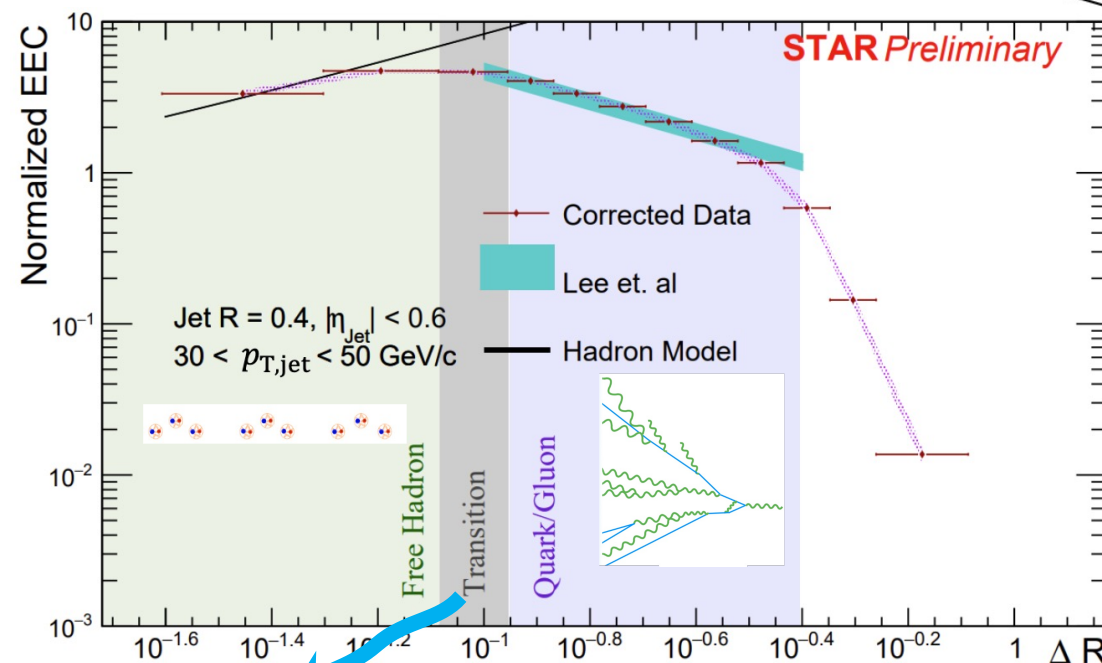
$$\text{Normalized EEC} = \frac{1}{\sum_{\text{Jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T, \text{Jet}}^2}} \frac{d \left( \sum_{\text{Jets}} \sum_{i \neq j} \frac{E_i E_j}{p_{T, \text{Jet}}^2} \right)}{d(\Delta R)}$$

$$\text{Formation Time: } t_f \propto \frac{1}{\Delta R^2}$$

*Apolinário, Cordeiro, & Zapp EPJC 81 (2021)*



*Lee, Mecaj, Mout, arxiv:2205.03414*



Transition Region at  $\Delta R_{\text{Turnover}} \times p_T^{\text{Jet}} \sim 2 - 3$  GeV  $\rightarrow$  No  $p_T^{\text{Jet}}$  dependence

**Universal scale for confinement of quark/gluon to hadrons**



# Probing npQCD region in vacuum

## Charge Correlators

$$r_c(X) = \frac{d\sigma_{h_1 h_2}/dX - d\sigma_{h_1 \bar{h}_2}/dX}{d\sigma_{h_1 h_2}/dX + d\sigma_{h_1 \bar{h}_2}/dX}$$

$h_1 h_2$ : same charge hadrons,  
 $h_1 \bar{h}_2$ : opposite charge hadrons

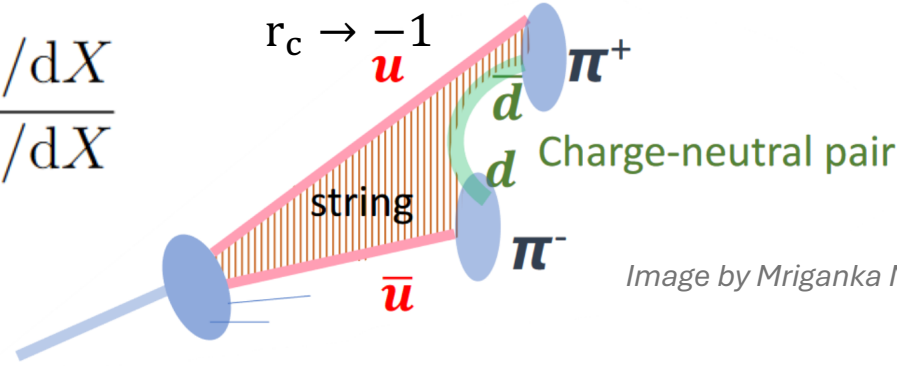


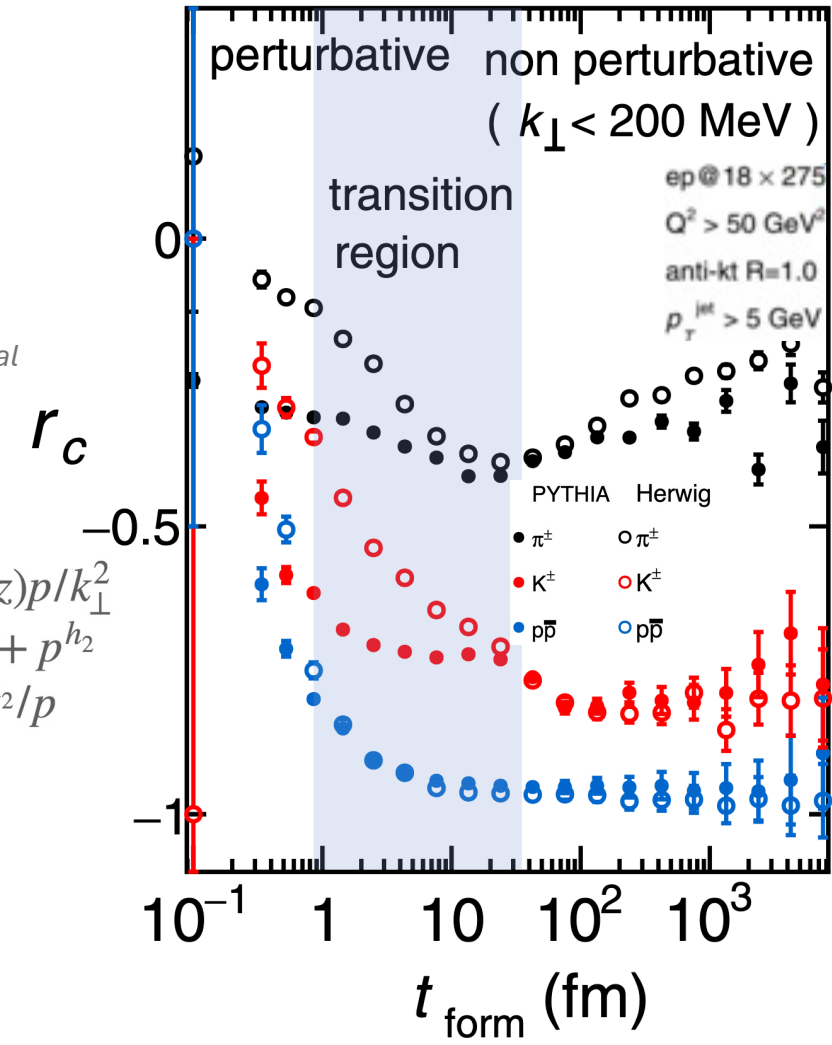
Image by Mriganka Mondal

Chien, Deshpande, Mondal, Sterman, *PRD* 105 (2022) 5, L051502

$$t_{\text{form}} = z(1-z)p/k_{\perp}^2$$

with  $p = p^{h_1} + p^{h_2}$   
 and  $z = p^{h_2}/p$

- $r_c$  can probe for evidence of string-like fragmentation
- In vacuum, can establish baseline for studying medium modification of hadronization



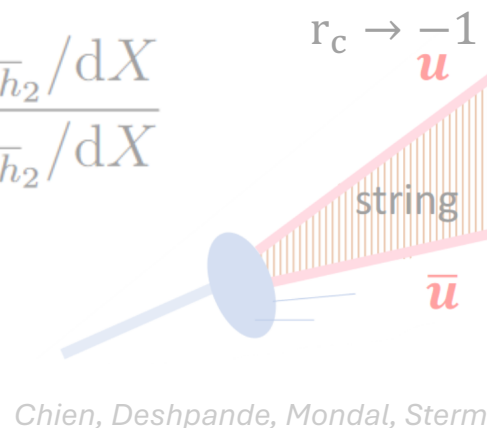


# Probing npQCD region in vacuum

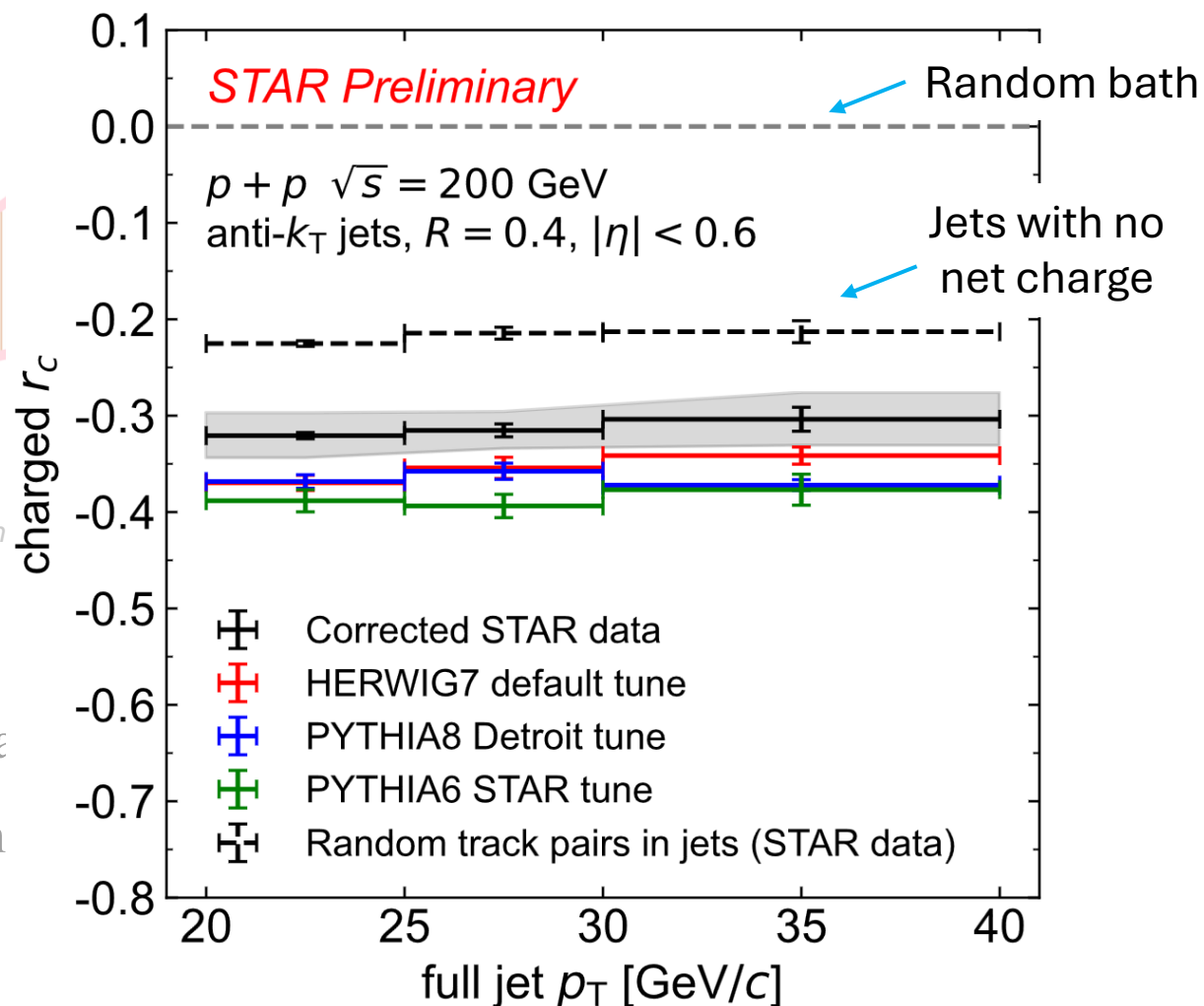
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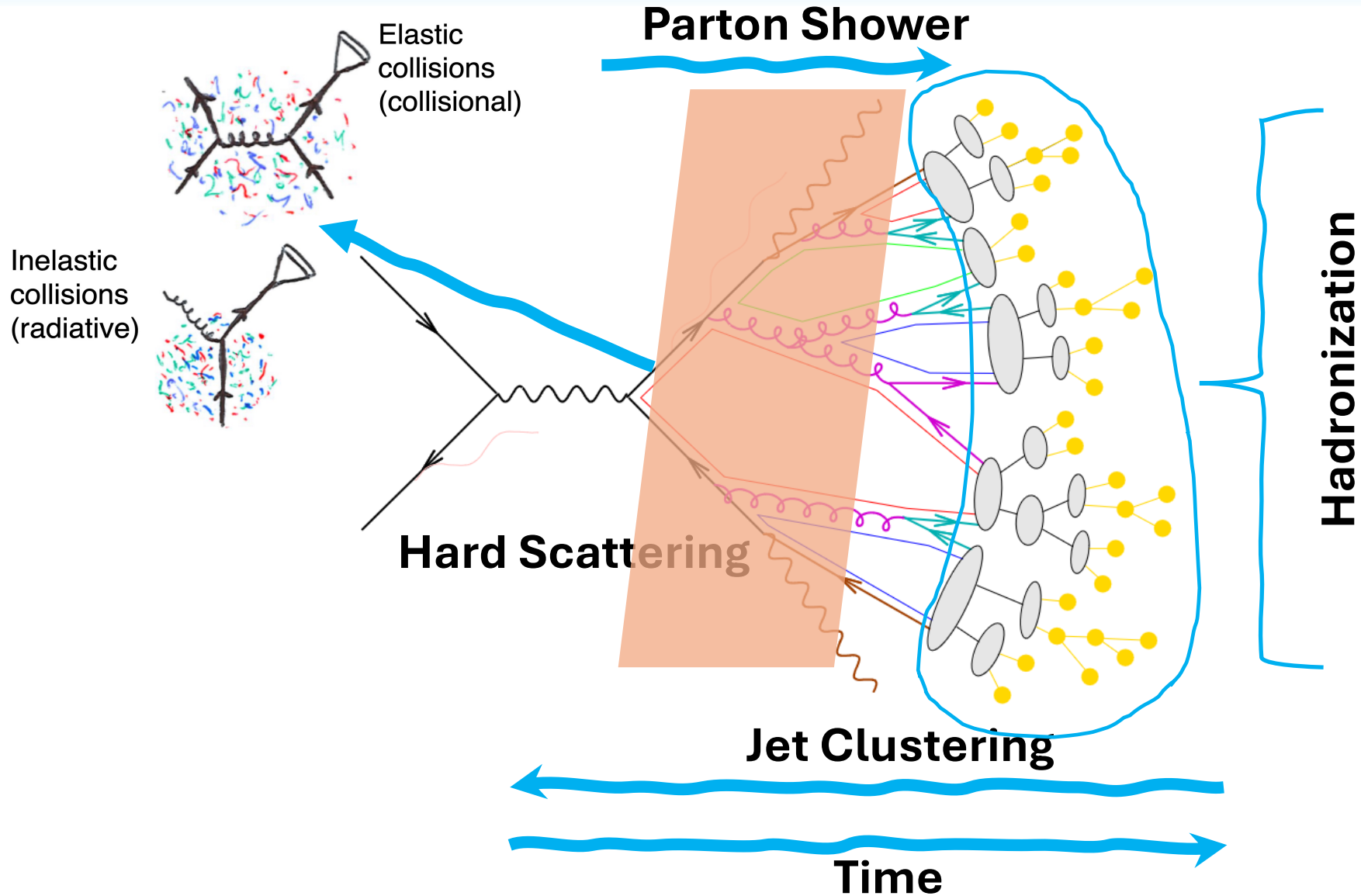
- $r_c$  can probe for evidence of string-like fragmentation
- In vacuum, can establish baseline for studying modification of hadronization



**First measurement in p+p: Both string-like and cluster hadronization underpredict STAR data**  
 More model tuning required.



# Jets in medium



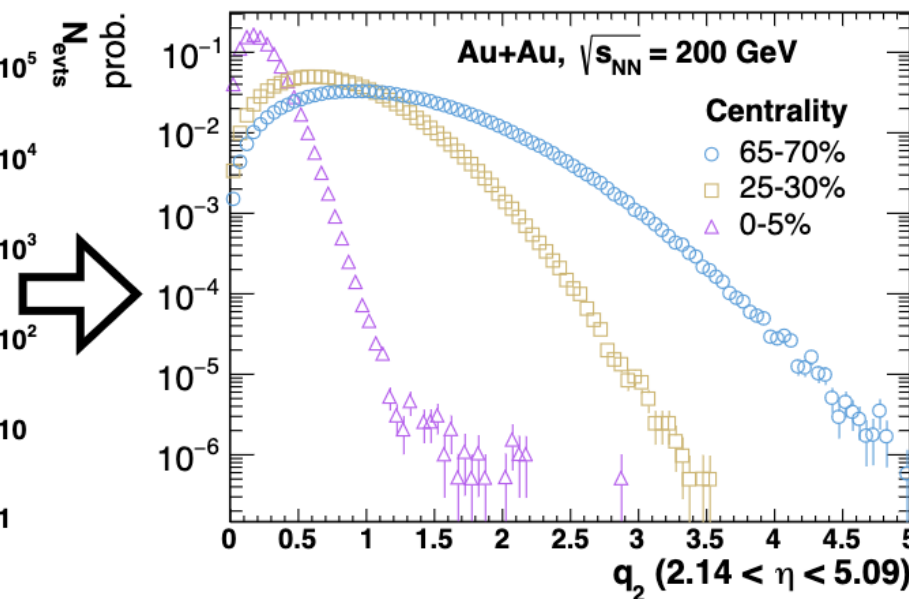
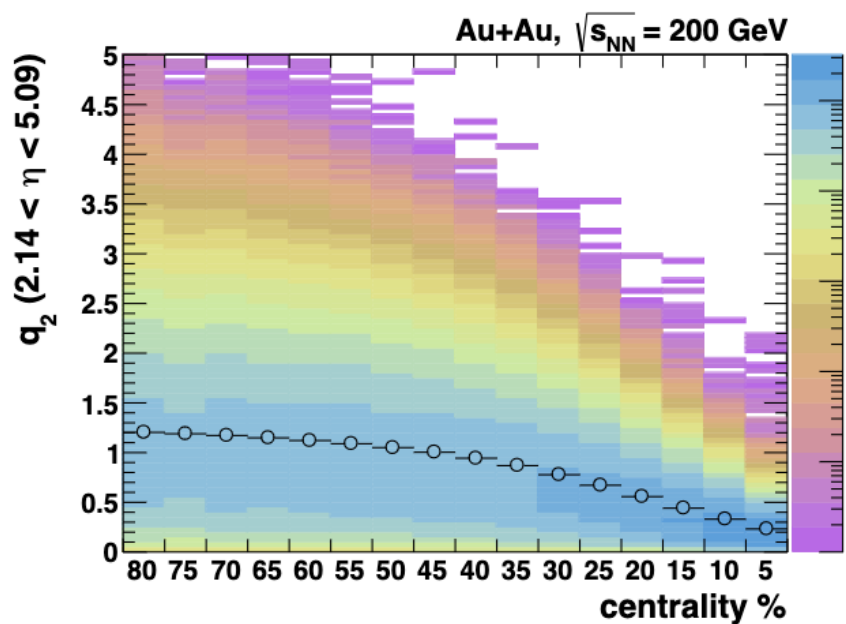
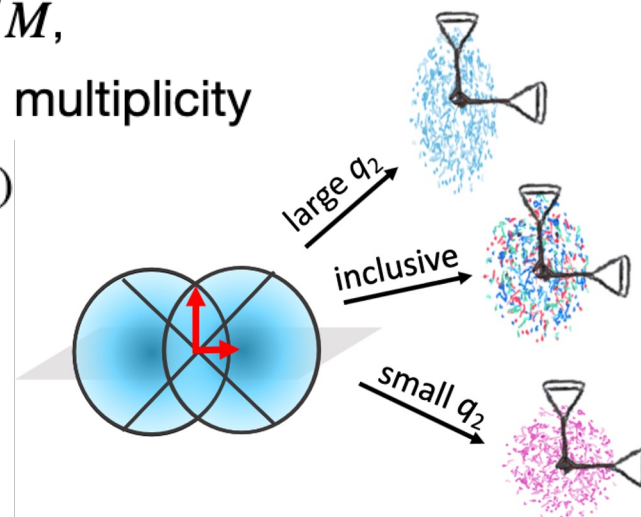
# Event shape engineering

$$Q_2 = \left( \sum_{i=1}^M w_i \cos(2\phi_i), \sum_{i=1}^M w_i \sin(2\phi_i) \right), \quad q_2 = |Q_2|/\sqrt{M},$$

$w_i$ : nMIP weight,  $M$ : multiplicity

$$v_2 = \langle \cos(2(\phi - \Psi_2)) \rangle$$

Images by Caitie Beattie



- Centrality and  $q_2$  are correlated, event selection based on both centrality and  $q_2$
- Charged particle spectra from TPC,  $q_2$  from EPD-W to avoid autocorrelation



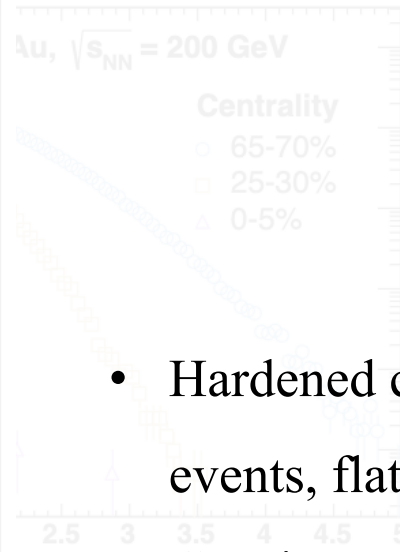
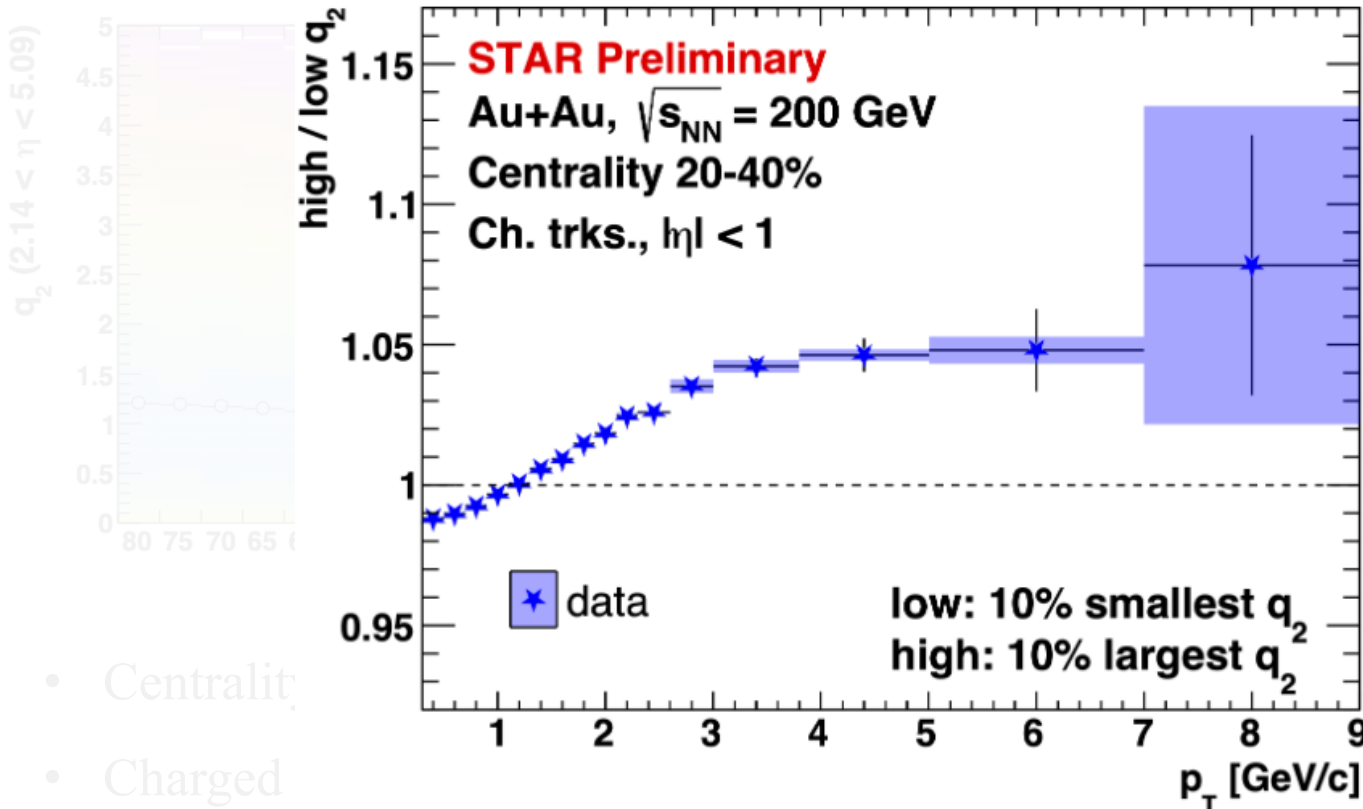
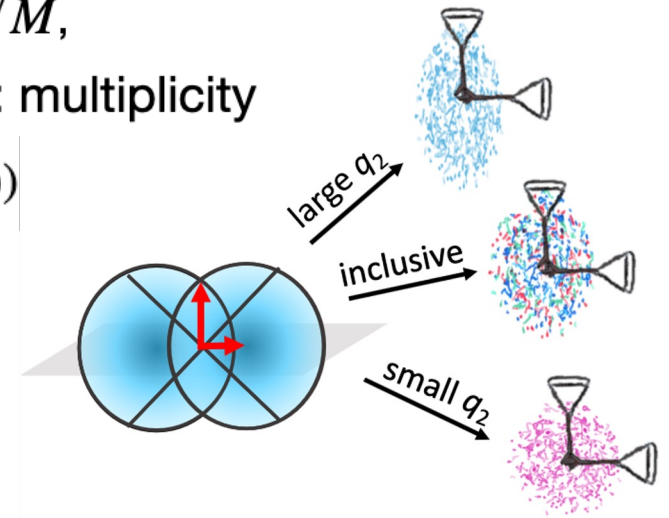
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Images by Caitie Beattie



- Hardened charged particle spectra in high- $q_2$  events, flattened at high  $p_T$
- Consistent with ALICE results at 2.76 TeV

*ALICE, PRC 93 (2016) 3, 034916*

- Centrality
- Charged

th centrality and  $q_2$   
d autocorrelation



# Jet substructure in medium

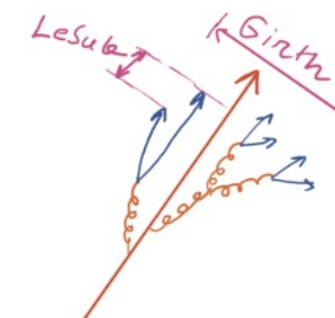
## Generalized Angularities

$$\lambda_{\beta}^{\kappa} = \sum_{\text{const} \in \text{jet}} \overbrace{\left( \frac{p_{T,\text{const}}}{p_{T,\text{jet}}} \right)^{\kappa}}^{\text{soft/hard radiation}} \times \overbrace{r(\text{const}, \text{jet})^{\beta}}^{\text{collinearity sensitive}}$$

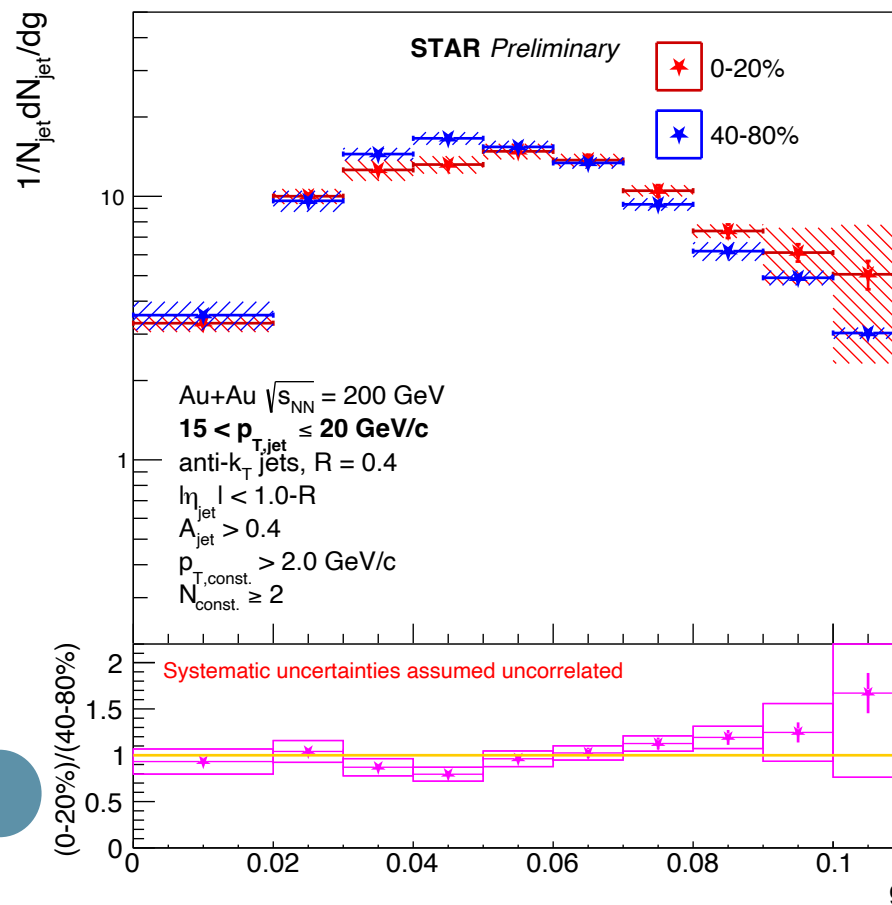
$$\lambda_1^1 = g - \text{girth}$$

$$\lambda_2^1 - \text{thrust}$$

$$\lambda_0^2 = (p_T^D)^2 - \text{momentum dispersion}$$



- Angularities, tunable sensitivities to energy, angular scales – some of them are IRC safe
- Can probe the modification of radiation pattern in medium



**Girth consistent in central and peripheral collisions**

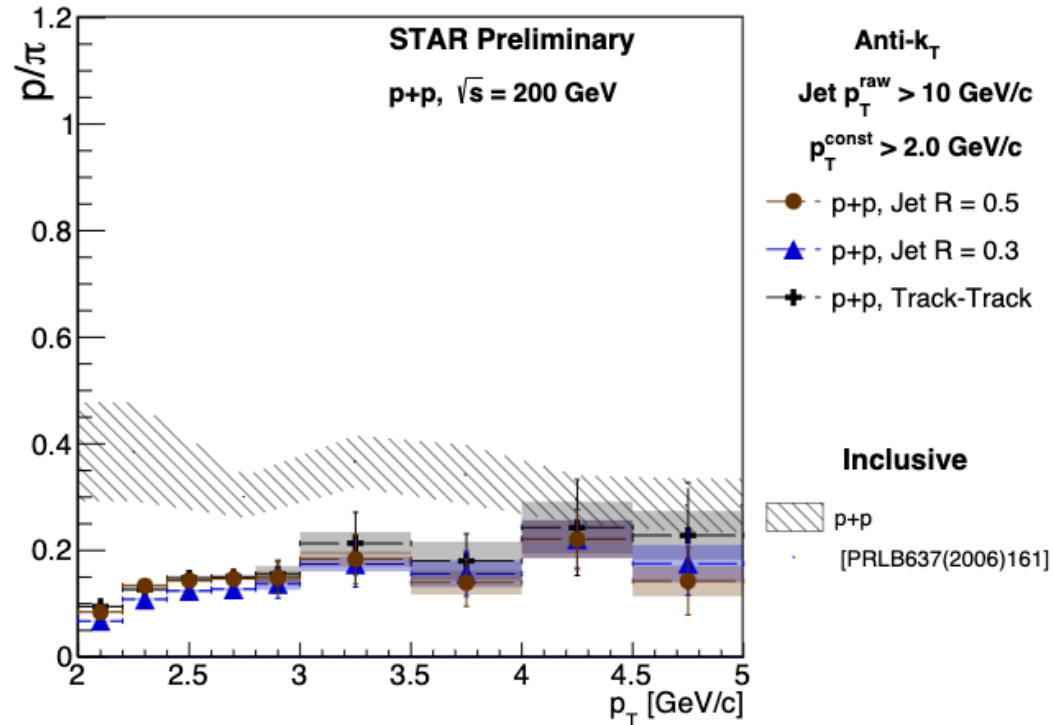
Better handle on systematics and sample bias being explored



# Jet chemistry

## Baryon-To-Meson Ratio in p+p and Au+Au

$$\frac{p^+ + p^-}{\pi^+ + \pi^-}$$



Pion production preferred over proton in jets

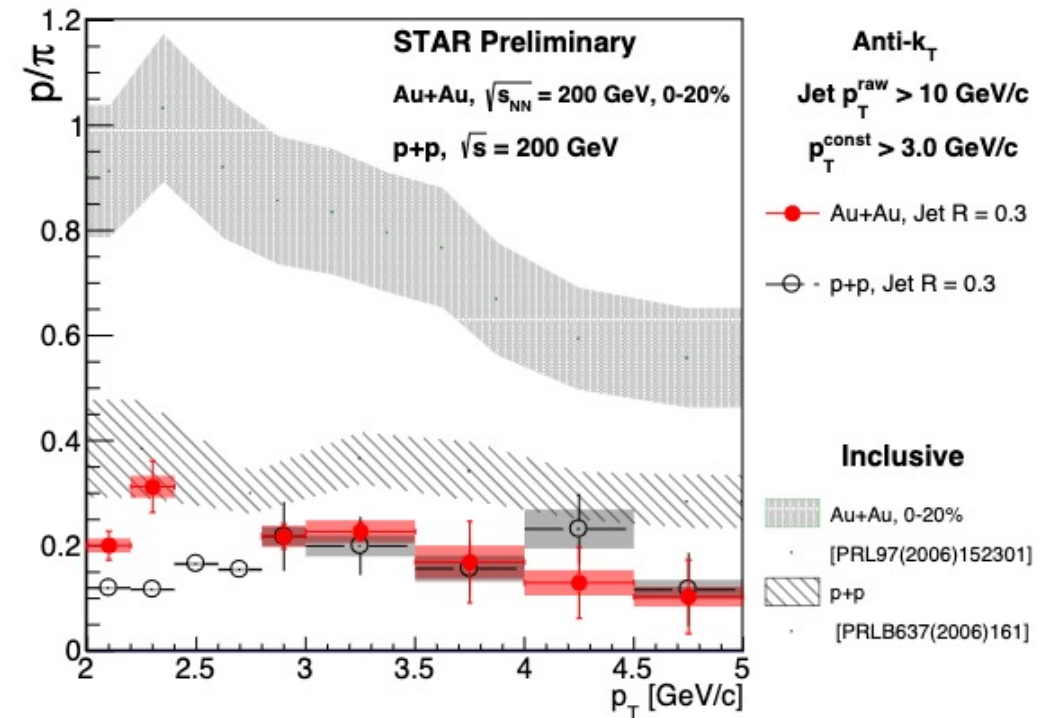
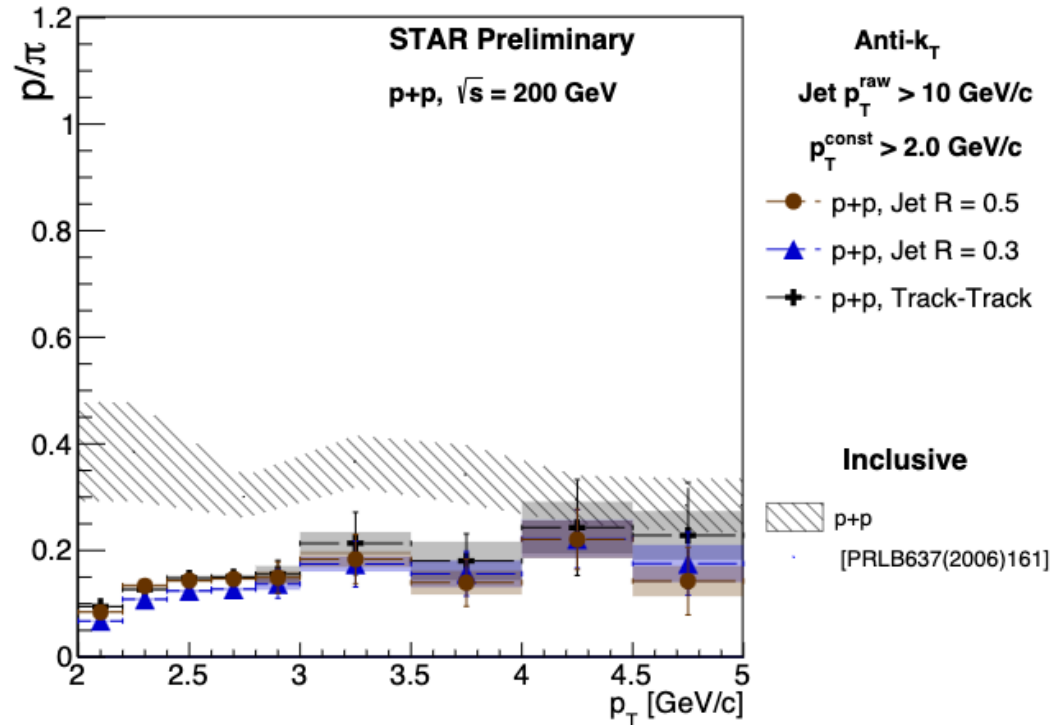
Stronger preference for pions in jets compared to inclusive p+p



# Jet chemistry

## Baryon-To-Meson Ratio in p+p and Au+Au

$$\frac{p^+ + p^-}{\pi^+ + \pi^-}$$



Pion production preferred over proton in jets

**No significant difference in Au+Au p/π ratio compared to p+p**

Stronger preference for pions in jets compared to inclusive p+p

Hard-core selection bias (?) Survivor bias(?)

**Studies ongoing with jets with different hard-core definitions**



# Flavor dependence – D<sup>0</sup> jets

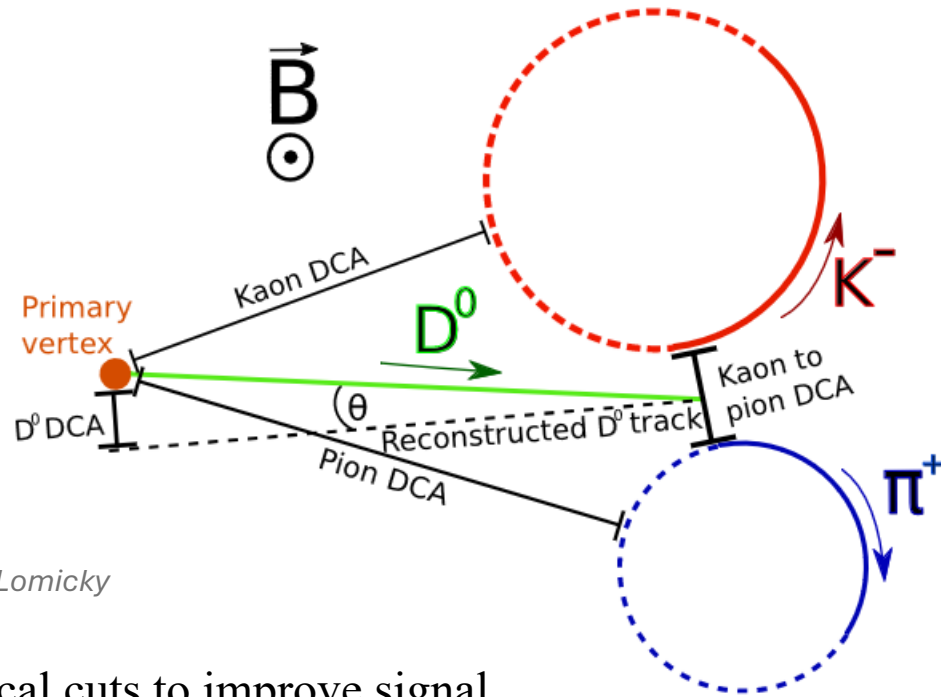


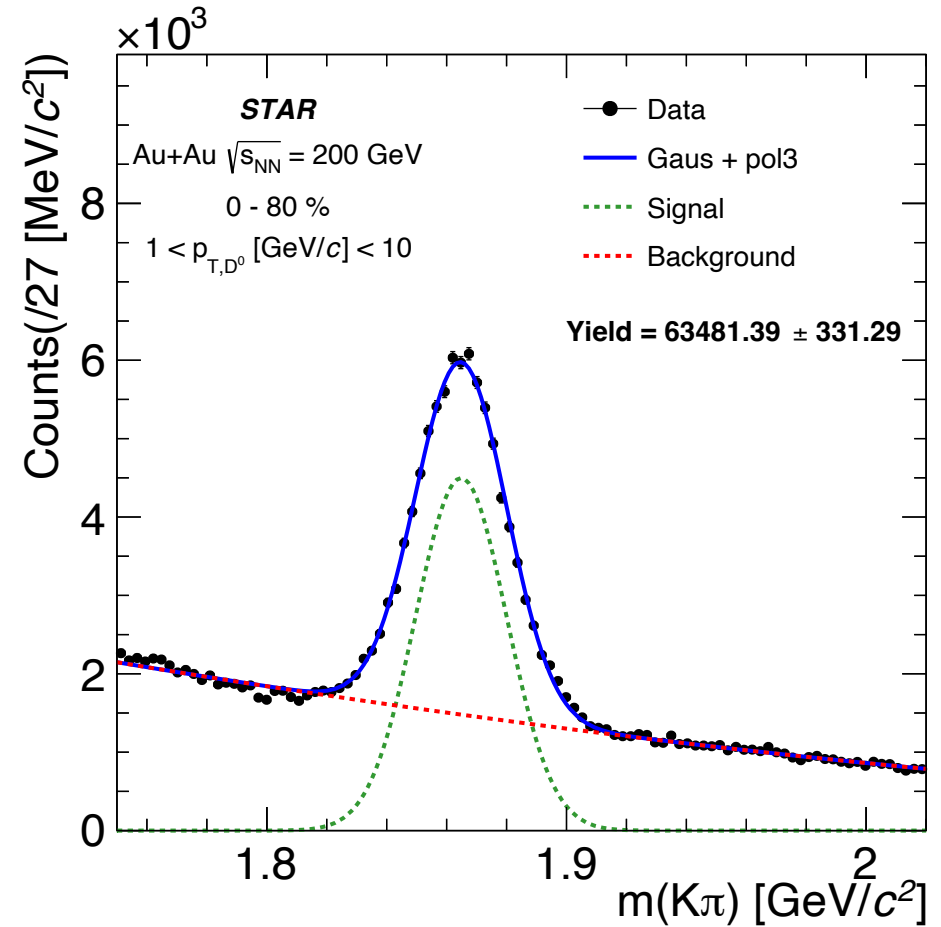
Image by Ondrej Lomicky

- Topological cuts to improve signal significance of D<sup>0</sup>
- Yield calculation using sPlot method

$$sPlot \quad s\mathcal{P}_n(m_{K\pi,i}) = \frac{\sum_{j=1}^{N_T} V_{nj} f_j(m_{K\pi,i})}{\sum_{k=1}^{N_T} N_k f_k(m_{K\pi,i})}$$

Efficiency Correction

$$s\mathcal{P}_n(m_{K\pi,i}) \rightarrow \frac{s\mathcal{P}_n(m_{K\pi,i})}{\varepsilon(m_{K\pi,i})}$$



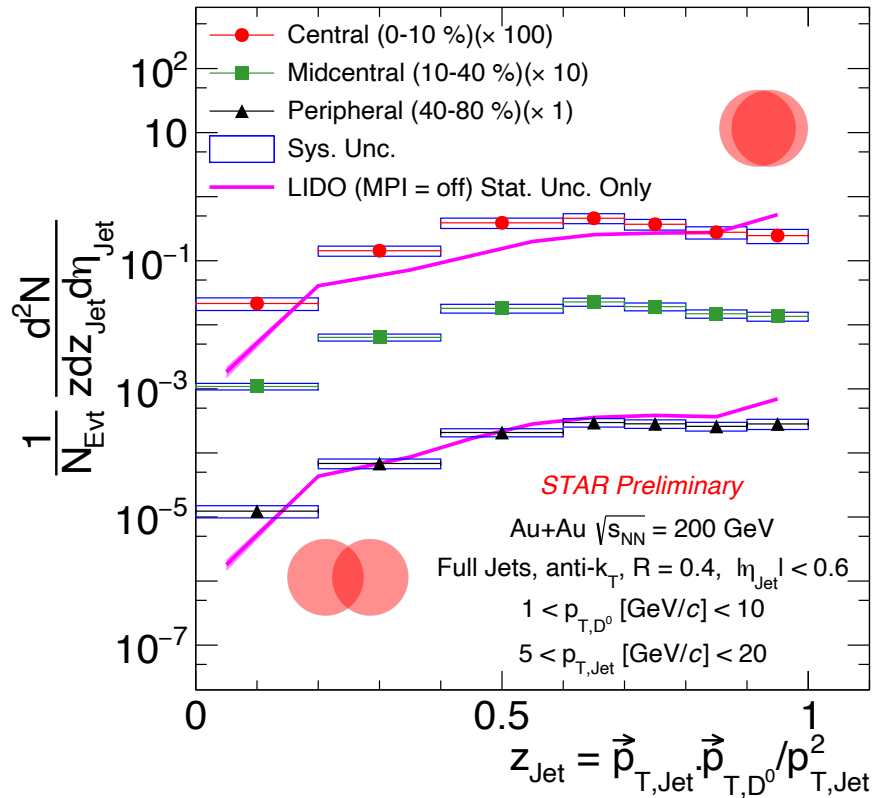
Pivk, Diberder, Nucl. Instrum. Methods Phys. Res., A (2005) 555





# Flavor dependence – Fragmentation of $D^0$ jets

LIDO, Phys. Rev. C 98, 064901



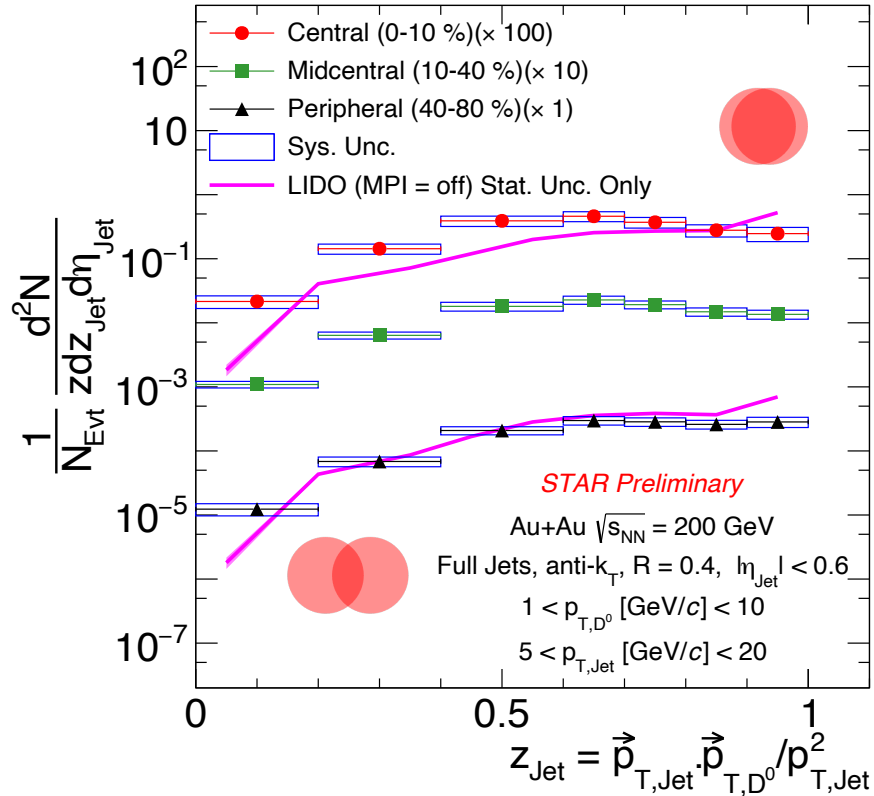
$$z_{\text{Jet}} = \frac{\vec{p}_{T,\text{Jet}} \cdot \vec{p}_{T,D^0}}{|\vec{p}_{T,\text{Jet}}|^2}$$

- 2D unfolded with  $p_{T,\text{Jet}}$
- LIDO overestimates hard fragmented  $D^0$  jets  $\rightarrow$  Data shows softer fragmentation

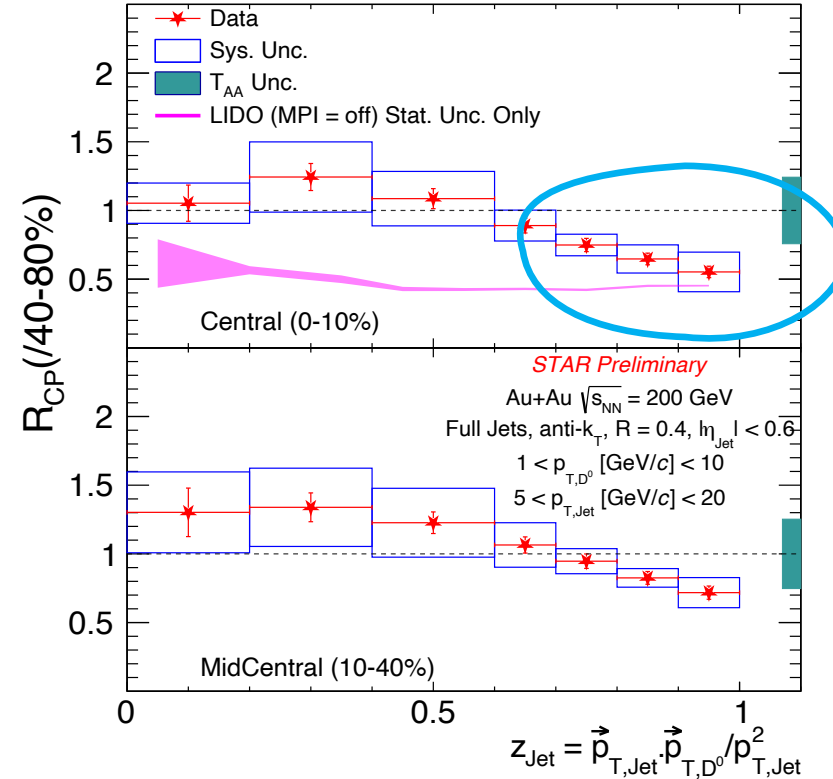
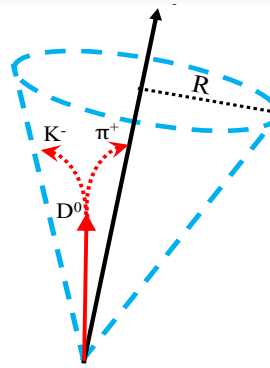


# Flavor dependence – Fragmentation of $D^0$ jets

LIDO, Phys. Rev. C 98, 064901



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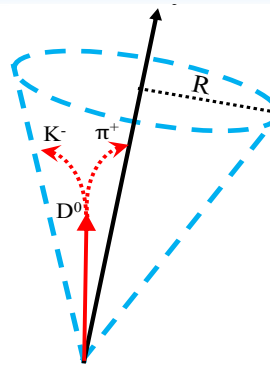


## Suppression for hard fragmented $D^0$ jets in central collisions

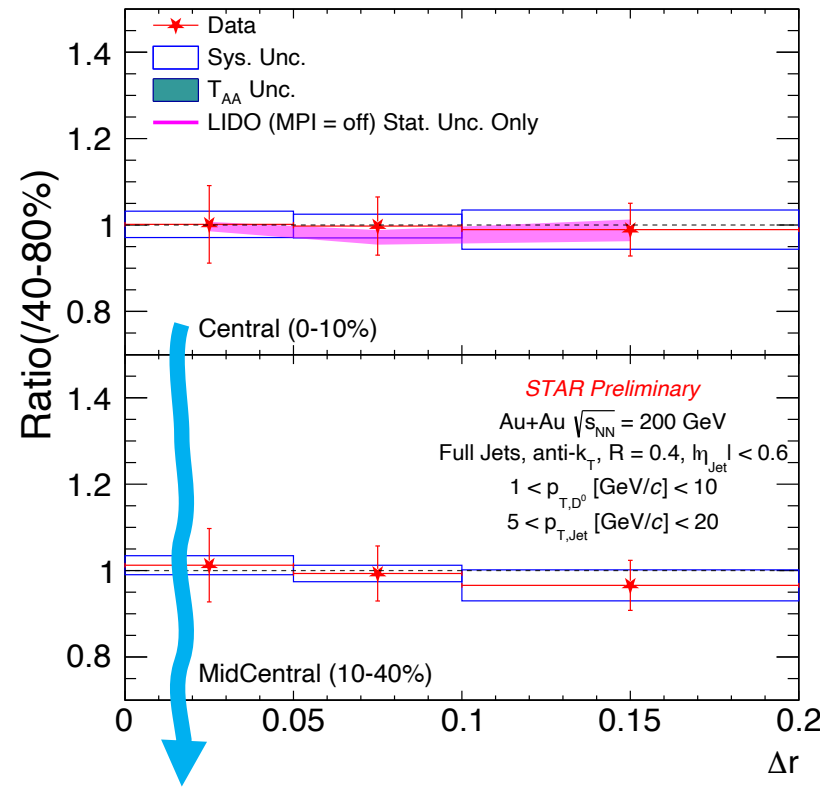
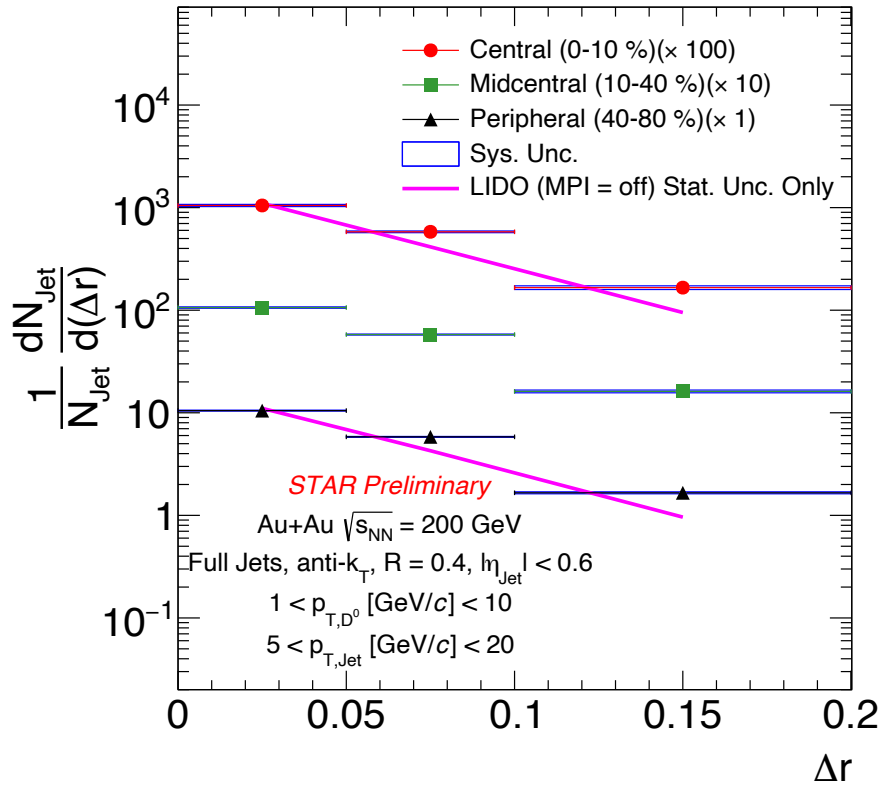
- 2D unfolded with  $p_{T,\text{Jet}}$
- LIDO overestimates hard fragmented  $D^0$  jets  $\rightarrow$  Data shows softer fragmentation



# Flavor dependence – Radial profile of $D^0$ jets



LIDO, Phys. Rev. C 98, 064901



- 2D unfolded with  $p_{T,\text{Jet}}$
- LIDO qualitatively explains radial profile trends, along with ratio of central and peripheral

**Ratio of radial profile consistent with 1**



# What's next for STAR?

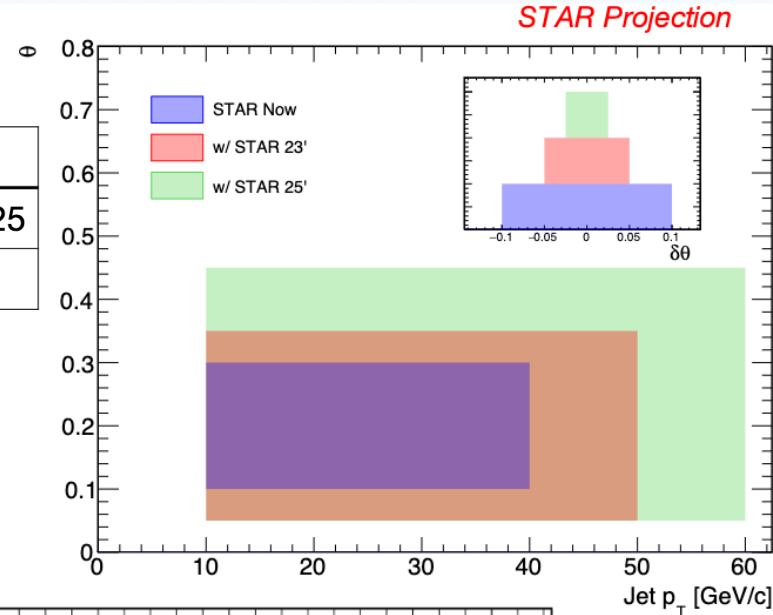
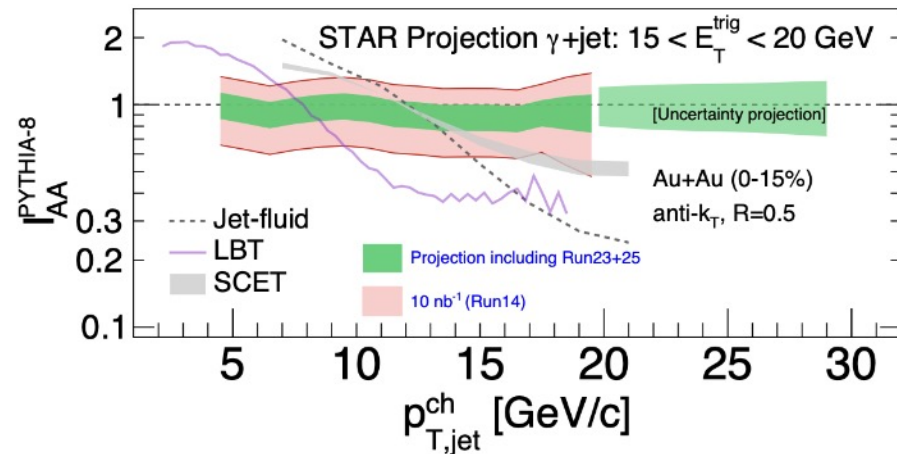
$\sqrt{s_{NN}}$ (GeV)	Species	Sampled Luminosity	Year
200	Au+Au&p+Au	AuAu 32.7 nb <sup>-1</sup> / pAu 0.69 pb <sup>-1</sup>	2023+2025
200	p+p	142 pb <sup>-1</sup>	2024

EPD for triggering  
Independent event-plane determination

## Event shape Engineering

Recoil Jets triggered by  $\gamma^{\text{dir}}$

$$I_{AA} = \text{Yield in Au+Au} / \text{Yield in p+p}$$



Increased statistics for jet-substructure measurements

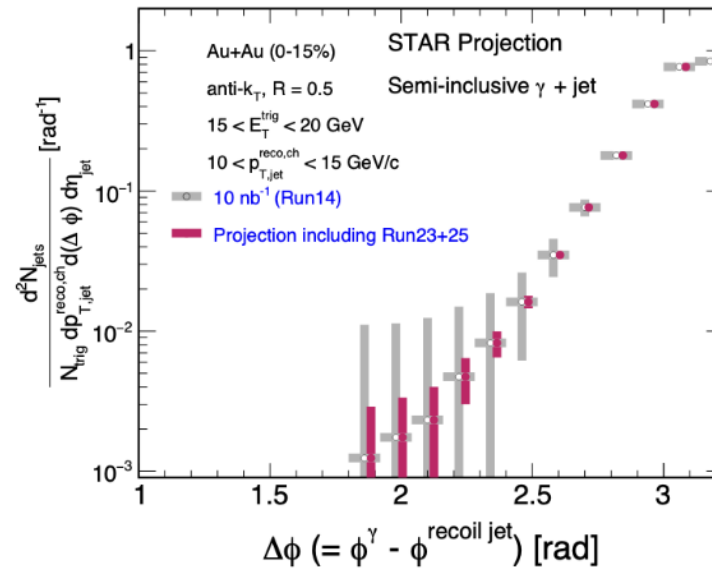
Use jet substructure as taggers

Access to high  $p_T$  jets

Access to wide-angle radiation

Increased resolution in angular

scale



Wider kinematic range for  $I_{AA}$  and acoplanarity measurements

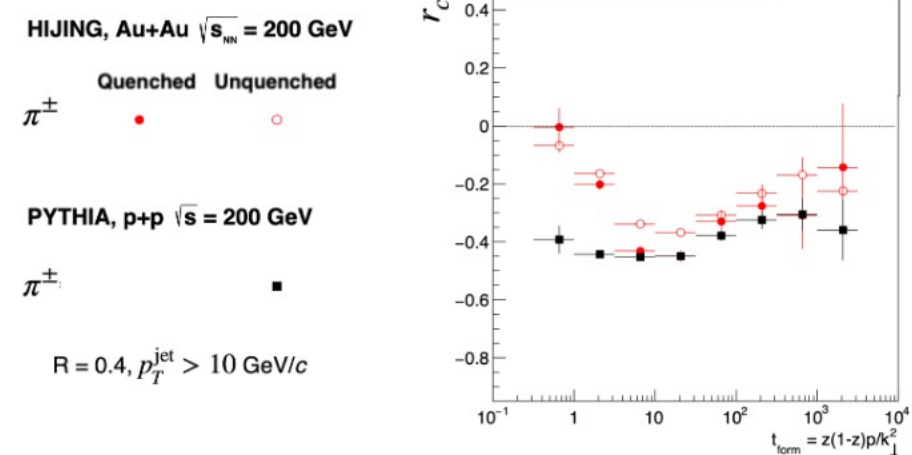
Access to forward rapidity

Larger statistics  $\rightarrow$  Improved uncertainty

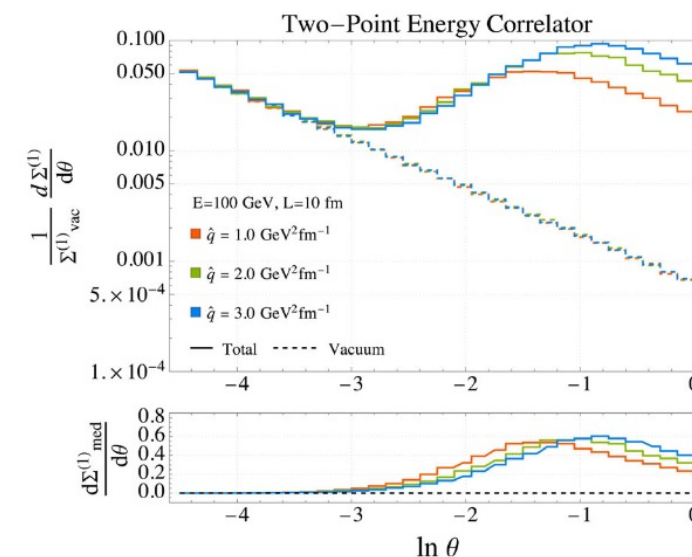


# What's next for STAR?

- Charge correlators in heavy ion collisions
- HERWIG tune to RHIC kinematics ongoing
- Jet chemistry in unbiased sample (constituent  $p_T$  dependence)
- Generalized angularities for  $D^0$ -jets
- Higher order EECs, charge dependence, medium modifications
- Event shape engineering: Probing event-plane angle dependence



Esha, *Hard Probes 2023*

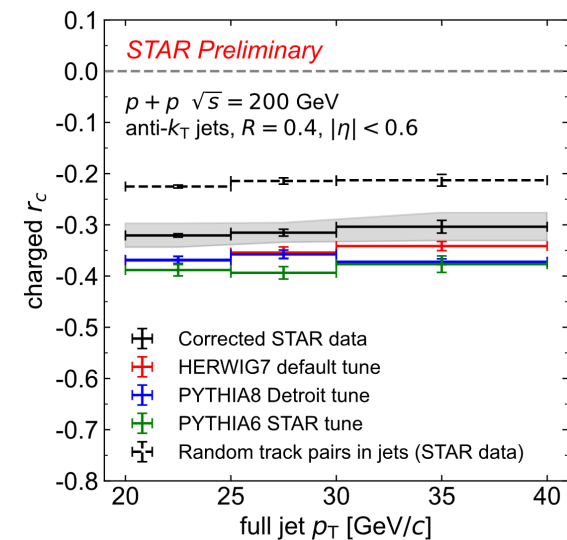
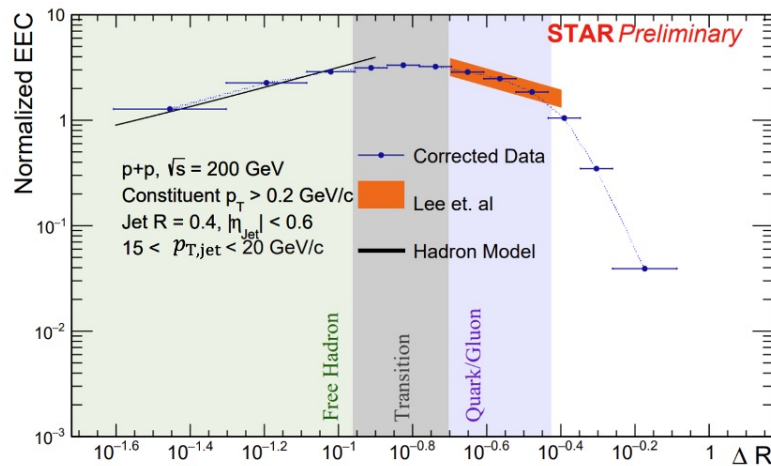
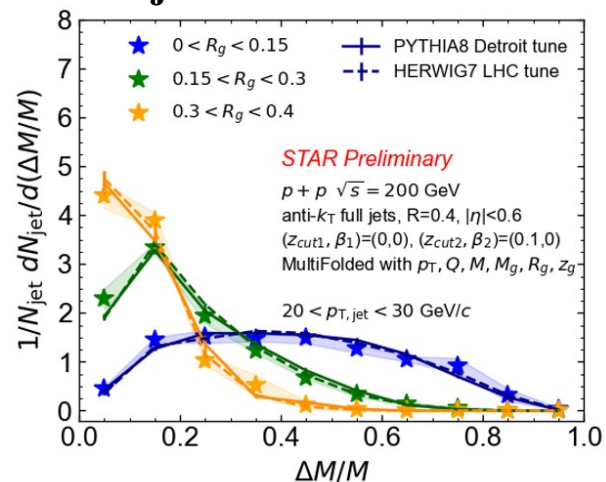


Andres, Dominguez, Kunnawalkam Elayavalli, Holguin, Marquet, Moutl, *PRL 130 (2023) 26, 262301*

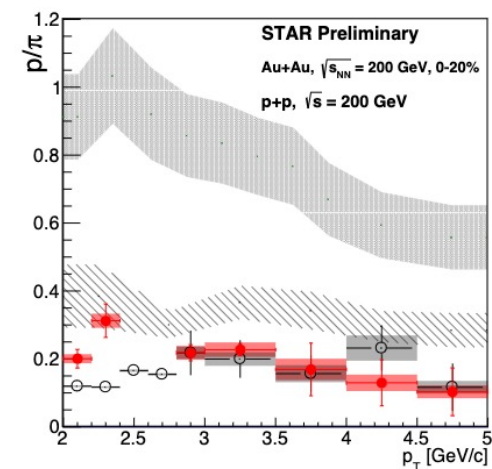
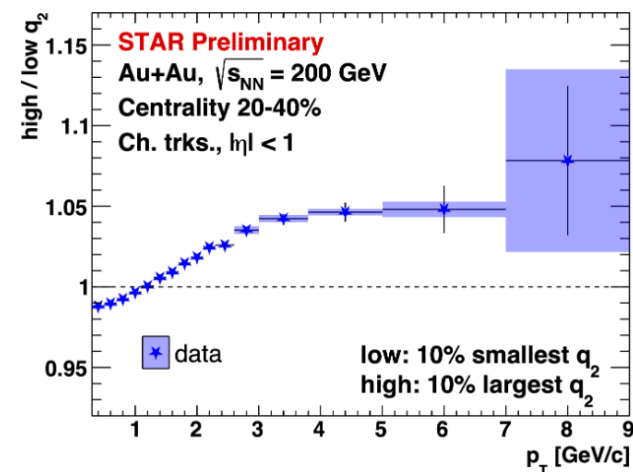
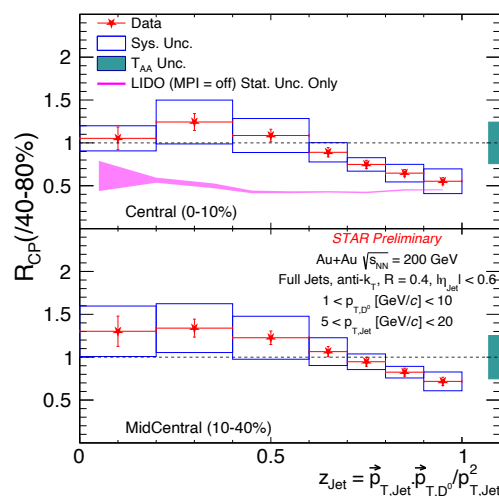
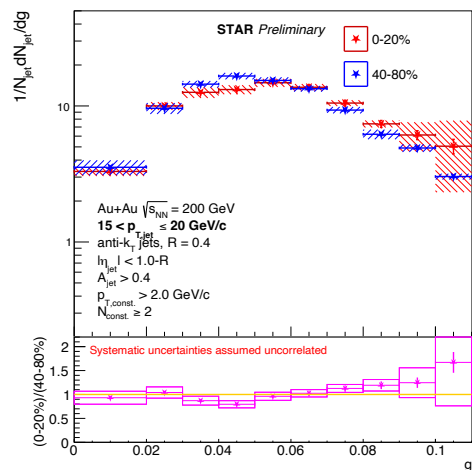


# Summary

## Precision era of jet substructure



## Substructure measurements in medium



*Thank You*



# Flavor dependence – Generalized angularities

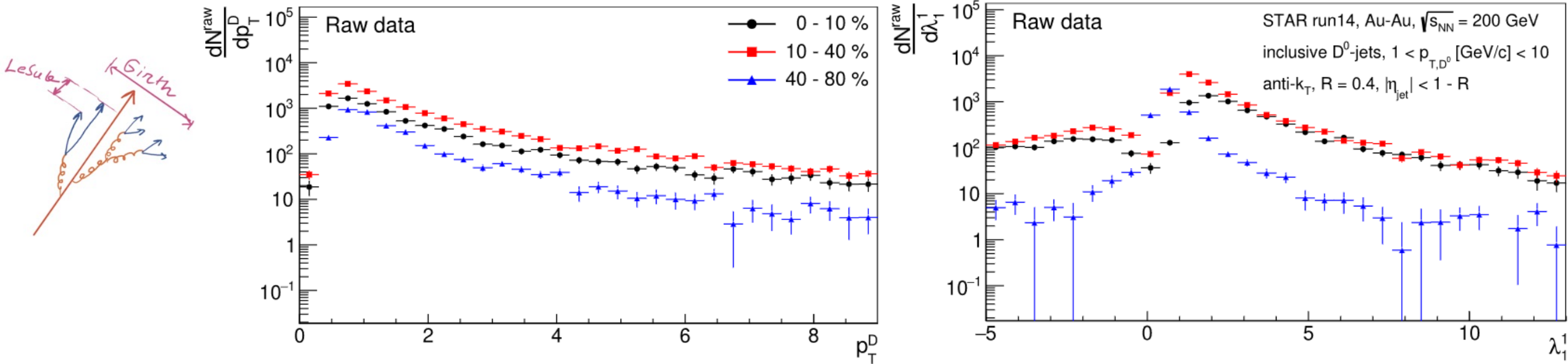
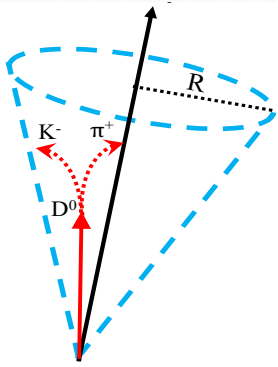
## Generalized Angularities

$\lambda_1^1 = g$  - girth

$\lambda_2^1$  - thrust

$\lambda_0^2 = (p_T^D)^2$  - momentum dispersion

$$\lambda_\beta^\kappa = \sum_{\text{const} \in \text{jet}} \overbrace{\left( \frac{p_{T,\text{const}}}{p_{T,\text{jet}}} \right)^\kappa}^{\text{soft/hard radiation}} \times \overbrace{r(\text{const}, \text{jet})^\beta}^{\text{collinearity sensitive}}$$



## Ongoing measurements for angularities

- Unphysical results caused by median background subtraction → Unfolding required