

# BOOST 2024: Genova, Italy

## Interpretation of the MC top mass parameter with the soft-dropped groomed jet mass

**Naseem Bouchhar**  
**31 July 2024**

Collaborators:  
Andre Hoang  
Matt LeBlanc

Johannes Michel  
Aditya Pathak

Iain Stewart  
Miguel Villaplana  
Marcel Vos

# TOP QUARK MASS INTERPRETATION

- The physical mass of the top quark is that found in the Lagrangian. However, MC top quark mass has much lower uncertainties.

- **Want to interpret direct mass measurements within a field-theoretical renormalisation scheme.**

- ▶ **Test the relation between the two mass parameters:**

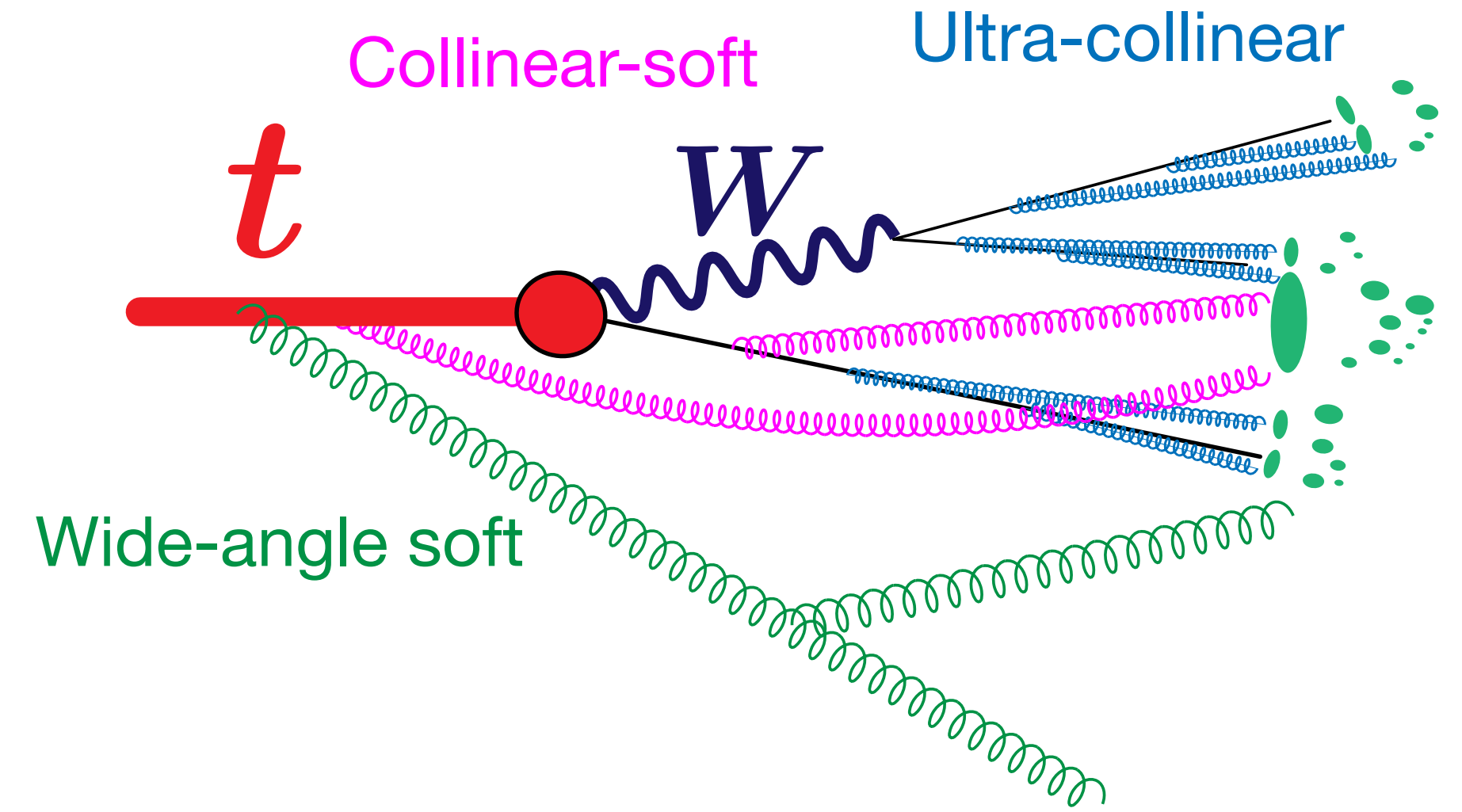
$$m_t^{MC} = m_t^{theo} + \Delta_{t,MC}$$

- The ambiguity can be reduced through dedicated 'calibration' studies.
- ➔ Is  $m_t^{MC}$  only effective in matching experimental data, using mass parameters that don't directly correspond to fundamental QCD parameters?
- ➔ Or, if closely aligned with QCD, can  $m_t^{MC}$  represent the physical mass in a given scheme of the top quark as in the QCD Lagrangian.

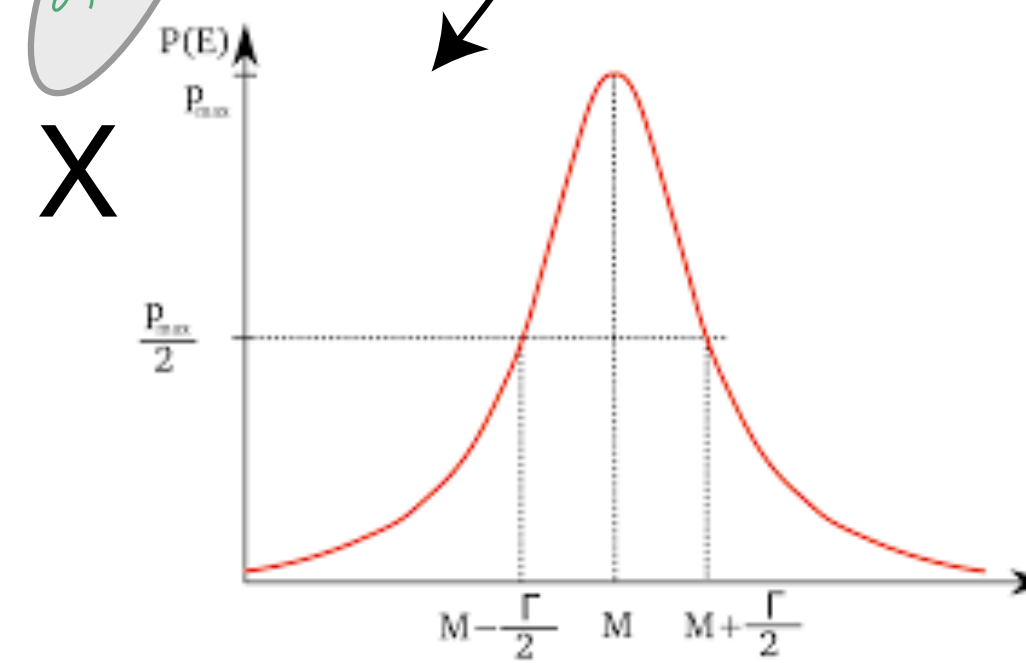
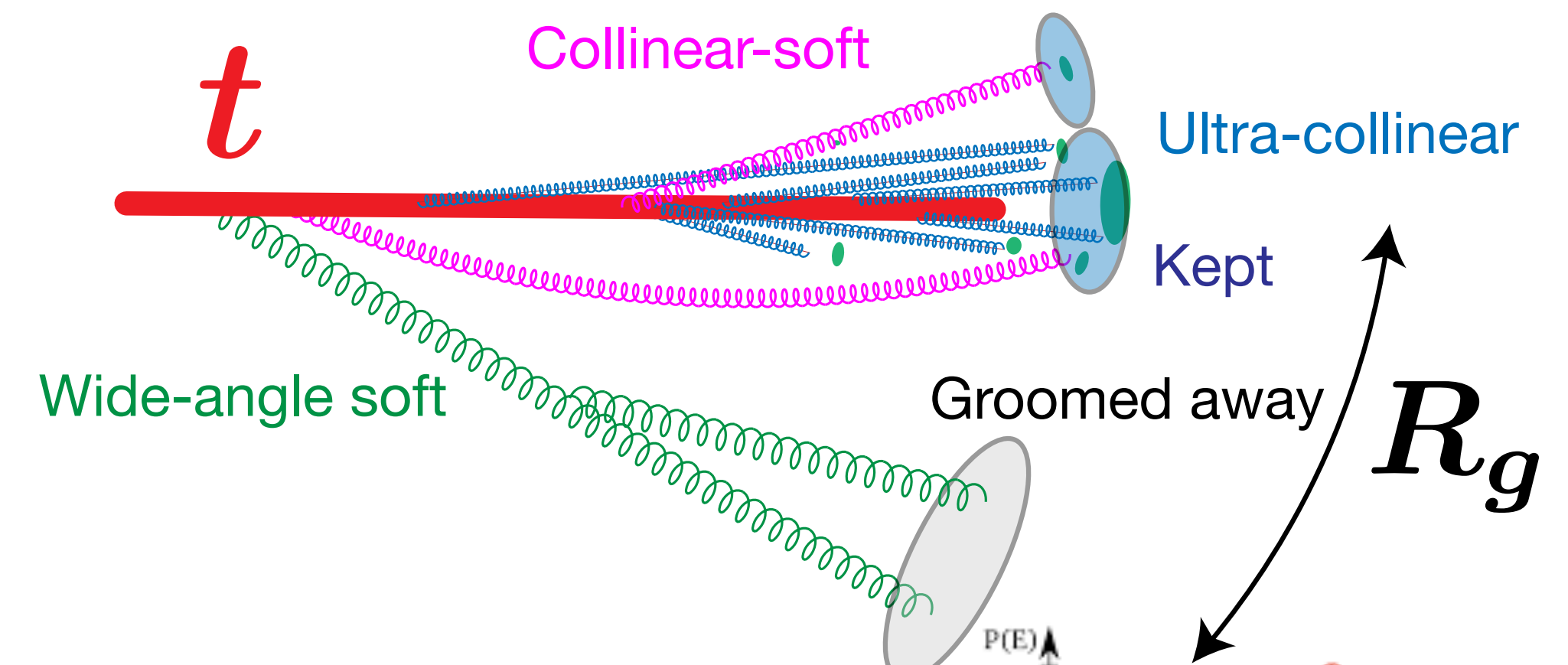
# INTRODUCTION

The top quark presents unique challenges with jet substructure with its intricate three-body electroweak decay.

- For observables that do not resolve the details of the top decay (inclusive over the top quark decay) the jet substructure can be described by considering a stable top quark that radiates and including the decay effects via Breit Wigner distribution. [Phys.Rev.D 77 \(2008\) 074010](#)
- Currently no **theoretical framework describing the decay of the top in an exclusive way** i.e. when the decay products are resolved.
- Top quark mass studies with soft-dropped grooming typically focuses on **light grooming region**.
  - Enables a relatively simple inclusive description of the top quark decay.
- In this talk we test how far we can push the limit of inclusive description
- **The first step towards trying to build a more robust groomed observable for  $pp$  collisions.**



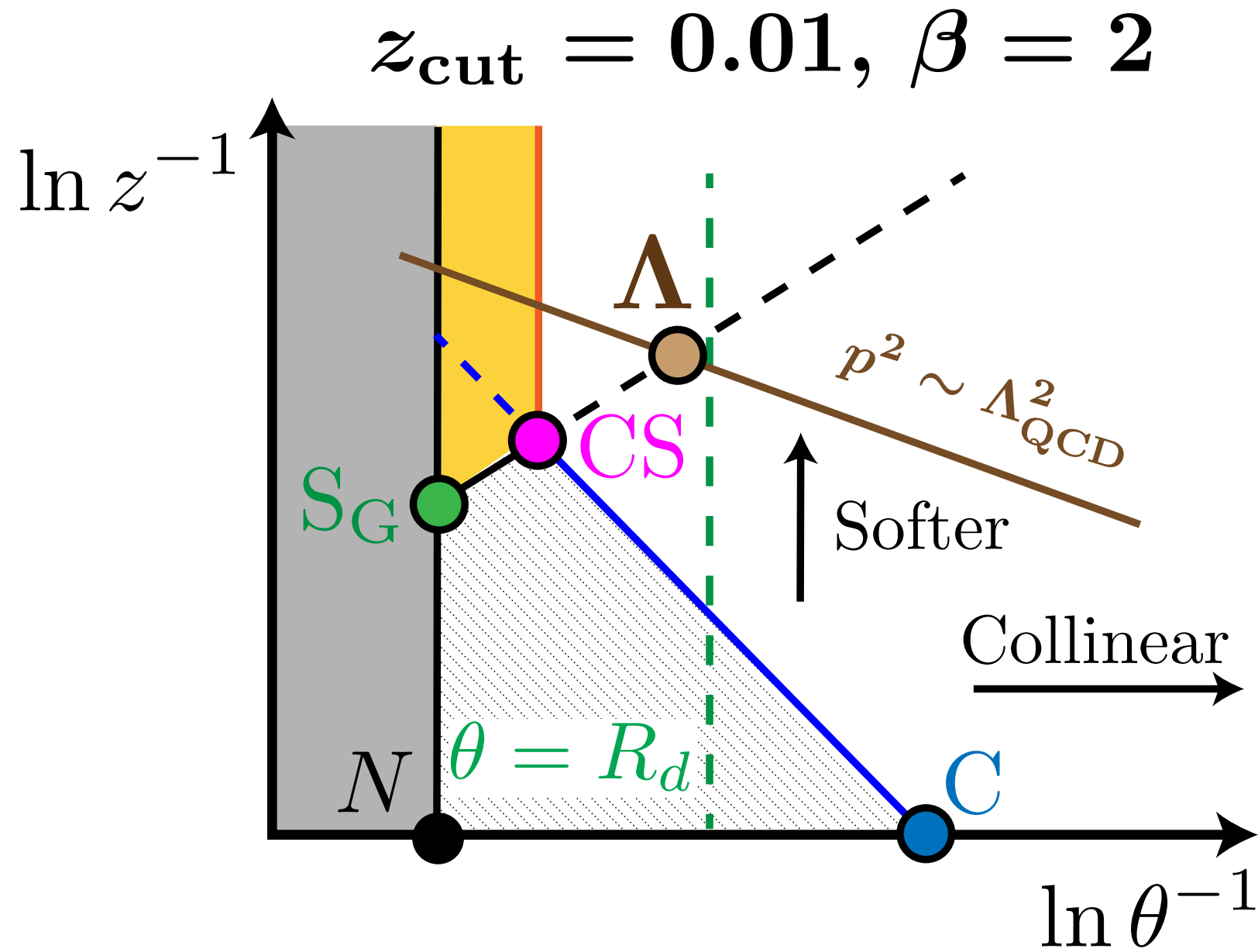
Inclusive treatment of the decay products



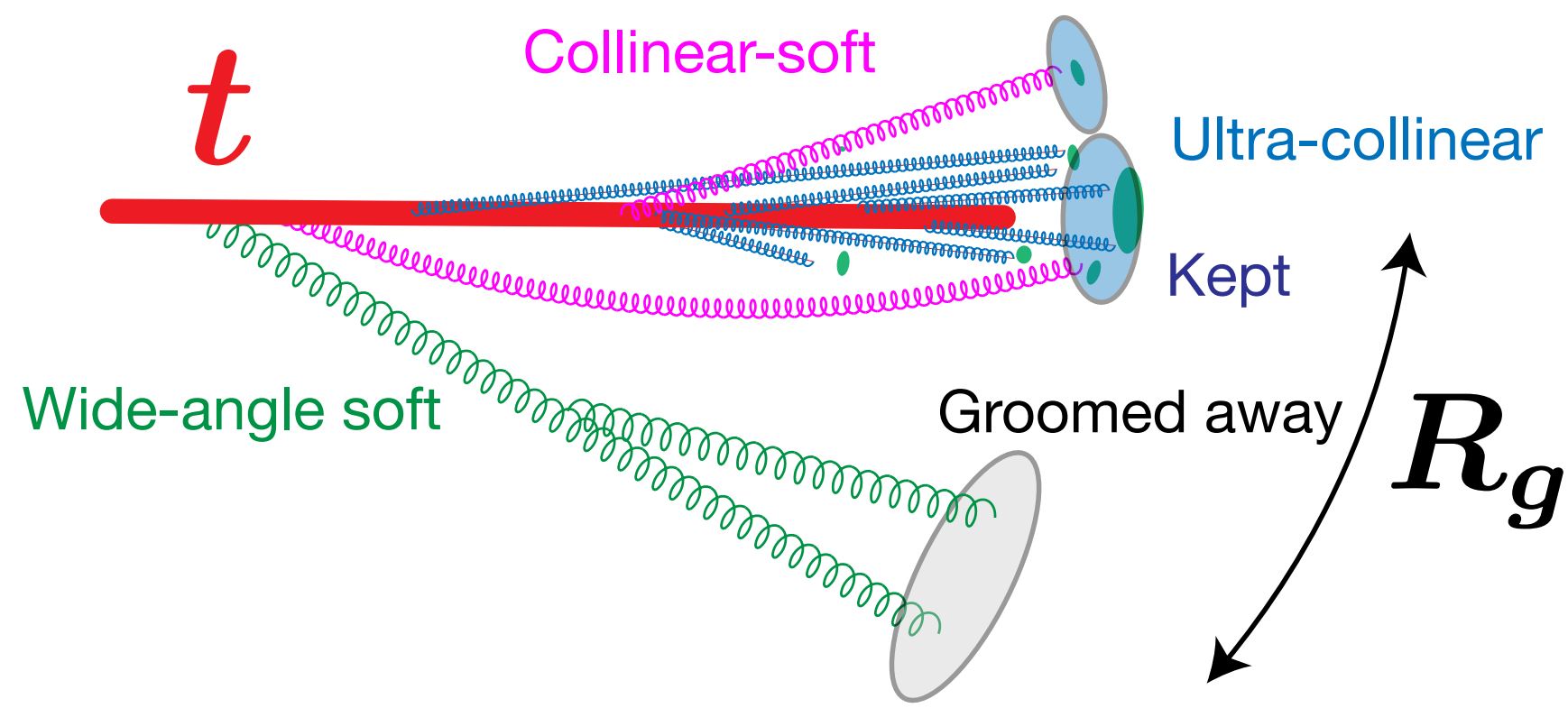
# SOFT DROP JET MASS OF BOOSTED TOP QUARKS

[JHEP 12 (2019) 002].

Effective field theory modes  $S_g$ ,  $CS$ ,  $C$  capture physics at different phase space points in the Lund plane.



Inclusive treatment of the decay products

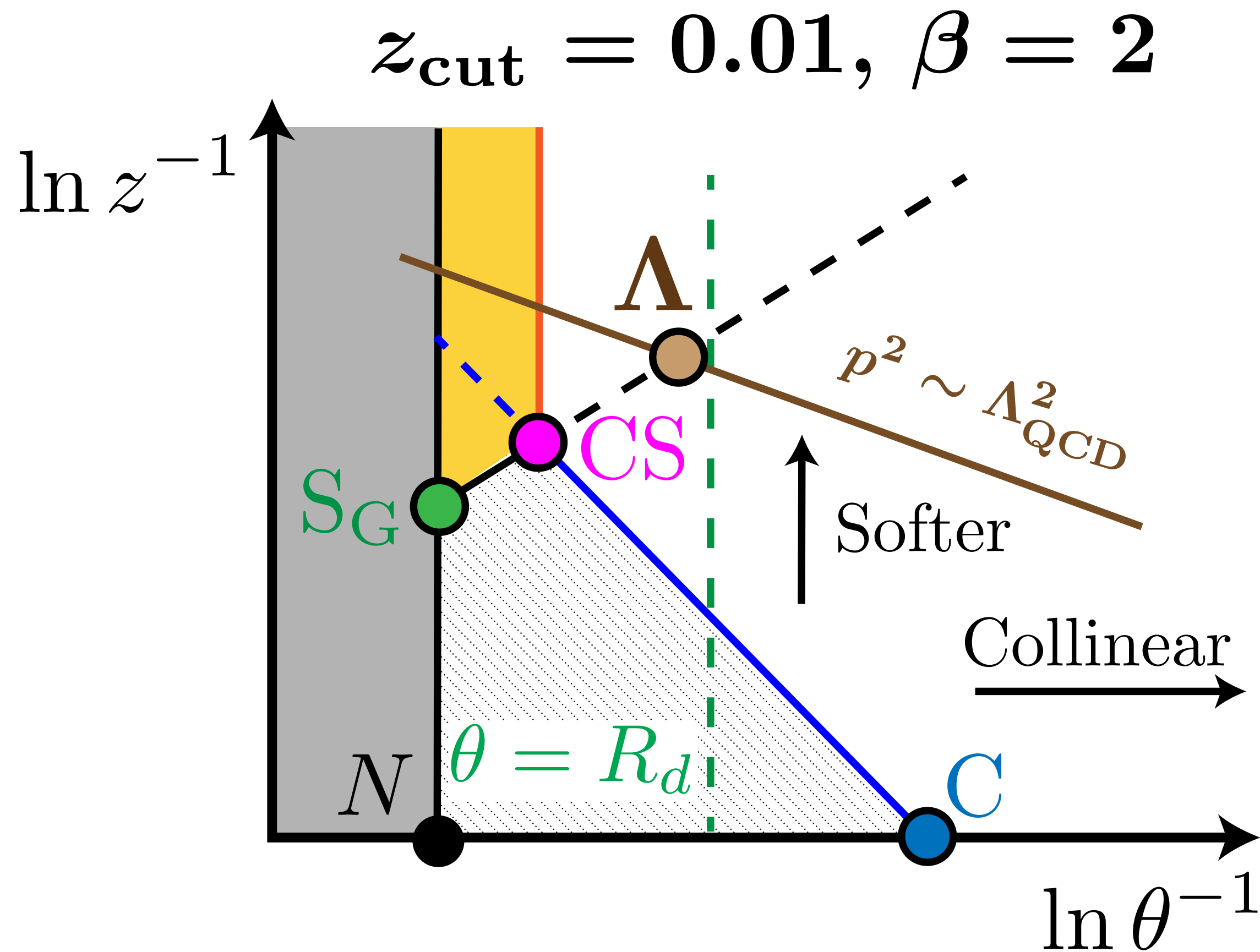


- This Lund plane picture is the same as for light quark and gluon jets.
  - Straightforward generalisation to the case of top quarks called the **light grooming region**.
    - Enables an “inclusive” description of the top quark decay.
  - For light grooming, the soft drop never encounters the ultracollinear radiation at small angular scales where the  $C$  mode is.
    - For light grooming ( $z_{cut} = 0.01$  and  $\beta = 2$ ) the soft drop stops at angles larger than the decay product opening angle:  $r_g > r_d$
    - Retain control over the mass scheme probed by this observable (main goal of this study).
- The mode  $\Lambda$  sits on the line  $p^2 \sim \Lambda_{QCD}^2$  and captures the leading non-perturbative effects on the groomed jet mass.

# SOFT DROP JET MASS OF BOOSTED TOP QUARKS

[JHEP 12 (2019) 002].

Effective field theory modes  $S_g$ ,  $CS$ ,  $C$  capture physics at different phase space points in the Lund plane.



• However, there are some undesirable features of this light grooming:

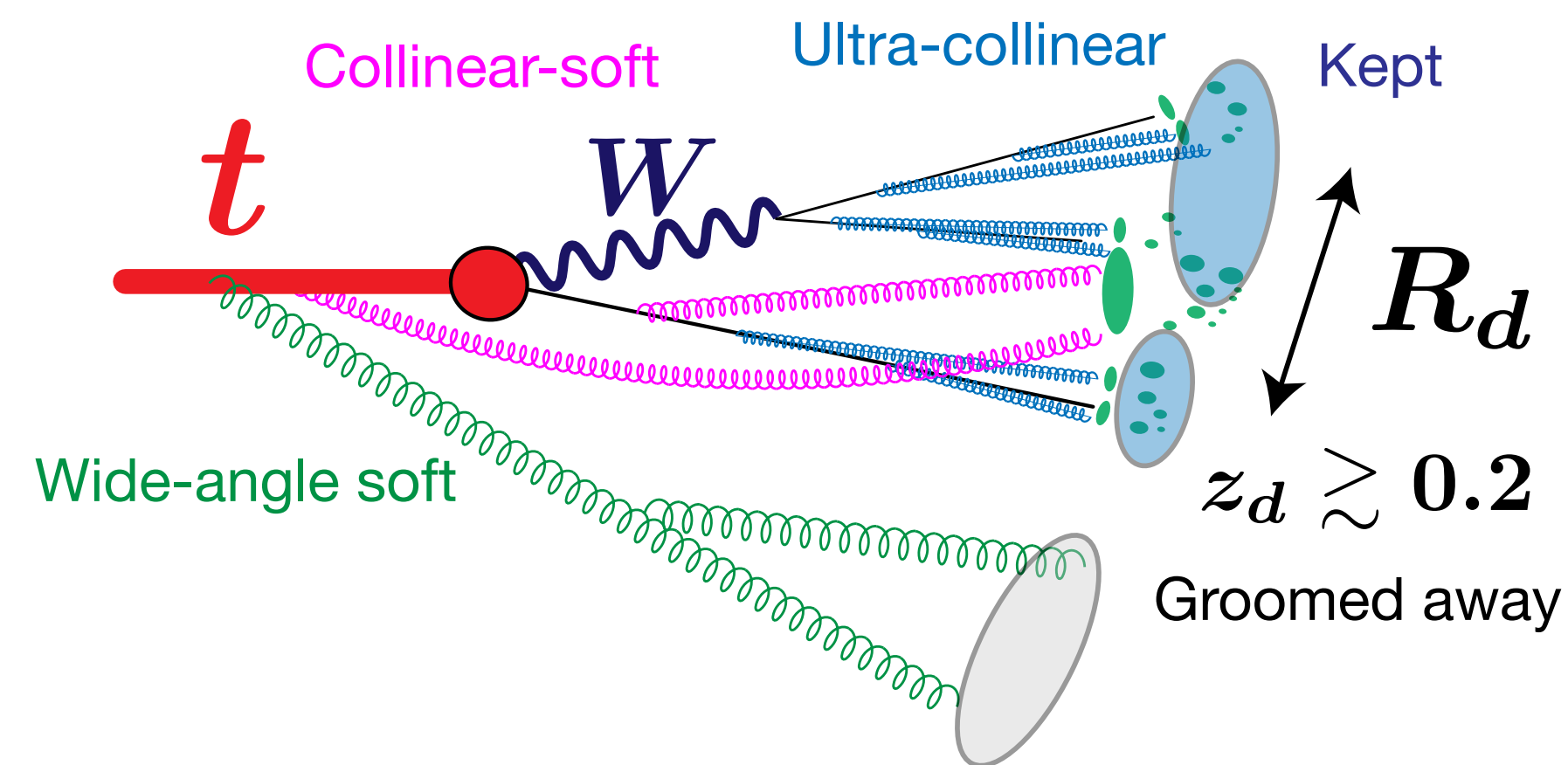
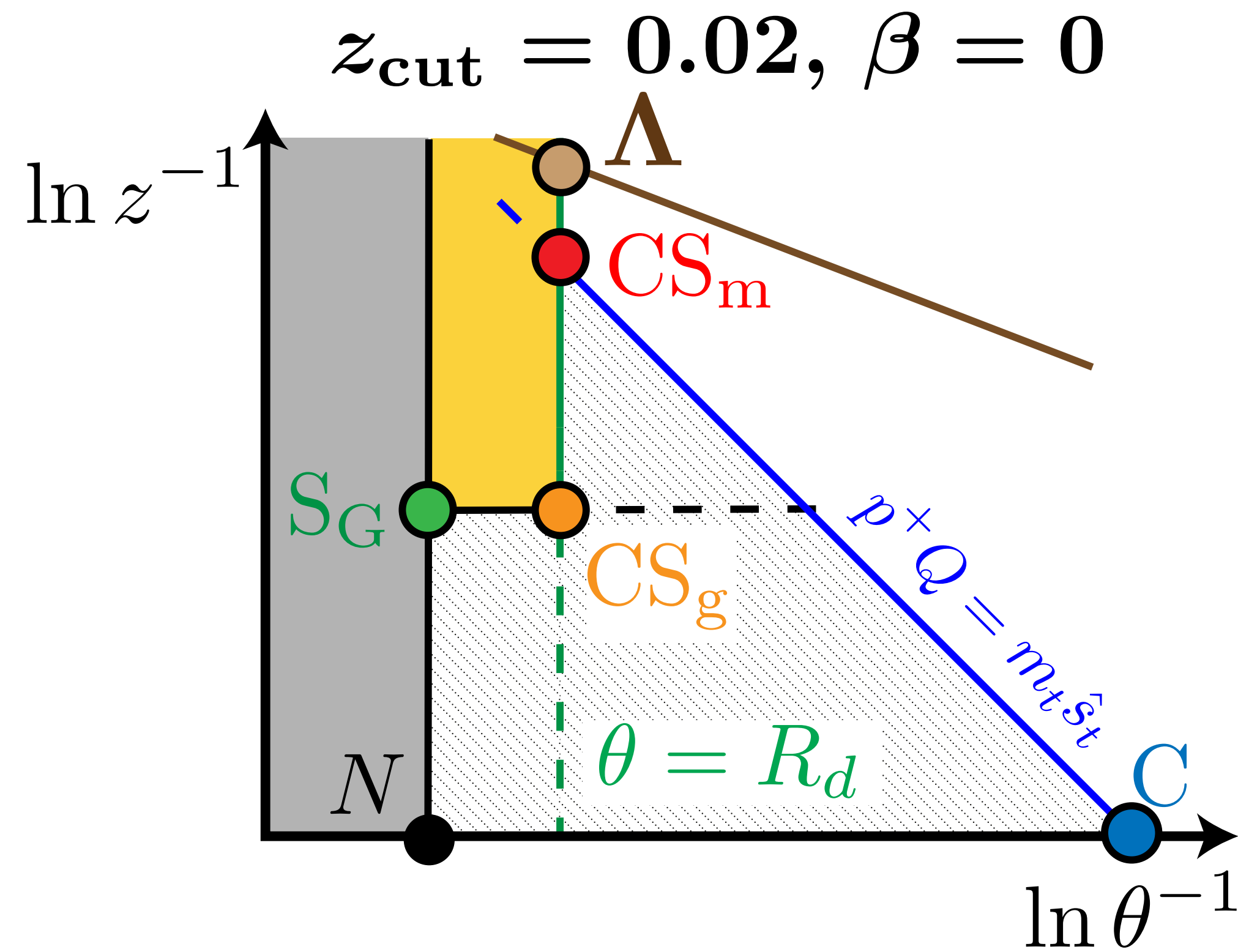
- Firstly, the grooming is small and it is not very effective in removing contamination.
- Secondly, a more theoretical reason is that the  $\Lambda$  mode sits on the intersection of the jet mass and the soft drop line.

➔ Means leading nonperturbative corrections are governed by the dynamics of soft drop.

▸ Not like in the case of ungroomed jets and makes things complicated.

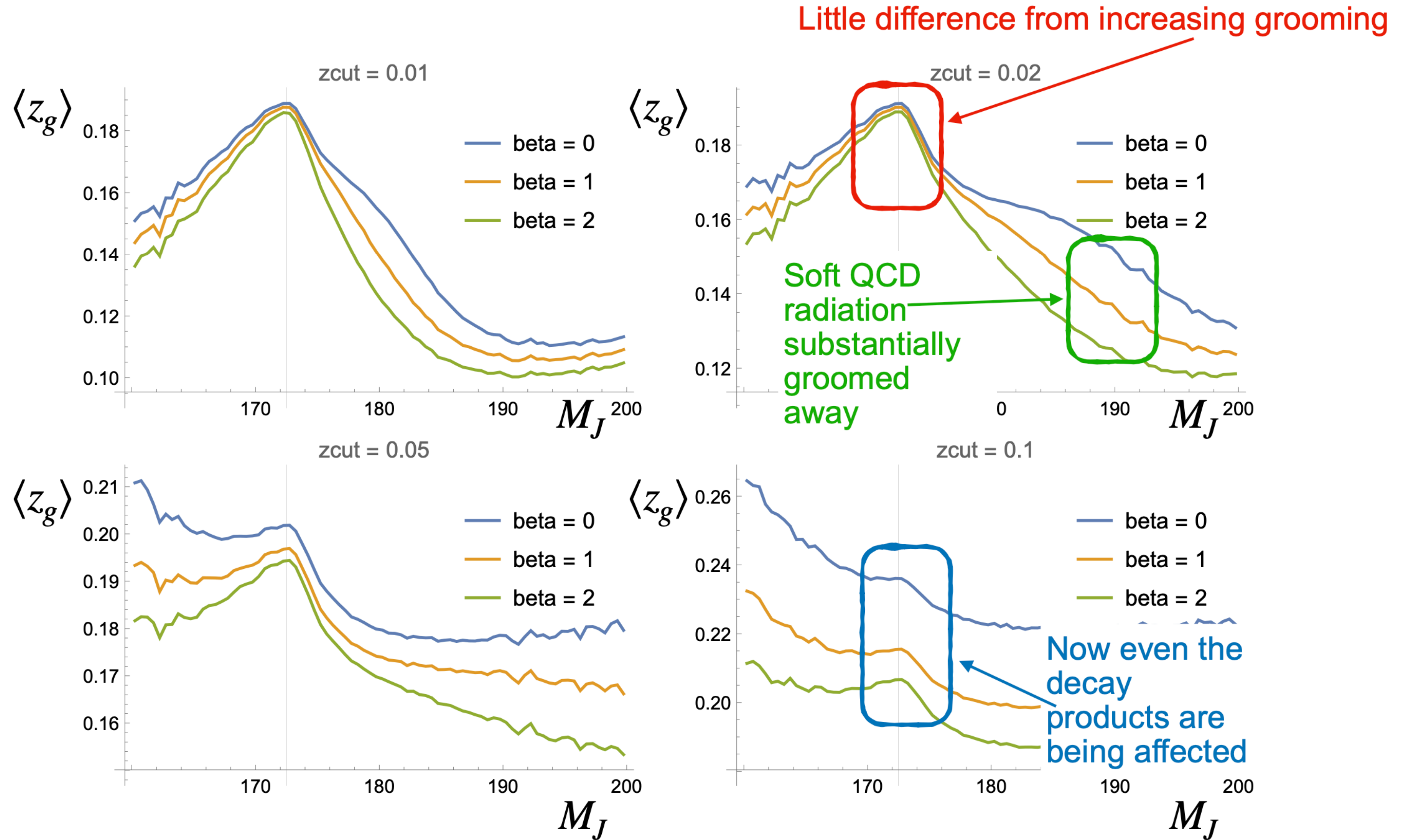
# CAN WE PUSH THIS FURTHER?

- How can we apply more aggressive grooming but still retain the inclusive description of the top decay.
- The top decay products effectively shield the ultra collinear radiation at small angular scales
  - The ultracollinear radiation remains clustered with the decay products.
- In this case a different set of modes are present in the factorisation.
  - Instead of the CS mode, the soft sector is now split into CS<sub>g</sub> and CS<sub>m</sub> modes.
- Here, the decay products effectively render interior of the jet as ungroomed.
  - The  $\Lambda$  mode location is now determined by the opening angle of the decay product  $r_d$  and is not influenced by the soft drop.
  - Enables a simpler description of the nonperturbative effects like ungroomed jets.



# PROOF OF CONCEPT

The decay products protect the ultra-collinear radiation.



# GOAL OF ANALYSIS

- **The interpretation of the top mass in an MC generator, in terms of a renormalised mass in the pole scheme:**

$$m_t^{\text{MC}} = m_t^{\text{pole}} + \Delta m_t^{\text{pole}}$$

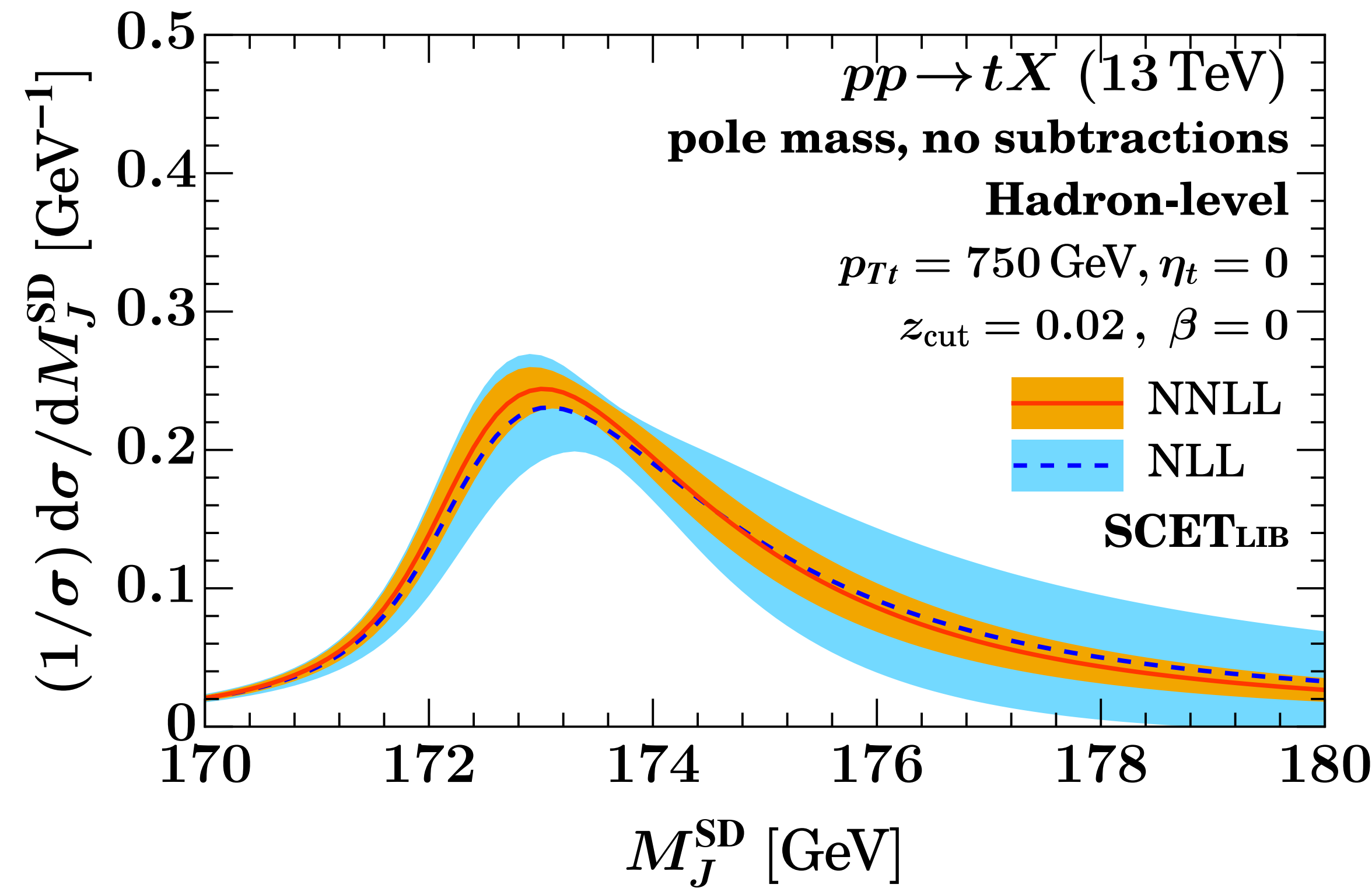
- Calibration performed with **NNLL calculation** compared against **Pythia MC** predictions with **NNPDF3.0 NLO PDF set** and **A14 set of tuned parameters**.

$m_t^{\text{MC}}$  is set to **172.5 GeV**.



# THEORETICAL CALCULATION

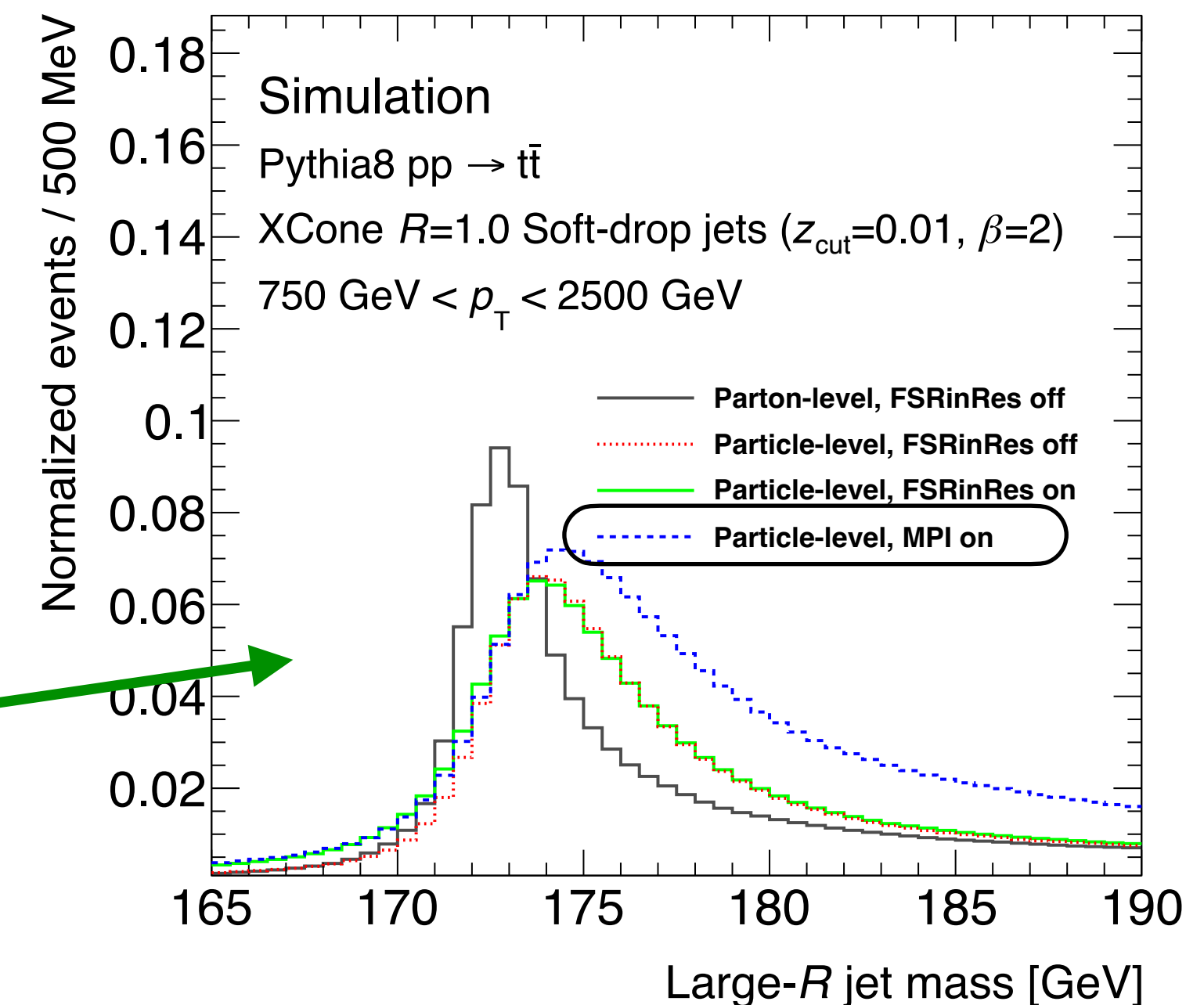
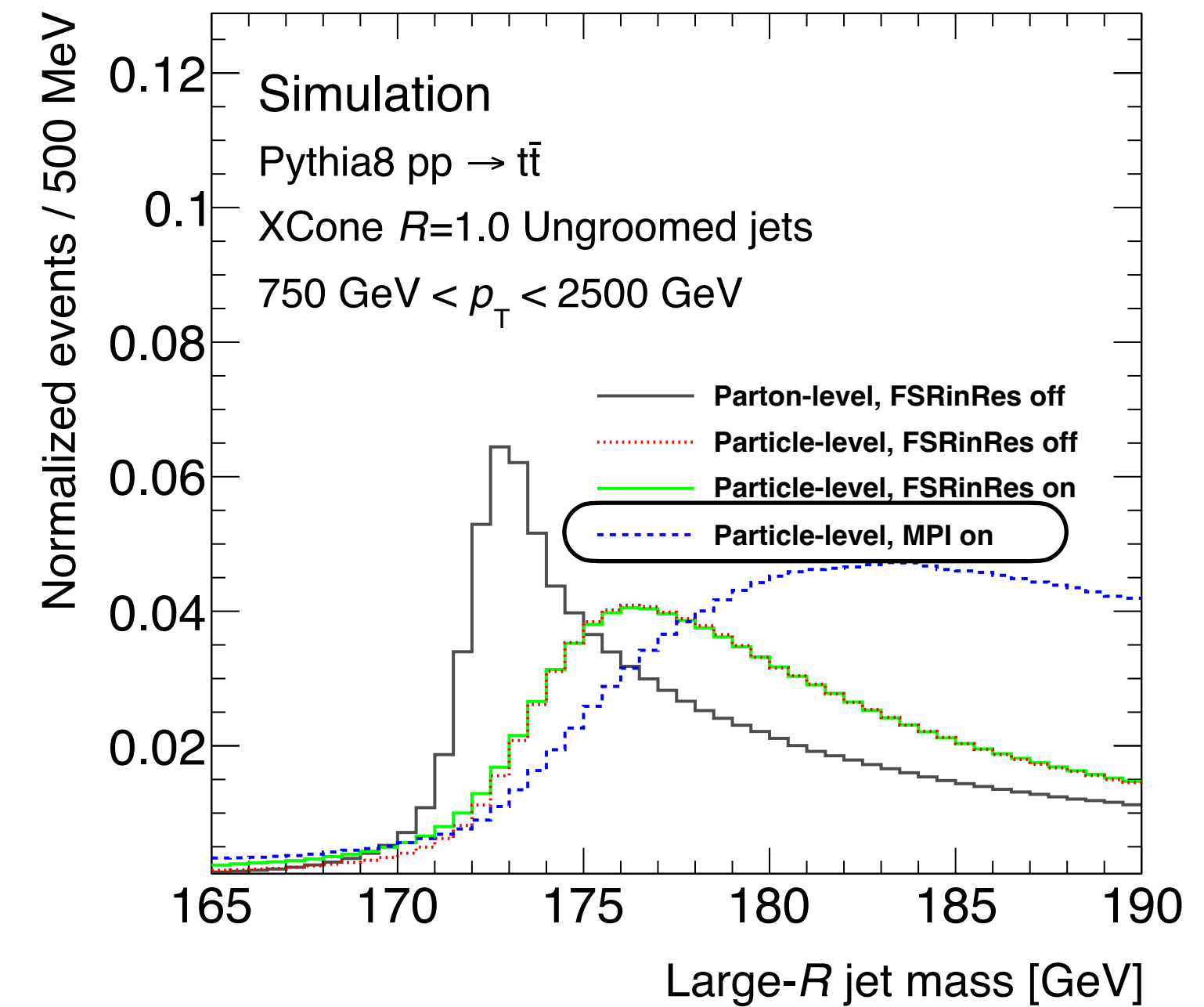
- **First-principle calculation with good control over the mass scheme.**
  - Yields particle-level predictions that can be compared directly to MC for a limited set of inclusive observables in boosted top production.
- **Continuation of top mass interpretations:**
  - $e^+e^- \rightarrow t\bar{t}$  processes NLL and NNLL accuracy.  
[Phys.Rev.Lett. 117 (2016) 23, 232001]  
[JHEP 12 (2023) 065]
  - $pp \rightarrow t\bar{t}$  processes at NLL accuracy.  
[ATL-PHYS-PUB-2021-034]
- **Using SCET-based calculation with NNLL accuracy**
  - Improved perturbative stability.



MICHEL, PATHAK, STEWART  
IN PREPARATION

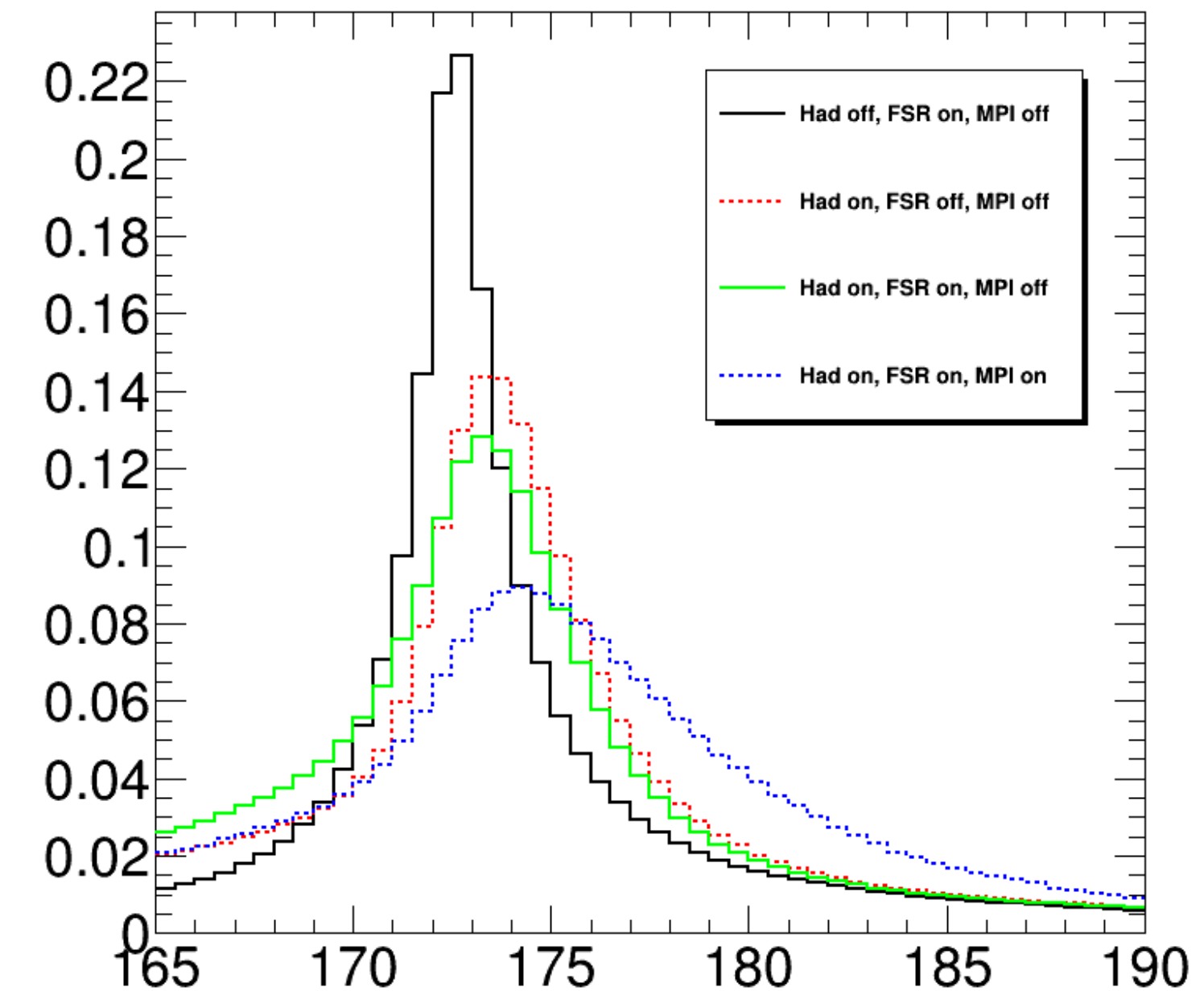
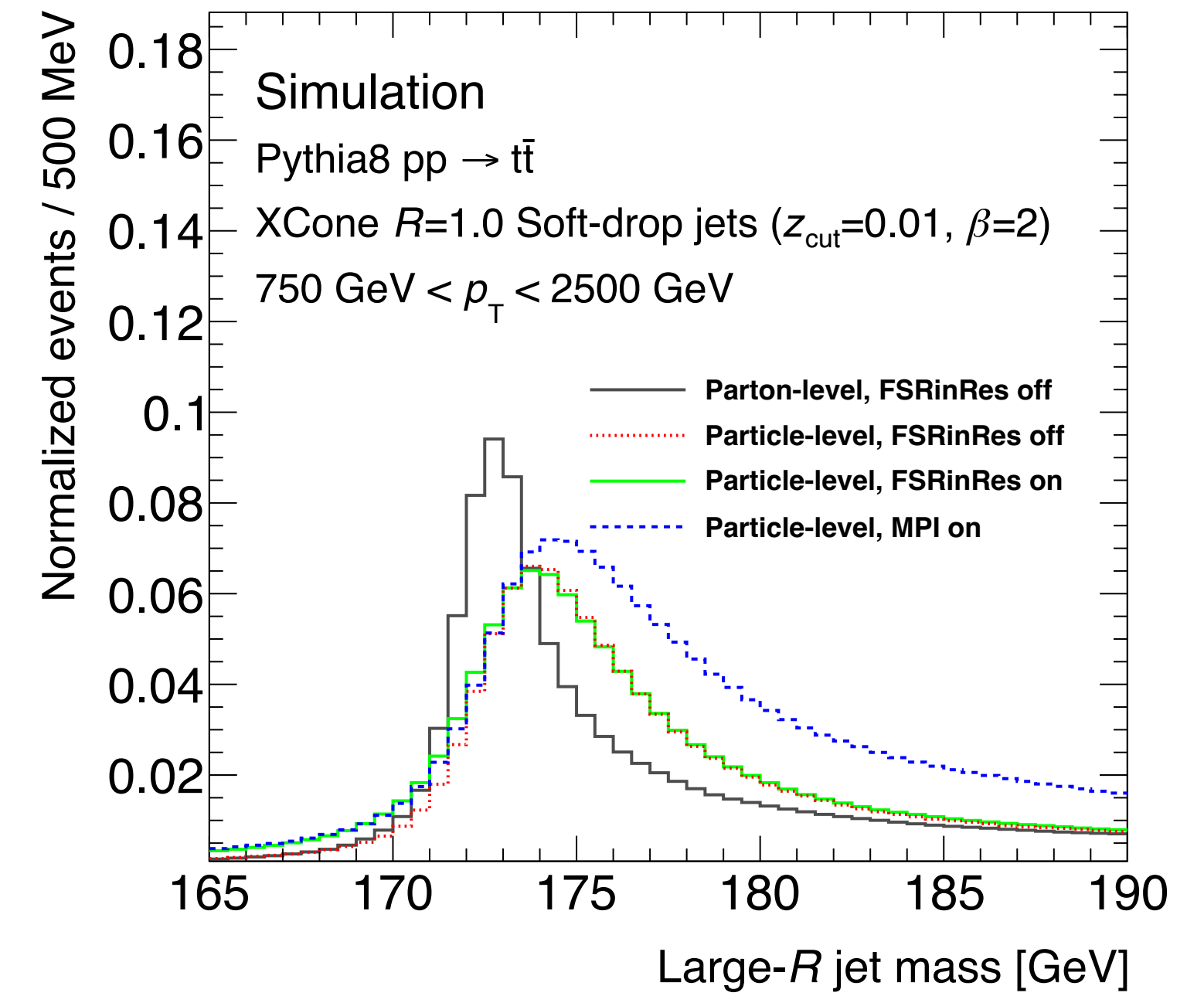
# JET BUILDING

- Focus on particle-level hadronic top quark decay in  $pp \rightarrow t\bar{t}$  and  $e^+e^- \rightarrow t\bar{t}$  processes.
  - Top mass determined by fitting large-R jet mass containing hadronic top.
  - Mass reconstructed using information from decay products of top quark within large-R jet.
- Large-R jets built with:
  - **XCone** jet algorithm with **R = 1**.
    - Jet algorithm minimising N-jettiness. Useful filtering out unwanted jets in densely populated events (useful in boosted regime where signal jets may partially overlap.)
  - **Soft-drop light grooming** applied to remove soft-wide radiation ( $z_{cut} = 0.01$ ,  $\beta = 2$ ).
    - Considerably reduces UE impact. Shift of  $\sim 5$  GeV down to  $\sim 1$  GeV.



# NEW JET BUILDING

- **Soft-drop light grooming** applied to remove soft-wide radiation ( $z_{cut} = 0.01, \beta = 2$ ).
  - Considerably **reduces UE impact**. Shift of  $\sim 5$  GeV down to  $\sim 1$  GeV.
- **Now possible to use a more aggressive grooming scheme** ( $z_{cut} = 0.02, \beta = 0$ ).
  - **Reduces UE impact even further.**



# HADRONIZATION AND UNDERLYING EVENT: EXPECTATION

- Peak position in new grooming scheme:

$$M_J^{Peak} \sim m_t + \underbrace{\Gamma_t(1 + \alpha_s \dots)}_{\text{Hadronisation effects}} + h \times \Omega_1 + \underbrace{\langle R_d^4 \rangle_{M_J}}_{\text{UE effects}} \Lambda_{UE}$$

- The top quark jet only depends on  $h$ , defined as  $R_d = \frac{m_t}{p_T} h$ 
  - We now also bin in  $h$  in addition to  $p_T$  (an independent variable).
  - Underlying event contribution  $\Lambda_{UE}$  depends on  $p_T$  (depends on the catchment area). **Can disentangle underlying event by considering different  $p_T$  bins.**

# JET BUILDING

- Focus on particle-level hadronic top quark decay in  $pp \rightarrow t\bar{t}$  and  $e^+e^- \rightarrow t\bar{t}$  processes.
  - Top mass determined by fitting large-R jet mass containing hadronic top.
  - Mass reconstructed using information from decay products of top quark within large-R jet.

- Boosted jet - Inclusive treatment of decay products:

- Previously used four orthogonal jet  $p_T$  bins:

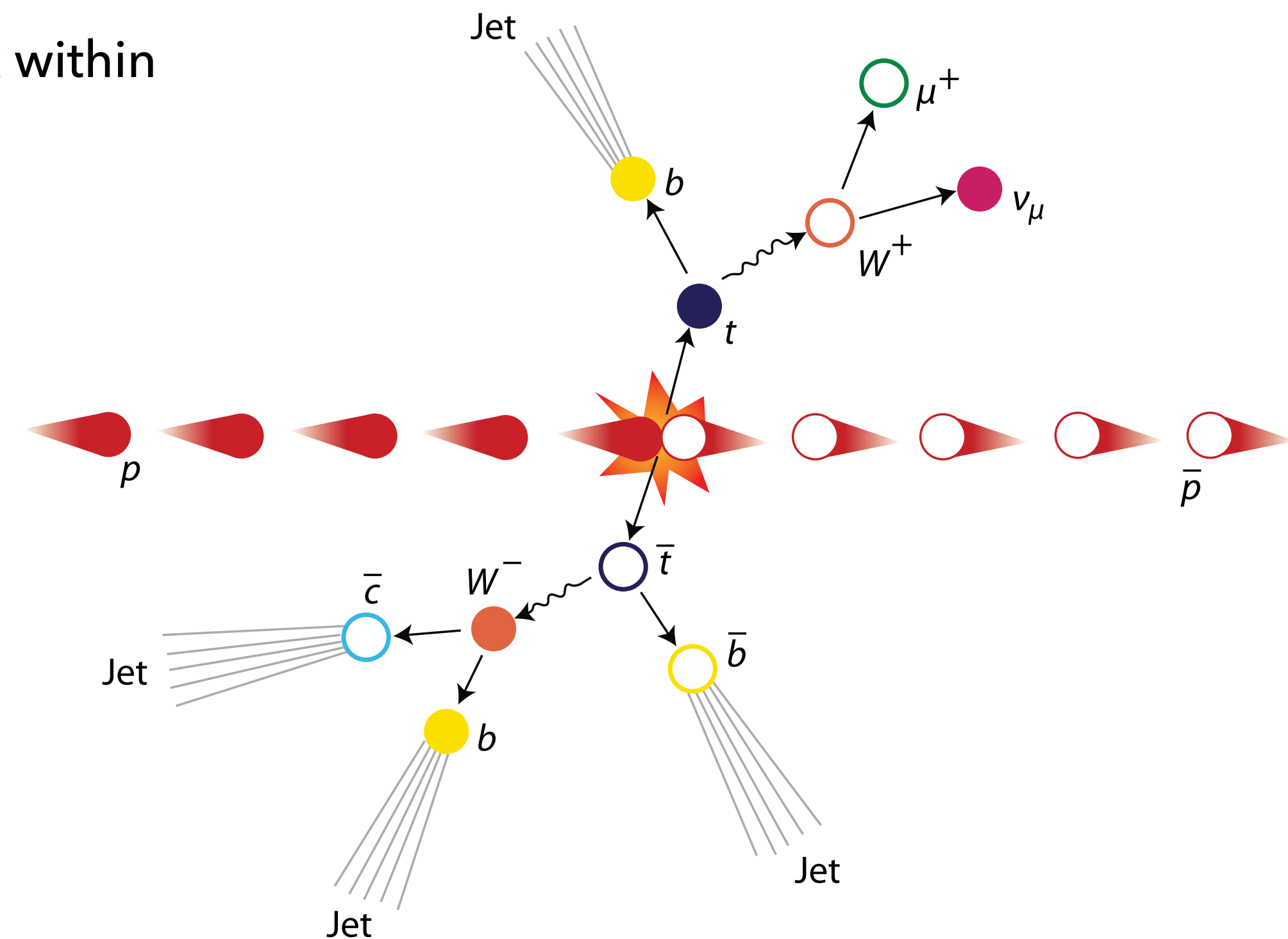
$$p_T^{jet} \in \{750, 1000, 1500, 2000, 2500\} \text{ GeV.}$$

- Now use two orthogonal jet  $p_T$  bins:

$$p_T^{jet} \in \{750, 1250, 2500\} \text{ GeV.}$$

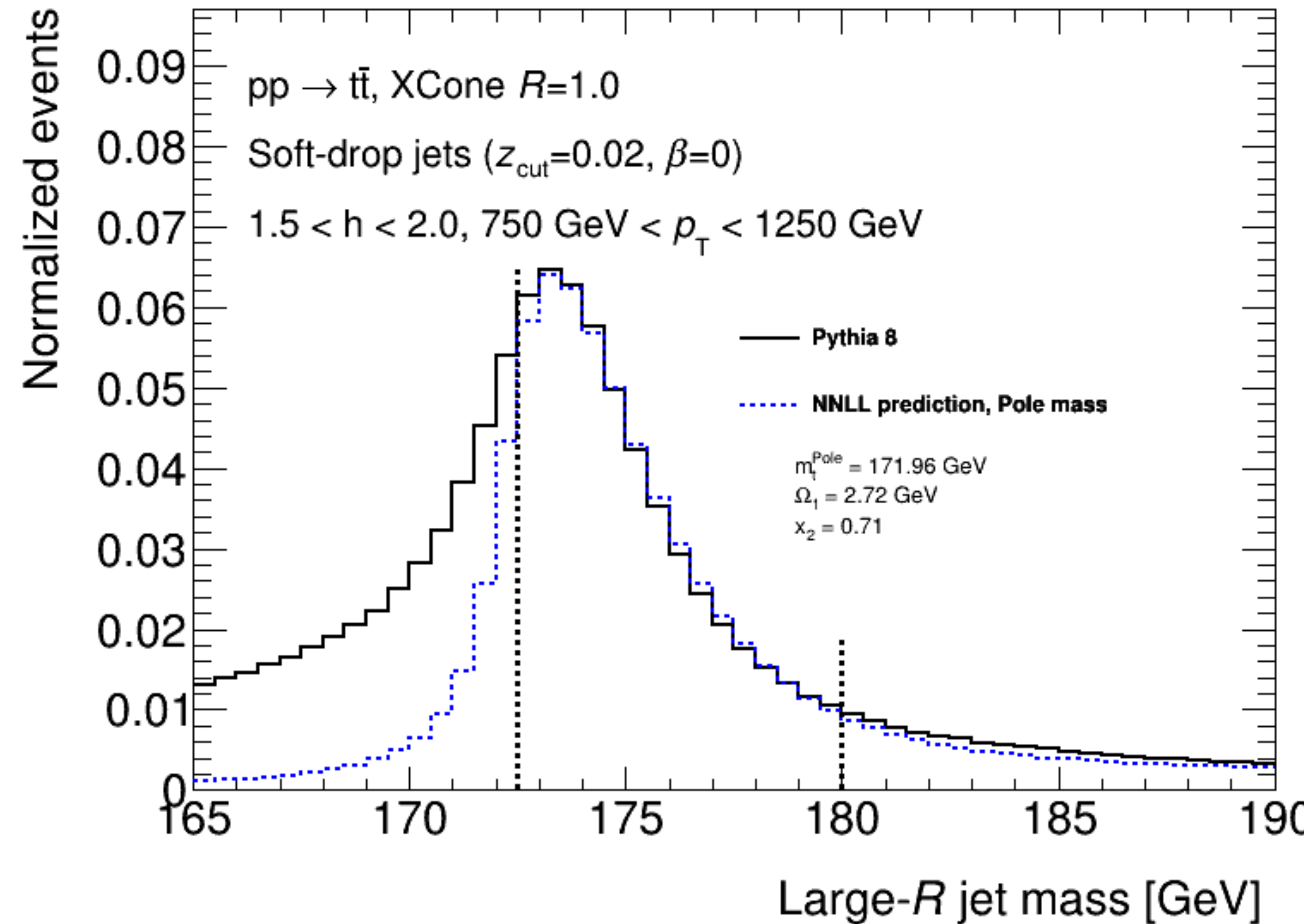
- and four jet h bins:

$$h^{jet} \in \{1.5, 2, 2.5, 3, 3.5\}.$$



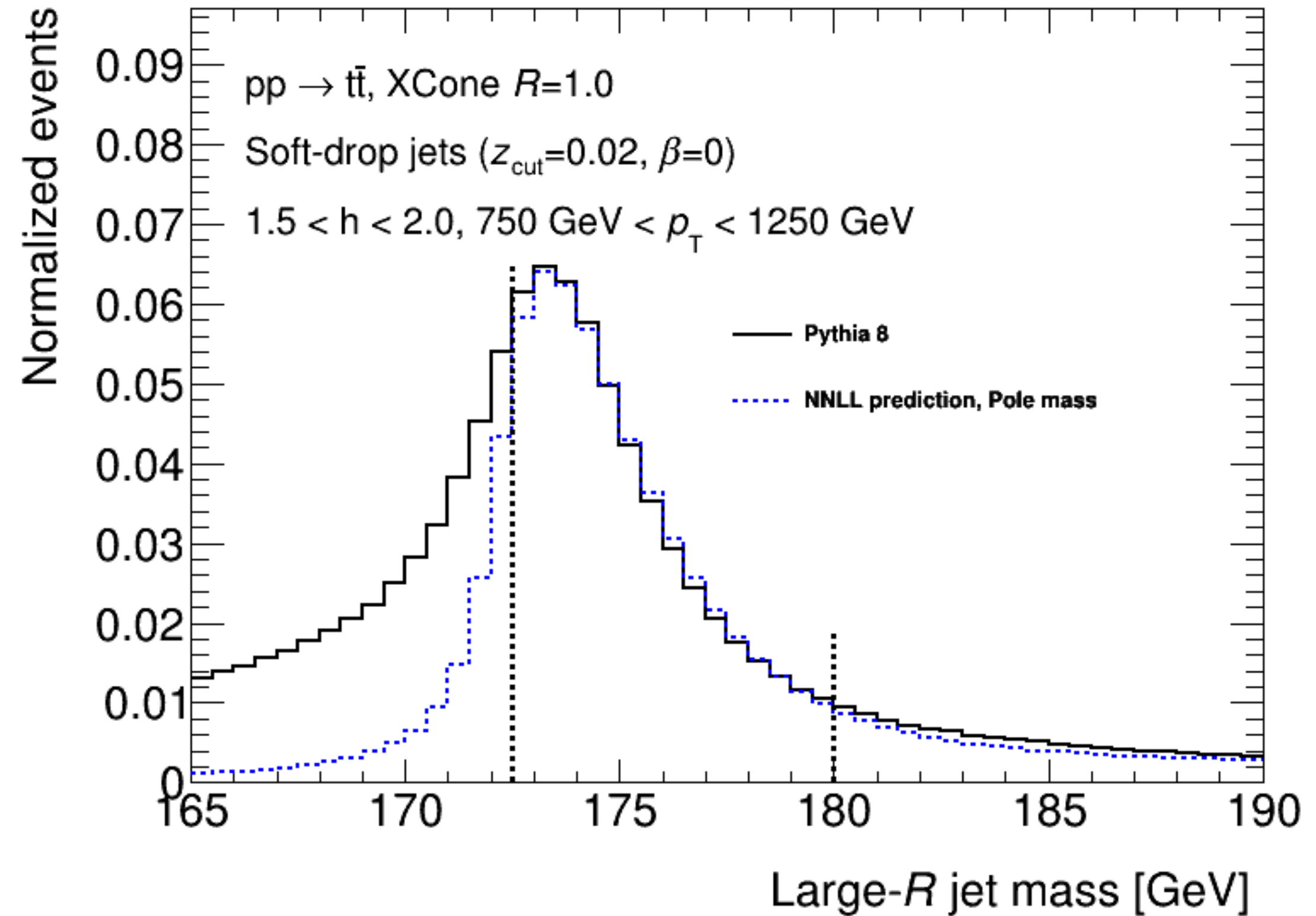
# FITTING DETAILS

- Model uses three parameters,  $m_t^{Pole}$ ,  $\Omega_1^{had}$ , and  $x_2$  associated with **first-** and **second-moment non-perturbative corrections**.
- Idea is to obtain **value of parameters in NNLL theory** calculation that **best describe MC prediction**.
- $m_t^{Pole}$ ,  $\Omega_1^{had}$ , and  $x_2$  varied:
  - Best fit of MC-to-theory distributions found for variations of the three parameters.
  - $\chi^2$  minimisation fit applied to the three parameters to find the **global minimum**.
  - **This is how we extract the top quark pole mass.**



# FITTING DETAILS

- Idea is to obtain **value of parameters in NNLL theory calculation that best describe MC prediction.**
  - **Decay product FSR effects are not yet included in calculation where we treat decay inclusively.**
    - In grooming procedure, **theory does not accurately describe the low-mass tail** present in the generator prediction.
    - Must **restrict fit range to avoid the low jet-mass tail**, that would **bias the extracted top mass to lower values.**
- ➔ **Fit range set to 172.5-180 GeV.**

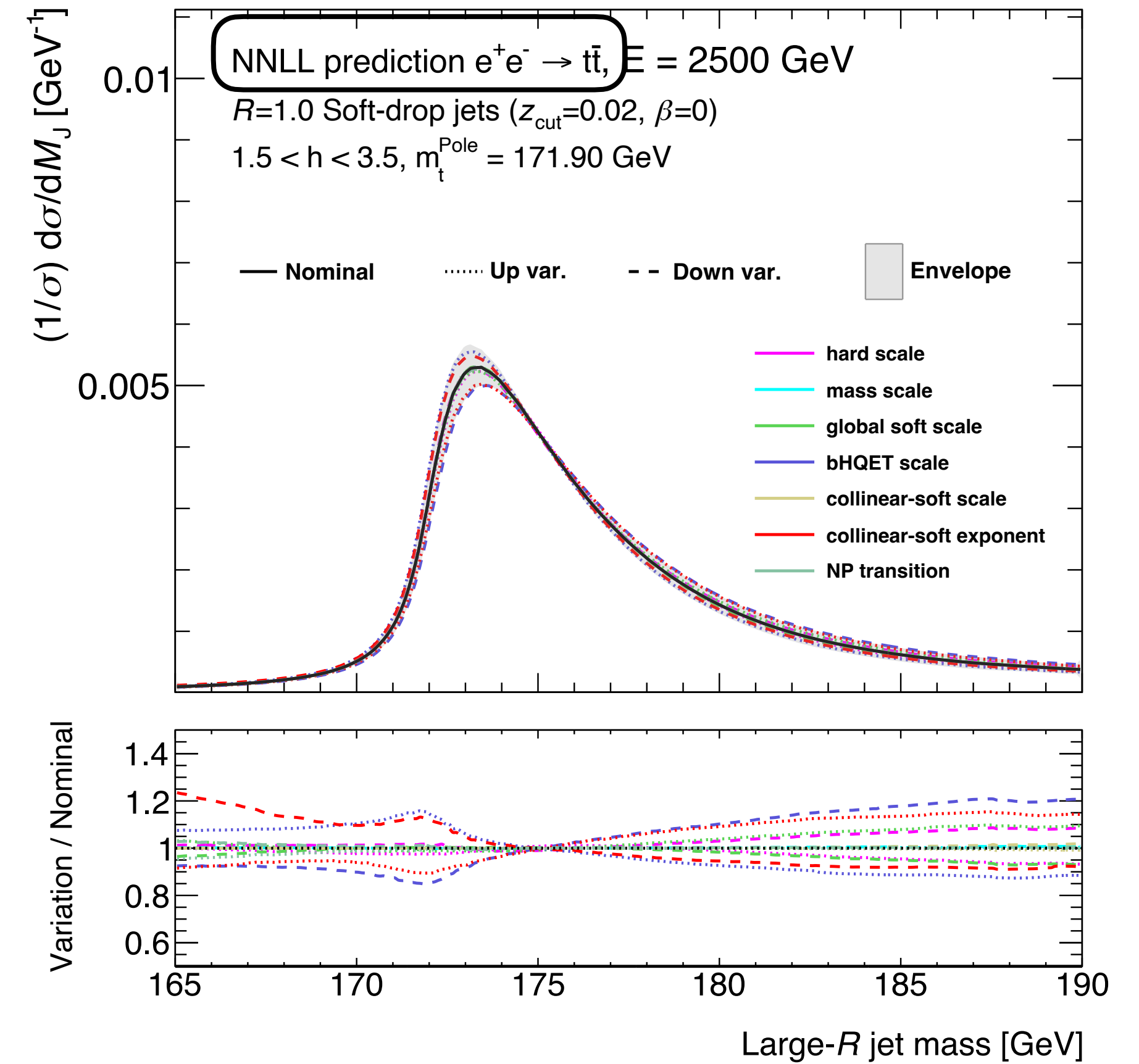
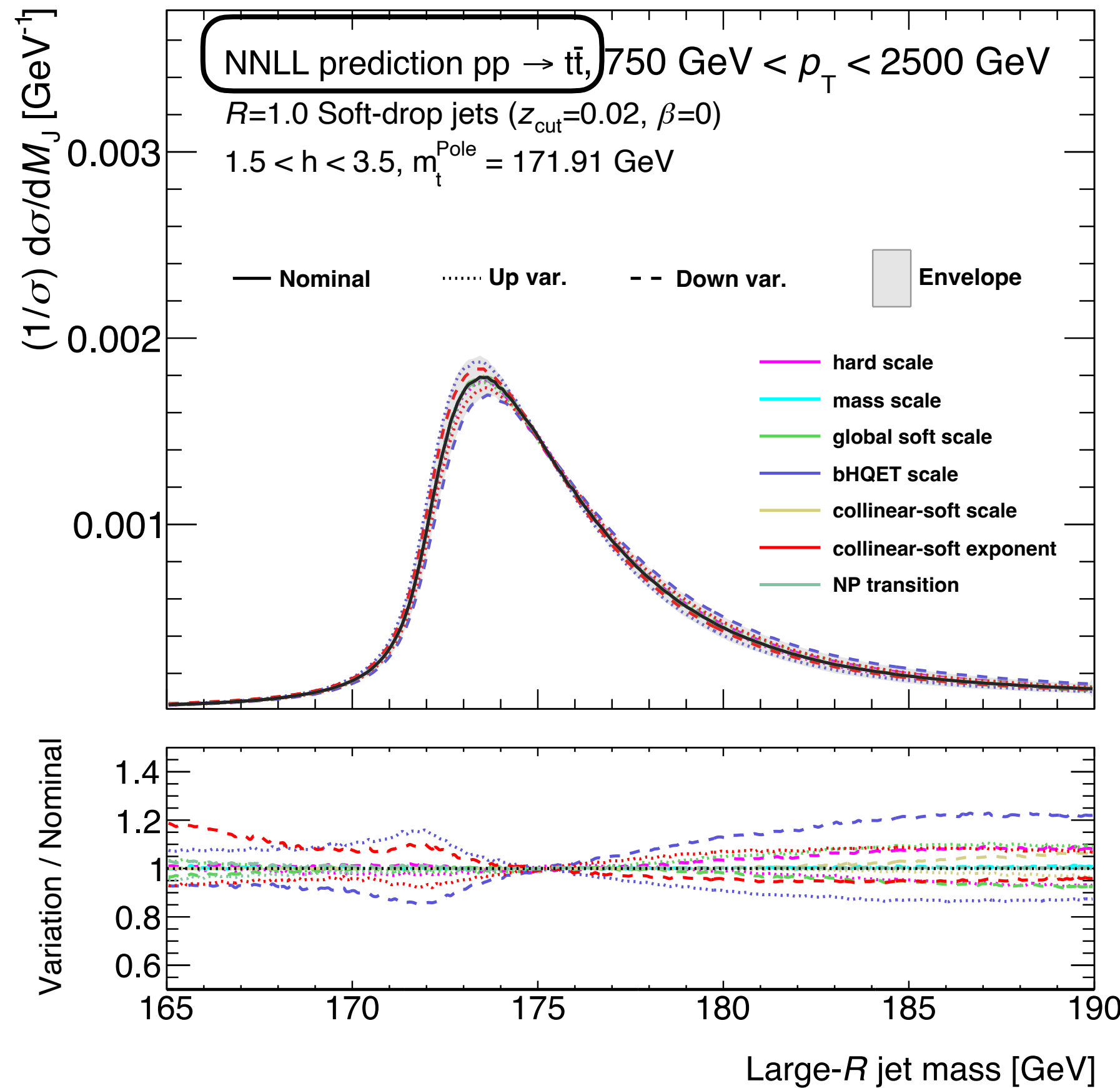


- Perform a **fit range study.**
  - Measure the top quark mass value at **172-180 GeV** and **173-180 GeV**
- ➔ Estimate an uncertainty.

# THEORETICAL UNCERTAINTIES

- Theoretical uncertainty determined by jet mass dependent renormalisation scales that to estimate the perturbative uncertainty.

Scale variations on these dependencies measured and compared to central value.



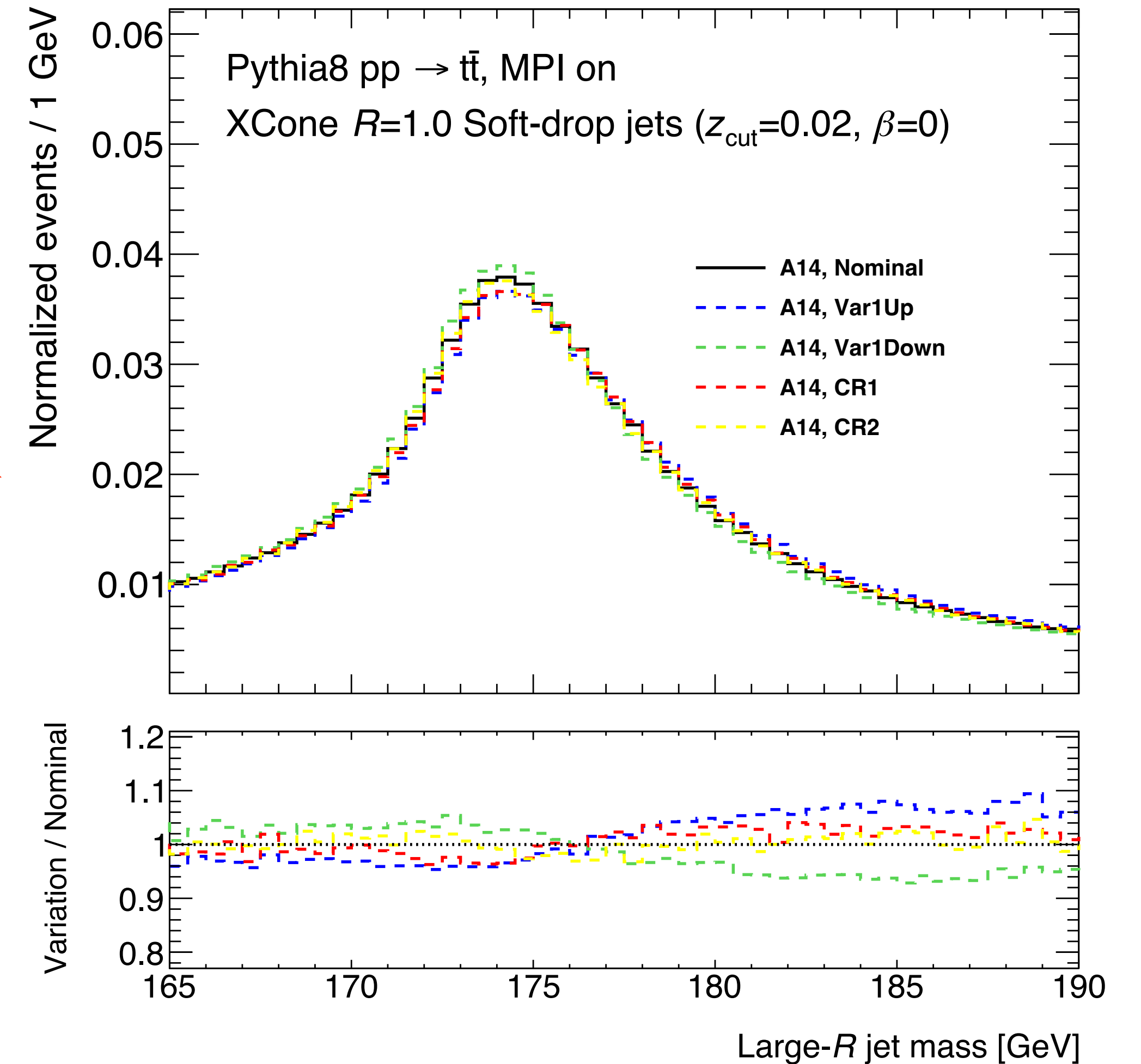


# FITTING DETAILS

- Need to cover any bias of the influence of the kinematic ranges on the mass relation.
  - **Impact of the choice of large-R jet  $h$  is evaluated.**
  - **Compare fits on sub-sets of three  $h$  intervals for all permutations of the set of the 4  $h$  bins.**
  - **Maximum variation taken as the uncertainty.**
    - Also apply this for  $p_T$  impact.

# UNDERLYING EVENT UNCERTAINTIES

- UE for now must be estimated through simulation parameter changes in MC-to-MC fits.
  - Comparing nominal MPI-on Pythia against **A14 eigentune variations** (coverage of **UE variations modelling uncertainties**).
- Not yet testing the UE peak shift hypothesis
- UE contribution of the jet mass peak is disentangled from hadronisation effects:
  - Can be possible to calculate the UE effect on the relation result.
  - Future work.



# UNCERTAINTY

- Uncertainties are applied to account for:
  - **Estimation of perturbative uncertainty in calculation.**
  - **Fitting methodology (FSR estimation not present in calculation).**
  - **$h$  and  $p_T$  influence of large-R jet.**
  - **UE not yet present in the calculation.**

	Uncertainty (ee) [MeV]	Uncertainty (pp) [MeV]
Theoretical	+160/-210	+150/-250
Fitting	220	260
Kinematic range	-30/-50	+20/-90
Underlying event	N/A	+225/-180
<b>Total</b>	<b>+275/-215</b>	<b>+375/-320</b>

# RELATION RESULTS

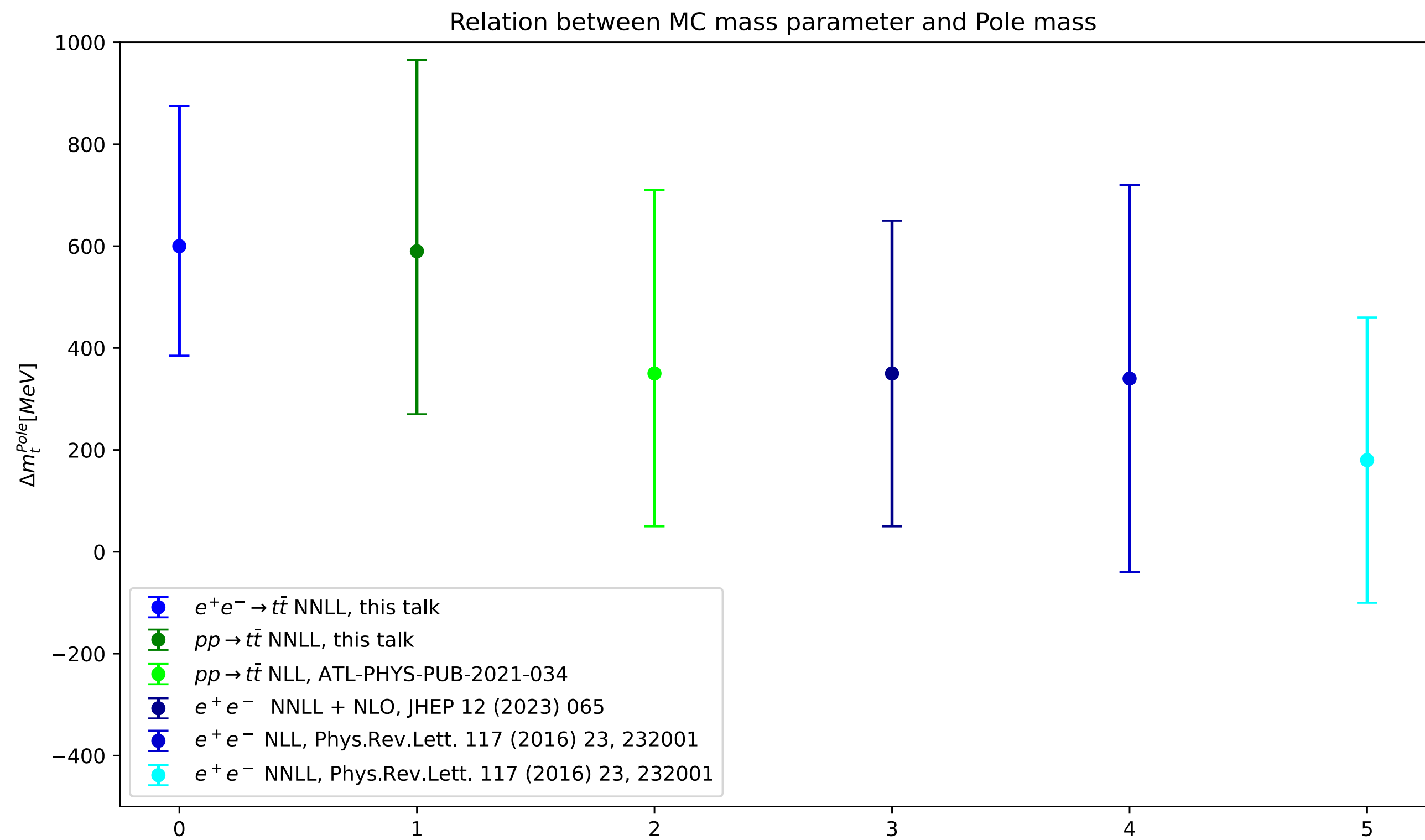
Mass relation (pp):

$$\Delta^{Pole} = m_t^{MC} - m_t^{Pole} = 590^{+375}_{-320} \text{ MeV}$$

Mass relation ( $e^+e^-$ ):

$$\Delta^{Pole} = m_t^{MC} - m_t^{Pole} = 600^{+275}_{-215} \text{ MeV}$$

**Suggests universality between the  $e^+e^-$  and  $pp$  processes.**



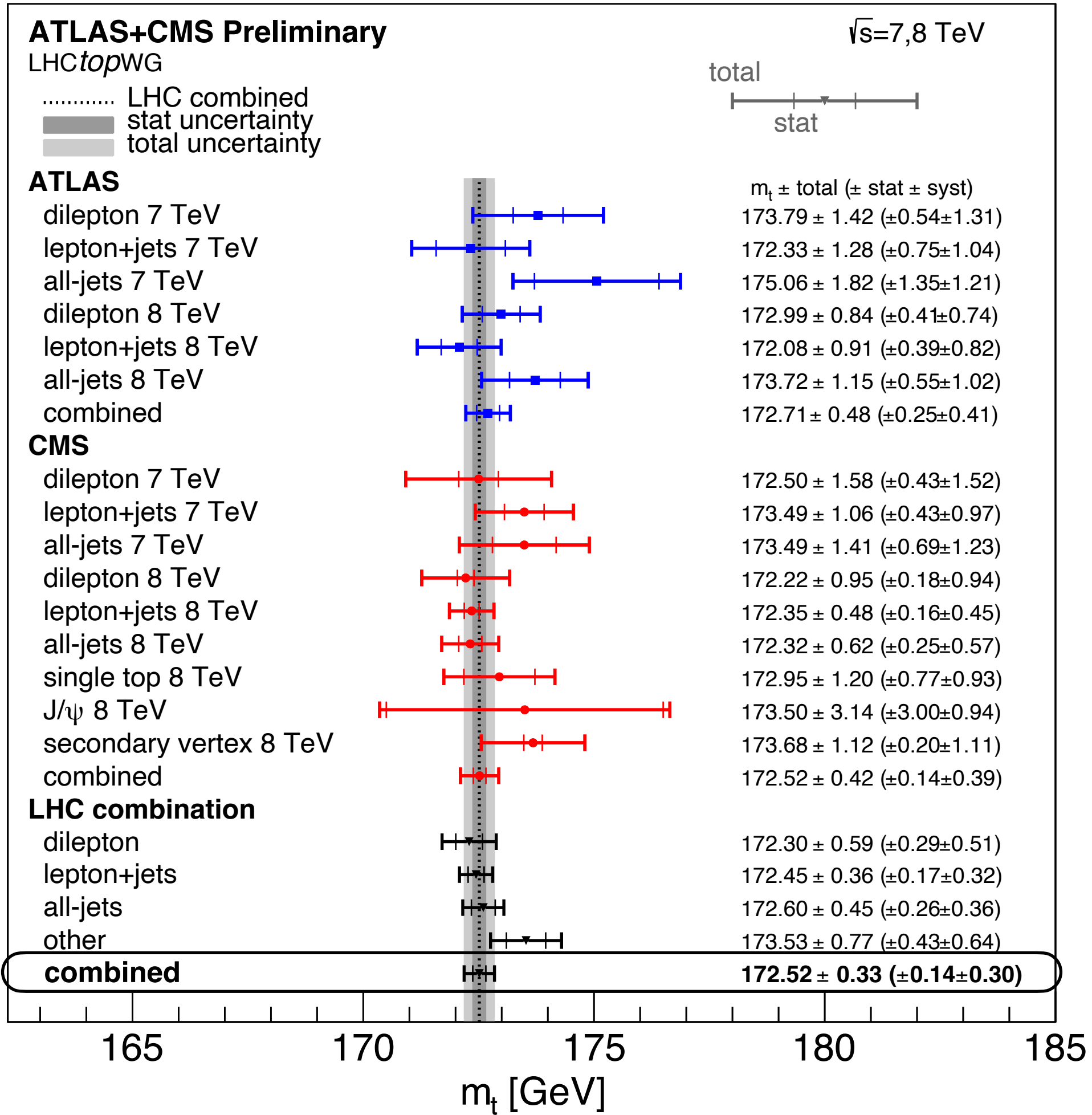
- Comparisons with previous top quark mass interpretations in  $pp$  and  $e^+e^-$  collision regimes.
  - All results compatible within uncertainties.

# CONCLUSION

- **Preliminary results have found a working procedure for the new grooming method described in this talk.**
- Accounts for more physics effects while giving a relatively accurate result.
- Relation results of 590 and 600 MeV for  $pp$  and  $e^+e^-$ , respectively are promising considering their compatibility.
  - ▶ Suggests universality for different processes of  $t\bar{t}$  interactions.
- Work to be done with Herwig samples that should closer reflect the hadronisation effects of the theory and give more accurate results.
- More time necessary to fully understand these results and write into a comprehensive paper.
- Future would be applying this to MSR mass for more accurate and applicable results.

**Thanks for listening**

# TOP QUARK MASS MEASUREMENTS – DIRECT



- $m_t$  determined **experimentally** by studying **top quark decay products**.

- **MC-based** templates at **detector level** - combination of first principle QCD calculations and modelling techniques (e.g. hadronisation and parton shower).

- Hard reactions at **high energy + low energy QCD effects**.

- **Most precise** determination of top quark mass  $m_t^{MC} \mathcal{O}$  **(330) MeV precision**.

Average  $m_t^{MC}$  from LHC top WG combination:

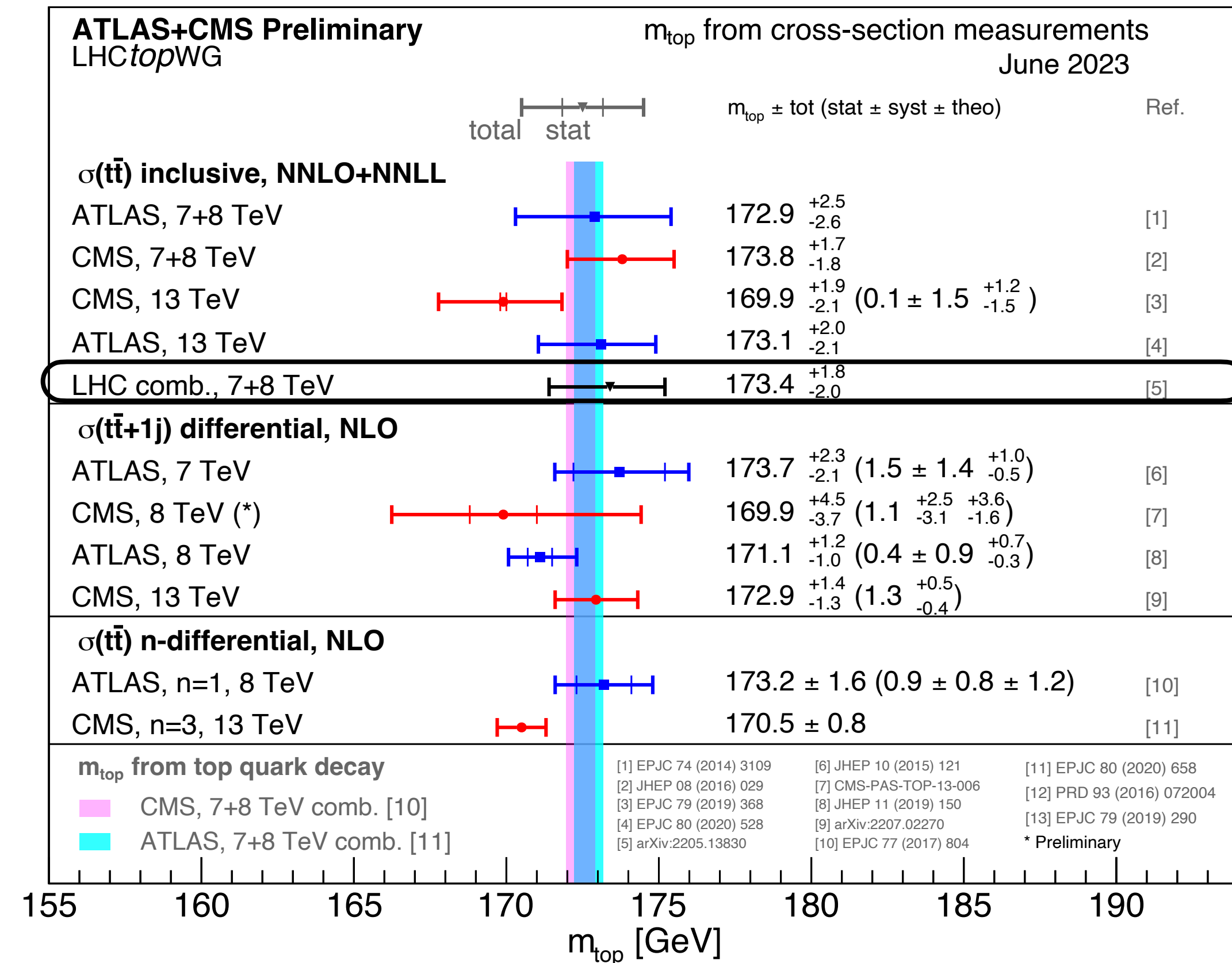
$$m_t^{MC} = 172.52 \pm 0.33 \text{ GeV}$$

# TOP QUARK MASS MEASUREMENTS – INDIRECT

- Aims for accurate measurements in terms of the **Lagrangian mass parameter**.
- $m_t$  determined by analysing **parton-level production cross-sections** (inclusive and differential).
  - Calculations sensitive to **hard scatter at high energy scales**, where top quarks are produced.
- **Most precise measurements  $\mathcal{O}(1)$  GeV precision**.
  - $m_t^{pole}$  value from **inclusive** cross section measurements  
LHC top WG combination:

$$m_t^{pole} = 173.4^{+1.8}_{-2.0} \text{ GeV}$$

Phys. Rev. Lett. 132, 261902



Differential cross sections yield more precise results.



# PREVIOUS RESULT

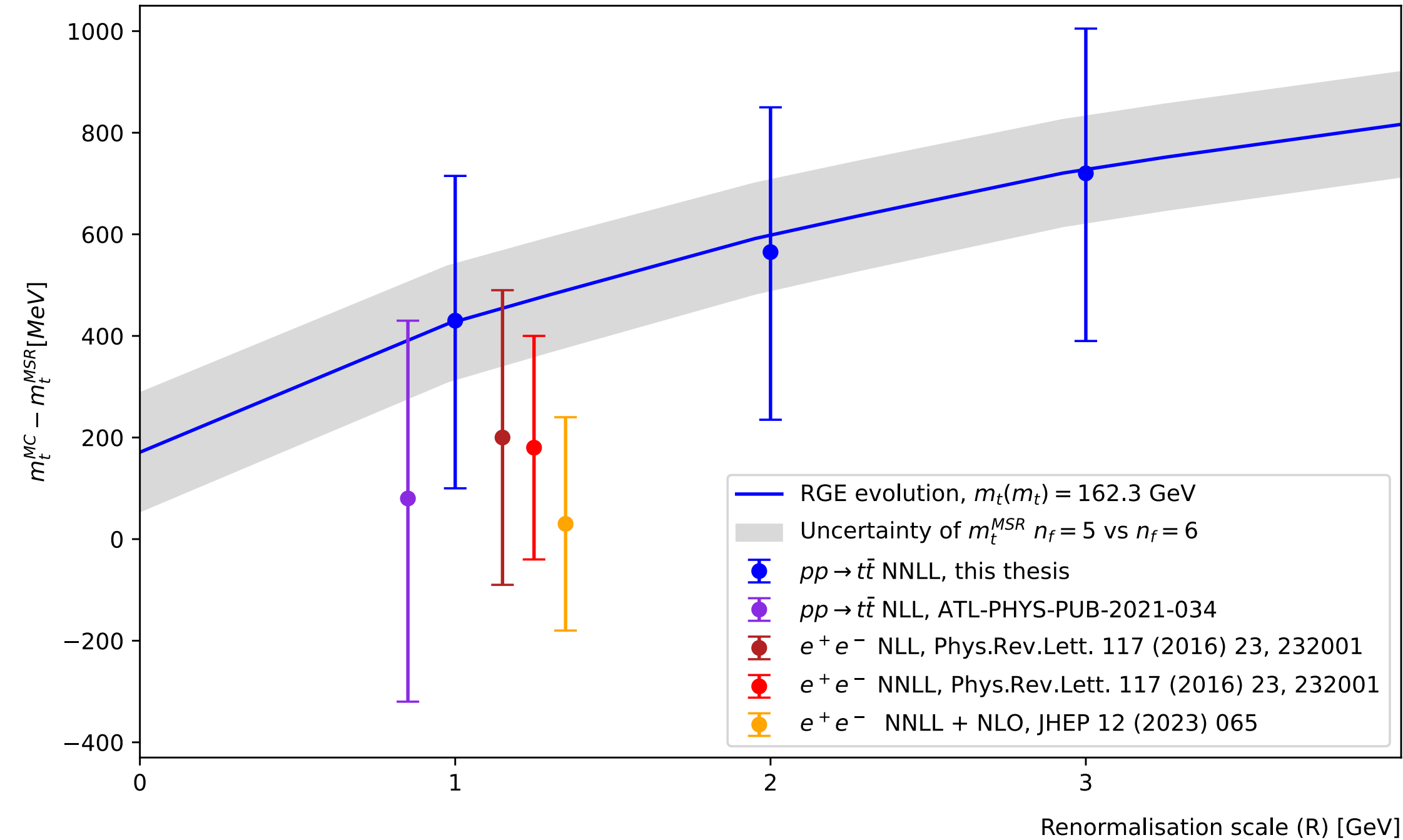
- Previously, we performed a study on the MSR mass for calibration in  $pp \rightarrow t\bar{t}$  processes.

Mass relation of:

$$\Delta^{MSR} = m_t^{MC} - m_t^{MSR}(R = 3 \text{ GeV}) = 720^{+285}_{-330} \text{ MeV}$$

- Uncertainties decreased significantly from previous relation with  $pp$  processes.
- Optimised  $R$  value and improved NNLL calculation.

Relation between MC mass parameter and MSR mass



- Comparisons with previous top quark mass interpretations in  $pp$  and  $e^+e^-$  collision regimes.
  - All results compatible within uncertainties.

# Comps eB = 2 vs normal

Nominal result  $p_T = 750-1250$  GeV:

$$m_t^{pole} = 171.92$$

$$\Omega_1 = 2.58$$

$$x_2 = 0.84$$

Nominal result  $p_T = 750-1250$  GeV (eB=2):

$$m_t^{pole} = 171.54$$

$$\Omega_1 = 2.33$$

$$x_2 = 0.60$$

Nominal result  $E = 2500$  GeV:

$$m_t^{pole} = 171.90$$

$$\Omega_1 = 2.72$$

$$x_2 = 0.89$$

Nominal result  $E = 2500$  GeV (eB=2):

$$m_t^{pole} = 171.75$$

$$\Omega_1 = 2.42$$

$$x_2 = 0.76$$

# pp Nominal results

Nominal result  $p_T = 750-1250$  GeV (eB=2):

$$m_t^{pole} = 171.54$$

$$\Omega_1 = 2.33$$

$$x_2 = 0.60$$

Nominal result  $p_T = 750-1250$  GeV:

$$m_t^{pole} = 171.92$$

$$\Omega_1 = 2.58$$

$$x_2 = 0.84$$

Nominal result combined  $p_T$ :

$$m_t^{pole} = 171.91$$

$$\Omega_1 = 2.59$$

$$x_2 = 0.75$$

Nominal result  $p_T = 1250-2500$  GeV:

$$m_t^{pole} = 171.91$$

$$\Omega_1 = 2.56$$

$$x_2 = 0.72$$

# ee Nominal results

Nominal result combined  $E$ :

$$m_t^{pole} = 171.90$$

$$\Omega_1 = 2.79$$

$$x_2 = 0.88$$

Nominal result  $E = 2500$  GeV:

$$m_t^{pole} = 171.90$$

$$\Omega_1 = 2.72$$

$$x_2 = 0.89$$

Nominal result  $E = 3500$  GeV:

$$m_t^{pole} = 171.90$$

$$\Omega_1 = 2.82$$

$$x_2 = 0.89$$

# h bin comparison

